



Superallowed Nuclear β Decay

The Precision Frontier of Low-Energy Nuclear Physics

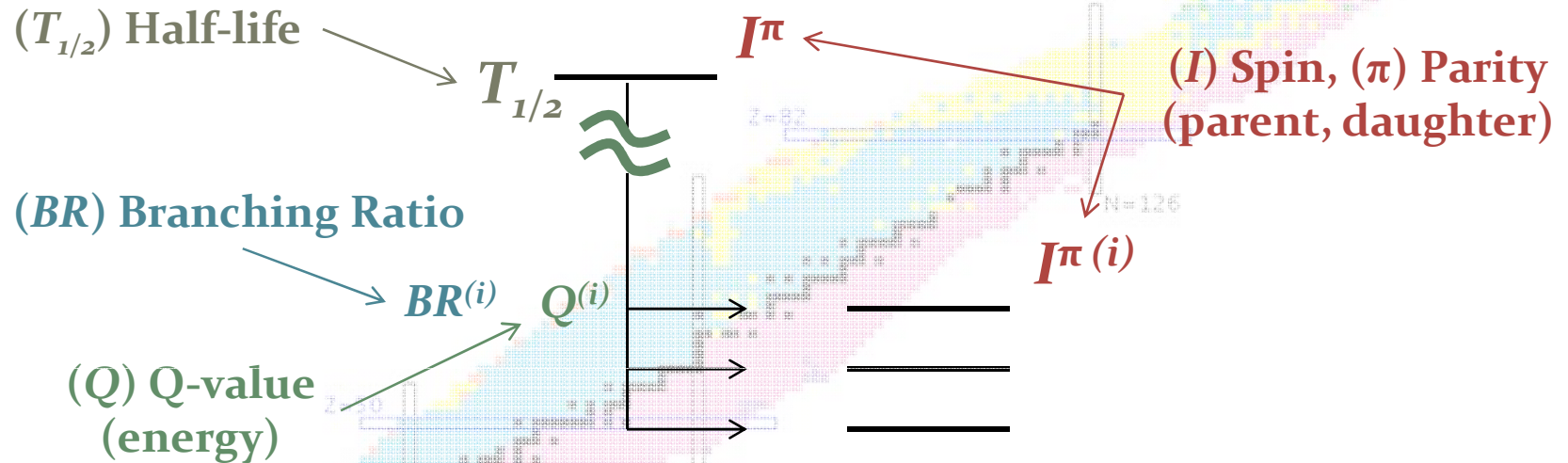
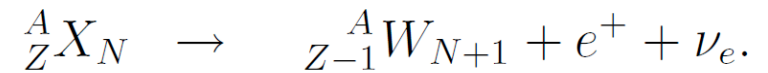
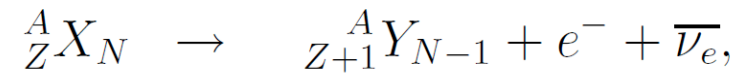
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NNPSS-TSI June 23 2010*

Outline

- Introduction and Review
 - Nuclear β decay
 - Isospin symmetry and isospin symmetry breaking
 - Conserved Vector Current hypothesis (CVC)
- Status of Superallowed Fermi β Decay
 - V_{ud} and CKM unitarity
- Case Study ^{62}Ga : “How to measure the ft value”
 - State-of-the-art facilities and experimental techniques
 - Comparison to theoretical isospin symmetry breaking corrections
- Summary and Conclusions

Nuclear β Decay



- Selection Rules:

$$\pi_P = \pi_D (-1)^L$$

$$\vec{I}_P = \vec{I}_D + \vec{L} + \vec{S}$$

- Allowed Decay

- L=0

- “Forbidden” Decay

- L=1, 2, 3, ...

- Fermi (Vector)

- S=0 $\beta \uparrow\downarrow \nu_e$

- Gamow-Teller (Axial-Vector)

- S=1 $\beta \uparrow\uparrow \nu_e$

β Decay Classification

B. Singh, J.L. Rodriguez, S.S.M.Wong, and J.K.Tuli, Nuclear Data Sheets 84, 487 (1998)

Case	$I^\pi (P \rightarrow D)$	Classification	$T_{1/2}$	Fraction
$^{18}\text{N} \rightarrow ^{18}\text{C}$	$1^- \rightarrow 1^-$	Allowed (GT&F)	624 ms	64%
$^6\text{He} \rightarrow ^6\text{Li}$	$0^+ \rightarrow 1^+$	Allowed (GT only)	807 ms	
$^{10}\text{C} \rightarrow ^{10}\text{B}$	$0^+ \rightarrow 0^+$	Allowed (F only)	19 s	1%
$^{38}\text{Cl} \rightarrow ^{38}\text{Ar}$	$2^- \rightarrow 2^+$	1 st Forbidden	37 min	33%
$^{36}\text{Cl} \rightarrow ^{36}\text{Ar}$	$2^+ \rightarrow 0^+$	2 nd Forbidden	3 10^5 years	1%
$^{40}\text{K} \rightarrow ^{40}\text{Ca}$	$4^- \rightarrow 0^+$	3 rd Forbidden	1 10^9 years	0.1%
$^{50}\text{V} \rightarrow ^{50}\text{Cr}$	$6^+ \rightarrow 2^+$	4 th Forbidden	1 10^{17} years	0.1%

- The ft value is a powerful way to characterize β decay

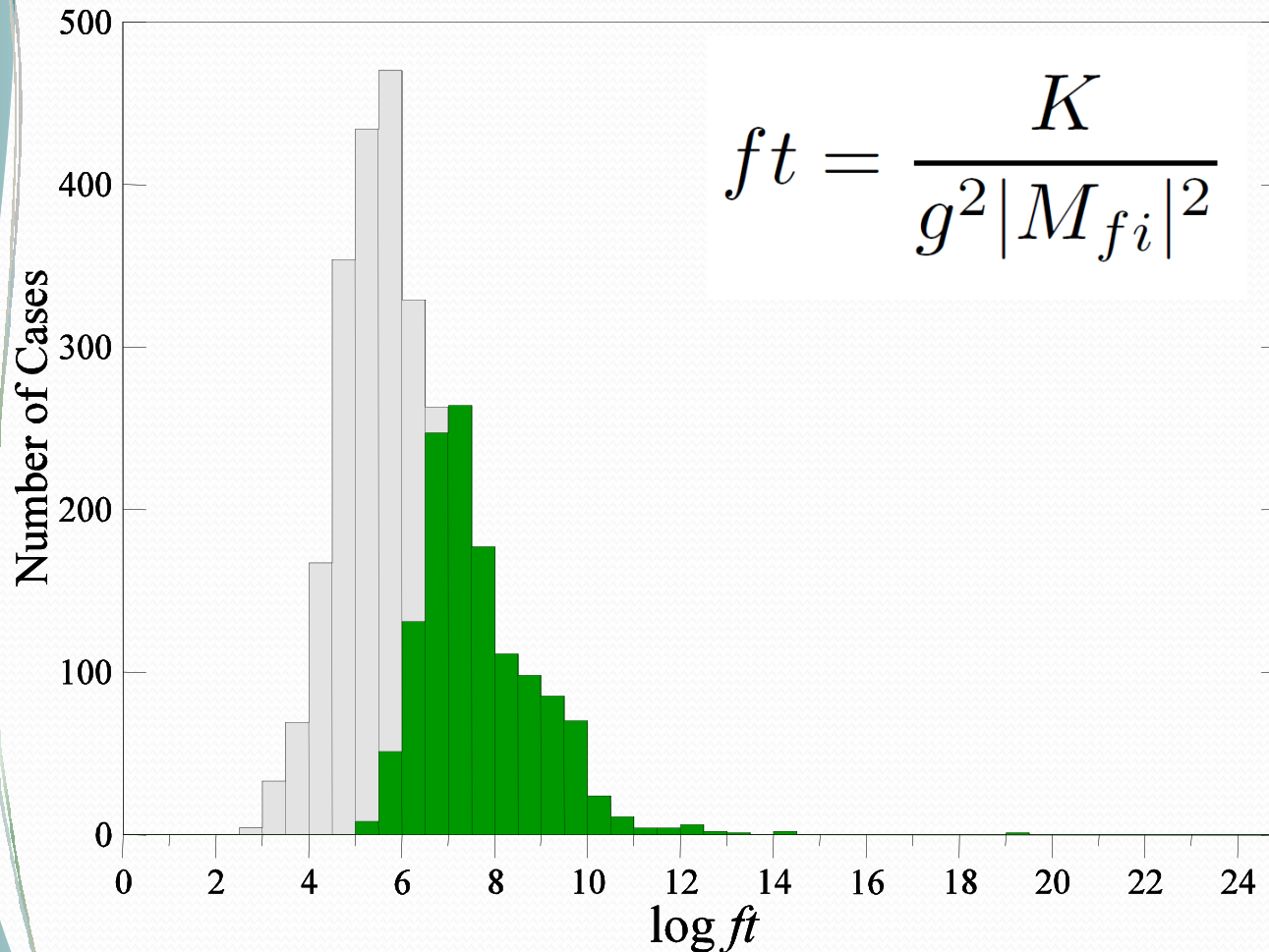
$$ft = \frac{f T_{1/2}}{BR} = \frac{K}{g^2 |M_{fi}|^2}$$

