

Lectures I : Neutrino Theory Basics

M.J. Ramsey-Musolf

U Mass Amherst



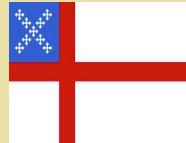
AMHERST CENTER FOR FUNDAMENTAL INTERACTIONS

Physics at the interface: Energy, Intensity, and Cosmic frontiers

University of Massachusetts Amherst

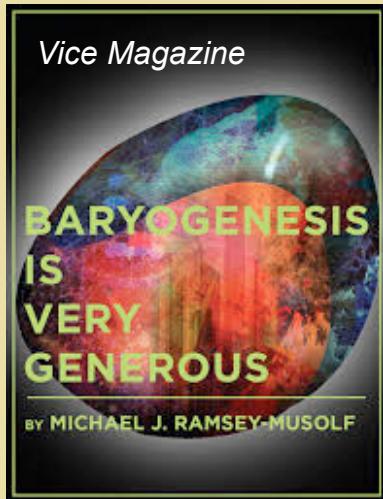
<http://www.physics.umass.edu/acfi/>

NNPSS, Wright Laboratory
Yale 6/18-29/18



About MJRM

Theoretical Physics



- *Why does the Universe contain more matter than antimatter ?*
- *What are the laws of nature beyond those of the Standard Model & General Relativity ?*
- *How do quantum field theories work ?*

My pronouns: he/him/his

Fundamental Symmetries & Neutrinos

<p><i>EDM searches:</i> <i>BSM CPV, Origin of Matter</i></p>	<p><i>0νββ decay searches:</i> <i>Nature of neutrino, Lepton number violation, Origin of Matter</i></p>
<p><i>Electron & muon prop's & interactions:</i> <i>SM Precision Tests, BSM "diagnostic" probes</i></p>	<p><i>Radioactive decays & other tests</i> <i>SM Precision Tests, BSM "diagnostic" probes</i></p>

Fundamental Symmetries & Neutrinos

EDM searches:

BSM CPV, Origin of Matter

Lecture III

$0\nu\beta\beta$ decay searches:

Nature of neutrino, Lepton number violation, Origin of Matter

Lectures I & II

Electron & muon prop's & interactions:

SM Precision Tests, BSM “diagnostic” probes

Radioactive decays & other tests

SM Precision Tests, BSM “diagnostic” probes

Lecture I Goals

- *Review the basic theoretical formulation of neutrino oscillation phenomenology*
- *Review some of the open questions in neutrino physics*
- *Provide a simple overview of classes of neutrino mass models with example illustrations*
- *Invite questions !*

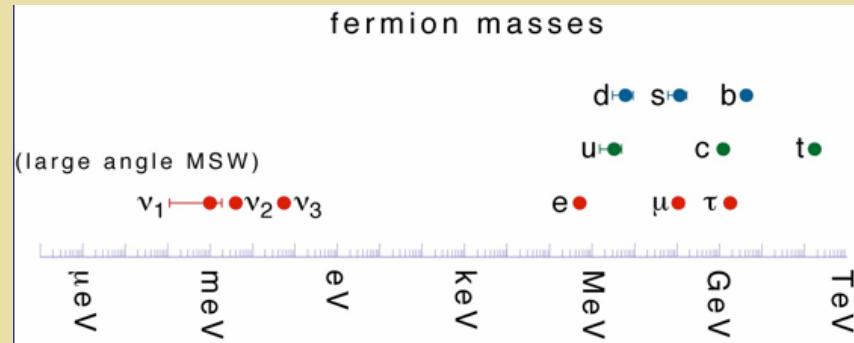
Lecture I Outline

- I. *Overview*
- II. *Neutrino oscillations imply non-zero m_ν*
- III. *Open questions*
- IV. *Neutrino Mass Models*
- V. *Discussion questions*

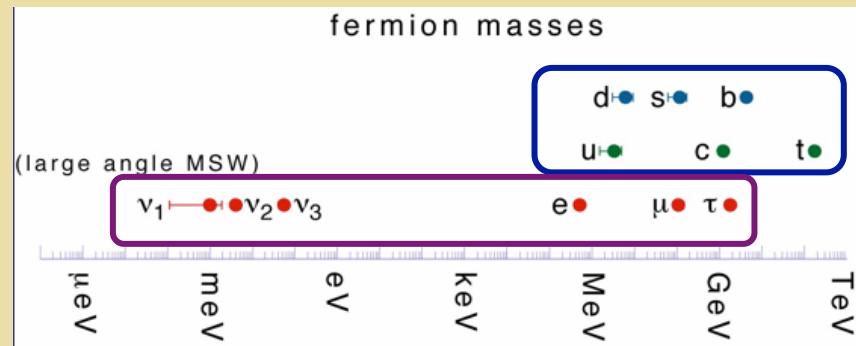
I. Overview

*Theoretical complement to D. Parno's
excellent experimental overview*

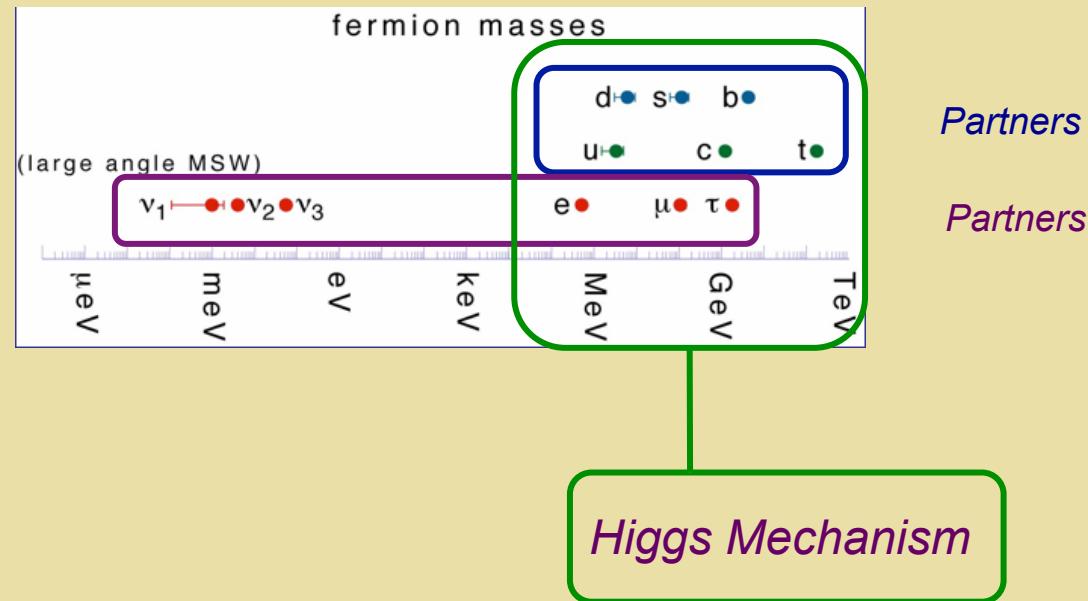
Neutrino Masses



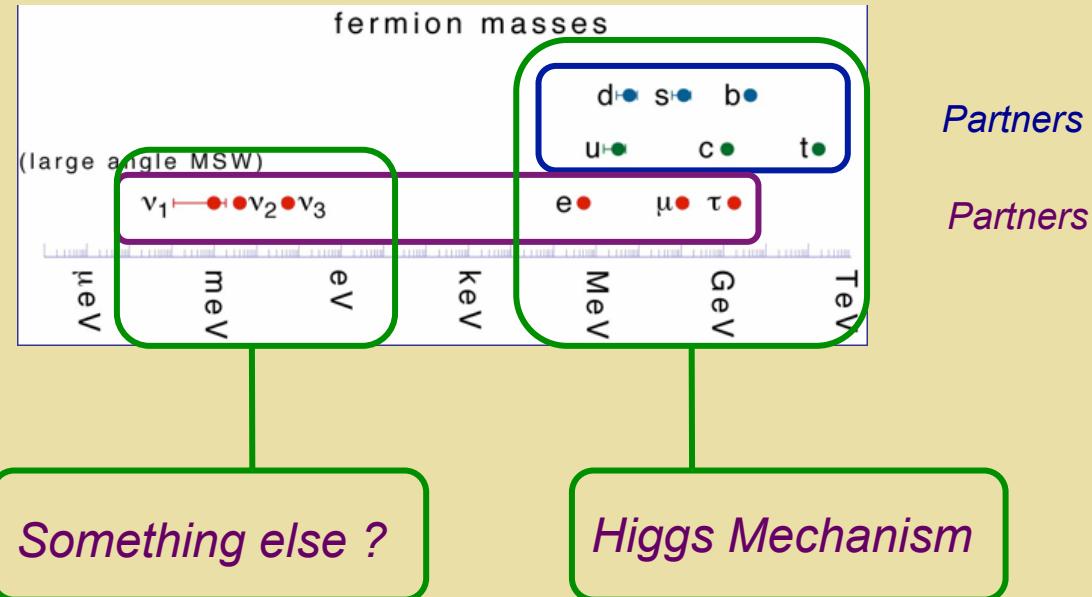
Neutrino Masses



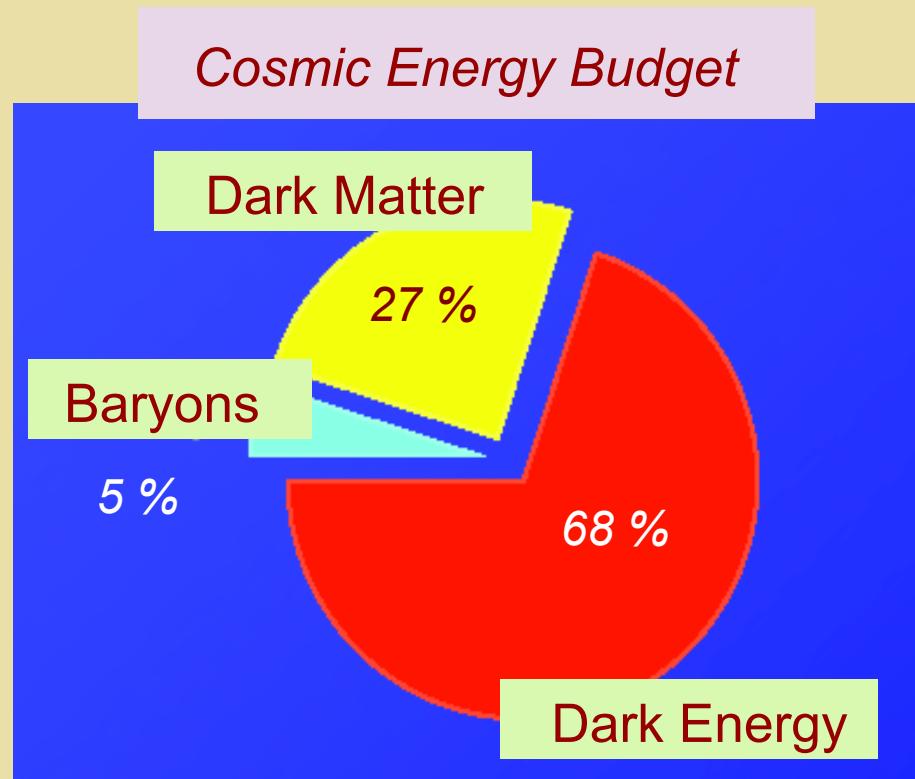
Neutrino Masses



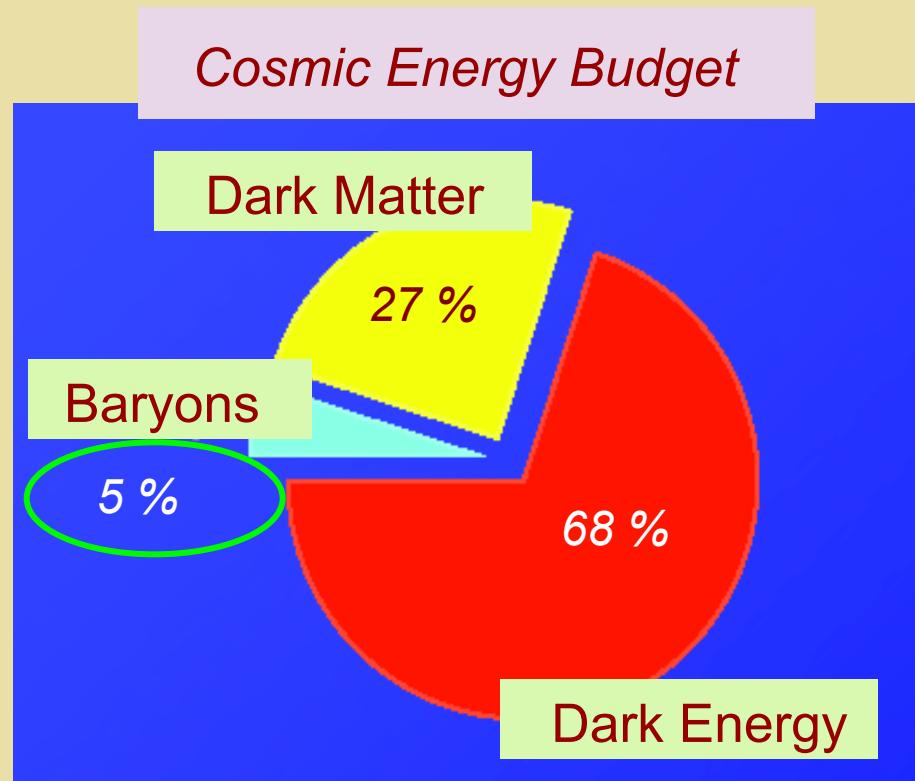
Neutrino Masses



The Origin of Matter



The Origin of Matter



Ingredients for Baryogenesis



- *B violation (sphalerons)*
- *C & CP violation*
- *Out-of-equilibrium or CPT violation*

Ingredients for Baryogenesis



	<i>Standard Model</i>	<i>BSM</i>
• <i>B violation (sphalerons)</i>	✓	✓
• <i>C & CP violation</i>	✗	✓
• <i>Out-of-equilibrium or CPT violation</i>	✗	✓

