

Partonic Hadron Structure II

Paul E Reimer Physics Division Argonne National Laboratory July 2017

- I. Where are we now?
- II. Longitudinal Parton Distributions
- III. Nuclear Effects
- IV. Spin



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Partonic Hadron Structure | Review

- 1. Physicists seek organization and order
- 2. The quark model can explain many of the properties of the observed hadronic spectra.
- 3. Elastic scattering shows that the proton is not a point particle



- 4. Richard Feynman was a genius.
 - Hadron-hadron scattering is a collision of many point-like particles (partons)
 - Each parton carries a fraction of the hadron's momentum
 - Parton distributions can be described in terms of a probability distribution of a parton existing with momentum fraction in [x, x+dx]
- 5. Deep Inelastic Scattering cam be described in terms of a summation over point-like scattering from partons.
- 6. Parton distributions may be extracted from hard scattering data.
 - Generally requires data from multiple measurements
 - Care must be taken to avoid false assumptions





FIG. 6. This figure shows the experimental points at 236 Mev and the attempts to fit the shape of the experimental curve. The best fit lies near 0.78×10^{-13} cm.

 $F_2^{\mu p}(x) \propto \sum_{q \in \{u, d, \dots\}} e_q^2 x \left[q(x, Q^2) + \bar{q}(x, Q^2) \right]$

 $F_2^{\nu p}(x) + F_2^{\nu n} \propto \sum_{q \in \{u, d, \dots\}} x \left[q(x, Q^2) + \bar{q}(x, Q^2) \right]$

 $xF_3^{\nu N}(x) \propto \sum_{q \in \{u,d,\dots\}} x \left[q(x,Q^2) - \bar{q}(x,Q^2)\right]$

Becky is now Happy!

TAN



Infinite Momentum Frame

Feynman:

"By Lorentz transformation, the fields to be radiated are becoming narrower and narrower in the z direction as W rises. The energy in this field is therefore distributed as a δ function in z."

- The concept of longitudinal parton distributions exists in a Lorentz frame in which the hadron is moving fast enough that transverse momentum of the partons within the hadron can be neglected, or P_z >>> M.
- The proton is Lorentz contracted which implies that during the ∆t of the reaction, the partons do not communicate with each other. That is, they are quasi-free particles

Gluons are partons too!

Gluons also form partons and carry a fraction (possibly large) of the proton's momentum

Momentum conservation
$$\int_0^1 x \left[u(x) + \bar{u}(x) + d(x) + \bar{d}(x) + \cdots \right] dx = 1$$

- Part of the "..." in the momentum constraint. Sometimes written out explicitly as g(x).
- Difficult to measure because
 - They do not couple to the electromagnetic force (hard to probe)
 - They are suppressed by 1/ α_{EM}
- Observation of direct photons



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Scaling

Bjorken limit

 $Q^2 \rightarrow \infty$ and $\nu \rightarrow \infty$ at fixed x

The structure functions were measured to be relatively constant over many decades of Q²

OBSERVED BEHAVIOR OF HIGHLY INELASTIC ELECTRON-PROTON SCATTERING

M. Breidenbach, J. I. Friedman, and H. W. Kendall Department of Physics and Laboratory for Nuclear Science,* Massachusetts Institute of Technology, Cambridge, Massachusetts 02139

and

E. D. Bloom, D. H. Coward, H. DeStaebler, J. Drees, L. W. Mo, and R. E. Taylor Stanford Linear Accelerator Center,[†] Stanford, California 94305 (Received 22 August 1969)



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Scaling

Bjorken limit

 $Q^2 \to \infty$ and $\nu \to \infty$ at fixed xThe structure functions were measured to be relatively constant over many decades of Q²

 Field theories predicted gross violation of scaling except in a special case:

Asymptotically Free

theories



What parton distributions are available?

CTEQ

- Coordinated Theor.-Exp. Project on QCD
- https://www.physics.smu.edu/scalise/cteq/
- MRSTW—Martin, Roberts, Stirling, Thorne, Watt
- GRV—Gluck, Reya, Vogt
- NNPDF—Neural Network PDF
 - Goal: reduce assumption/theory bias
 - https://nnpdf.hepforge.org/

HERA

 using only data from the H1 and Zeus experiments at DESY HERA

BS15

- Statistical model to constrain PDFs
- Others. . .

For "black box" use, LHAPDF compiles most available PDF sets into a common interface

- Currently 773 different PDF sets available
- Many superseded by newer versions as more data becomes available
- Many exploring different constraints on QCD
- <u>https://lhapdf.hepforge.org/</u>

If you want to fit your own, there is a data archive at Durham called HEPDAT

- <u>http://hepdat.net</u>
- Also, be sure that your data makes it into this archive. . .

