

Lecture II : $0\nu\beta\beta$ -Decay

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U Mass Amherst



AMHERST CENTER FOR FUNDAMENTAL INTERACTIONS

Physics at the interface: Energy, Intensity, and Cosmic frontiers

University of Massachusetts Amherst

<http://www.physics.umass.edu/acfi/>

NNPSS, Wright Laboratory
Yale 6/18-29/18

Lecture II Goals

- *Give a theoretical overview of $0\nu\beta\beta$ decay*
- *Connect $0\nu\beta\beta$ decay to the origin of matter*
- *Provide a framework for interpreting $0\nu\beta\beta$ decay results:
the mechanisms*
- *Discuss the interplay with other experiments*
- *Invite questions !*

Lecture II Outline

- I. *Overview*
- II. *“Standard Mechanism” for $0\nu\beta\beta$ -Decay*
- III. *TeV Scale LNV*
- IV. *Sub-weak scale LNV*
- V. *Discussion questions*

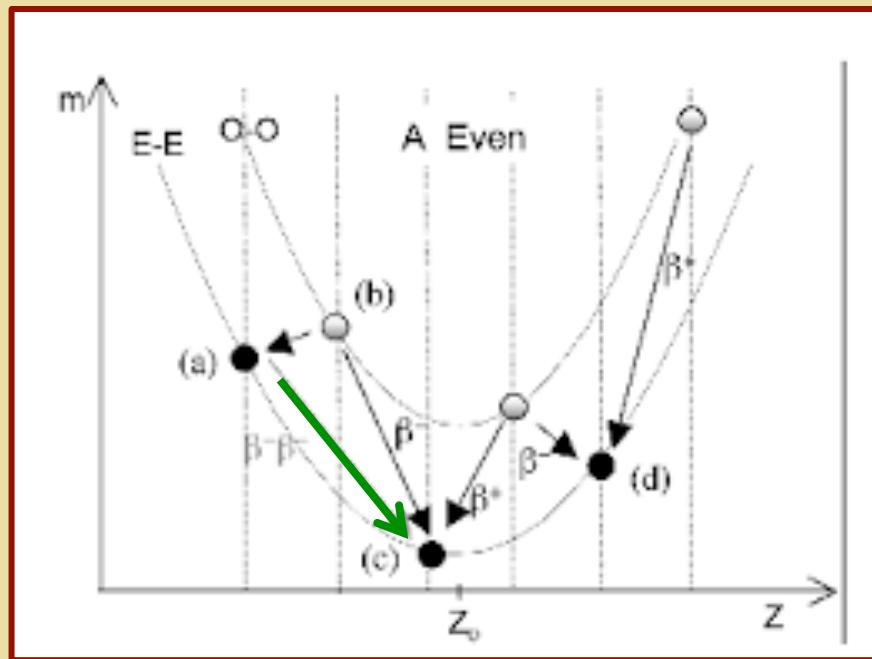
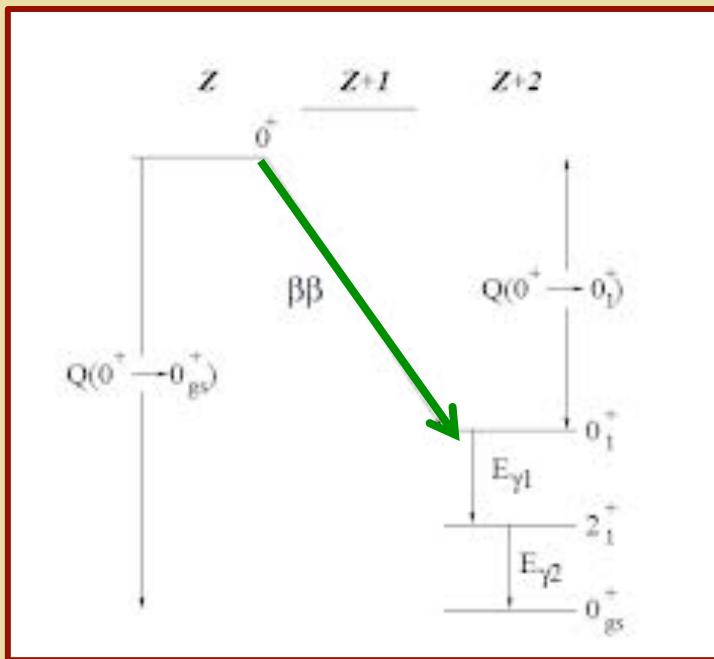
I. $0\nu\beta\beta$ -Decay Overview

What is Neutrinoless Double Beta Decay ?



Test of total lepton number conservation

Why Do Nuclei Double Beta Decay ?



$2\nu DBD:$ $A(Z,N) \rightarrow A(Z+2, N-2) + e^- e^- \bar{\nu} \bar{\nu}$

Observed

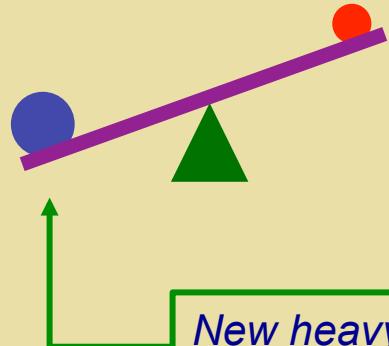
What Questions Does It Address ?

- *Is the neutrino its own antiparticle ?*
- *Why is there more matter than antimatter ?*
- *Why are neutrino masses so small?*

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“See saw mechanism”



“Leptogenesis”

$$\nu = \bar{\nu}$$

Heavy neutrino decays in early universe generate baryon asym

New heavy neutrino-like particle = its own anti-particle

Neutrinos and the Origin of Matter

- *Heavy neutrinos decay out of equilibrium in early universe*
- *Majorana neutrinos can decay to particles and antiparticles*
- *Rates can be slightly different (CP violation)*

$$\Gamma(N \rightarrow \ell H) \neq \Gamma(N \rightarrow \bar{\ell} H^*)$$

- *Resulting excess of leptons over anti-leptons partially converted into excess of quarks over anti-quarks by Standard Model sphalerons*

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0νββ-Decay: LNV? Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

What Questions Does It Address ?

- *Is the neutrino its own antiparticle ?*

2ν DBD: $A(Z,N) \rightarrow A(Z+2, N-2) + e^- e^- \bar{\nu} \bar{\nu}$

0ν DBD: $A(Z,N) \rightarrow A(Z+2, N-2) + e^- e^-$

What Questions Does It Address ?

- *Is the neutrino its own antiparticle ?*

2ν DBD:



*If own antiparticle, can be emitted
then absorbed during decay*

0ν DBD:



What Questions Does It Address ?

- *Is the neutrino its own antiparticle ?*

Yes → “*Majorana neutrino*”:

Theoretically favored explanation of the matter-antimatter asymmetry & small scale of neutrino masses

No → “*Dirac neutrino*”:

Points to alternate origin of matter-antimatter asymmetry & some string theory underpinnings of neutrino masses

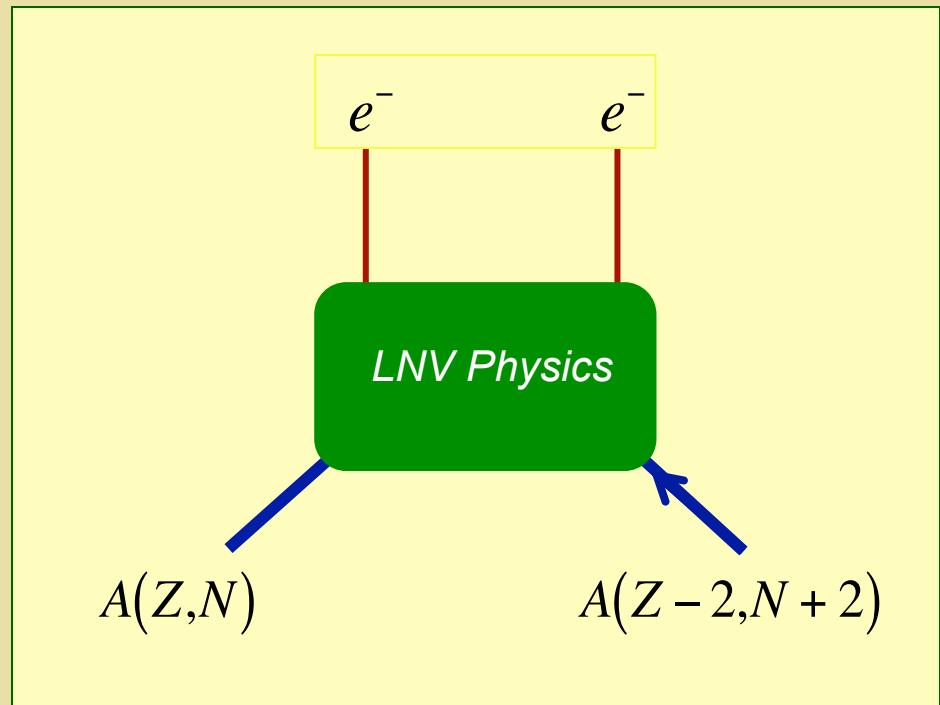
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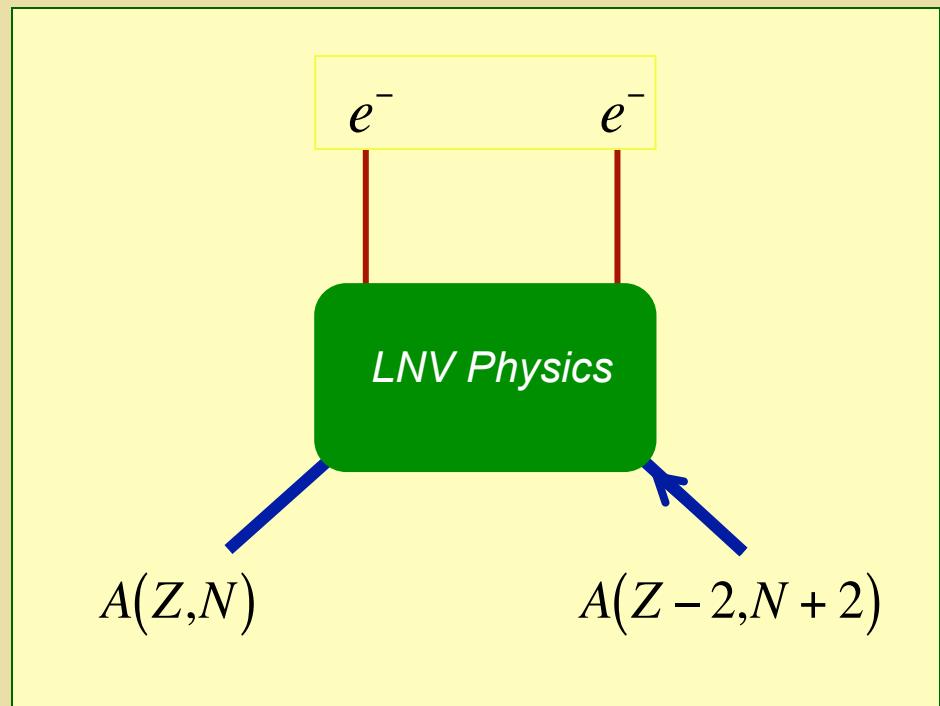
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Impact of observation

- Total lepton number not conserved at classical level
- New mass scale in nature, Λ
- Key ingredient for standard baryogenesis via leptogenesis

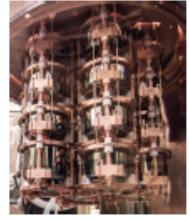


Ton Scale Experiments: Worldwide Quest

0νββ decay Experiments - Efforts Underway

CUORE					
Collaboration	Isotope	Technique	mass (0νββ isotope)	Status	
CANDLES	Ca-48	305 kg CaF ₂ crystals - liq. scint	0.3 kg	Construction	
CARVEL	Ca-48	⁴⁸ CaWO ₄ crystal scint.	~ ton	R&D	
GERDA I	Ge-76	Ge diodes in LAr	15 kg	Complete	
GERDA II	Ge-76	Point contact Ge in LAr	31	Operating	
MAJORANA DEMONSTRATOR	Ge-76	Point contact Ge	25 kg	Operating	
LEGEND	Ge-76	Point contact	~ ton	R&D	
NEMO3	Mo-100 Se-82	Foil with tracking	6.9 kg 0.9 kg	Complete	
SuperNEMO Demonstrator	Se-82	Foil with tracking	7 kg	Construction	
SuperNEMO	Se-82	Foil with tracking	100 kg	R&D	
LUCIFER (CUPID)	Se-82	ZnSe scint. bolometer	18 kg	R&D	
AMORE	Mo-100	CaMoO ₄ scint. bolometer	1.5 - 200 kg	R&D	
LUMINEU (CUPID)	Mo-100	ZnMoO ₄ / Li ₂ MoO ₄ scint. bolometer	1.5 - 5 kg	R&D	
COBRA	Cd-114,116	CdZnTe detectors	10 kg	R&D	
CUORICINO, CUORE-0	Te-130	TeO ₂ Bolometer	10 kg, 11 kg	Complete	
CUORE	Te-130	TeO ₂ Bolometer	206 kg	Operating	
CUPID	Te-130	TeO ₂ Bolometer & scint.	~ ton	R&D	
SNO+	Te-130	0.3% ⁸⁰ Te suspended in Scint	160 kg	Construction	
EXO200	Xe-136	Xe liquid TPC	79 kg	Operating	
nEXO	Xe-136	Xe liquid TPC	~ ton	R&D	
KamLAND Zen	Xe-136	2.7% in liquid scint.	380 kg	Complete	
KamLAND-Zen (I, II)	Xe-136	2.7% in liquid scint.	750 kg	Upgrade	
KamLAND2-Zen	Xe-136	High pressure Xe TPC	5 kg	Operating	
NEXT-NEW	Xe-136	High pressure Xe TPC	100 kg - ton	R&D	
NEXT	Xe-136	High pressure Xe TPC	~ ton	R&D	
PandaX - 1k	Xe-136	High pressure Xe TPC	20 kg	R&D	
DCBA	Nd-150	Nd foils & tracking chambers			

GERDA					
					

MAJORANA					
					

SNO+					
					

The U.S. Context

2015 NSAC Long Range Plan

RECOMMENDATION II

The excess of matter over antimatter in the universe is one of the most compelling mysteries in all of science. The observation of neutrinoless double beta decay in nuclei would immediately demonstrate that neutrinos are their own antiparticles and would have profound implications for our understanding of the matter-antimatter mystery.

We recommend the timely development and deployment of a U.S.-led ton-scale neutrinoless double beta decay experiment.

Why Is It So Challenging to Observe ?

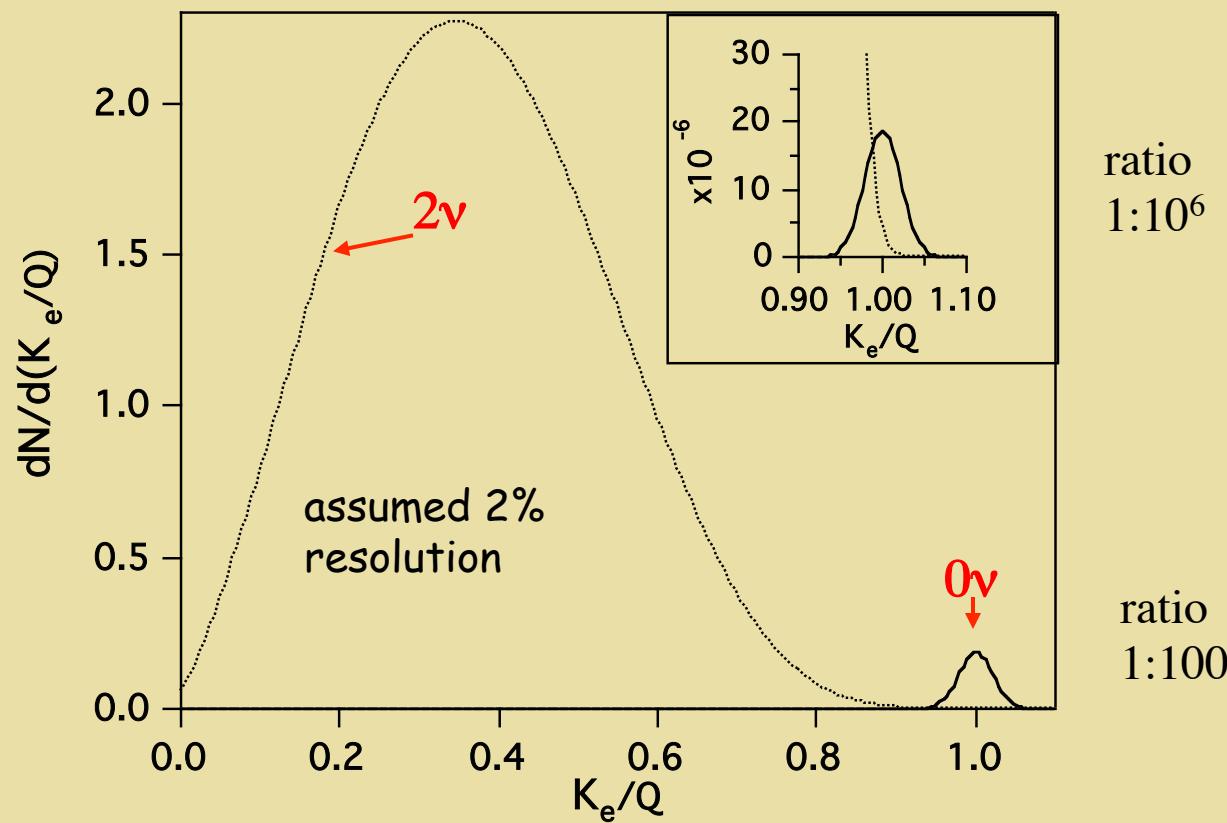
- *The rate is exceptionally tiny*

$$\Gamma \sim (m_{\text{eff}})^2$$

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*Experimental details: See
D. Parno second lecture !*

$0\nu\beta\beta$ -Decay: LNV? Mass Term?

