

Nuclear Astrophysics

Jeff Blackmon (LSU)

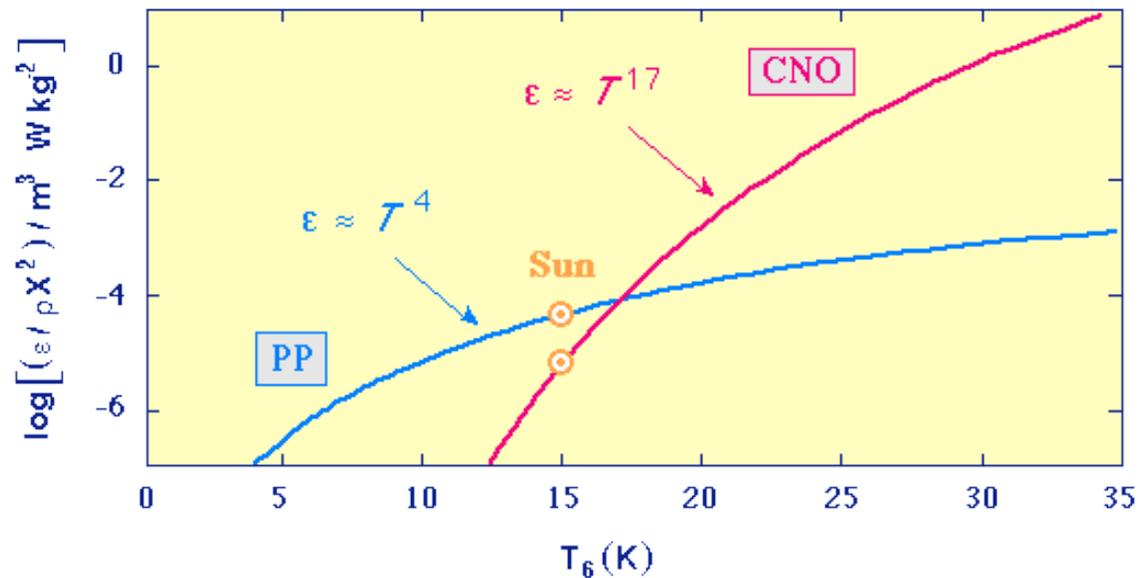
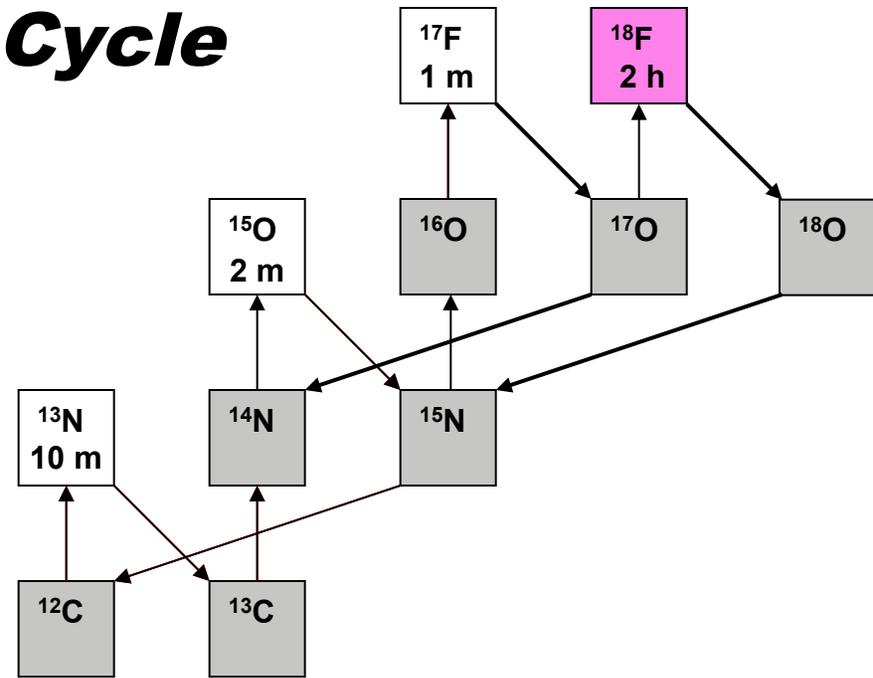
1. Introduction, Formalism, Big Bang and H burning
2. He burning, Heavy elements & s process
3. Stellar Explosions

1 H																	2 He
3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
55 Cs	56 Ba	57 La	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
87 Fr	88 Ra	89 Ac	104 Unq	105 Unp	106 Unh	107 Uns	108 Uno	109 Une	110 Unn								

58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu
90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr

CNO Cycle

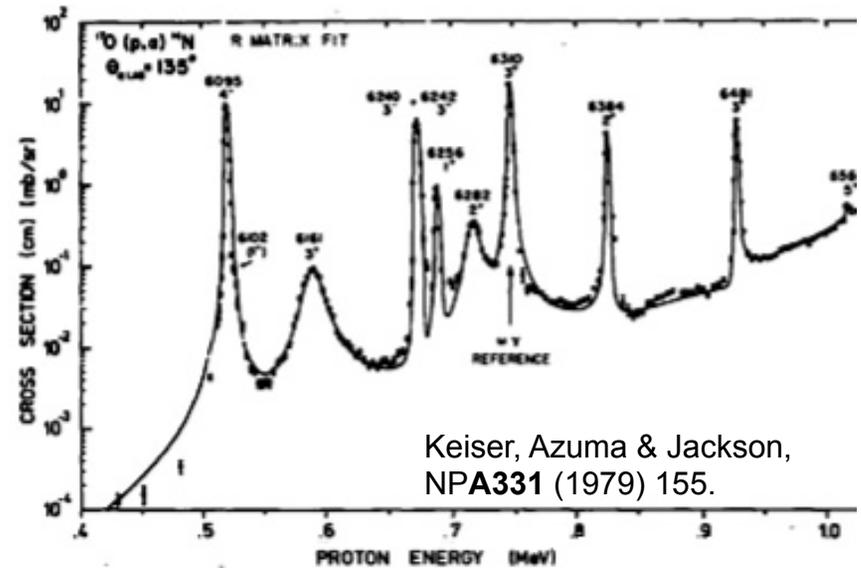
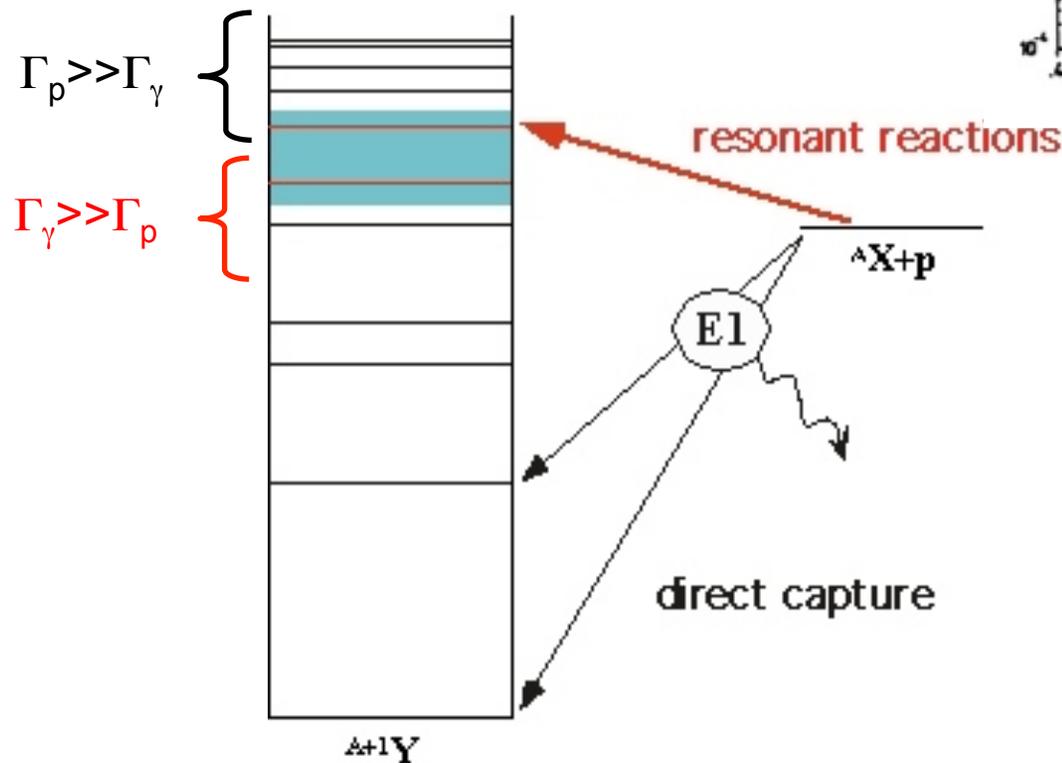
- Dominant source of energy generation in stars heavier than the sun
- What is CNO contribution to energy production in the sun? Few% ?
- CNO abundances in sun uncertain
- Stellar photospheric metallicity disagrees with helioseismology



Resonances are important

$$\langle \sigma v \rangle = \sqrt{\frac{8}{\pi \mu}} (kT)^{3/2} \int_0^{\infty} \sigma E e^{-E/(kT)} dE$$

$$\sigma(E) = \pi \hat{\lambda}^2 \frac{2J+1}{(2J_x+1)(2J_y+1)} \frac{\Gamma_x \Gamma_y}{(E - E_r)^2 + (\Gamma/2)^2}$$



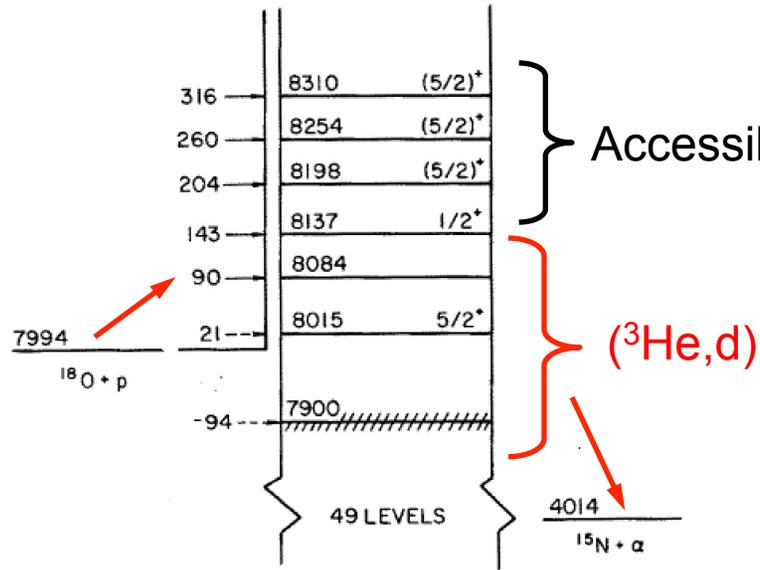
If resonance is narrow

$$\langle \sigma v \rangle = \left(\frac{2\pi}{\mu} kT \right)^{3/2} \hbar^2 (\omega \gamma) e^{-E_r/kT}$$

$$\omega \gamma = \frac{2J+1}{(2J_x+1)(2J_y+1)} \frac{\Gamma_x \Gamma_y}{\Gamma}$$

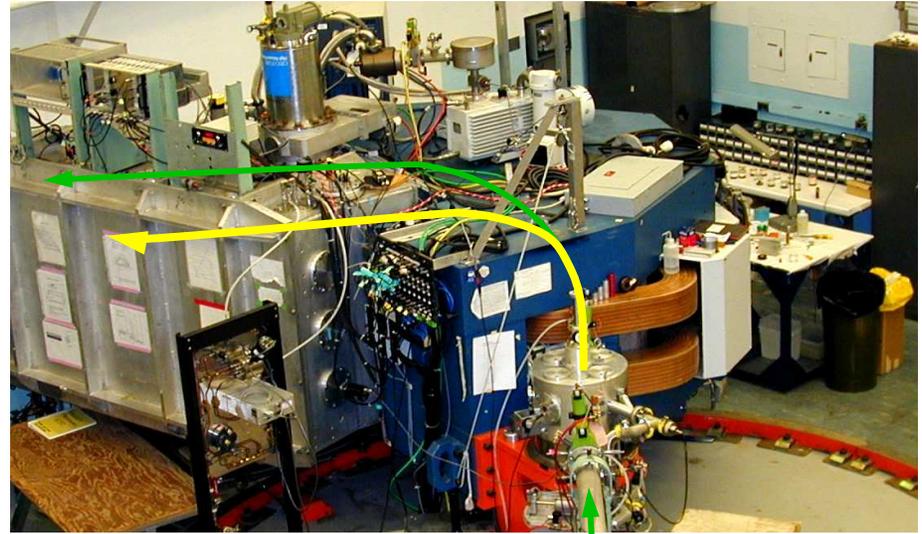
“resonance strength”

Example: $^{18}\text{O}(p,\alpha)^{15}\text{N}$

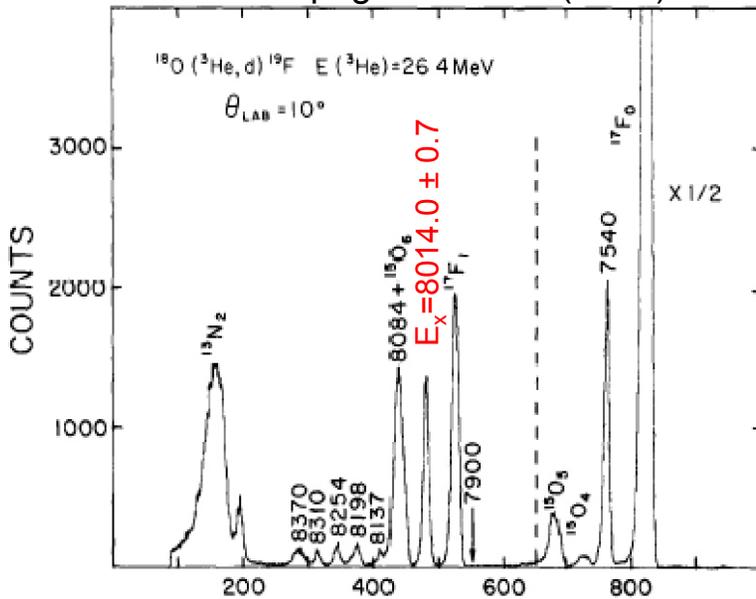


Accessible with high intensity proton beams

Magnetic Spectrograph



Champagne and Pitt (1986)

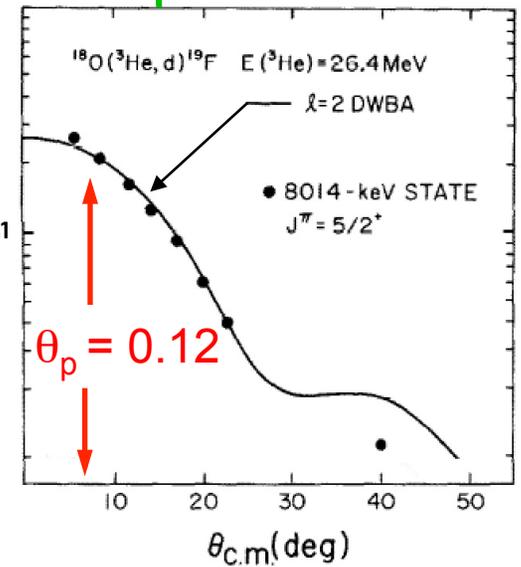


➤ Accurate E_x

➤ ℓ, J^π inferred

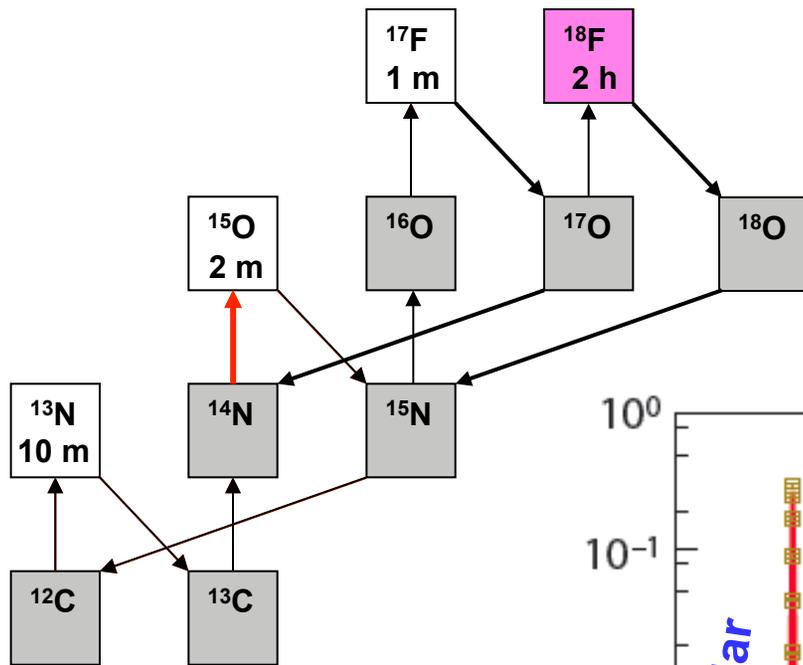
$$\Gamma_p = 2 \left(\frac{\hbar^2}{\lambda \mu R} \right) \left(\frac{\theta_p^2}{F_\ell^2 + G_\ell^2} \right)$$