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Example fits: CTEQ14NNLO



FIG. 5: The CT14 parton distribution functions at Q = 2 GeV and Q = 100 GeV for $u, \overline{u}, d, \overline{d}, s = \overline{s}$, and g.

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Nuclear effects

Hadrons

Guggenheim, Bilbao, Spain

- Interna

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Free proton vs. bound proton

Are the parton distributions of nucleons within a nucleus the same as free nucleons?

- Hard scattering assumption includes an implicit assumption that the interaction is so energetic that the binding of quarks in a proton is negligible
 - so surely, the binding of protons in the nucleus is also small
- Nuclear targets are used in many experiments
 - \Box v-DIS cross sections are small -> require dense targets
- Do the quarks change configuration?



The EMC Effect (European Muon Collaboration)

Volume 123B, number 3,4 PHYSICS LETTERS THE RATIO OF THE NUCLEON STRUCTURE FUNCTIONS F_2^N FOR IRON AND DEUTERIUM The European Muon Collaboration Attempt to increase integrated 13 luminosity using a denser target. g 12 (Fe) / F₂^N Comparison iron data with earlier 11 deuterium data found a striking × ~ difference 10 Systematics: 09 Normalization uncertainty ²H and Fe x-dependent uncertainties in slope 08 Additional data helped clarify effect 02 04 06 Х 0

31 March 1983

Additional data revealed

- Larger-x depletion in strength was correct
- Low-x structure was significantly different.

Are the parton distributions actually different?

$$F_2(x) = \sum_{q \in \{u, d...\}} e_q^2 \left[\frac{q(x)}{q(x)} + \bar{q}(x) \right]_{0.85}$$



- Effect generally divided into 4 regions
- x < 0.1 Shadowing Region</p>
- 0.1 < x < 0.3 Anti-Shadowing</p>
- 0.3 < x < 0.6 EMC effect</p>
- 0.6 < x Fermi motion</p>

What do we understand? (or what are we guessing at?)



Fermi motion

 Intrinsic motion of nucleons in a nucleus at rest.



Fermi motion

 Intrinsic motion of nucleons in a nucleus at rest.

Shadowing

- Small-x partons (gluons) from one nucleus overlap with those of a neighboring nucleus
- Expect to start at some x_{onset} and saturate at x_{sat}~ ½ rM



Fermi motion

 Intrinsic motion of nucleons in a nucleus at rest.

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- Small-x partons (gluons) from one nucleus overlap with those of a neighboring nucleus
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- Anti-Shadowing
- Small-x parton (gluons) overlap?



Fermi motion

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- Small-x parton (gluons) overlap?

EMC Effect

Many theories or models exist. None are satisfactory.





Droll-Van Cross Section				
Diell-Tall Clus	S SECLIOI		MRST	
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\mathbf{I} \\ \mathbf{I} \end{bmatrix} = \begin{bmatrix} \mathbf{I} & \mathbf{I} \\ \mathbf{I} \\ \mathbf{I} \end{bmatrix} = \begin{bmatrix} I$		
 Cross section is a convolution of beam and 10 0 0.2 0.4 0.6 0.8 1 target parton distributions 				
$\frac{d^2\sigma}{dx_{\rm b}dx_{\rm t}} = \frac{4\pi\alpha^2}{x_{\rm b}x_{\rm t}s} \sum_{q \in \{u,d,s,\ldots\}} e_q^2 [\bar{q}_{\rm t}(x_{\rm t})q_{\rm b}(x_{\rm b}) + \bar{q}_{\rm b}(x_{\rm b})q_{\rm t}(x_{\rm t})]$ Acceptance limited				
(Fixed Target, Hadron Beam)				
$(2/3)^2$ vs. $(1/3)^2$	Beam	Sensitivity	Experiment	
	Hadron	Beam quarks target antiquarks	Fermilab, J-PARC RHIC (forward acpt.)	
	Anti-Hadron	Beam antiquarks Target quarks	J-PARC, GSI-FAIR Fermilab Collider	
Paul E Reimer Partonic Structure	Meson	Beam antiquarks Target quarks	COMPASS, J-PARC	

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Structure of nucleonic matter: How do DIS and Drell-Yan data compare?



Structure of nucleonic matter: Where are the nuclear pions?

- The binding of nucleons in a nucleus is expected to be governed by the exchange of virtual "Nuclear" mesons.
- No antiquark enhancement seen in Drell-Yan (Fermilab E772) data
- Contemporary models predict large effects to antiquark distributions as x increases.