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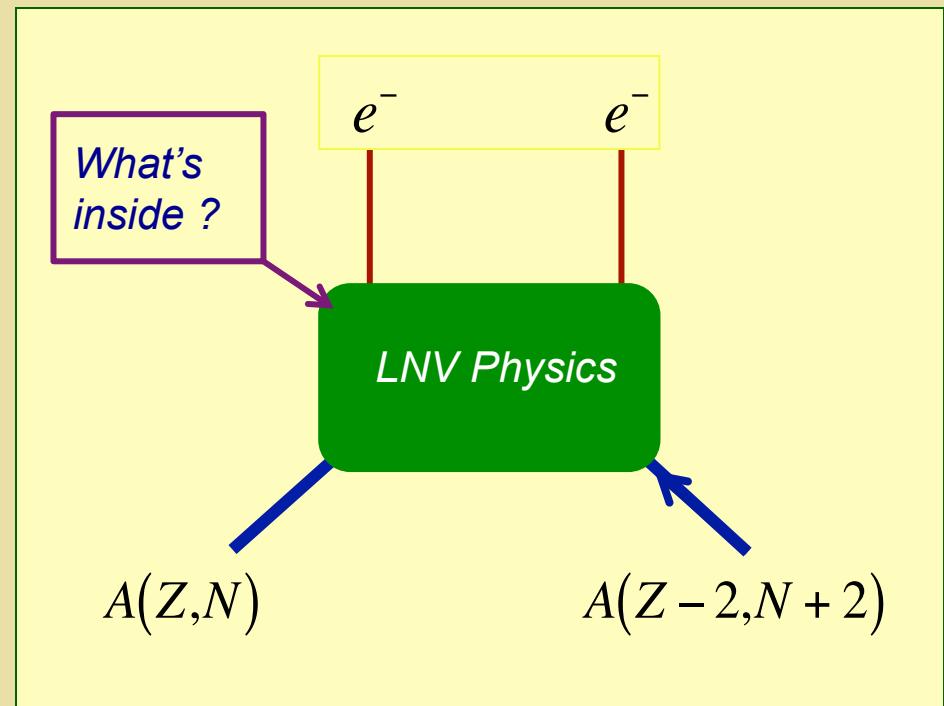
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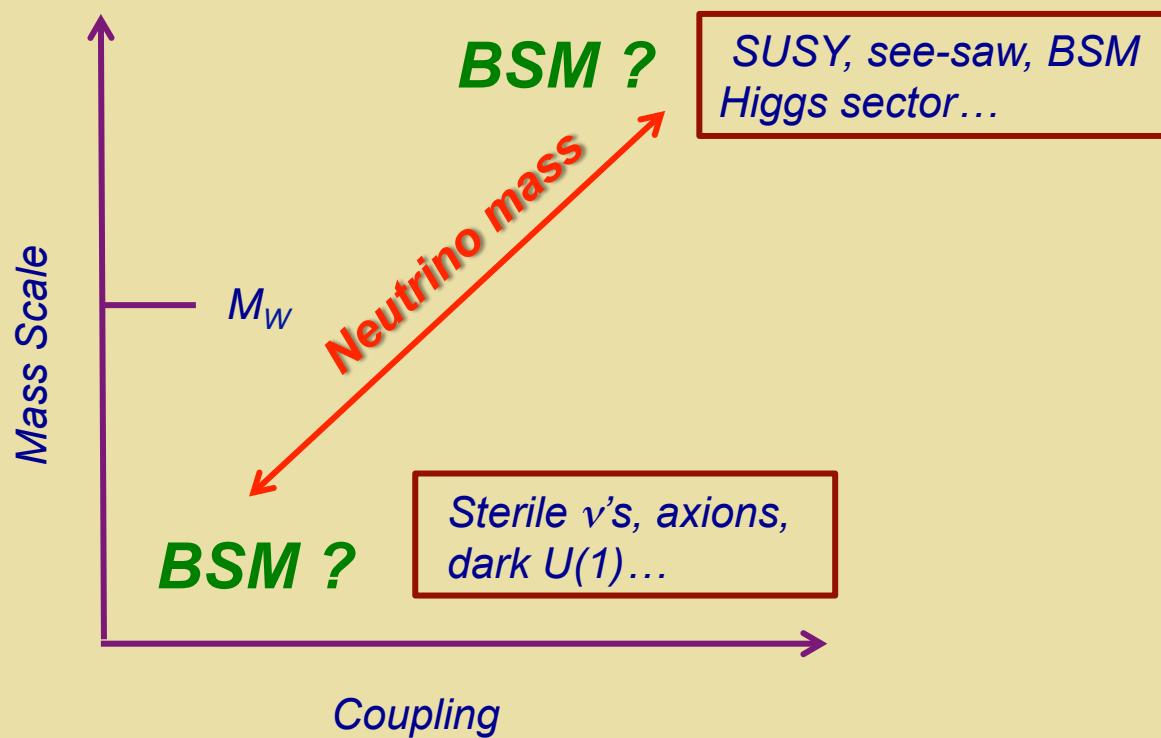
Majorana

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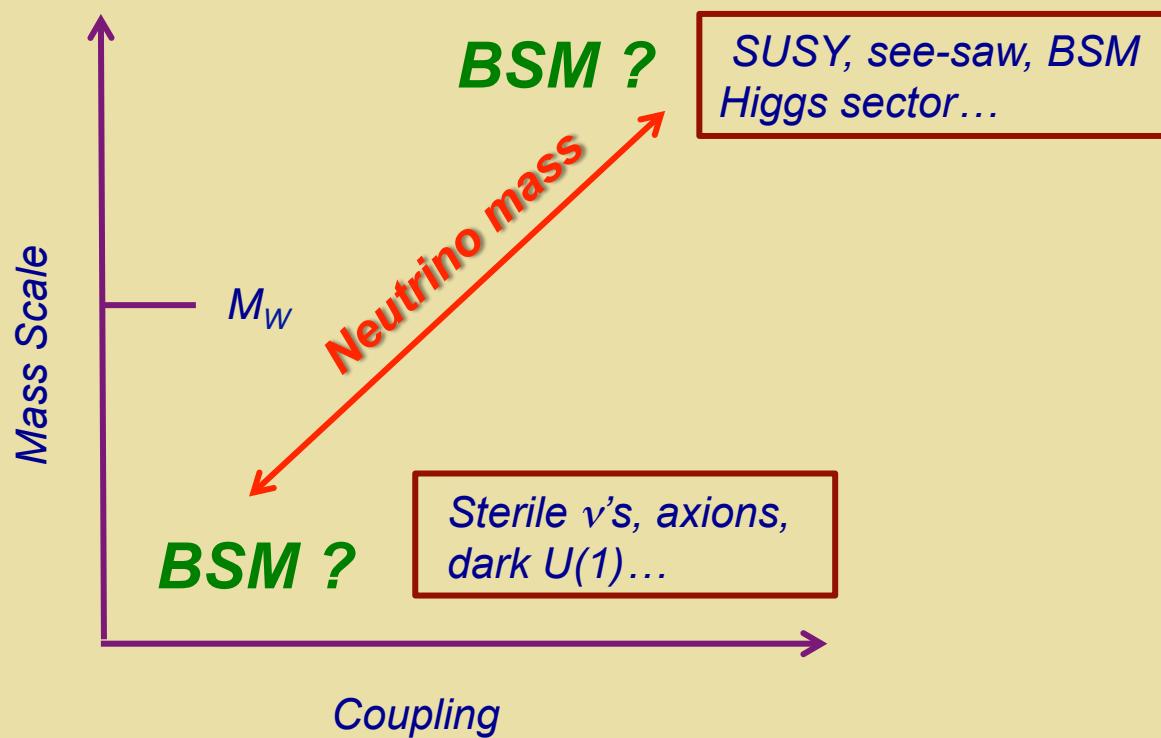
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BSM Physics: Where Does it Live ?



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Is the mass scale associated with m_ν far above M_W ? Near M_W ? Well below M_W ?

Why Might A “Ton-Scale” Expt See It?

$$A(Z, N) \rightarrow \text{Underlying Physics} \rightarrow A(Z+2, N-2) + e^- e^-$$

- *3 light neutrinos only: source of neutrino mass at the very high see-saw scale*
- *3 light neutrinos with TeV scale source of neutrino mass*
- *> 3 light neutrinos*

II. The “Standard Mechanism”

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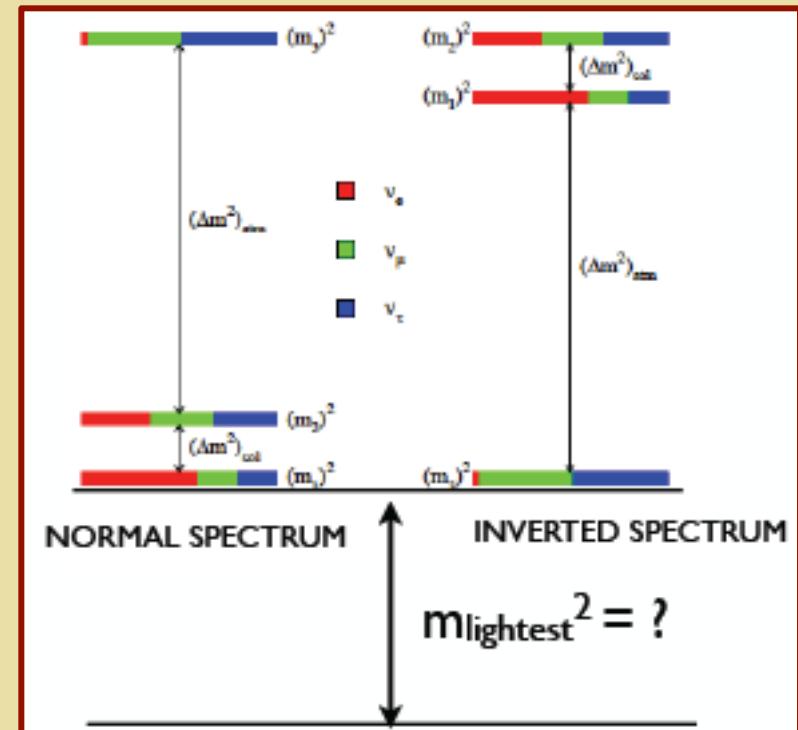
Three Light Neutrinos: What Do We Know ?

Neutrinos mix

$$\begin{pmatrix} \nu_1 & \nu_2 & \nu_3 \end{pmatrix} = U \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

Physical neutrinos

e, μ , τ partners



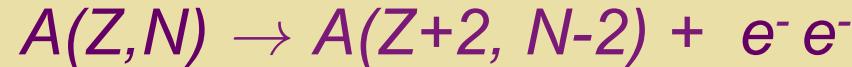
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If own antiparticle, can be emitted
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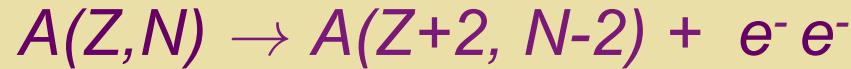
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All three light neutrinos participate →
Rate governed by an **effective mass**

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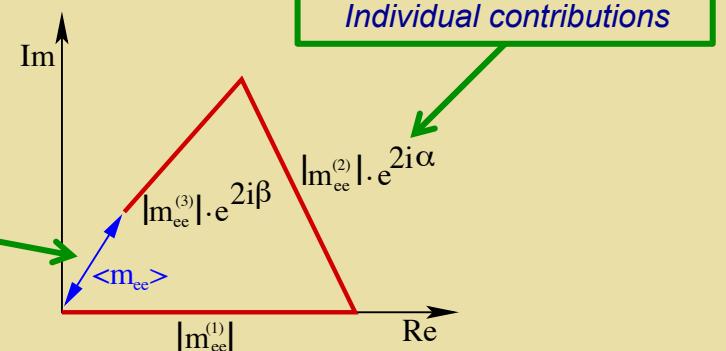
$$\bar{\nu} \bar{\nu}$$

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$0\nu\beta\beta$ -Decay: LNV? Mass Term?

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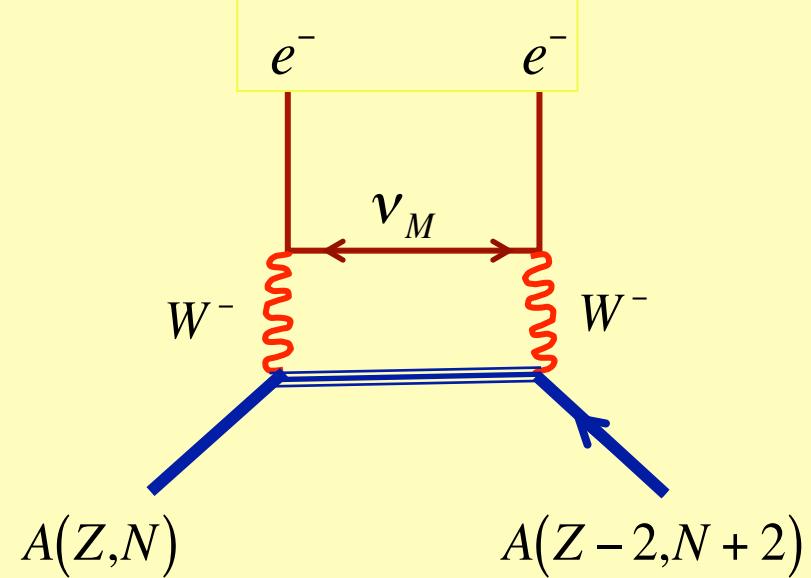
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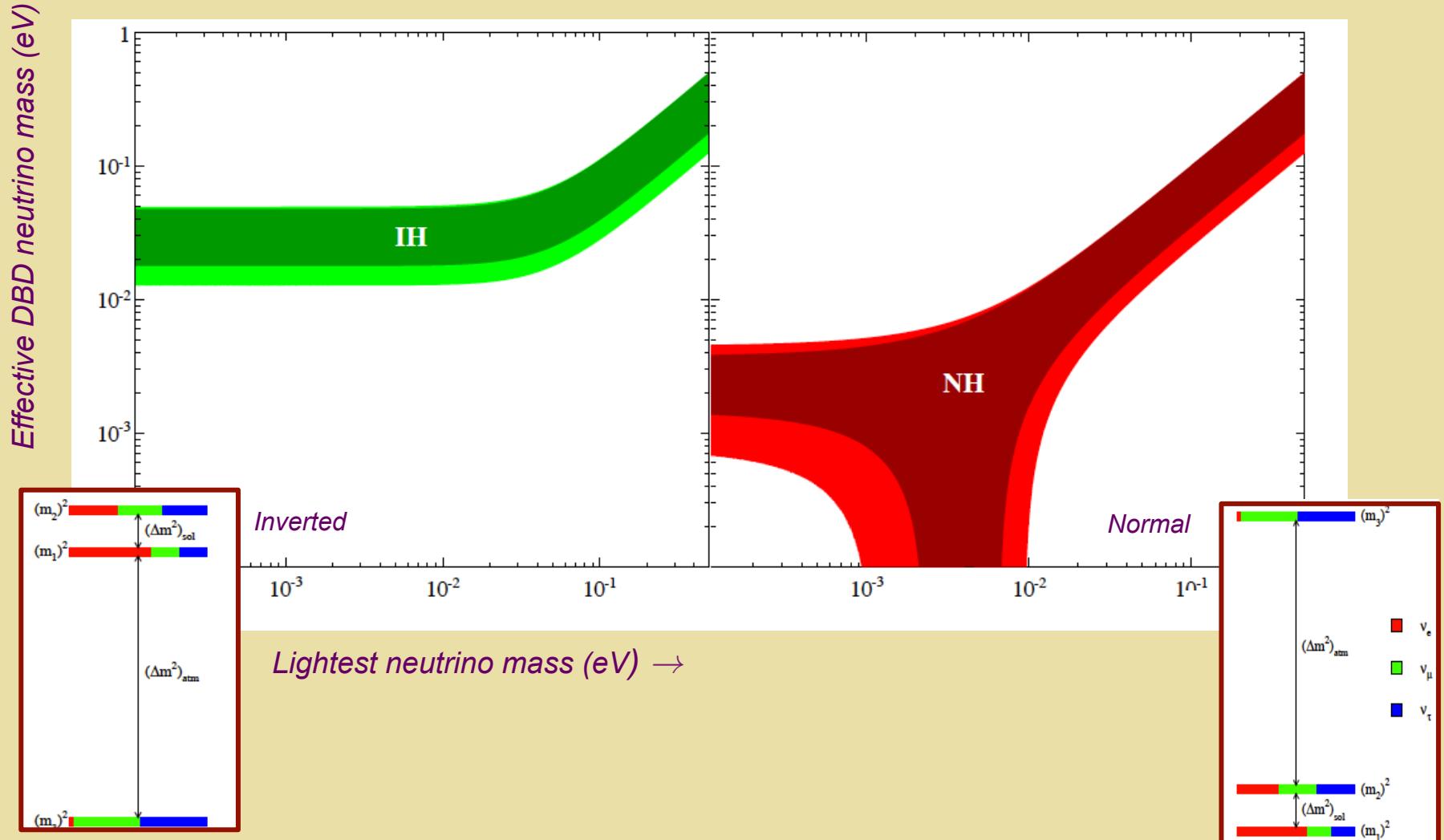
“Standard” Mechanism

- Light Majorana mass generated at the conventional see-saw scale: $\Lambda \sim 10^{12} - 10^{15}$ GeV
- 3 light Majorana neutrinos mediate decay process