with 1 mA p + ^{18}O
1 event / 3×10^5 years

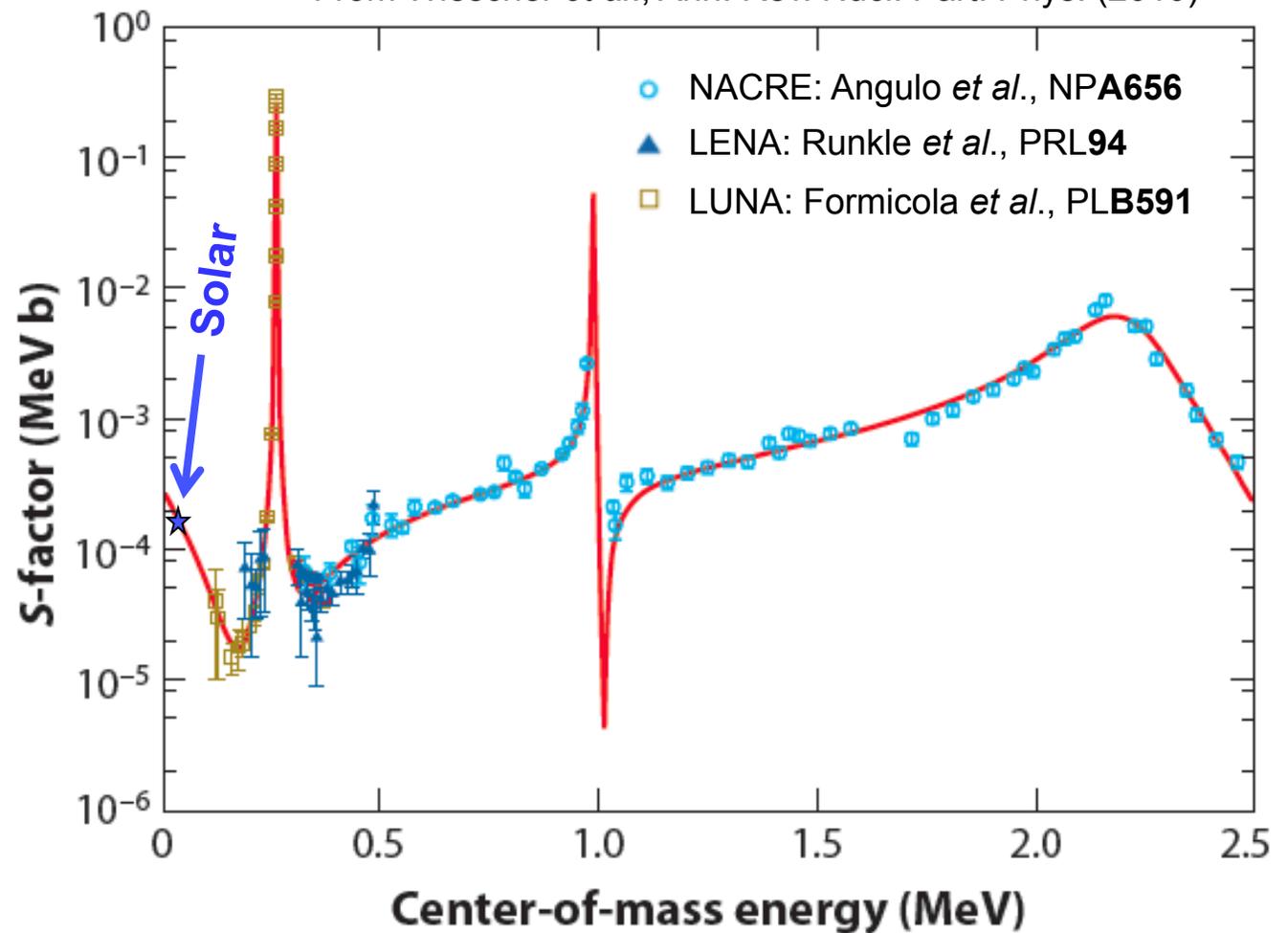


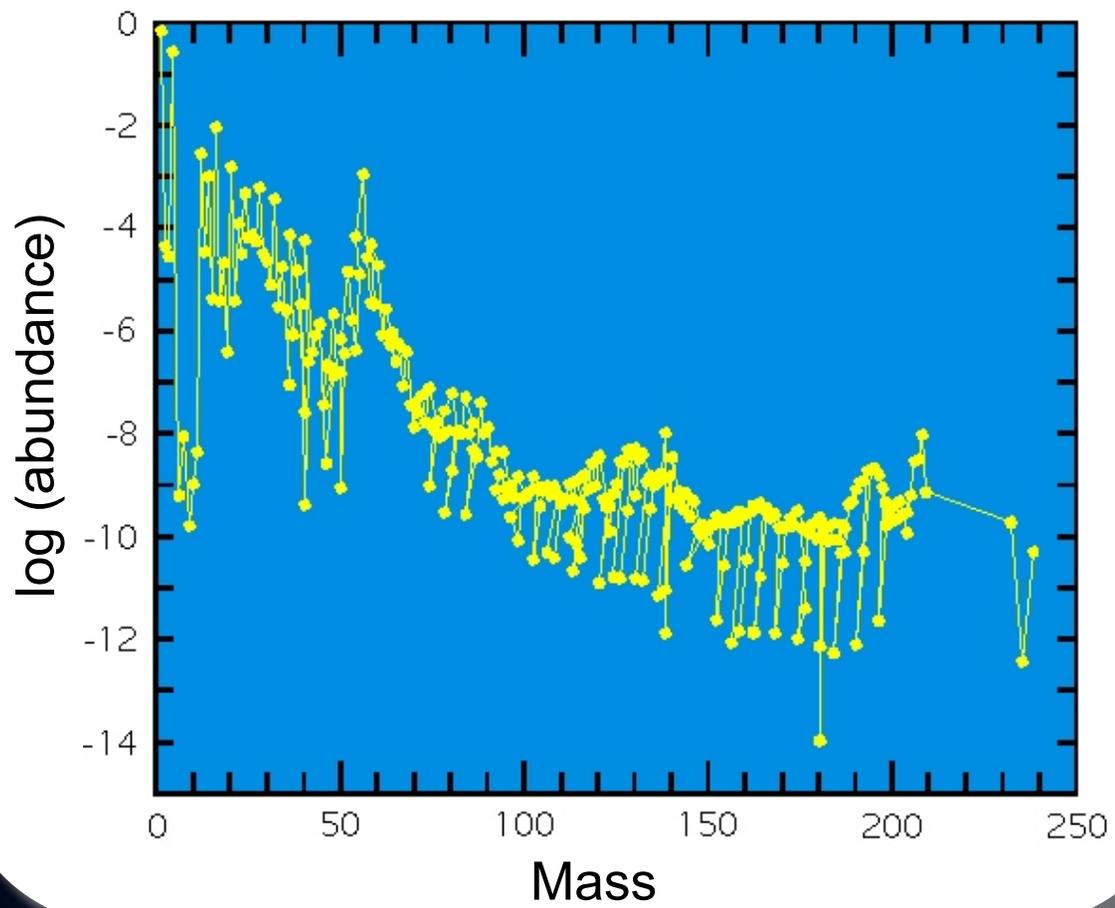
$^{14}\text{N}(p,\gamma)^{15}\text{O}$

- Slowest reaction in CN cycle
- Determines rate of energy generation and relative abundances



From Wiescher *et al.*, Ann. Rev. Nucl. Part. Phys. (2010)





Can the sun synthesize heavier elements?



No atoms exist in nature with an $A = 5$ or 8

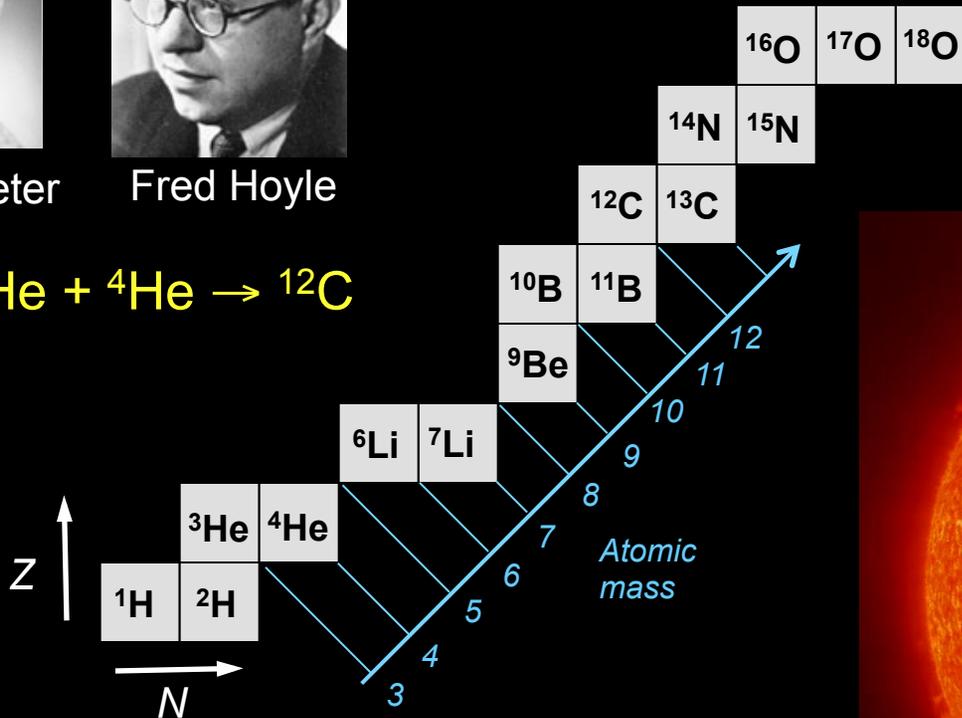
${}^8\text{Be}$ lifetime $\sim 10^{-16}$ s



Edwin Salpeter

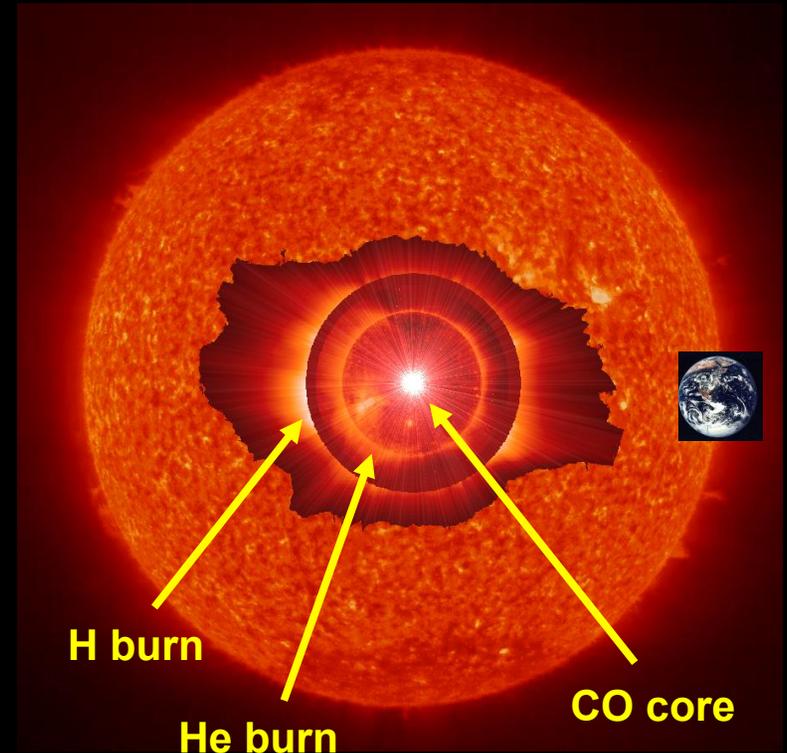


Fred Hoyle



Only possible if ${}^{12}\text{C}$ has a very large resonance at perfect energy

Red Giant Star



He burning & the “Hoyle” state

$$t_{1/2}({}^8\text{Be}) = 9.7 \times 10^{-17} \text{ s}$$

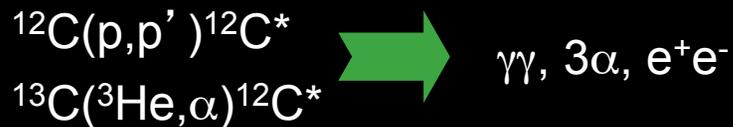


$$\frac{N({}^8\text{B})}{N(\alpha)} \approx 5 \times 10^{-10}$$

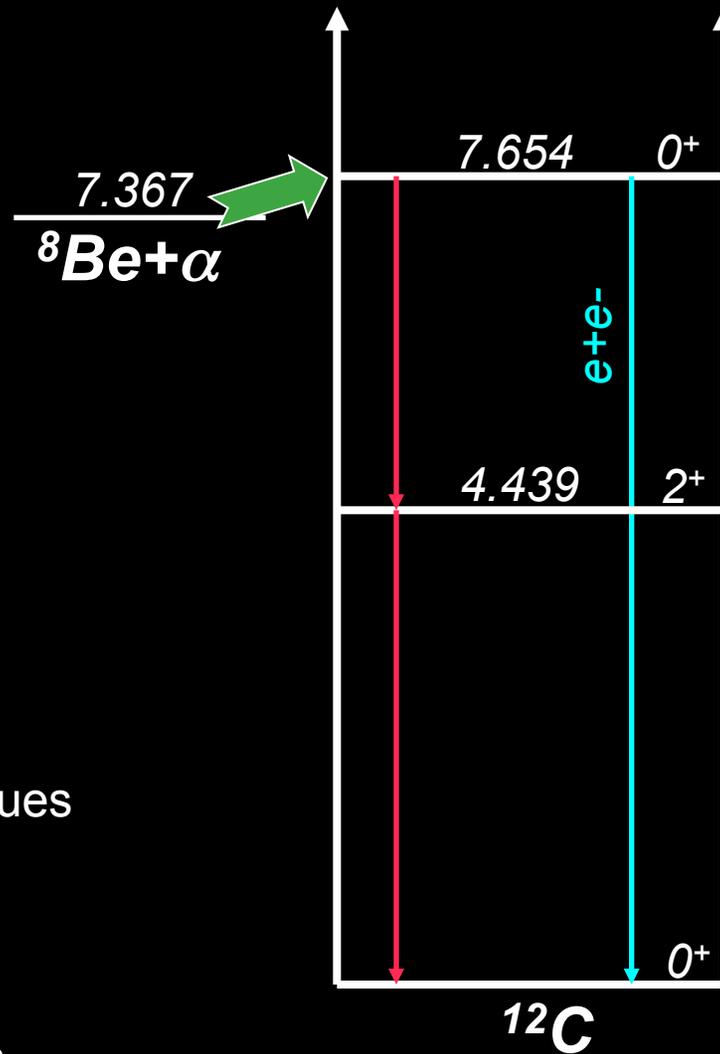
0⁺ resonance near the Gamow energy was predicted by Hoyle

Phys Rev 92 (1953) 1095.

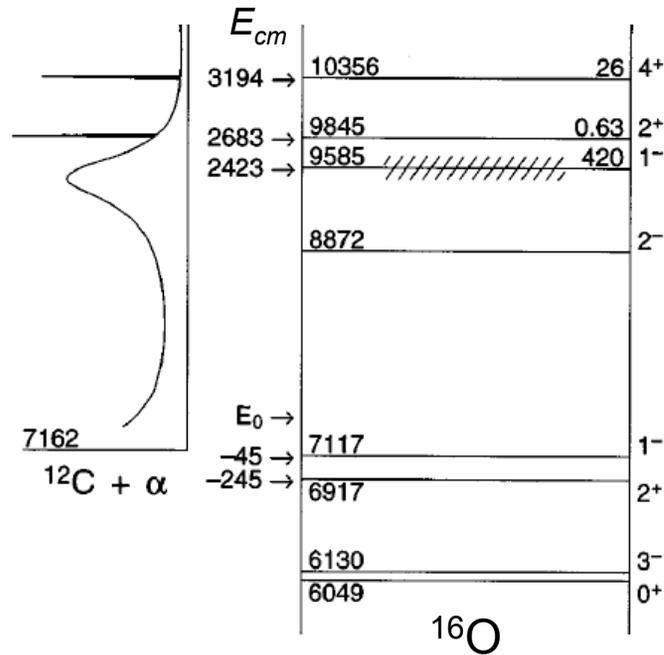
Numerous complementary techniques



Largest uncertainty $\Gamma_{ee} \sim 12\%$

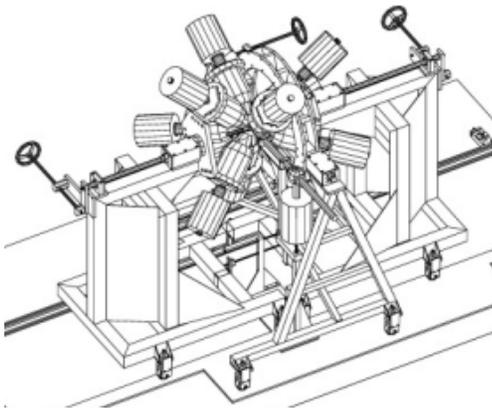


$^{12}\text{C}(\alpha,\gamma)^{16}\text{O}$

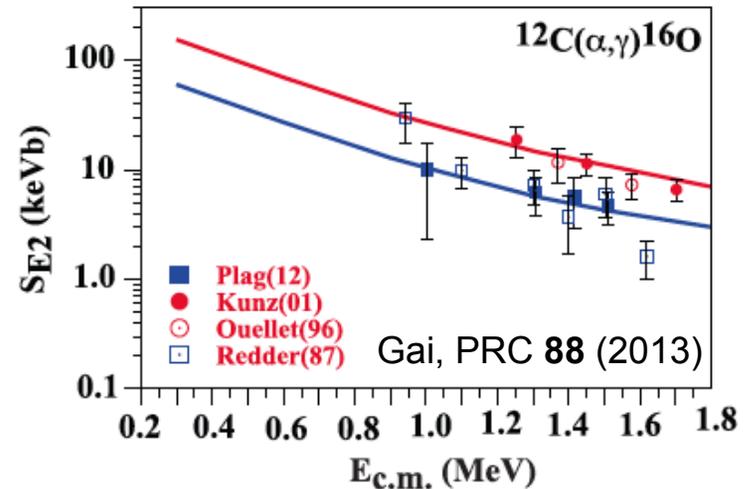
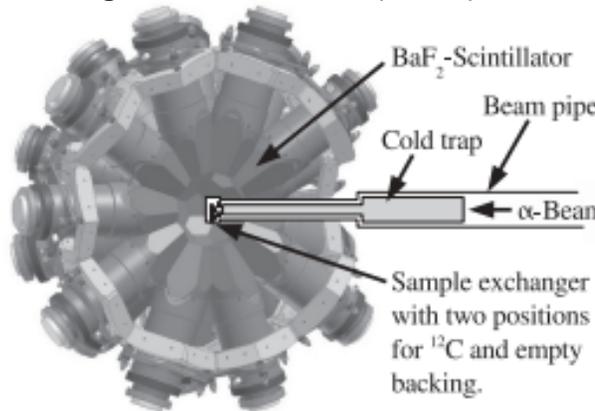


