

Fundamental Symmetries and Precision Physics

David Hertzog
University of Washington

- **Lecture 1**
 - Motivations
 - Symmetries, Parity, and the Weak Interaction
 - The Fermi Constant
 - Muon Decay as a test of V-A theory
- **Lecture 2**
 - Neutron beta decay
 - Parity as a tool to probe matter: PVES
 - Highly sensitive low-energy probes of New Physics
 - CPV and Electric Dipole Moments
- **Lecture 3 (transition here at some point ...)**
 - Charged Lepton Flavor Violation
 - Muon $g-2$
 - It's a wrap ...

A question from Day 1

CPT Violation Implies Violation of Lorentz Invariance

O. W. Greenberg

Phys. Rev. Lett. **89**, 231602 – Published 18 November 2002

A interacting theory that **violates CPT invariance necessarily violates Lorentz invariance**. On the other hand, CPT invariance is not sufficient for out-of-cone Lorentz invariance. Theories that violate CPT by having different particle and antiparticle masses must be nonlocal.

Topic 6

Charged Lepton Flavor Violation

Impressive sensitivity to new physics
when the SM theory “is zero” (or sort of)

Progress in Particle and Nuclear Physics 71 (2013) 75–92



Contents lists available at SciVerse ScienceDirect

Progress in Particle and Nuclear Physics

journal homepage: www.elsevier.com/locate/ppnp



Review

Lepton flavor and number conservation, and physics beyond
the standard model



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Review

Flavour violating muon decays

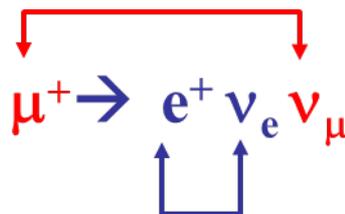
T. Mori^{*}, W. Ootani

International Centre for Elementary Particle Physics, The University of Tokyo, 7-3-1 Hongo, Bunkyo-ku, Tokyo 113-0033, Japan



... those conservation laws again ...

- Non-observation of **decay** $\mu \rightarrow e\gamma$ established that the muon was a distinct particle
- We have been stating for some time the reason ...
 - **Lepton Flavor is Conserved ...**



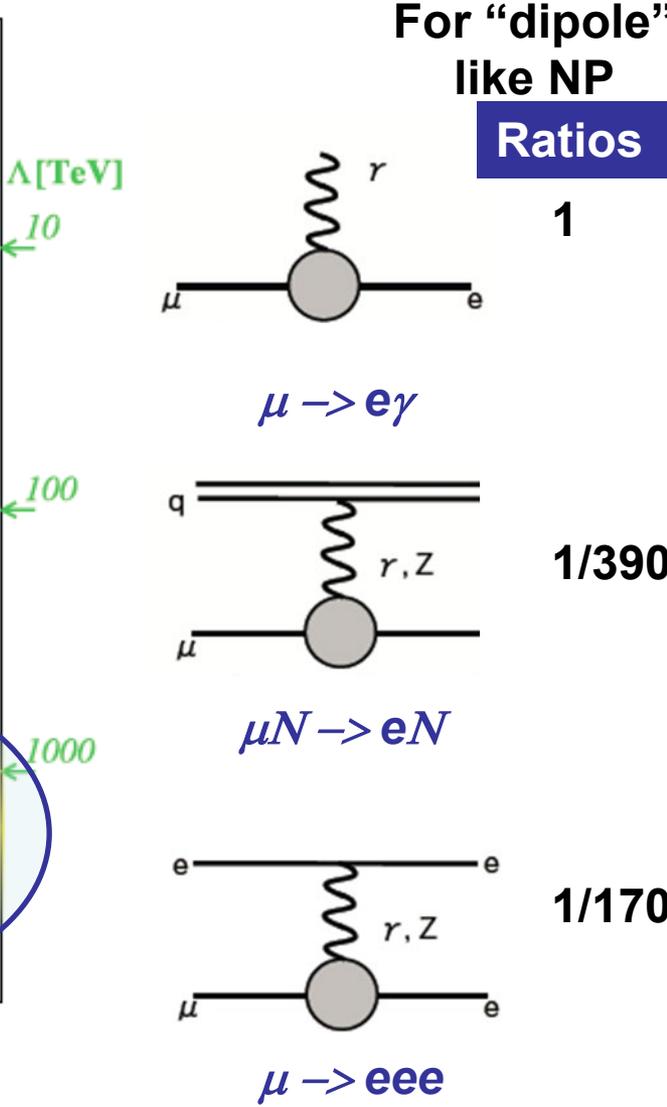
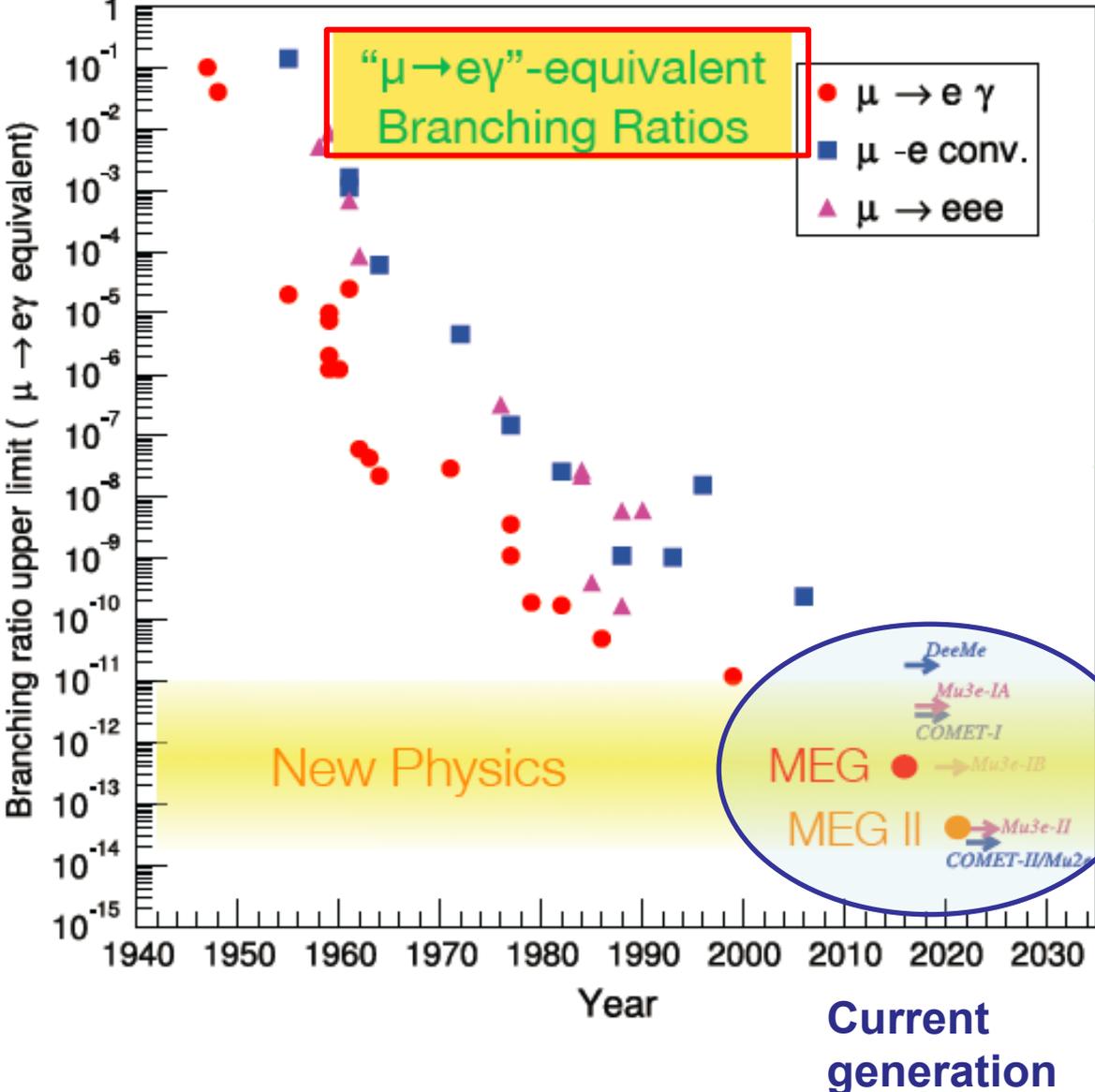
- **But we (now) know that neutrino flavors DO mix so ...it must be the case that $\mu \rightarrow e\gamma$ is not truly forbidden, right?**

$$\mathcal{B}(\mu \rightarrow e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_{i=2,3} U_{\mu i}^* U_{ei} \frac{\Delta m_{i1}^2}{M_W^2} \right|^2 \sim 10^{-54}$$

Right. It's just impossibly small

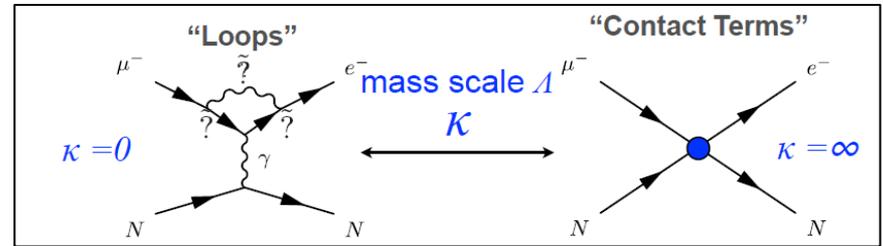
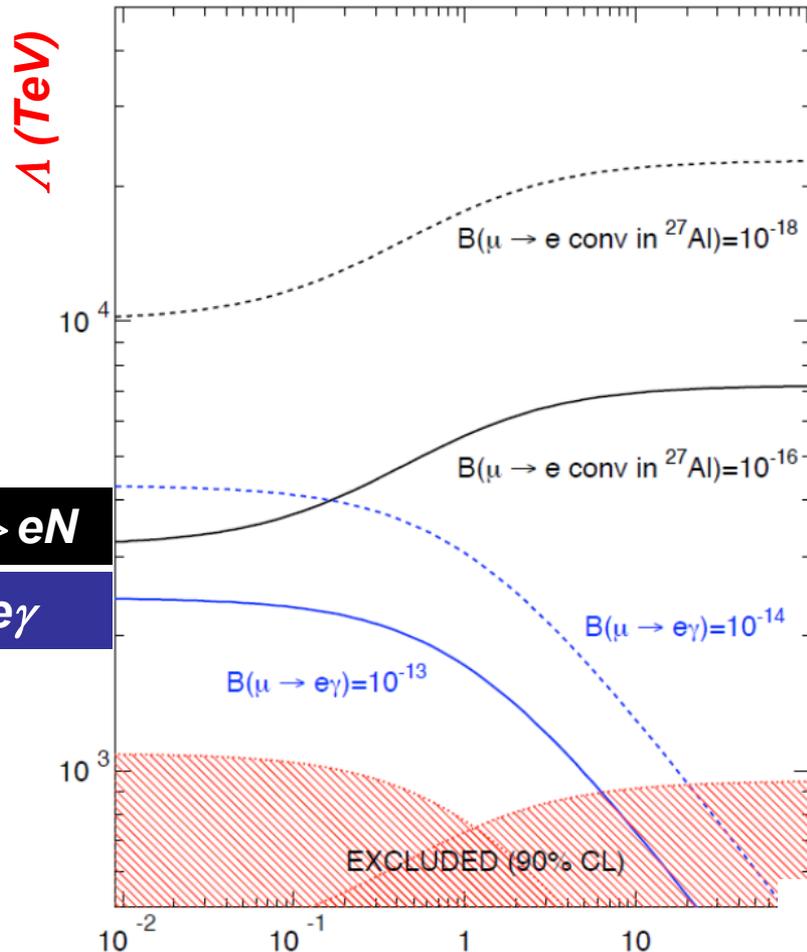
Charged lepton flavor violation limits are impressive

SM “allowed” but **Unobservable** e.g., $\mu \rightarrow e\gamma$ BR: 10^{-54}

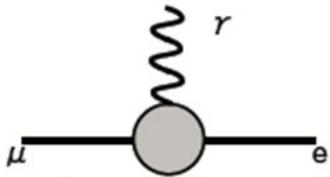


... but, life may not be just dipole

$$\mathcal{L}_{\text{CLFV}} = \frac{m_\mu}{(\kappa + 1)\Lambda^2} \bar{\mu}_R \sigma_{\mu\nu} e_L F^{\mu\nu} + \frac{\kappa}{(1 + \kappa)\Lambda^2} \bar{\mu}_L \gamma_\mu e_L (\bar{u}_L \gamma_\mu u_L + \bar{d}_L \gamma_\mu d_L)$$

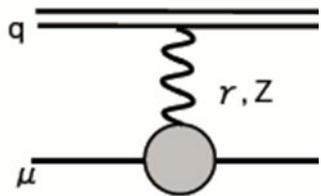


A comment about experimental design



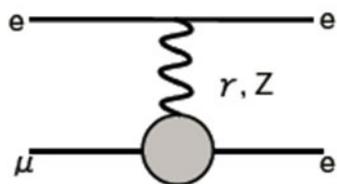
$\mu \rightarrow e\gamma$

- $\mu \rightarrow e\gamma$ is a coincidence experiment; you measure **e** and γ



$\mu N \rightarrow eN$

- $\mu N \rightarrow eN$ is a “singles” experiment. You only look for outgoing **e**



$\mu \rightarrow eee$

- $\mu \rightarrow eee$ is a triple coincidence experiment. You need to see all three **e** at once

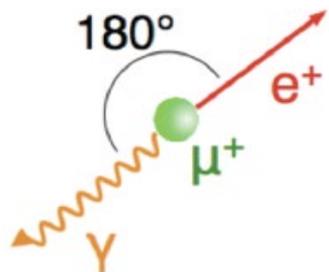
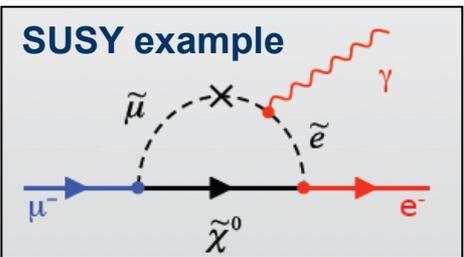
- **ALL THREE NEED HIGH STATISTICS**

A comment about experimental design

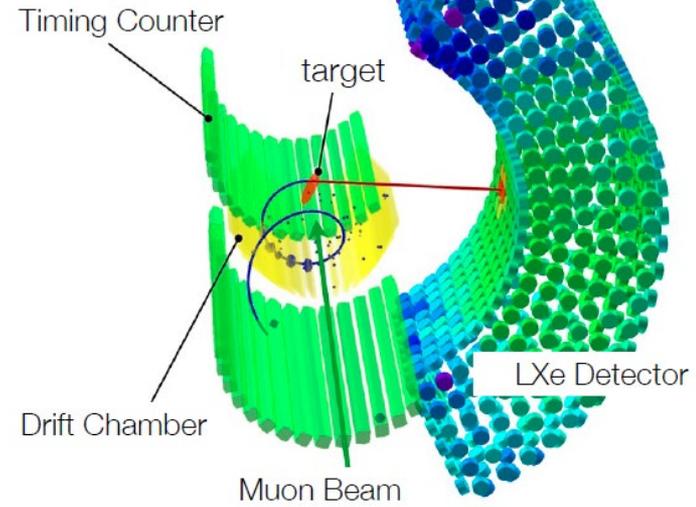
- Coincidence experiments need DC (continuous) beams to minimize PILEUP (event overlaps). Each event must be scrutinized to see if it follows the unique pattern: **PSI for $\sim 10^7 - 10^8 \mu^+$ /s beams**
- Monoenergetic emission experiment needs PULSED beam to avoid overlap of “beam” related background with “quiet” measuring period background; eg. FNAL pulsed every $1.7 \mu\text{s}$
 - It also allows for natural background subtraction periods
 - Examples from pulsed neutrino physics here ...

MEG: $\mu \rightarrow e\gamma$

Signal is back-to-back 53 MeV γ and e^+ from positive muons at rest



an example of events inside the blinded box



NEW 2016: $BR(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13}$ @ 90% CL

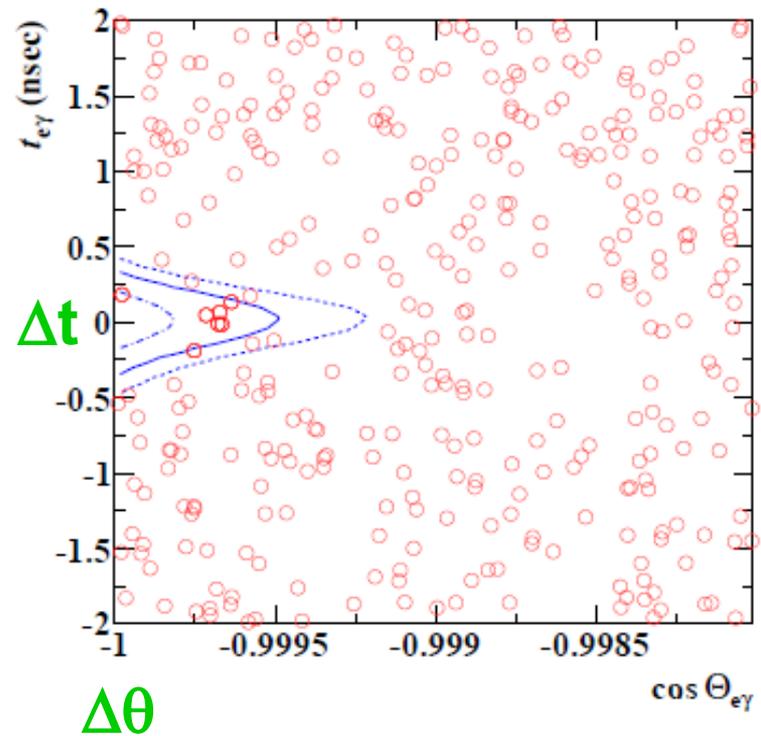
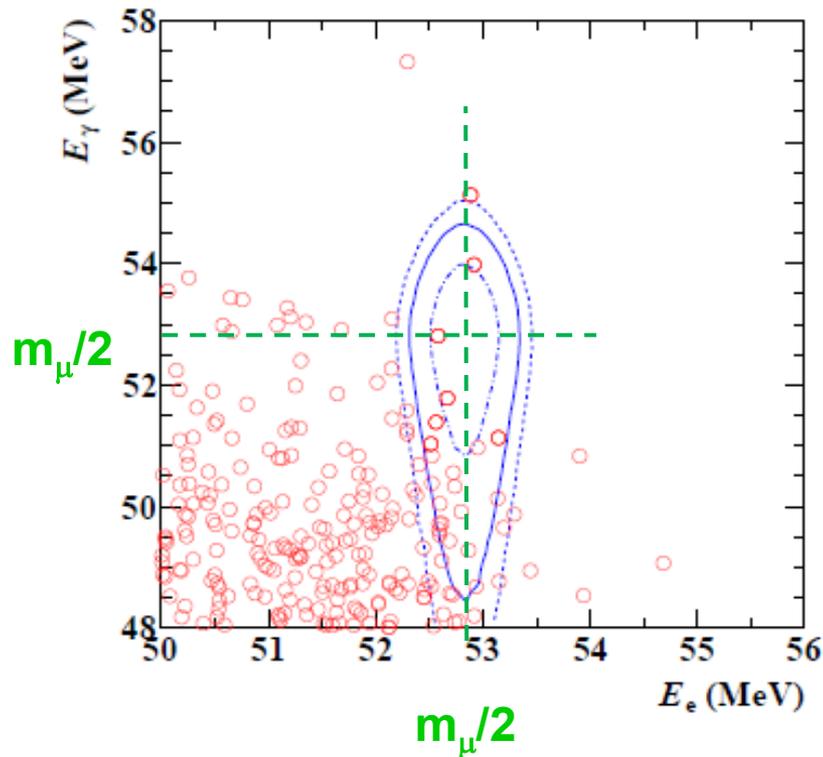
x30 improvement compared to pre-MEG

MEG II Upgrade approved at PSI: Expect to improve by another factor of 10 !