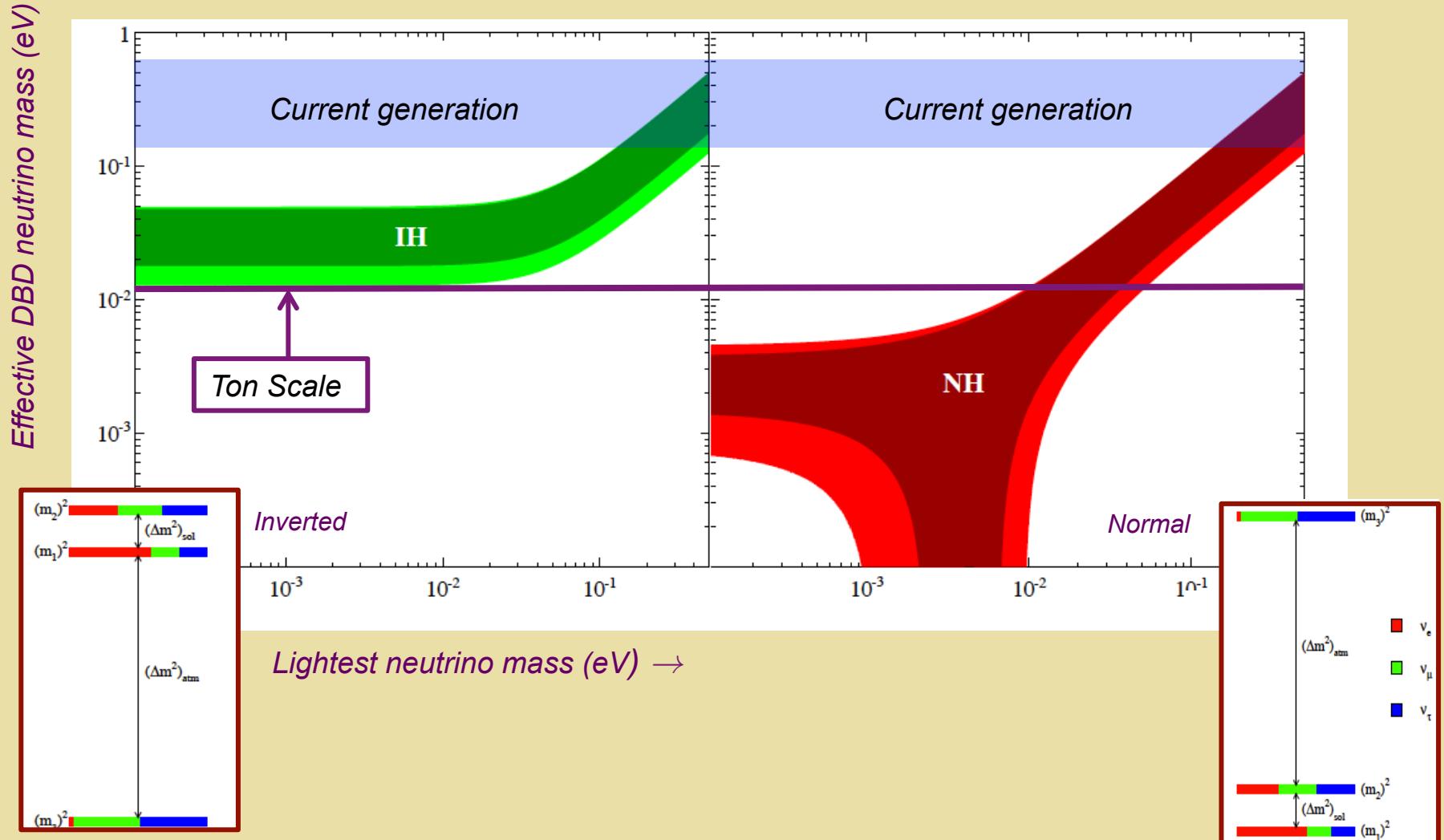
Why Might A “Ton-Scale” Expt See It?

Three active light neutrinos



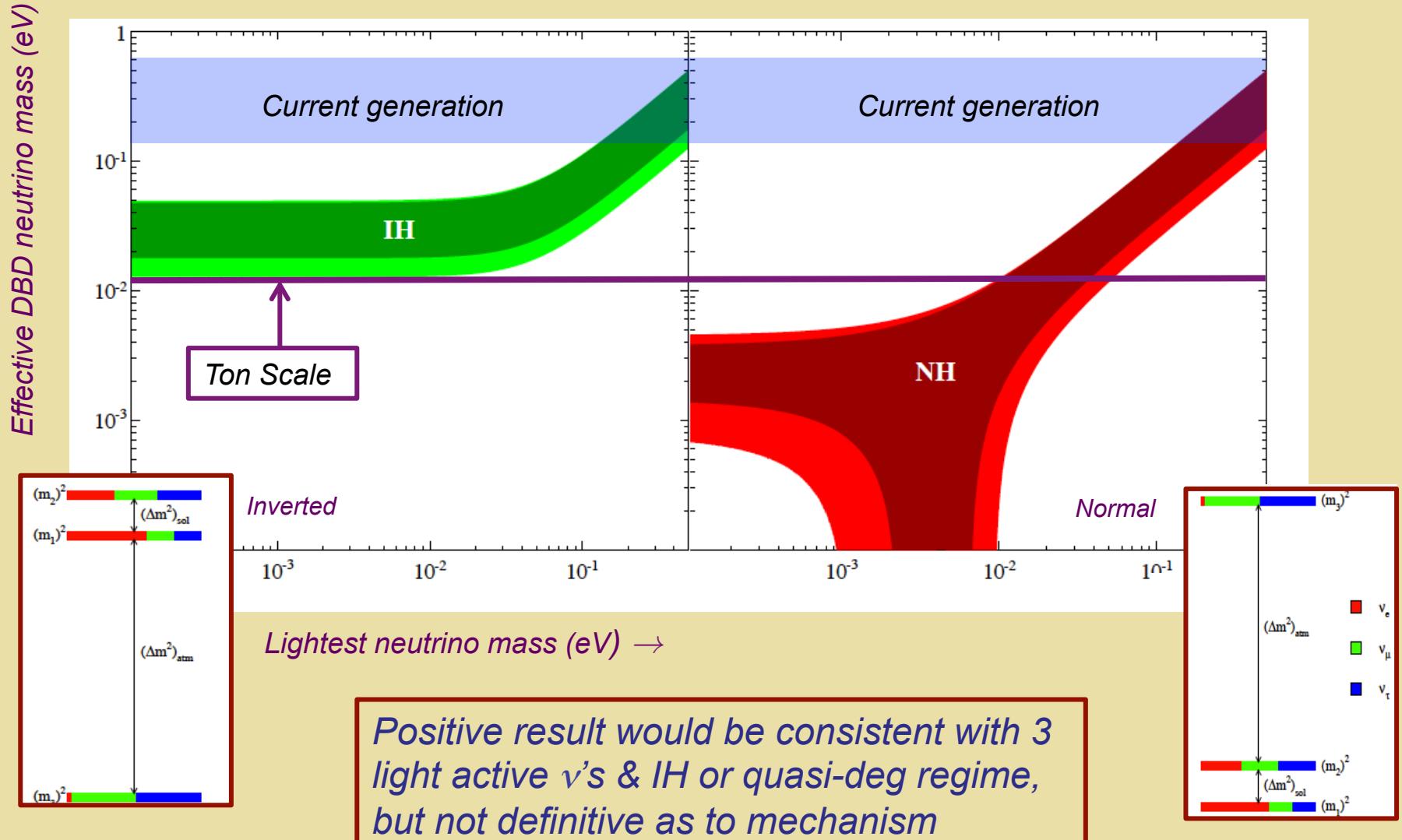
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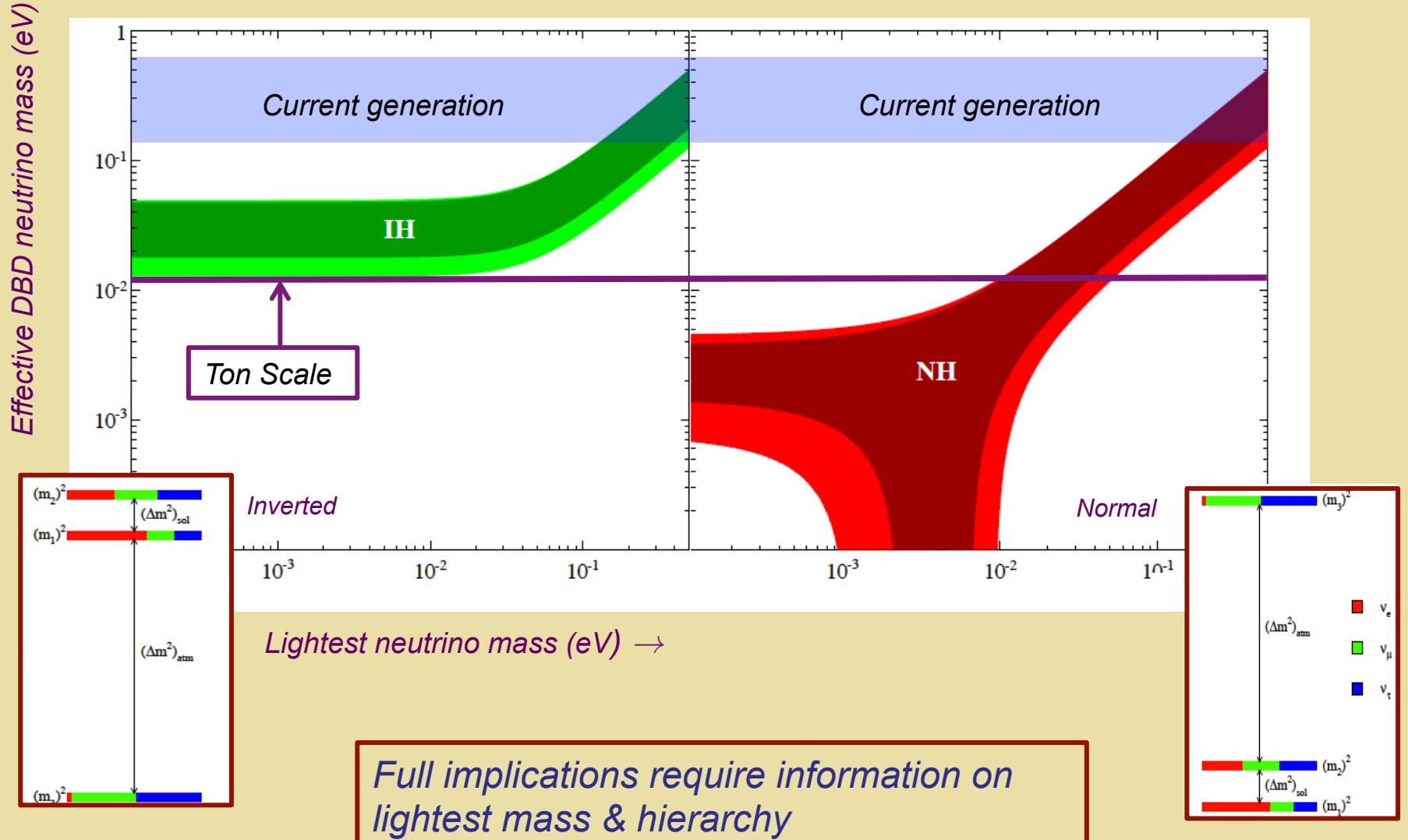
Interpreting a Positive Result