- The $^{12}\text{C}(\alpha,\gamma)^{16}\text{O}$ reaction rate fixes ratio of $^{12}\text{C}/^{16}\text{O}$ following helium burning
 - ➔ Abundance ratio of $^{12}\text{C}/^{16}\text{O}$
- The $^{12}\text{C}/^{16}\text{O}$ ratio governs subsequent evolution of the star:
 - ➔ Size of Fe core & supernova dynamics
- Subthreshold states influence → uncertain interference with other states
- Contribution of both E1 and E2 contributions are important → need to measure angular distributions

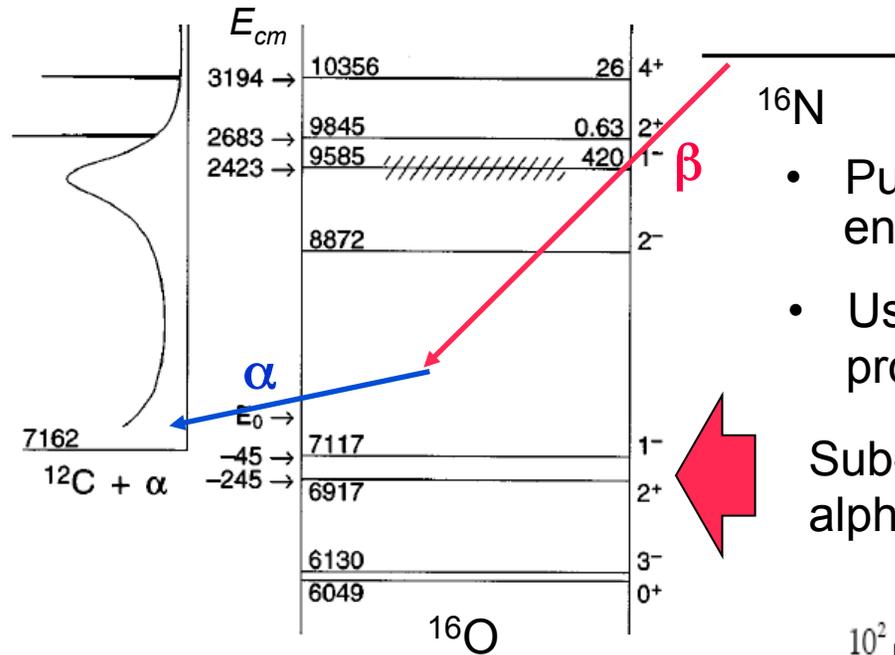
Stuttgart: Eurogam (HPGe)
Assunção *et al.*, PRC 73 (2006)



Karlsruhe: BaF_2
Plag *et al.*, PRC 86 (2012)



$^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$

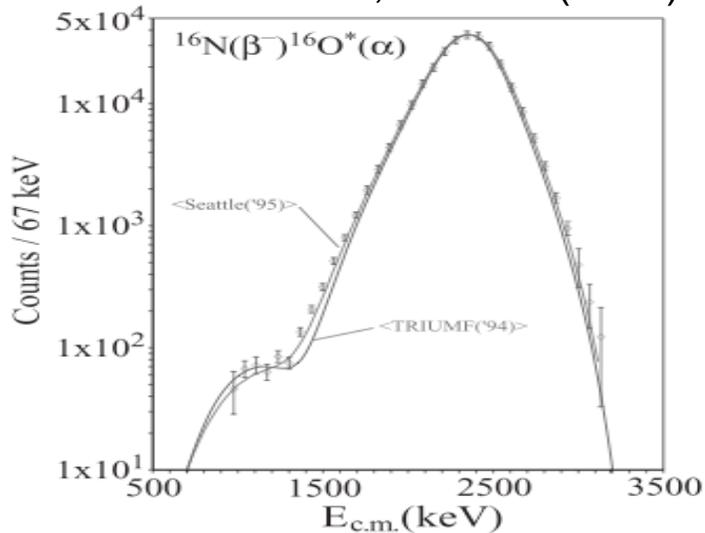


Inverse reaction studied from ^{16}N beta-delayed alpha emission

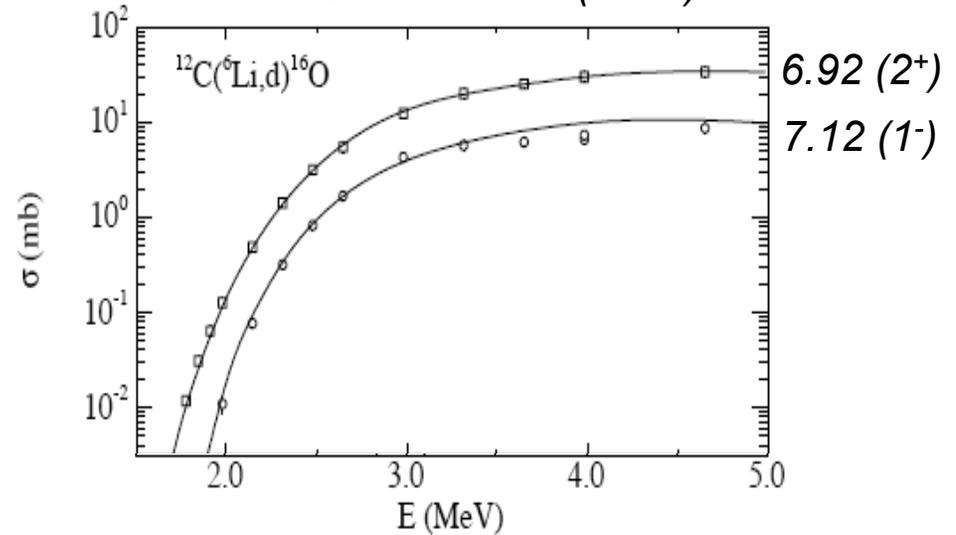
- Pushing direct measurements to lower energies is crucial but challenging
- Use indirect techniques to study properties of subthreshold states

Sub-Coulomb alpha transfer determines alpha-like part of asymptotic wavefunction

France *et al.*, PRC 75 (2007)



Brune *et al.* PRL 83 (1999)



$$S_{E_2}(300\text{keV}) = 40 \pm 20\text{keV} \cdot b$$

AGB Stars – Fate for $M < 8 M_{\odot}$

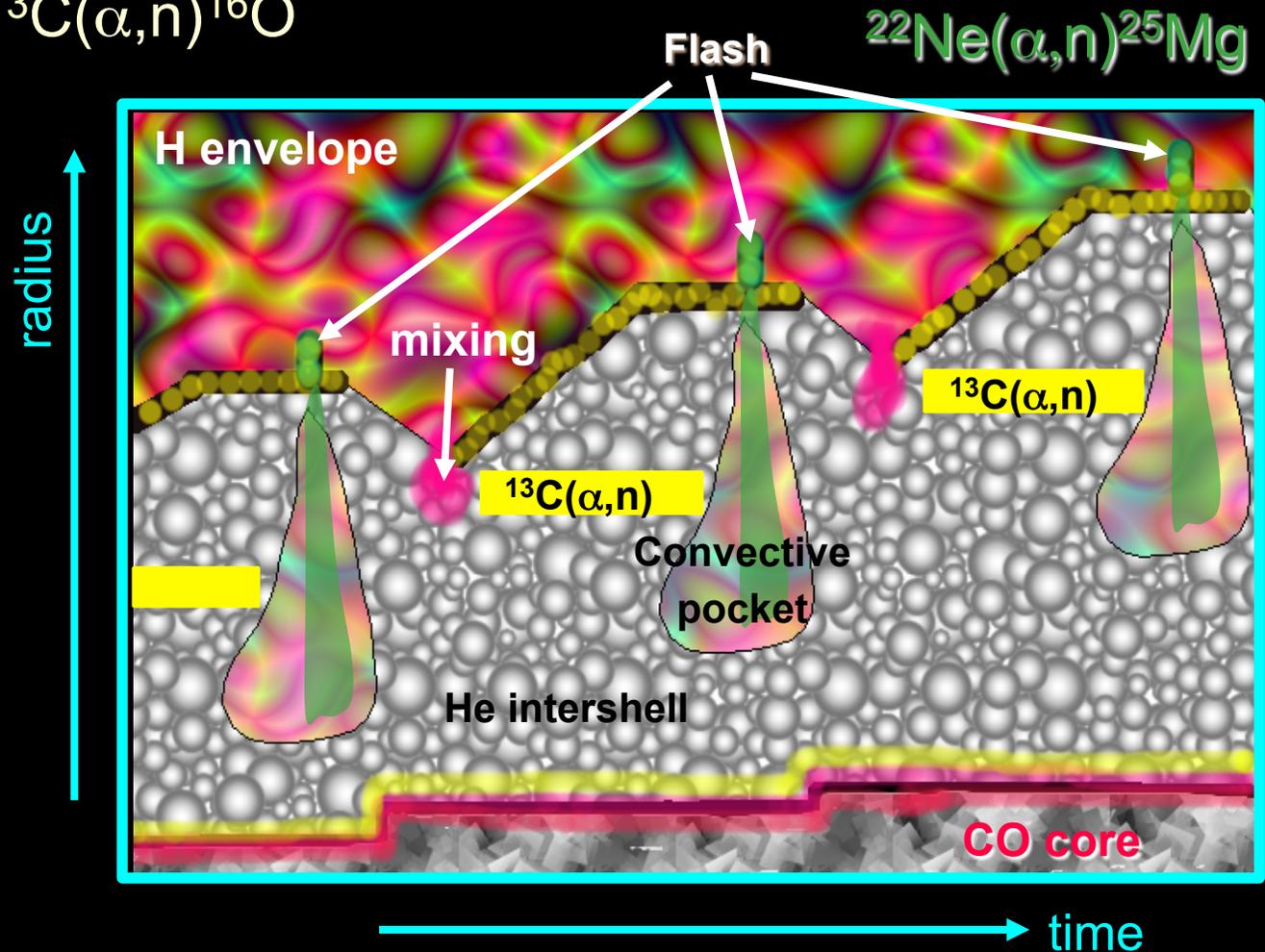
Thermally unstable: mixing, convection, mass loss → nebulae



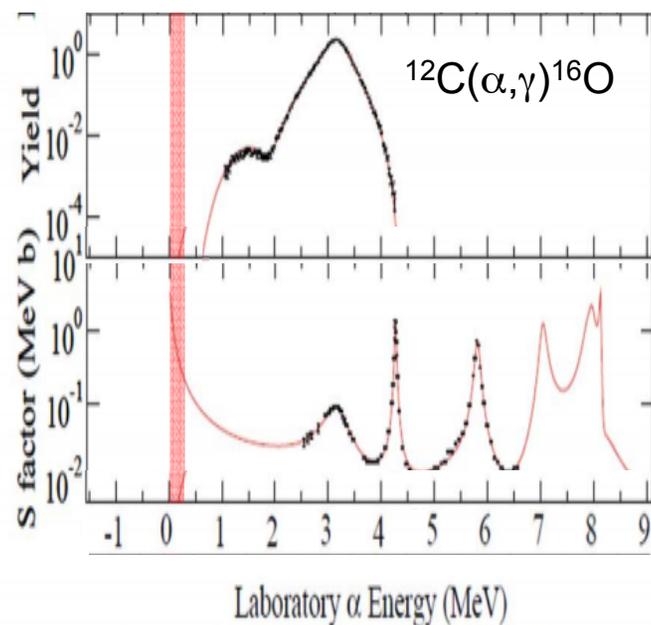
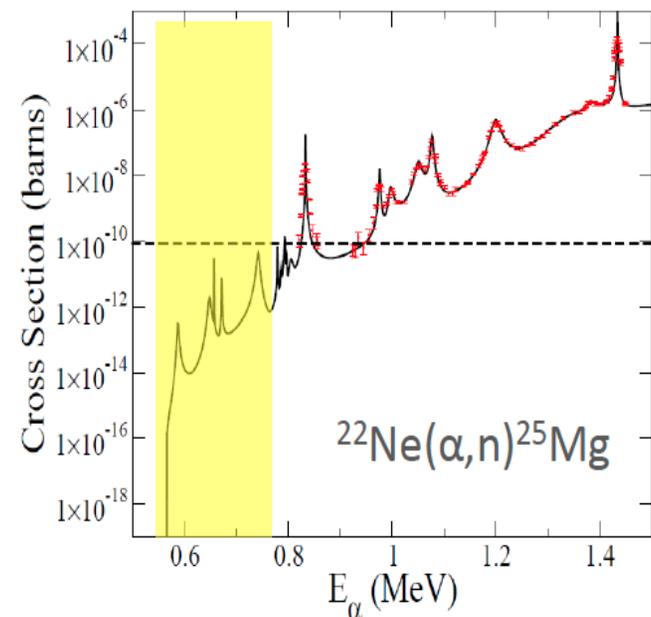
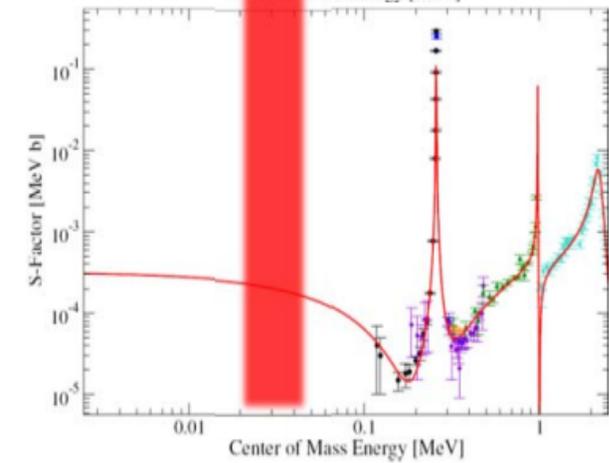
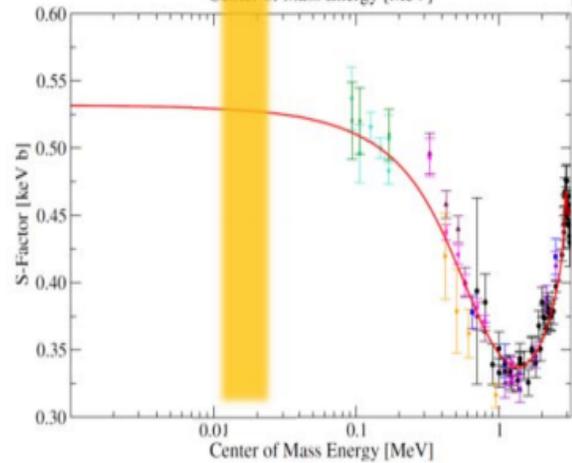
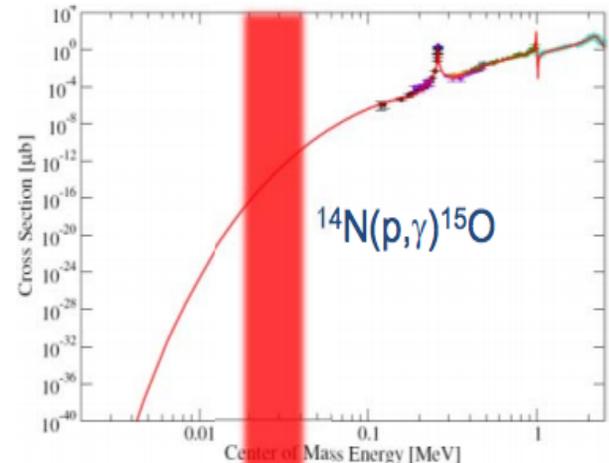
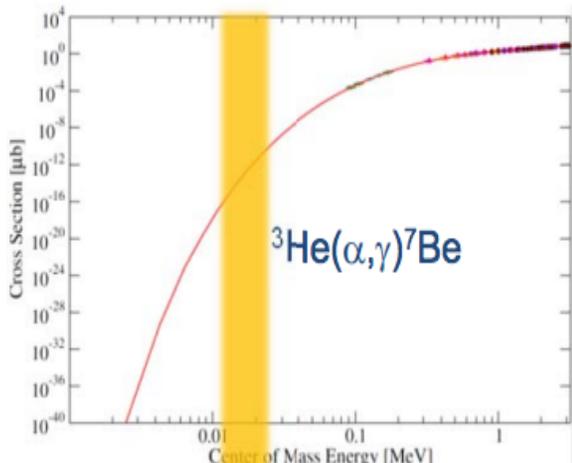
Neutrons drive synthesis of heavy elements



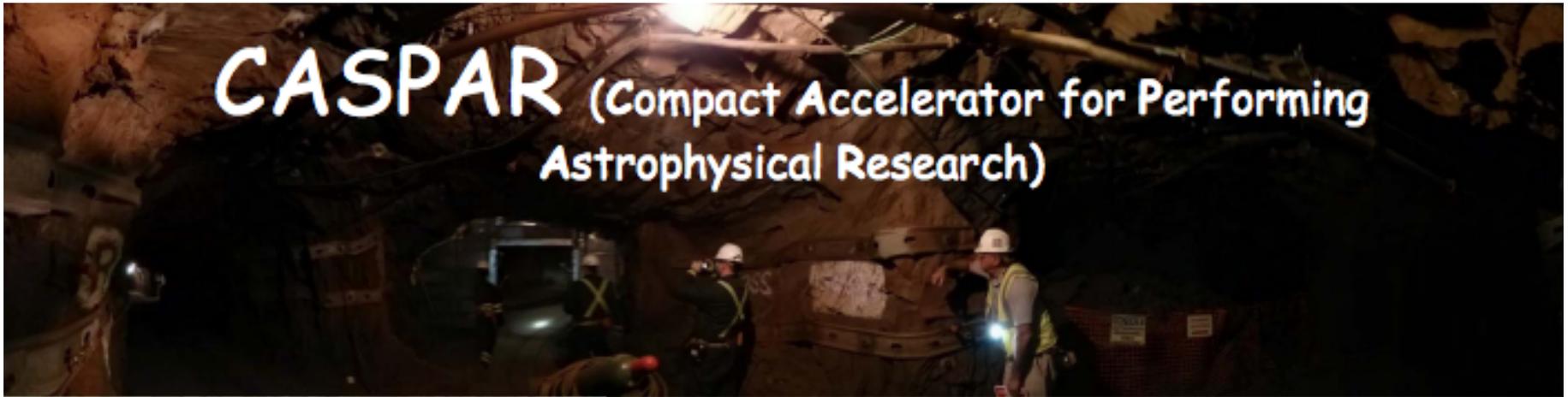
CO core remnant (white dwarf)