Q-value \swarrow Half-life \swarrow Constants \swarrow
 Branching Ratio \nearrow Matrix Element \swarrow
 Coupling Strength \swarrow

β decay ft values

B. Singh, J.L. Rodriguez, S.S.M.Wong, and J.K.Tuli, Nuclear Data Sheets 84, 487 (1998)

- Survey of 3840 β decay ft values



Allowed:

$1^+ \rightarrow 0^+, 2^- \rightarrow 2^-$
(2497)

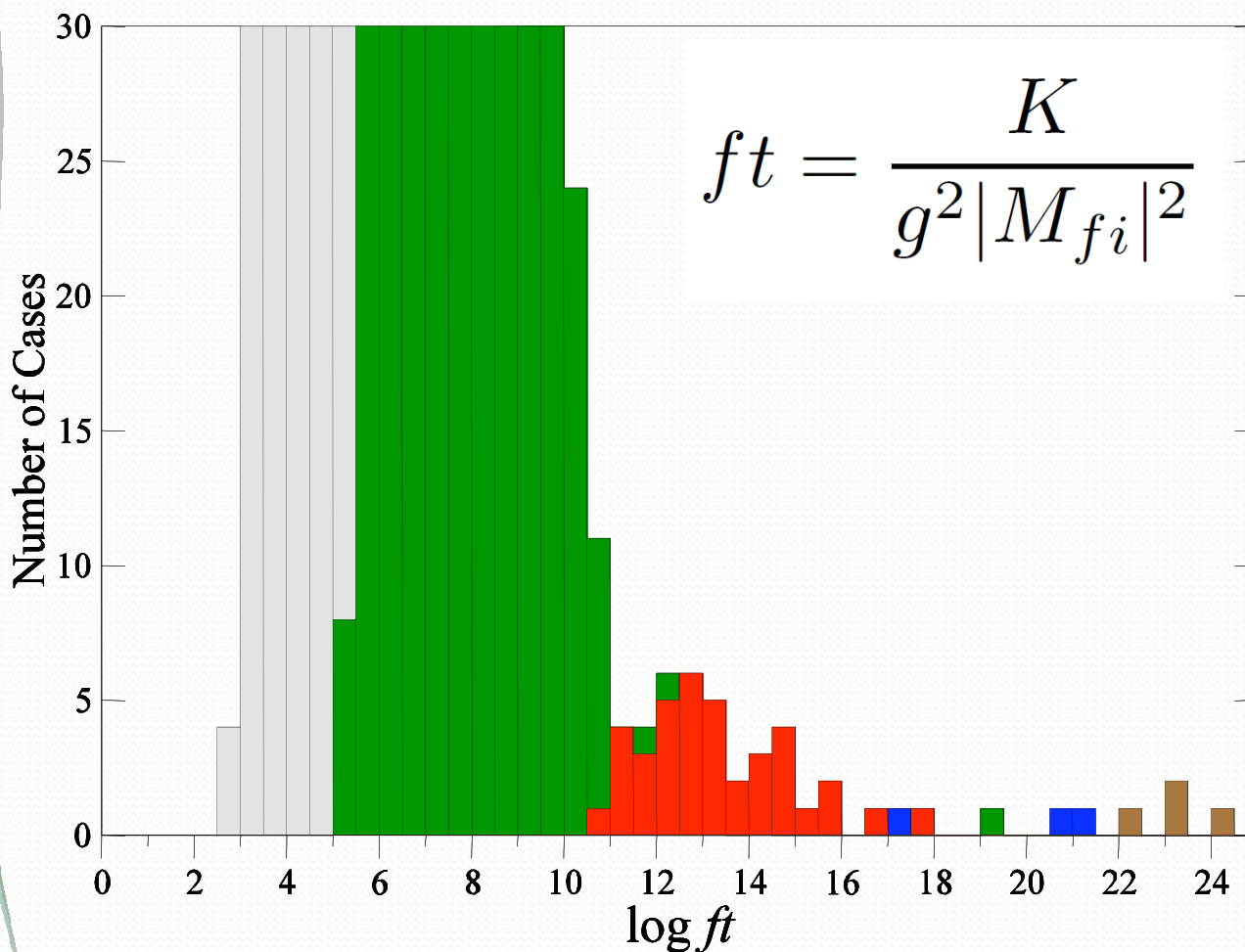
1st Forbidden:

$2^+ \rightarrow 0^+, 2^- \rightarrow 2^+$
(1297)

β decay ft values

B. Singh, J.L. Rodriguez, S.S.M.Wong, and J.K.Tuli, Nuclear Data Sheets 84, 487 (1998)

- β decay ft values span ~ 21 orders of magnitude



← 21 orders of magnitude! →

Allowed:

$1^+ \rightarrow 0^+, 2^- \rightarrow 2^-$
(2497)

1st Forbidden:

$2^+ \rightarrow 0^+, 2^- \rightarrow 2^+$
(1297)

2nd Forbidden:

$2^+ \rightarrow 0^+, 5^+ \rightarrow 2^+$
(39)