Three active light neutrinos

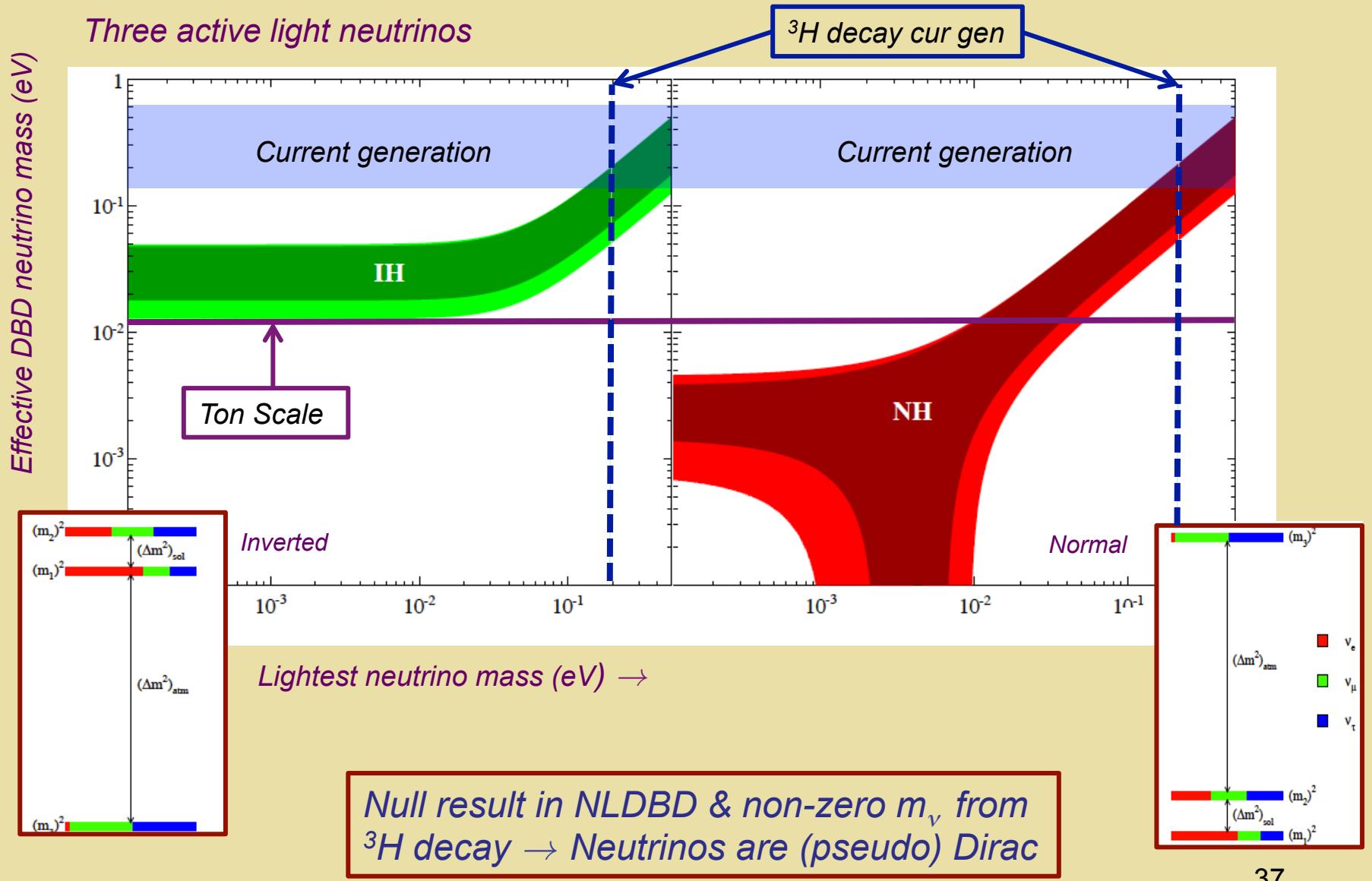


Interpreting a Null Result

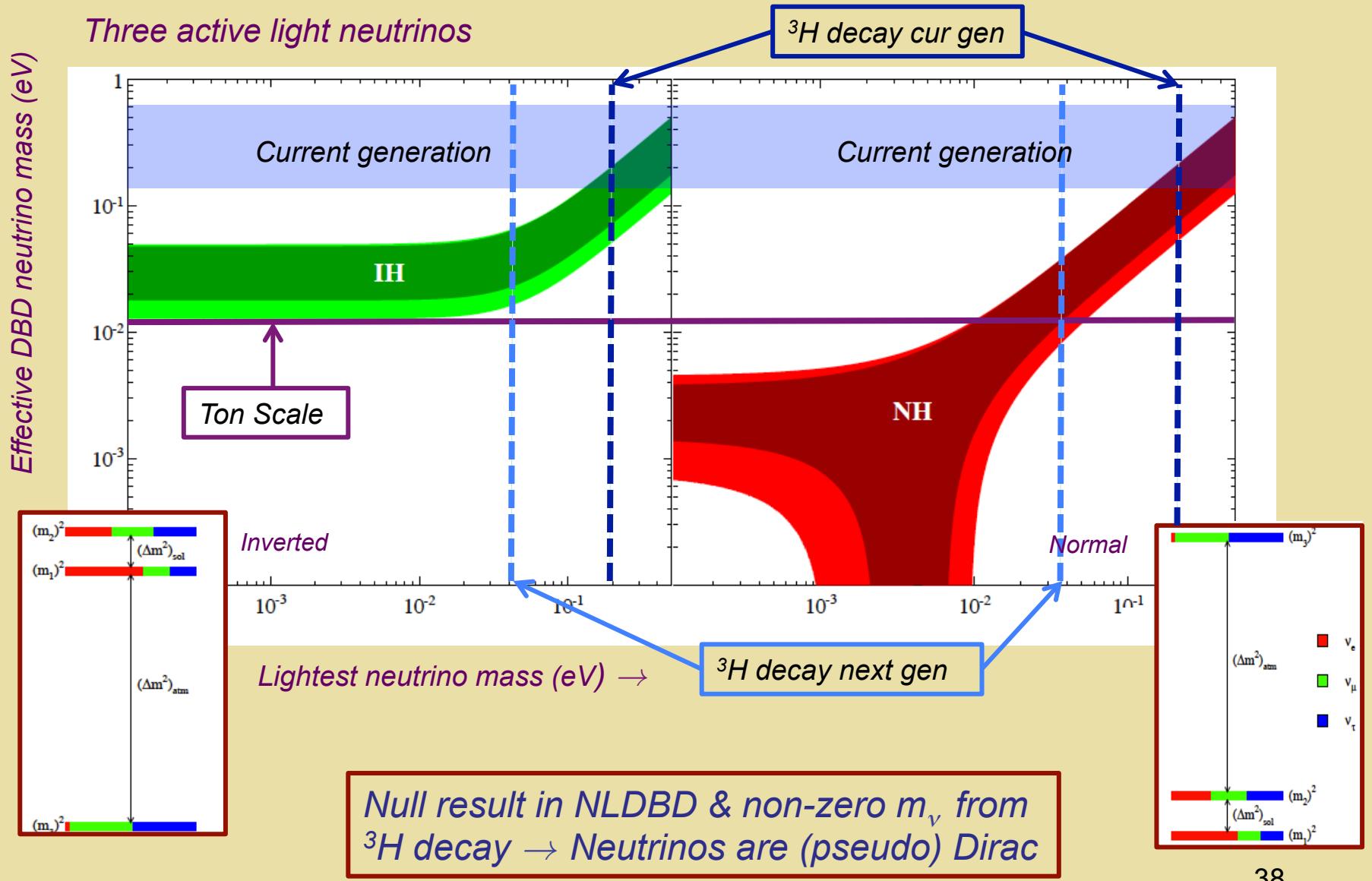
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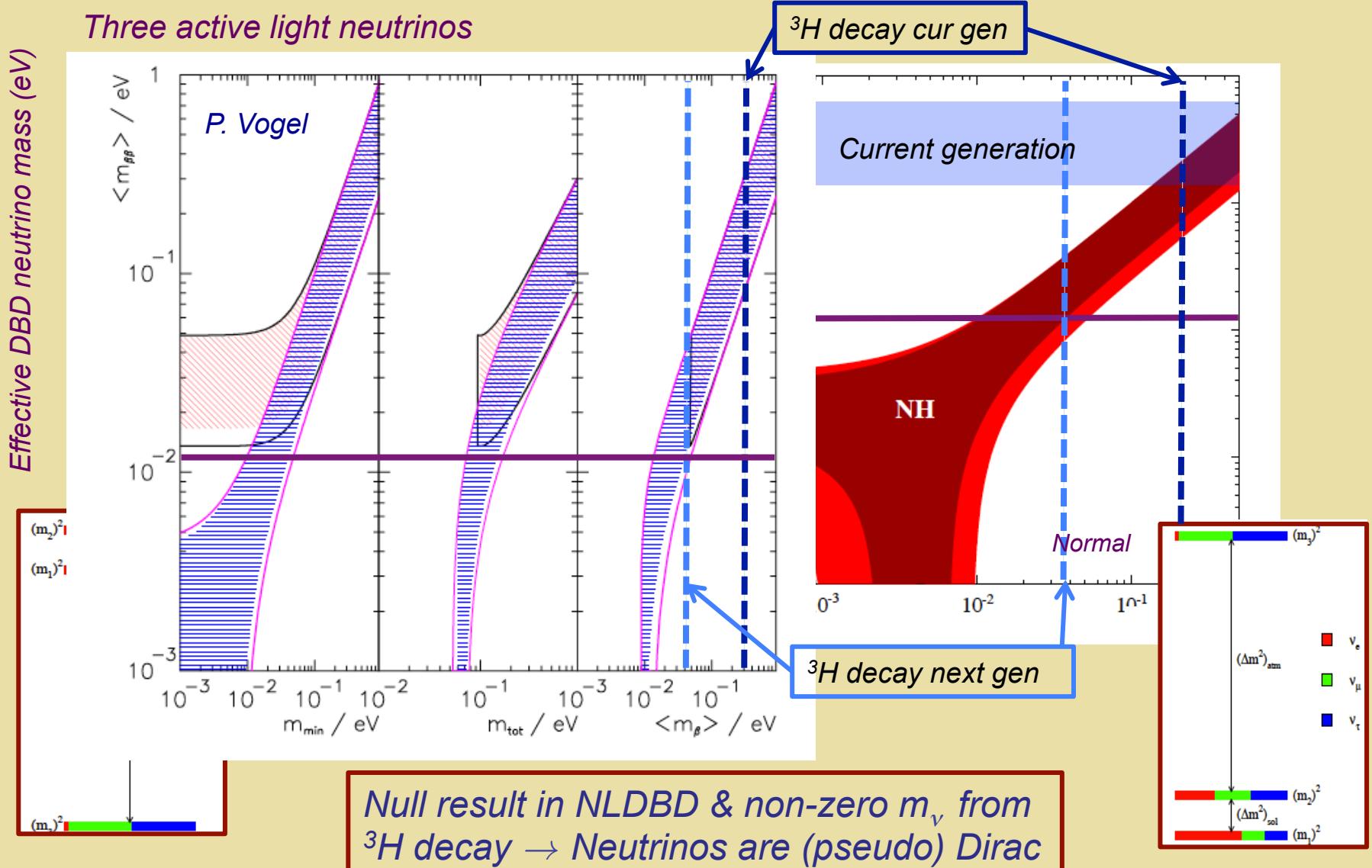
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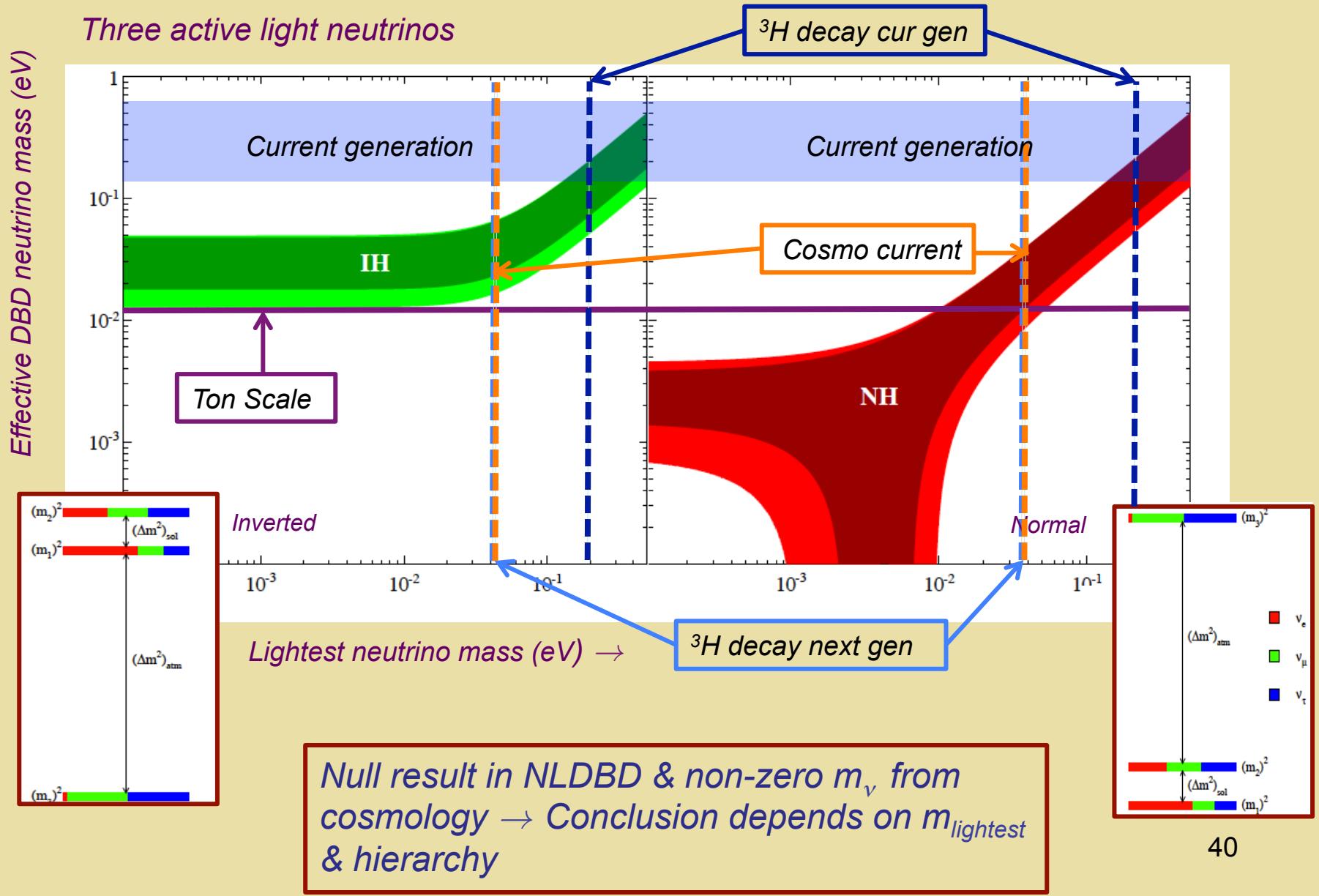
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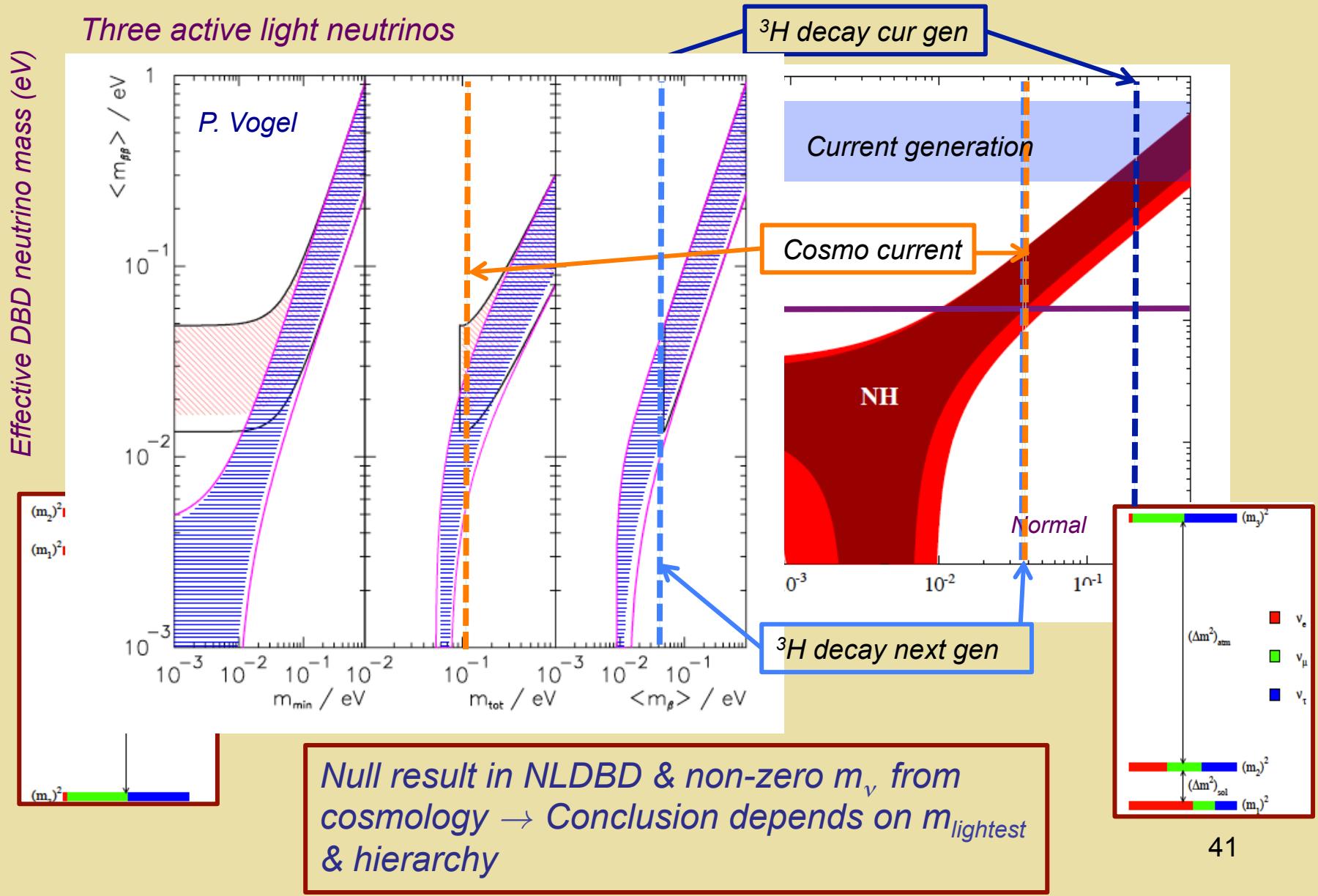
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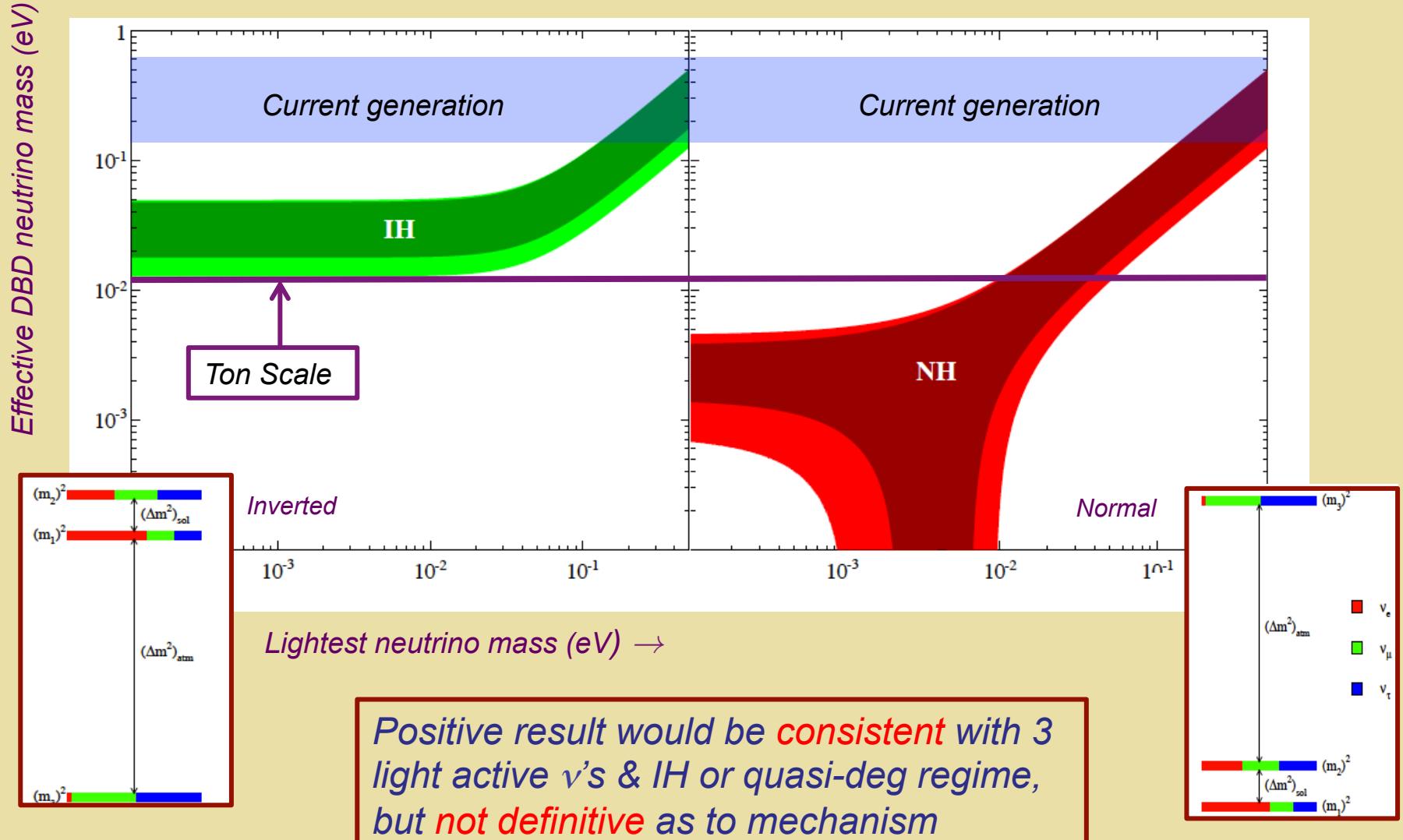


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Interpreting a Positive Result

Three active light neutrinos



$0\nu\beta\beta$ -Decay: Nuclear Matrix Element

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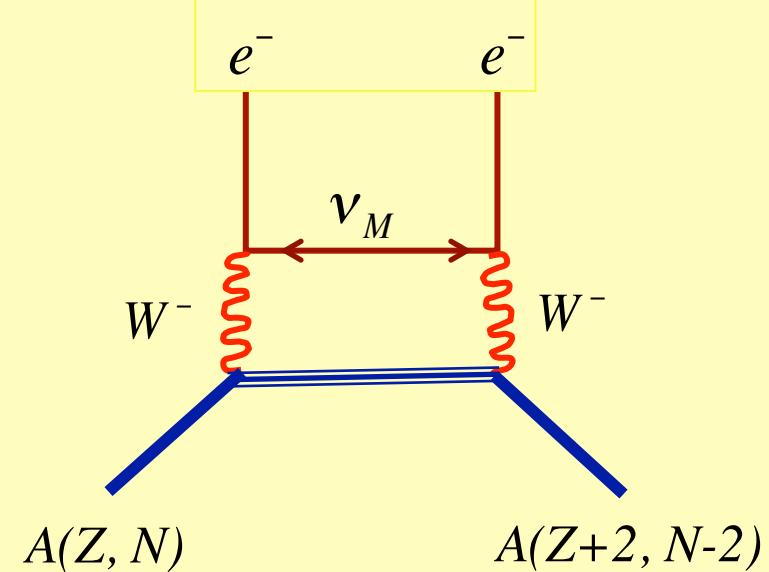
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Majorana

Light ν exchange

$$\frac{1}{T_{1/2}} = G^{0\nu}(E, Z) |M_{0\nu}|^2 |\langle m_{\beta\beta} \rangle|^2$$



