Time-Dependent Green's Functions approach to nuclear reactions

Understanding the 1D mean-field dynamics

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1

Time-dependent formalism for nuclear reactions

Our "ambitious" goal:

- Simulate time evolution of nuclear phenomena in 3D
- Use Time-Dependent Green's Functions formalism
 - Good in the static case
 - Fully quantal
 - Correlations in initial state and in dynamics
 - Microscopic conservation laws are preserved

Where we are right now:

- Mean-field simulation of collisions of 1D slabs
- Understanding TDGF formalism before...
 - Including correlations in 1D (NN collisions)
 - Extending to higher dimensions (2D, 3D)

Kadanoff-Baym equations

Eventually:

$$\begin{split} \rho(\mathbf{x}_{1}t_{1},\mathbf{x}_{1'}t_{1'}) &= \mathcal{G}^{<}(\mathbf{11}') = i \left\langle \Phi_{0} \middle| \hat{a}^{\dagger}(\mathbf{x}_{1'}t_{1'}) \hat{a}(\mathbf{x}_{1}t_{1}) \middle| \Phi_{0} \right\rangle \\ \mathcal{G}^{>}(\mathbf{11}') &= i \left\langle \Phi_{0} \middle| \hat{a}(\mathbf{x}_{1}t_{1}) \hat{a}^{\dagger}(\mathbf{x}_{1'}t_{1'}) \middle| \Phi_{0} \right\rangle \end{split}$$

$$\begin{cases} i\frac{\partial}{\partial t_{1}} + \frac{\nabla_{1}^{2}}{2m} \end{cases} \mathcal{G}^{\leq}(\mathbf{1}\mathbf{1}') = \int d\bar{\mathbf{r}}_{1} \Sigma_{HF}(\mathbf{I}\bar{\mathbf{I}}) \mathcal{G}^{\leq}(\bar{\mathbf{I}}\mathbf{1}') \\ + \int_{t_{0}}^{t_{0}} d\bar{\mathbf{I}} \left[\Sigma^{>}(\mathbf{I}\bar{\mathbf{I}}) - \Sigma^{<}(\mathbf{I}\bar{\mathbf{I}}) \right] \mathcal{G}^{\leq}(\bar{\mathbf{I}}\mathbf{1}') - \int_{t_{0}}^{t_{1}'} d\bar{\mathbf{I}} \Sigma^{\leq}(\mathbf{I}\bar{\mathbf{I}}) \left[\mathcal{G}^{>}(\bar{\mathbf{I}}\mathbf{1}') - \mathcal{G}^{<}(\bar{\mathbf{I}}\mathbf{1}') \right] \\ -i\frac{\partial}{\partial t_{1'}} + \frac{\nabla_{1'}^{2}}{2m} \end{cases} \mathcal{G}^{\leq}(\mathbf{I}\mathbf{1}') = \int d\bar{\mathbf{r}}_{1} \mathcal{G}^{\leq}(\mathbf{I}\bar{\mathbf{I}}) \Sigma_{HF}(\bar{\mathbf{I}}\mathbf{1}') \\ + \int_{t_{0}}^{t_{0}} d\bar{\mathbf{I}} \left[\mathcal{G}^{>}(\mathbf{I}\bar{\mathbf{I}}) - \mathcal{G}^{<}(\mathbf{I}\bar{\mathbf{I}}) \right] \Sigma^{\leq}(\bar{\mathbf{I}}\mathbf{1}') - \int_{t_{0}}^{t_{1}'} d\bar{\mathbf{I}} \mathcal{G}^{\leq}(\mathbf{I}\bar{\mathbf{I}}) \left[\Sigma^{>}(\bar{\mathbf{I}}\mathbf{1}') - \Sigma^{<}(\bar{\mathbf{I}}\mathbf{1}') \right] \end{cases}$$

- Evolution equations for non-equilibrium systems (Keldysh)
- Can be derived from general principles
- Include correlation and memory effects
- Complicated numerical solution!

Kadanoff & Baym, Quantum Statistical Mechanics (1962).