Challenging measurements at lower energies needed to understand important reaction rates



CASPAR (Compact Accelerator for Performing Astrophysical Research)

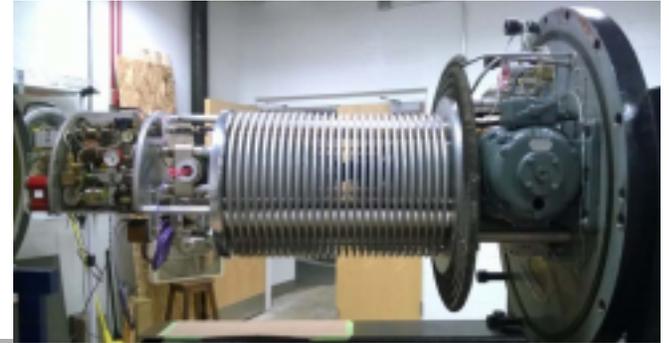


ACCELERATOR CAVERN

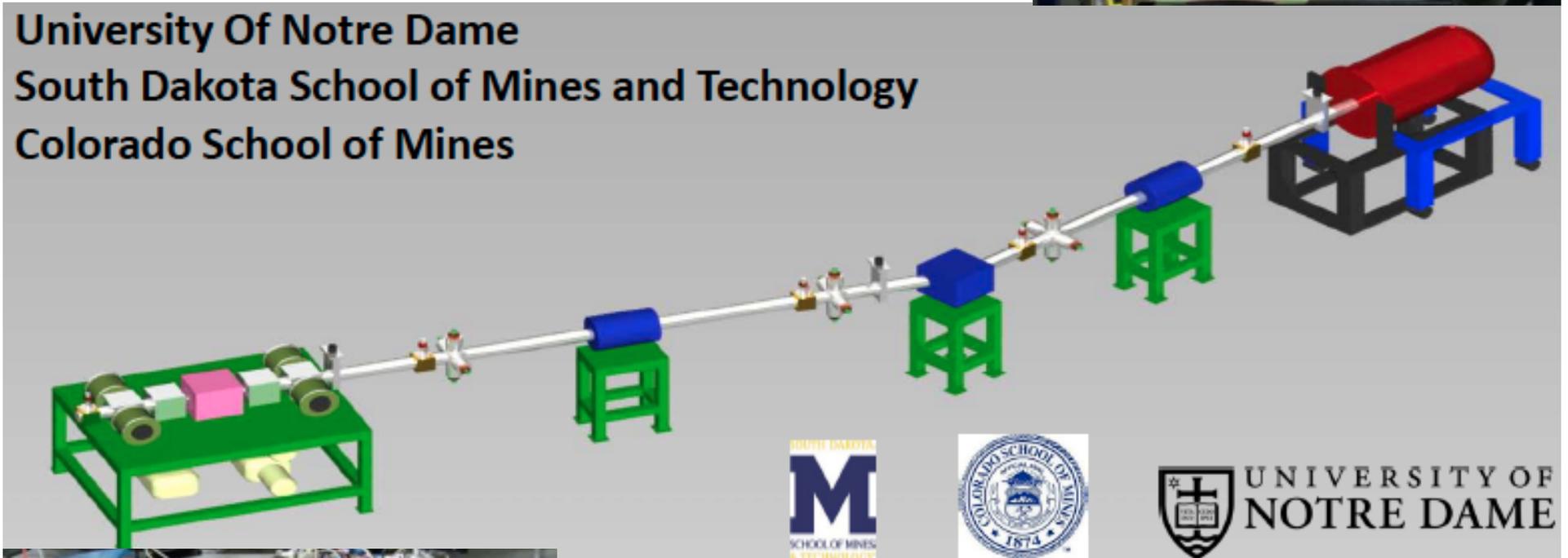


Sanford Underground Facility
in Homestake Mine, South Dakota

High current JN accelerator



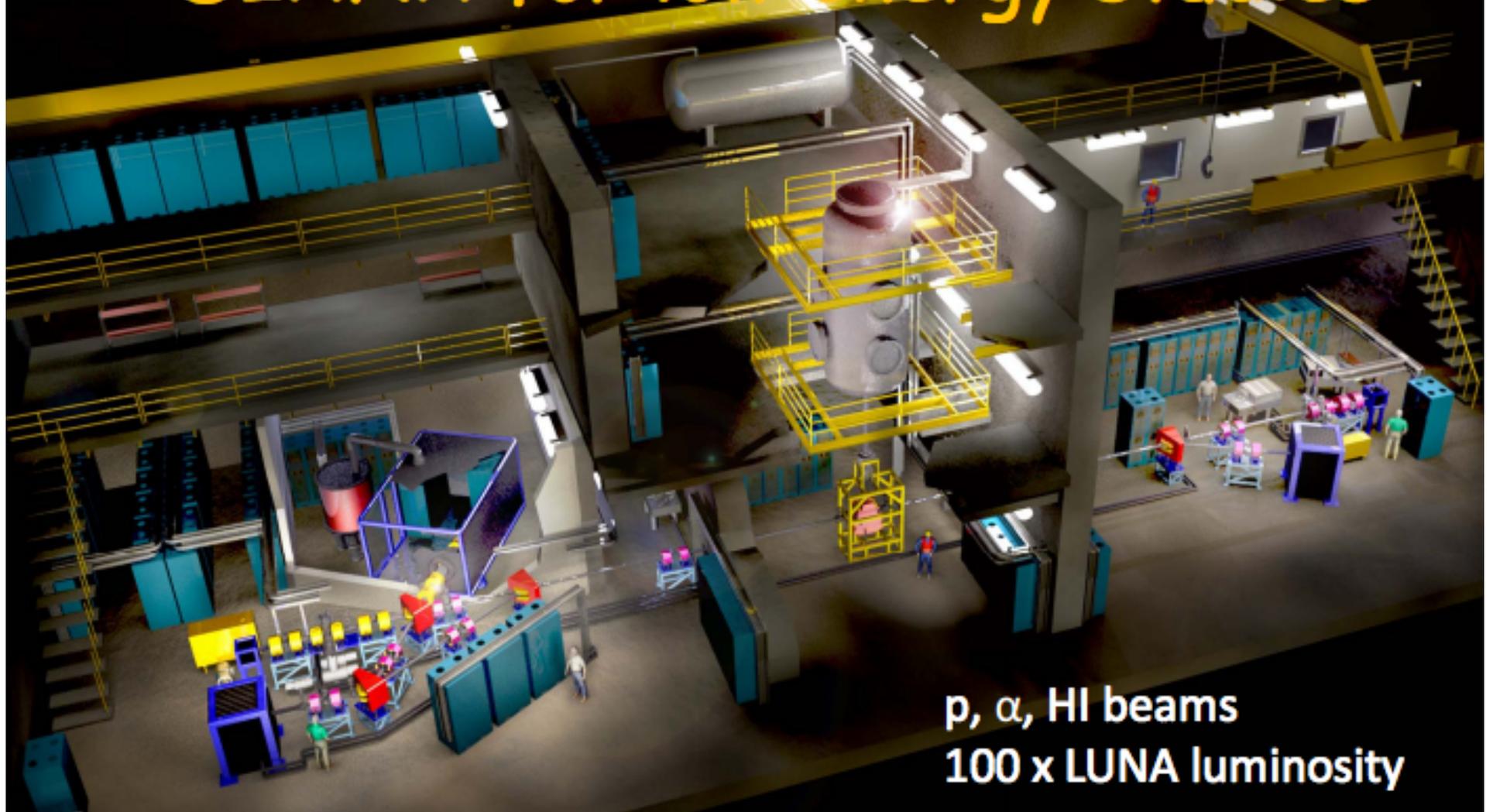
University Of Notre Dame
South Dakota School of Mines and Technology
Colorado School of Mines



Windowless gas target



Underground accelerator project DIANA for low energy studies

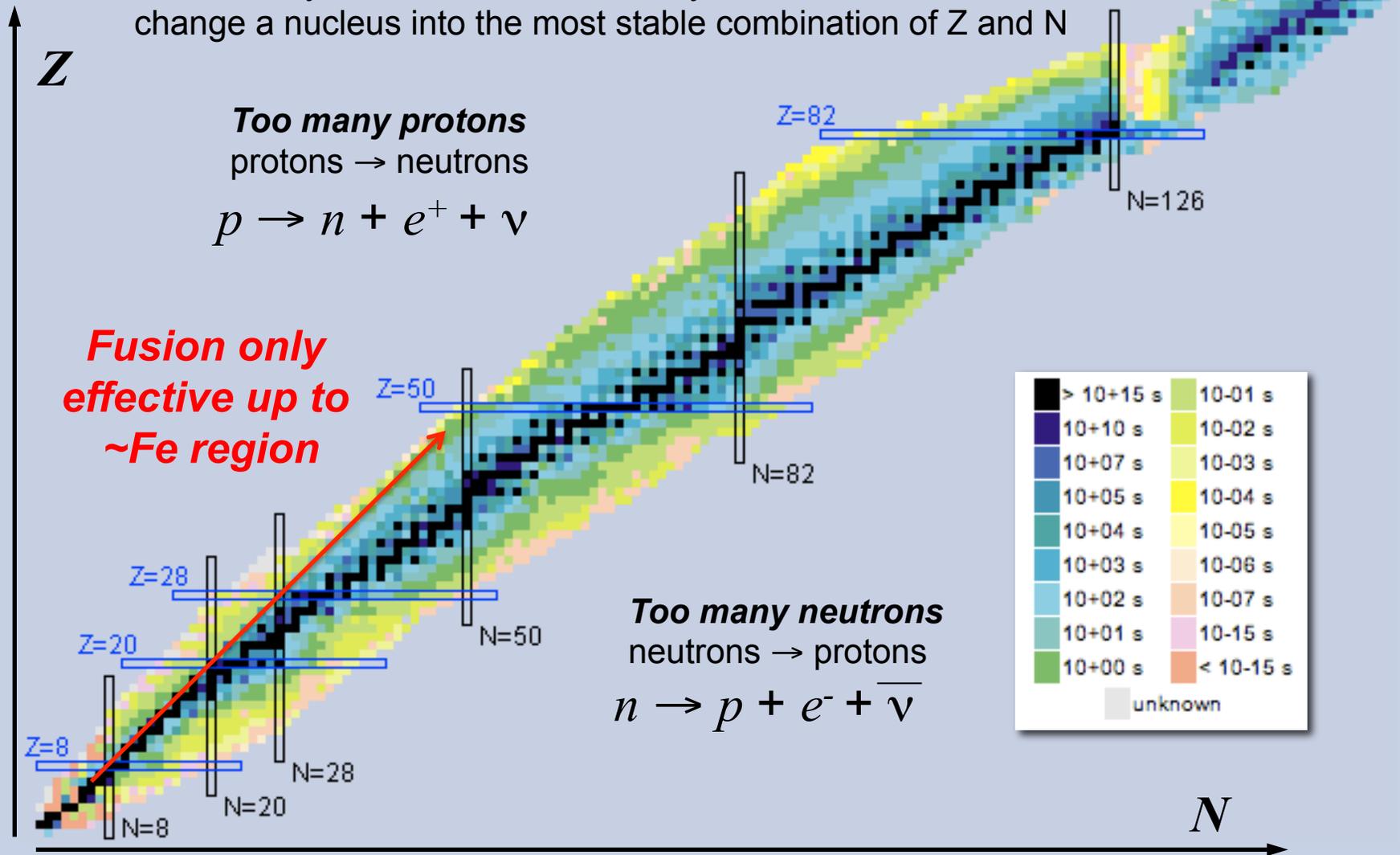


p, α , HI beams
100 x LUNA luminosity

High luminosity, low background experiments

Generally only one combination of protons and neutrons is stable for each A (smallest mass)

“Beta decay” is a form of radioactivity that nature uses to change a nucleus into the most stable combination of Z and N



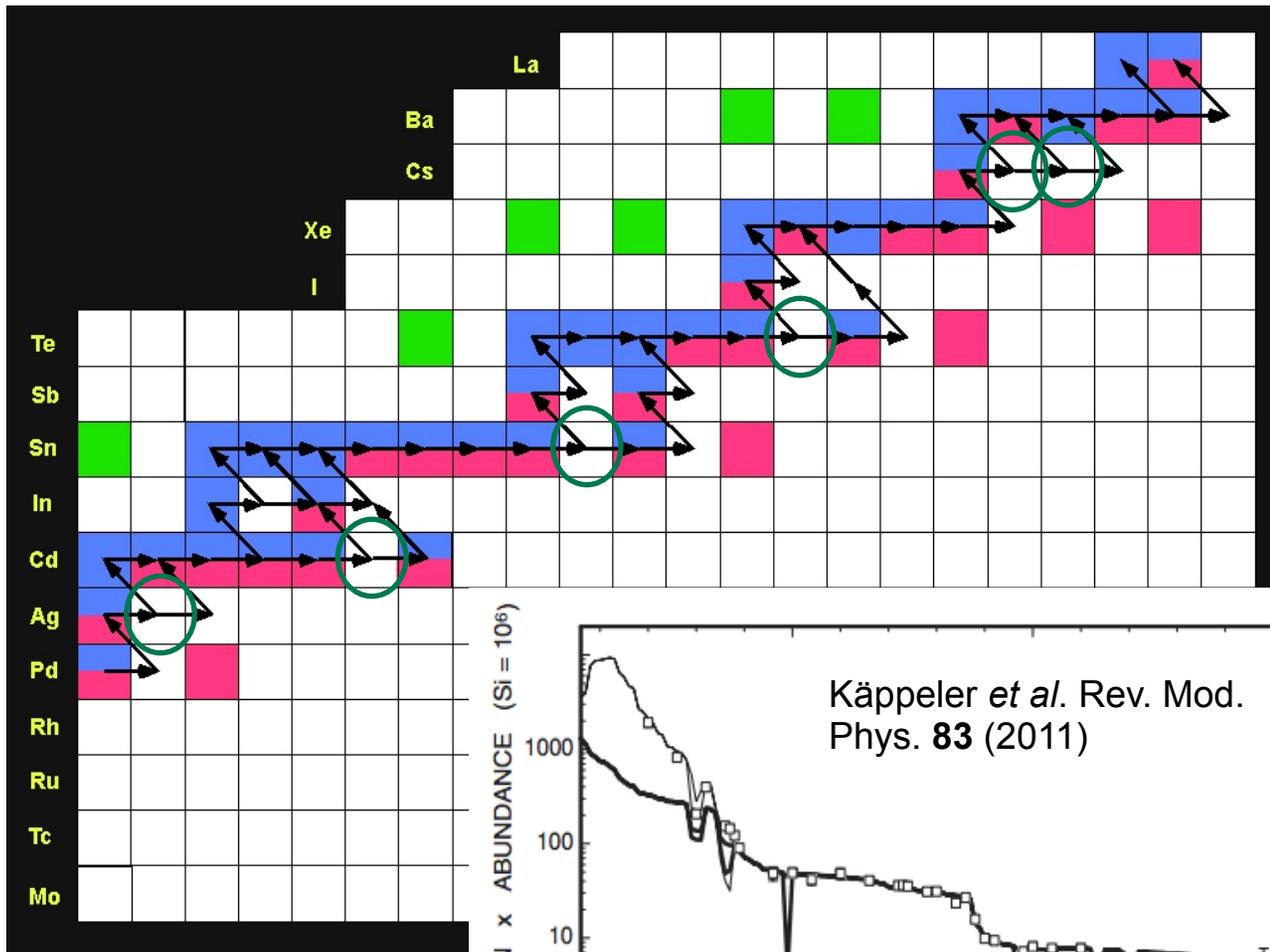
Too many protons
 protons → neutrons
 $p \rightarrow n + e^+ + \nu$

Fusion only effective up to ~Fe region

Too many neutrons
 neutrons → protons
 $n \rightarrow p + e^- + \bar{\nu}$

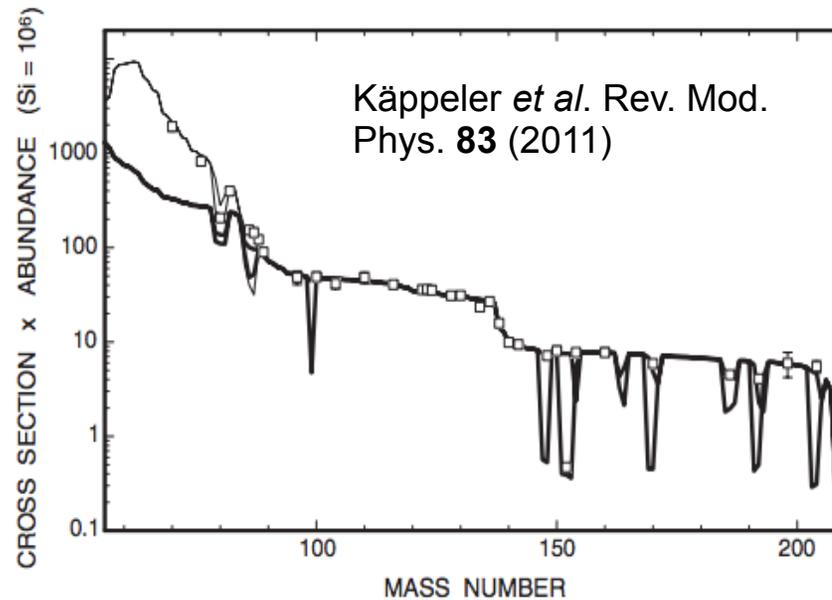
> 10 ⁺¹⁵ s	10 ⁻⁰¹ s
10 ⁺¹⁰ s	10 ⁻⁰² s
10 ⁺⁰⁷ s	10 ⁻⁰³ s
10 ⁺⁰⁵ s	10 ⁻⁰⁴ s
10 ⁺⁰⁴ s	10 ⁻⁰⁵ s
10 ⁺⁰³ s	10 ⁻⁰⁶ s
10 ⁺⁰² s	10 ⁻⁰⁷ s
10 ⁺⁰¹ s	10 ⁻¹⁵ s
10 ⁺⁰⁰ s	< 10 ⁻¹⁵ s
unknown	