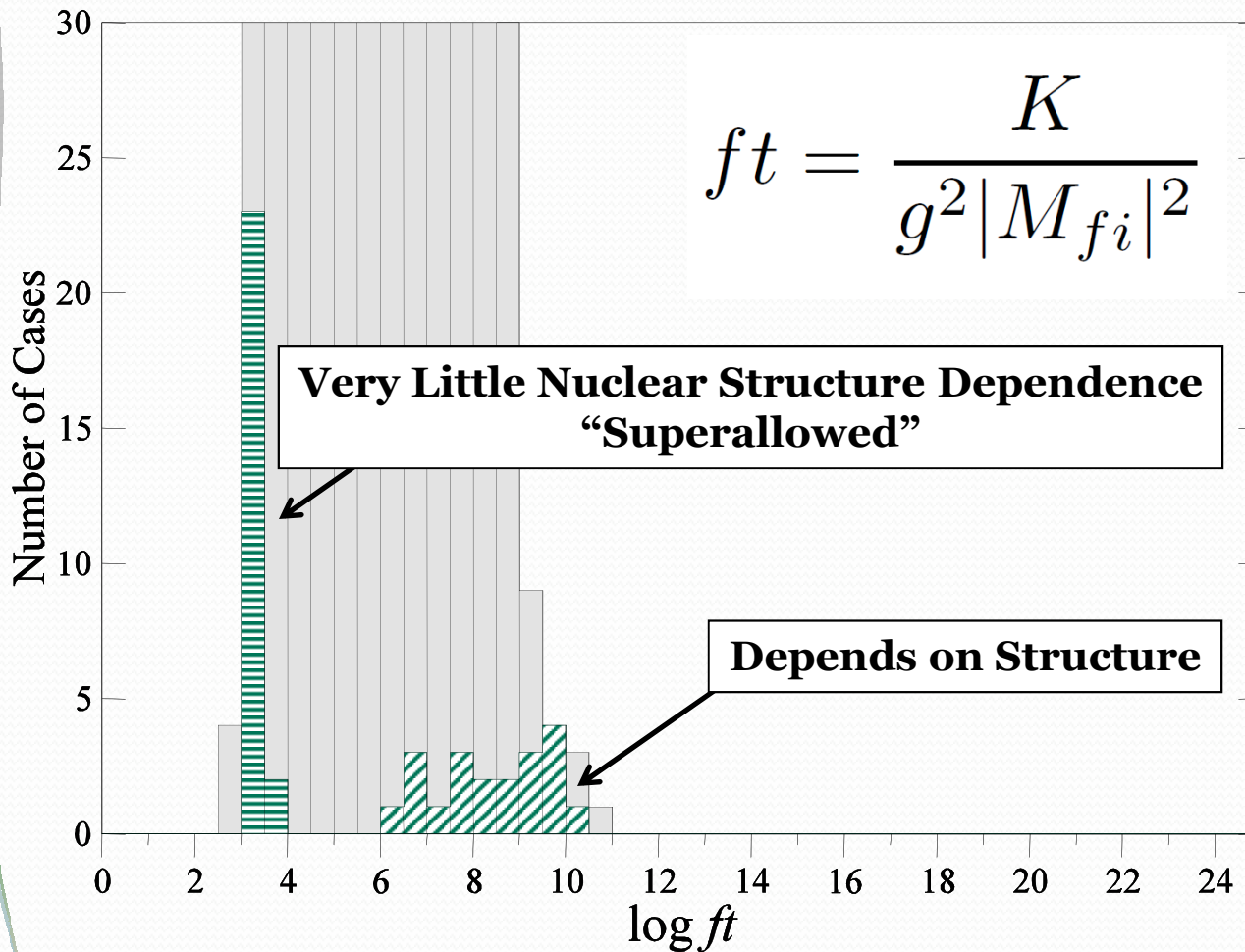
3rd & 4th Forbidden

$4^- \rightarrow 0^+, 6^+ \rightarrow 2^+$
(3), (4)

$0^+ \rightarrow 0^+$ Pure Fermi Decay

- $0^+ \rightarrow 0^+$ is *pure* Fermi decay (no GT)

$$\vec{I}_P = \vec{I}_D + \vec{L} + \vec{S}$$



Allowed:
 $1^+ \rightarrow 0^+, 2^- \rightarrow 2^-$
 (2497)

Allowed (pure F):
 $0^+ \rightarrow 0^+$
 (25 + 20)

Nuclear Isospin

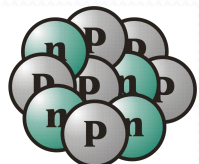
- In 1932, months after the discovery of the neutron Heisenberg proposes the concept of “nucleons”
 - n’s and p’s are “spin” projections of nucleons

$$\textcircled{\mathbf{p}} \quad t_z(p) = -\frac{1}{2} \qquad \textcircled{\mathbf{n}} \quad t_z(n) = +\frac{1}{2}$$



1932 - Heisenberg

- Isospin projection and total isospin of a nucleus



$$T_z = \frac{1}{2}(N - Z) \qquad \mathbf{T} = |T_z|, |T_z| + 1, \dots, \frac{N + Z}{2}$$

- For $0^+ \rightarrow 0^+$ β decay ($n \leftrightarrow p$) between *isobaric analogue states* (same T , different T_z) the M.E. is an isospin ladder operator

$$|M_F|^2 = (T \mp T_z)(T \pm T_z + 1)$$

For $T = 1$ decays $|M_F|^2 = 2$

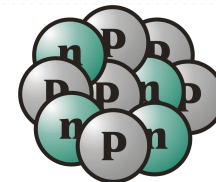
EXACT!
(to the extent that
isospin is a valid
symmetry) 8

CVC Hypothesis

R.P. Feynman and M. Gell-Mann Phys. Rev. 109 193 (1958)

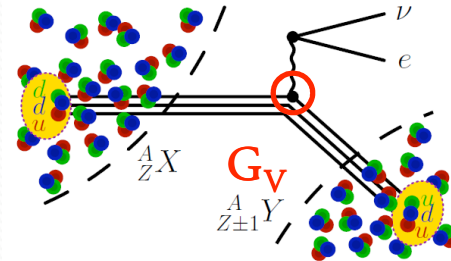
- Electric charge of the proton is not altered inside a nucleus
 - E&M is governed by a conserved vector current (CVC)

$$e = 1.602\,176\,487 \times 10^{-19} \text{ C}$$



- Hypothesize that a CVC also exists for weak interaction
 - Does the weak interaction have a universal coupling strength?

$$G_V = 1.13621 \times 10^{-5} \text{ GeV}^{-2}$$



- Consequence of CVC:
 - The ft values for ALL superallowed Fermi decays are identical!

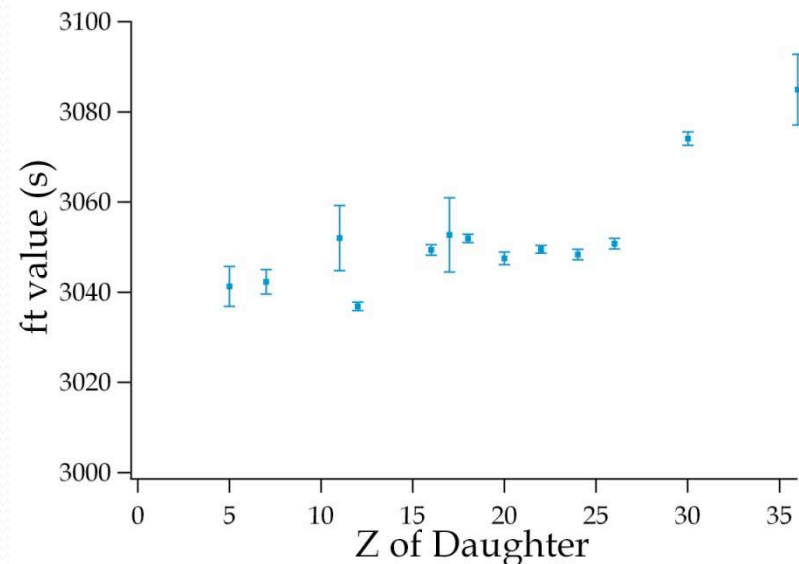
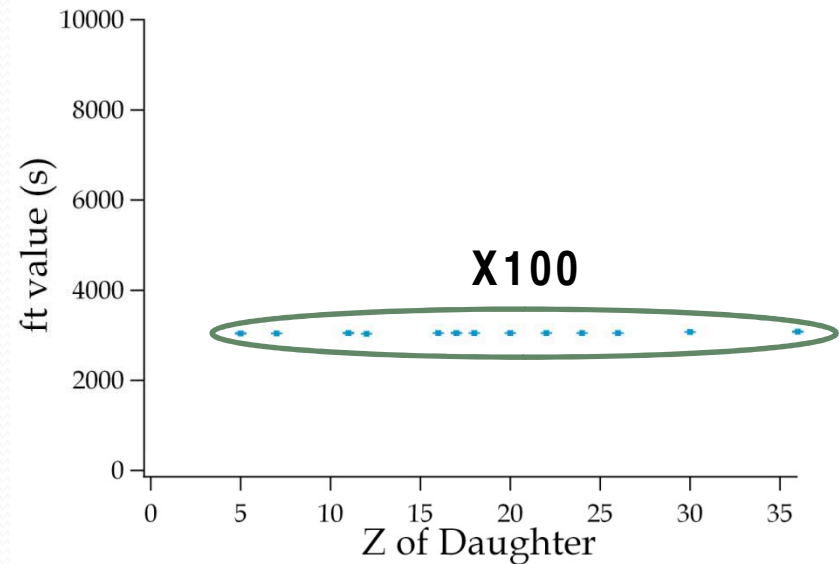
$$ft = \frac{fT_{1/2}}{BR} = \frac{K}{2G_V^2} = \text{constant}$$