Kadanoff-Baym equations

Right now:

$$\begin{split} \rho(\mathbf{x}_{1}t_{1},\mathbf{x}_{1'}t_{1'}) &= \mathcal{G}^{<}(\mathbf{11}') = i \left\langle \Phi_{0} \middle| \hat{a}^{\dagger}(\mathbf{x}_{1'}t_{1'}) \hat{a}(\mathbf{x}_{1}t_{1}) \middle| \Phi_{0} \right\rangle \\ \mathcal{G}^{>}(\mathbf{11}') &= i \left\langle \Phi_{0} \middle| \hat{a}(\mathbf{x}_{1}t_{1}) \hat{a}^{\dagger}(\mathbf{x}_{1'}t_{1'}) \middle| \Phi_{0} \right\rangle \end{split}$$

$$\left\{i\frac{\partial}{\partial t_1}+\frac{\nabla_1^2}{2m}\right\}\mathcal{G}^{\lessgtr}(\mathbf{1}\mathbf{1}')=\int\!\mathrm{d}\bar{\mathbf{r}}_1\Sigma_{HF}(\mathbf{1}\bar{\mathbf{1}})\mathcal{G}^{\lessgtr}(\bar{\mathbf{1}}\mathbf{1}')$$

$$\left\{-i\frac{\partial}{\partial t_{1'}}+\frac{\nabla_{1'}^2}{2m}\right\}\mathcal{G}^{\lessgtr}(\mathbf{1}\mathbf{I}')=\int\!\mathrm{d}\bar{\mathbf{r}}_1\mathcal{G}^{\lessgtr}(\mathbf{1}\bar{\mathbf{I}})\Sigma_{HF}(\bar{\mathbf{1}}\mathbf{I}')$$

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TDGF in the MF approximation vs. TDHF

- TDGF and TDHF give exactly the same results...
- but are expressed in rather different terms!

Time Dependent Green's Functions	Time Dependent Hartree-Fock
$i\frac{\partial}{\partial t}\mathcal{G}^{<}(x,x';t) = \left\{-\frac{1}{2m}\frac{\partial^{2}}{\partial x^{2}} + U(x)\right\}\mathcal{G}^{<}(x,x';t)$ $-i\frac{\partial}{\partial t}\mathcal{G}^{<}(x,x';t) = \left\{-\frac{1}{2m}\frac{\partial^{2}}{\partial x'^{2}} + U(x')\right\}\mathcal{G}^{<}(x,x';t)$	for $\alpha = 1,, N_{\alpha}$ $i \frac{\partial}{\partial t} \varphi_{\alpha}(x, t) = \left\{ -\frac{1}{2m} \frac{\partial^2}{\partial x^2} + U(x) \right\} \varphi_{\alpha}(x, t)$ end
 2 equations N_x × N_x matrix Testing ground & new perspective 	 1 equation N_x × N_α matrix Limited to mean-field!

- First calculation: mean-field time evolution of 1D slabs
- Understand time evolution of uncorrelated density matrices

Collisions of 1D slabs

- Frozen *y*, *z* coordinates, dynamics in *x*
- Simplified Skyrme mean field:

$$U(x) = \frac{3}{4}t_0 n(x) + \frac{2+\sigma}{16}t_3 [n(x)]^{(\sigma+1)}$$

- Initial state from adiabatic theorem or static HF
- Ground state of $\mathcal{G}^{<}$ with A = 8 ($N_{\alpha} = 2$)

Two aspects:

- Phenomenology: dependence on bombarding energy?
- New perspective from Green's functions \Rightarrow off-diagonal elements

Mean-field approximation

Diagonal & off-diagonal matrix elements

• Diagonal elements yield physical properties:

$$n(x) = \mathcal{G}^{<}(x, x' = x) = \sum_{\alpha=0}^{N_{\alpha}} n_{\alpha} |\varphi_{\alpha}(x)|^{2}$$
$$K = \sum_{k} \frac{k^{2}}{2m} \mathcal{G}^{<}(k, k' = k)$$

• What happens off the diagonal?

$$\mathcal{G}^{<}(x,x') = \sum_{\alpha=0}^{N_{\alpha}} n_{\alpha} \varphi_{\alpha}(x) \varphi_{\alpha}^{*}(x')$$

- What is their meaning? $\Rightarrow \mathcal{G}^{<} \sim \text{correlation function}$
- Are all $x \neq x'$ necessary for the time-evolution?
- Eliminate them at higher D's to avoid $N_x^D \times N_x^D \times N_t$ matrices?