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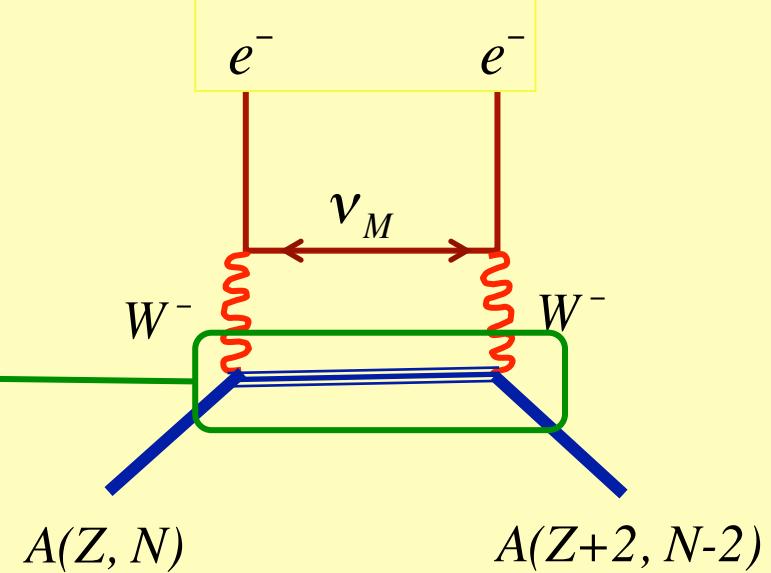
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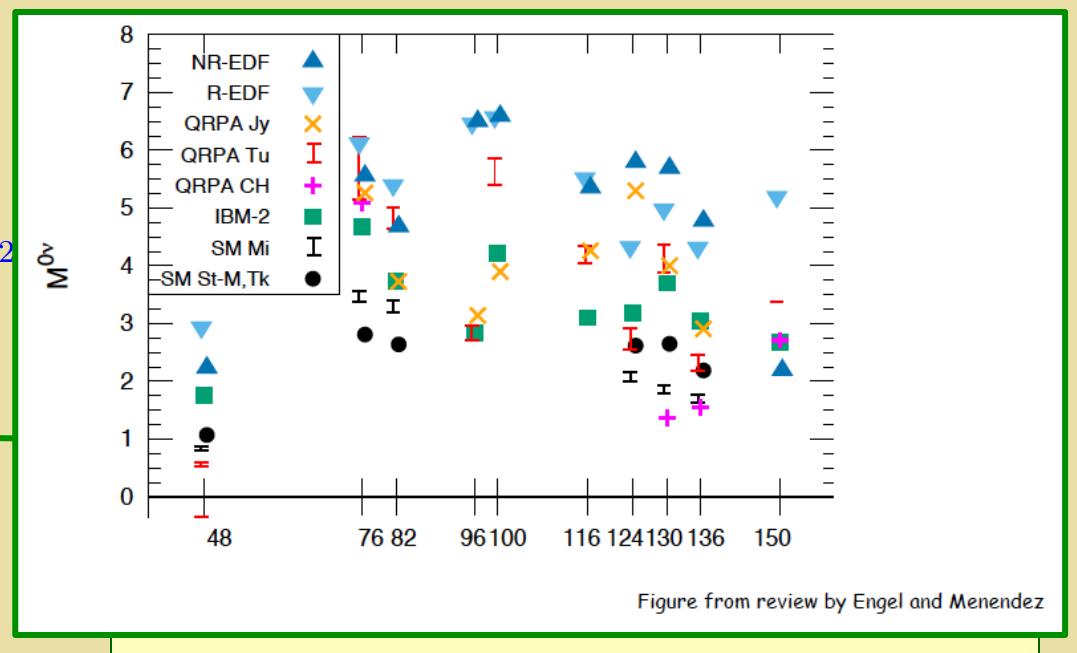
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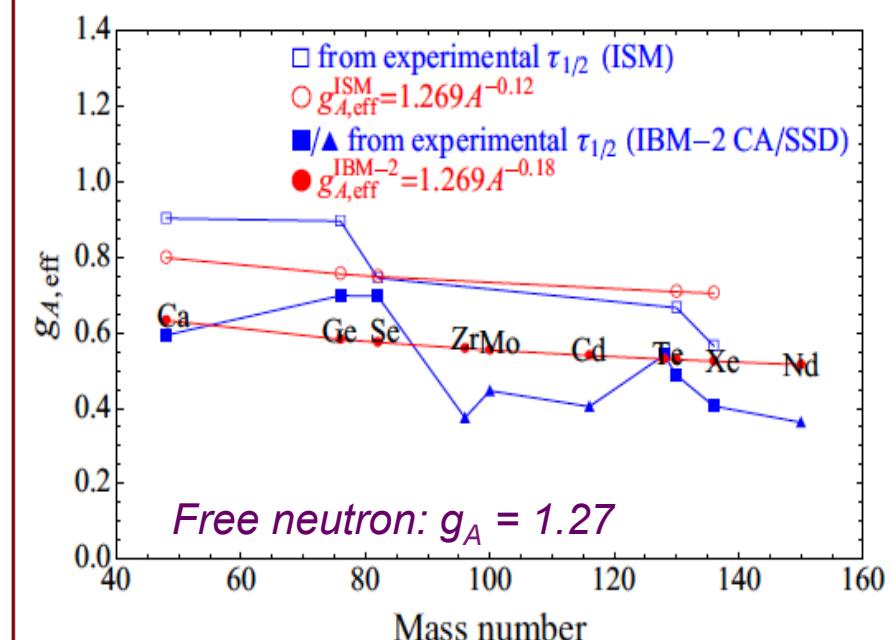
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$M_{0\nu}$: Quadratic dependence on g_A



Free neutron: $g_A = 1.27$

from F. Iachello

$0\nu\beta\beta$ -Decay: Nuclear Matrix Element

Problem must be due to some combination of:

J. Engel

1. Truncation of model space.

Should be fixable in ab-initio shell model, which compensates effects of truncation via effective operators.

2. Many-body weak currents.

Size still not clear, particularly for $0\nu\beta\beta$ decay, where current is needed at finite momentum transfer q .

Leading terms in chiral EFT for finite q only recently worked out. Careful fits and use in decay computations will happen in next year or two.

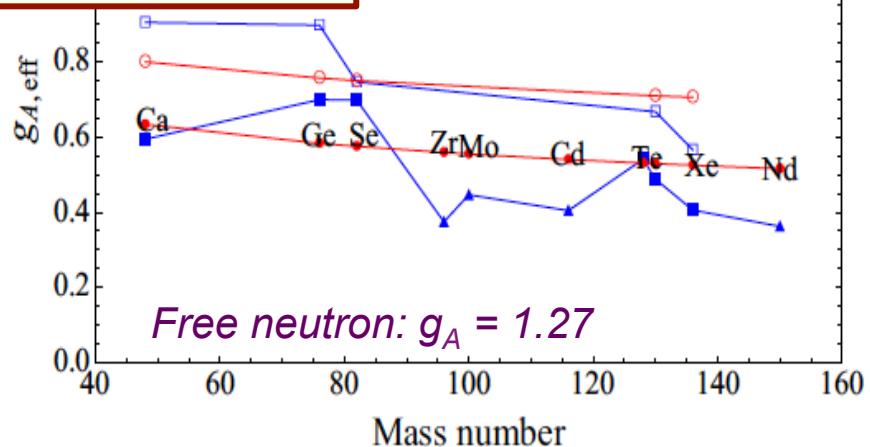
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$$\begin{aligned} \text{experimental } \tau_{1/2} \text{ (ISM)} &= 1.269 A^{-0.12} \\ \text{m experimental } \tau_{1/2} \text{ (IBM-2 CA/SSD)} &= 1.269 A^{-0.18} \end{aligned}$$



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III. TeV Scale LNV

LNV Mass Scale & $0\nu\beta\beta$ -Decay

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Two parameters: Effective coupling & effective heavy particle mass

$0\nu\beta\beta$ -Decay: LNV? Mass Term?

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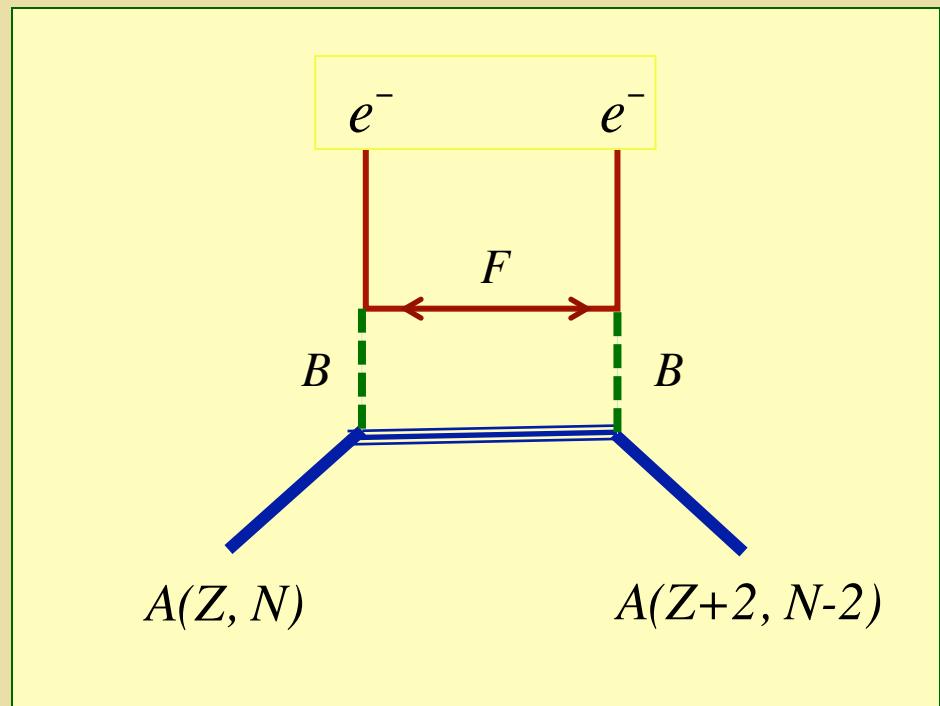
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Majorana

TeV LNV Mechanism

- Majorana mass generated at the TeV scale
 - Low-scale see-saw
 - Radiative m_ν
- $m_{\text{MIN}} \ll 0.01 \text{ eV}$ but $0\nu\beta\beta$ -signal accessible with tonne-scale exp'ts due to heavy Majorana particle exchange



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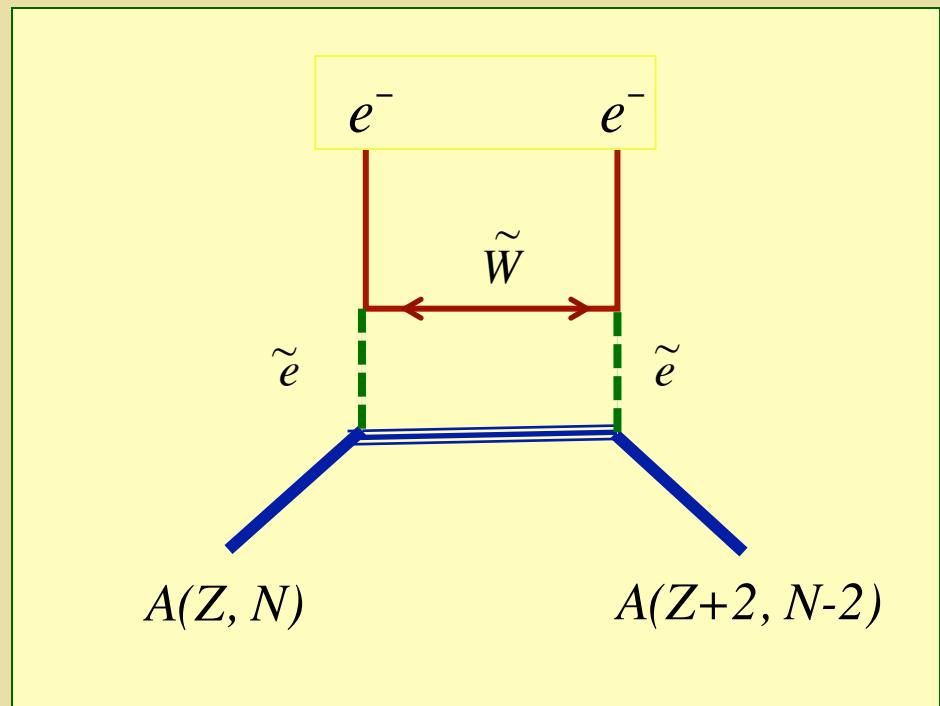
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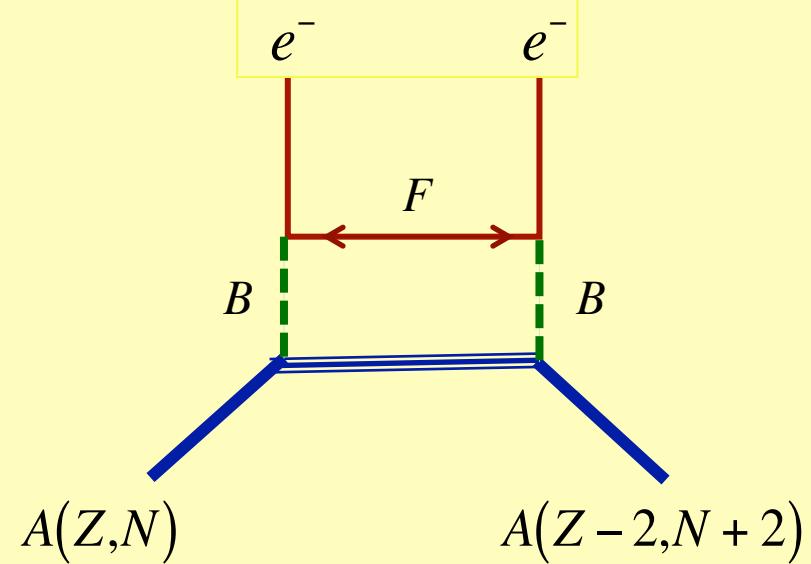
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TeV LNV Mechanism