Ingredients for Baryogenesis



Scenarios: *leptogenesis, EW baryogenesis, Affleck-Dine, asymmetric DM, cold baryogenesis, post-sphaleron baryogenesis...*

	<i>Standard Model</i>	<i>BSM</i>
• <i>B violation (sphalerons)</i>	✓	✓
• <i>C & CP violation</i>	✗	✓
• <i>Out-of-equilibrium or CPT violation</i>	✗	✓

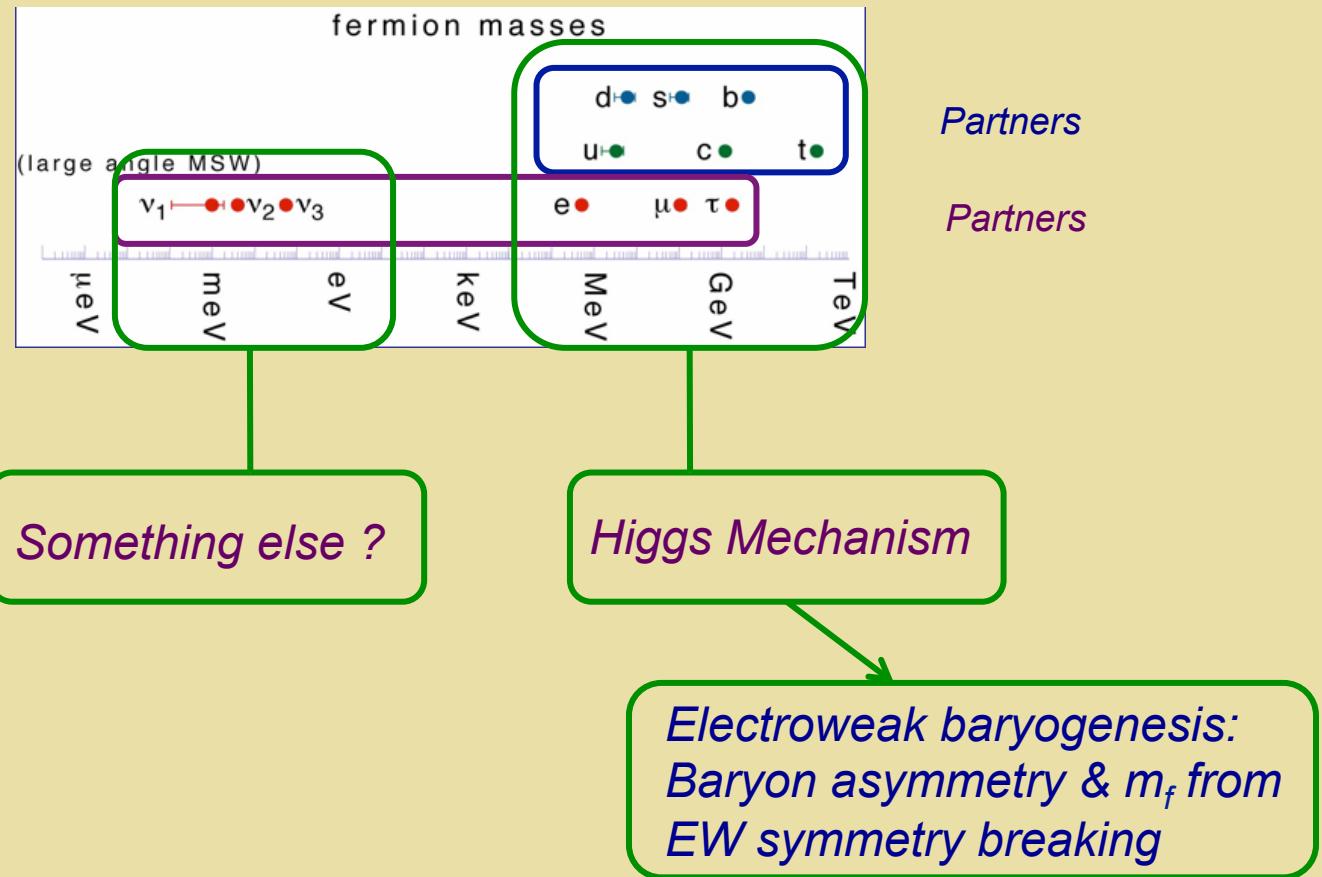
Ingredients for Baryogenesis



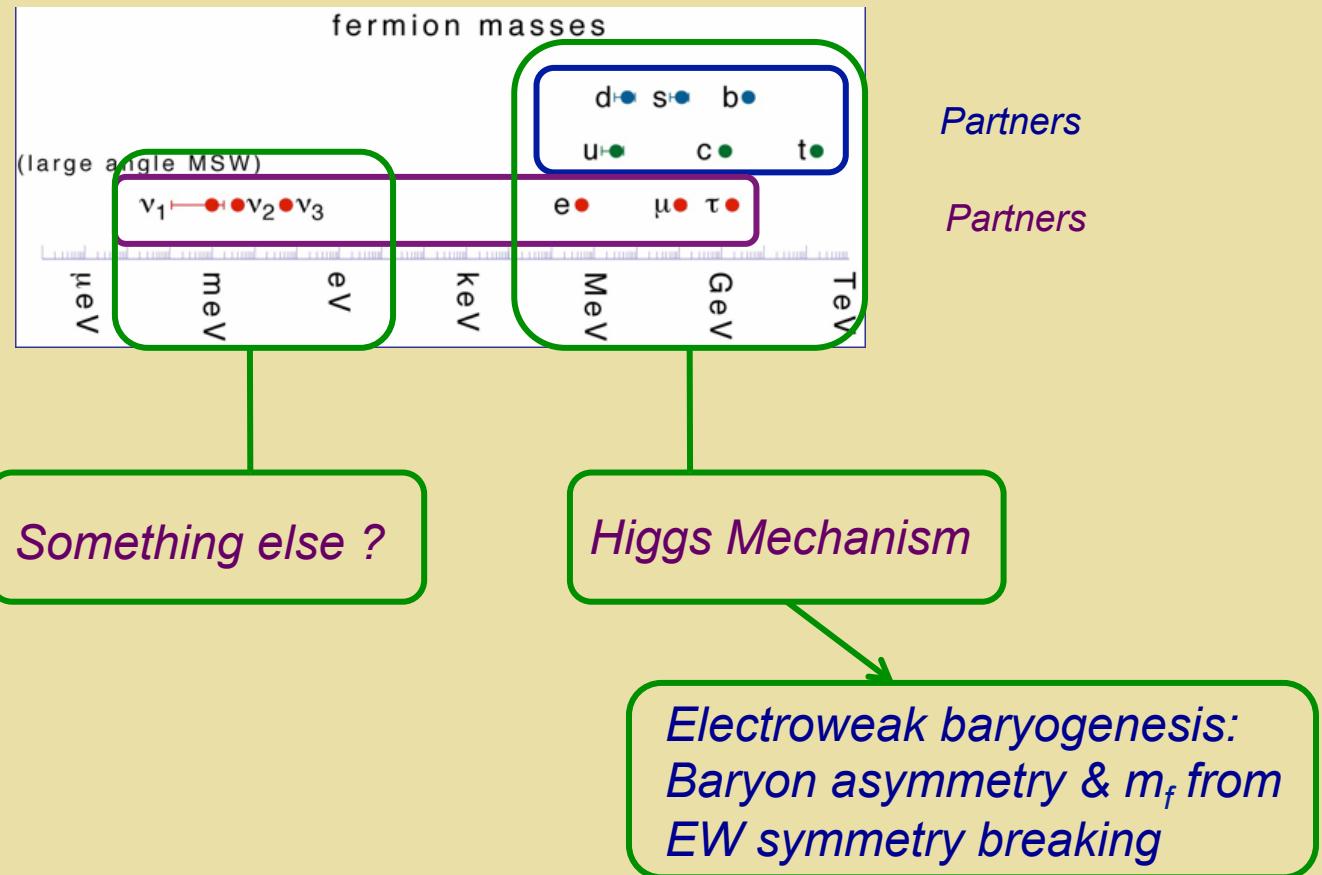
Scenarios: leptogenesis,
EW baryogenesis, Afflek-
Dine, asymmetric DM, cold
baryogenesis, post-
sphaleron baryogenesis...

	Standard Model	BSM
• <i>B violation (sphalerons)</i>	✓	✓
• <i>C & CP violation</i>	✗	✓
• <i>Out-of-equilibrium or CPT violation</i>	✗	✓

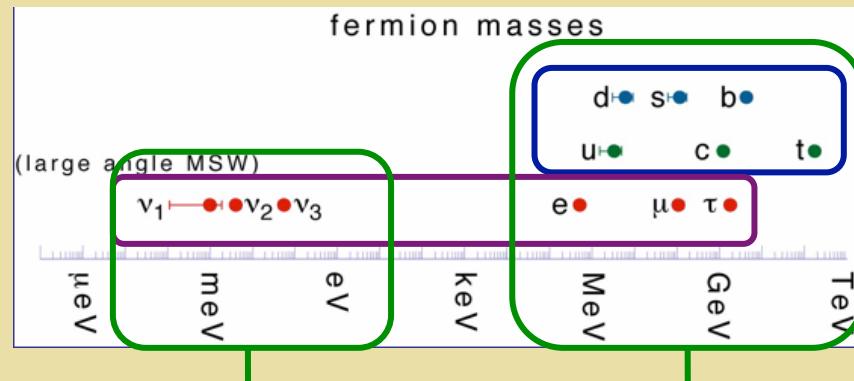
Fermion Masses & Baryon Asymmetry



Fermion Masses & Baryon Asymmetry



Fermion Masses & Baryon Asymmetry



Partners

Partners

Something else ?

Higgs Mechanism

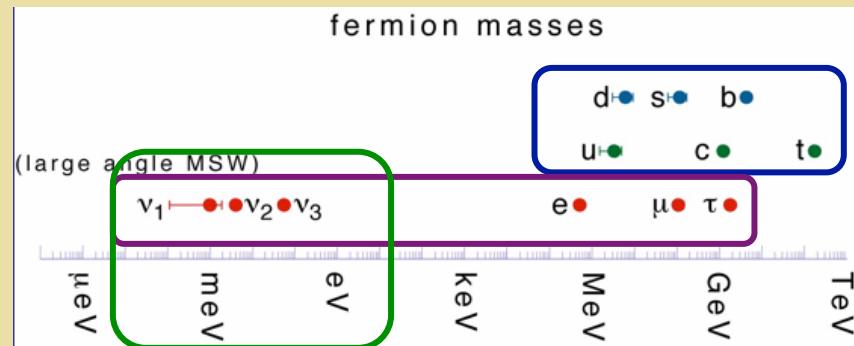
Leptogenesis: Baryon asymmetry & m_ν from lepton number violation

Electroweak baryogenesis: Baryon asymmetry & m_f from EW symmetry breaking

Lecture II

Lecture III

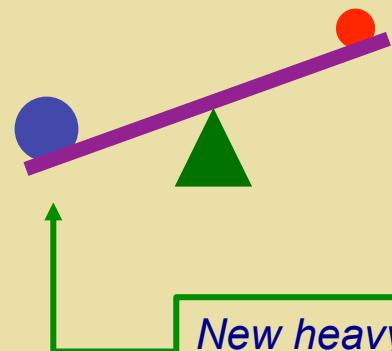
Neutrino Masses



Partners

Partners

“See saw mechanism”



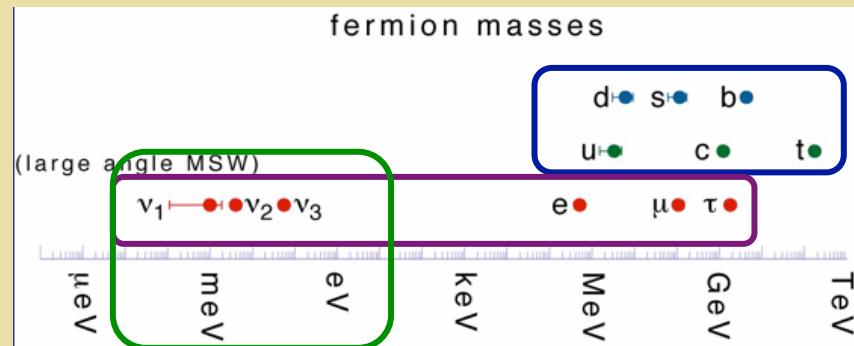
Physical state masses

$$m_1 \approx \frac{m_D^2}{M_N} \quad \sim \text{eV}$$

$$m_2 \approx M_N \quad \sim 10^{12} - 10^{15} \text{ GeV}$$

New heavy neutrino-like particle =
its own anti-particle

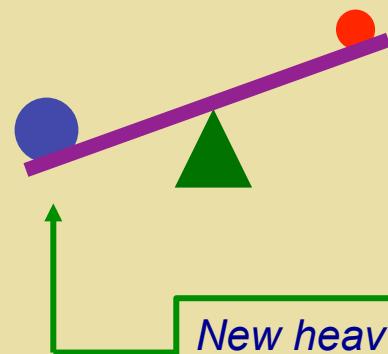
Neutrino Masses



Partners

Partners

“See saw mechanism”



“Leptogenesis”

Heavy neutrino decays in early universe generate baryon asym

New heavy neutrino-like particle =
its own anti-particle

II. Neutrino Oscillations Implies $m_\nu \neq 0$

A. Two Level System

*Flavor (weak interaction)
eigenstates:*

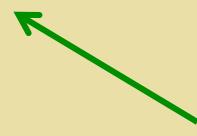
$$|\nu_A\rangle, |\nu_B\rangle$$

Mass eigenstates:

$$|\nu_1\rangle, |\nu_2\rangle$$

Unitary transformation:

$$\begin{pmatrix} |\nu_A\rangle \\ |\nu_B\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_{12} & \sin \theta_{12} \\ -\sin \theta_{12} & \cos \theta_{12} \end{pmatrix} \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \end{pmatrix}$$


$$V$$

A. Two Level System

Initial state, created by weak interaction (e.g. β -decay):

$$|\psi(0)\rangle = |\nu_A\rangle = \cos \theta_{12} |\nu_1\rangle + \sin \theta_{12} |\nu_2\rangle \quad \text{e.g., } \nu_A = \nu_e$$