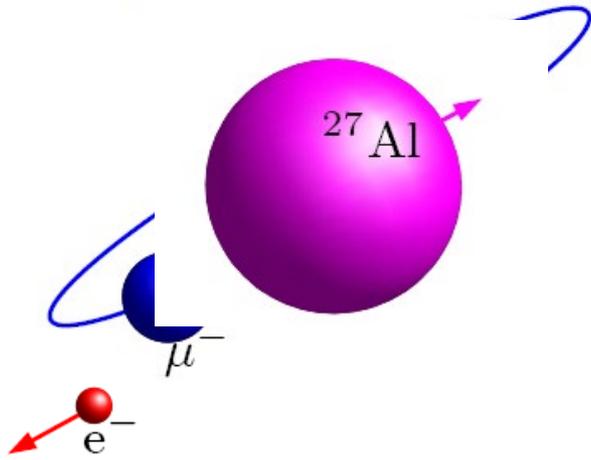
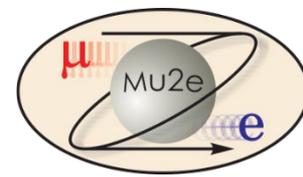
To give you a feeling: Event Search

- **e** and **γ** are back-to-back, $\Delta\theta = 180^\circ$
- **e** and **γ** are simultaneous, $\Delta t = 0$
- **$E_e = E_\gamma = m_\mu/2$**

4



Perhaps the most sensitive approach is coherent μ -to- e conversion



105 MeV
monoenergetic e-

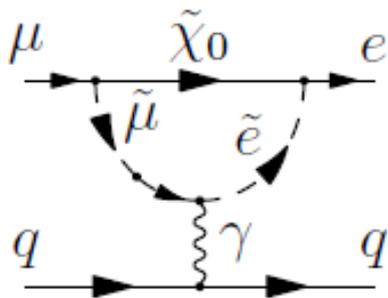
$$\mu^- N \rightarrow e^- N$$
$$R_{\mu e} = \frac{\Gamma(\mu^- + (A, Z) \rightarrow e^- + (A, Z))}{\Gamma(\mu^- + (A, Z) \rightarrow \nu_\mu + (A, Z - 1))}$$

- This signature is quite unique
- Goal $R_{\mu e}$ to $< 6 \times 10^{-17}$ (90% C.L.)
 - Present is $< 7 \times 10^{-13}$ → So this is very ambitious

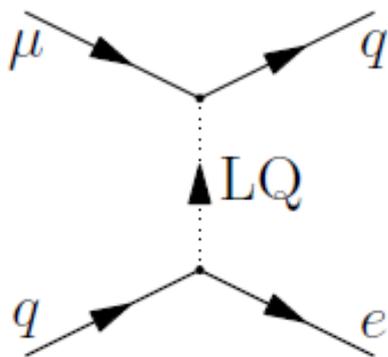
4 order of magnitude gain !!

What can Mu2e discover?

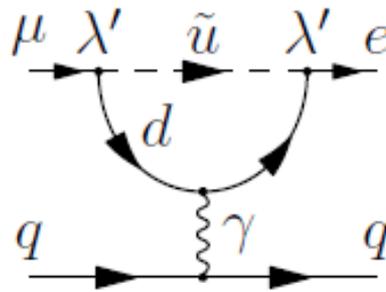
SUSY



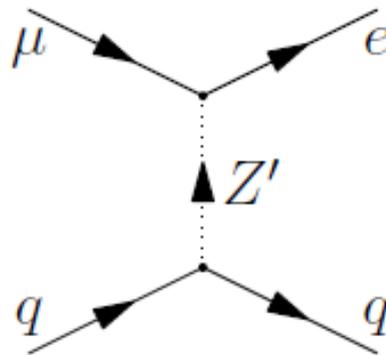
Leptoquarks



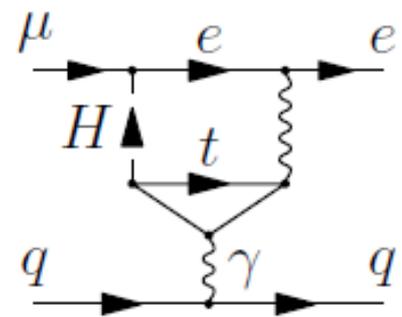
RPV SUSY



Z' /anomalous couplings



Second Higgs doublet



Extra dimensions, etc.

Theory reviews:

Y. Kuno, Y. Okada, 2001

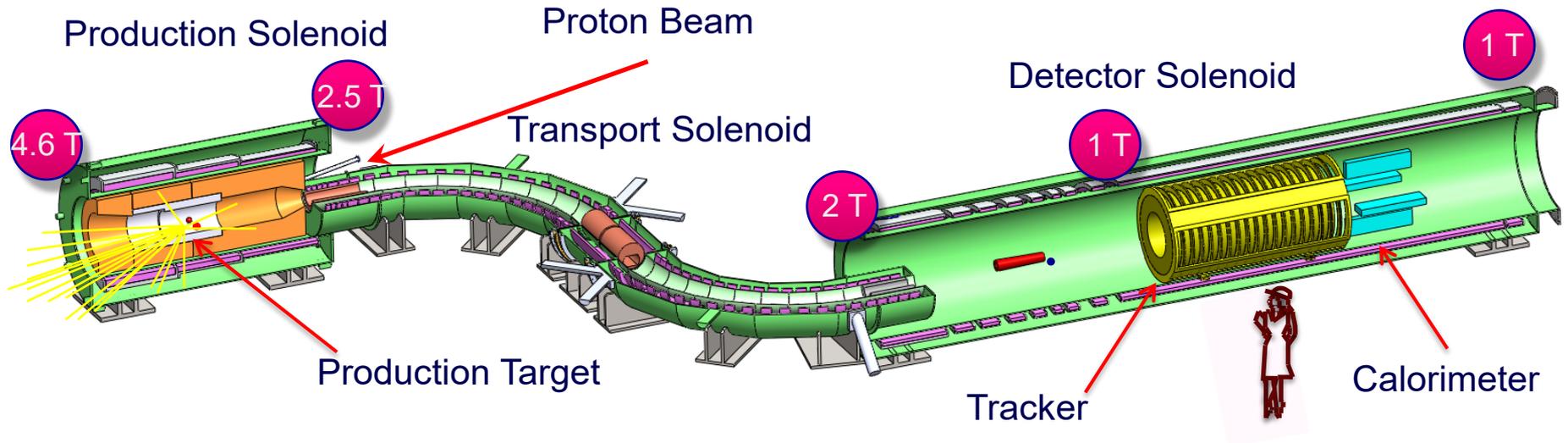
M. Raidal *et al.*, 2008

A. de Gouvêa, P. Vogel, 2013

How it is done

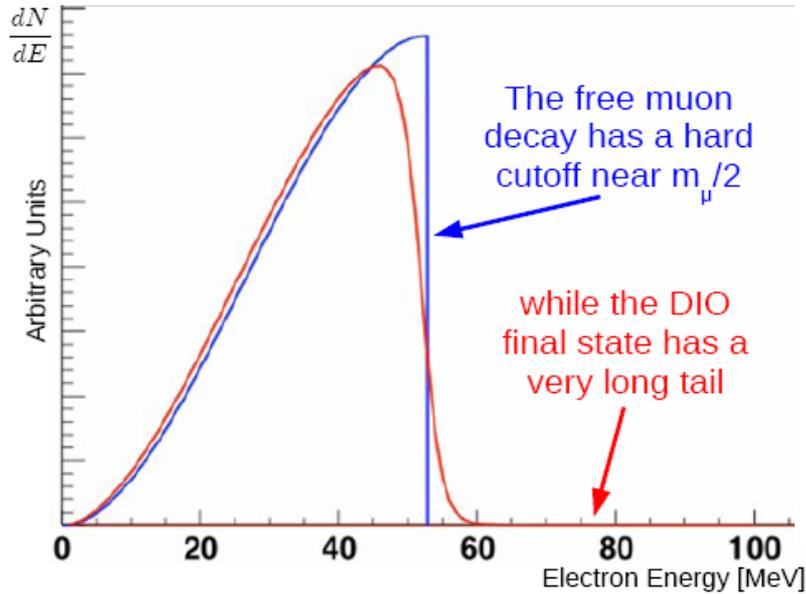


- Need intense pulsed source of low-energy muons
- Stop in thin Al target
- Form **muonic Al** atoms.
- Observe
 - 40% will decay “in orbit”;
 - 60% will capture (hadronic junk emitted)

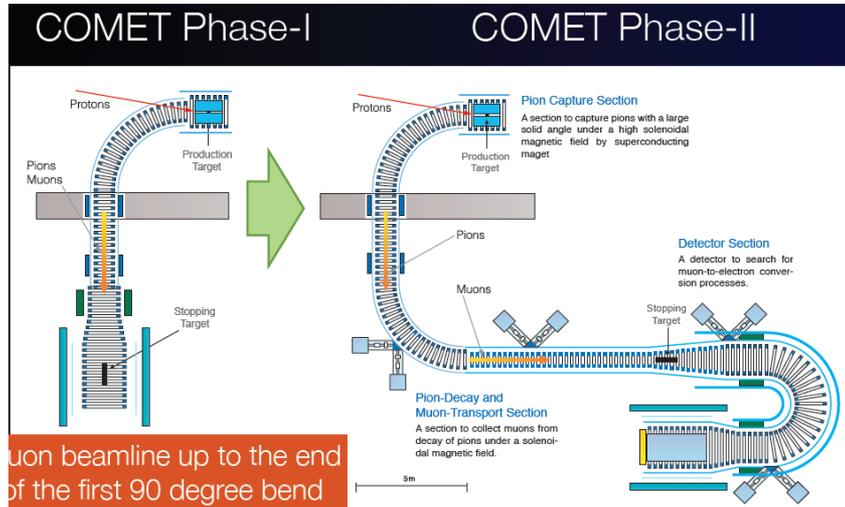
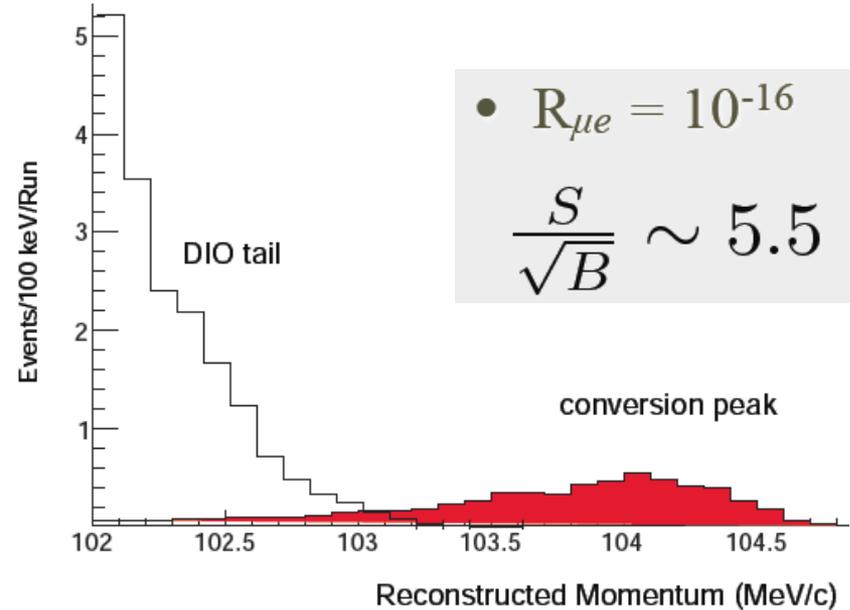


Challenge: find signal above "Decay in Orbit" tail

Tricky calculation; solved **Czarnecki et al**



Resolution and Redundancy critical

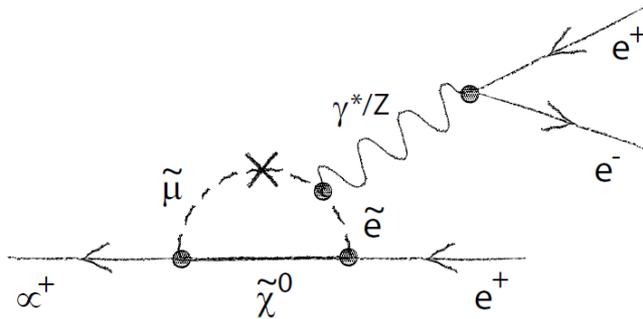
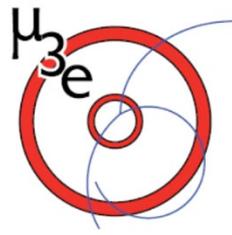


Similar: COMET in Japan

- Staged approach.
- Approved for Phase-1
 - Sensitivity: $< 7 \times 10^{-15}$
- Full phase later, similar to Mu2e

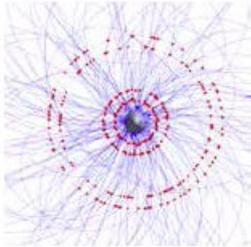
Next-generation: $\mu \rightarrow eee$

(2013: approved at PSI)

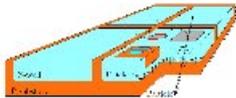


Typical comparison to $\mu \rightarrow e\gamma$ *without enhancement*

$$\frac{B(\mu \rightarrow eee)}{B(\mu \rightarrow e\gamma)} = 0.006 \quad (\text{essentially } \alpha_{em})$$



- **Goal:**
 - Finding **1 in 10^{16}** muon decays



- **Special technique**
 - High-voltage monolithic active pixel sensors

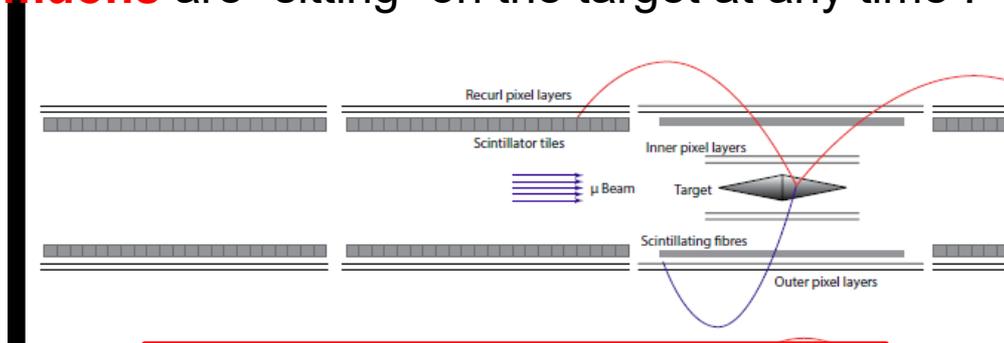


- **The detector**
 - Minimum material, maximum precision

Again, a unique and challenging signature

A staged approach is starting now

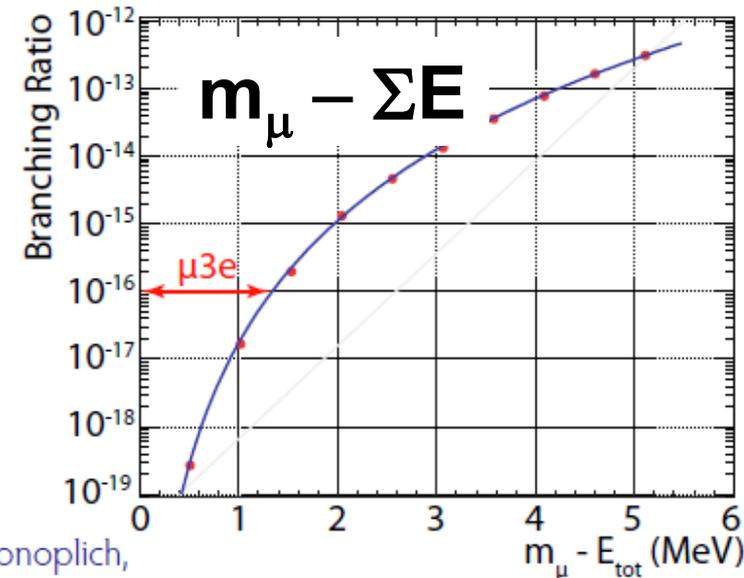
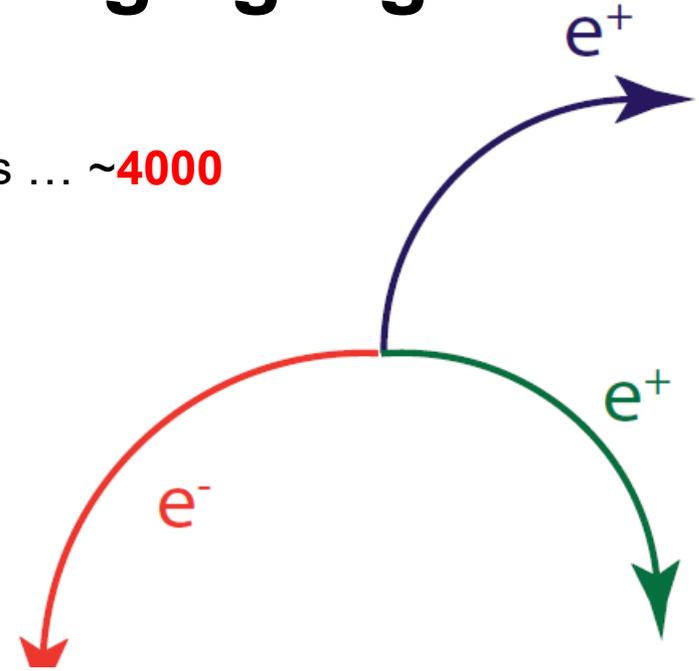
To achieve final statistics, extraordinary high rates ... **~4000 muons** are "sitting" on the target at any time !



10,000 times MuLan statistics

- **200 MHz muon rate**
- **Wow!!!**

- **2 e⁺, 1 e⁻**
- **Common vertex**
- **Common time**
- **Σ energies = m_μ**
- **No energy > m_μ / 2**



(onoplich, 2008) 0730004)

Topic 7

The Muon's Anomalous Magnetic Moment

(finally, something I am doing)

... *"our future discoveries must be looked for in the sixth place of decimals."*

It follows that every means which facilitates accuracy in measurement is a possible factor in a future discovery, and this will, I trust, be a sufficient excuse for bringing to your notice the various methods and results which form the subject matter of these lectures.

- Albert Abraham Michelson -

Dirac and beyond ...

$$i(\partial_\mu - ieA_\mu(x))\gamma^\mu\psi(x) = m\psi(x)$$

- 4-component (spinor) electron wave function Ψ in an EM potential (A_μ)
- Anticipates **antiparticles** (later found)
- Predicts **$g = 2$** , as observed in atomic fine-structure experiments for the **spin-1/2** electron magnetic moment (whereas an orbital picture $\rightarrow g = 1$)

$$\vec{\mu} = g \left(\frac{Qe}{2m} \right) \vec{s}, \quad e > 0$$

- Allows for a so-called **Pauli interaction** term to accommodate **deviations of g from 2** (as we will see are very important !)

At first, $g \approx 2$ was observed. But later, the proton ...

$$g_p = 5.59$$

and then the neutron

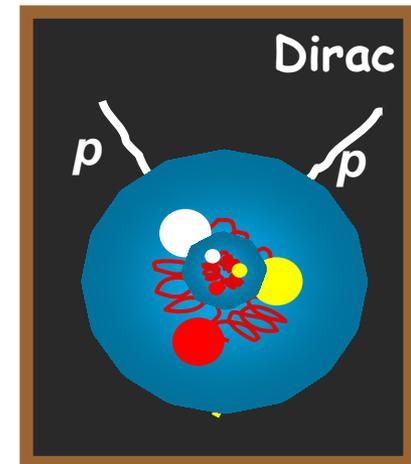
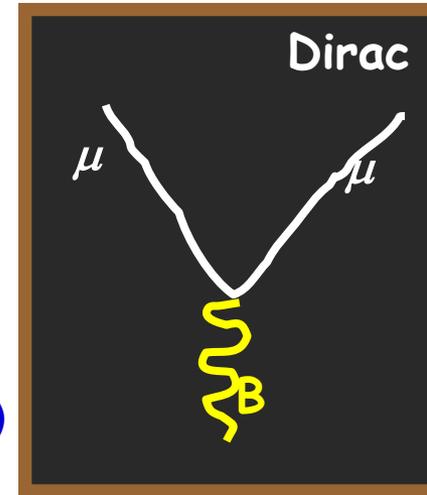
$$g_n = -3.8$$

each showed large magnetic moments ($g \neq 2$ by a lot)

The neutron? It's not even charged!