Slow neutron capture (s) process



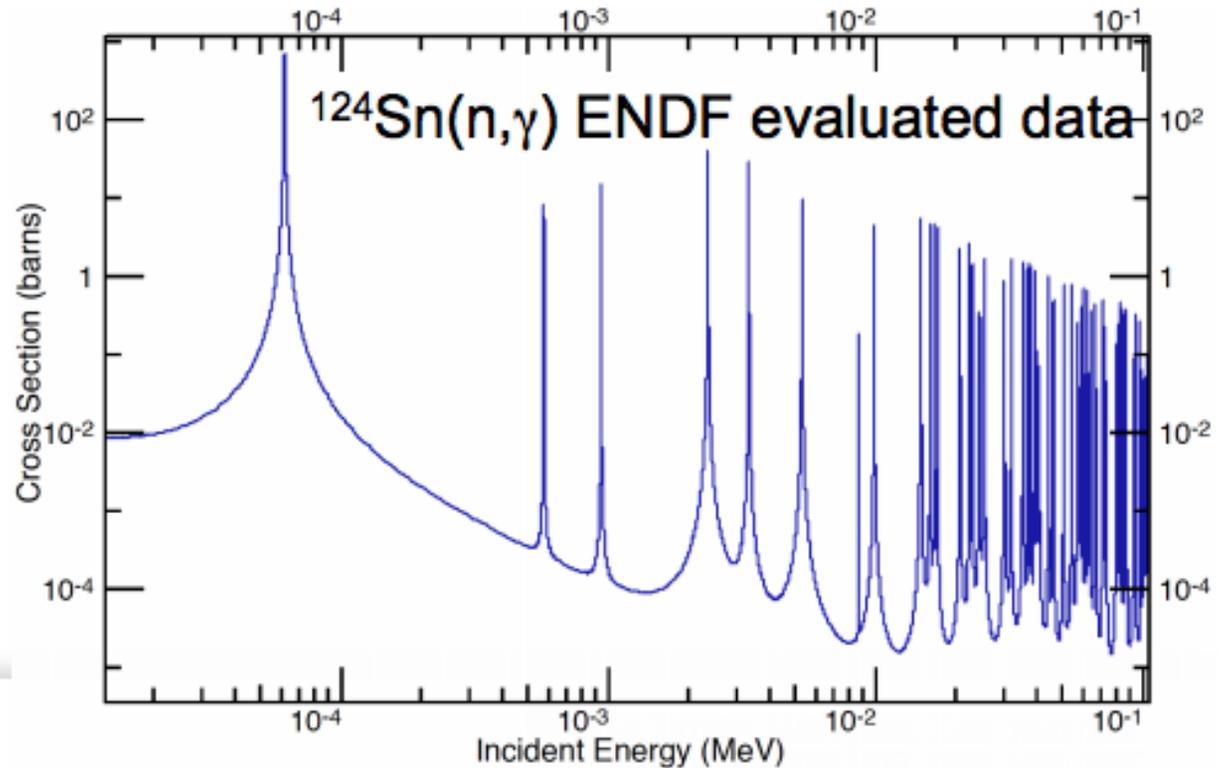
- **s process**

- Produces about half of matter that is heavier than iron
- Series of slow neutron captures
- Pattern of isotopes produced is generally well understood
- Most σ 's measured



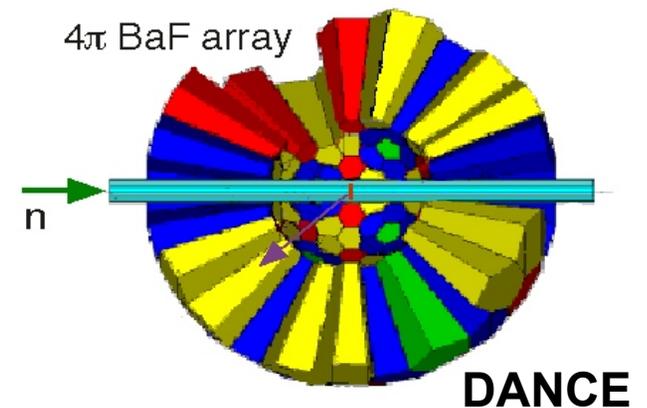
(n,γ) cross sections for the s process

- Good data on most stable isotopes
- Spallation n sources
- TOF techniques
- Good energy resolution
- Usually high level densities



Some outstanding issues

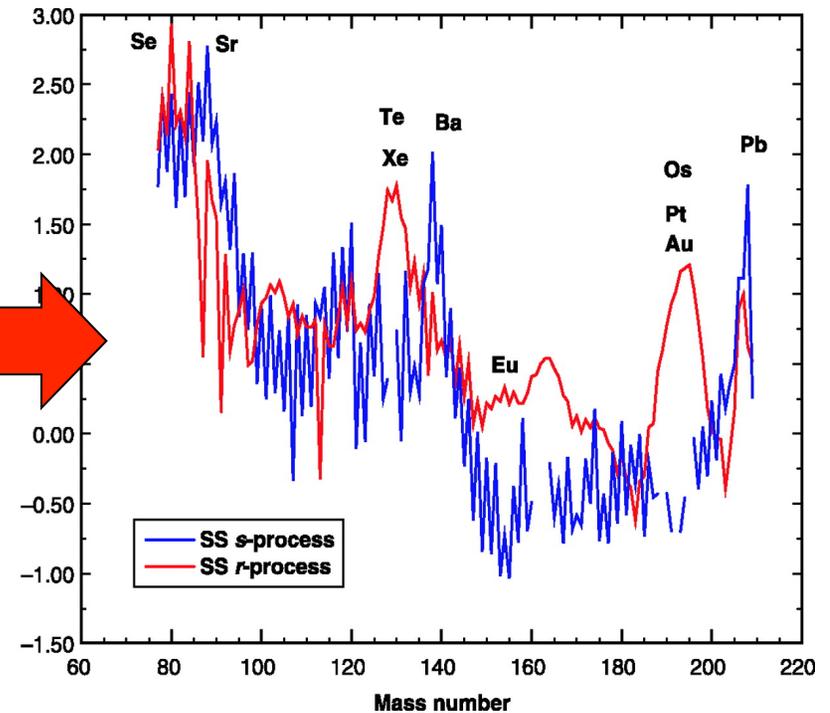
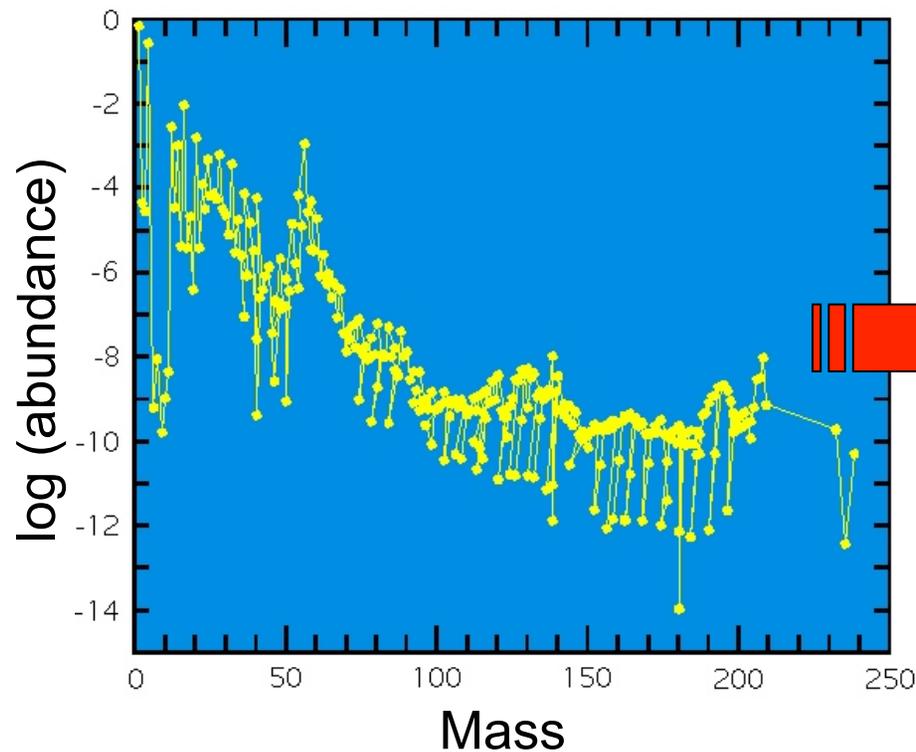
- Influence of low-energy levels at low temp
- Direct capture near closed shells
- Effect of thermal excitations in stellar environmen
- Branch point isotopes
- Lighter elements – “weak s process”



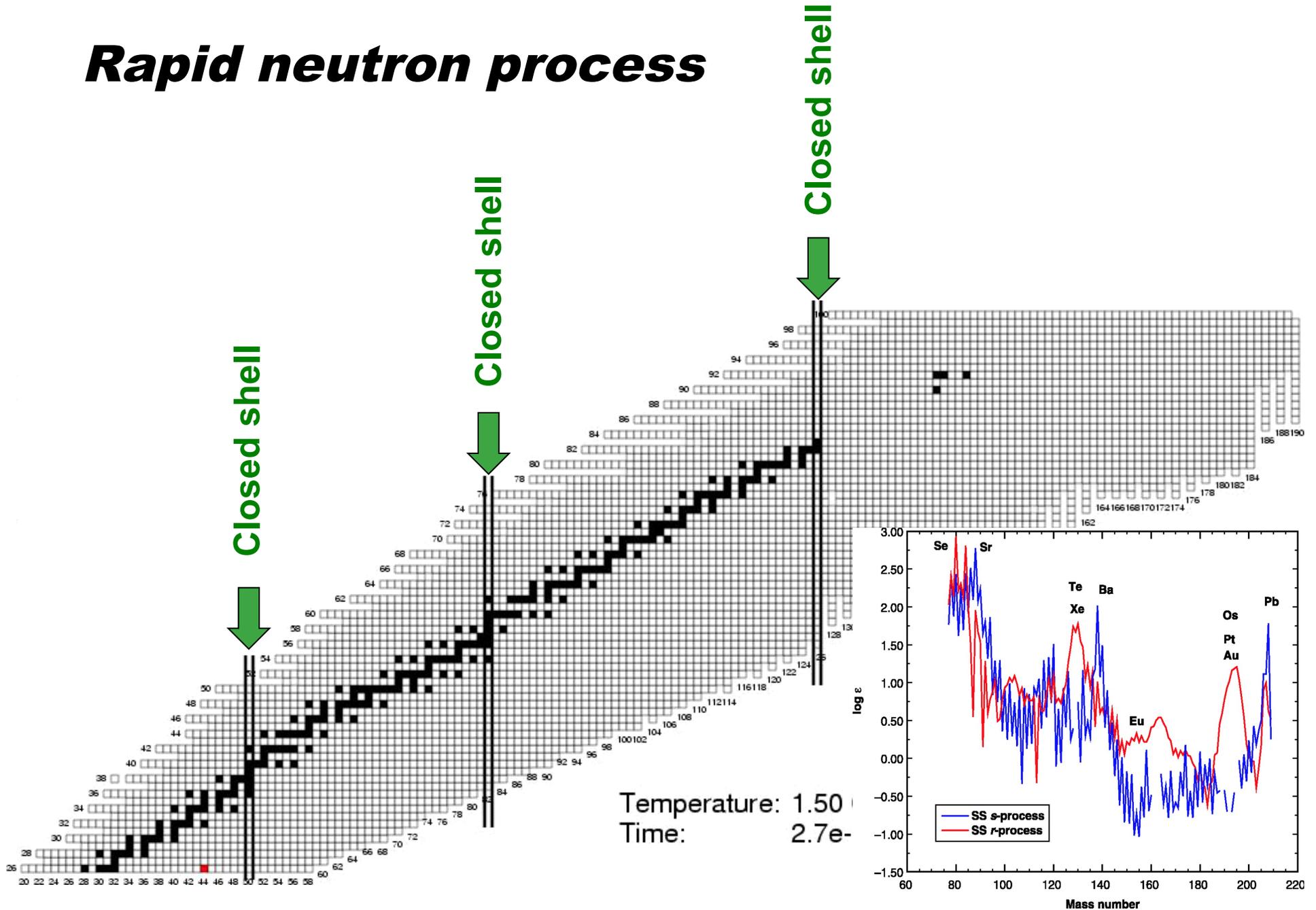
What can AGB stars produce?

Relative abundances of s process isotopes well understood

Subtract s abundances from solar system to get remainder

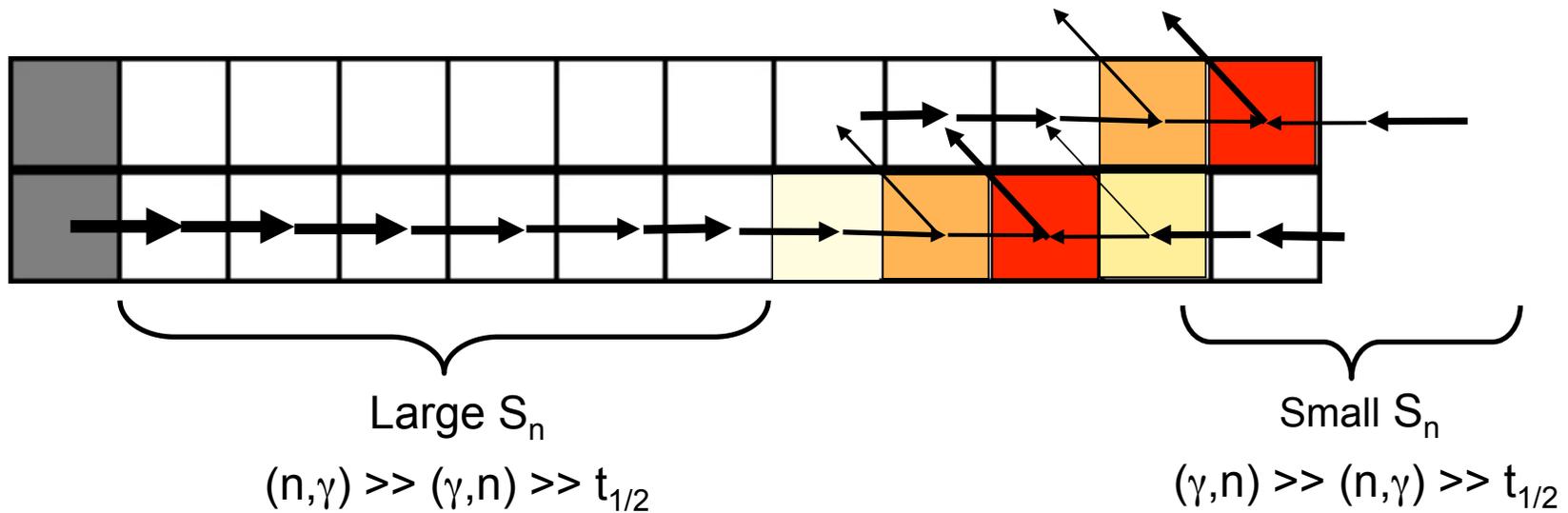


Rapid neutron process



Cartoon r process

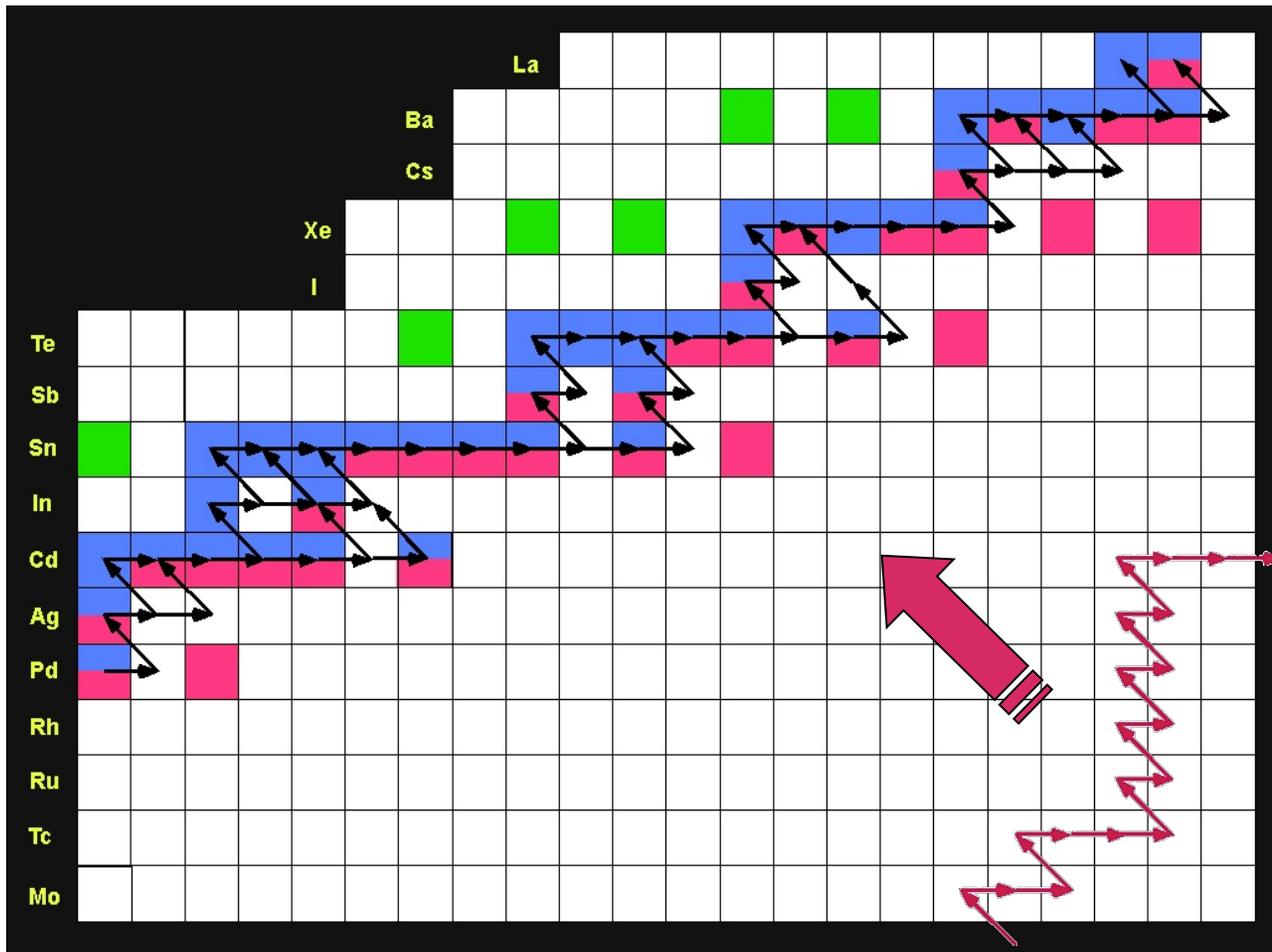
$$\frac{Y(A+1)}{Y(A)} \approx \frac{1}{2} \left(\frac{2\pi\hbar^2}{m_u kT} \right) n_n e^{S_n/(kT)}$$