Time evolution:

$$|\psi(t)\rangle = e^{-iE_1 t} \cos \theta_{12} |\nu_1\rangle + e^{-iE_2 t} \sin \theta_{12} |\nu_2\rangle$$

What is probability of being in state $|\nu_A\rangle$ at time t after creation ?

$$\mathcal{P}(\nu_A \rightarrow \nu_A) \quad \text{“Survival probability”}$$

A. Two Level System

Survival amplitude:

$$\begin{aligned}\langle \nu_A | \psi(t) \rangle &= e^{-iE_1 t} \cos \theta_{12} \langle \nu_A | \nu_1 \rangle + e^{-iE_2 t} \sin \theta_{12} \langle \nu_A | \nu_2 \rangle \\ &= e^{-iE_1 t} [\cos^2 \theta_{12} + \sin^2 \theta_{12} e^{-i(E_2 - E_1)t}]\end{aligned}$$

Survival probability:

$$\mathcal{P}(\nu_A \rightarrow \nu_A) = 1 - 4 \cos^2 \theta_{12} \sin^2 \theta_{12} \sin^2 [(E_2 - E_1)t/2]$$

A. Two Level System

Survival probability:

$$\mathcal{P}(\nu_A \rightarrow \nu_A) = 1 - 4 \cos^2 \theta_{12} \sin^2 \theta_{12} \sin^2 [(E_2 - E_1)t/2]$$

Massive, relativistic neutrinos (why?)

$$E_2 - E_1 = \sqrt{p^2 + m_2^2} - \sqrt{p^2 + m_1^2} \approx \frac{m_2^2 - m_1^2}{2p} \approx \frac{m_2^2 - m_1^2}{2E} \quad t = L/v \approx L$$

A. Two Level System

Survival probability:

$$\mathcal{P}(\nu_A \rightarrow \nu_A) = 1 - 4 \cos^2 \theta_{12} \sin^2 \theta_{12} \sin^2 [(E_2 - E_1)t/2]$$

Massive, relativistic neutrinos (why?)

$$E_2 - E_1 = \sqrt{p^2 + m_2^2} - \sqrt{p^2 + m_1^2} \approx \frac{m_2^2 - m_1^2}{2p} \approx \frac{m_2^2 - m_1^2}{2E} \quad t = L/v \approx L$$

$$\mathcal{P}(\nu_A \rightarrow \nu_A) = 1 - 4 \cos^2 \theta_{12} \sin^2 \theta_{12} \sin^2 \left[\frac{(m_2^2 - m_1^2)L}{4E} \right]$$

$$= 1 - \sin^2 2\theta_{12} \sin^2 \left[\frac{(m_2^2 - m_1^2)L}{4E} \right]$$

A. Two Level System

$$\begin{aligned}\mathcal{P}(\nu_A \rightarrow \nu_A) &= 1 - 4 \cos^2 \theta_{12} \sin^2 \theta_{12} \sin^2 \left[\frac{(m_2^2 - m_1^2)L}{4E} \right] \\ &= 1 - \sin^2 2\theta_{12} \sin^2 \left[\frac{(m_2^2 - m_1^2)L}{4E} \right]\end{aligned}$$

- *Two massless neutrinos:* $\theta_{12} = 0$
- *At least one massive neutrino:* $\theta_{12} \neq 0$ and $\mathcal{P}(\nu_A \rightarrow \nu_A) < 1$
- *Dependence on Δm^2 x (L/E)*
- *Transition probability:* $\mathcal{P}(\nu_A \rightarrow \nu_B) = 1 - \mathcal{P}(\nu_A \rightarrow \nu_A)$

B. Three Light Neutrinos

Lepton mixing:

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} N_R + \text{h.c.} \quad \tilde{H}_a = \epsilon_{ab} H_b^*$$

Pontecorvo-Maki-Nakagawa-Sakata

$$\nu_{Li}^I = (S_\nu)_{ij} \nu_{Lj}^{\text{diag}}$$

$$N_{Ri}^I = (T_N)_{ij} N_{Rj}^{\text{diag}}$$

$$\ell_{Li}^I = (S_\ell)_{ij} \ell_{Lj}^{\text{diag}}$$

$$\ell_{Ri}^I = (T_\ell)_{ij} \ell_{Rj}^{\text{diag}}$$

$$V_{\text{PMNS}} = S_\ell^\dagger S_\nu$$

$$J_\mu^{W-} = \bar{L} \gamma_\mu \tau^- V_{\text{PMNS}} L$$

B. Three Light Neutrinos

Pontecorvo-Maki-Nakagawa-Sakata

$$V_{\text{PMNS}} =$$

$$\begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{bmatrix} \times \text{diag}(1, e^{i\frac{\alpha_{21}}{2}}, e^{i\frac{\alpha_{31}}{2}}) .$$

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

“Atmospheric”

“Reactor”

“Solar”

B. Three Light Neutrinos

Oscillation probability (vacuum)

$$\mathcal{P}(\nu_\alpha \rightarrow \nu_\beta) =$$

$$\begin{aligned} & \delta_{\alpha\beta} - 4 \sum_{i>j} \Re(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin^2 \left(\Delta m_{ij}^2 \frac{L}{4E} \right) \\ & + 2 \sum_{i>j} \Im(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin \left(\Delta m_{ij}^2 \frac{L}{2E} \right) \end{aligned}$$

$$V_{PMNS} = U$$

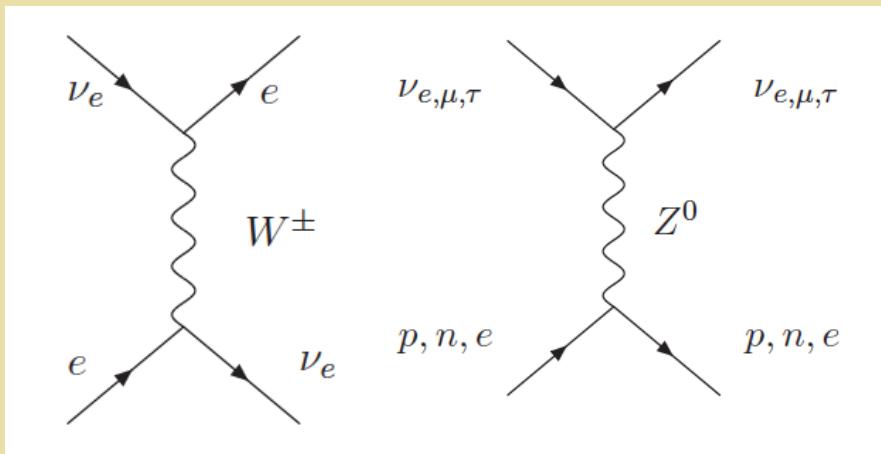
C. Oscillations in Matter

References:

- *P. Hernandez, CERN-2016-005*

Matter Effects: MSW

Forward scattering in matter



CC Hamiltonian

$$\mathcal{H}_{CC} = 2\sqrt{2}G_F [\bar{e}\gamma_\mu P_L \nu_e] [\bar{\nu}_e \gamma^\mu P_L e]$$

Fierz transf

$$= 2\sqrt{2}G_F [\bar{e}\gamma_\mu P_L e] [\bar{\nu}_e \gamma^\mu P_L \nu_e]$$

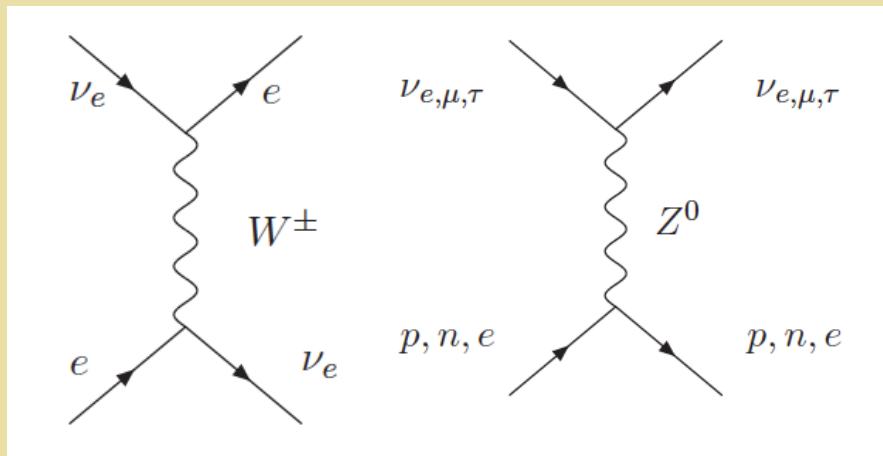
$$\langle \bar{e}\gamma_\mu P_L e \rangle_{\text{unpol. med.}} = \delta_{\mu 0} \frac{N_e}{2}$$

$$\langle \mathcal{H}_{CC} + \mathcal{H}_{NC} \rangle_{\text{unpol. med.}} = \bar{\nu} V_m \gamma^0 (1 - \gamma_5) \nu$$

$$V_m = \begin{pmatrix} \frac{G_F}{\sqrt{2}} \left(N_e - \frac{N_n}{2} \right) & 0 & 0 \\ 0 & \frac{G_F}{\sqrt{2}} \left(-\frac{N_n}{2} \right) & 0 \\ 0 & 0 & \frac{G_F}{\sqrt{2}} \left(-\frac{N_n}{2} \right) \end{pmatrix}$$

Matter Effects: MSW

Forward scattering in matter



CC Hamiltonian

$$\mathcal{H}_{CC} = 2\sqrt{2}G_F [\bar{e}\gamma_\mu P_L \nu_e] [\bar{\nu}_e \gamma^\mu P_L e]$$

Fierz transf

$$= 2\sqrt{2}G_F [\bar{e}\gamma_\mu P_L e] [\bar{\nu}_e \gamma^\mu P_L \nu_e]$$

$$\langle \bar{e}\gamma_\mu P_L e \rangle_{\text{unpol. med.}} = \delta_{\mu 0} \frac{N_e}{2}$$

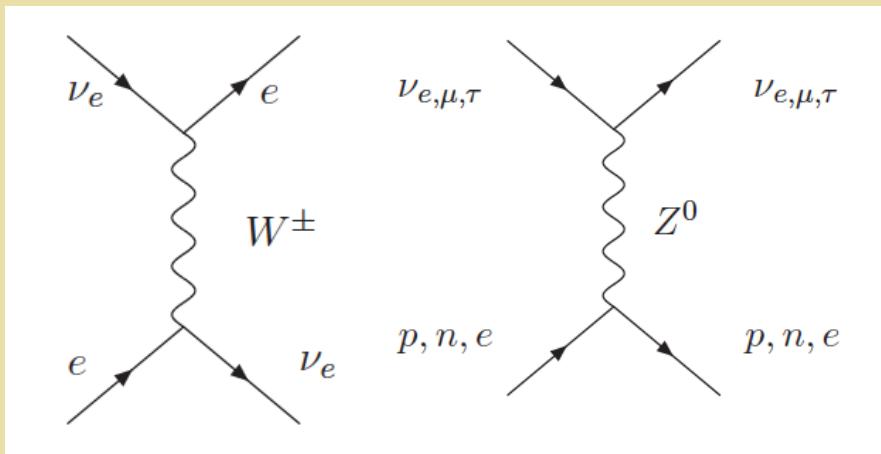
$$\langle \mathcal{H}_{CC} + \mathcal{H}_{NC} \rangle_{\text{unpol. med.}} = \bar{\nu} V_m \gamma^0 (1 - \gamma_5) \nu$$

$$V_m = \begin{pmatrix} \frac{G_F}{\sqrt{2}} \left(N_e - \frac{N_n}{2} \right) & 0 & 0 \\ 0 & \frac{G_F}{\sqrt{2}} \left(-\frac{N_n}{2} \right) & 0 \\ 0 & 0 & \frac{G_F}{\sqrt{2}} \left(-\frac{N_n}{2} \right) \end{pmatrix}$$

Neutral current contribution

Matter Effects: MSW

Forward scattering in matter



CC Hamiltonian

$$\mathcal{H}_{CC} = 2\sqrt{2}G_F [\bar{e}\gamma_\mu P_L \nu_e] [\bar{\nu}_e \gamma^\mu P_L e]$$

Fierz transf

$$= 2\sqrt{2}G_F [\bar{e}\gamma_\mu P_L e] [\bar{\nu}_e \gamma^\mu P_L \nu_e]$$

$$\langle \bar{e}\gamma_\mu P_L e \rangle_{\text{unpol. med.}} = \delta_{\mu 0} \frac{N_e}{2}$$

$$\langle \mathcal{H}_{CC} + \mathcal{H}_{NC} \rangle_{\text{unpol. med.}} = \bar{\nu} V_m \gamma^0 (1 - \gamma_5) \nu$$

$$V_m = \begin{pmatrix} \frac{G_F}{\sqrt{2}} \left(N_e - \frac{N_n}{2} \right) & 0 & 0 \\ 0 & \frac{G_F}{\sqrt{2}} \left(-\frac{N_n}{2} \right) & 0 \\ 0 & 0 & \frac{G_F}{\sqrt{2}} \left(-\frac{N_n}{2} \right) \end{pmatrix}$$

Q: why no e,p NC contributions ?