↑ Isospin symmetry ↑ CVC

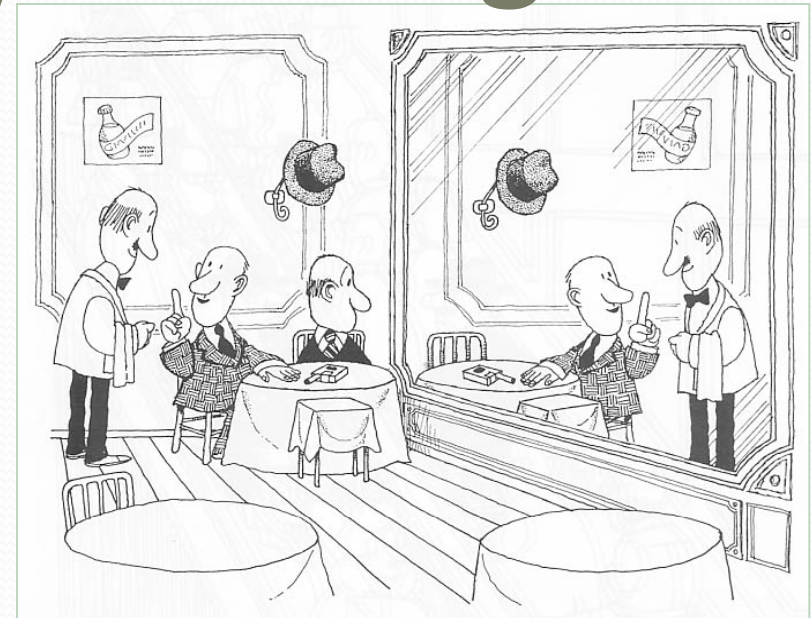
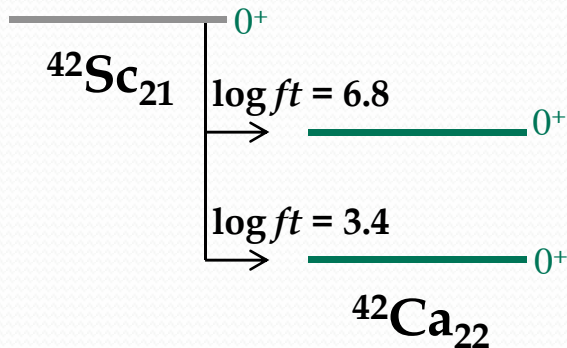
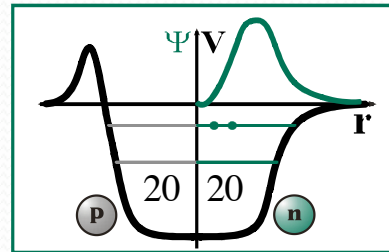
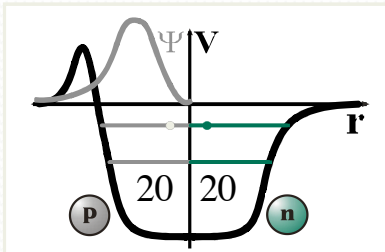
Superaligned Survey (2009)

J.C. Hardy and I.S. Towner Phys. Rev. C 79 055502 (2009)

- Survey of world data
 - 130 independent measurements
 - 13 ft values known to $\leq 0.4\%$
 - 8 ft values known to $\leq 0.05\%$
- All known ft values
 - Span $ft = 10^3$ to 10^{24} s
- All allowed ft values
 - Span $ft = 10^3$ to 10^{10}
- Superaligned ft values
 - Span $ft = 3040$ to 3100 s (2%)
 - Assumes isospin symmetry



Isospin Symmetry Breaking

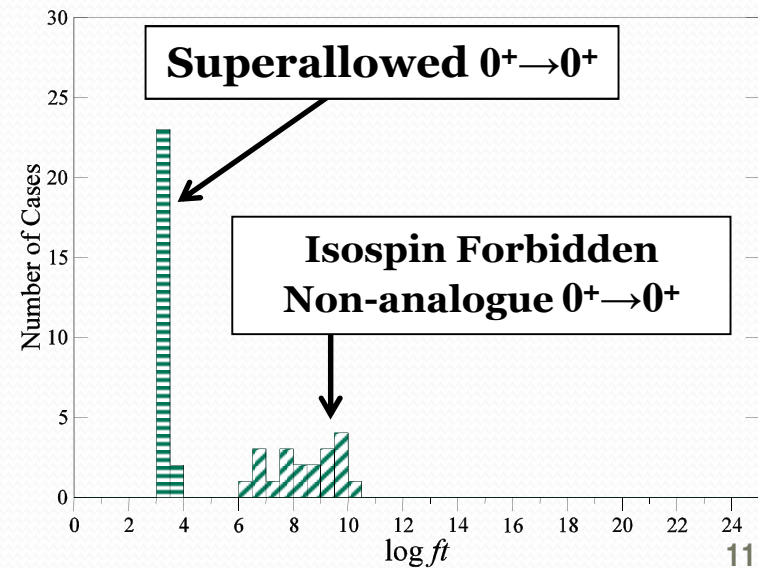


- If isospin symmetry were *exact*

- Only the S.A.F. decay would occur
- Non-analogue decays have been observed
 - Suppressed by factors of > 1000

$$|M_F^{ias}|^2 = 2(1 - \delta_C)$$

$$|M_F^{nas}|^2 = 2\delta_C$$

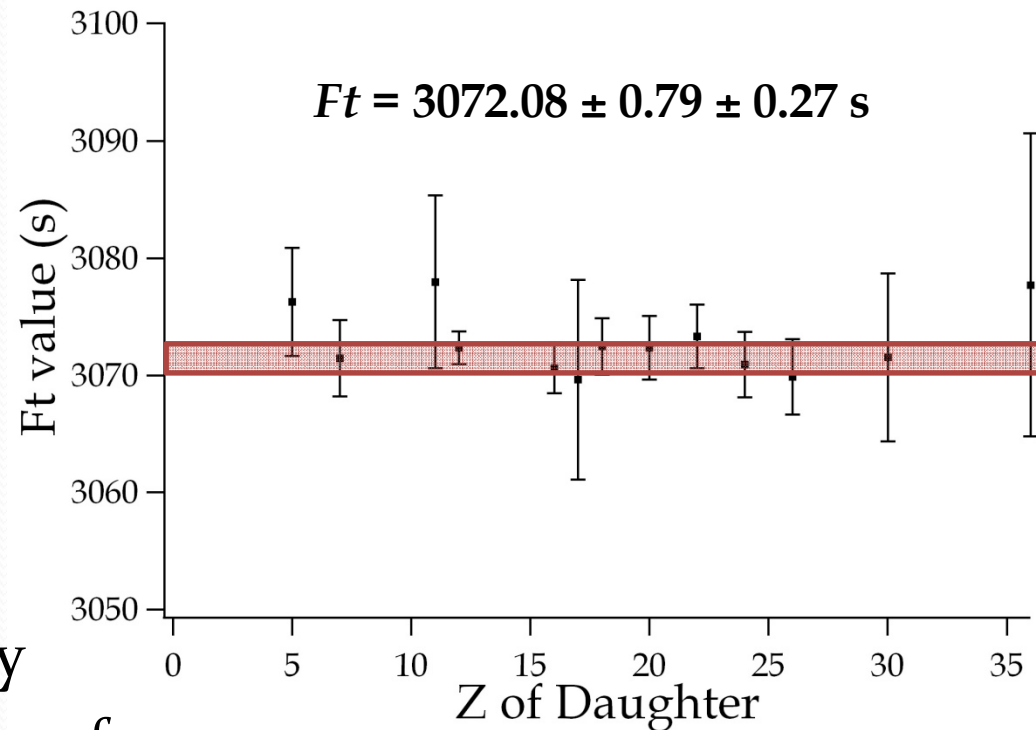


World Status (2009)