- > Free parameters n_n, kT, t
- > freezeout relatively fast & decay back to stability

➡ Most important: masses, $t_{1/2}$, and P_n

Synthesis of heavy elements



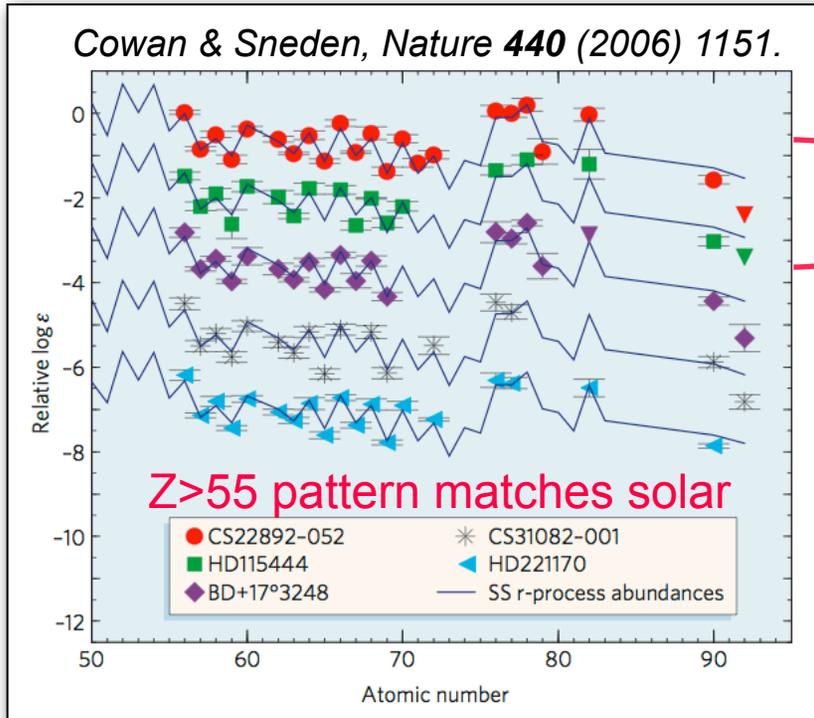
- s process

- r process

- Produces about half of matter heavier than iron
- High neutron flux
- Reactions on unstable isotopes
- Site unknown

r process in the early Galaxy

New observations of unmixed abundances early in the Galactic halo



CS22892-052

Fe/H = (8×10^{-4}) solar = very old
r/Fe = 50 solar

Only 2 known in 2000

Now extensive surveys

e.g. see Frebel et al., *ApJ* **652** (2006) 1585

SEGUE (Sloan DSS)

Spectra of $> 2 \times 10^5$ selected halo stars

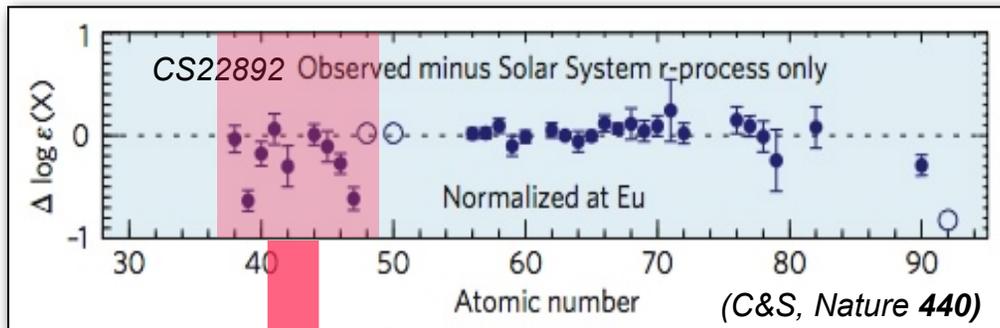
Expect $\sim 1\%$ with Fe/H < 0.001 solar

~ 36 known r process stars

11 with r/Fe > 10 solar

Distribution Fe/H puzzling

Lowest Fe/H stars intriguing



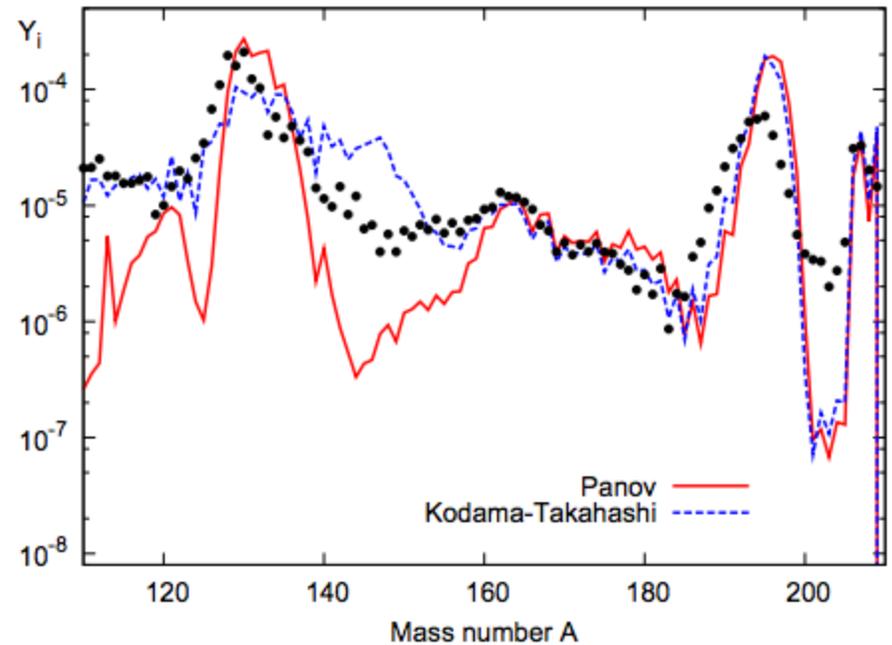
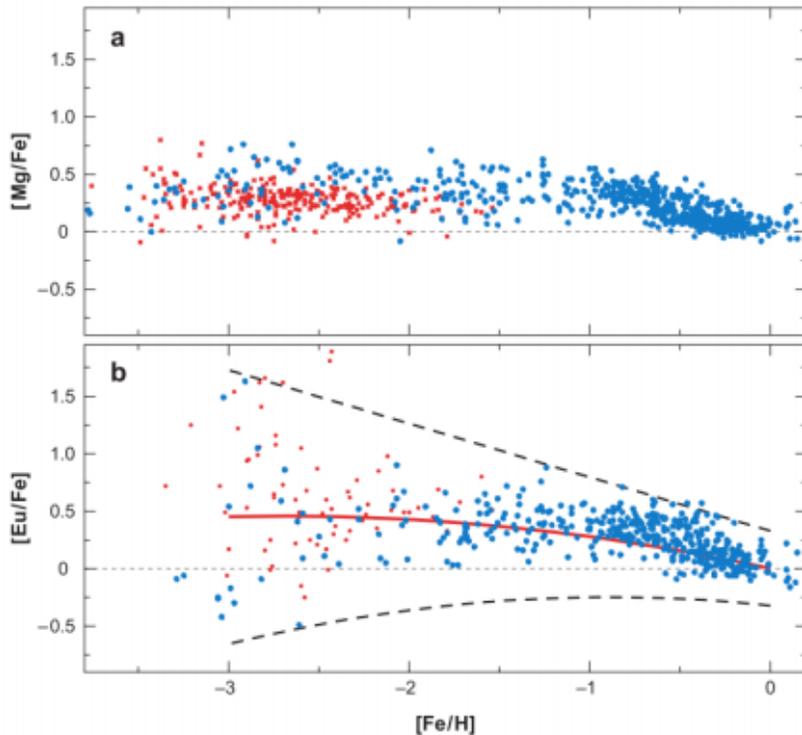
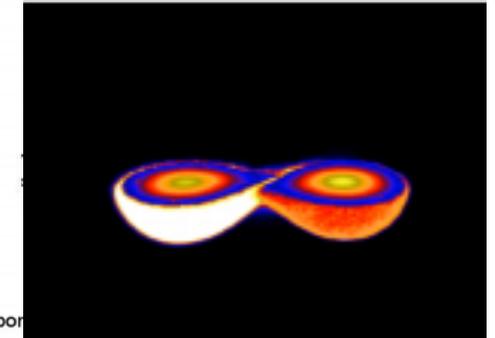
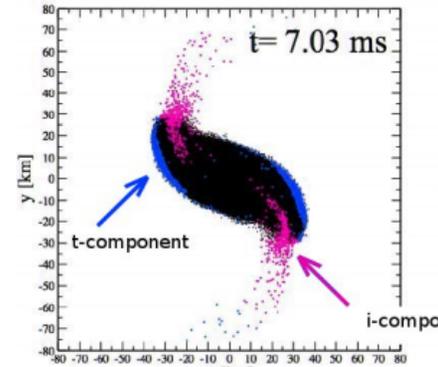
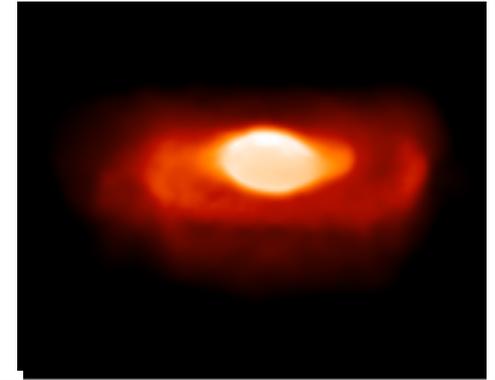
Z < 50 abundances vary

**An additional process
(besides r/s) must
contribute significantly to
elements from Fe-Sn**

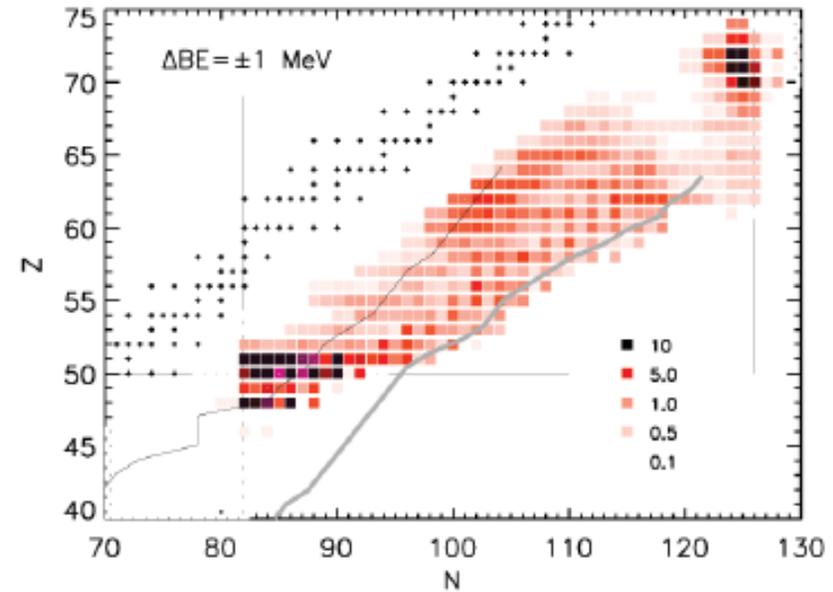
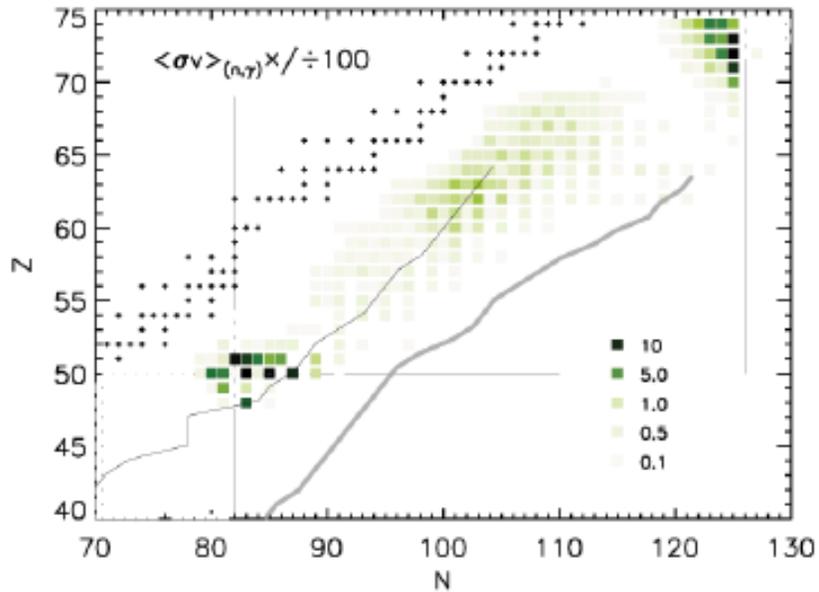
r process, where?

Neutron star mergers?

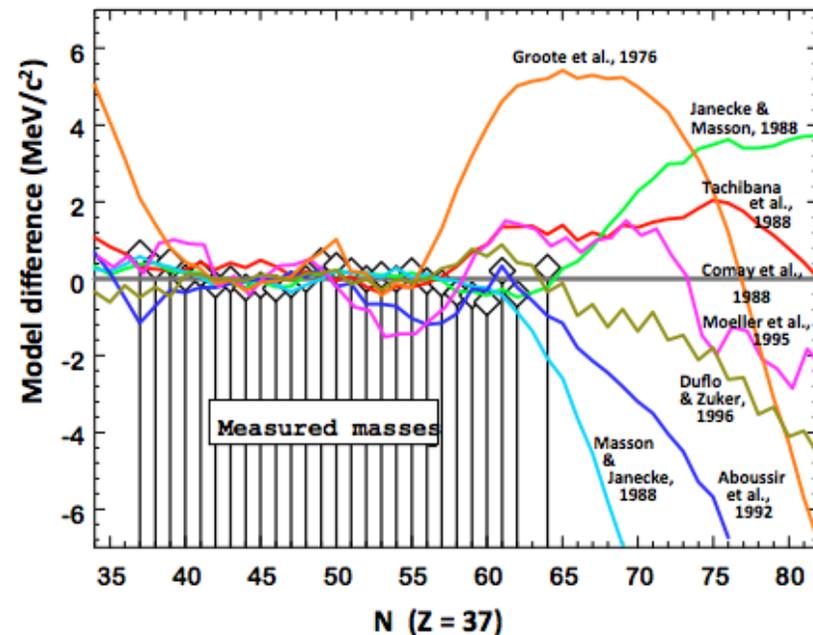
- 10 systems known in our Galaxy
- Believed to be origin of gamma-ray bursts
- Potentially explains dispersion in Eu
- High output in rare event
- Calculations produce a robust r process
- Huge uncertainties due to nuclear data



r process: masses and reaction rates



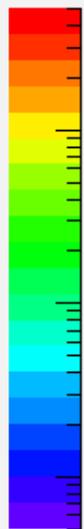
- Some (n,γ) capture rates are important
- Abundances are very sensitive to most atomic masses and decay properties
- Most mass models do not reliably extrapolate away from stability
- Need measurements of nuclear properties in neutron-rich nuclei



Atomic masses

- About 2400 isotopes have measured masses
- Average precision better than 0.1 ppm