Matter Effects: MSW

Forward scattering in matter

$$\langle \mathcal{H}_{CC} + \mathcal{H}_{CC} \rangle_{\text{unpol. med.}} = \bar{\nu} V_m \gamma^0 (1 - \gamma_5) \nu$$

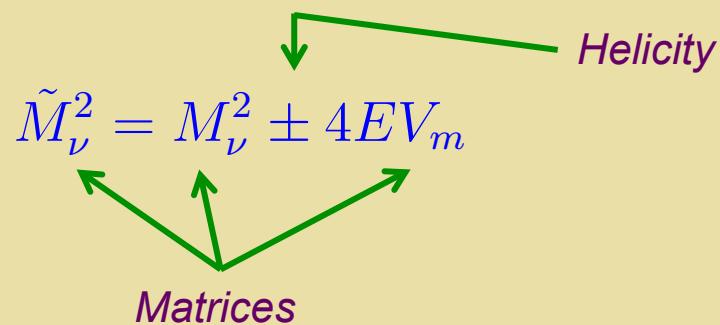
$$V_m = \begin{pmatrix} \frac{G_F}{\sqrt{2}} \left(N_e - \frac{N_n}{2} \right) & 0 & 0 \\ 0 & \frac{G_F}{\sqrt{2}} \left(-\frac{N_n}{2} \right) & 0 \\ 0 & 0 & \frac{G_F}{\sqrt{2}} \left(-\frac{N_n}{2} \right) \end{pmatrix}$$

Dirac Eq: Effective mass = $f(E, h)$

$$\tilde{M}_\nu^2 = M_\nu^2 \pm 4EV_m$$

Matrices

Helicity



Δm^2 & mixing angles
depend on E & N_e

Matter Effects: MSW

Δm^2 & mixing angles
depend on E & N_e

Two-flavor example:

$$\Delta \tilde{m}^2 = \sqrt{\left(\Delta m^2 \cos 2\theta \mp 2\sqrt{2}E G_F N_e\right)^2 + (\Delta m^2 \sin 2\theta)^2},$$

$$\sin^2 2\tilde{\theta} = \frac{(\Delta m^2 \sin 2\theta)^2}{(\Delta \tilde{m}^2)^2}$$

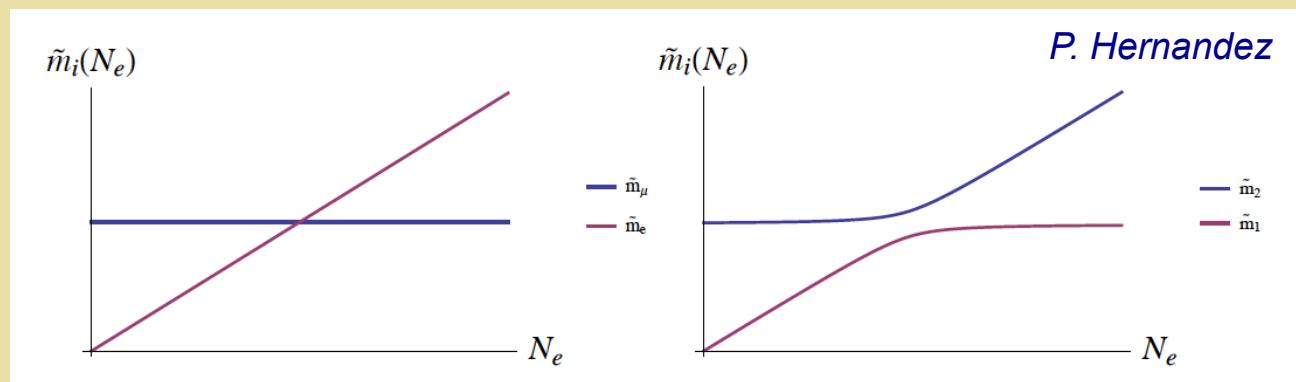
Resonance:

$$\sin^2 2\tilde{\theta} = 1$$

$$\sqrt{2} G_F N_e \mp \frac{\Delta m^2}{2E} \cos 2\theta = 0$$

Matter Effects: MSW

Resonance: would be level crossing



Matter Effects: MSW

Variable matter density (sun)

Consider adiabatic variation of N_e

$$\begin{aligned} |\tilde{\nu}_1\rangle &= |\nu_e\rangle \cos \tilde{\theta} - |\nu_\mu\rangle \sin \tilde{\theta}, \\ |\tilde{\nu}_2\rangle &= |\nu_e\rangle \sin \tilde{\theta} + |\nu_\mu\rangle \cos \tilde{\theta}. \end{aligned}$$

Electron neutrino produced at center of sun

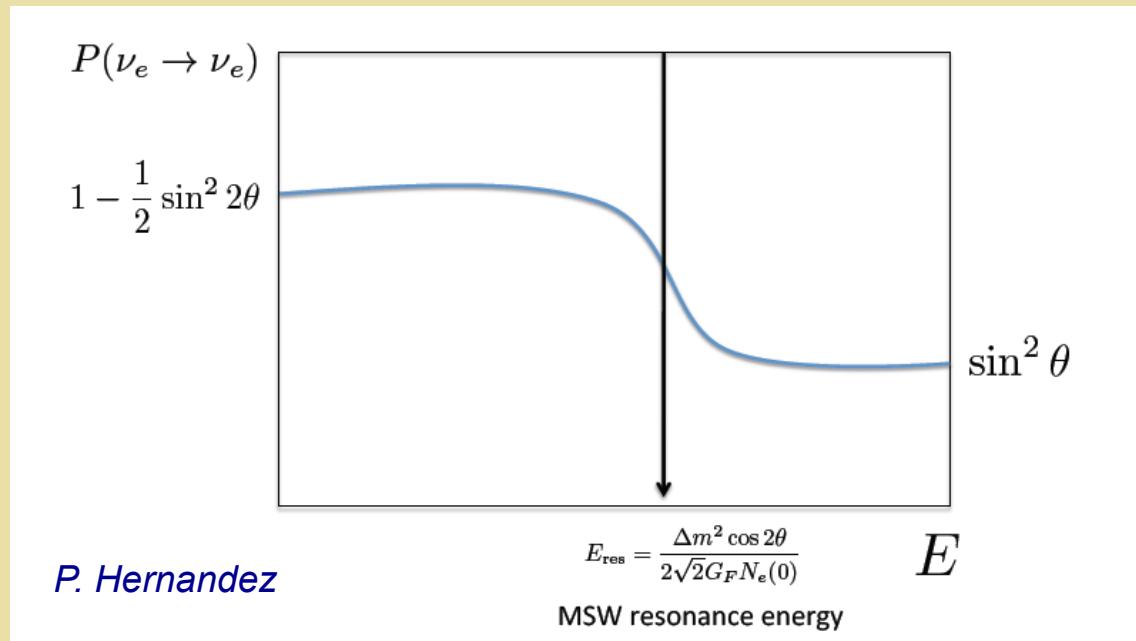
$$2\sqrt{2}G_F N_e(0) \xrightarrow{E} \gg \Delta m^2 \cos 2\theta \quad \tilde{\theta} \simeq \frac{\pi}{2} \Rightarrow |\nu_e\rangle \simeq |\tilde{\nu}_2\rangle$$

Exiting sun: $N_e = 0$

$$\tilde{\theta} \rightarrow \theta_{\text{vac}}$$

$$\nu_e \rightarrow \nu_\mu$$

Matter Effects: MSW



$$\mathcal{P}(\nu_A \rightarrow \nu_A) = 1 - \sin^2 2\theta_{12} \sin^2 \left[\frac{(m_2^2 - m_1^2)L}{4E} \right]$$

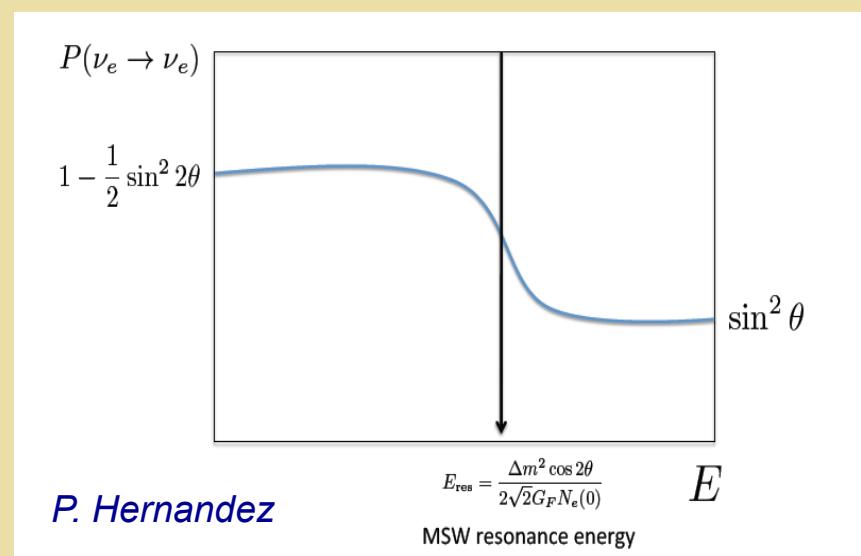
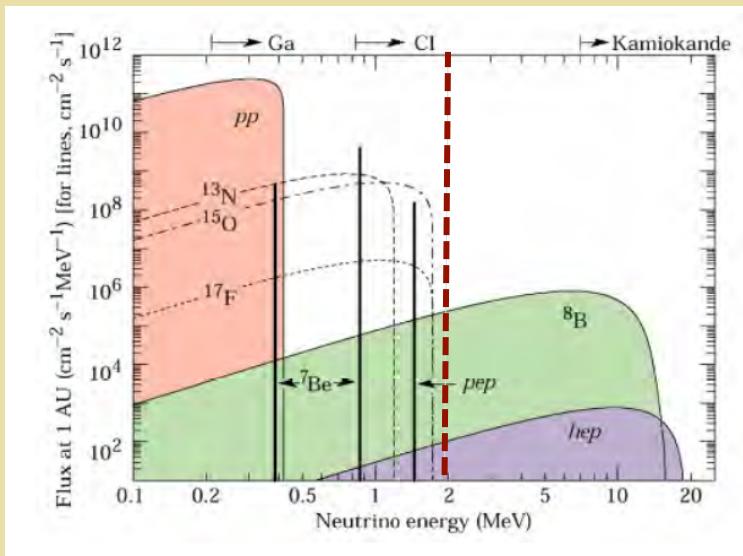


$L/E \gg 1$: average to 1/2

Neutrino spectrum includes both regions: need to take the “MSW effect” into account

Solar Neutrinos

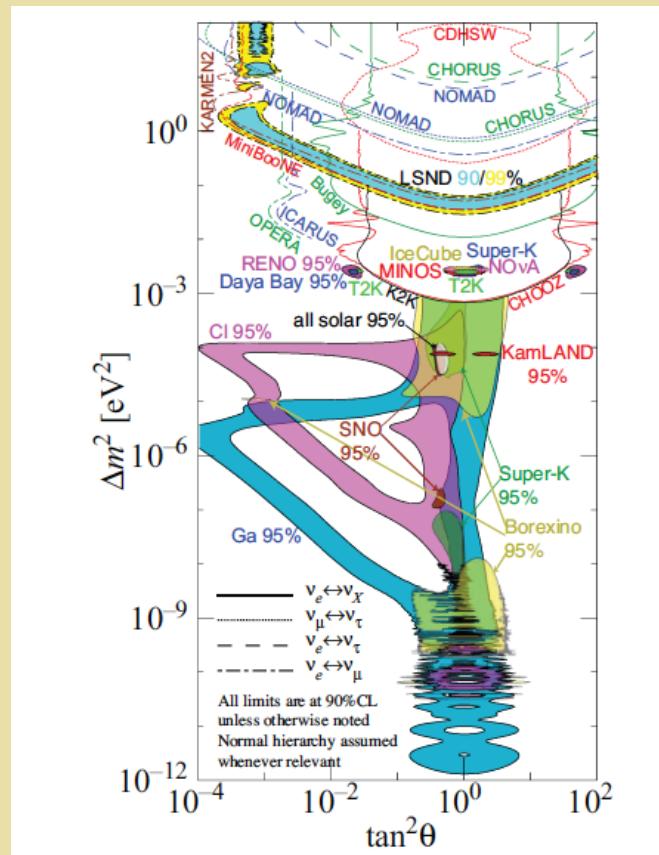
Standard Solar Model (SSM)



- *Analysis of terrestrial spectrum requires convolution of SSM predictions w/ MSW effect*

Oscillation Parameters

Particle Data Group & H. Murayama



Parameter	best-fit	3σ
Δm_{21}^2 [10 ⁻⁵ eV ²]	7.37	6.93 – 7.96
$\Delta m_{31(23)}^2$ [10 ⁻³ eV ²]	2.56 (2.54)	2.45 – 2.69 (2.42 – 2.66)
$\sin^2 \theta_{12}$	0.297	0.250 – 0.354
$\sin^2 \theta_{23}, \Delta m_{31(32)}^2 > 0$	0.425	0.381 – 0.615
$\sin^2 \theta_{23}, \Delta m_{32(31)}^2 < 0$	0.589	0.384 – 0.636
$\sin^2 \theta_{13}, \Delta m_{31(32)}^2 > 0$	0.0215	0.0190 – 0.0240
$\sin^2 \theta_{13}, \Delta m_{32(31)}^2 < 0$	0.0216	0.0190 – 0.0242
δ/π	1.38 (1.31)	2 σ : (1.0 - 1.9) (2 σ : (0.92-1.88))

Oscillation Parameters

See D. Parno slides

<https://globalfit.astroparticles.es/>

parameter	best fit $\pm 1\sigma$	3σ range	relative 1σ uncertainty
Δm_{21}^2 [10^{-5} eV 2]	$7.55^{+0.20}_{-0.16}$	7.05–8.14	2.4%
$ \Delta m_{31}^2 $ [10^{-3} eV 2] (NO)	2.50 ± 0.03	2.41–2.60	1.3%
$ \Delta m_{31}^2 $ [10^{-3} eV 2] (IO)	$2.42^{+0.03}_{-0.04}$	2.31–2.51	
$\sin^2 \theta_{12}/10^{-1}$	$3.20^{+0.20}_{-0.16}$	2.73–3.79	5.5%
$\sin^2 \theta_{23}/10^{-1}$ (NO)	$5.47^{+0.20}_{-0.30}$	4.45–5.99	4.7%
$\sin^2 \theta_{23}/10^{-1}$ (IO)	$5.51^{+0.18}_{-0.30}$	4.53–5.98	4.4%
$\sin^2 \theta_{13}/10^{-2}$ (NO)	$2.160^{+0.083}_{-0.069}$	1.96–2.41	3.5%
$\sin^2 \theta_{13}/10^{-2}$ (IO)	$2.220^{+0.074}_{-0.076}$	1.99–2.44	
δ/π (NO)	$1.32^{+0.21}_{-0.15}$	0.87–1.94	10%
δ/π (IO)	$1.56^{+0.13}_{-0.15}$	1.12–1.94	9%

deSalas et al, 1708.01186 (May 2018)