Models must explain both DIS-EMC effect and Drell-Yan



Kulagin and Petti sea vs. valence nuclear effects



FMB—Fermi Motion and Nuclear Binding OS—Off shell effects NS—nuclear shadowing PI—nuclear pions

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Aside: Problem for PDF fits

Many experiments used nuclear targets

- Does this data need to be thrown out now?
- Information of d-quark distributions comes from Deuterium and isospin symmetry
- Neutrino DIS data?
 - Old H2 bubble chamber data OK
 - Modern experiments use iron target
 - Magnitude of Sea Quark distributions dominated by neutrino data

$$q \in \{u, d, \dots\}$$
$$x F_3^{\nu N}(x) \propto \sum_{q \in \{u, d, \dots\}} x \left[q(x, Q^2) - \overline{q}(x, Q^2) \right]$$

 $F_2^{\nu p}(x) + F_2^{\nu n} \propto \sum x \left[q(x, Q^2) + \bar{q}(x, Q^2) \right]$

$$q \in \{u, u, \dots\}$$

- Parameterize measurements?
- K. J. Eskola, V. J. Kolhinen, and P. V. Ruuskanen, Nucl. Phys. B535, 351 (1998);

Becky has fallen asleep



Now add Spin

Dynamics make things messy ... Or more interesting?

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Polarized parton Polarized Proton

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Parton polarized transversely and Proton longitudinally polarized.



$$\mathbf{q}(\mathbf{x}) \equiv \textcircled{} \Rightarrow \Rightarrow \Rightarrow \textcircled{}$$

$$\begin{split} \Delta \mathbf{q}(\mathbf{x}) &\equiv & \textcircled{} \rightarrow & \textcircled{} \rightarrow & \textcircled{} \rightarrow & \swarrow \\ \Delta q(x) &= q_{\rightarrow}(x) - q_{\leftarrow}(x) \\ q &\in \left\{ u, \bar{u}, d, \bar{d}, s, \bar{s}, \dots g \right\} \end{split} \text{Note: Some authors separate While others denote their sum by} \\ \Delta q \end{split}$$

Spin dependent structure functions

$$g_{1}(x) = \frac{1}{2} \sum_{i \in \{u, \bar{u}, d, \dots\}} e_{i}^{2} \left[\Delta q(x) + \Delta \bar{q}(x) \right] \xrightarrow{\text{http://www.scholarpedia.org/article/Helicit}}{\underline{y_dependent_parton_distributions}}$$

$$= \frac{1}{2} \left[\frac{4}{9} \Delta u(x) + \frac{4}{9} \Delta \bar{u}(x) + \frac{1}{9} \Delta d(x) + \frac{1}{9} \Delta \bar{d}(x) + \frac{1}{9} \Delta s(x) + \frac{1}{9} \Delta \bar{s}(x) + \cdots \right]$$

$$\Delta q_{3} = (\Delta u + \Delta \bar{u}) - (\Delta d + \Delta \bar{d}) \qquad g_{A}^{3} = \int_{0}^{1} \Delta q_{3} dx$$

$$\Delta q_{8} = (\Delta u + \Delta \bar{u}) + (\Delta d + \Delta \bar{d}) - 2(\Delta s + \Delta \bar{s}) \qquad g_{A}^{8} = \int_{0}^{1} \Delta q_{8} dx$$

$$\Delta \Sigma = (\Delta u + \Delta \bar{u}) + (\Delta d + \Delta \bar{d}) + (\Delta s + \Delta \bar{s}) \qquad g_{A}^{0} = \int_{0}^{1} \Delta \Sigma dx$$

$$g_{1}(x) = \frac{1}{9} \left[\frac{3}{4} \Delta q_{3}(x) + \frac{1}{4} \Delta q_{8}(x) + \Delta \Sigma(x) \right]$$

Spin dependent structure functions

Bjorken sum rule:

relates g_A^3 and g_A^8 to the axial and vector coupling constants G_A and G_V in β decay

 $g_A^3 \equiv g_A = 1.2670 \pm 0.0035$ and $g_A^8 = 0.585 \pm 0.025$

Elis-Jaffe Sum Rule:
$$g^0_A = g^8_A + 3\left(\Delta s + \Delta ar s
ight)$$

$$\Gamma_{1} = \int_{0}^{1} g_{1}(x)dx$$

$$= \frac{1}{9} \left[\frac{3}{4} \int_{0}^{1} \Delta q_{3}(x)dx + \frac{1}{4} \int_{0}^{1} \Delta g_{8}(x)dx + \int_{0}^{1} \Delta \Sigma(x)dx \right]$$

$$= \frac{1}{9} \left[\frac{3}{4} g_{A} + \frac{1}{4} g_{A}^{8} + \int_{0}^{1} \Delta \Sigma(x)dx \right]$$
The spin carried by the partons

Why do I care—I'm happy nibbling on clover

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The proton's spin

- The proton is a spin-½ particle
- The quarks are spin-½ particles (this was also established in the parton model.)
- The gluons are spin-1 particles

How do the quarks' and gluons' angular momentum add to form a spin-1/2 proton?