These are “Anomalous” magnetic moments owing to substructure

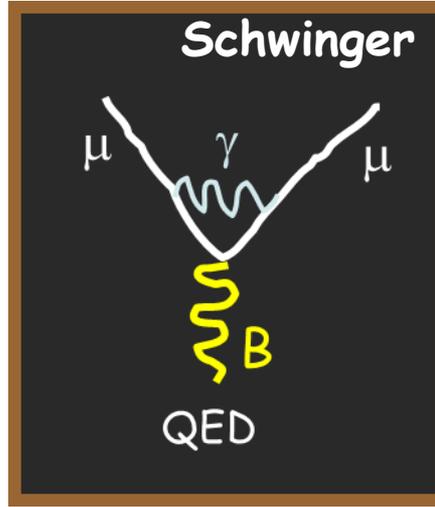
$$g = 2(1 + a) \quad \text{or} \quad a = \frac{(g - 2)}{2}$$



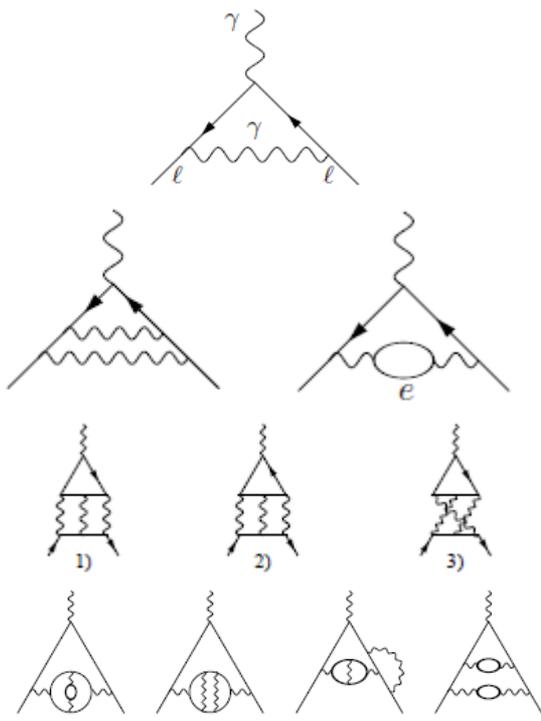
In 1947, small deviations from $g = 2$ for the “pointlike” electron were observed at about the $\sim 0.1\%$ level

What is that ?? $a_e = \frac{(g-2)}{2} \approx \frac{1}{2} \frac{\alpha}{\pi} \approx \frac{1}{800}$ ----->

- Schwinger calculates 1st order radiative correction
- It agrees with experiment
- Higher-order terms are expansions in powers of α/π
- The set of radiative terms, represents the QED anomalous magnetic moment contribution for the leptons



$$a = \sum_{j=1} C_j \left(\frac{\alpha}{\pi} \right)^j$$



Another story, but a_e is calculated so precisely (and accurately) that we obtain the best α from it:

$$\frac{1}{\alpha}(a_e) = 137.035\,999\,085\,(12)(8x)(33)$$

QED recent update, including tenth-order terms ! 12,672 diagrams

Complete Tenth-Order QED Contribution to the Muon $g - 2$

Tatsumi Aoyama,^{1,2} Masashi Hayakawa,^{3,2} Toichiro Kinoshita,^{4,2} and Makiko Nio²

¹*Kobayashi-Maskawa Institute for the Origin of Particles and the Universe (KMI), Nagoya University, Nagoya, 464-8602, Japan*

²*Nishina Center, RIKEN, Wako, Japan 351-0198*

³*Department of Physics, Nagoya University, Nagoya, Japan 464-8602*

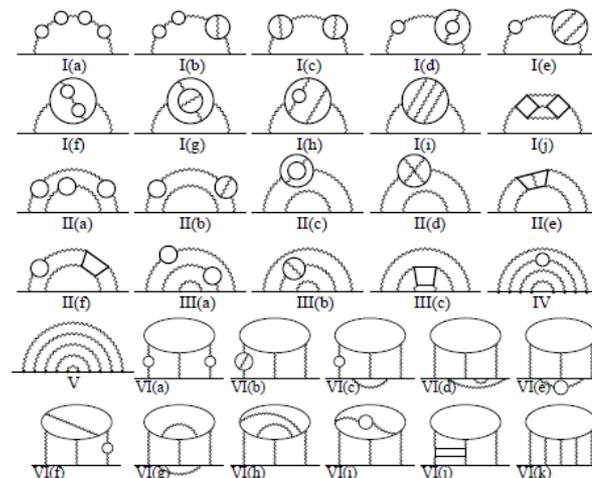
⁴*Laboratory for Elementary Particle Physics, Cornell University, Ithaca, New York, 14853, U.S.A*

(Dated: May 29, 2012)

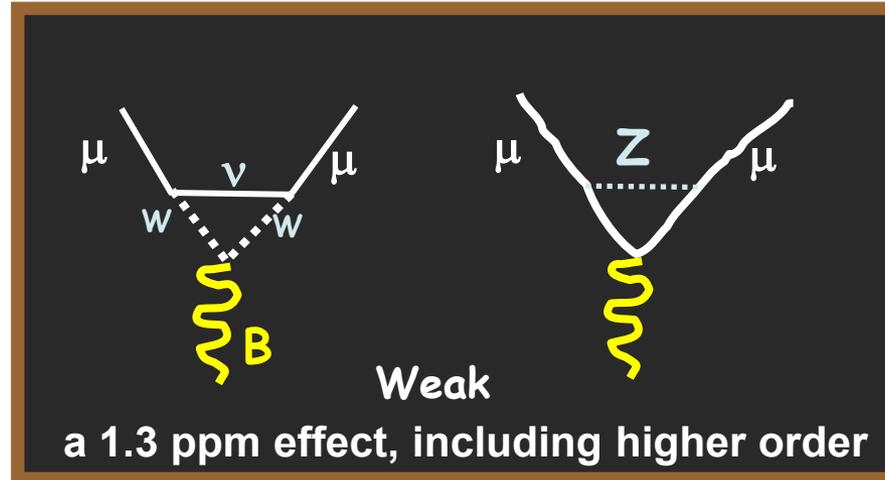
$$a_{\mu}(QED)^* = 116\,584\,718.09(14)(4)_{\alpha} \times 10^{-11}$$

Note: way better than expt.

Do not try to calculate these at home:



The Electroweak theory says, e.g., we can replace any γ with a Z ... and compute the **Weak** contribution to the anomaly

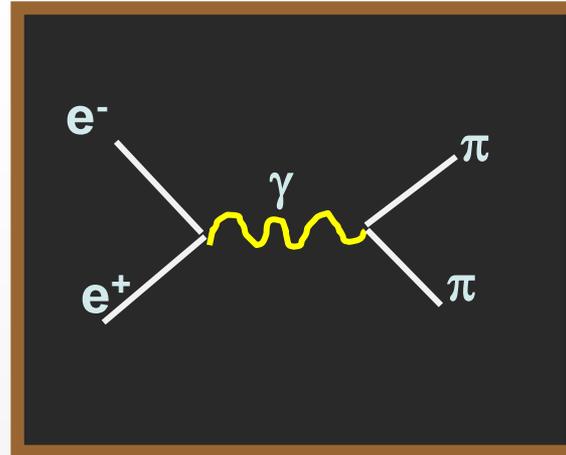
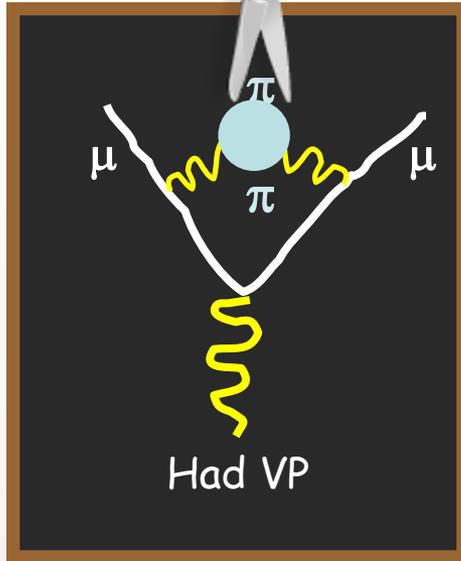


Known well, but wasn't easy

$$a_{\mu}(\text{Weak}) = 152(2)(1) \times 10^{-11}$$

Note: also way better than expt.

Hadronic vacuum polarization cannot be calculated using perturbation theory. **The strong coupling is too large**



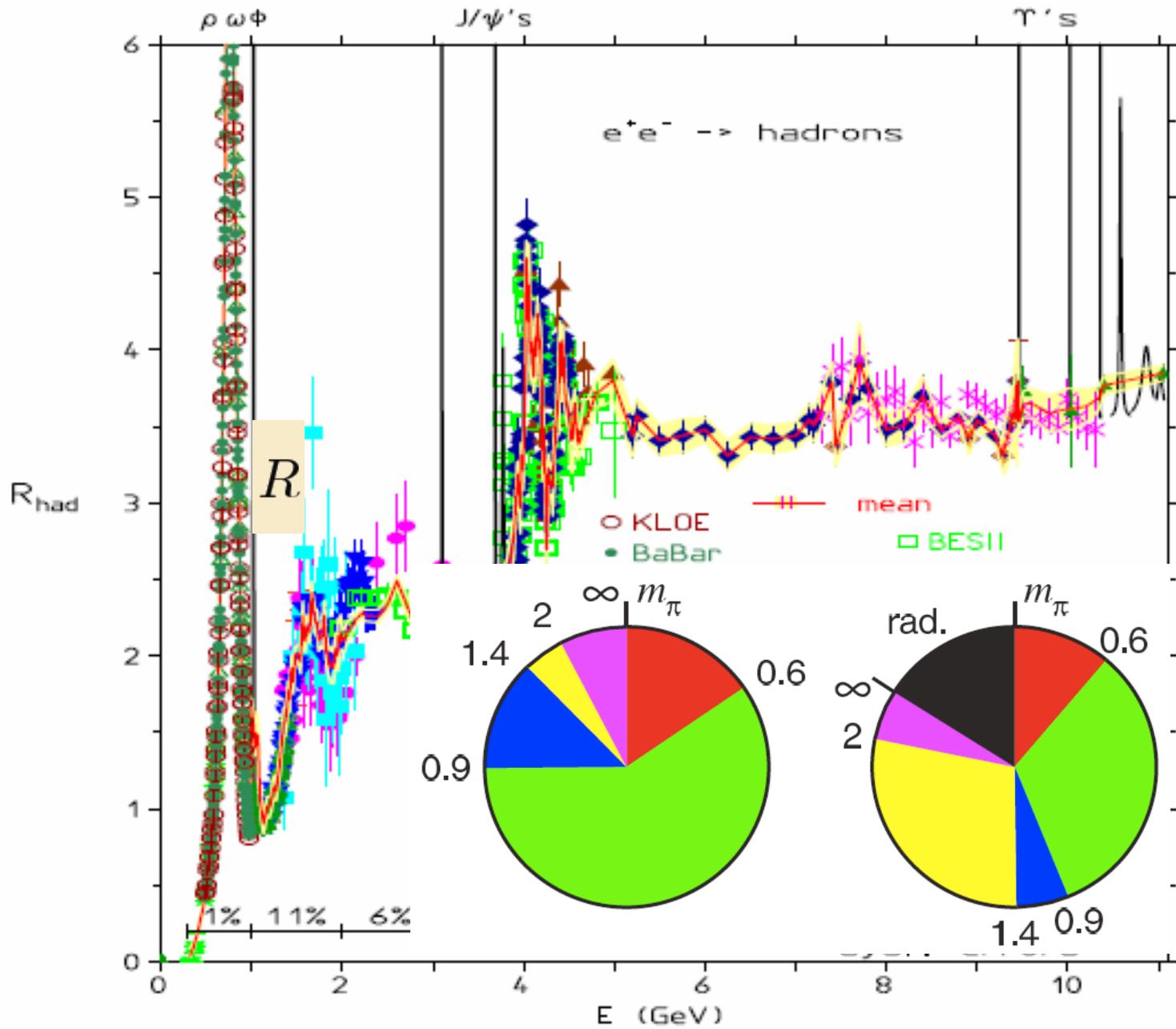
This contribution can be exactly linked to experimental data

1. Cut diagram down middle
2. Looks like $\gamma \rightarrow \pi\pi$
3. Dispersion relation connects $e^+e^- \rightarrow \pi\pi$ cross section measurement to anomaly contribution of 1st-order HVP

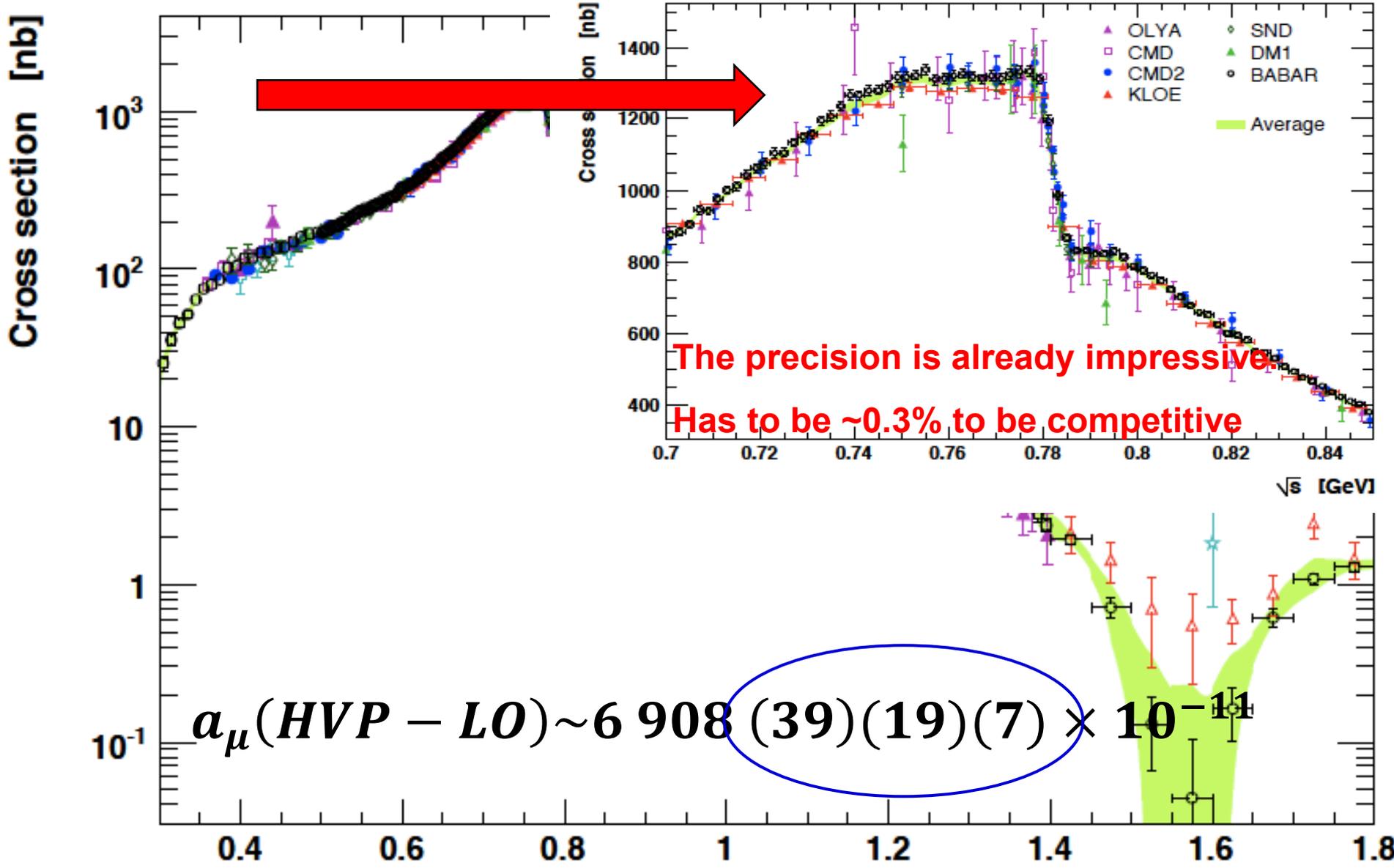
$$a_{\mu}^{\text{had,LO}} = \frac{\alpha^2(0)}{3\pi^2} \int_{4m_{\pi}^2}^{\infty} ds \frac{K(s)}{s} R(s)$$

$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \text{muons})}$$

The cross sections scan a wide range in energy

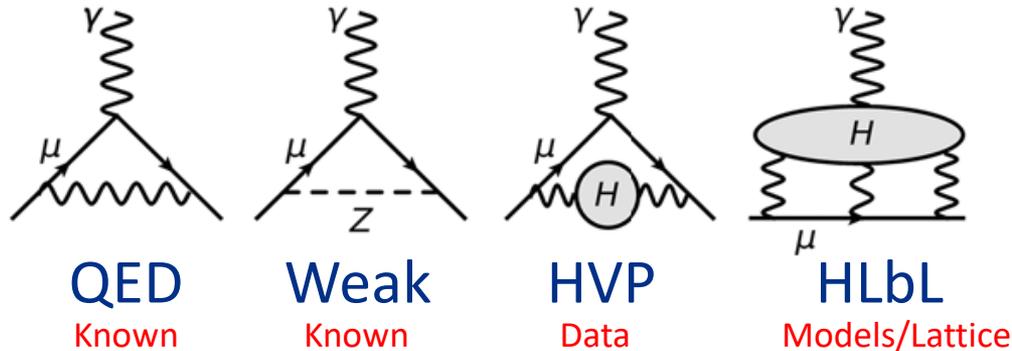


Most important is low energy: $e^+e^- \rightarrow \pi^+ \pi^-$



Note: Largest error, but can be reduced by new experiments

Standard Model contributions to a_μ ... updates $\rightarrow 3.6 \sigma$

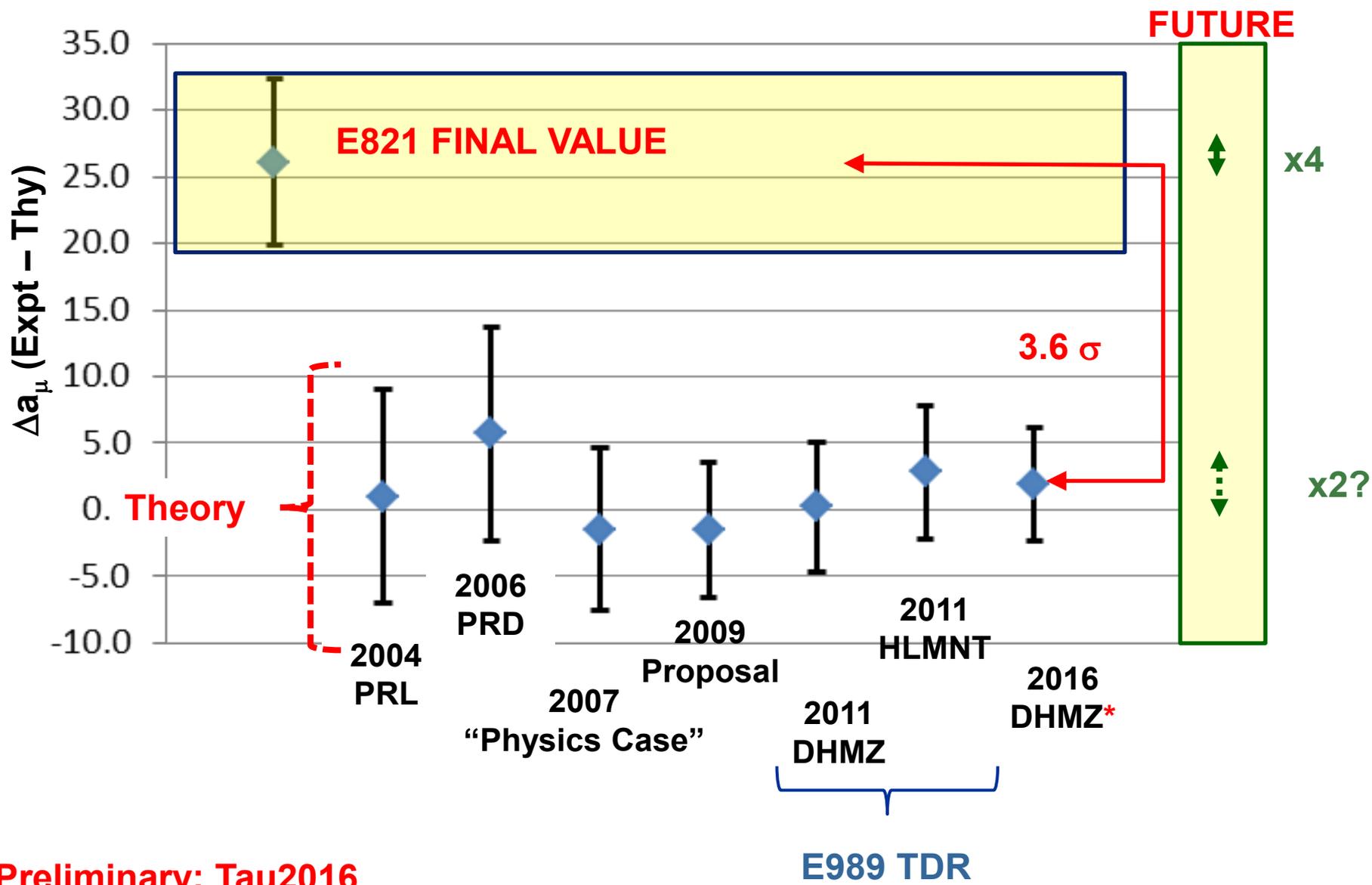


	VALUE ($\times 10^{-10}$) UNITS
QED ($\gamma + \ell$)	$11\,658\,471.8951 \pm 0.0009 \pm 0.0019 \pm 0.0007 \pm 0.0077_\alpha$
HVP(lo) Davier17	692.6 ± 3.33
HVP(lo)KNT2017	693.9 ± 2.6
HVP(ho) KNT2017	-9.84 ± 0.07
HLbL Glasgow	10.5 ± 2.6
EW	15.4 ± 0.1
Total SM Davier17	$11\,659\,181.7 \pm 4.2$
Total SM KNT17	$11\,659\,182.7 \pm 3.7$

\leftarrow This is a fancy guess; it will change \rightarrow

BNL E821 $\delta a_\mu(\text{Expt}) = \pm 6.3$

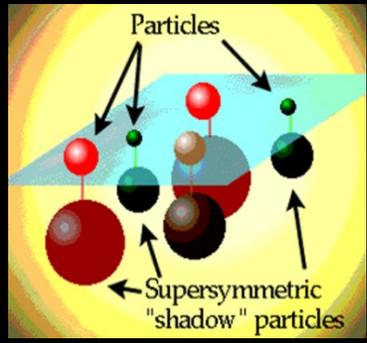
In the 12 years since BNL E821, the “g-2 Test” has continued to point to something interesting



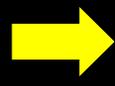
*Preliminary; Tau2016

What could it mean now?