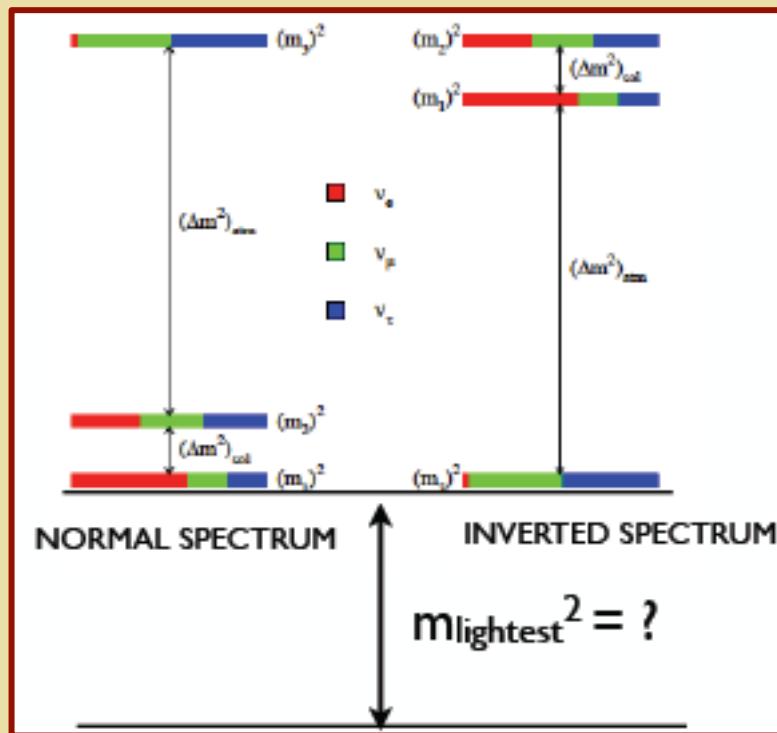
III. Open Questions

- *Majorana or Dirac ?*
- *Mass hierarchy ?*
- *Absolute mass ?*
- *CPV ?*
- *Light sterile neutrinos ?*
- *Neutrino vs. quark mixing ?*
- *Theoretical origin of m_ν*

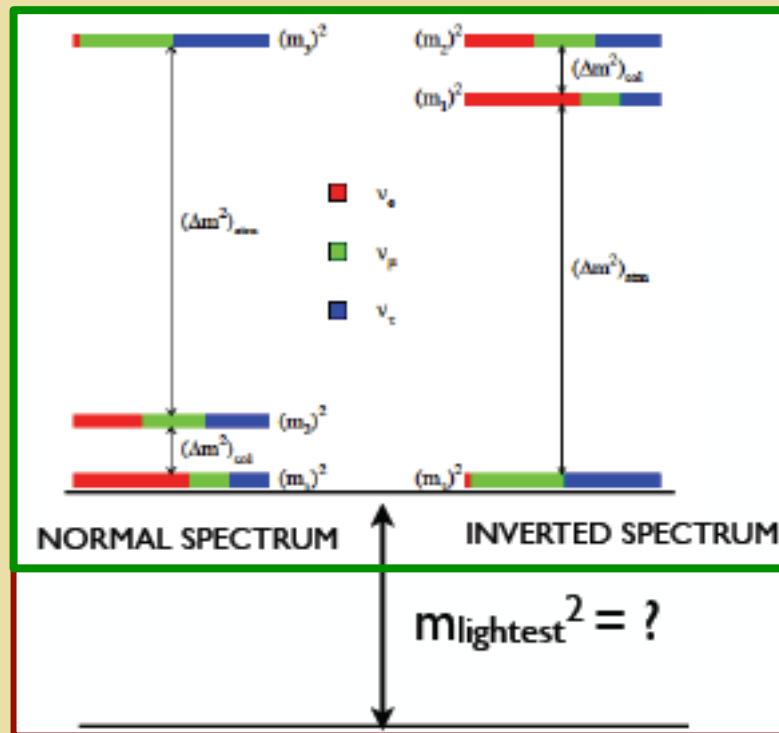
III. Open Questions

- *Majorana or Dirac ?*
- *Mass hierarchy ?*
- *Absolute mass ?*
- *CPV ?*
- *Light sterile neutrinos ?*
- *Neutrino vs. quark mixing ?*
- *Theoretical origin of m_ν*

Absolute Mass & Mass Hierarchy

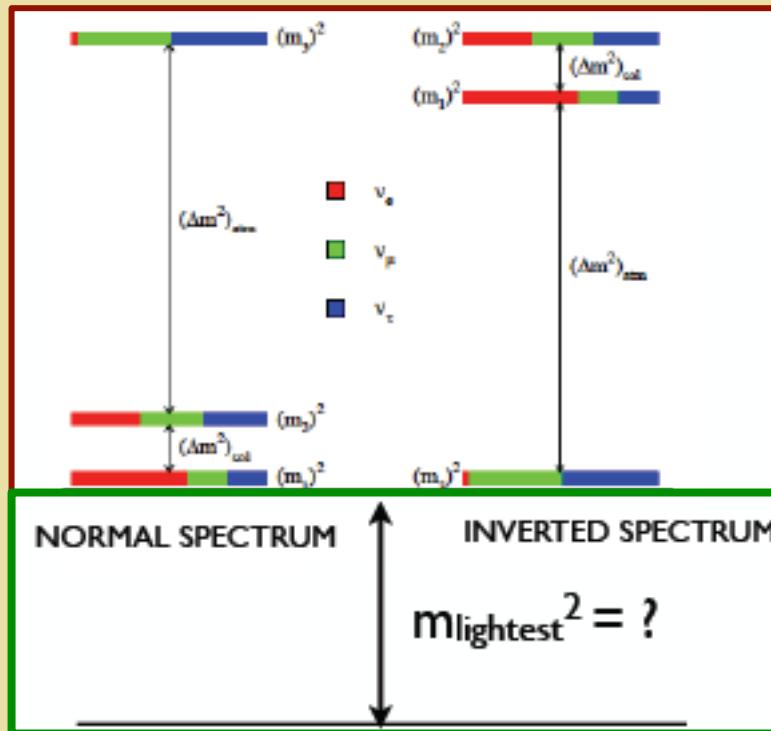


Absolute Mass & Mass Hierarchy



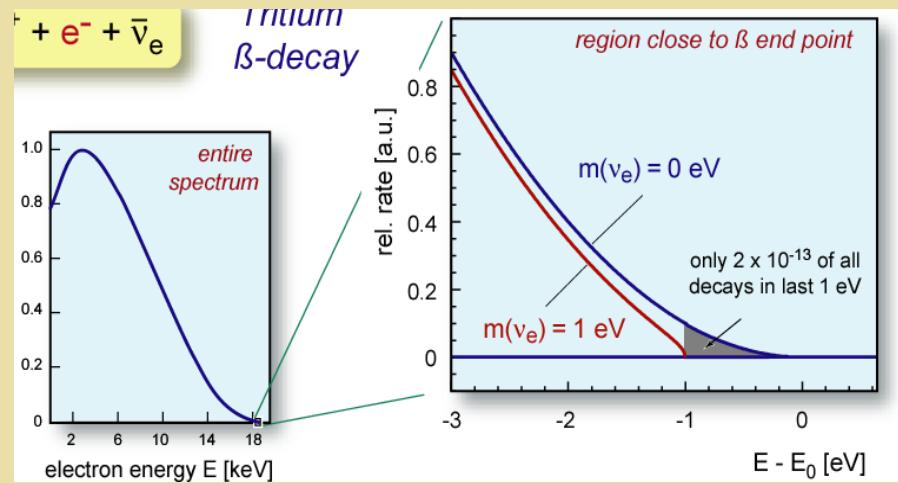
Oscillation
Expts

Absolute Mass & Mass Hierarchy



- ${}^3H \beta\text{-decay}$
- *Cosmology & astrophysics*

Kinematic Neutrino Mass Measurements



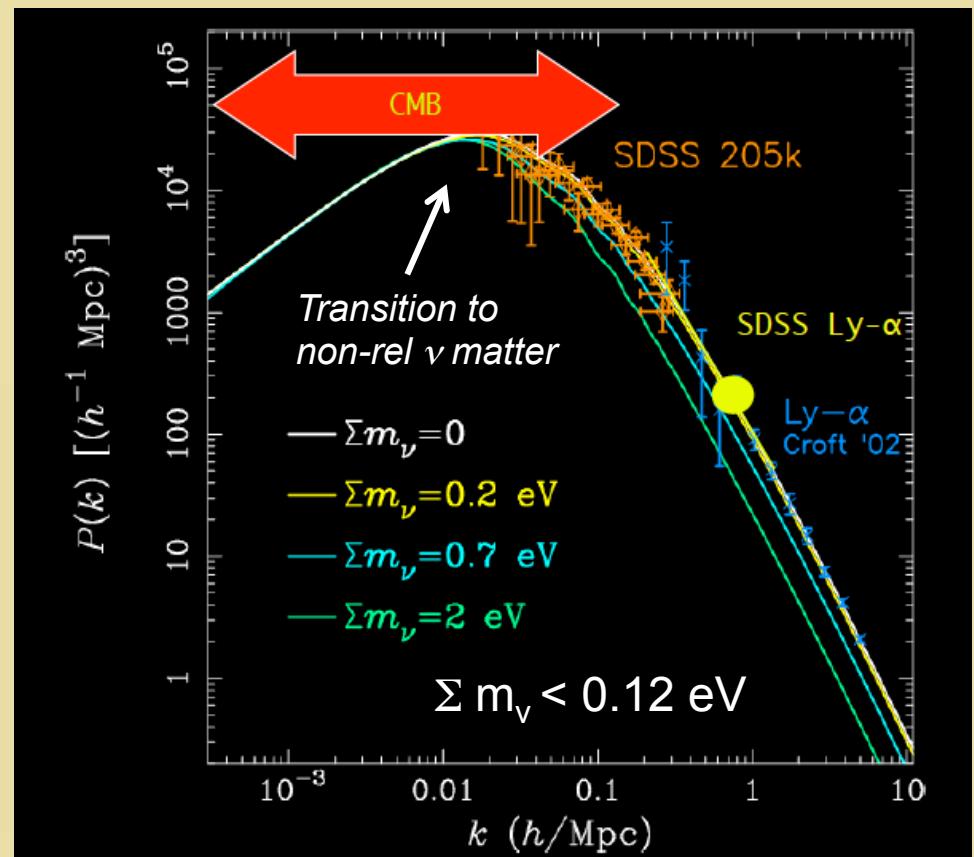
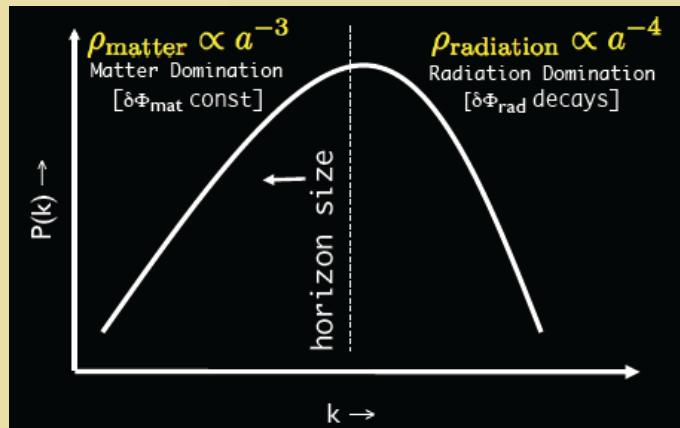
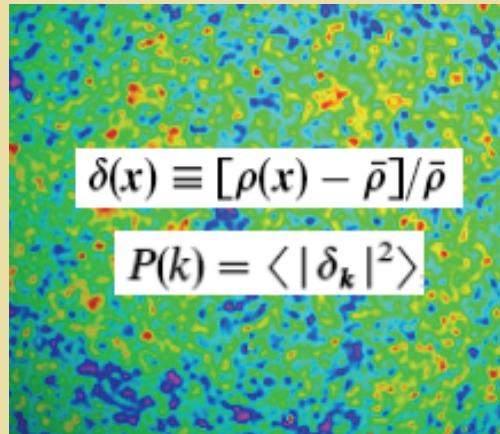
KATRIN



$$\frac{dN}{dp_e} \propto (E_0 - E_e)^2 \left[1 - \frac{m_\nu^2}{(E_0 - E_e)^2} \right]$$