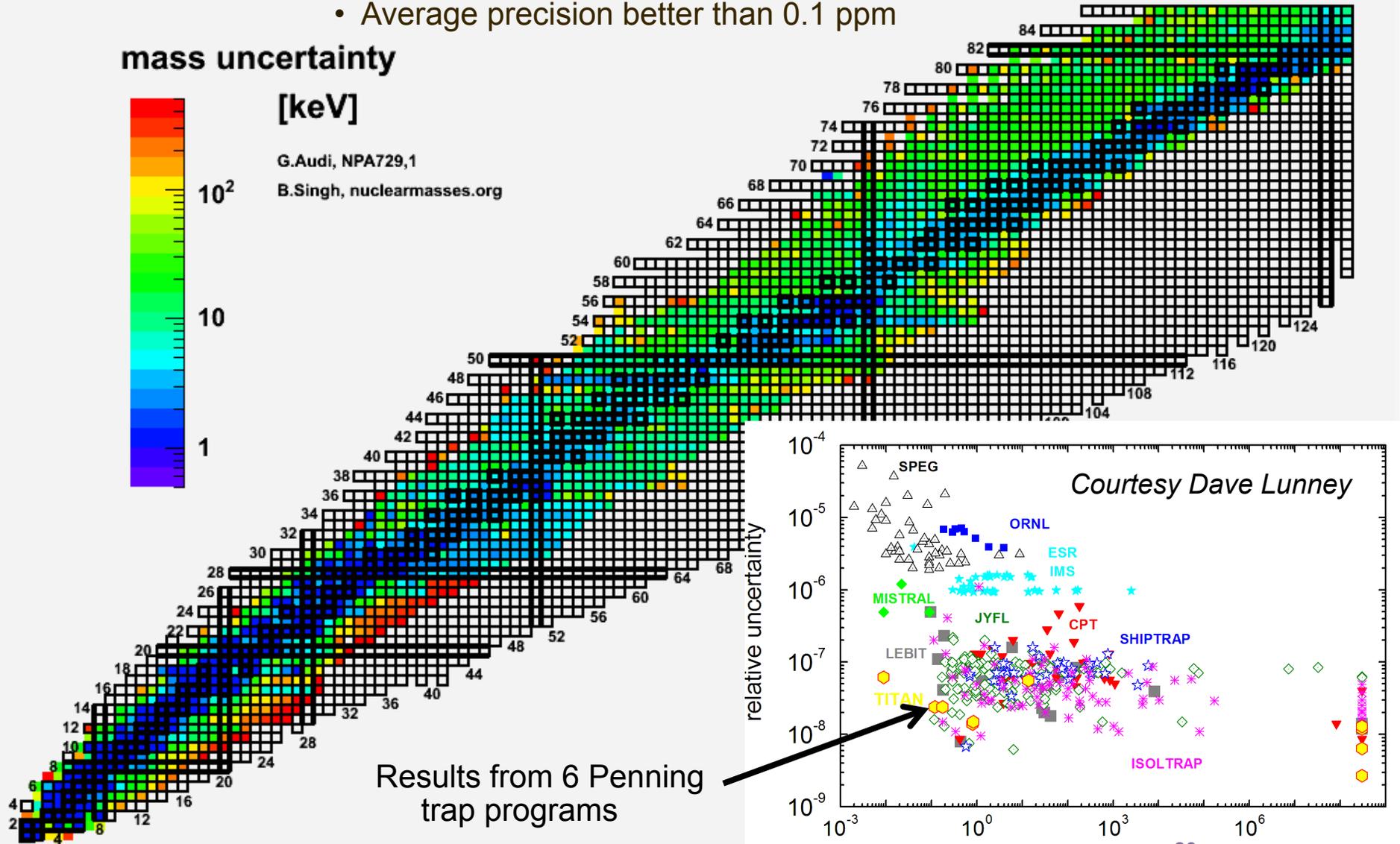
mass uncertainty



[keV]

G.Audi, NPA729,1

B.Singh, nuclearmasses.org



Results from 6 Penning trap programs

Courtesy Dave Lunney

half life (seconds)

26

Penning Traps: In a Nutshell

Trap electrodes in a hyperboloid geometry ($r_0/z_0 \approx 1.16$)

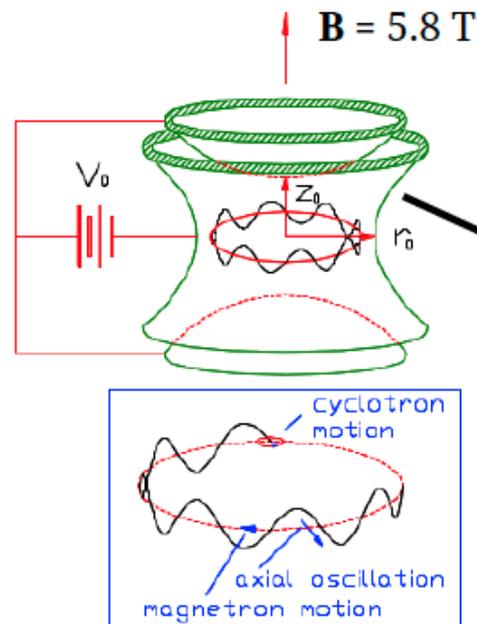
Placed inside uniform magnetic field

Three types of motion:
 - Axial (ω_z)
 - Reduced cyclotron (ω_+)
 - Magnetron motion (ω_-)

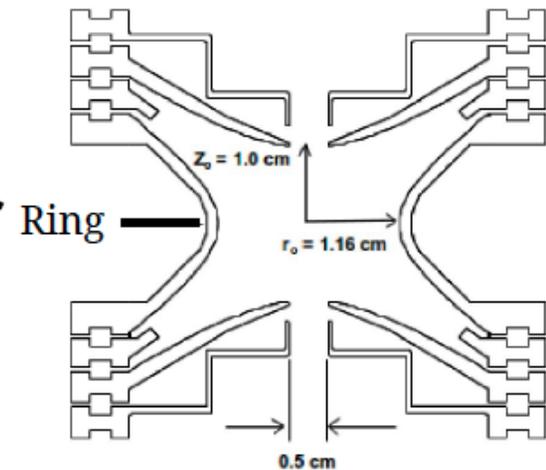
Drive trapped ions into an excitation using RF signal

Ions are ejected from the trap and measure their time-of-flight to a detector

Resonant enhancement at the ion cyclotron frequency: ω_c

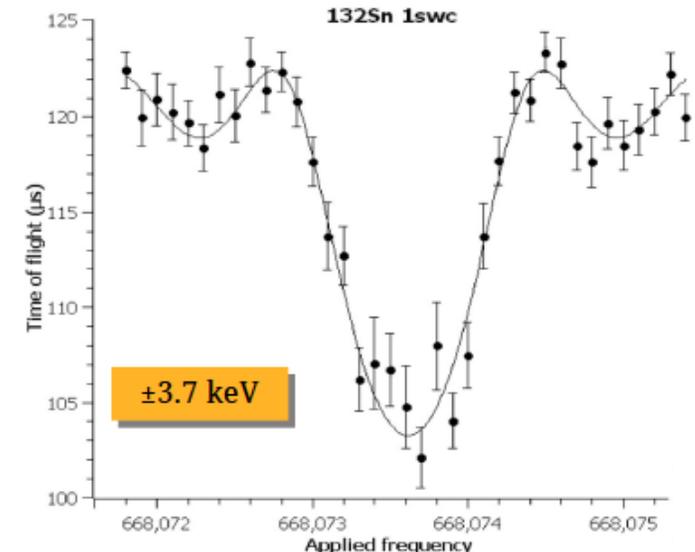


Cross section of CPT

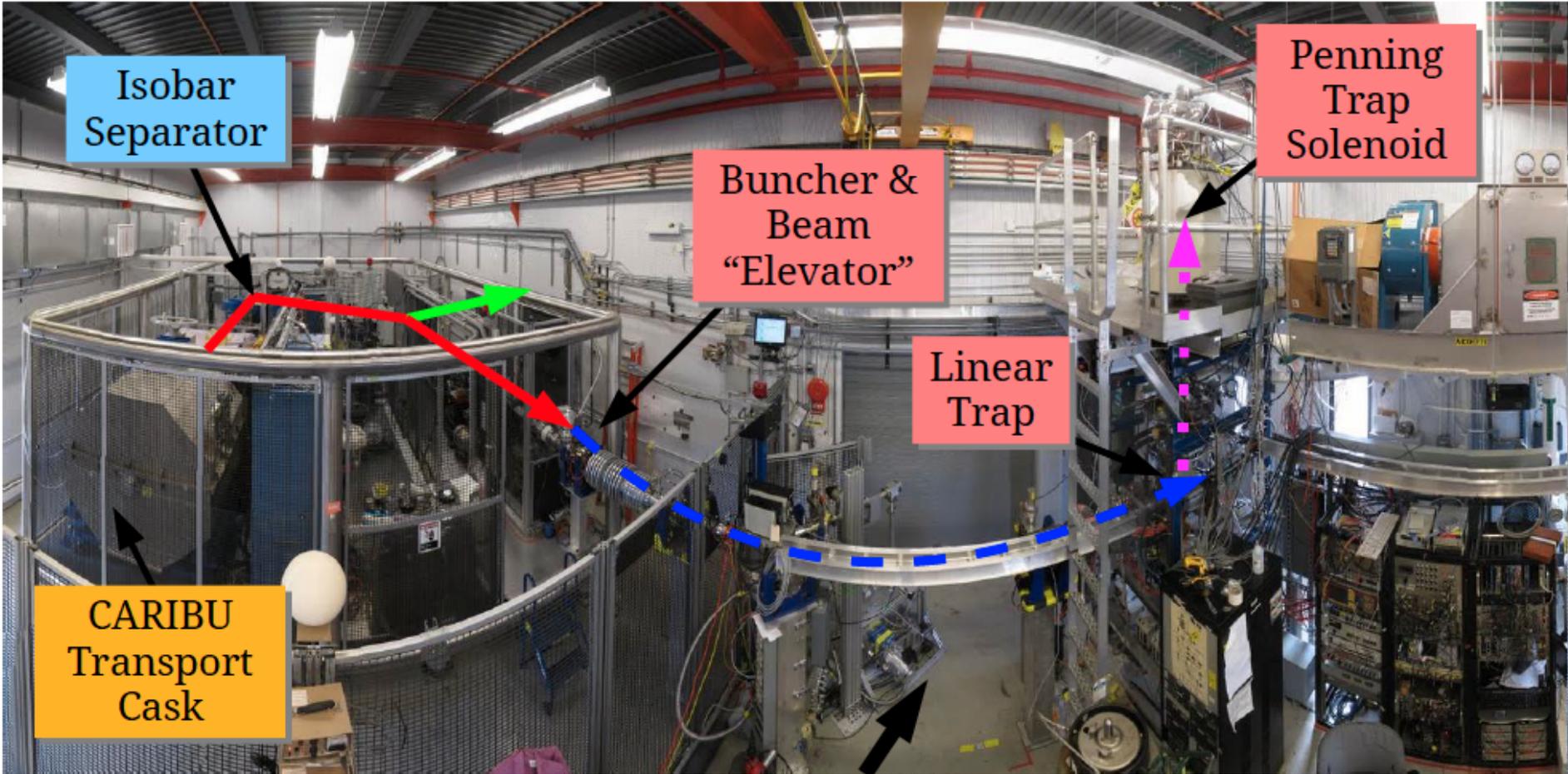


The mass of the trapped ions are measured indirectly by determining the cyclotron frequency:
 $\omega_c = qB/m$

Time-of-Flight vs. Applied Frequency



Canadian Penning Trap
Argonne National Laboratory
ATLAS Accelerator Facility



Precision Mass Measurements

$$\text{Precision} = \frac{\Delta m}{m} \propto \frac{m}{TqB\sqrt{N}}$$

Conversion Time
Time spent exciting the ion in the trap.
Limited by isotope half-life!!!

Ion Charge State
Typically $q=1^+$
For some isotopes $q=2^+$; More precise measurement

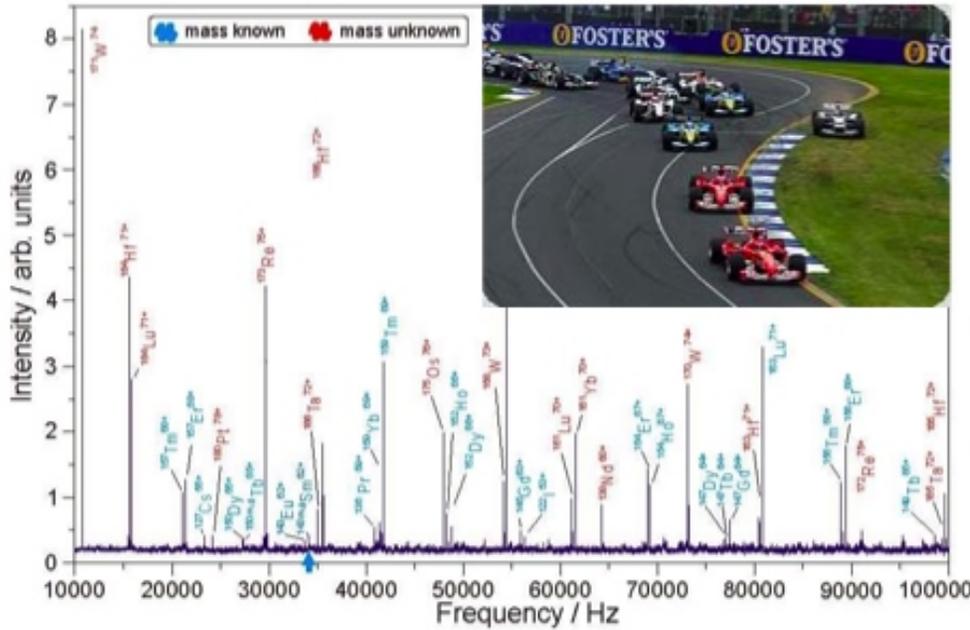
Magnetic Field
Measured with a reference ion ($^{133}\text{Cs}^+$) of known mass in between mass measurements

Statistics
The longer you acquire data, the better the measurement.
CARIBU transport efficiency!

Many of the most neutron-rich nuclei produced at CARIBU have short half lives (≤ 150 ms) as well as small fission branches from ^{252}Cf ($\leq 10^{-4}$ %)

Improvements in both CARIBU and the CPT systems are required to perform mass measurements of influential r -process nuclei to $\Delta m/m \leq 10^{-7}$

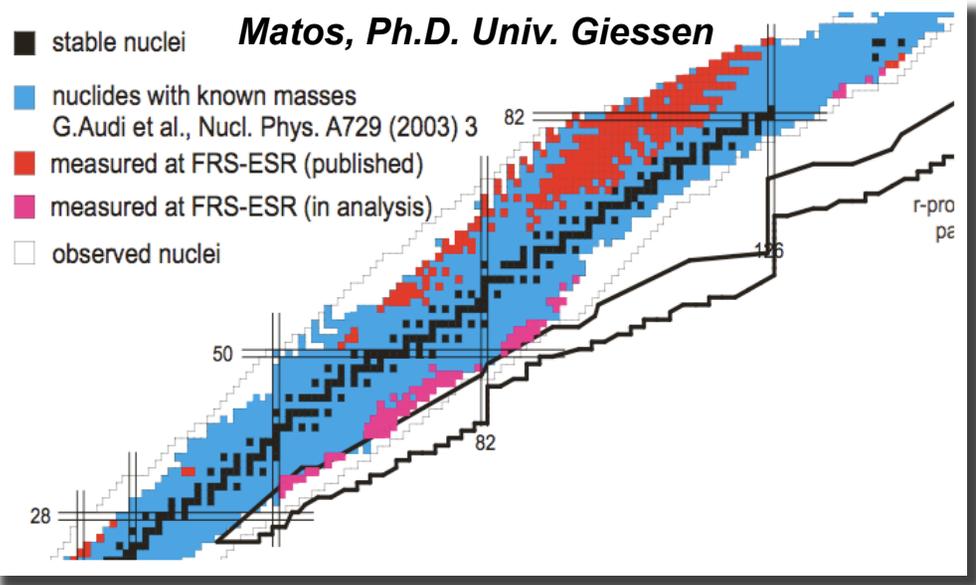
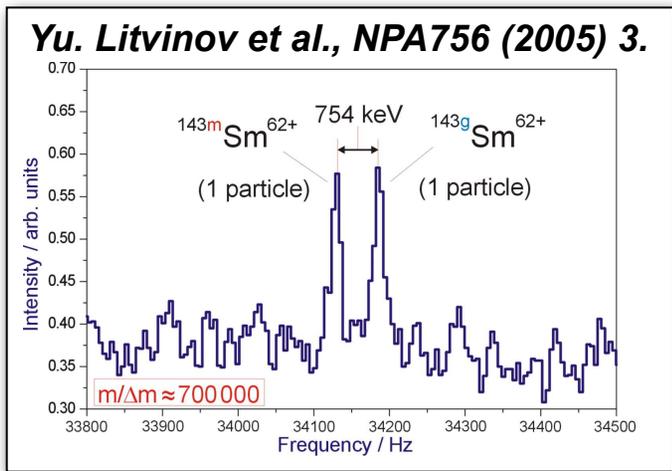
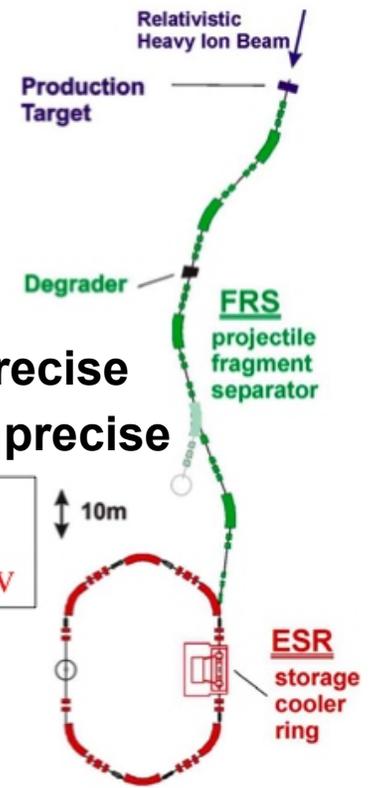
Mass measurements – storage rings