J.C. Hardy and I.S. Towner Phys. Rev. C 79 055502 (2009)

- Apply δ_C, δ_R corrections
- World average Ft
 - $Ft = 3072.08(79) \text{ s}$
 - $\chi^2/\nu = 0.28$
- CVC hypothesis
 - Validated (10^{-4} level)
- Systematic uncertainty
 - $\pm 0.27 \text{ s}$ from comparison of two sets of δ_C calculations

$$\mathcal{F}t = ft(1 + \delta_R)(1 - \delta_C)$$

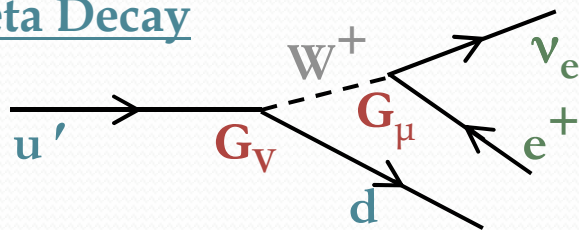


Evaluation of G_V and V_{ud}

I.S. Towner and J.C. Hardy Rep. Prog. Phys. 73 046301 (2010)

- With CVC satisfied from 13 superallowed decays, derive G_V

Beta Decay

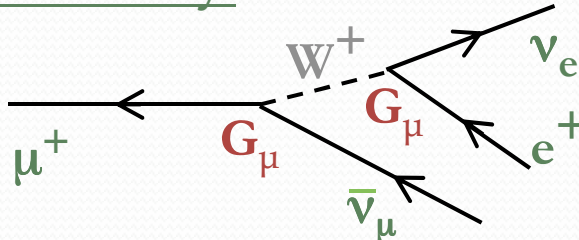


$$\mathcal{F}t = \frac{K}{2G_V^2} = \text{constant}$$

$$G_V = 1.13621 \times 10^{-5} \text{ GeV}^{-2}$$

- Compare to the corresponding value G_μ from μ decay

Muon Decay



$$G_\mu = 1.16637 \times 10^{-5} \text{ GeV}^{-2}$$

$$V_{ud} = G_V/G_\mu = 0.97425(22)$$

Quark Mixing

- 1958 R.P. Feynman and M. Gell-Mann Phys. Rev. 109 193 (1958)
 - CVC Hypothesis
- 1960 R.K. Bardin et al. Phys. Rev. Lett. 5 323 (1960)
 - ft value for ^{14}O : $3069(13) \text{ s}$ [0.4% !]
 - $G_V = 0.982G_\mu!$ Expected $G_V = 1.0G_\mu$
- 1963 N.Cabibbo Phys. Rev. Lett. 10 531 (1963)
 - Weak eigenstates \neq mass eigenstates
 - Included “mixing” between 2 generations
 - Expressed as a unitary rotation
 - Cabibbo angle $\theta_C = 11^\circ$ explains ^{14}O decay
- 1973 M.Kobayashi and T.Maskawa Prog. Theor. Phys. 49 652 (1973)
 - Require a third quark generation to explain CP violation
 - Cabibbo’s 2x2 matrix is expanded to 3x3

Leptons	Quarks	u up	c charm	t top
		d down	s strange	b bottom
		ν_e e- Neutrino	ν_μ μ - Neutrino	ν_τ τ - Neutrino
	e electron	μ muon	τ tau	
	I	II	III	
	The Generations of Matter			

$$\begin{pmatrix} d' \\ s' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} \\ V_{cd} & V_{cs} \end{pmatrix} \begin{pmatrix} d \\ s \end{pmatrix}$$

$$\begin{aligned} V_{ud} &= \cos \theta_C & V_{us} &= \sin \theta_C \\ V_{cd} &= -\sin \theta_C & V_{cs} &= \cos \theta_C \end{aligned}$$

$$|V_{ud}|^2 + |V_{us}|^2 = 1$$



2008 Nobel Prize in Physics
 1/4 Kobayashi & 1/4 Maskawa

CKM Quark Mixing Matrix

I.S. Towner and J.C. Hardy Rep. Prog. Phys. 73 046301 (2010)

- Three generation quark mixing matrix

Leptons	Quarks		
	u up	c charm	t top
	d down	s strange	b bottom
	ν_e e- Neutrino	ν_μ μ - Neutrino	ν_τ τ - Neutrino
	e electron	μ muon	τ tau
	I	II	III
The Generations of Matter			

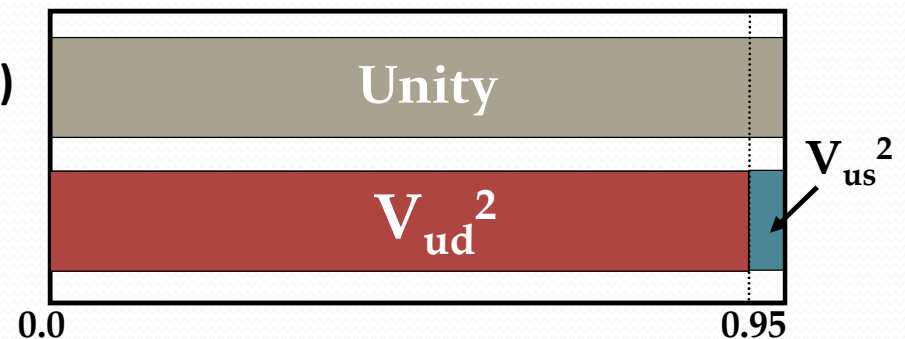
$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$\begin{aligned} V_{ud} &= 0.97 & V_{us} &= 0.22 & V_{ub} &< 0.01 \\ V_{cd} &= 0.23 & V_{cs} &= 1.04 & V_{cb} &< 0.01 \\ V_{td} &< 0.01 & V_{ts} &= 0.04 & V_{tb} &> 0.74 \end{aligned}$$

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1 \quad \text{Unitarity Condition}$$