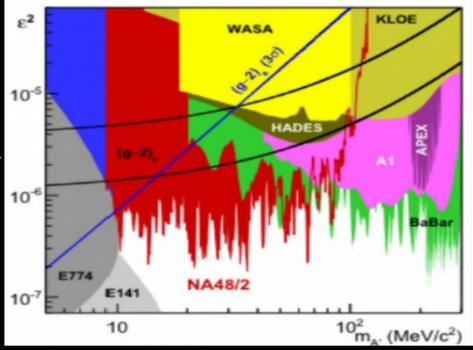
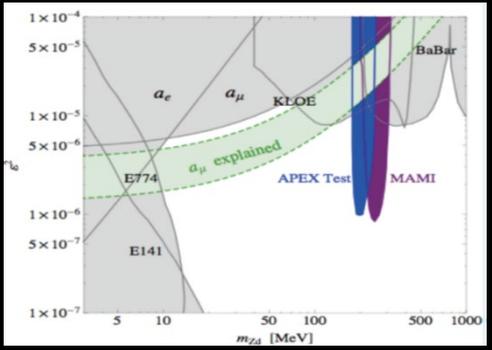
Some things "seen" just wash away ... And some things are under Tension



LHC limits growing, but SUSY, if exists, is hiding well



And some things don't seem to be so ...

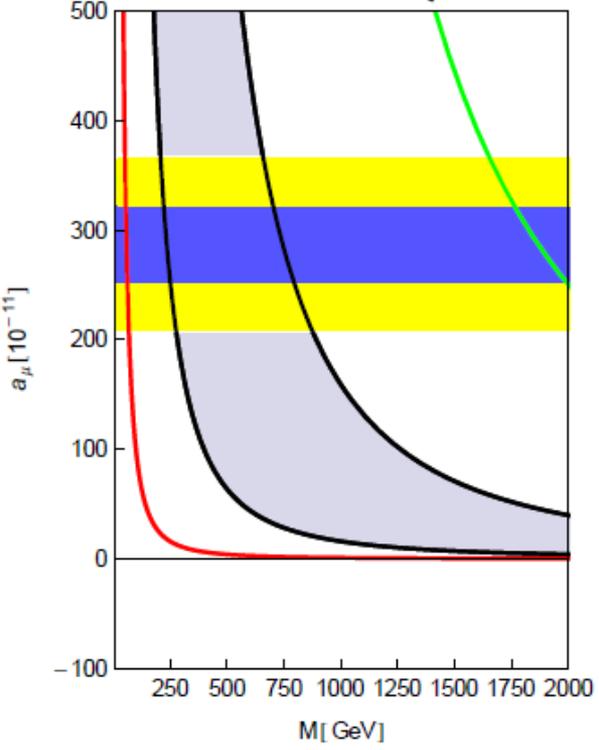
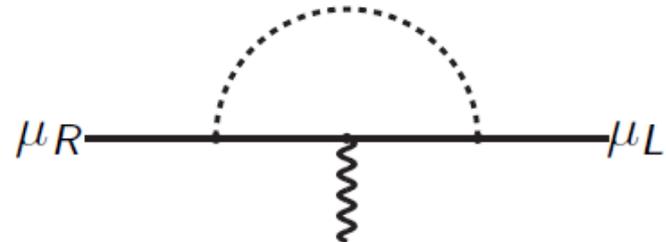


g-2: An uncomfortably lonely search for a Crack in the SM

In a generic sense, these are “loop effects” that couple to the muon mass and moment in similar fashion, characterized by C , a coupling:

$$\mathcal{O}(C) \left(\frac{m_\mu}{M}\right)^2$$

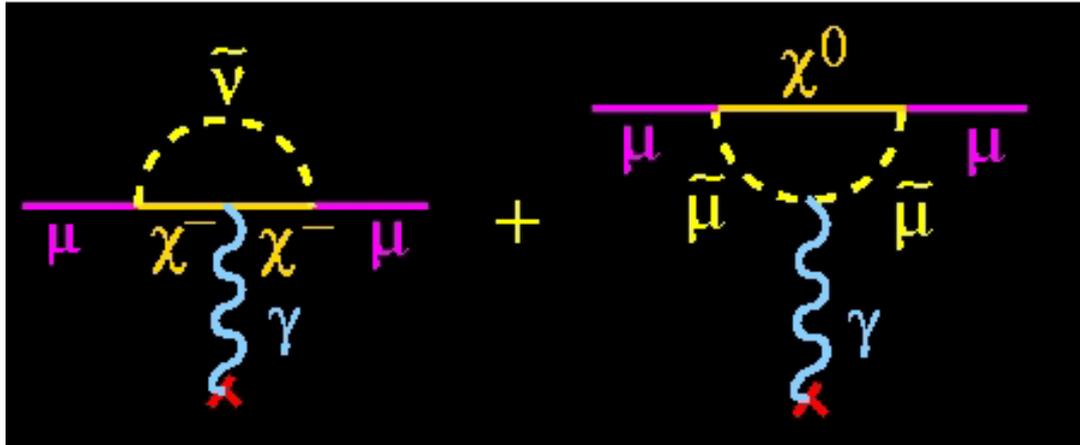
$$C = \frac{\delta m_\mu(\text{N.P.})}{m_\mu}$$



$\mathcal{O}(1)$	radiative muon mass generation ... [Czarnecki, Marciano '01] [Crivellin, Girschbach, Nierste '11][Dobrescu, Fox '10]
$\mathcal{O}\left(\frac{\alpha}{4\pi} \dots\right)$	supersymmetry ($\tan \beta$) vectorlike fermions ...
$\mathcal{O}\left(\frac{\alpha}{4\pi}\right)$	SM: Z, W . New physics: $Z', W' \dots$
$< \frac{\alpha}{4\pi}$	2-Higgs doublet model, dark photon .

Following Czarnecki, Marciano, and Stockinger

SUSY contribution to a_μ :



Difficulty to measure at the LHC

$$a_\mu^{\text{SUSY-11}} \approx 130 \times 10 \tan \beta \operatorname{sign} \mu \left(\frac{100 \text{ GeV}}{M_{\text{SUSY}}} \right)^2 \quad ()$$

Recall, the deviation between Experiment and Theory is $\sim 280 \times 10^{-11}$, so the above calculation is interesting if you put in M_{SUSY} , and $\tan \beta$

$\tan \beta$? Ratio of the two vacuum values of the 2 neutral Higgses, typically estimated in range from **3 to 55**

A few key numbers determine the precision of the $g-2$ Test:

$$a_\mu(\text{New Physics}) \equiv a_\mu(\text{Expt}) - a_\mu(\text{SM})$$

- $a_\mu(\text{SM}) = a_\mu(\text{QED}) + a_\mu(\text{Weak}) + a_\mu(\text{HVP}) + a_\mu(\text{Had HO}) + a_\mu(\text{HLbL})$

✓ ✓ a_μ(HVP) ✓ a_μ(HLbL)

A few remarks here

In E821 $\equiv \mathcal{R}_\mu(\text{E821}) = 0.003\,707\,206\,4(20)$ [0.54 ppm]

- $a_\mu(\text{Expt}) = \frac{g_e}{2} \frac{\omega_a}{\tilde{\omega}_p} \frac{m_\mu}{m_e} \frac{\mu_p}{\mu_e}$

ω_a
m_μ
μ_p
ω̃_p
m_e
μ_e

-2.002 319 304 361 53(53) [0.26 ppt]
Electron $g-2$ + QED

-.001519270384(12) [8 ppb]
206.768 2843(52) [25 ppb]

Spin motion for a particle *moving* in a magnetic field

$$\omega_C = \frac{eB}{mc\gamma}$$

Momentum turns at ω_c
cyclotron frequency

$$\omega_S = \frac{geB}{2mc} + (1 - \gamma)\frac{eB}{\gamma mc}$$

Spin turn depends on g
and on ω_c with $1-\gamma$ factor

If $g = 2$ exactly, then the difference between SPIN and MOMENTUM vectors

$$\omega_S - \omega_C = \left[\frac{eB}{mc} + \frac{eB}{mc\gamma} - \frac{eB}{mc} \right] - \frac{eB}{mc\gamma} = 0$$

Spin motion for a particle *moving* in a magnetic field

$$\omega_S = \frac{geB}{2mc} + (1 - \gamma) \frac{eB}{\gamma mc} \qquad \omega_C = \frac{eB}{mc\gamma}$$

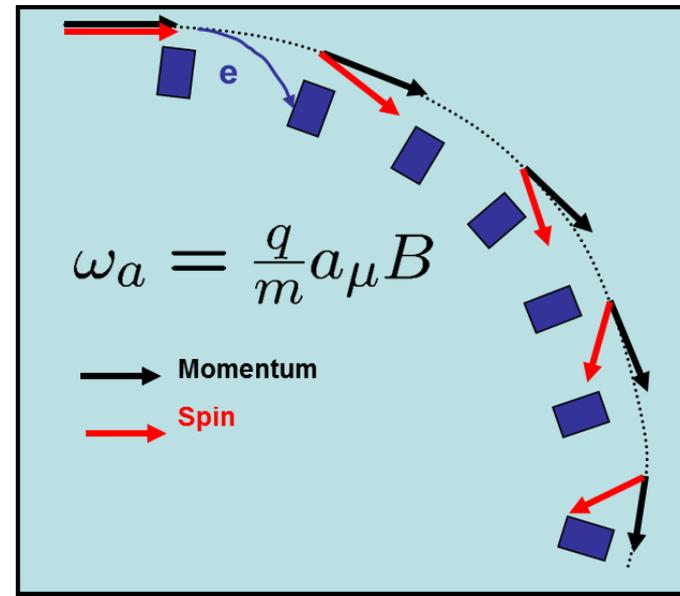
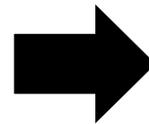
The **Spin** frequency relative to the **Cyclotron** frequency is the “anomalous precession frequency”, ω_a

Does **NOT** depend on γ !

Proportional to $g - 2$ and B !

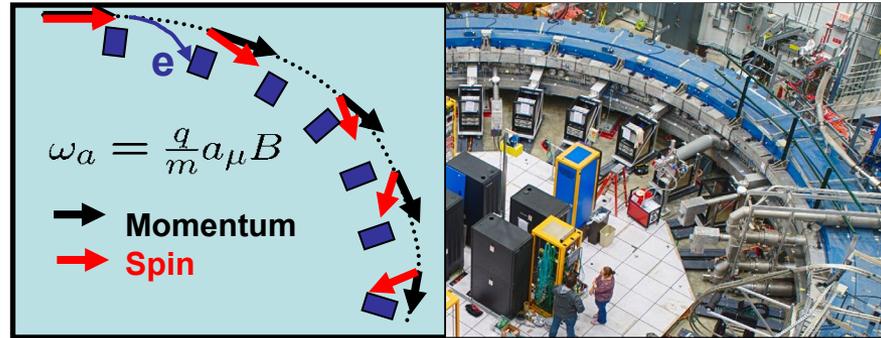
$$\omega_a = \omega_S - \omega_C$$

$$= \left(\frac{g - 2}{2} \right) \frac{eB}{mc} = a \frac{eB}{mc}$$



Measurement of Muon g-2 and μ EDM

Goal: 140 ppb
X 4 improvement



Determine difference between spin precession and cyclotron motion for a muon moving in a magnetic field:

The expression including E-field focusing and possible EDM

$$\vec{\omega}_{net} = -\frac{q}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} + \frac{\eta}{2} \left(\vec{\beta} \times \vec{B} + \frac{\vec{E}}{c} \right) \right]$$

↑ Measure these

Get a_μ

Magic γ

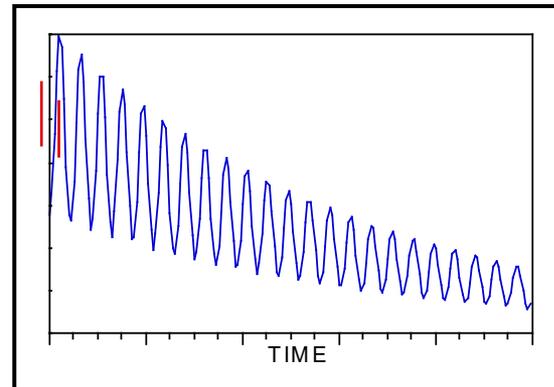
EDM

$$\vec{\omega}_{net} = \vec{\omega}_a + \vec{\omega}_{EDM}$$

Two “blinded” frequency measurements are made. The ratio gives $a_\mu \equiv (g-2)/2$

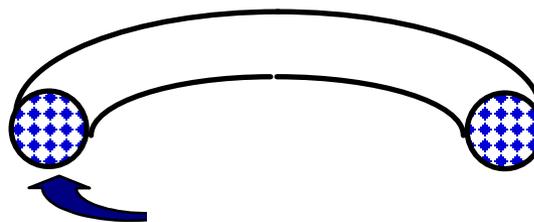
(1) Precession frequency

(1) Calorimeters



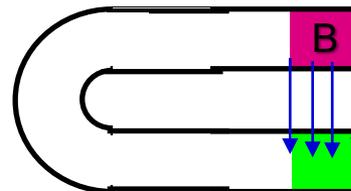
(2) Muon distribution

(2) Trackers & Models



(3) Magnetic field

(3) proton pNMR



$$(g-2) \propto \frac{(1)}{\langle \int (2)(3) \rangle}$$

How do we get each of these?

Systematic error projections in-line with statistical goal

Precession

E821 Error	Size [ppm]	Plan for the New ($g - 2$) Experiment	Goal [ppm]
Gain changes	0.12	Better laser calibration and low-energy threshold	0.02
Lost muons	0.09	Long beamline eliminates non-standard muons	0.02
Pileup	0.08	Low-energy samples recorded; calorimeter segmentation	0.04
CBO	0.07	New scraping scheme; damping scheme implemented	0.04
E and pitch	0.05	Improved measurement with traceback	0.03
Total	0.18	Quadrature sum	0.07

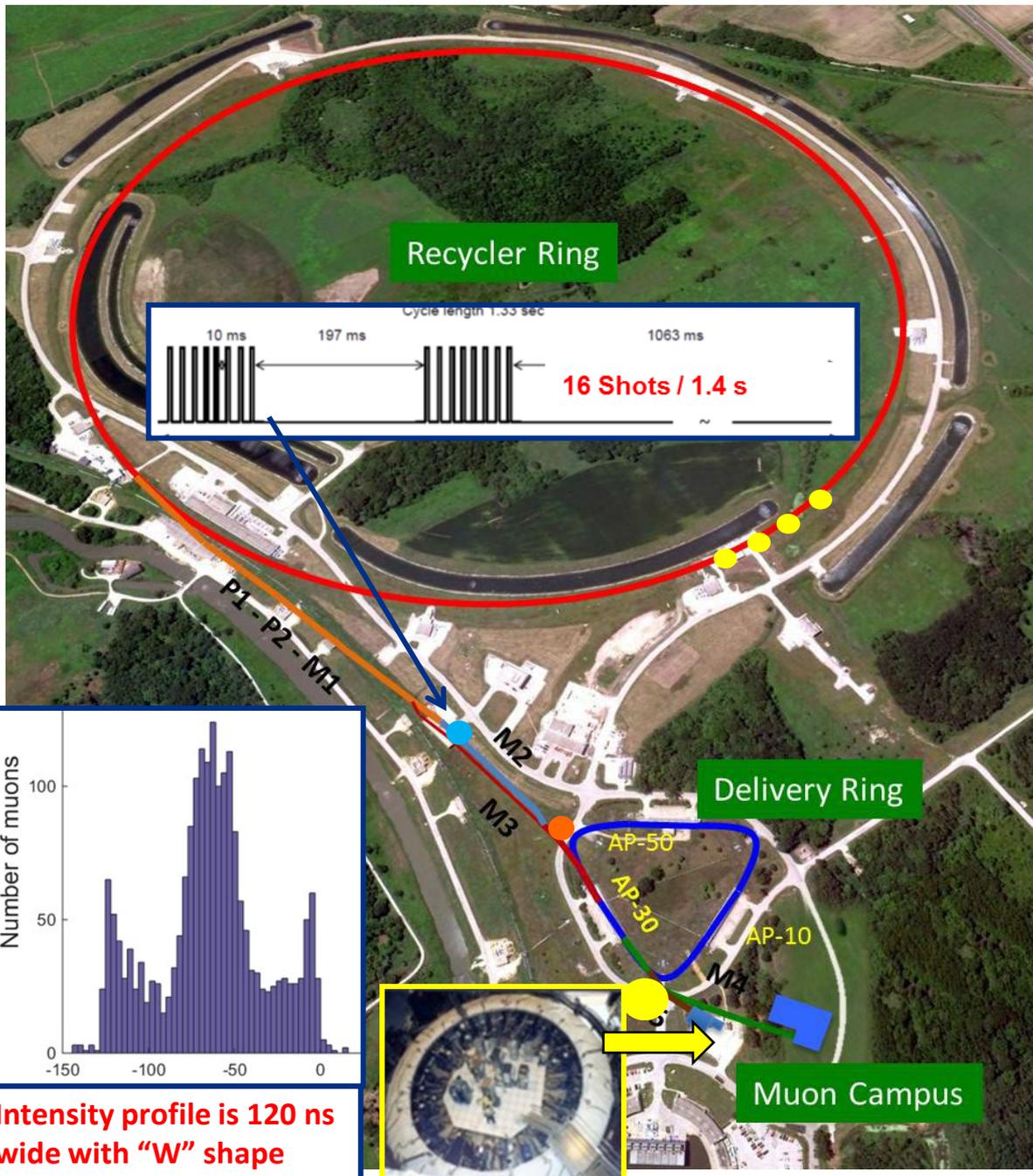
Magnetic field

Source of errors	Size [ppm]				
	1998	1999	2000	2001	future
Absolute calibration of standard probe	0.05	0.05	0.05	0.05	0.05
Calibration of trolley probe	0.3	0.20	0.15	0.09	0.06
Trolley measurements of B_0	0.1	0.10	0.10	0.05	0.02
Interpolation with fixed probes	0.3	0.15	0.10	0.07	0.06
Inflector fringe field	0.2	0.20	-	-	-
Uncertainty from muon distribution	0.1	0.12	0.03	0.03	0.02
Others		0.15	0.10	0.10	0.05
Total systematic error on ω_p	0.5	0.4	0.24	0.17	0.11

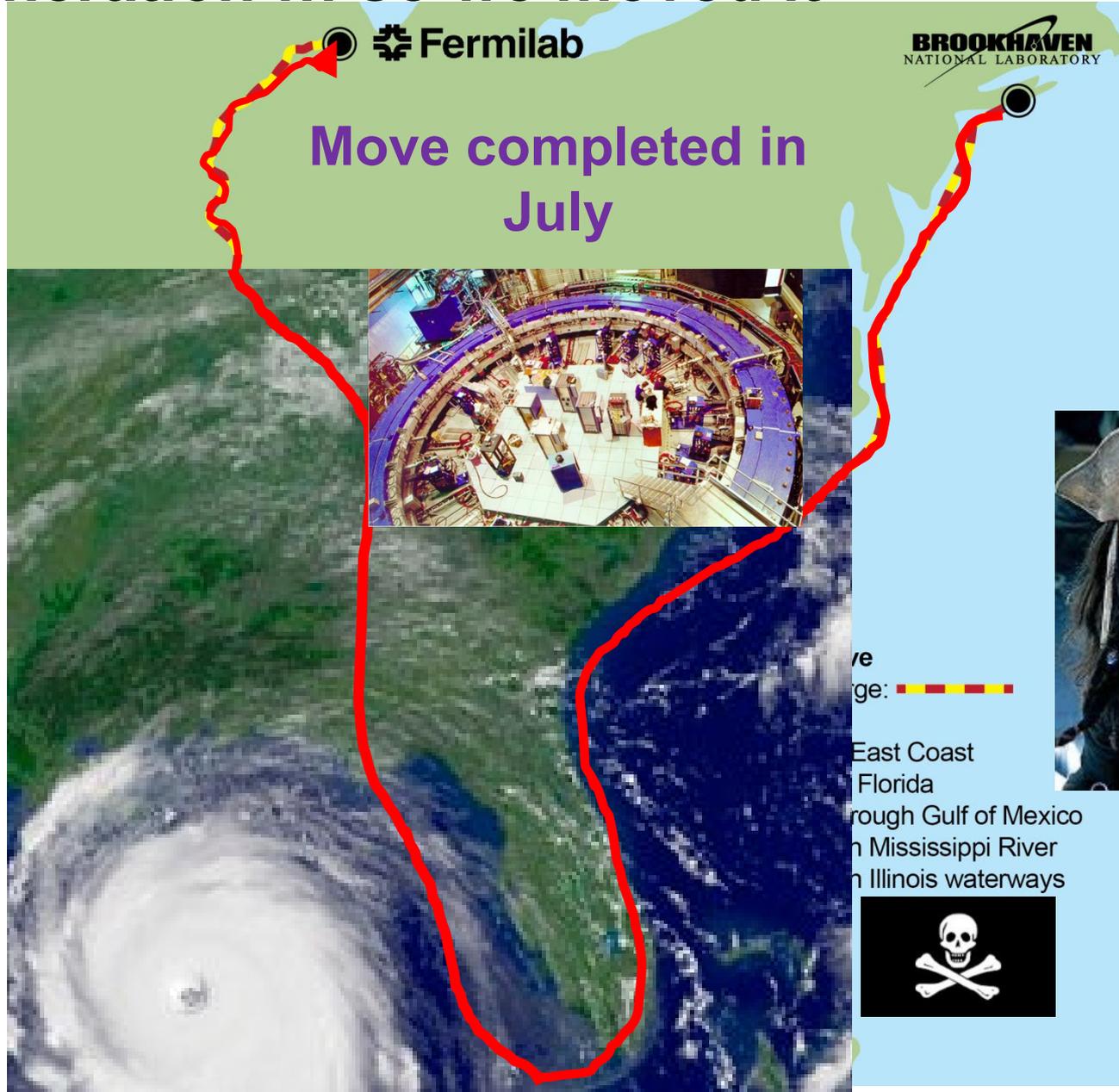
Improvement vs time →

Creating the Muon Beam for g-2

- 8 GeV p batch into Recycler
- Split into 4 bunches
- Extract 1 by 1 to strike target
- Long FODO channel to collect $\pi \rightarrow \mu\nu$
- $\rho/\pi/\mu$ beam enters DR; protons kicked out; π decay away
- μ enter storage ring