Neutrino Mass & Cosmology

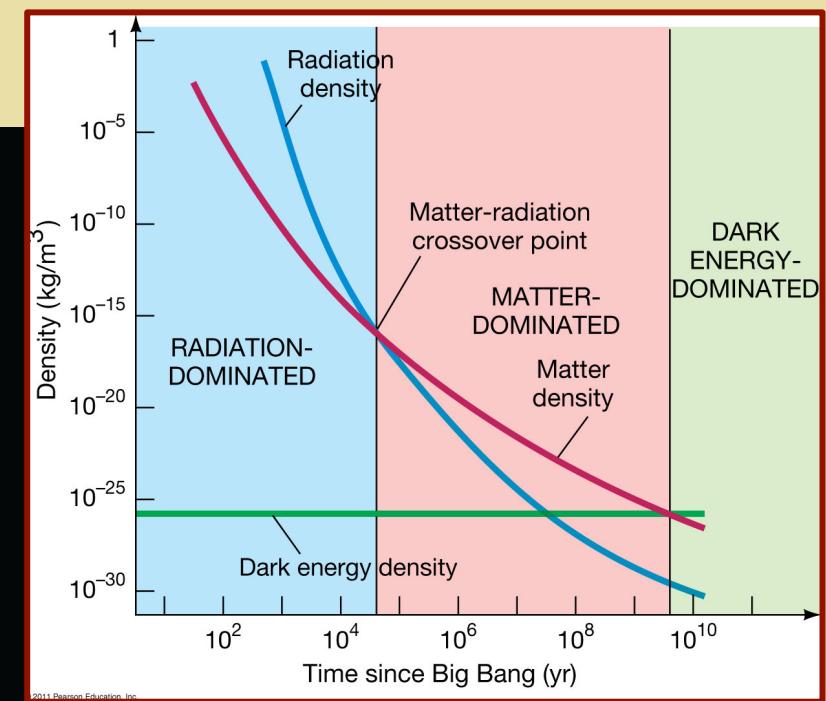
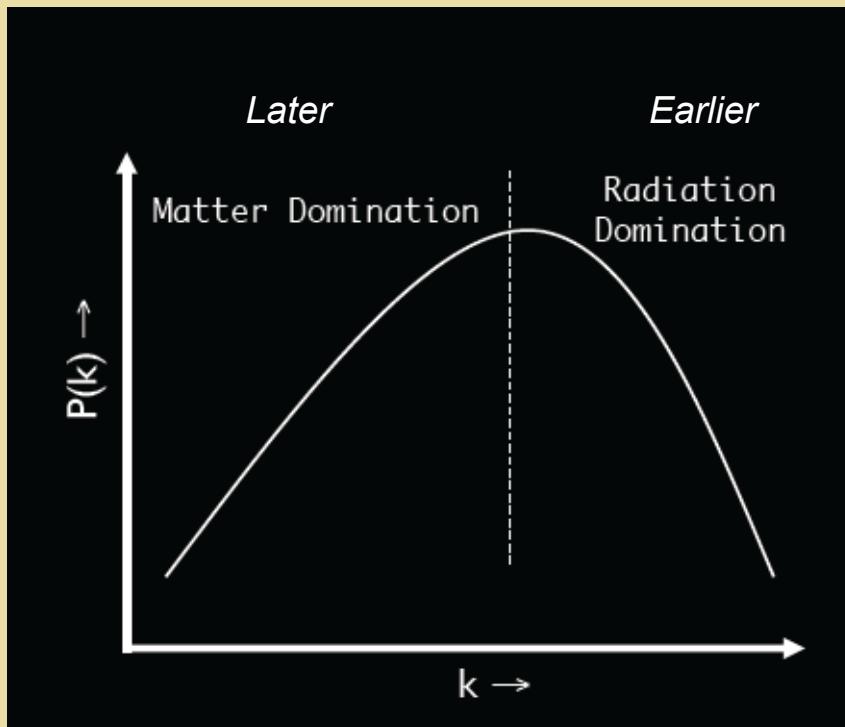
Matter Power Spectrum



Massive neutrinos suppress power (relative to large scale power) at scales below free streaming scale