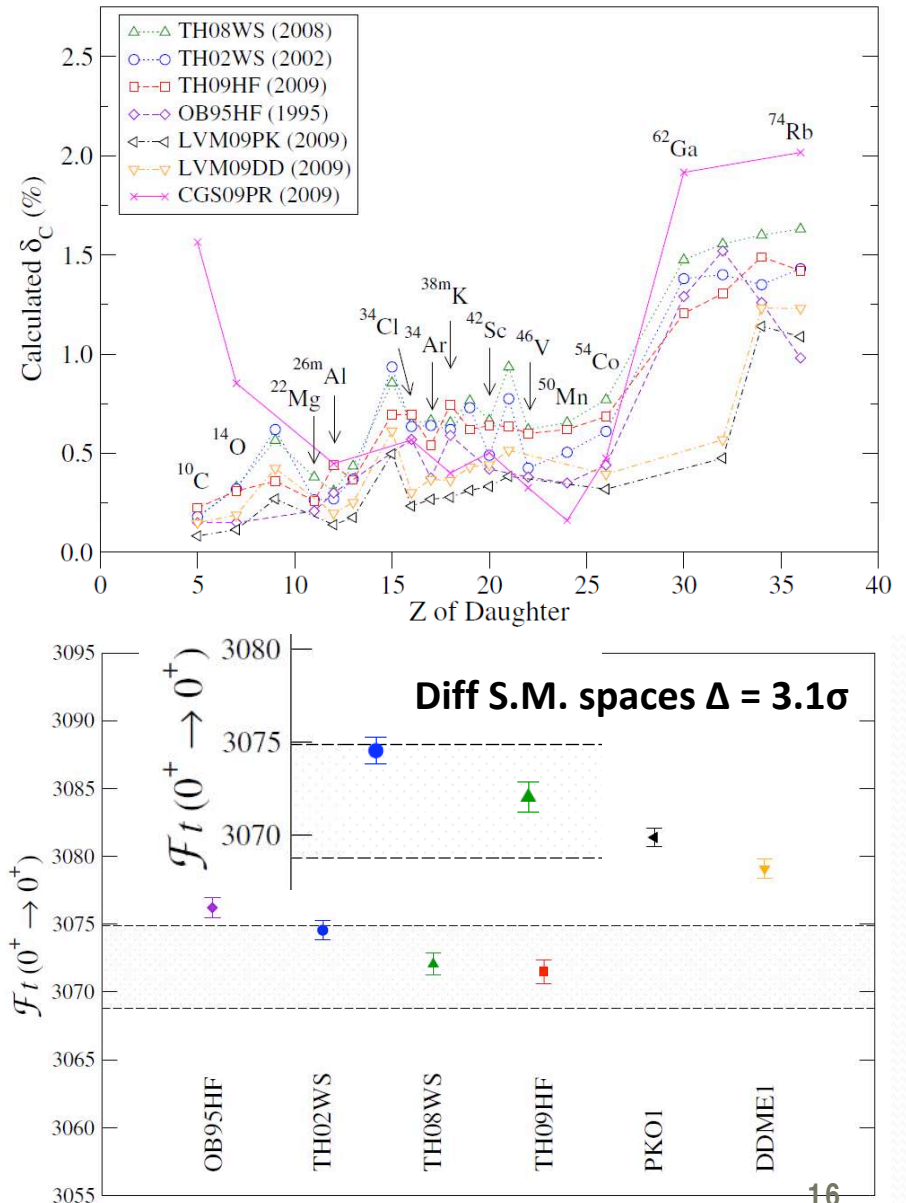
- Most precise experimental test of unitarity comes from the top row
 - Present Status*: $V_{ud} = 0.97425(22)$, $V_{us} = 0.2246(12)$, $V_{ub} = 0.0039(4)$

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 0.9999(6)$$



Recent I.S.B. Calculations

- Several “sets” of δ_C values
 - Different approaches
 - Shell Models
 - Relativistic HF + RPA
 - Collective
 - Results differ by factors of 2
- Impact on V_{ud} and CKM:
 - All use same ft , δ_R values
- Strategy:
 - Constrain via experiment
 - Nuclei that exhibit largest theoretical differences





Experimental Results
Superaligned Decay of ^{62}Ga

The case of ^{62}Ga

J.C. Hardy and I.S.Towner EPJA. 25 695 (2005)

- Status of the ^{62}Ga ft value in 2005...

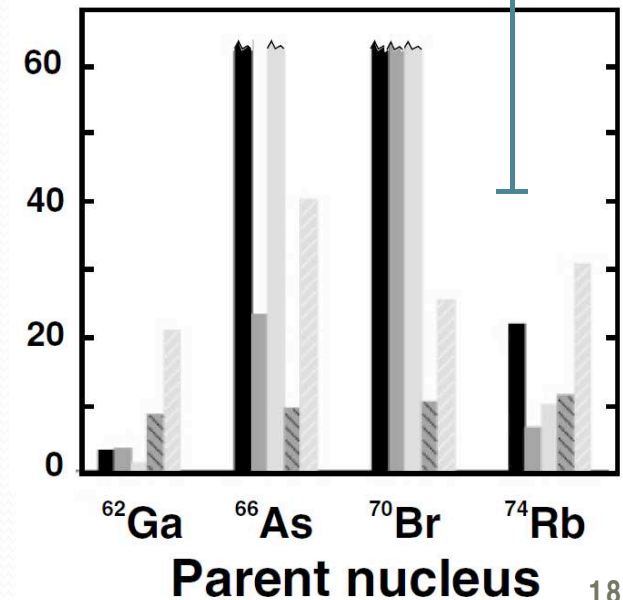
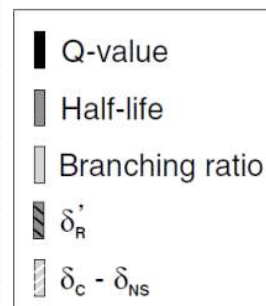
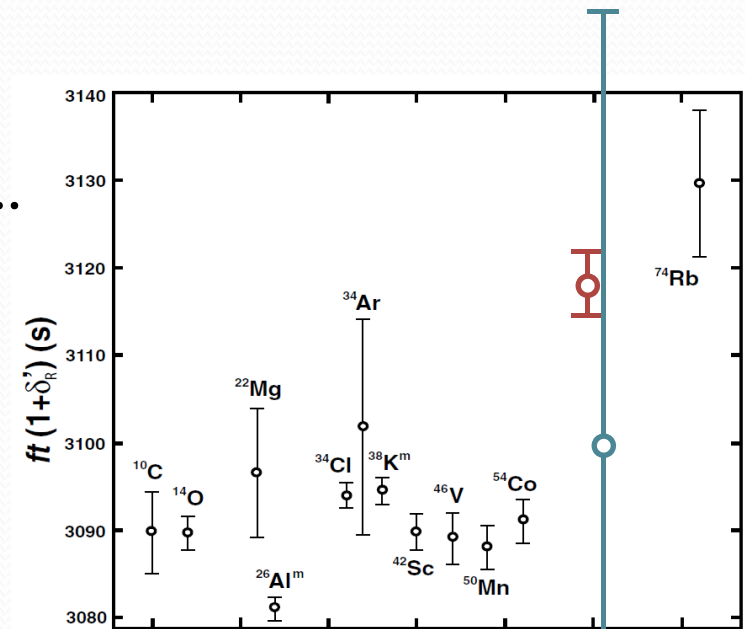
- $ft(1+\delta_R) = 3102(48) \text{ s}$

- Since: all 3 quantities measured

- Q-value (0.004%) Jyväskylä
 - $T_{1/2}$ (0.022%) TRIUMF
 - B.R. (0.001%) TRIUMF

- Present value

- $ft(1+\delta_R) = 3119(3) \text{ s}$
 - 16x improvement
 - Ft value now limited by δ_C



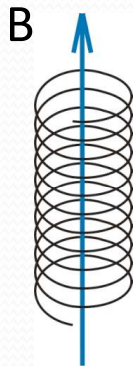
Penning Trap Basics

Figures courtesy of A.Kwiatkowski NSCL

- Ions with charge q mass m in a mag. field B
 - Circular motion with radius r and velocity v

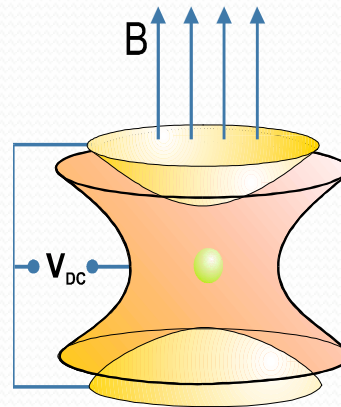
$$\mathbf{F} = q\mathbf{v} \times \mathbf{B} = \frac{mv^2}{r}$$