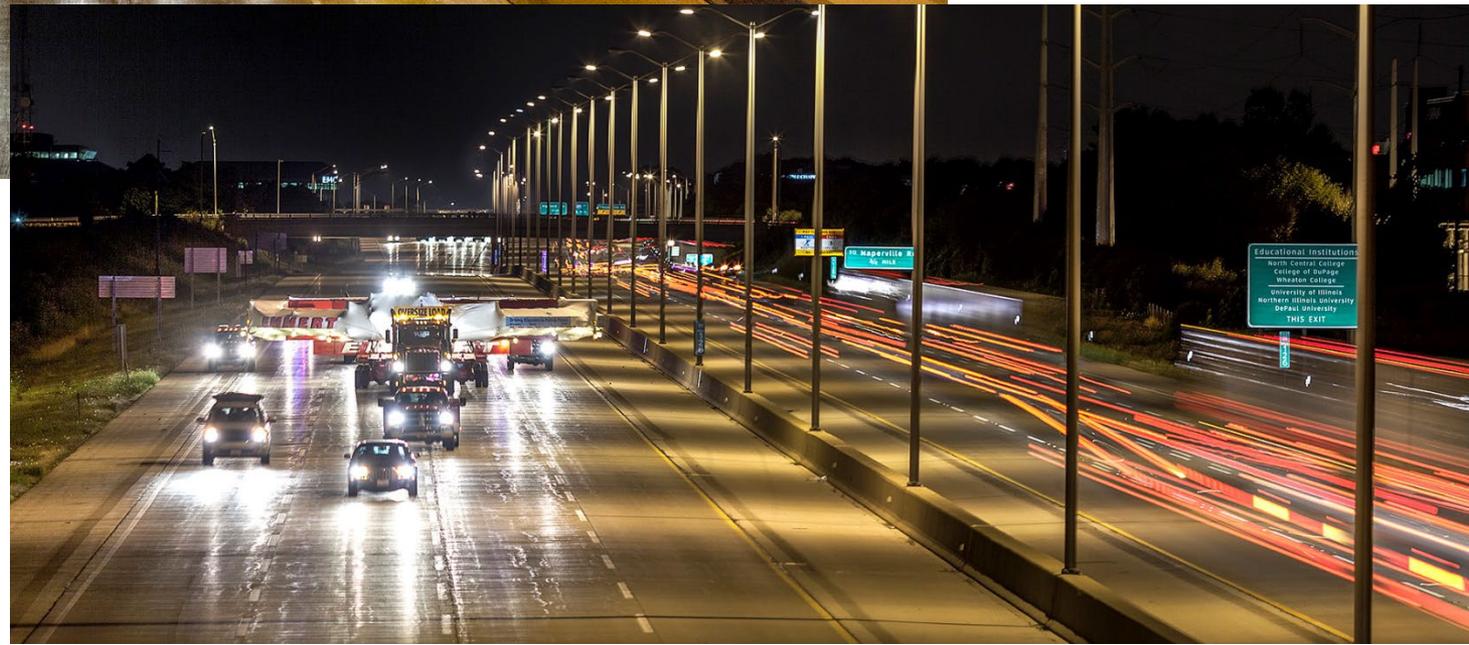
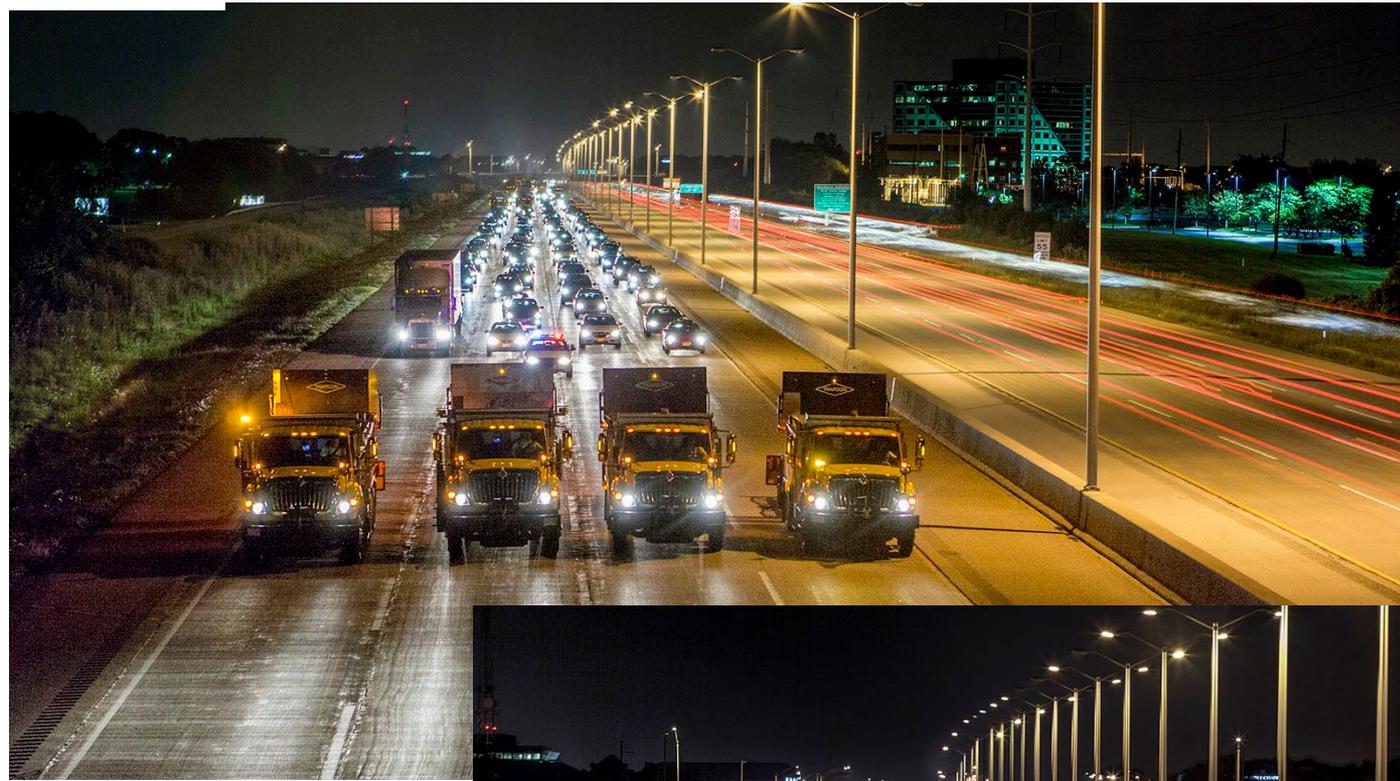
The **MAGNET** is the centerpiece and worthy of the next generation ... so we moved it



Leaving BNL and loading the barge



30 police cars escort it and close interstate

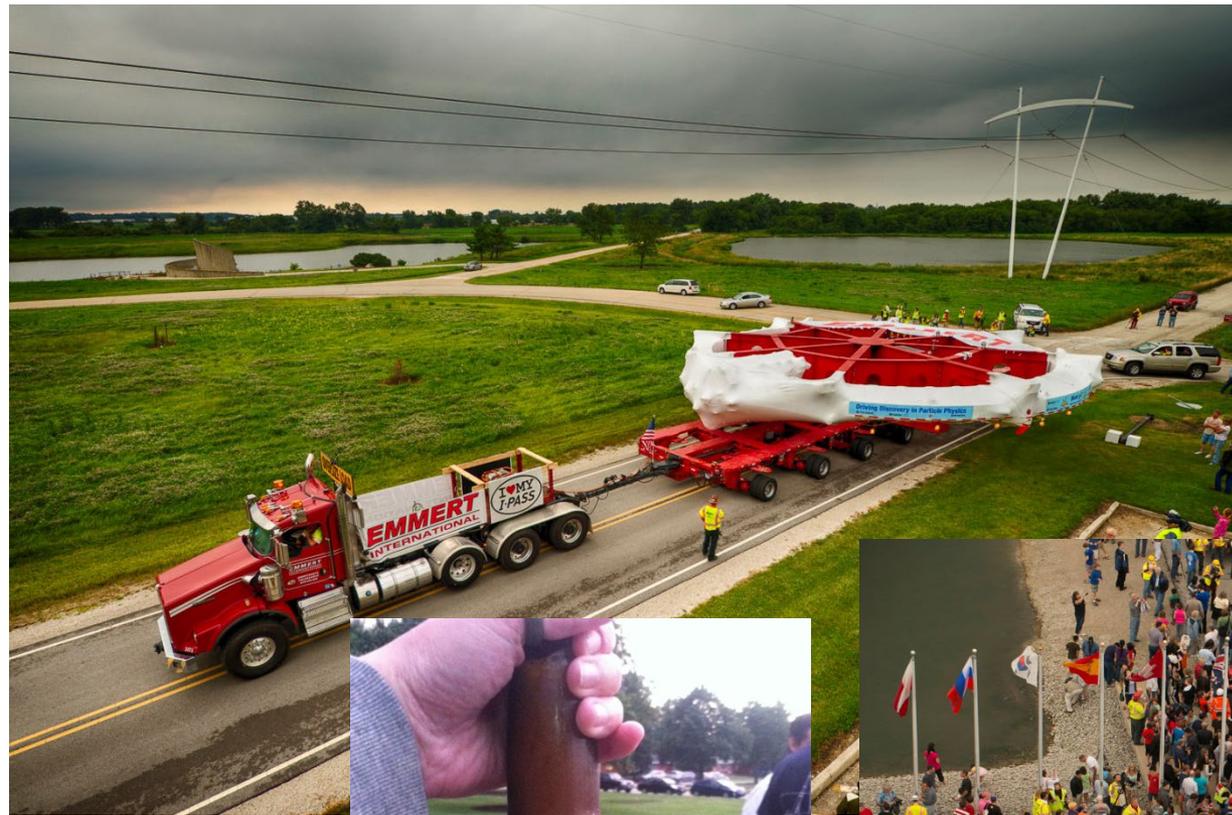
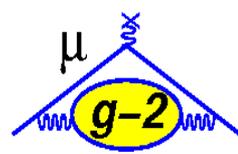


Squeezing through the I-355 tollbooth and a tight underpass





Arriving at FNAL to a huge crowd



Magnetic Moment Beer



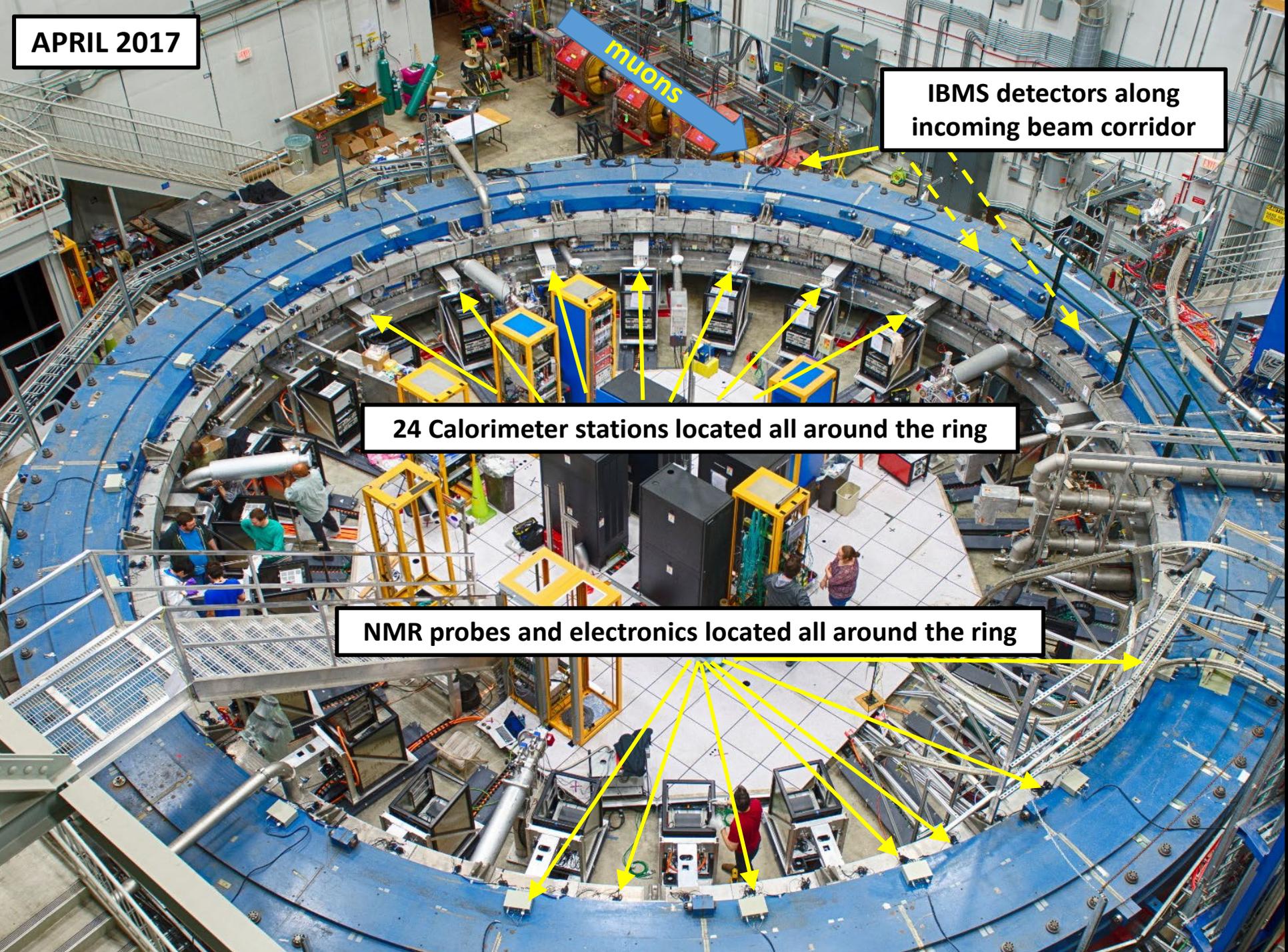
APRIL 2017

muons

IBMS detectors along incoming beam corridor

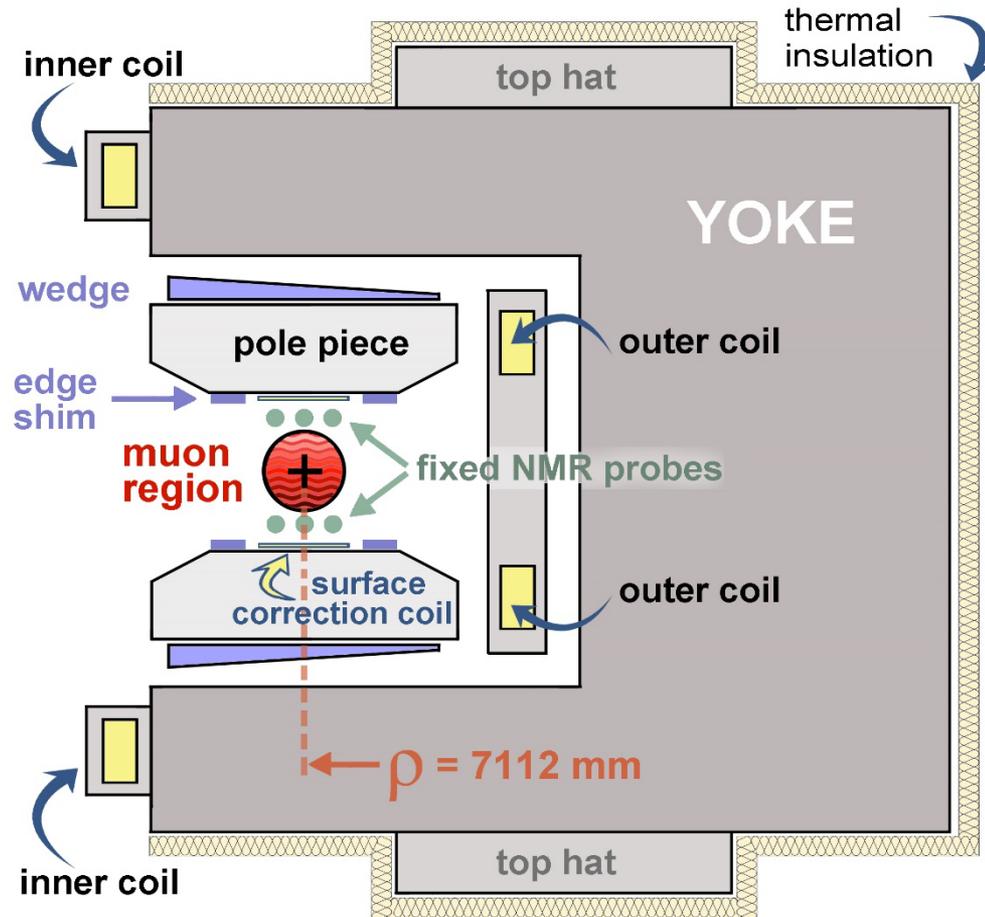
24 Calorimeter stations located all around the ring

NMR probes and electronics located all around the ring

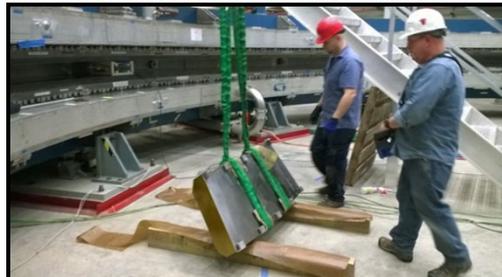
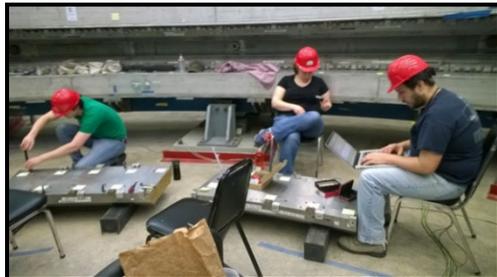