$$\frac{A_H}{A_L} \sim \frac{M_W^4 \bar{k}^2}{\Lambda^5 m_{\beta\beta}}$$

$O(1)$ for $\Lambda \sim 1$ TeV



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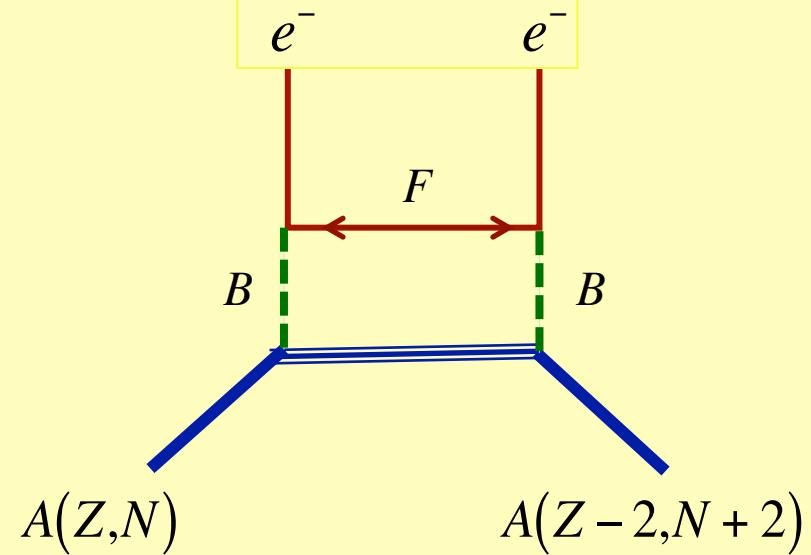
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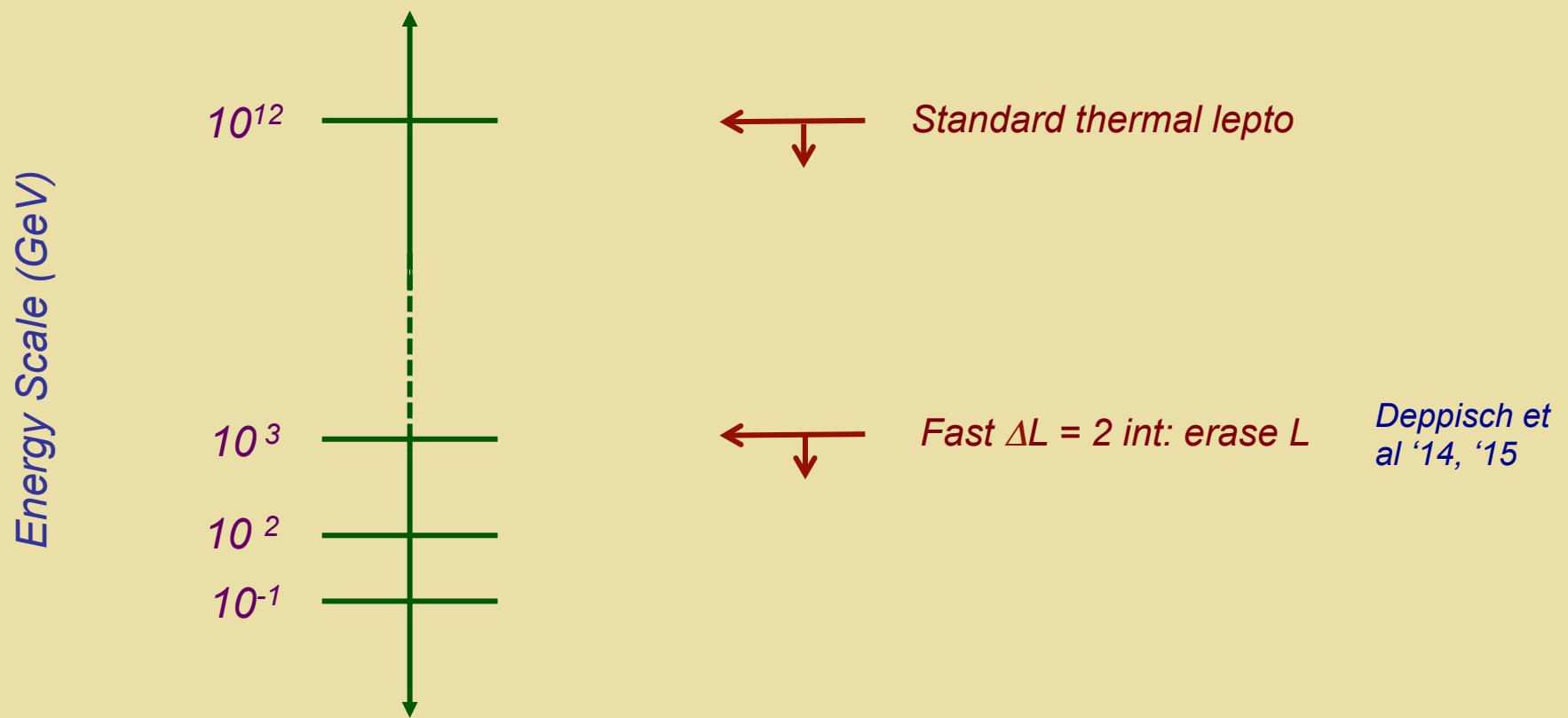
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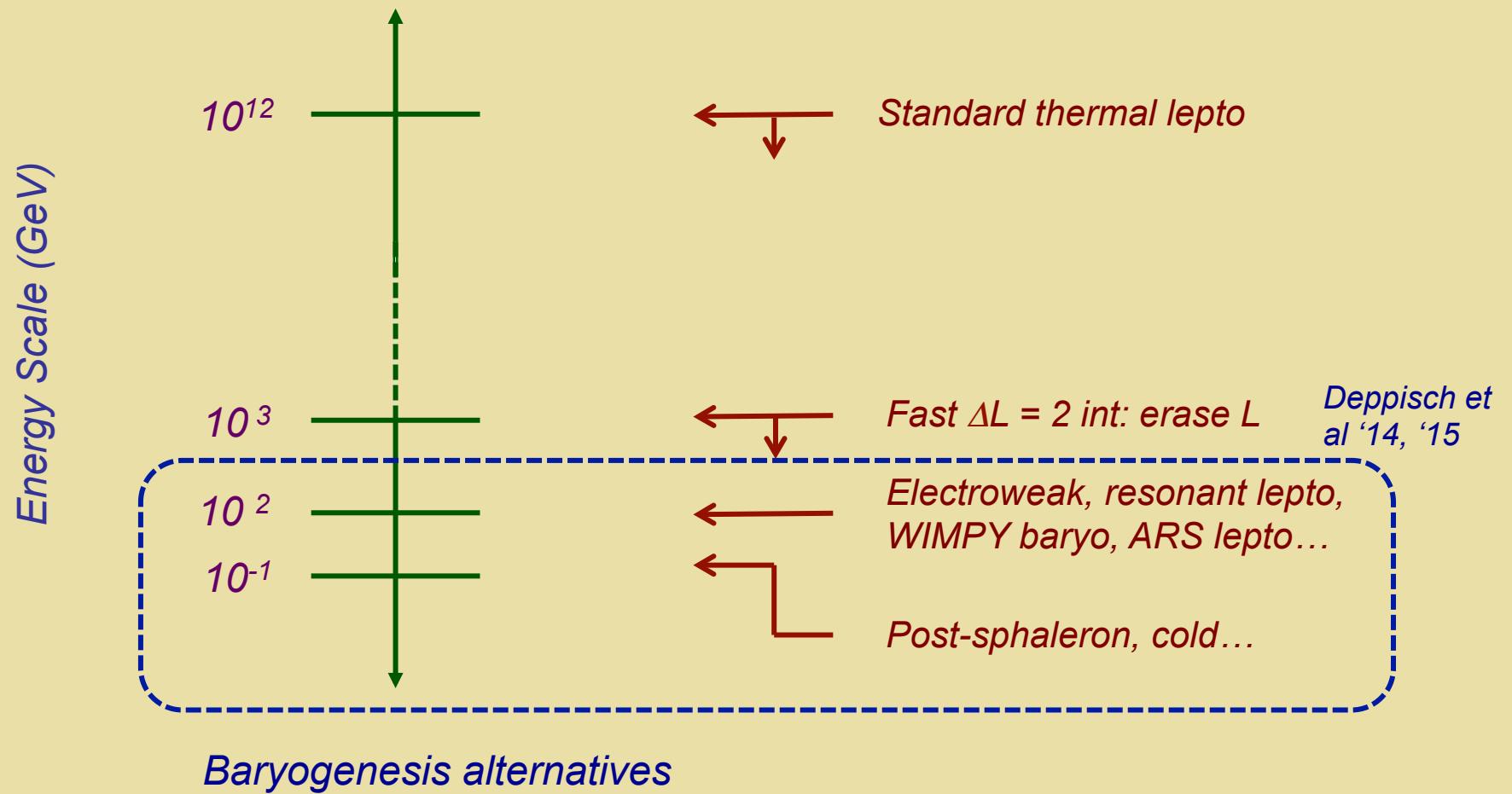
Implications



TeV LNV & Leptogenesis



TeV LNV & Leptogenesis



$0\nu\beta\beta$ -Decay: TeV Scale LNV

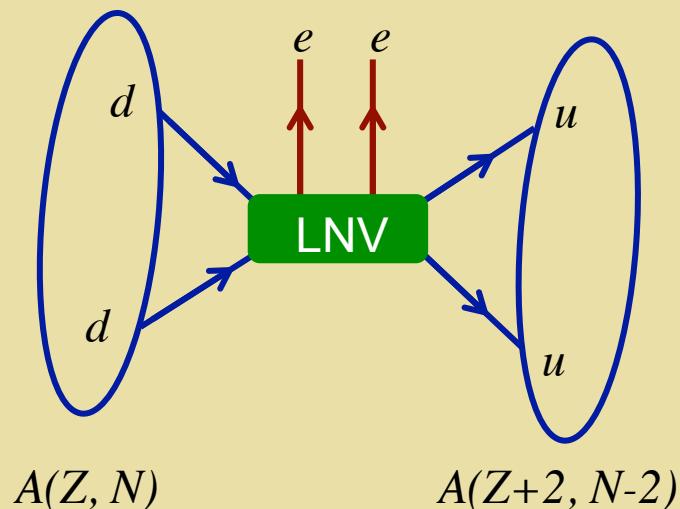
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

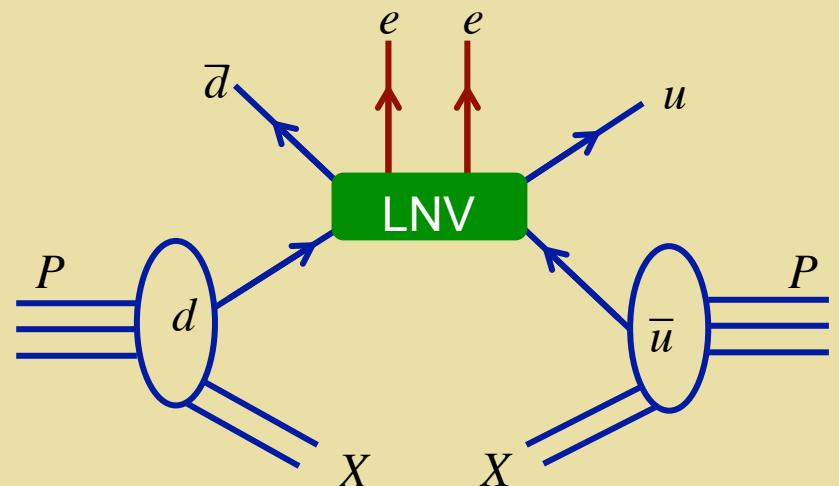
$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

$0\nu\beta\beta$ -Decay



pp Collisions



$0\nu\beta\beta$ -Decay: TeV Scale LNV

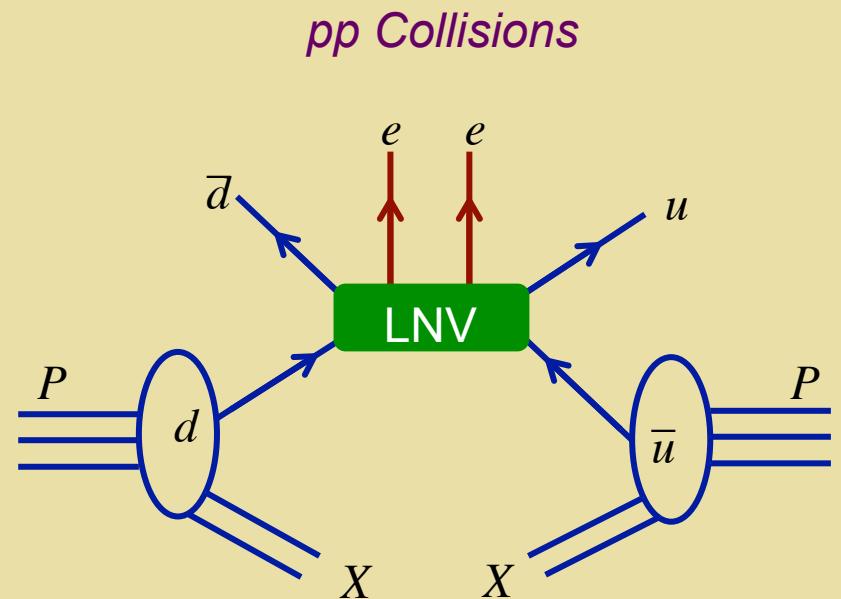
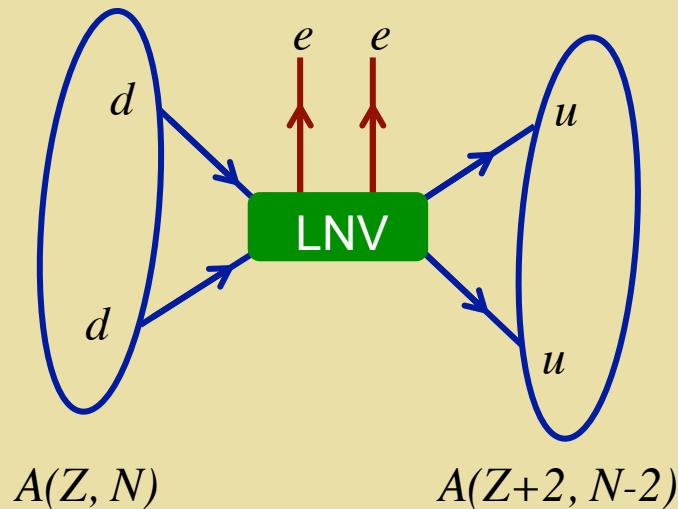
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

$$\mathcal{L}_{\text{mass}} = \frac{y_{\text{Dirac}}}{\Lambda} \bar{L} H H^T L + \text{h.c.}$$

Majorana

Dirac

$0\nu\beta\beta$ -Decay



$0\nu\beta\beta$ -Decay: TeV Scale LNV

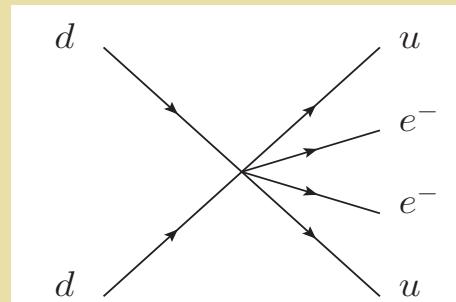
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

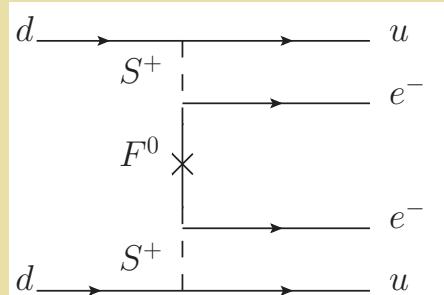
$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

$0\nu\beta\beta$ - decay



LHC: $pp \rightarrow jj e^-e^-$



TeV Scale LNV

Can it be discovered
with combination of
 $0\nu\beta\beta$ & LHC searches ?

Simplified models

$0\nu\beta\beta$ -Decay: TeV Scale LNV

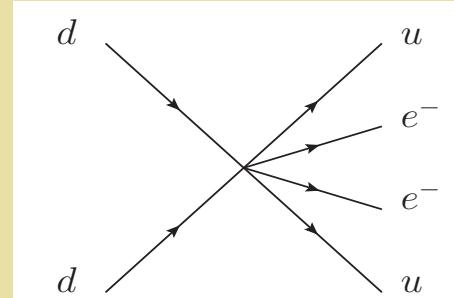
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

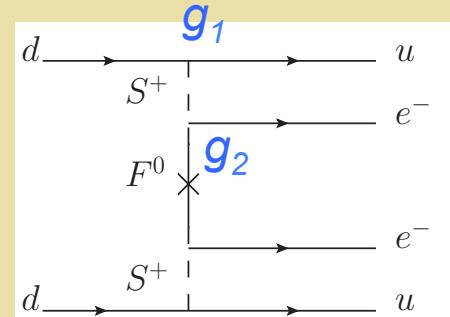
$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

$0\nu\beta\beta$ - decay



LHC: $pp \rightarrow jj e^-e^-$



TeV Scale LNV

Effective operators:

$$\mathcal{L}_{\text{LNV}}^{\text{eff}} = \frac{C_1}{\Lambda^5} \mathcal{O}_1 + \text{h.c.}$$

$$\mathcal{O}_1 = \bar{Q} \tau^+ d \bar{Q} \tau^+ d \bar{L} L^C$$

$$g_{\text{eff}} = \sqrt{g_1 g_2}$$

0νββ-Decay: Rate & Mass Dependence

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Light ν exchange

$$\frac{1}{T_{1/2}} = G^{0\nu}(E, Z) |M_{0\nu}| |\langle m_{\beta\beta} \rangle|^2$$

Quadratic dependence on $m_{\beta\beta}$

Heavy particle exchange

$$\begin{aligned} \frac{1}{T_{1/2}} = & G_{01} \left(\frac{\text{TeV}}{m_e} \right)^2 \left(\frac{\Lambda_H}{\text{TeV}} \right)^4 \left(\frac{1}{18} \right) \left(\frac{v}{\text{TeV}} \right)^8 \\ & \times \left(\frac{1}{g_A \cos \theta_C} \right)^4 |M_0|^2 \left[\frac{C_{\text{eff}}^2}{(\Lambda/\text{TeV})^{10}} \right], \end{aligned}$$

Scales as $1/M^{10}$

$0\nu\beta\beta$ -Decay: TeV Scale LNV

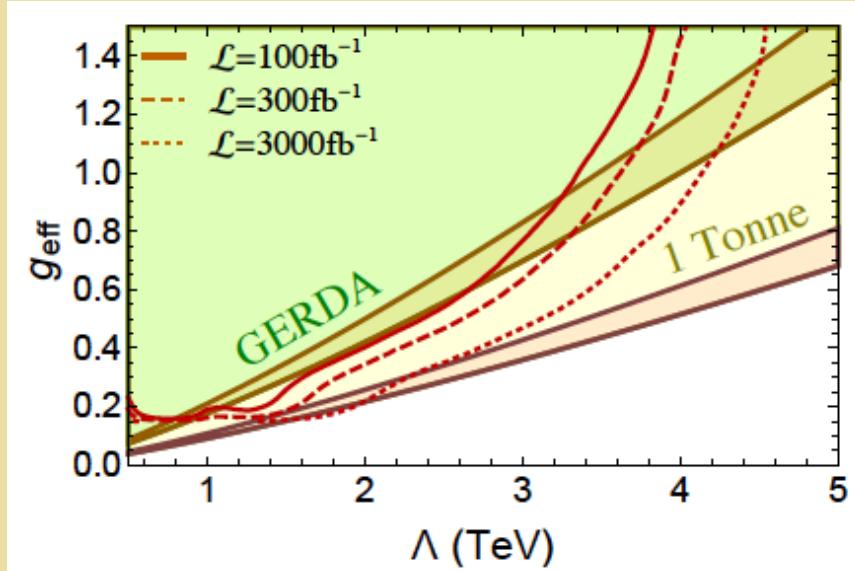
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

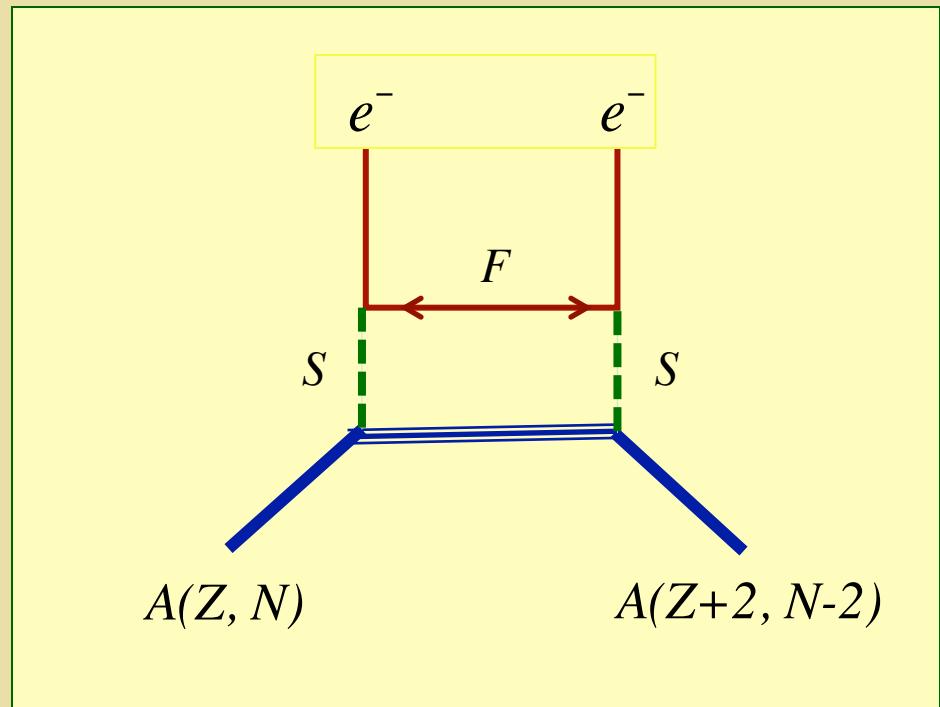
$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Benchmark Sensitivity: TeV LNV