Neutrino Mass & Cosmology

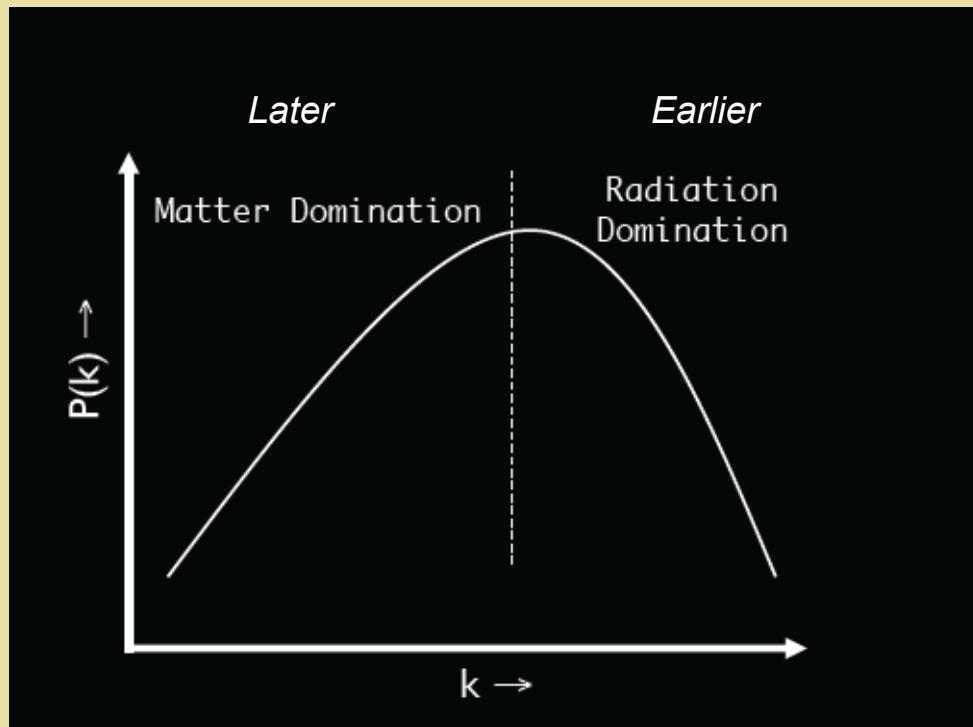
Matter Power Spectrum



J. Brau, U. Oregon

Neutrino Mass & Cosmology

Matter Power Spectrum



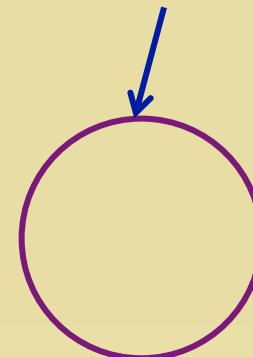
Neutrino Free Streaming

$$\Omega_M = \Omega_\nu + \Omega_{DM} + \Omega_B$$

$$\delta\rho_\nu \longleftrightarrow \delta\rho_{DM}$$

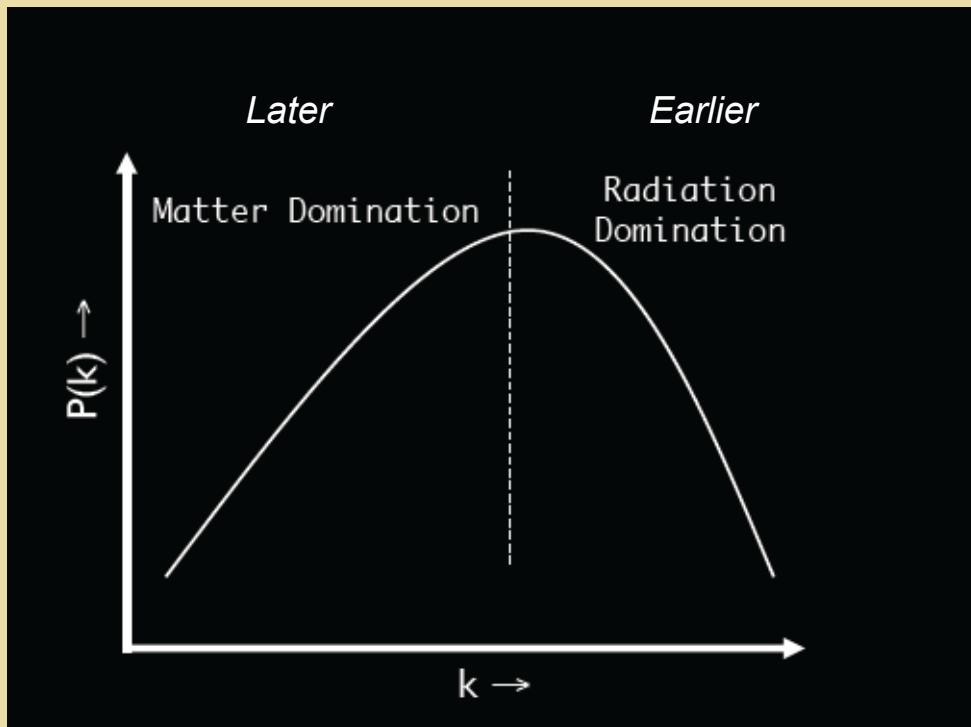
Free Streaming Scale

$$L_{fs} \propto m_\nu^{-1/2}$$



Neutrino Mass & Cosmology

Matter Power Spectrum



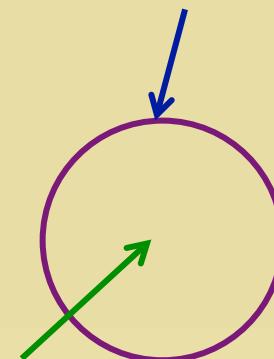
Neutrino Free Streaming

$$\Omega_M = \Omega_\nu + \Omega_{DM} + \Omega_B$$

$$\delta\rho_\nu \leftrightarrow \delta\rho_{DM}$$

Free Streaming Scale

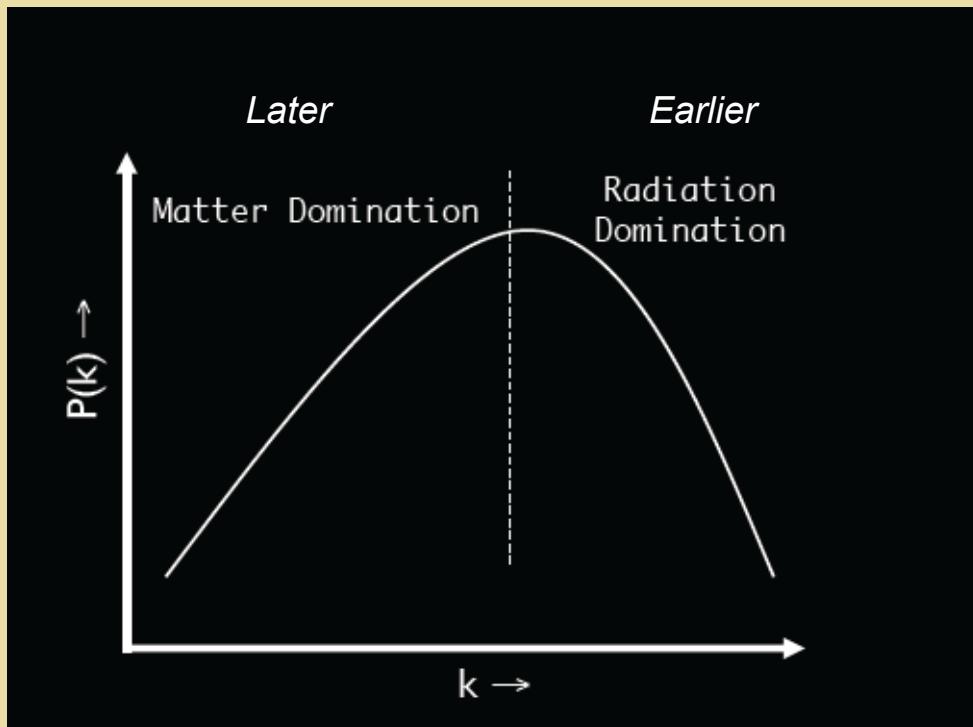
$$L_{fs} \propto m_\nu^{-1/2}$$



$\delta\rho_\nu$ (power) suppressed
for $L < L_{fs}$

Neutrino Mass & Cosmology

Matter Power Spectrum



$\delta\rho_\nu$ (power) suppressed
for $L < L_{fs}$

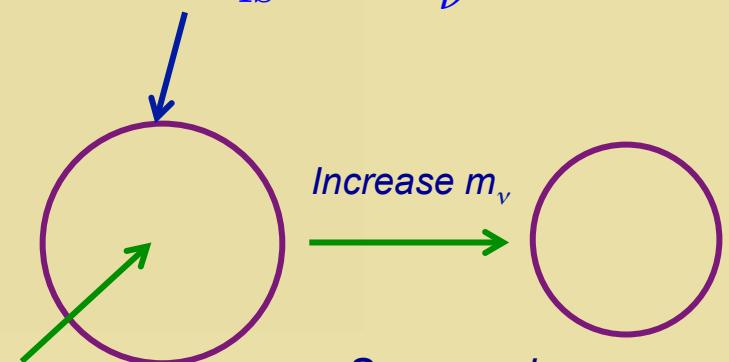
Neutrino Free Streaming

$$\Omega_M = \Omega_\nu + \Omega_{DM} + \Omega_B$$

$$\delta\rho_\nu \leftrightarrow \delta\rho_{DM}$$

Free Streaming Scale

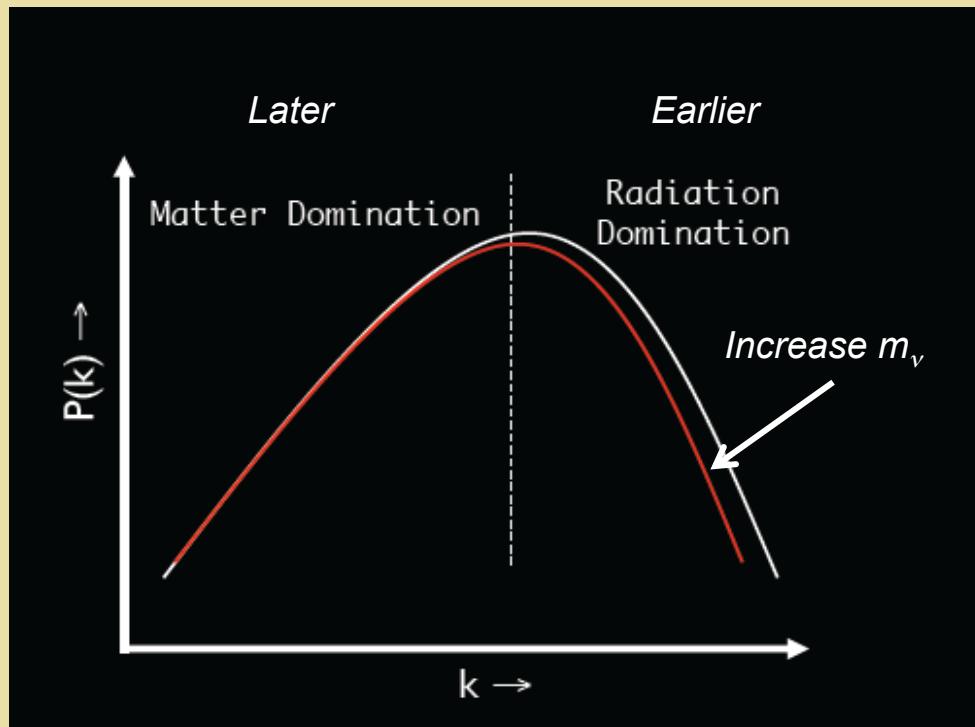
$$L_{fs} \propto m_\nu^{-1/2}$$



Suppression moves
to smaller scales →
Larger k

Neutrino Mass & Cosmology

Matter Power Spectrum



$$\Sigma m_\nu < 0.12 \text{ eV}$$

Palanque-Dalabrouille '15

$\delta\rho_\nu$ (power) suppressed
for $L < L_{fs}$

K. Abazajian ACFI neutrino mass workshop

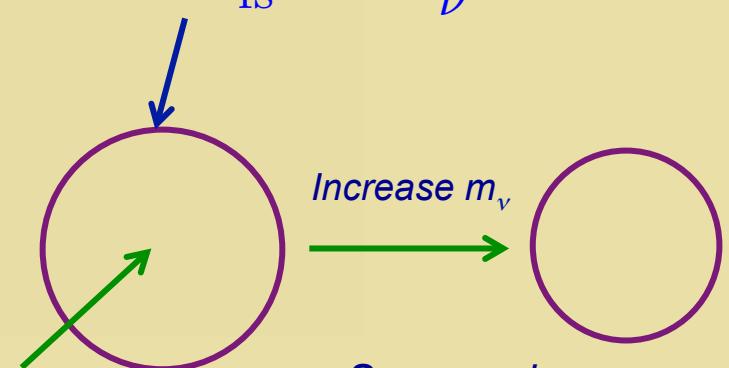
Neutrino Free Streaming

$$\Omega_M = \Omega_\nu + \Omega_{DM} + \Omega_B$$

$$\delta\rho_\nu \leftrightarrow \delta\rho_{DM}$$

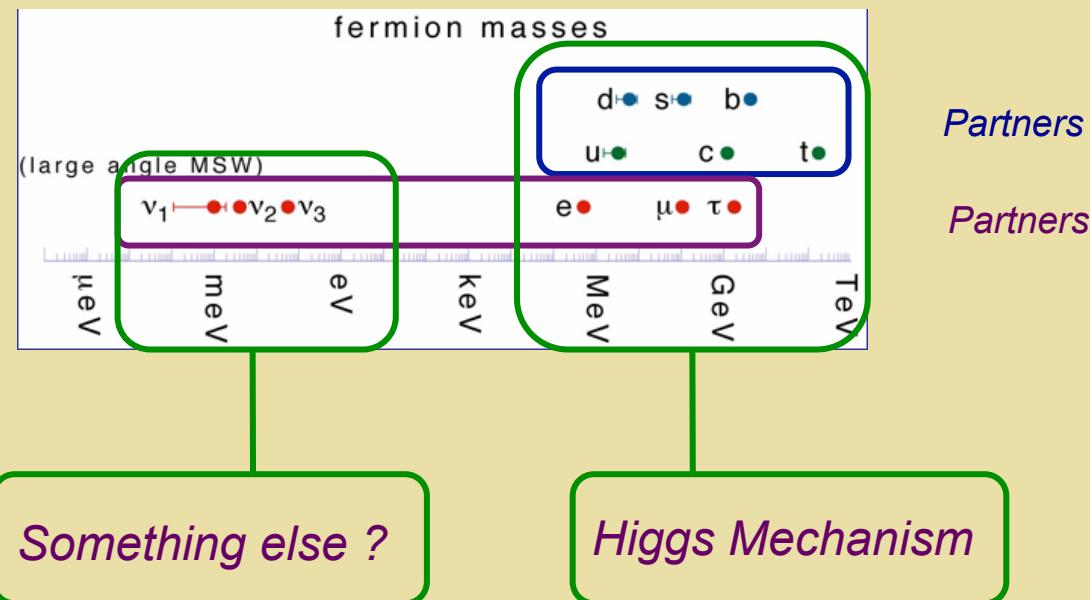
Free Streaming Scale

$$L_{fs} \propto m_\nu^{-1/2}$$



Suppression moves
to smaller scales →
Larger k

IV. Neutrino Mass Models



*How do we understand the origin
of m_ν theoretically ?*

IV. Neutrino Mass Models

- *Type I see-saw* “ ν SM”, “ ν MSSM”
 - *Type II see-saw* LRSM
 - *Type III see-saw* GUTs
 - *Inverse see-saw* LRSM
 - *Radiative* MSSM
- + combinations & many other examples

IV. Neutrino Mass Models

- *Type I see-saw* “ ν SM”, “ ν MSSM”
 - *Type II see-saw* LRSM
 - *Type III see-saw* GUTs
 - *Inverse see-saw* LRSM
 - *Radiative* MSSM
- + combinations & many other examples

Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Lepton number violating



Neutrino Mass Models

- *Type I see-saw* “ ν SM”, “ ν MSSM”
 - *Type II see-saw* LRSM
 - *Type III see-saw* GUTs
 - *Inverse see-saw* LRSM
 - *Radiative* MSSM
- + combinations & many other examples

Type I See-Saw

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

One generation: SM + one N_R

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} N_R + \text{h.c.} + M_N \bar{N}_R^C N_R$$



$$\mathcal{L}_{\text{mass}} = \begin{pmatrix} \bar{\nu}_L & \bar{N}_R^C \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D & M_N \end{pmatrix} \begin{pmatrix} \nu_L \\ N_R \end{pmatrix}$$

Type I See-Saw

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

One generation: SM + one N_R

Lepton number violating

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} N_R + \text{h.c.} + M_N \bar{N}_R^C N_R$$



$$\mathcal{L}_{\text{mass}} = \begin{pmatrix} \bar{\nu}_L & \bar{N}_R^C \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D & M_N \end{pmatrix} \begin{pmatrix} \nu_L \\ N_R \end{pmatrix}$$

Type I See-Saw

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

One generation: SM + one N_R

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} N_R + \text{h.c.} + M_N \bar{N}_R^C N_R$$



$$\mathcal{L}_{\text{mass}} = \begin{pmatrix} \bar{\nu}_L & \bar{N}_R^C \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D & M_N \end{pmatrix} \begin{pmatrix} \nu_L \\ N_R \end{pmatrix}$$

Lepton number violating

Eigenvalues

$$m_1 \approx \frac{m_D^2}{M_N}$$

$$m_2 \approx M_N$$

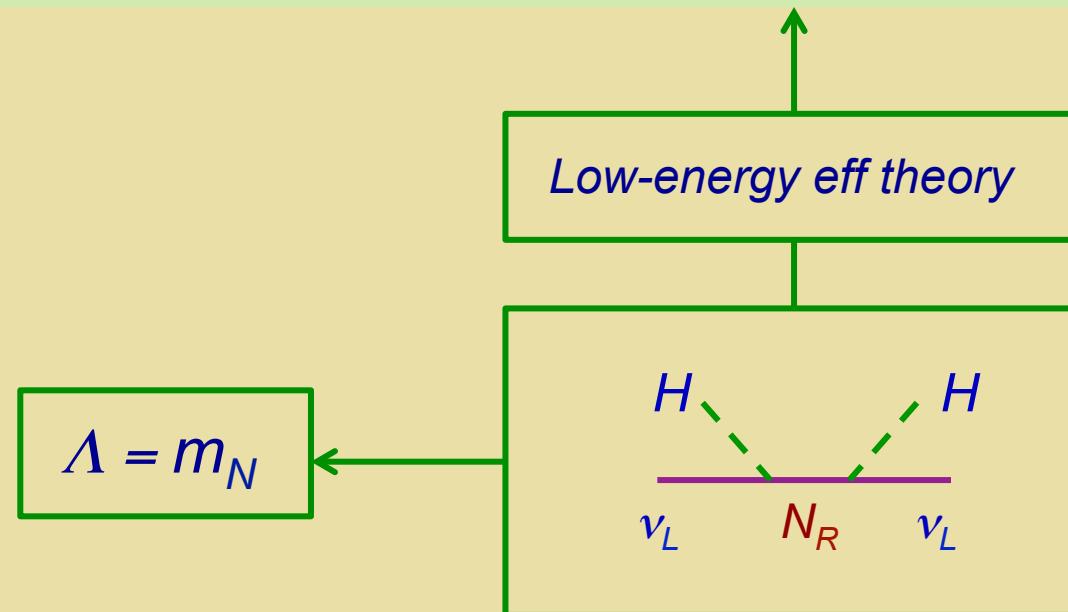
Type I See-Saw

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana



Type I See-Saw

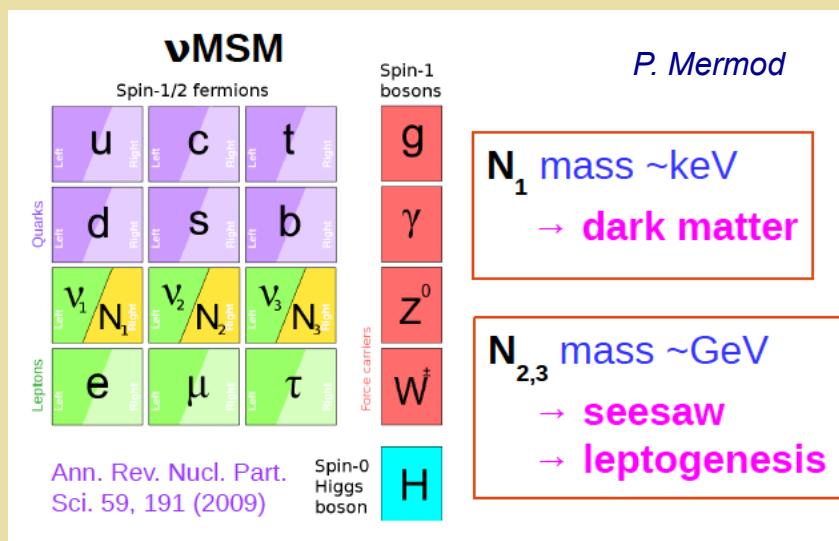
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

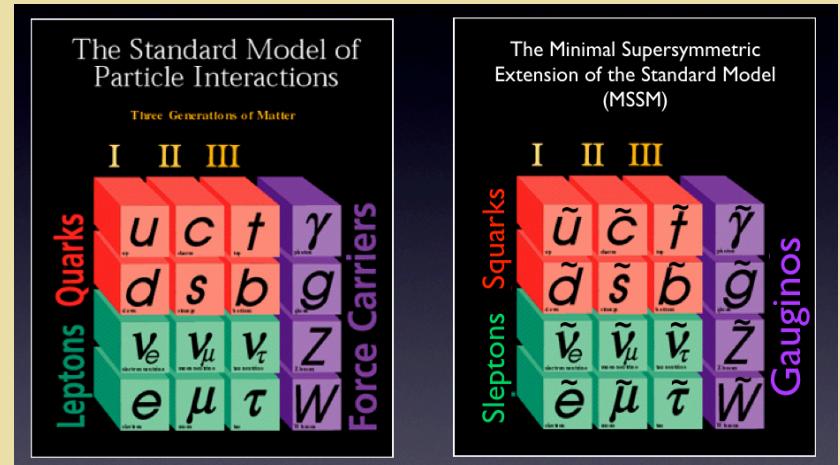
$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

“ν MSM”



“ν MSSM”



$$+ \left(N_R, \tilde{N}_R \right)$$

Neutrino Mass Models

- *Type I see-saw* “ ν SM”, “ ν MSSM”
 - *Type II see-saw* LRSM
 - *Type III see-saw* GUTs
 - *Inverse see-saw* LRSM
 - *Radiative* MSSM
- + combinations & many other examples

Type II See-Saw

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Introduce “Complex Triplet”: $\Delta_L \sim (1, 3, 2)$

$$\Delta_L = \begin{pmatrix} \Delta^+ \sqrt{2} & \Delta^+ \\ \Delta^0 & -\Delta^+ \sqrt{2} \end{pmatrix}$$

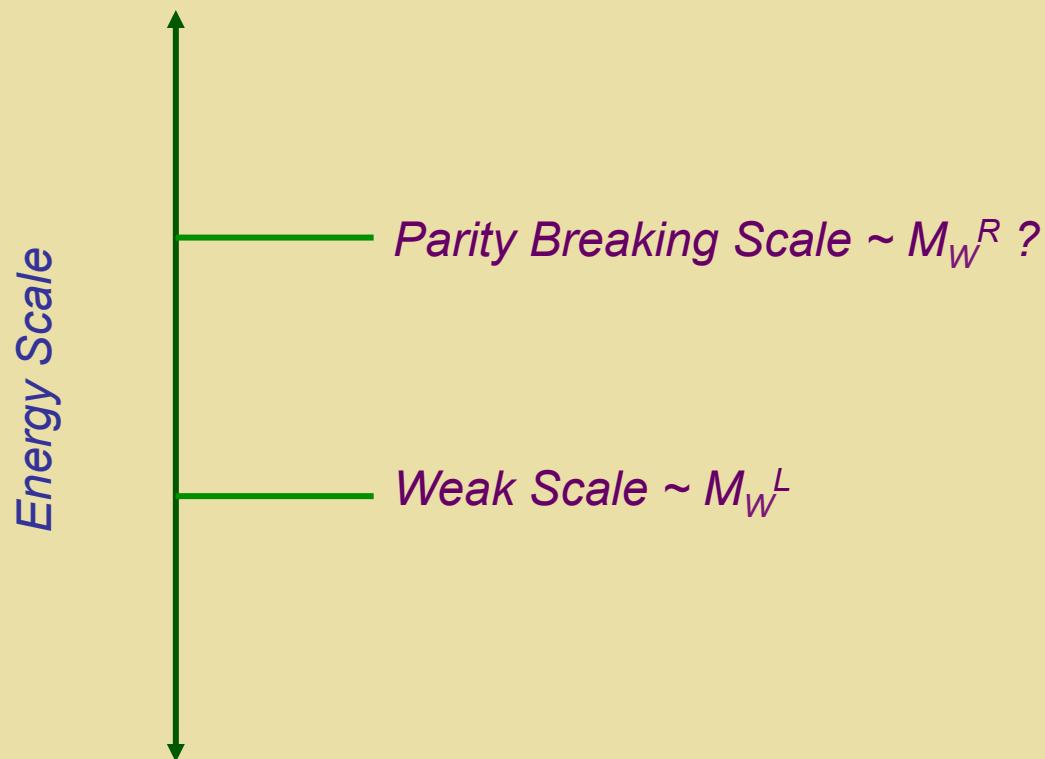
Δ^0 vev \rightarrow Majorana m_ν

$$\mathcal{L} = \frac{g}{2} h_{ij} [\bar{L}^{C_i} \varepsilon \Delta_L L^j] + \text{h.c.}$$