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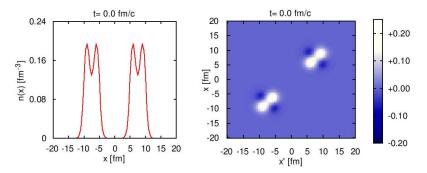
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Collisions of 1D slabs: fusion (off-diagonal)

$$\mathcal{G}^{<}(x,x',P) = e^{iPx} \mathcal{G}^{<}(x,x',P=0) e^{-iPx}$$

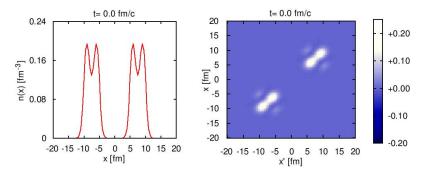
$$E_{CM}/A = 0.1 \text{ MeV}$$



Collisions of 1D slabs: break-up (off-diagonal)

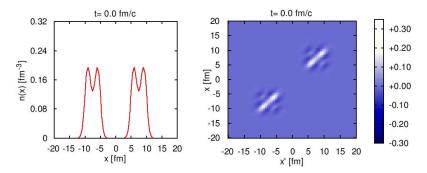
$$\mathcal{G}^{<}(x,x') = \sum_{\alpha} n_{\alpha} \varphi_{\alpha}(x) \varphi_{\alpha}(x')$$

$$E_{CM}/A = 4 \text{ MeV}$$

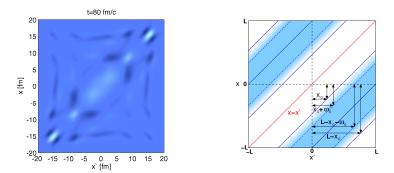


Collisions of 1D slabs: multifrag. (off-diagonal)

$$E_{CM}/A = 25 \text{ MeV}$$



Erasing off-diagonal elements

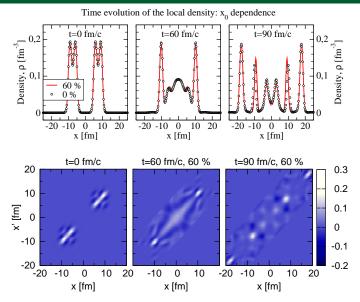


• Use a cut-off imaginary potential off the diagonal ~ superoperator

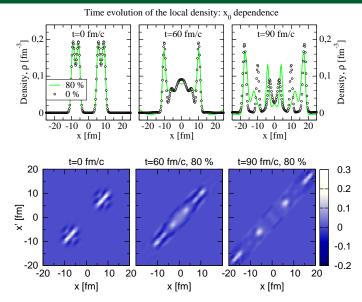
$$\mathcal{G}^{<}(x, x', t + \Delta t) \sim e^{i(\varepsilon(x) + iW(x, x'))\Delta t} \mathcal{G}^{<}(x, x', t) e^{-i(\varepsilon(x') - iW(x, x'))\Delta t}$$

Properties of the field set to preserve norm, periodicity, etc
Other choices do not yield good results!

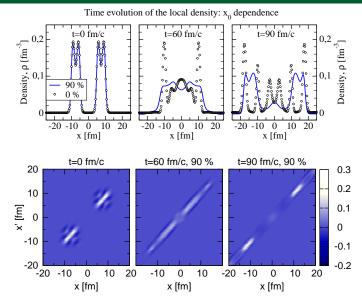
Erasing off-diagonal elements



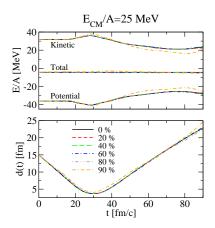
Erasing off-diagonal elements



Erasing off-diagonal elements



Erasing off-diagonal elements



• Up to 60% (safe) off-diagonal elements can be neglected safely?

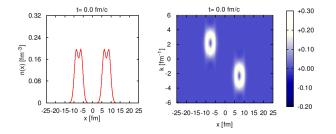
• Small effect of erasure for observables in high energy reactions!

Connection to transport phenomena

• Wigner transform:

$$f(x,P) = \int \frac{\mathrm{d}y}{2\pi} e^{-iPy} \mathcal{G}^{<}\left(x+\frac{y}{2},x-\frac{y}{2}\right)$$

- f(x, P) phase space density
- Obeys Boltzmann equation under assumptions
- Momentum average of $f(x, P) \sim$ erasure of off-diagonal elements

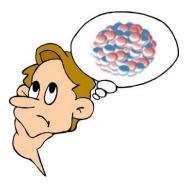


Conclusions

- KB equations offer a consistent framework for TD calculations
- Time evolution of 1D slabs in the mean-field
- New perspective is gained
- 60% off-diagonal elements are unimportant with appropiate cut-off
- Wigner functions are being analyzed
- Qualitative agreement with previous calculations*
- Beyond mean-field calculations are being implemented
- Extension to 2D and 3D ⇒ Promising approach to study correlations in time-dependent approaches!

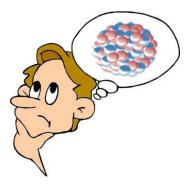
*Bonche, Koonin and Negele, PRC 13, 1226 (1976).

Thank you!



Let me know if you nide a ride back to Kellogg!

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