T. Peng, MRM, P. Winslow 1508.04444



$0\nu\beta\beta$ -Decay: TeV Scale LNV

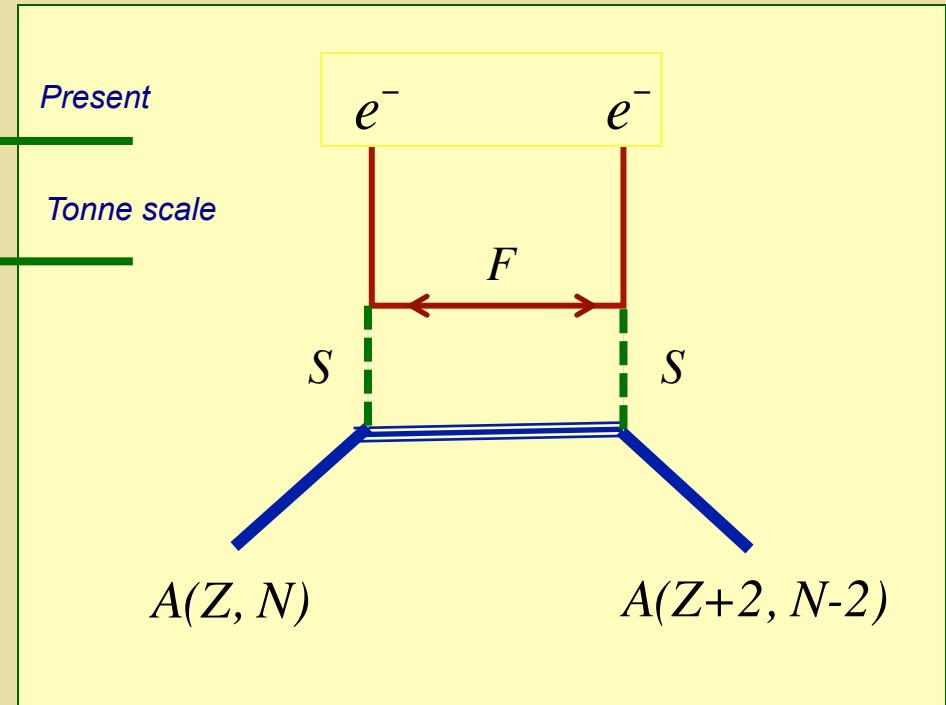
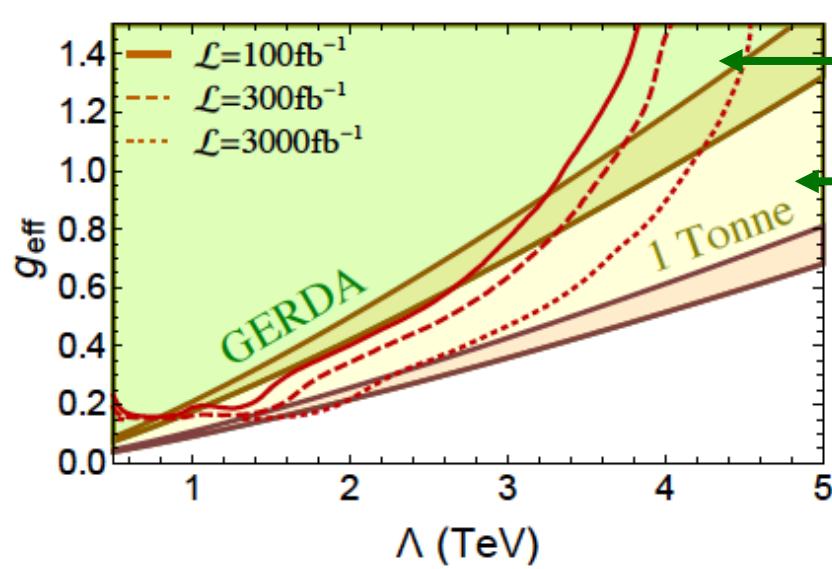
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Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Benchmark Sensitivity: TeV LNV



$0\nu\beta\beta$ -Decay: TeV Scale LNV

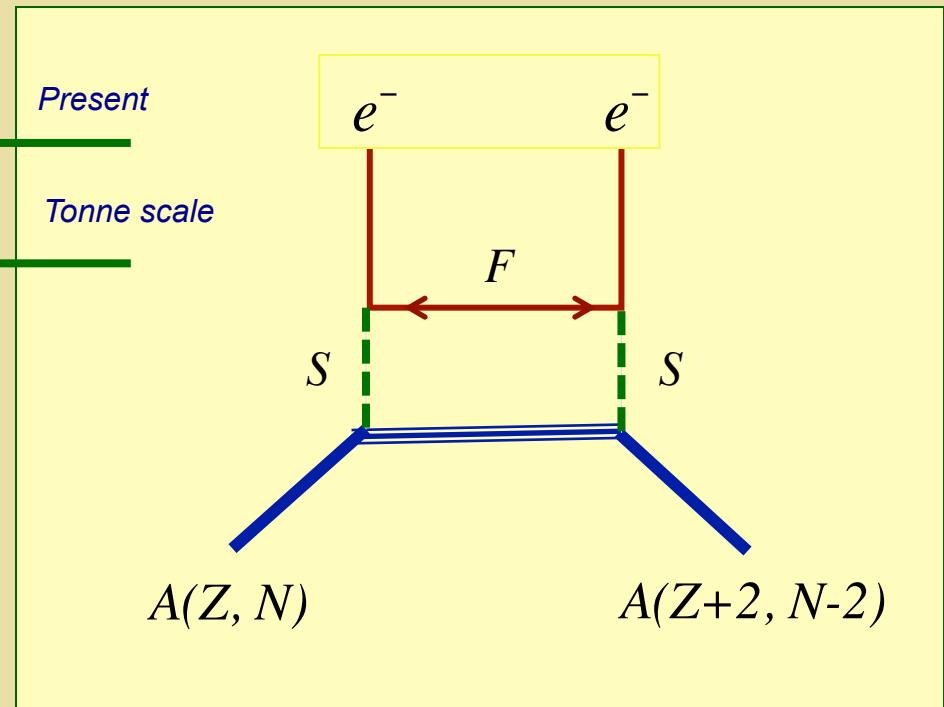
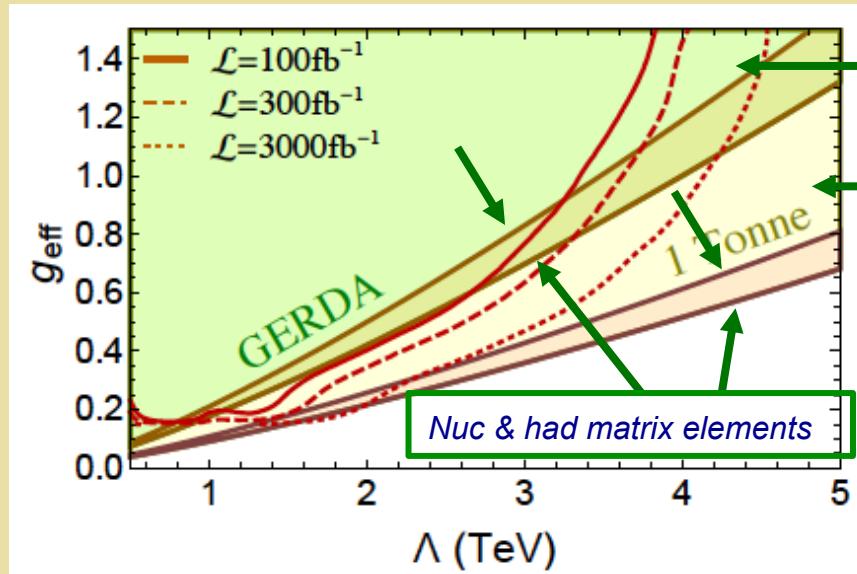
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Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Benchmark Sensitivity: TeV LNV



$0\nu\beta\beta$ -Decay: TeV Scale LNV

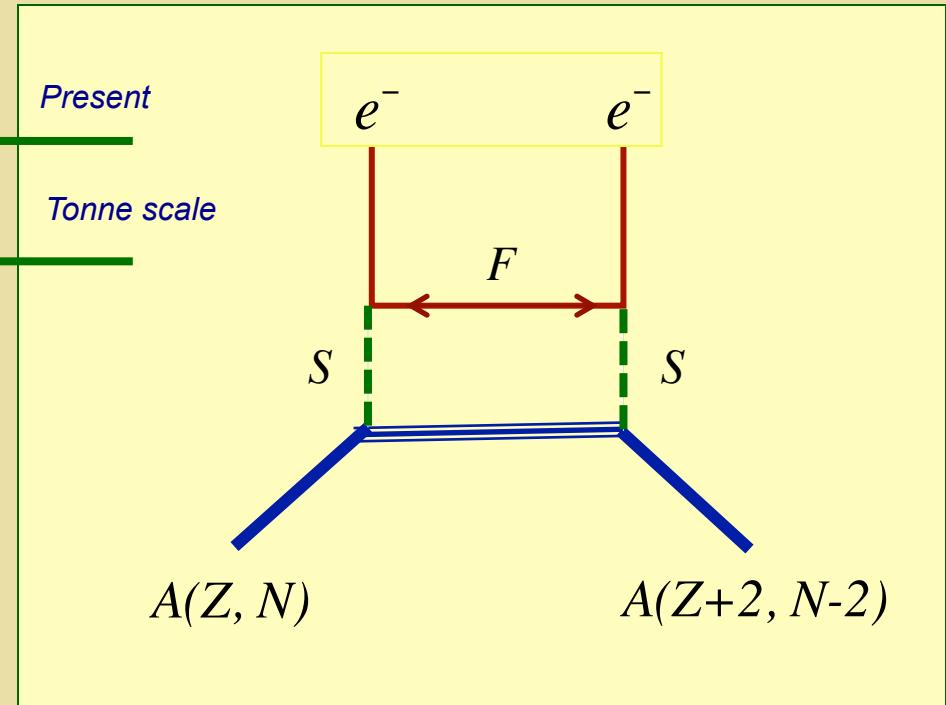
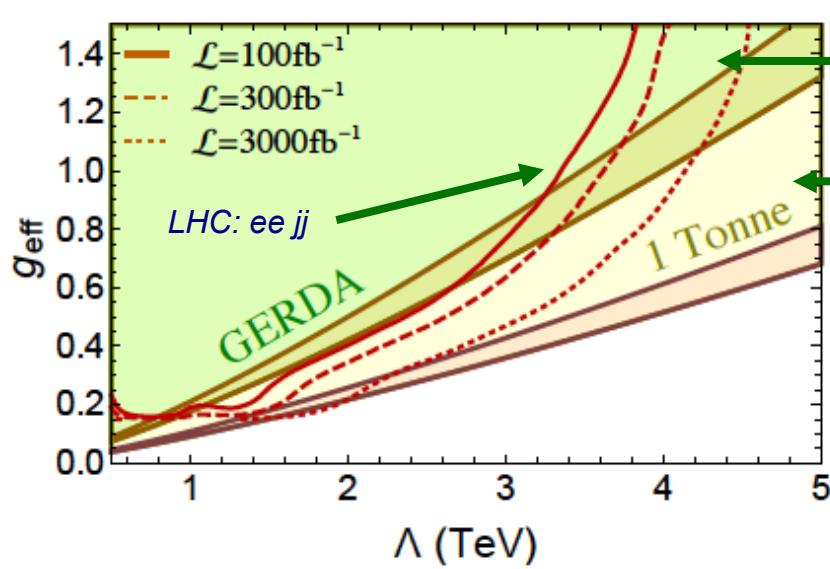
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Benchmark Sensitivity: TeV LNV



$0\nu\beta\beta$ -Decay: TeV Scale LNV & m_ν

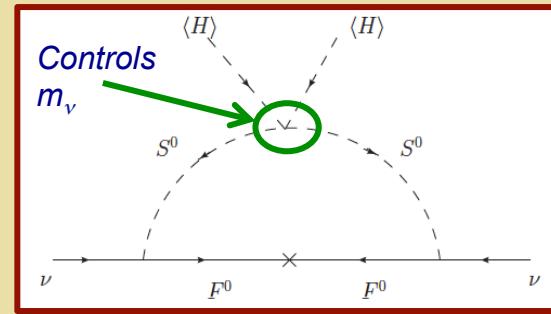
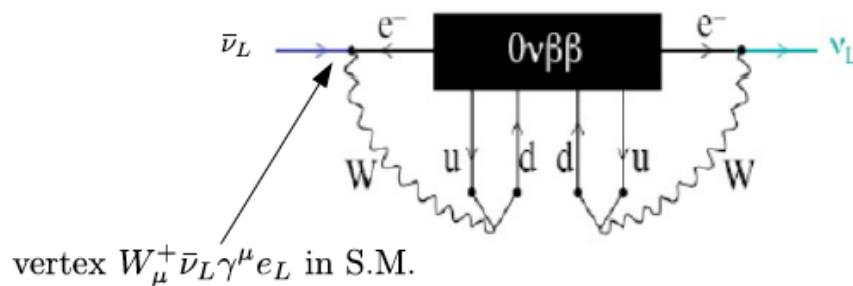
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Implications for m_ν :



Schechter-Valle: non-vanishing Majorana mass at (multi) loop level

Simplified model: possible (larger) one loop Majorana mass

$0\nu\beta\beta$ -Decay: TeV Scale LNV & m_ν

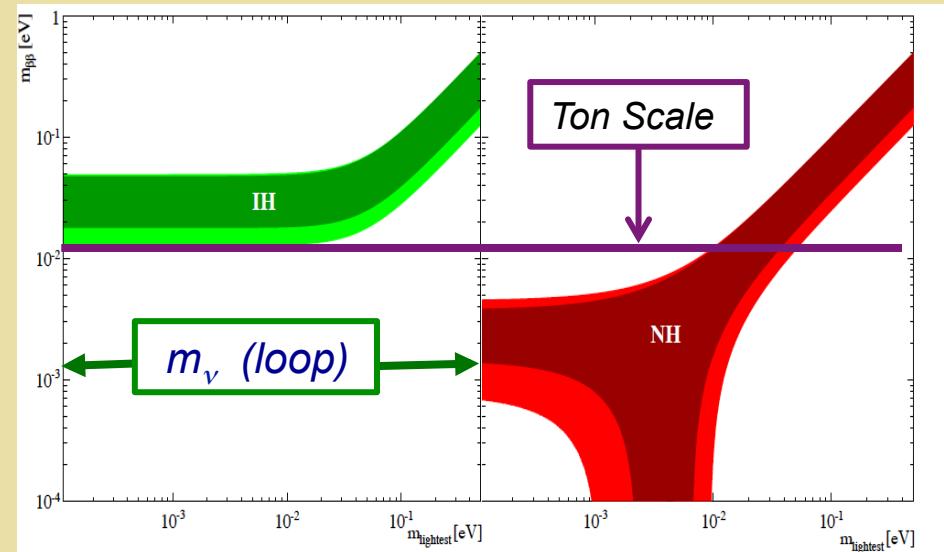
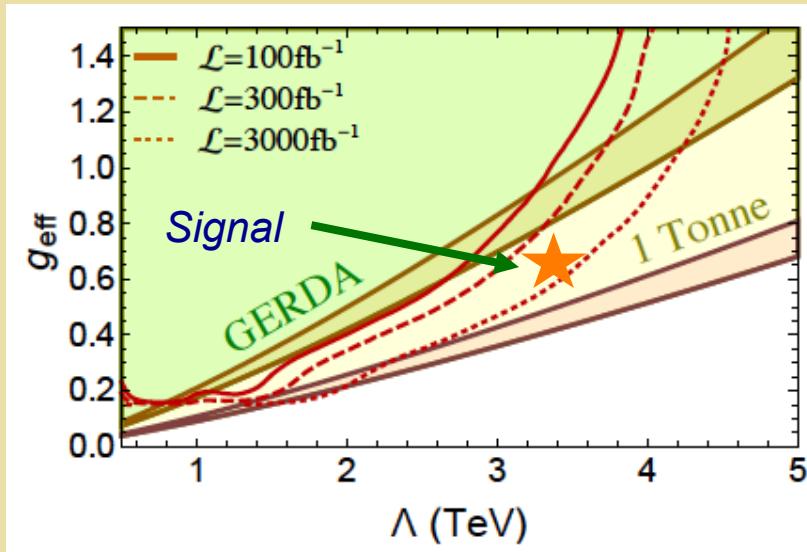
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Dirac