The storage ring magnet is built and shimmed

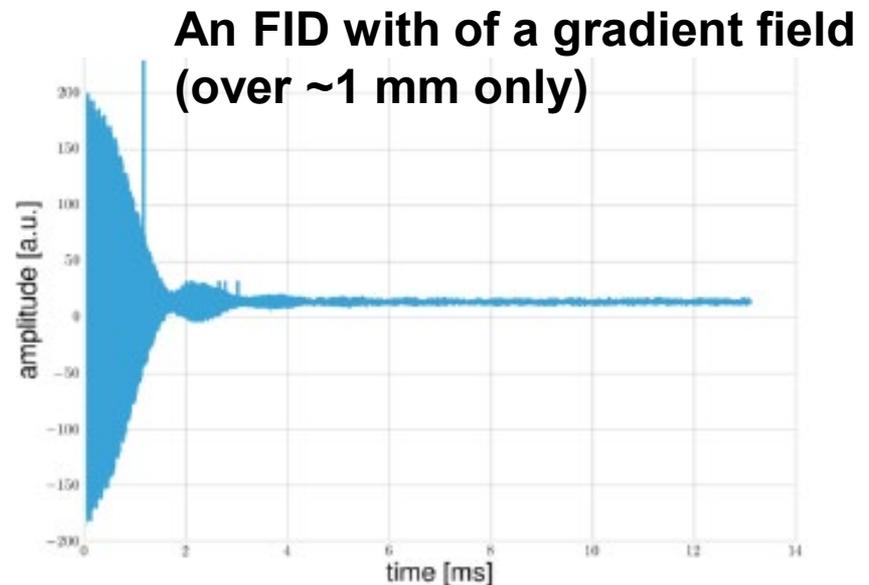
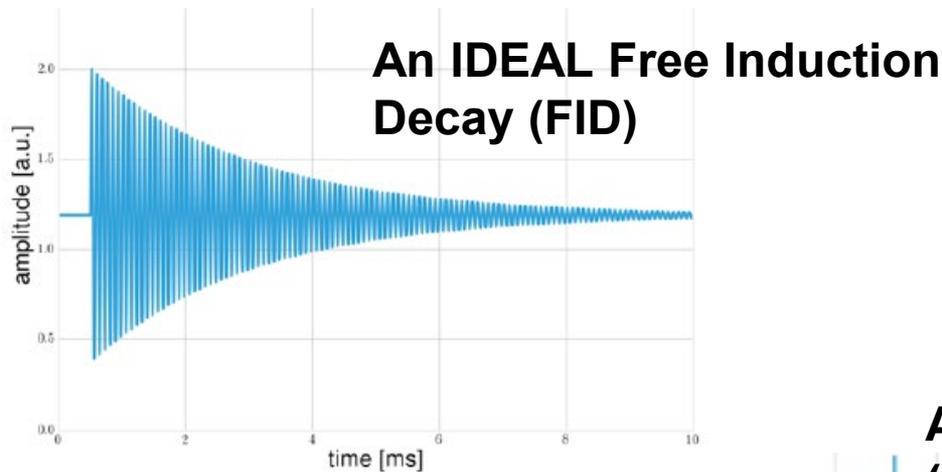
- **B Field 1.45T**
- **12 Yokes:** C shaped flux returns
- **72 Poles:** shape field
- **864 Wedges:** angle - quadrupole (QP)
- **24 Iron Top Hats:** change effective μ
- **Edge Shims:** QP, sextapole (SP)
- **8000 Surface iron foils:** change effective μ locally
- **Surface coils:** will add average field moments (360 deg)



g-2 Magnet in Cross Section



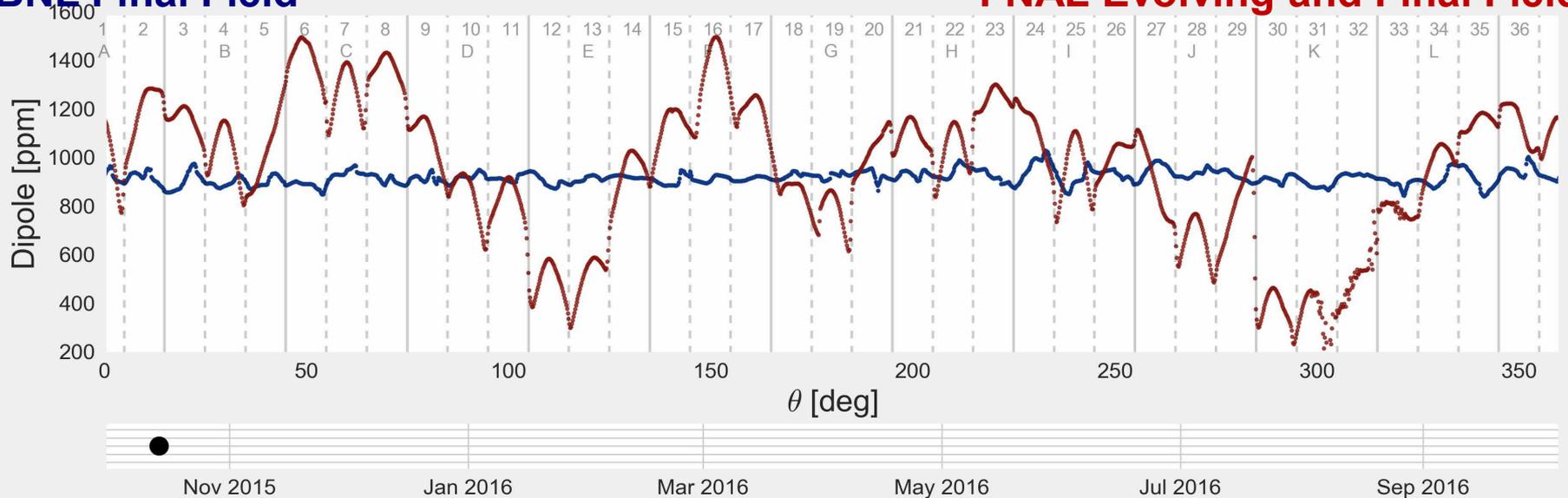
Field measured using a proxy: pulse NMR of protons



The x3 improved field uniformity compared to BNL was achieved by tuning knobs and *calculation*

BNL Final Field

FNAL Evolving and Final Field



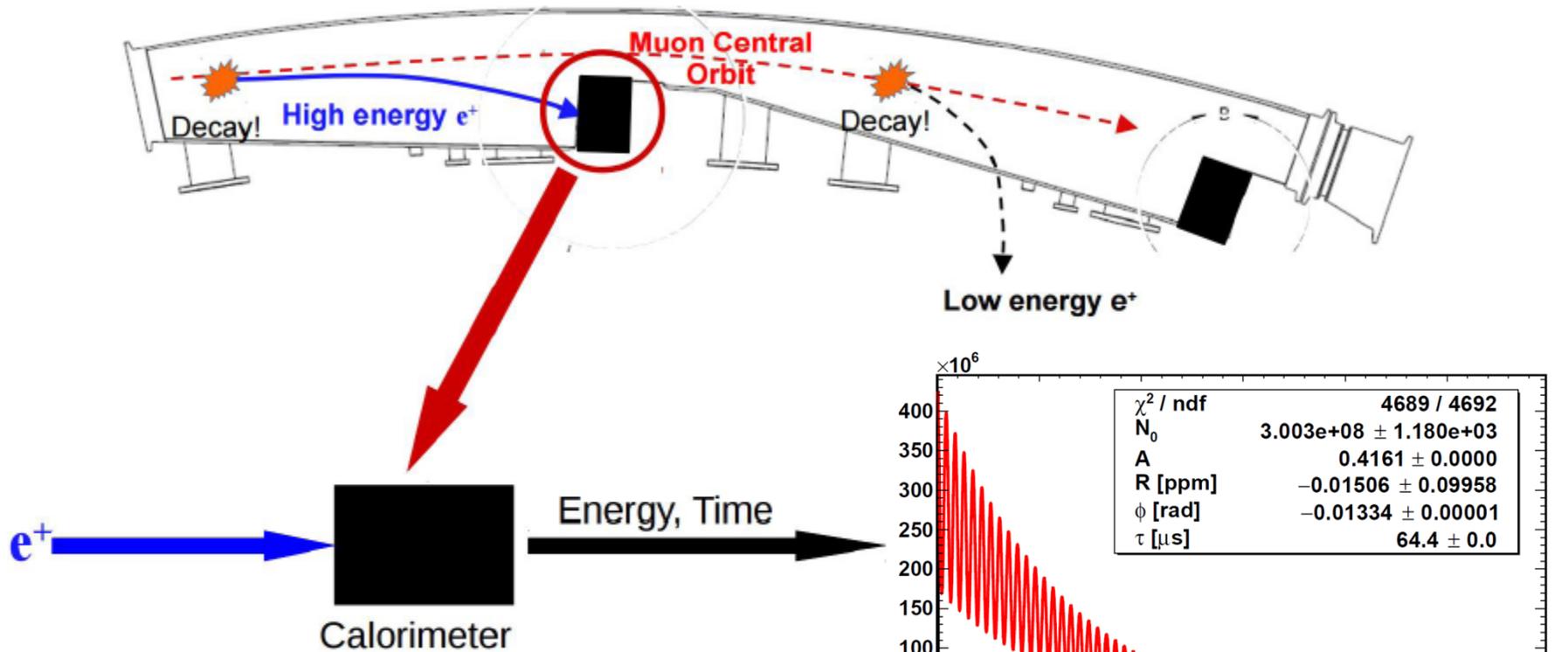
Evolution from “as built” → rough shimmed

Measures of the Average Dipole Field from 0 – 360 degrees vs MONTHs of effort

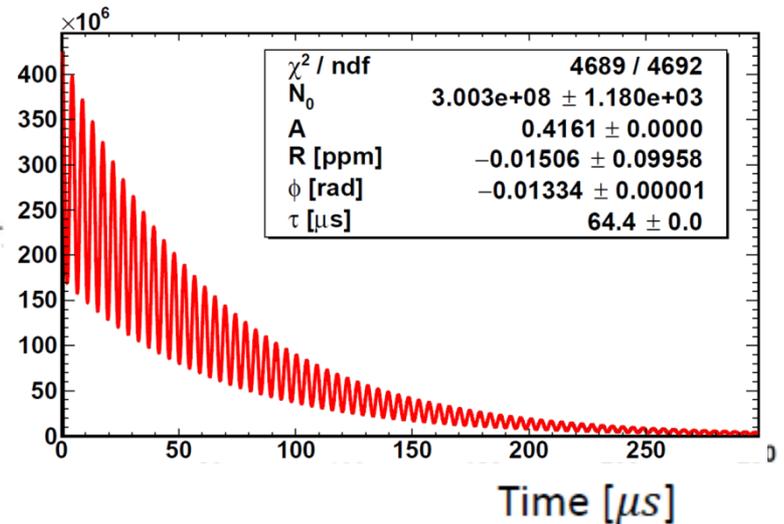
The result is 3 times better than BNL; +/- 10 ppm typically all the way around the 44 m circumference

Detector

... record muon decay times and energies, determine stored beam parameters ...



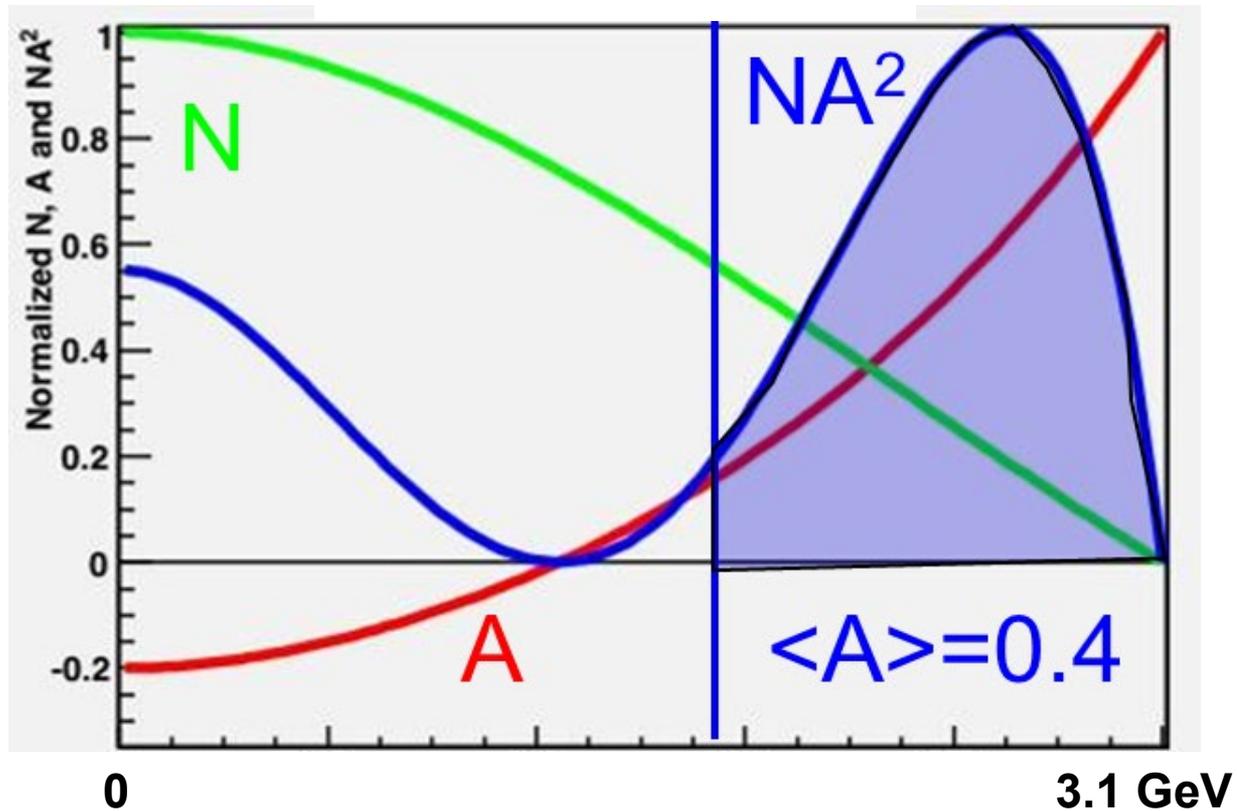
- Time and energy measurement
- Over-threshold events added to time histogram
- 700 μs muon "fills", $\sim 10,000$ stored muons/fill



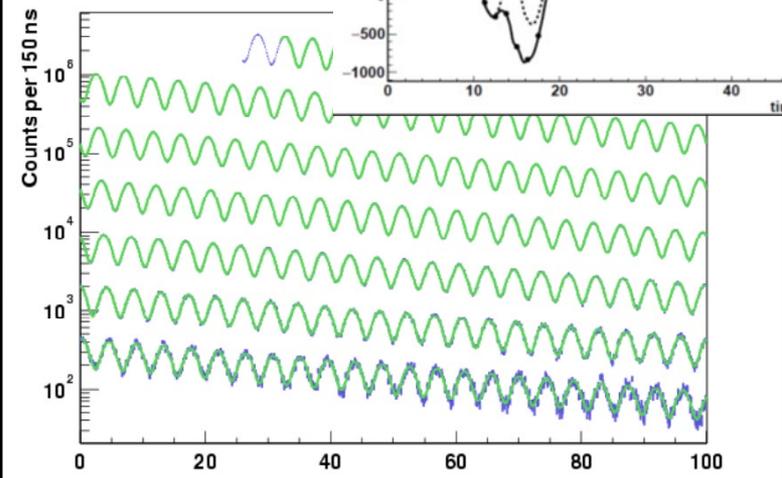
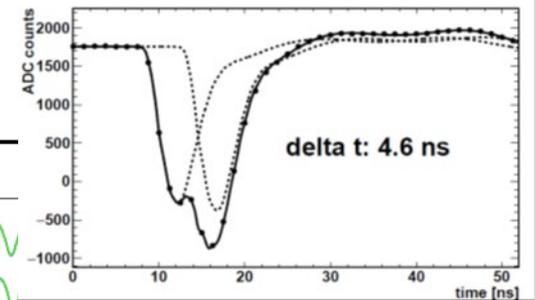
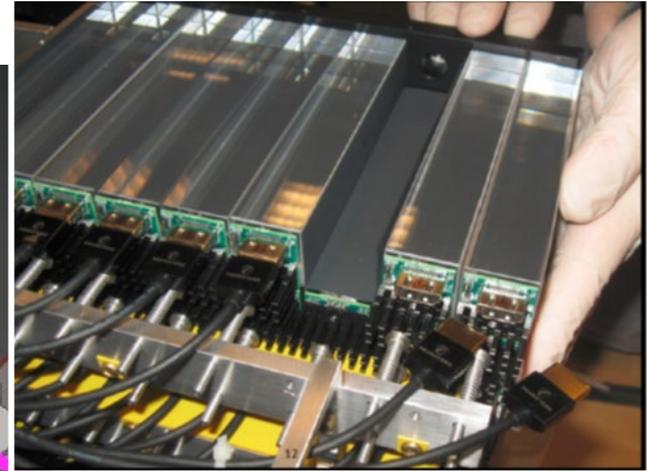
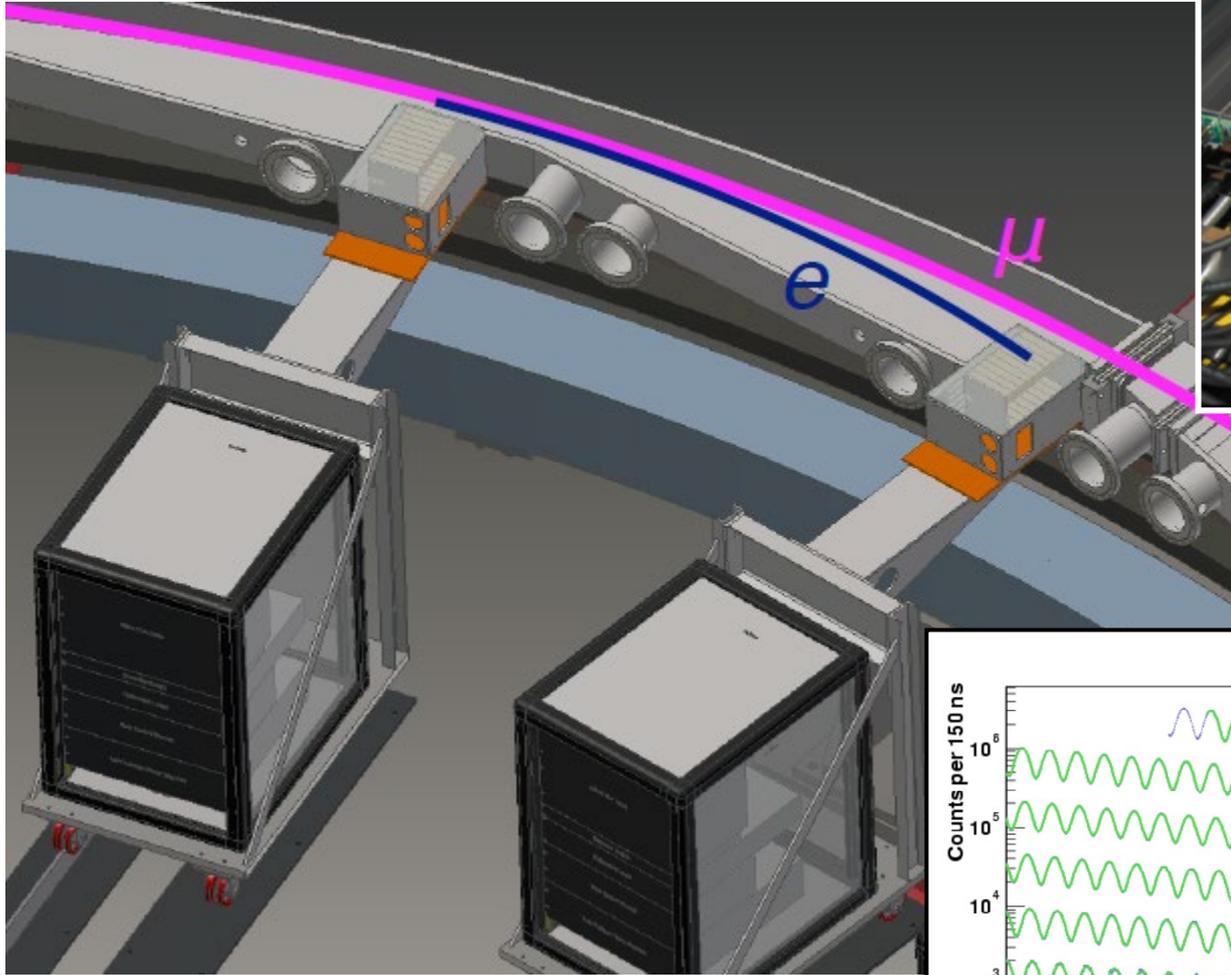
Time distribution of over-threshold decays for three thresholds