2 modes:
 Schottky - slow, more precise
 isochronous - fast, less precise

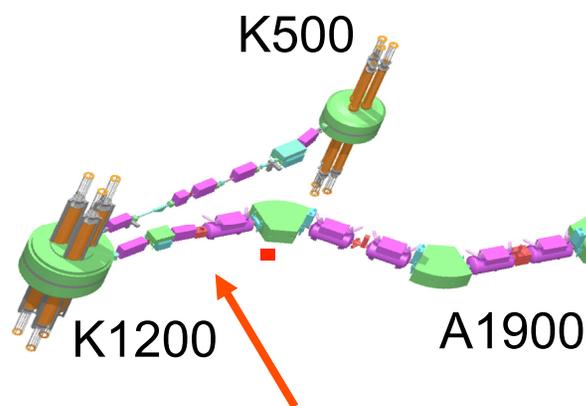
Experimental Storage Ring:

$$\Delta m/m = \gamma_t^2 \Delta f/f + (\gamma_t^2 - \gamma^2) \Delta v/v$$



TOF- $B\rho$ Mass Measurements

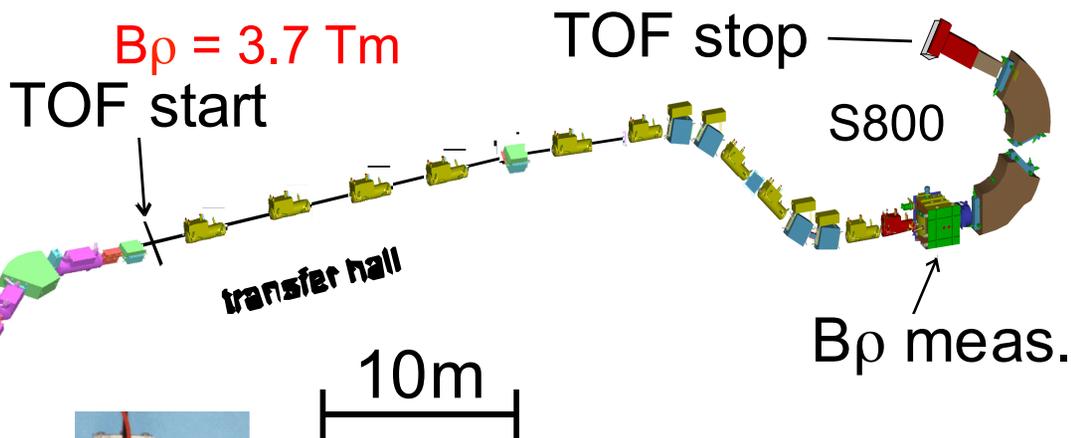
primary beam
 ^{86}Kr 100 MeV/u



production target
 Be 51 and 94 mg/cm²

$$B\rho = \frac{m}{q} v \quad \rightarrow \quad \text{mass}$$

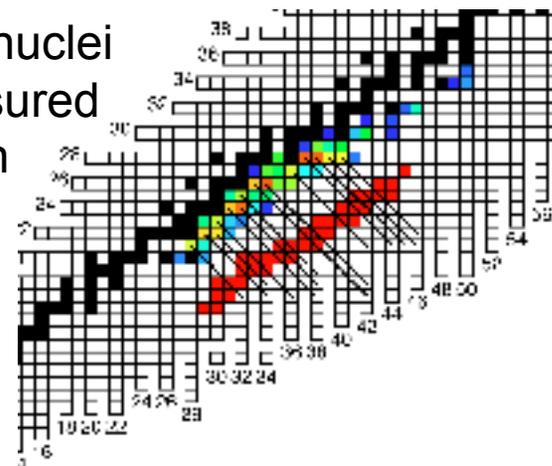
No dependence on lifetime!
 Path length ~ 58m
 Time resolution ~ 30ps



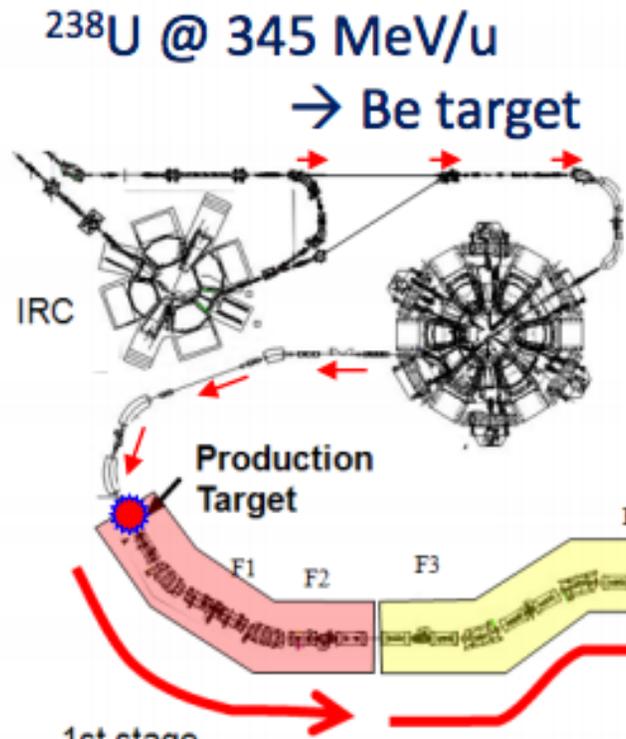
Nov 2011 – Fragmentation of ^{76}Ge
 Estrade, George, Matos, Schatz *et al.*

Masses of 30+
 neutron-rich nuclei
 near Ni measured
 with precision
 better than
 1:200,000

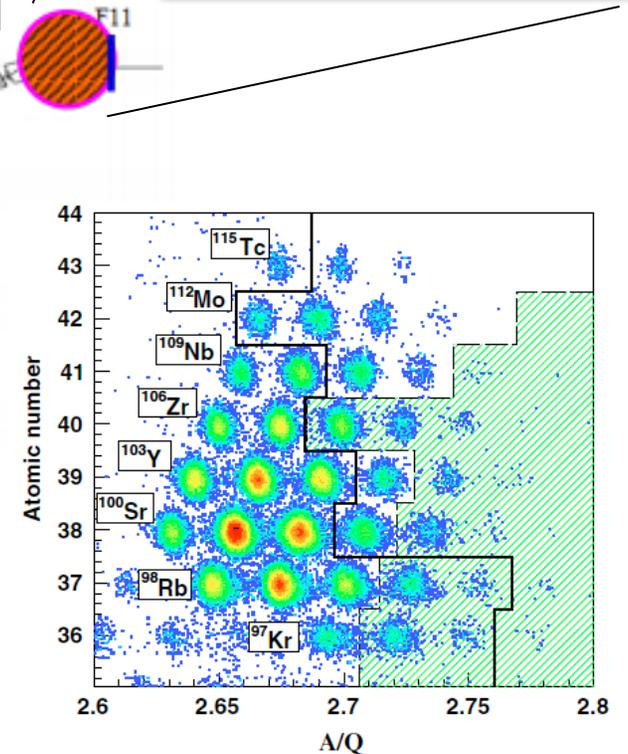
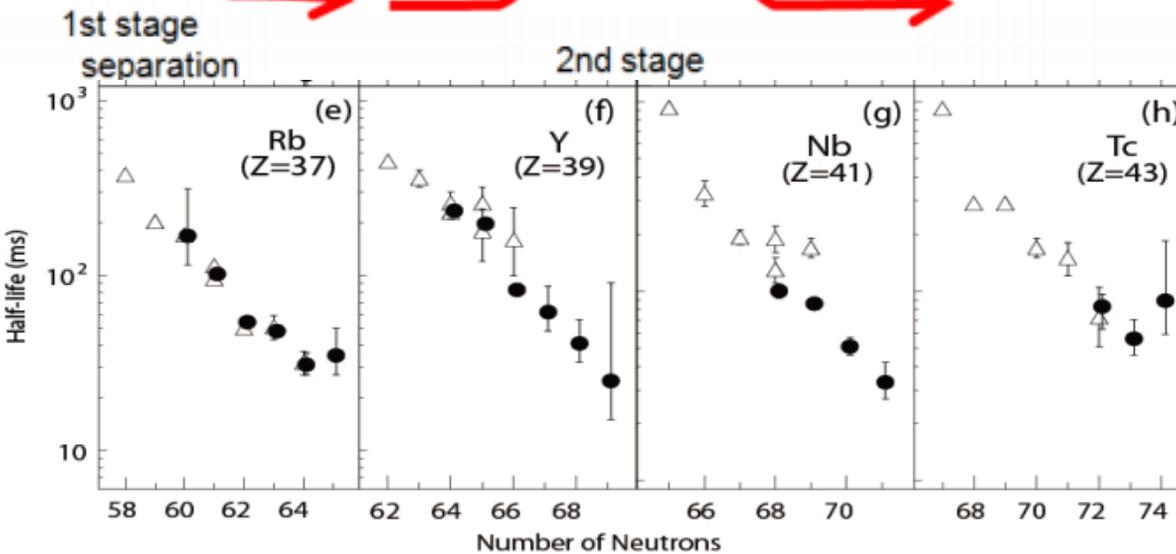
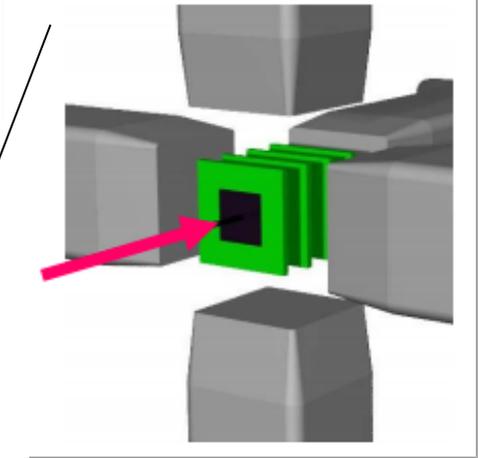
New approved exp.
Fragmentation of ^{124}Sn



Beta Decay Example: RIBF @ RIKEN

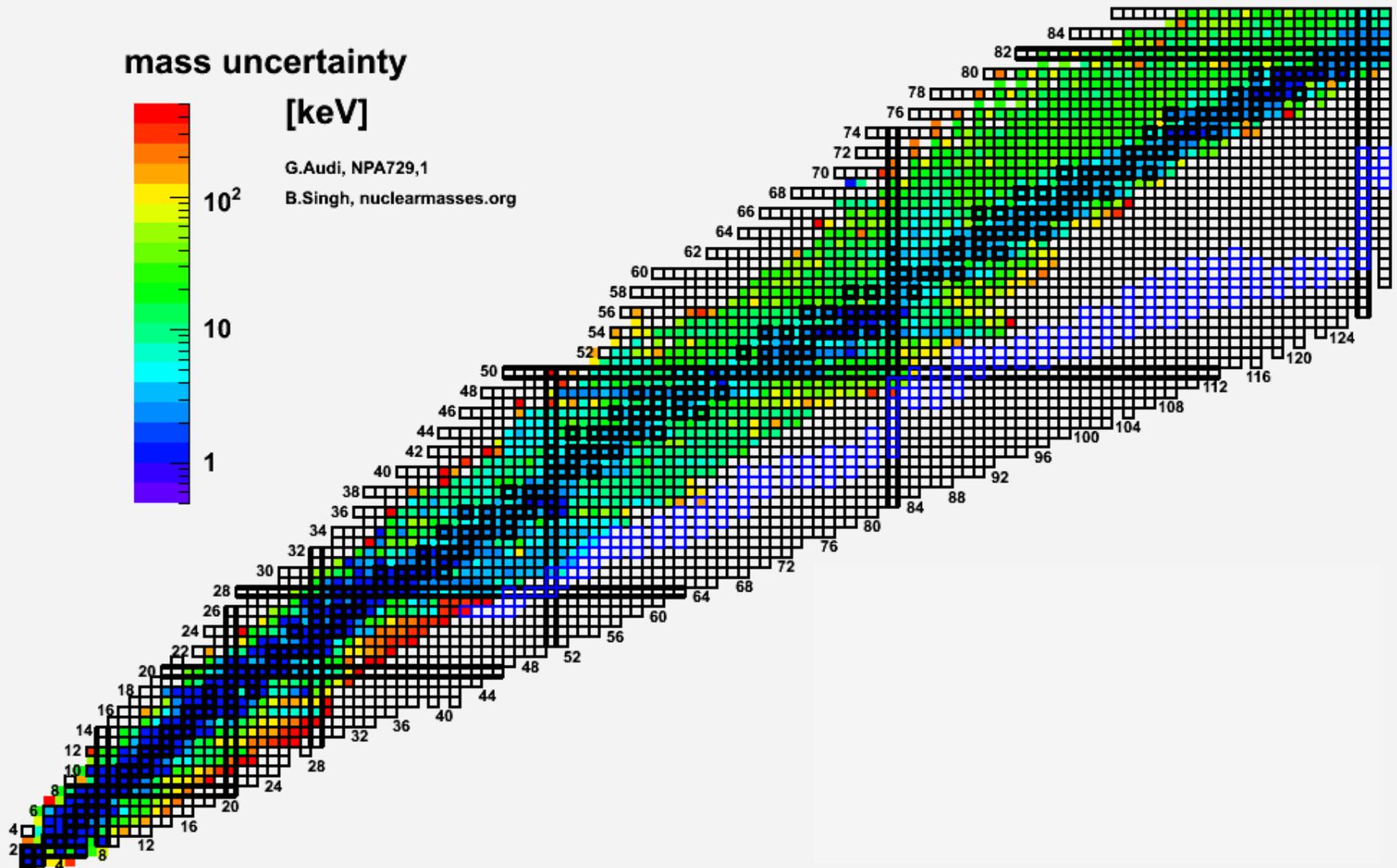


- Isotopes produced by fragmentation
- EM separated and implanted into Si detector stack (9)
- Identified by TOF and ΔE -E
- Decay β and γ measured
- Dozens of isotopes studied



S. Nishimura *et al.*, PRL **106**, 052502 (2011).

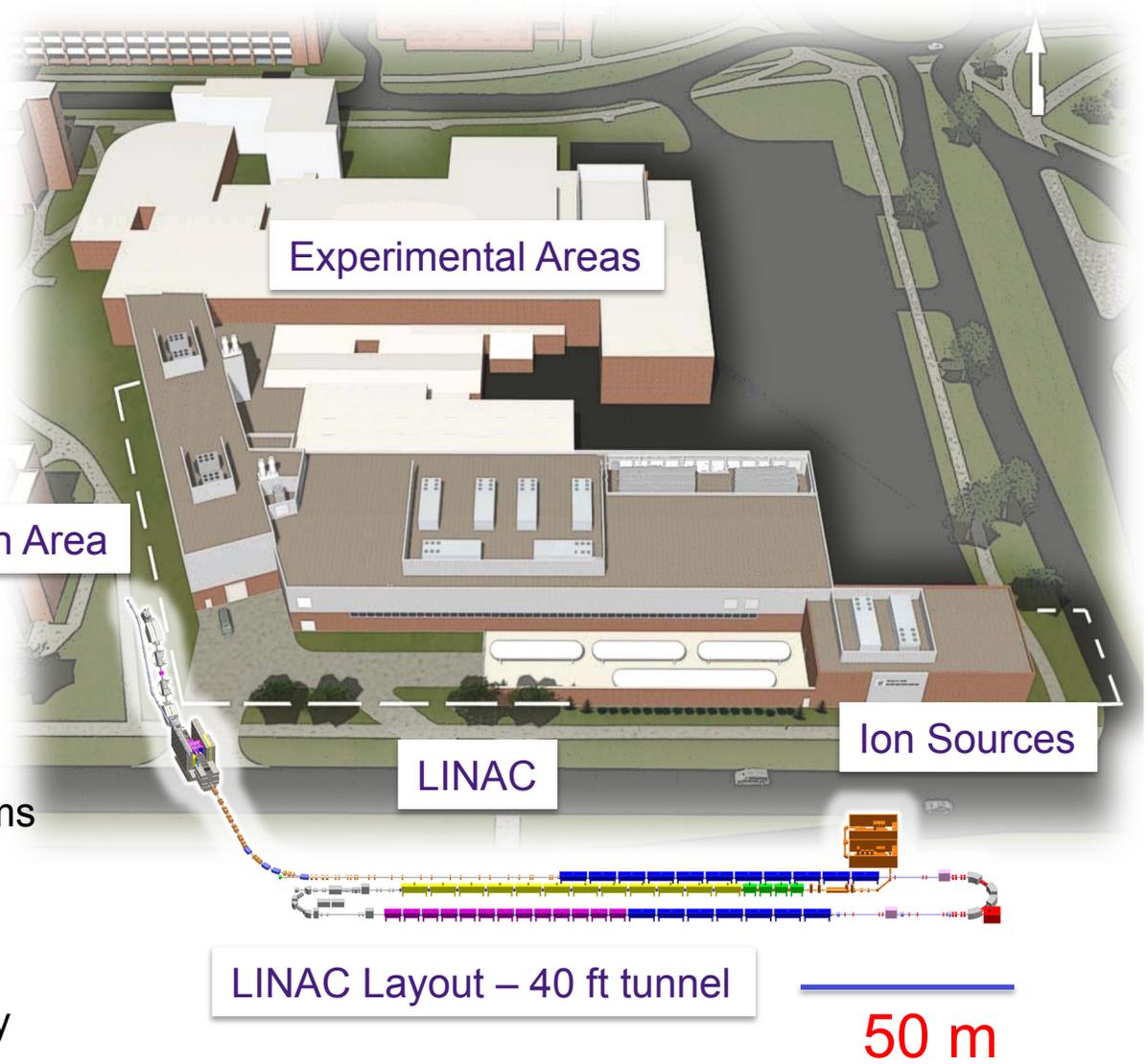
The r process





Facility for Rare Isotope Beams
U.S. Department of Energy Office of Science
Michigan State University

- On existing NSCL Site
- New gas stopping technology + post accelerator
- New Powerful driver LINAC
 - 200 MeV/u for U
 - 400 kW
- 9/01/10 CD-1
 - \$614M TPC
 - \$520 DOE
- CD-4 ~FY2020
- Wide **variety** of **intense** beams at **low energies**
- First access to many of the *r* process isotopes
- Direct measurements of many reactions for the first time





- Detailed study of all neutron-rich nuclei important for astrophysics possible up to Sn
- Masses and half-lives possible for most nuclei
 - Including $A \sim 190$ mass peak
 - Pin down r process site