Magnetic field

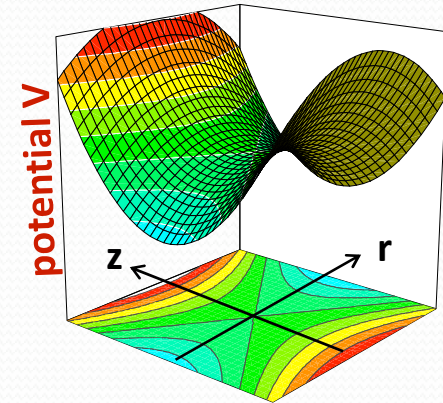


2D Radial confinement

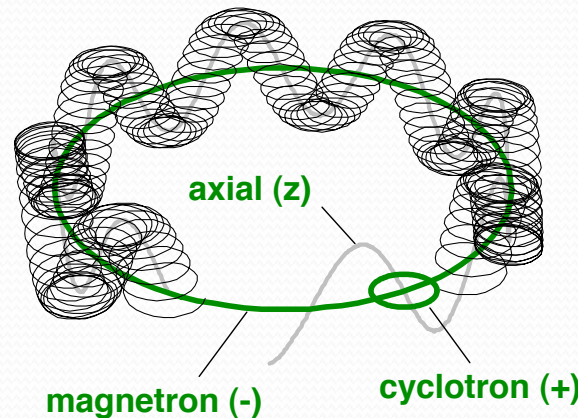
Electric quadrupole field



1D Axial confinement



3D confinement



3 Independent motions

Axial (z), Magnetron (-), reduced cyclotron (+)

$$2\pi\nu_c = \frac{v}{r} = \frac{qB}{m}$$

Mass of ^{62}Ga

T.Eronen et al. Phys. Lett. B 636 191 (2006)

- Beams of ^{62}Ga produced at IGISOL
 - Trap ^{62}Ga and scan rf frequency
 - Release ^{62}Ga from the trap
 - Measure time-of-flight (Trap/Detector)
 - Shortest TOF = resonance ν_c frequency
 - ^{62}Ga frequency relative to ^{62}Zn , ^{62}Cu , ^{62}Ni

$$r = \frac{\nu_{c,\text{ref}}}{\nu_c} \quad m = r(m_{\text{ref}} - m_e) + m_e$$

$$Q_{\text{EC}} = m_{\text{mother}} - m_{\text{daughter}} = (r - 1)(m_{\text{ref}} - m_e)$$



JYFLTRAP
Univ. Jyväskylä, Finland

R.I.B production
30 keV

Mass separation

Deflector

R.F.Q.
Cooler/buncher

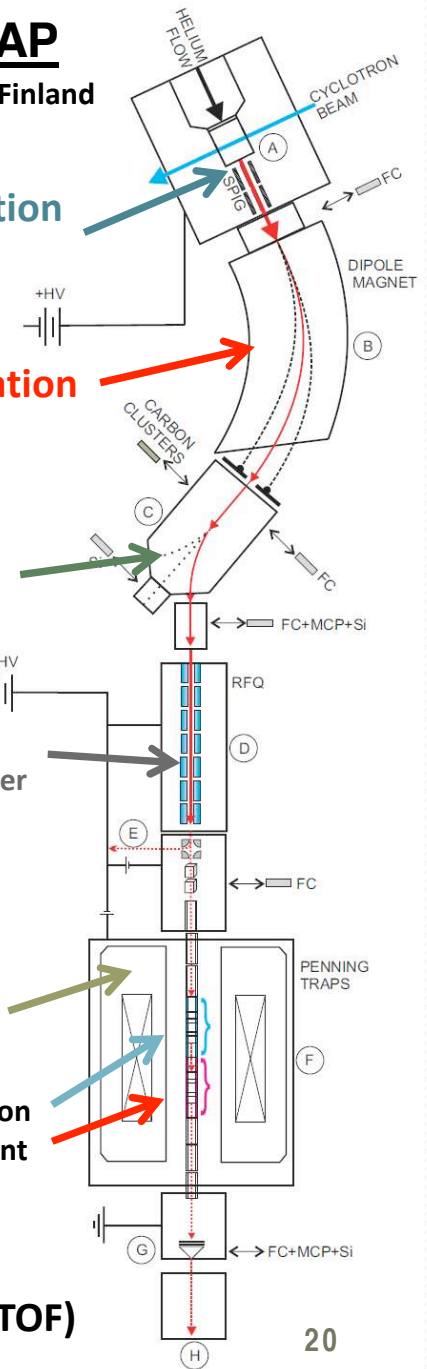
7T S.C. solenoid

Penning Traps

Trap 1 – Isobaric purification

Trap 2 – mass measurement

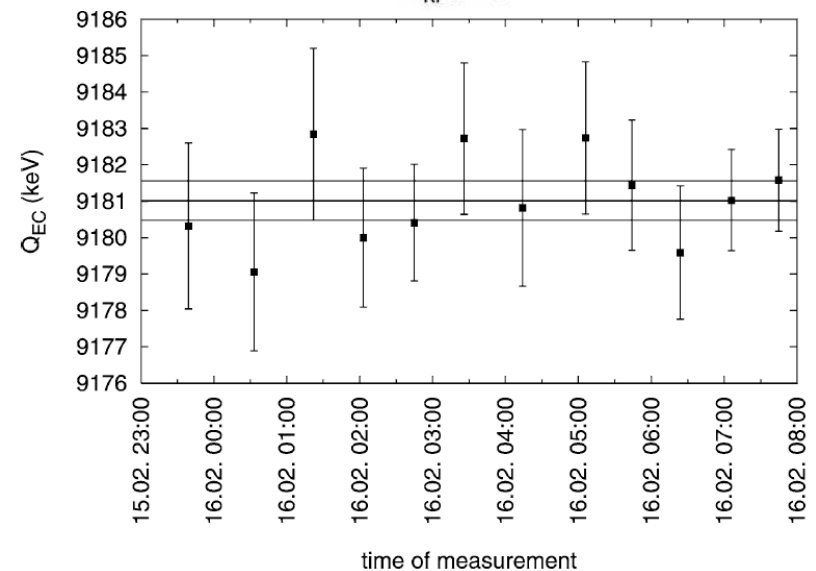
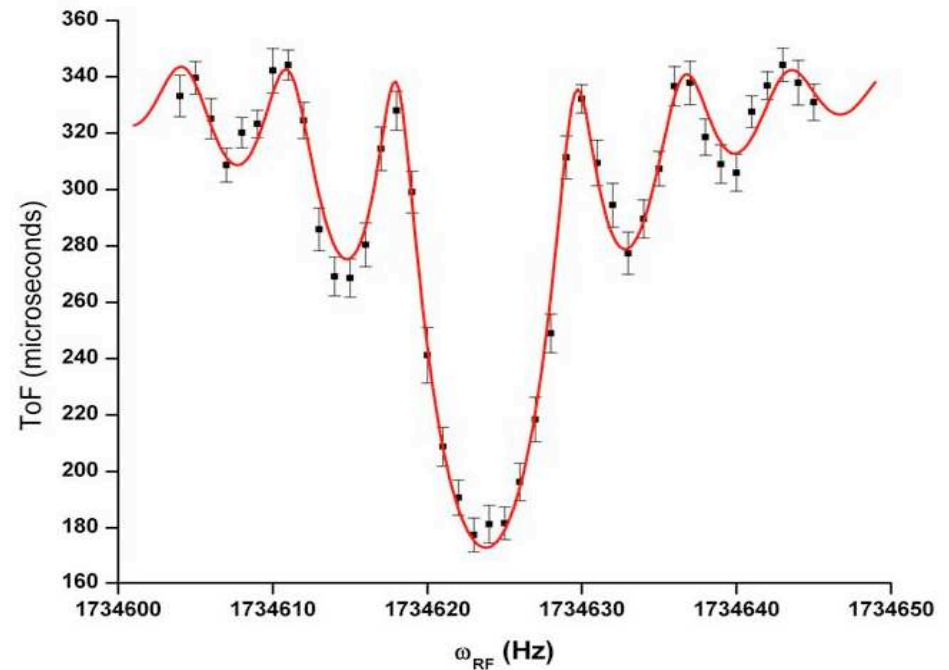
Detector Station (TOF)



Mass of ^{62}Ga

T.Eronen et al. Phys. Lett. B 636 191 (2006)