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Majorana

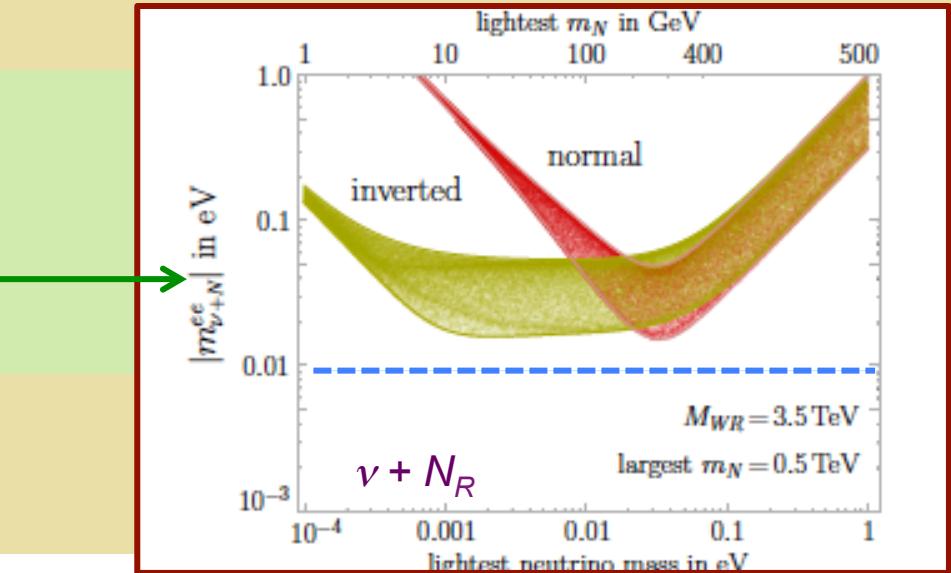
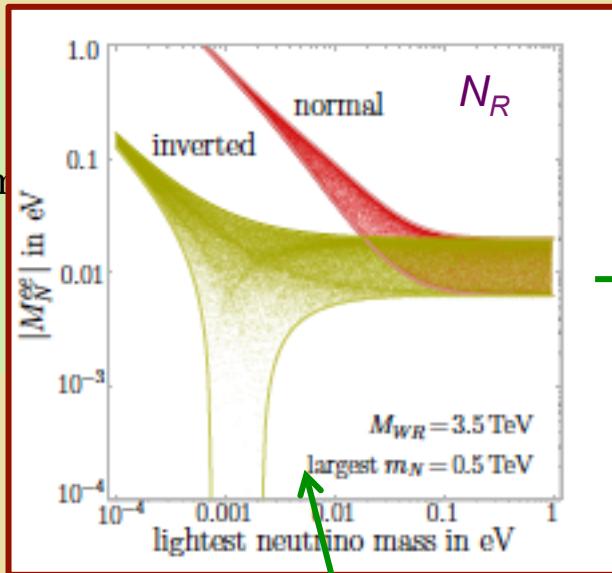
Implications for m_ν :



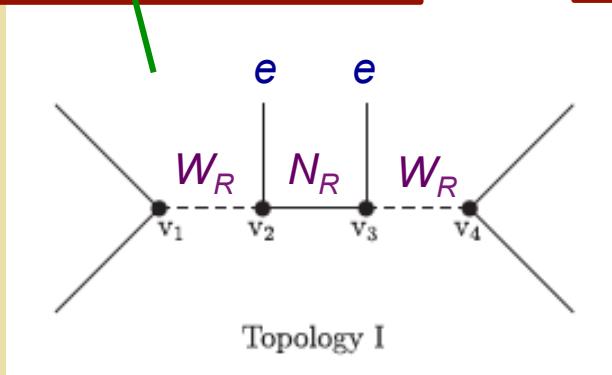
A hypothetical scenario

$0\nu\beta\beta$ -Decay: TeV Scale LNV

\mathcal{L}_n



C.



LRSM: Type I See-Saw

Mass: standard see-saw but TeV scale

Tello et al, 1011.3522

IV. Sub-Weak Scale LNV

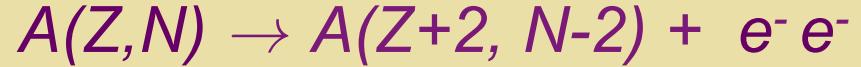
LNV Mass Scale & $0\nu\beta\beta$ -Decay

$$A(Z, N) \rightarrow \text{Underlying Physics} \rightarrow A(Z+2, N-2) + e^- e^-$$

- *3 light neutrinos only: source of neutrino mass at the very high see-saw scale*
- *3 light neutrinos with TeV scale source of neutrino mass*
- *> 3 light neutrinos*

$0\nu\beta\beta$ -Decay Sensitivity

2ν DBD:



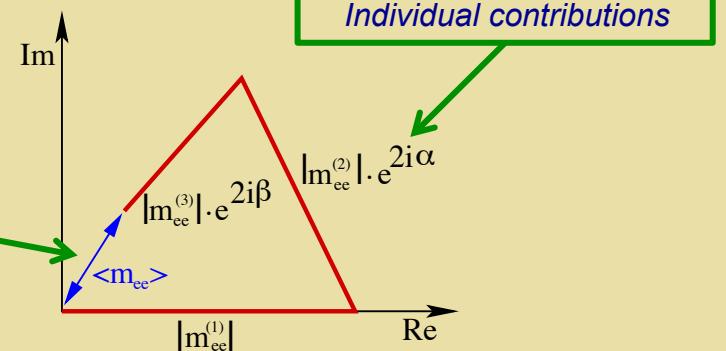
$\bar{\nu} \bar{\nu}$

If own antiparticle, can be emitted
then absorbed during decay

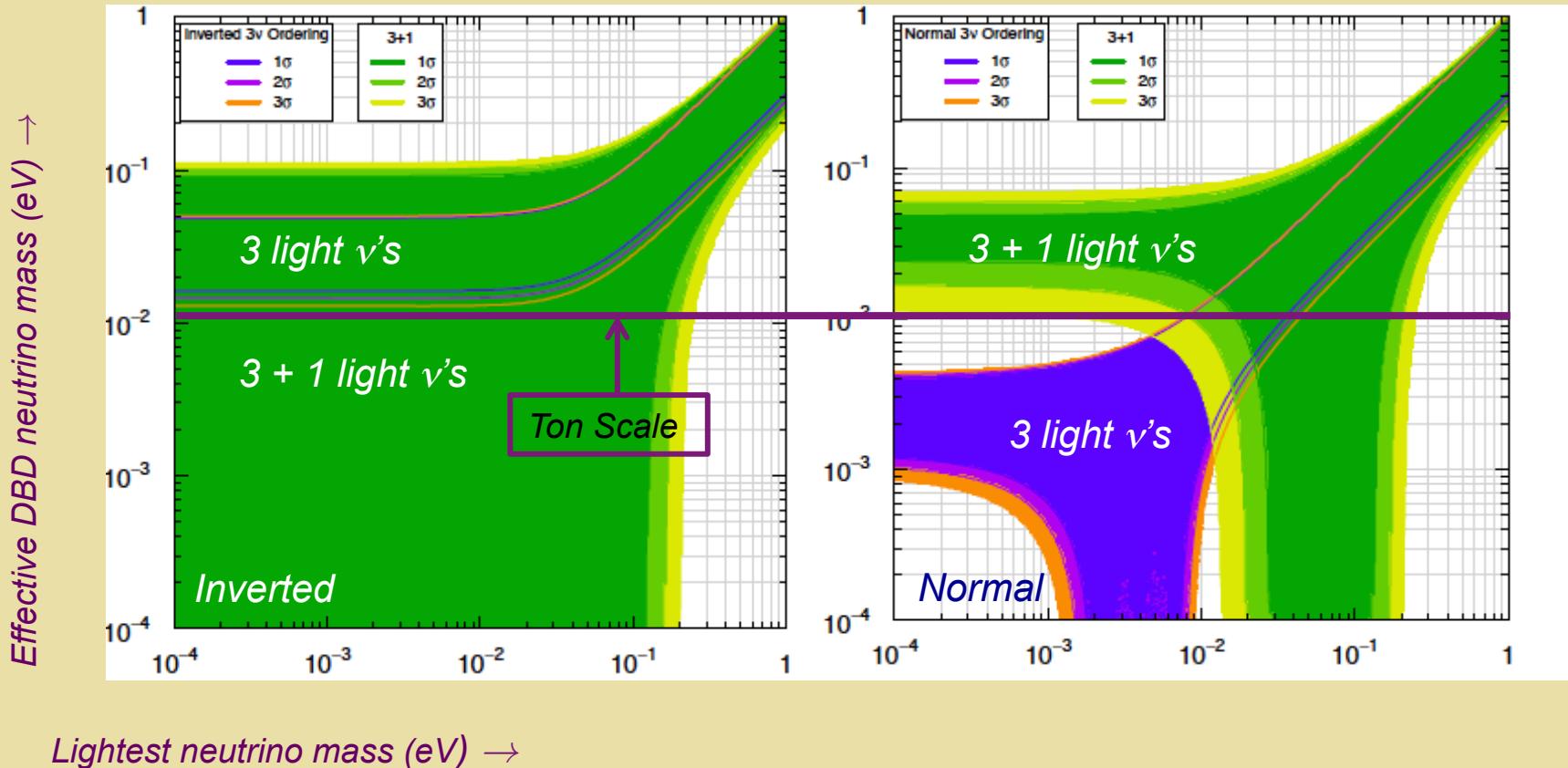
0ν DBD:



All three light neutrinos participate →
Rate governed by an **effective mass**

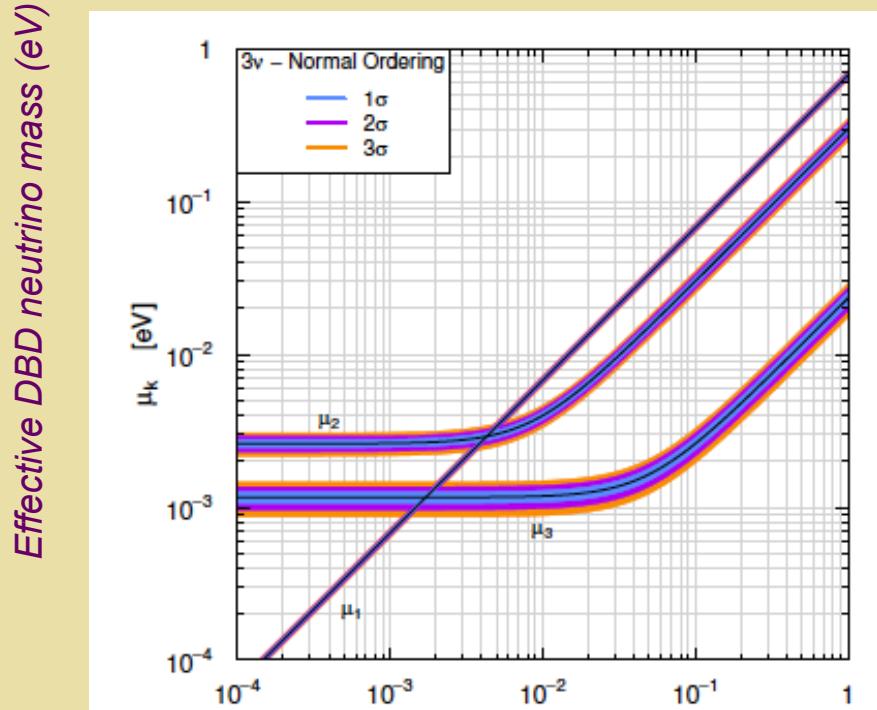


> 3 Light Neutrinos



Sterile Neutrinos & $0\nu\beta\beta$ -Decay

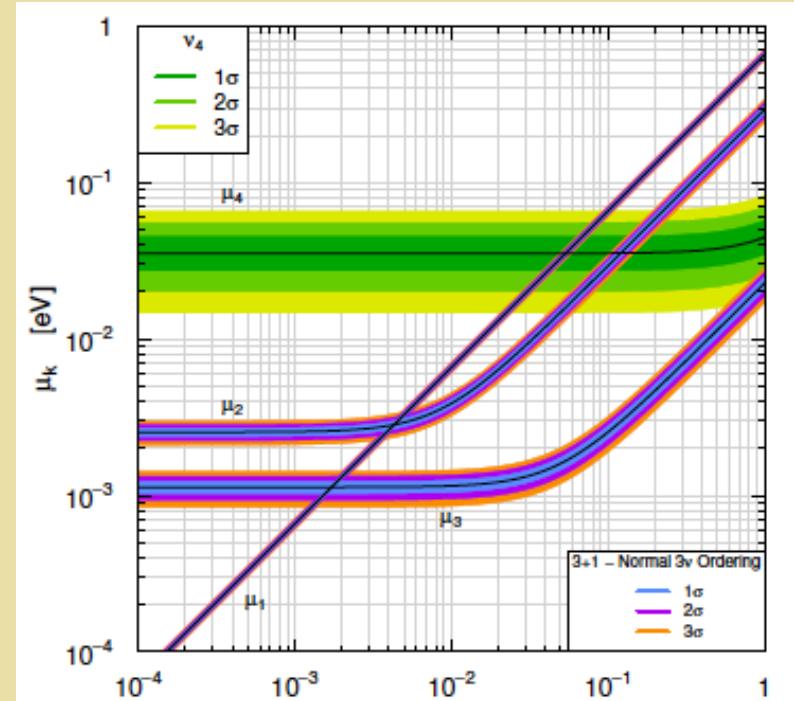
3 active light neutrinos



Lightest neutrino mass (eV) →

$$|m_{\beta\beta}| = |\mu_1 + \mu_2 e^{i\alpha_2} + \mu_3 e^{i\alpha_3}|$$

3+1 active light neutrinos

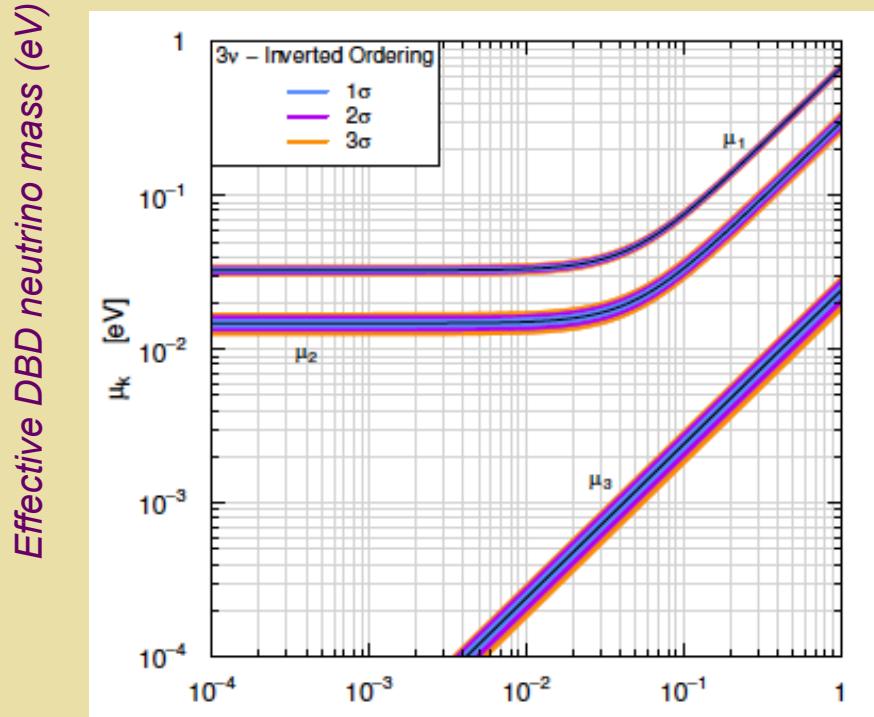


Lightest neutrino mass (eV) →

$$|m_{\beta\beta}| = |\mu_1 + \mu_2 e^{i\alpha_2} + \mu_3 e^{i\alpha_3} + \mu_4 e^{i\alpha_4}|$$

Sterile Neutrinos & $0\nu\beta\beta$ -Decay

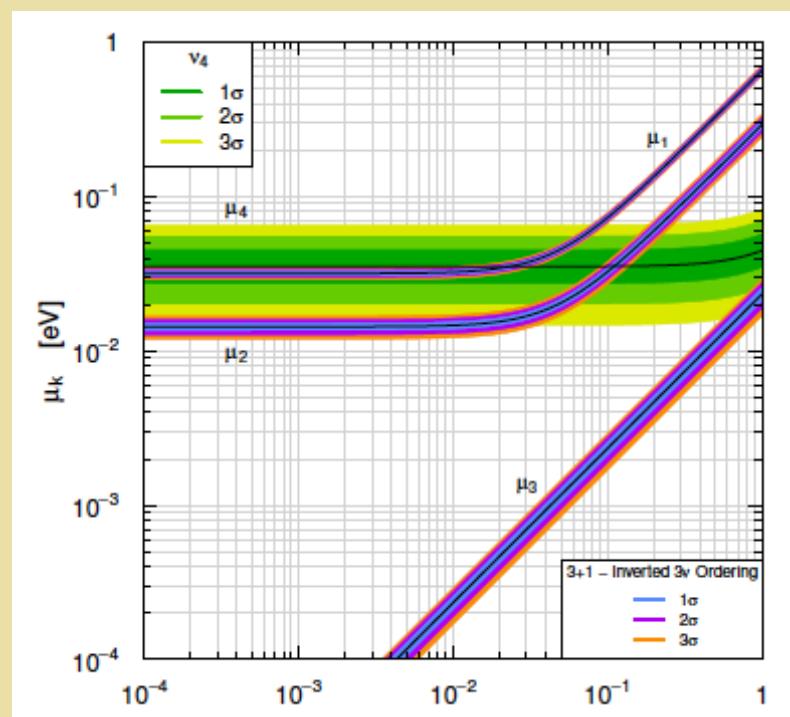
3 active light neutrinos



Lightest neutrino mass (eV) →

$$|m_{\beta\beta}| = |\mu_1 + \mu_2 e^{i\alpha_2} + \mu_3 e^{i\alpha_3}|$$

3+1 active light neutrinos



Lightest neutrino mass (eV) →

$$|m_{\beta\beta}| = |\mu_1 + \mu_2 e^{i\alpha_2} + \mu_3 e^{i\alpha_3} + \mu_4 e^{i\alpha_4}|$$

V. Discussion Questions

- *What is a sphaleron ?*
- *Is the CPV in V_{PMNS} the same as CPV for leptogenesis ?*
- *What is the conventional leptogenesis scale ?*
- *What is the Schecter-Valle (black box) theorem ?*