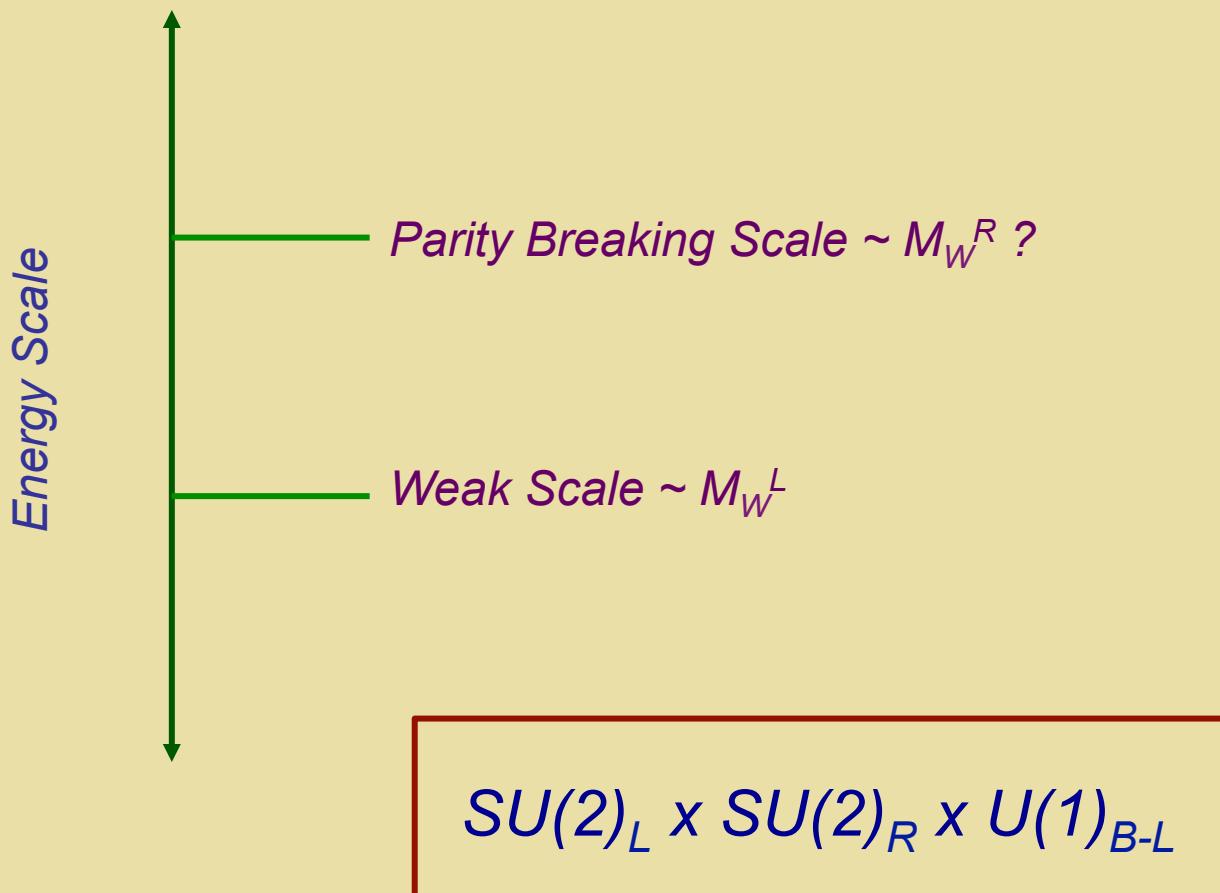
Lepton number violating

Types I & II: Left-Right Symmetric Model

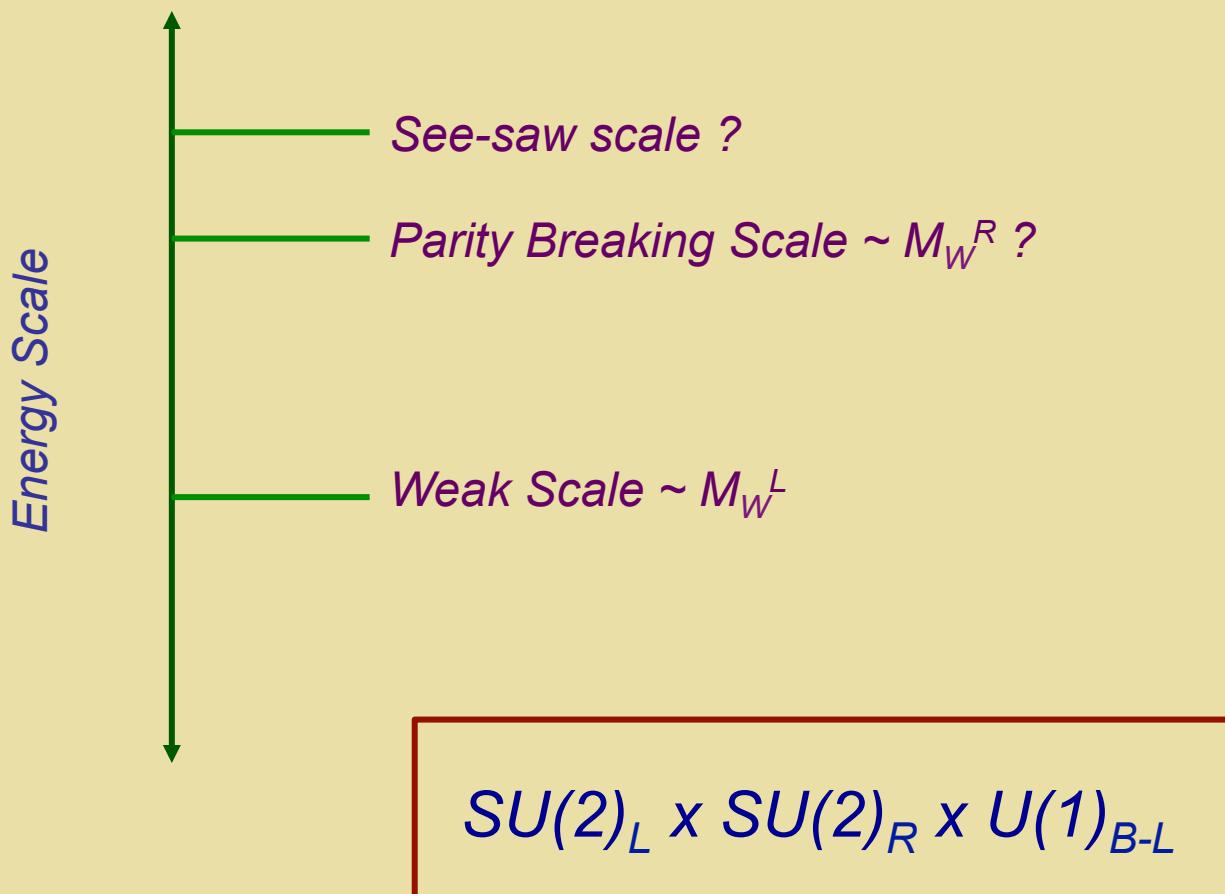
BSM Mass Scale



Left-Right Symmetric Model



Left-Right Symmetric Model



Left-Right Symmetric Model

Gauge boson mass eigenstates

$$W_1^+ = \cos \xi W_L^+ + \sin \xi e^{-i\alpha} W_R^+$$

$$W_2^+ = -\sin \xi e^{i\alpha} W_L^+ + \cos \xi W_R^+$$

CKM Matrices for LH & RH sectors: quarks

$$\begin{aligned} u_{Li}^I &= (S_u)_{ij} u_{Lj}^{\text{mass}} \\ u_{Ri}^I &= (T_u)_{ij} u_{Rj}^{\text{mass}} \\ d_{Li}^I &= (S_d)_{ij} d_{Lj}^{\text{mass}} \\ d_{Ri}^I &= (T_d)_{ij} d_{Rj}^{\text{mass}} \end{aligned}$$



$$V_{\text{CKM}}^L = S_u^\dagger S_d$$

$$V_{\text{CKM}}^R = T_u^\dagger T_d$$

Left-Right Symmetric Model

Gauge boson mass eigenstates

$$W_1^+ = \cos \xi W_L^+ + \sin \xi e^{-i\alpha} W_R^+$$

$$W_2^+ = -\sin \xi e^{i\alpha} W_L^+ + \cos \xi W_R^+$$

PMNS Matrices for LH & RH sectors: leptons

$$\begin{aligned}\nu_{Li}^I &= (S_\nu)_{ij} \nu_{Lj}^{\text{diag}} \\ N_{Ri}^I &= (T_N)_{ij} N_{Rj}^{\text{diag}} \\ \ell_{Li}^I &= (S_\ell)_{ij} \ell_{Lj}^{\text{diag}} \\ \ell_{Ri}^I &= (T_\ell)_{ij} \ell_{Rj}^{\text{diag}}\end{aligned}$$

$$V_{\text{PMNS}}^L = S_\nu^\dagger S_\ell$$

Left-Right Symmetric Model

Two sources of m_ν :

$$\mathcal{L} = \frac{g}{2} h_{ij} [L^{C_i} \varepsilon \Delta_L L^j] + (L \leftrightarrow R) + \text{h.c.}$$

Type I see-saw

$$\mathcal{L}_{\text{mass}} = \begin{pmatrix} \bar{\nu}_L & \bar{N}_R^C \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D & M_N \end{pmatrix} \begin{pmatrix} \nu_L \\ N_R \end{pmatrix} + \boxed{m_L \bar{\nu}_L^C \nu_L}$$

$$m_N \sim g h_R \langle \Delta_R^0 \rangle$$

Type II see-saw

$$m_L \sim g h_L \langle \Delta_L^0 \rangle$$

Neutrino Mass Models

- *Type I see-saw* “ ν SM”, “ ν MSSM”
 - *Type II see-saw* LRSM
 - *Type III see-saw* GUTs
 - *Inverse see-saw* LRSM
 - *Radiative* MSSM
- + combinations & many other examples

Type II See-Saw

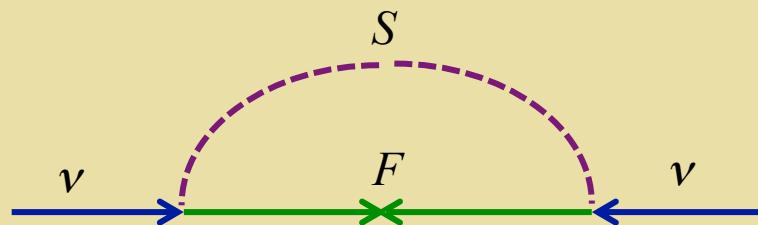
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Introduce new scalars (S) & Majorana fermions (F): “mediators”



Attach Higgs lines as appropriate to get Weinberg operator

Recent mini-review: H. Sugiyama, 1505.01738

Type II See-Saw

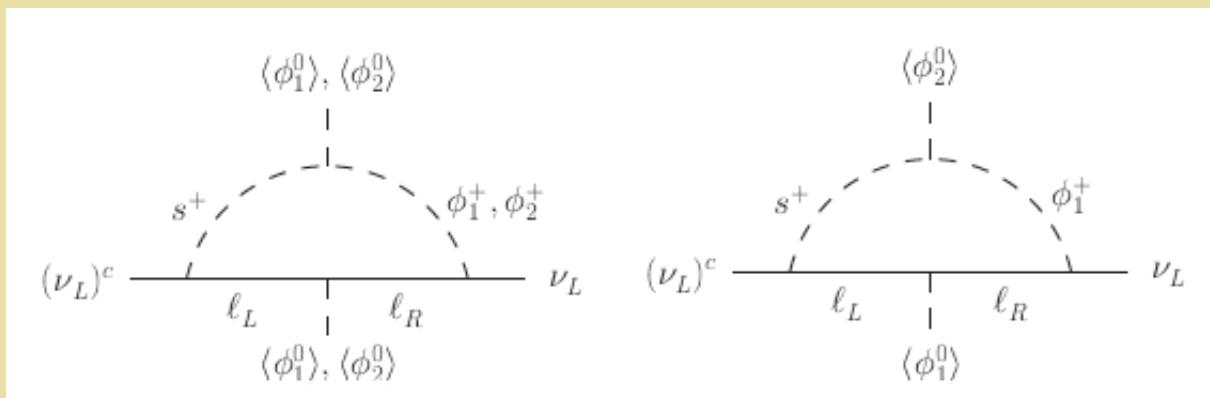
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Introduce new scalars (S) & Majorana fermions (F): “mediators”



“Zee Model”

Recent mini-review: H. Sugiyama, 1505.01738

Type II See-Saw

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

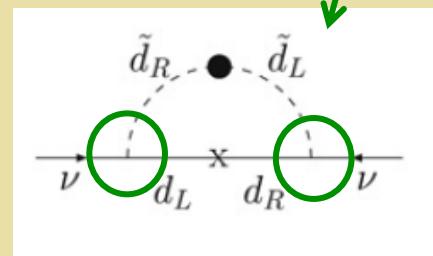
$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

SUSY with “R parity” violation $P_R = (-1)^{2S+3(B-L)}$

“Superpotential”

$$W_{\Delta L=1} = \frac{1}{2} \lambda_{ijk} L_i L_j \bar{e}_k + \lambda'_{ijk} L_i Q_j \bar{d}_k + \mu'_i L_i H_u,$$



V. Discussion Questions

- *What is the see-saw scale (M_N)?*
- *What might the comparison of m_ν from terrestrial & astrophysical determinations teach us?*
- *How do we know $N_\nu = 3$?*
- *How might we determine the correct neutrino mass model ?*