Optimizing Statistical Error

$$\delta\omega_a = \frac{1}{\gamma\tau_\mu} \sqrt{\frac{2}{NA^2}}$$



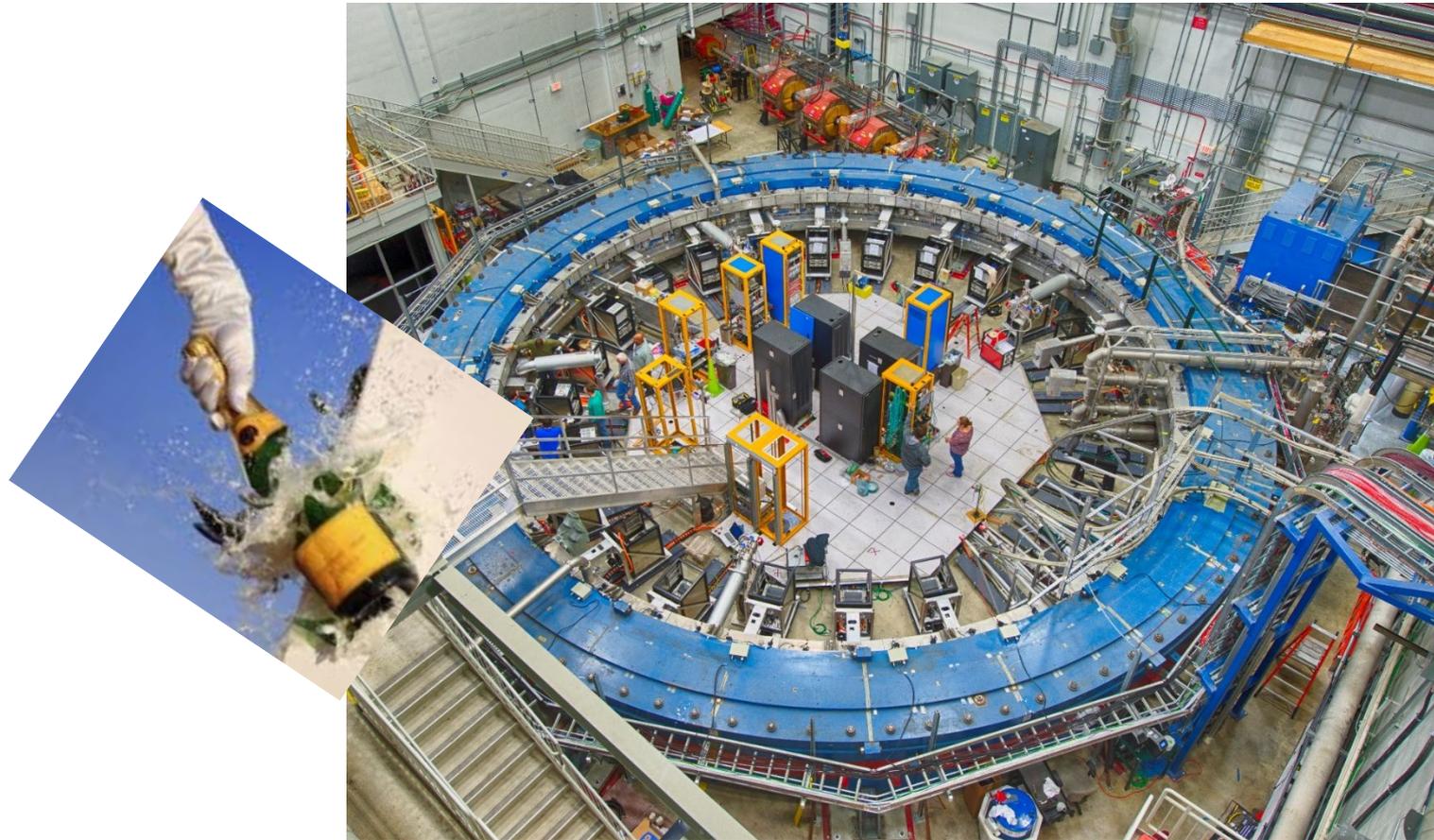
Ultra-fast PbF₂+SiPM calorimetry used to record e^+ times and energies; energy correlates with μ^+ spin



BREAKING NEWS

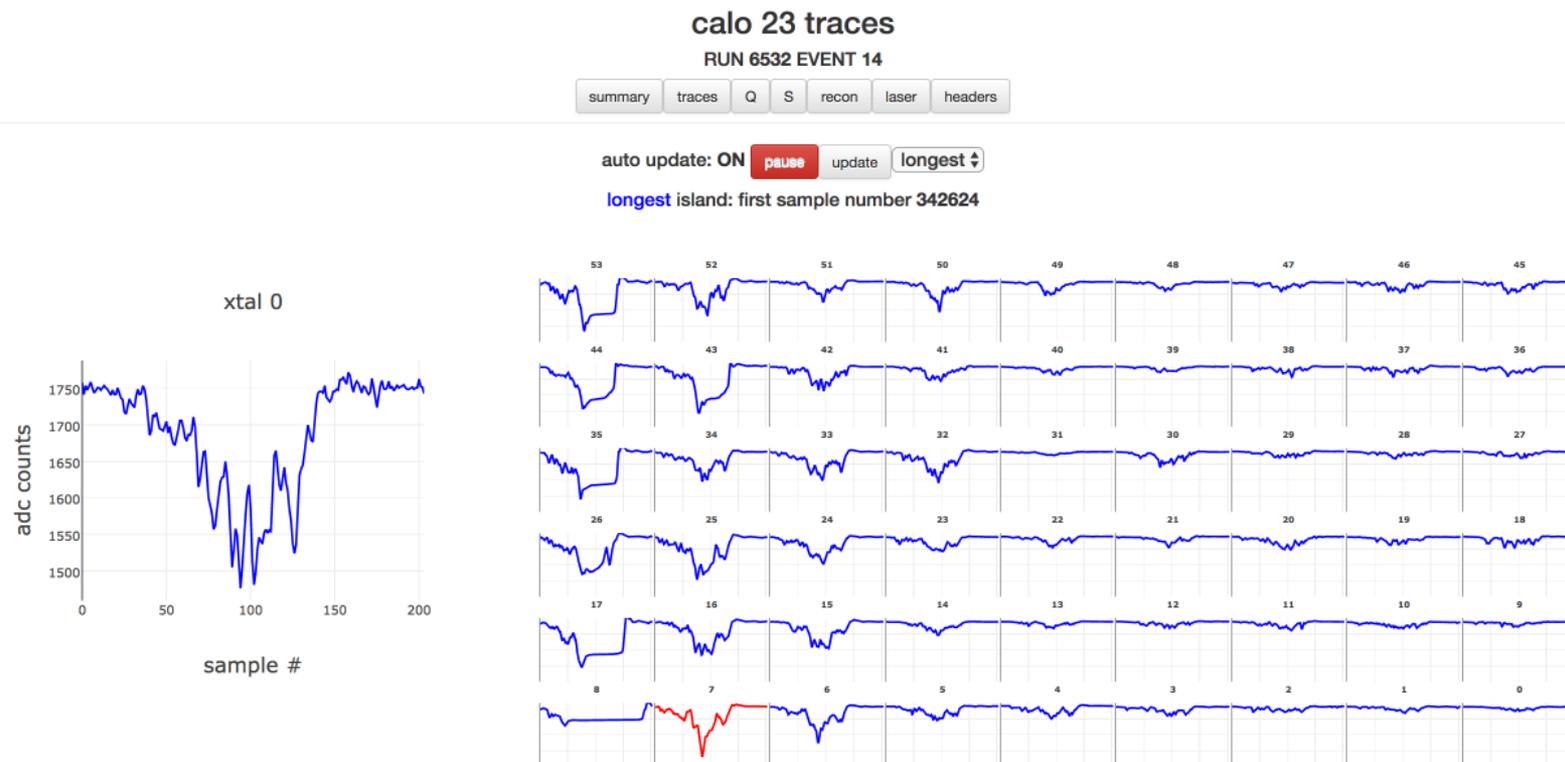
The experiment just finished a 6-week commissioning run

Let's change up a moment and let me describe the (rare) process of christening a battleship ... that is, launching a new experiment



Finding a beam is hard. When? How much?

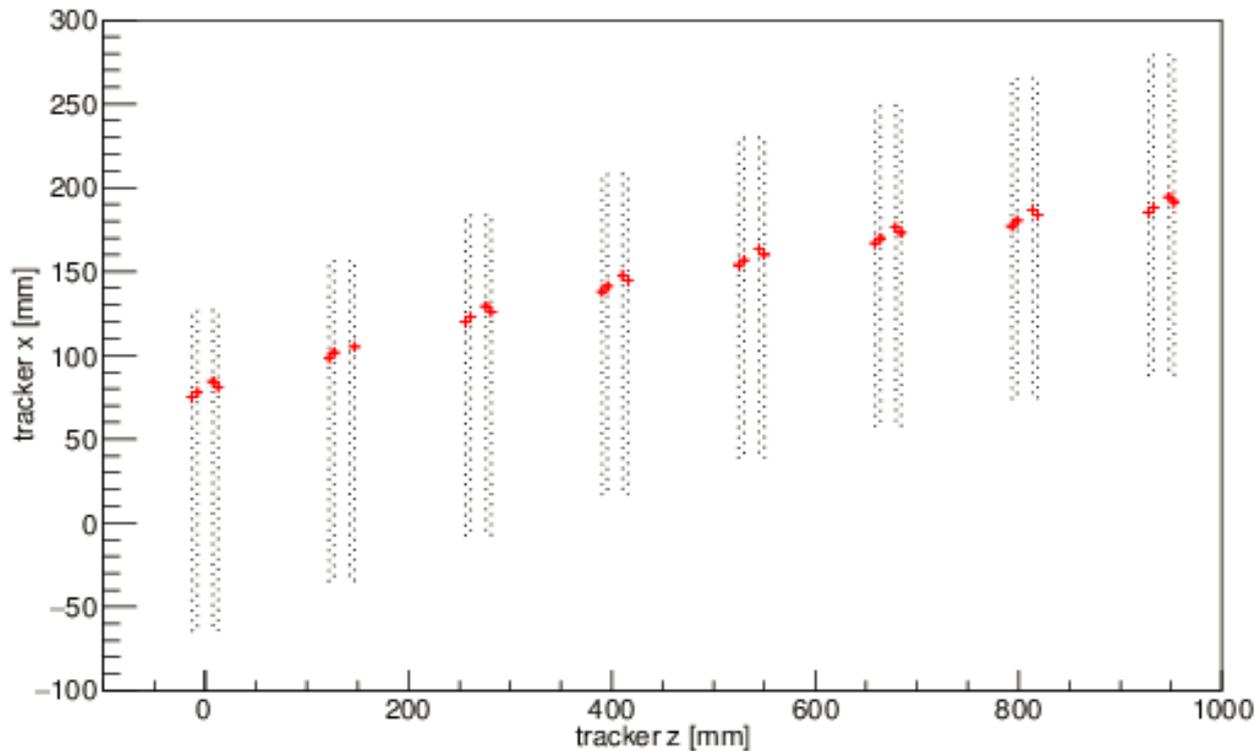
First Beam *Crashes* into Calorimeters



**Do any of the particles find stable orbits?
How about protons? Here's one that hung
around a long time**

First evidence of stored protons from some hand-selected events

Image from Tracker of escaping proton at late times

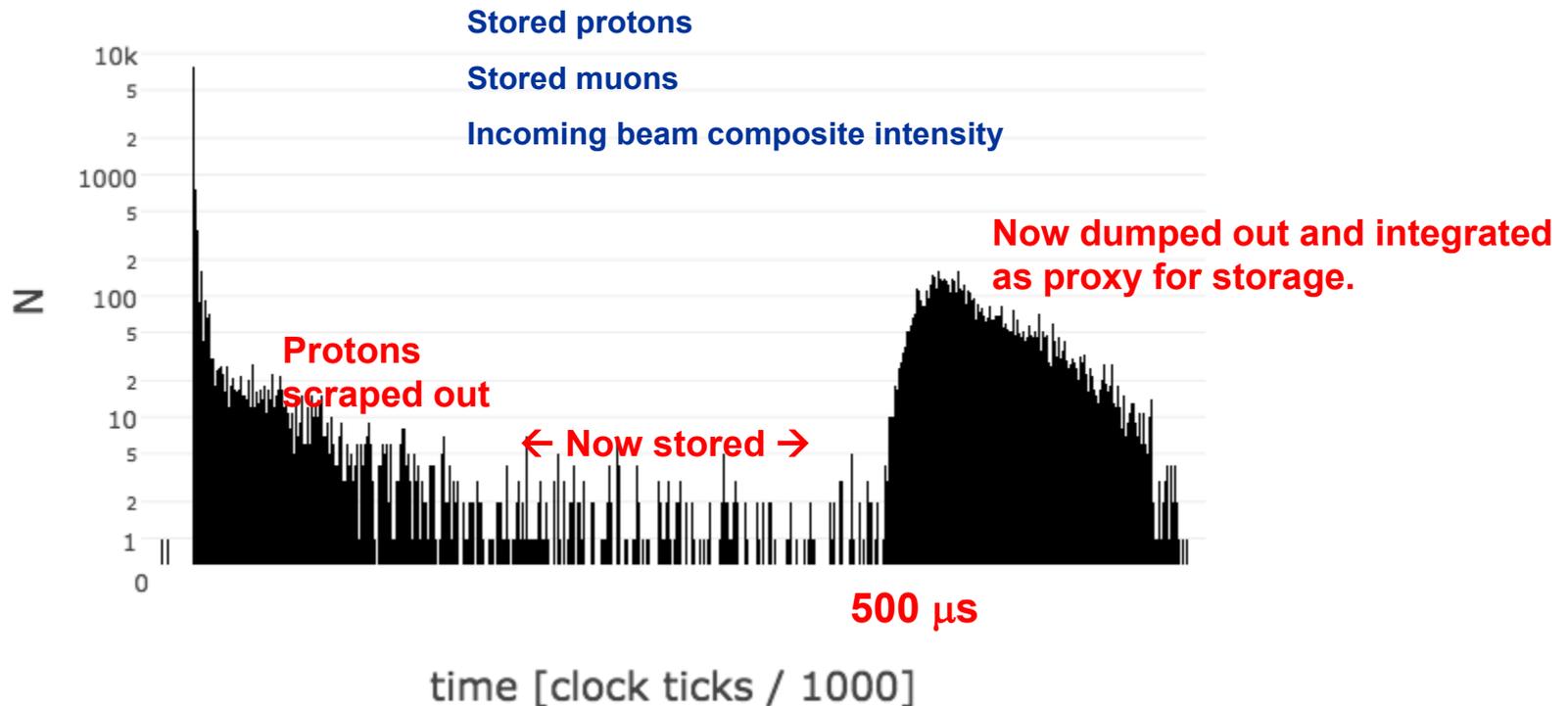


How do tune up the storage? In real time?

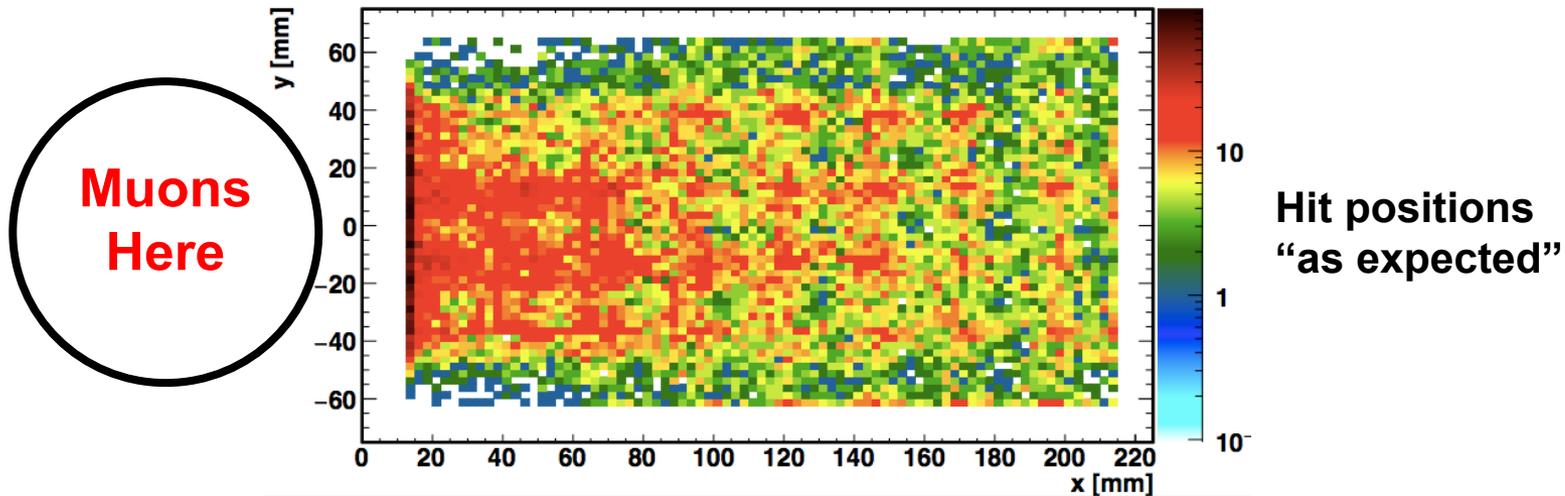
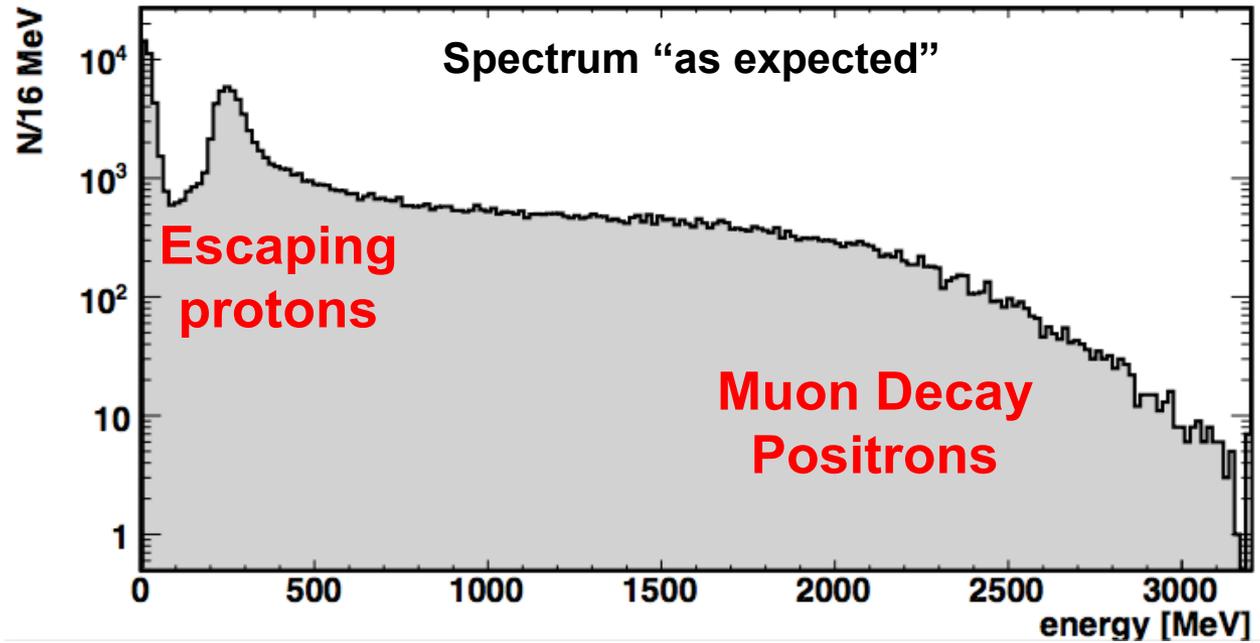
Sweep the following to optimize Storage of Protons (muons)

- Kicker timing
- Quad strength
- Inflector current
- Incoming beam x, x', y, y'
- Incoming beam focus parameters

We have online monitors of:



Do the energy spectra look like those beautiful Monte Carlo plots you made for years?