Should be able to sum over all partons with appropriate Clepsch-Gordon coefficients , right??

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + L_q + \Delta G$$

Now,

- $\Delta\Sigma$ is integrated over x with a normalization convention requiring $\frac{1}{2}$
- L_a is the orbital angular momentum of the quarks
- ΔG is the integral spin and orbital angular momentum of the glue

EMC Measures g₁

Volume 206, number 2

PHYSICS LETTERS B

19 May 1988

A MEASUREMENT OF THE SPIN ASYMMETRY AND DETERMINATION OF THE STRUCTURE FUNCTION g_1 IN DEEP INELASTIC MUON–PROTON SCATTERING

European Muon Collaboration

- scattered longitudinally polarized muons on a longitudinally polarized target
- measured the asymmetry in cross sections between the parallel and anti-parallel spins



The spin crisis

• Assuming that the missing spin for the Ellis-Jaffe sum rule is in the strange quarks:

$$\langle S_z \rangle_u = 0.373 \pm 0.019 \pm 0.039 \langle S_z \rangle_d = -0.254 \pm 0.019 \pm 0.039 \langle S_z \rangle_s = -0.113 \pm 0.019 \pm 0.039 \langle S_z \rangle_{u+d+s} = 0.006 \pm 0.058 \pm 0.117$$

Compilation of recent experiments gives

 $\Delta\Sigma\approx 0.40$

Nevertheless, where is the proton's spin?

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + L_q + \Delta G$$

Could the problem be with strange quarks and the Ellis-Jaffe sum rule?

Semi-Inclusive Deep Inelastic Scattering SIDIS

- Try to tag the struck quark by detecting a fast outgoing particle
- Measure Kaons to identify strange quarks







Spin dependent structure function g₁

World data compilation from PDG





Cautious, but still somewhat interested Could it be in the glue? **Polarized gluons**

• RHIC
$$\vec{p} + \vec{p} \to \text{Jets} + X$$
 $A_{LL} = \frac{\sigma^{\uparrow\downarrow} - \sigma^{\uparrow\uparrow}}{\sigma^{\uparrow\downarrow} + \sigma^{\uparrow\uparrow}}$



Polarized gluons

• RHIC $\vec{p} + \vec{p} \to \text{Jets} + X$ $A_{LL} = \frac{\sigma^{\uparrow\downarrow} - \sigma^{\uparrow\uparrow}}{\sigma^{\uparrow\downarrow} + \sigma^{\uparrow\uparrow}}$



Global fit to spin density--DSSV

PRL 113, 012001 (2014)

PHYSICAL REVIEW LETTERS

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Evidence for Polarization of Gluons in the Proton

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Werner Vogelsang[§]

Institute for Theoretical Physics, Tübingen University, Auf der Morgenstelle 14, 72076 Tübingen, Germany (Received 17 April 2014; published 2 July 2014)

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Global fit to spin density--DSSV

PRL 113, 012001 (2014)

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week ending 4 JULY 2014



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Global fit to spin density--DSSV

PRL 113, 012001 (2014)

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Evidence for Polarization of Gluons in the Proton



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Decerry City . New J. 1862

INDIAN RESERVATION

CHEVENNEN

Francis Ul ase surtenter to to

MAP OF COLORADO TERRITORY, Compiled from Government Mapskactual Surveys.

MADE IN 1961

$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + L_q + \Delta G$ $\frac{1}{2}\Delta\Sigma \approx 25\% \quad \Delta G \approx 0 - 15\%$

Where are we?

Quarks carry some of the Proton's spin

Gluons might carry some (a small amount) of the Proton's spin

Much of the spin is still lost treasure

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ABAPAHOES

The proton in terms of all parton distributions





aui ב Keimer The valence Sivers עוגדוסעדוסה and ערפוו-Yan

Transverse Momentum Distributions: Introduction



Survive k_T integration

Paul E Reimer The Valence Sivers Distribution and Drell-Yan

Transverse Momentum Distributions: Introduction



Paul E Reimer The Valence Sivers Distribution and Drell-Yan

Leave things spinning, but slowly getting there

Conclusions

- Data looks like scattering from many point particles (partons)
 - These distributions are universal and process independent
- Data allows the determination of the distributions of partons
 - But having the correct assumptions in interpreting data is critical!



FIG. 5: The CT14 parton distribution functions at Q = 2 GeV and Q = 100 GeV for $u, \overline{u}, d, \overline{d}, s = \overline{s}$, and g.



- Data says that these distributions change when the proton is contained in a nucleus
 - Models of these effects are not satisfactory

 The proton's spin is not where we are looking, but the options are narrowing

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + L_q + \Delta G$$

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$$\frac{1}{2}\Delta\Sigma \approx 25\% \quad \Delta G \approx 0 - 15\%$$