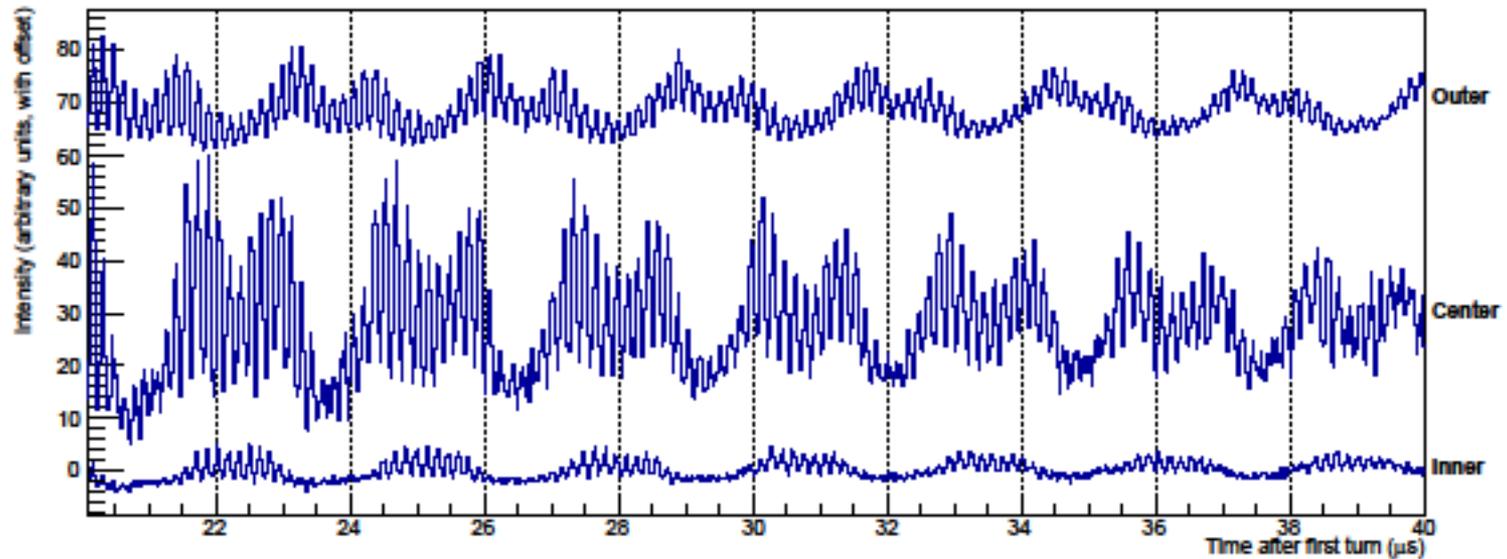


*easy to fix after run ends and we have access

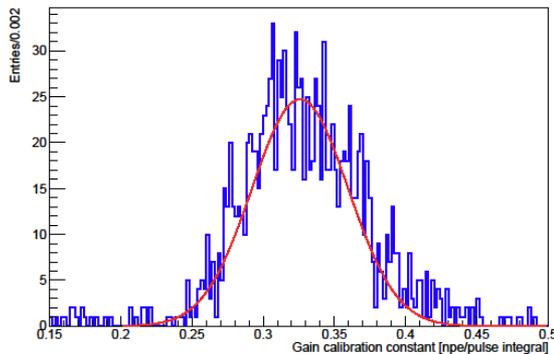
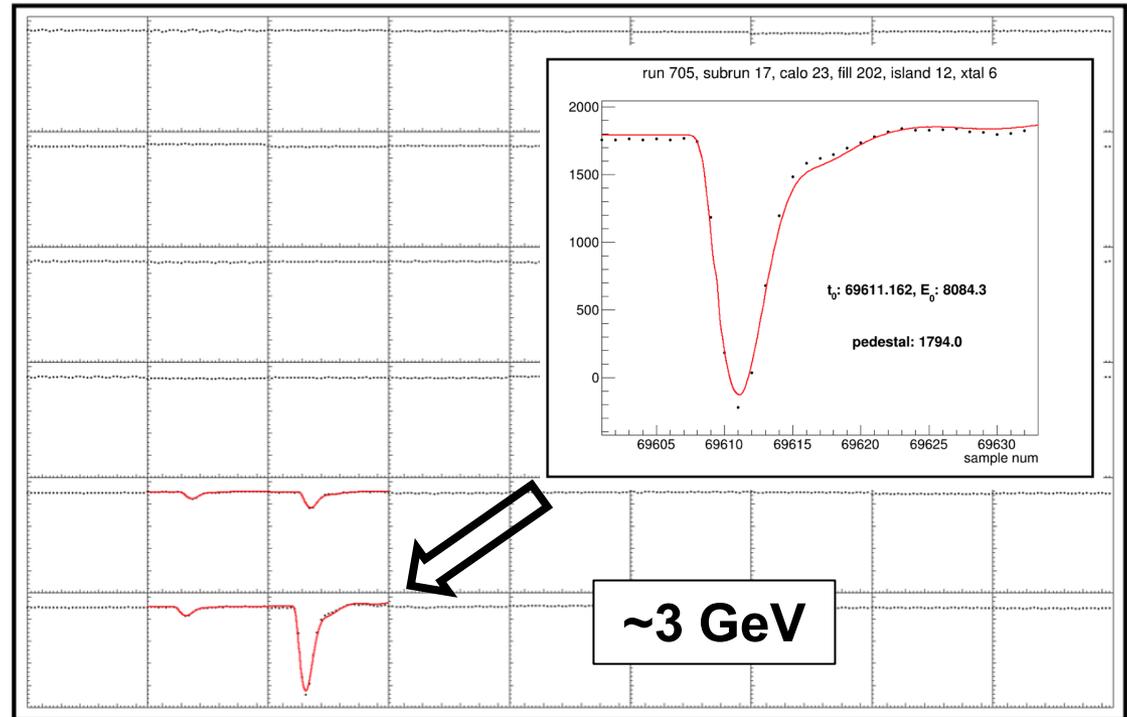
Imaging a beam is tricky; It also destroys it



Looking downstream as a Muon or Proton would



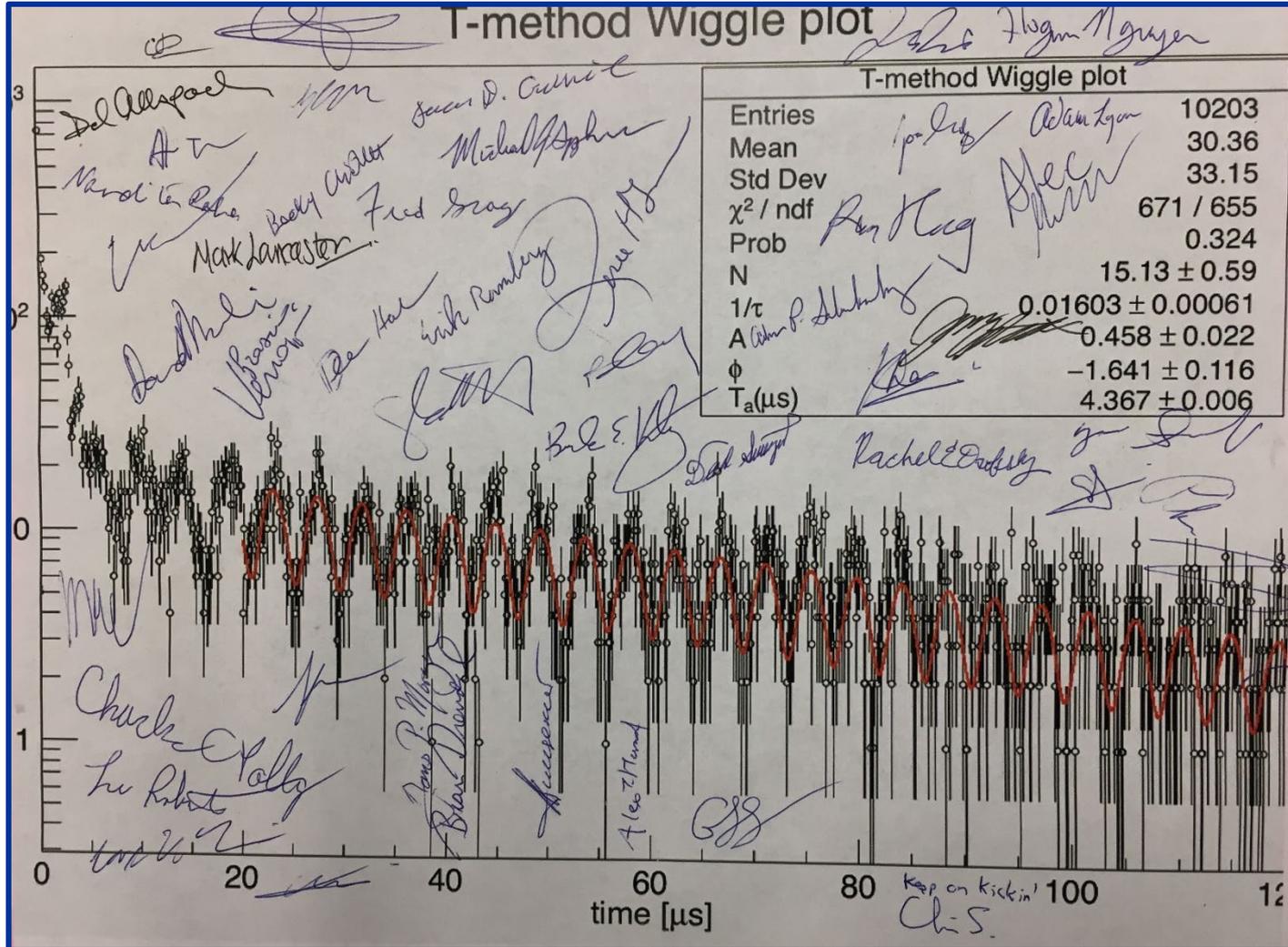
What a high-energy positron looks like in our calorimeter



Online pre-calibration gain of 1294 crystals using Laser system for absolute PE / pulse integral

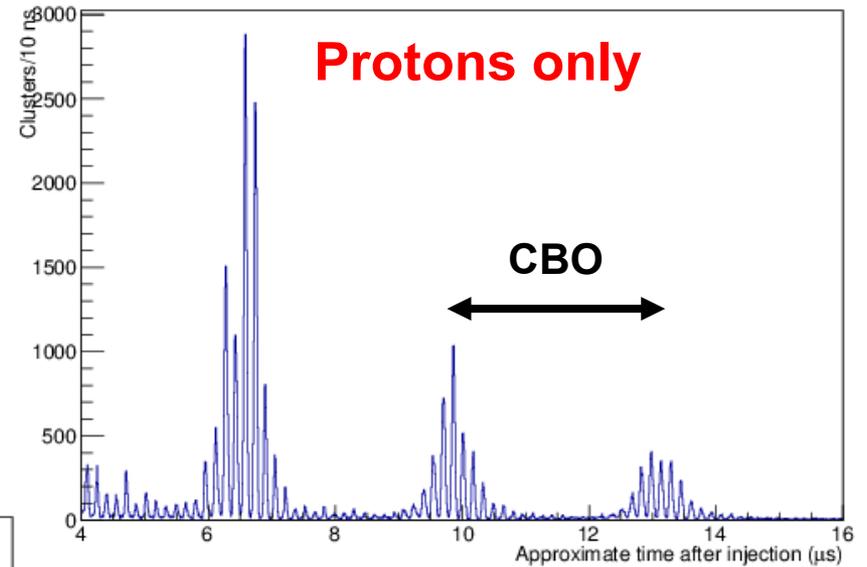
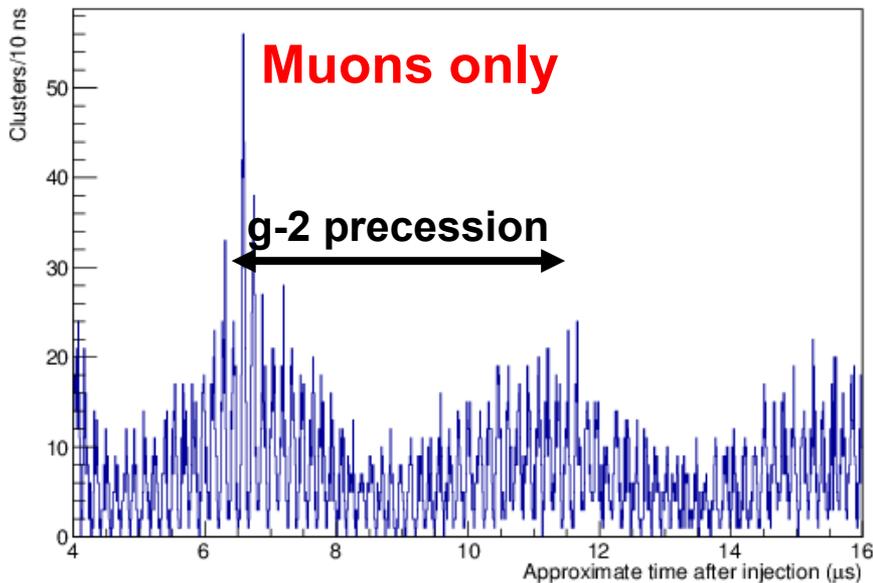
You have your moments

First evidence of stored muon precession

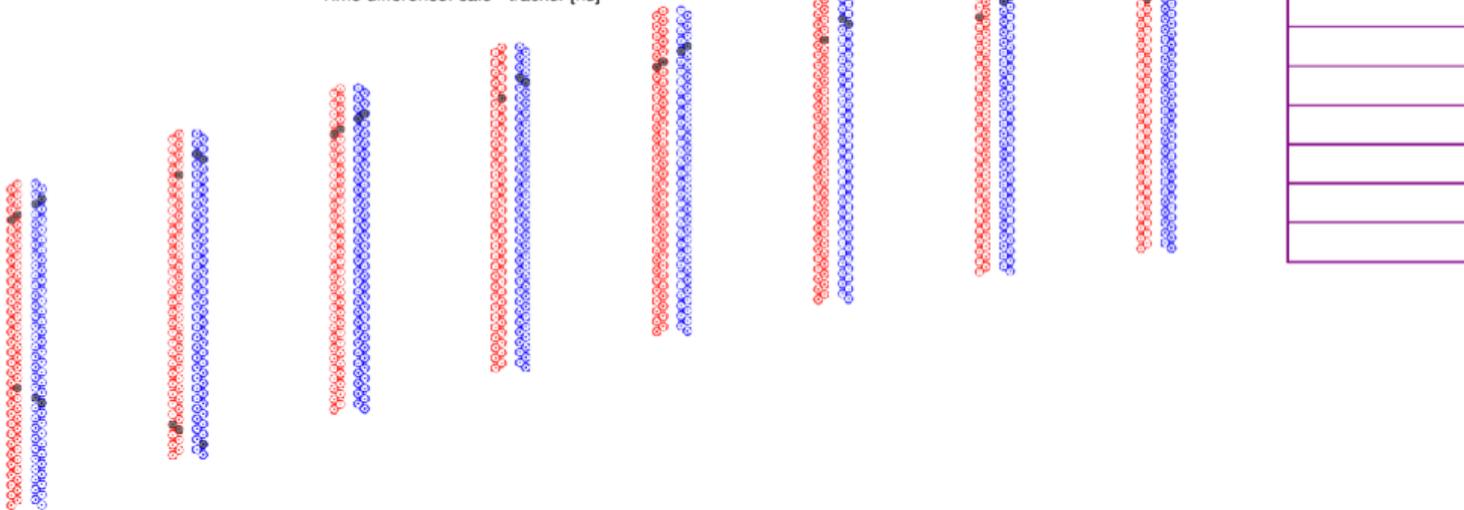
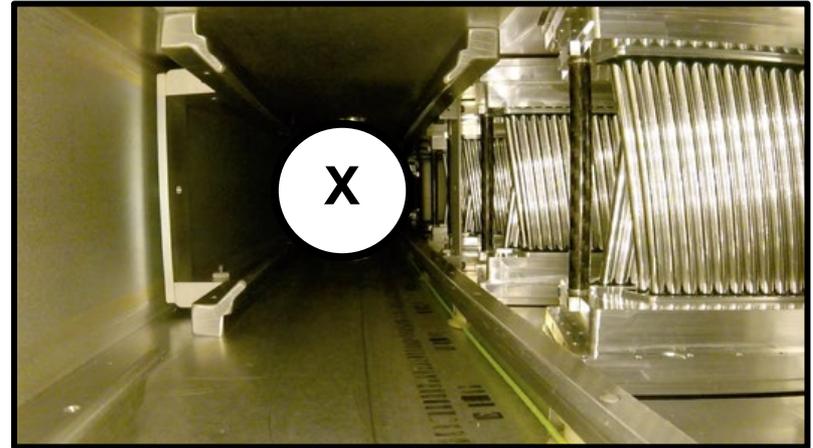
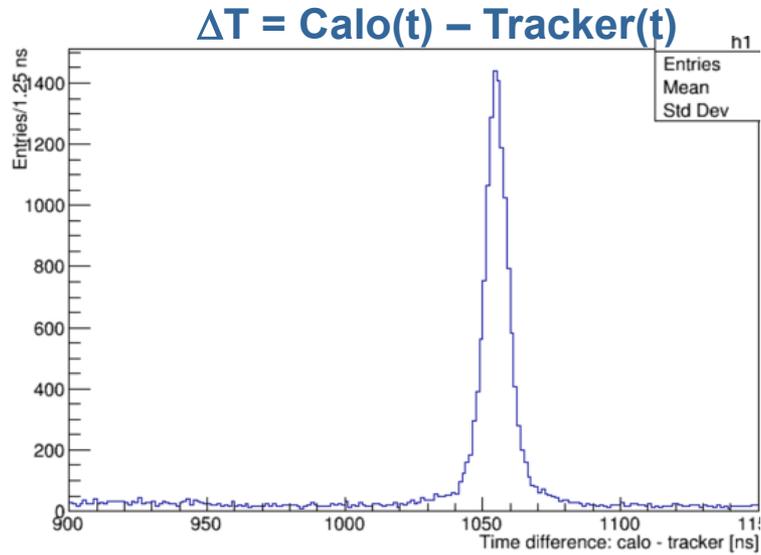


Proton and Muon Fast Rotation in calorimeters

(could not see this at BNL so easily)

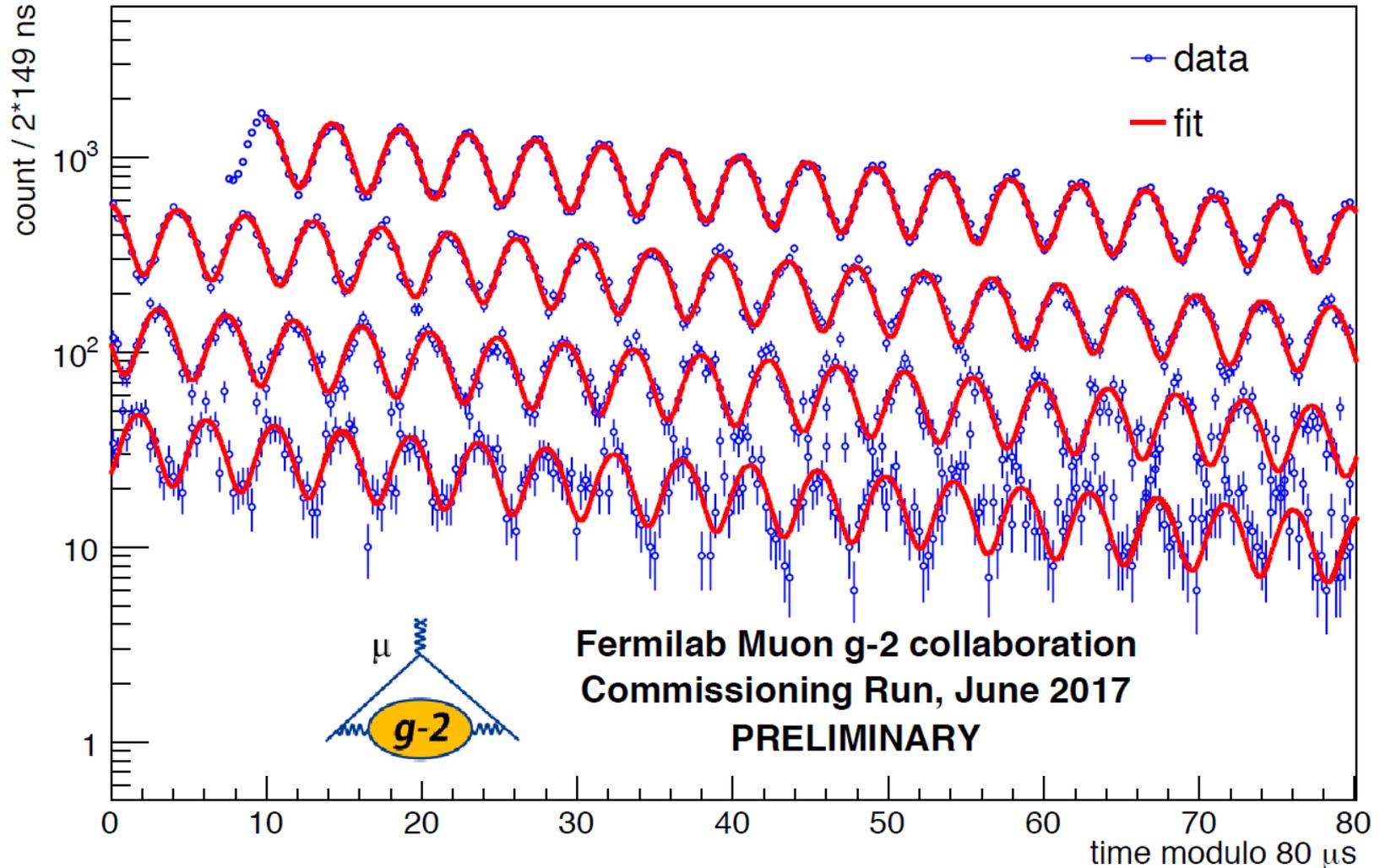


Tracker & Calorimeter working together



Getting better ... : June 25

Number of high energy positrons as a function of time



Okay, enough of that ...

- **Lessons (possibly learned)**

- The Field of Fundamental Symmetries (and later, neutrinos) has a finite number of rather specialized experiments that generally aim to do just one thing very well
- They take time
- They take ingenuity and patience
- They require a particular attention to systematics and details

- **The Physics case is rather profound**

- We aim to shake up the foundations of what is now just believed
- We KNOW there must be new physics out there ... back to Lecture 1, or else ???

- **THEORY plays a vital role in these missions**

- The known but hard: radiative corrections, hadronic effects
- The interpretations and vision: What if? And what else? And, does the idea survive the many tests as HIGH and LOW energies already?

My predictions ... (totally biased)

- **Muon $g-2$** is next most important one to watch.
 - The 3.6σ deviation is either a bad luck fluke or it's telling us something. The next experiment has started.
- **EDMs** are super promising.
 - Watch out for all systems, Hg, n, atoms, molecules, ...
- **cLFV** experiments are very sensitive to BSM now
 - New Mu2e here and MEG II in Switzerland to watch
 - Muons are much more sensitive than B factories (or future ones)
- **Neutrons** mostly “self consistency” issues
 - new generation experiments seem to be converging and the story will be looking good. There is little room for NP right now
- **I didn't mention nuclear beta decays.**
 - The He-6 system has promise but goal posts for this and other measurements $\sim x10$ beyond where experiments are now
- **PVES** has a unique reach
 - watch for imminent announcement by Qweak; keep eye on progress of MOLLER

The Final Stopping Point !

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