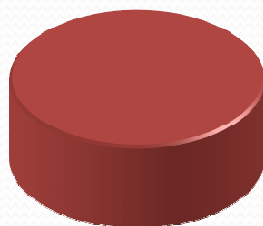
- Final result for ^{62}Ga :
 - Q-value = 9181.07(38) keV
 - f -value = 26401.6(83)
- Overall precision:
 - Q-value = 0.004%
 - f -value = 0.03%
- Previous values
 - Q-value = 9171(26) keV (\downarrow 70x)
 - f -value = 26250(400) (\downarrow 50x)



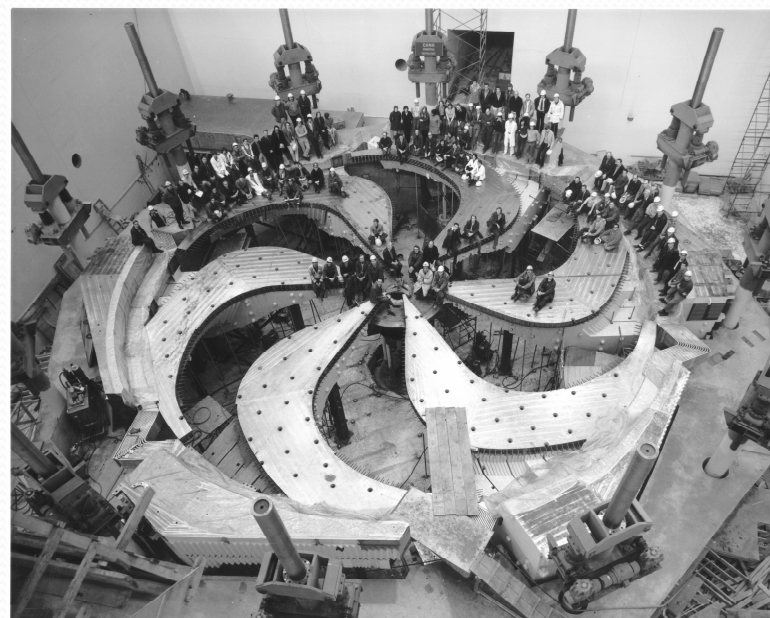
^{62}Ga $T_{1/2}$ and Branching Ratio

- Performed at TRIUMF-ISAC:

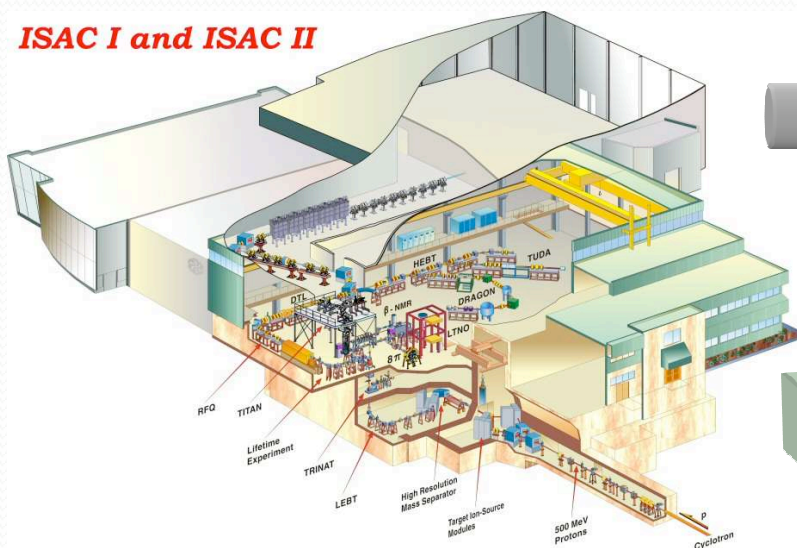
TRIUMF 500 MeV
Cyclotron 100 μA



proton beam



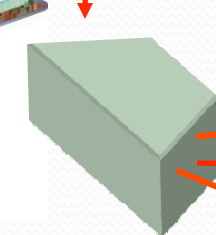
ISAC I and ISAC II



Production Target



Ion Source



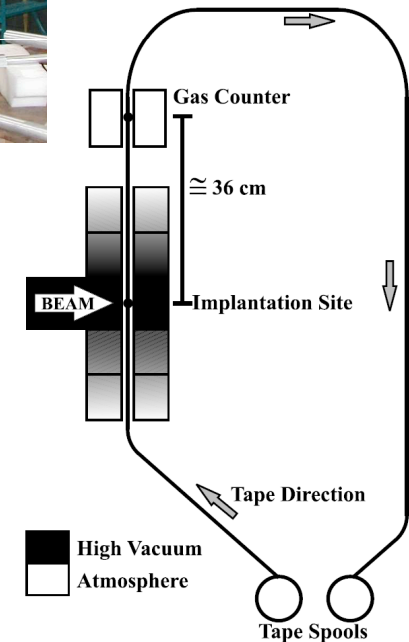
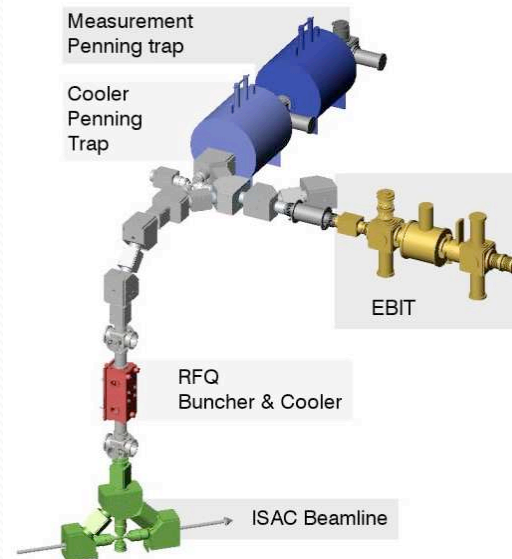
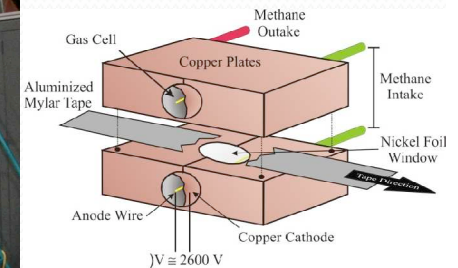
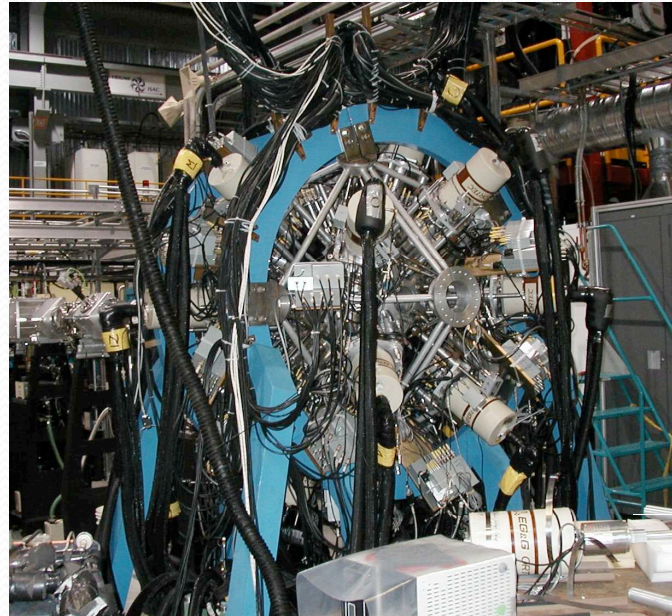
Mass Separator

60 keV ion beam

Experiments

Experiments at TRIUMF

- 8π γ -ray spectrometer
 - 20 HPGe detectors
 - β 's, β delayed γ -rays
 - β branching ratios
- Counter and fast tape
 - β particle counting
 - Half-life measurements
- Penning Trap (TITAN)
 - High-precision masses



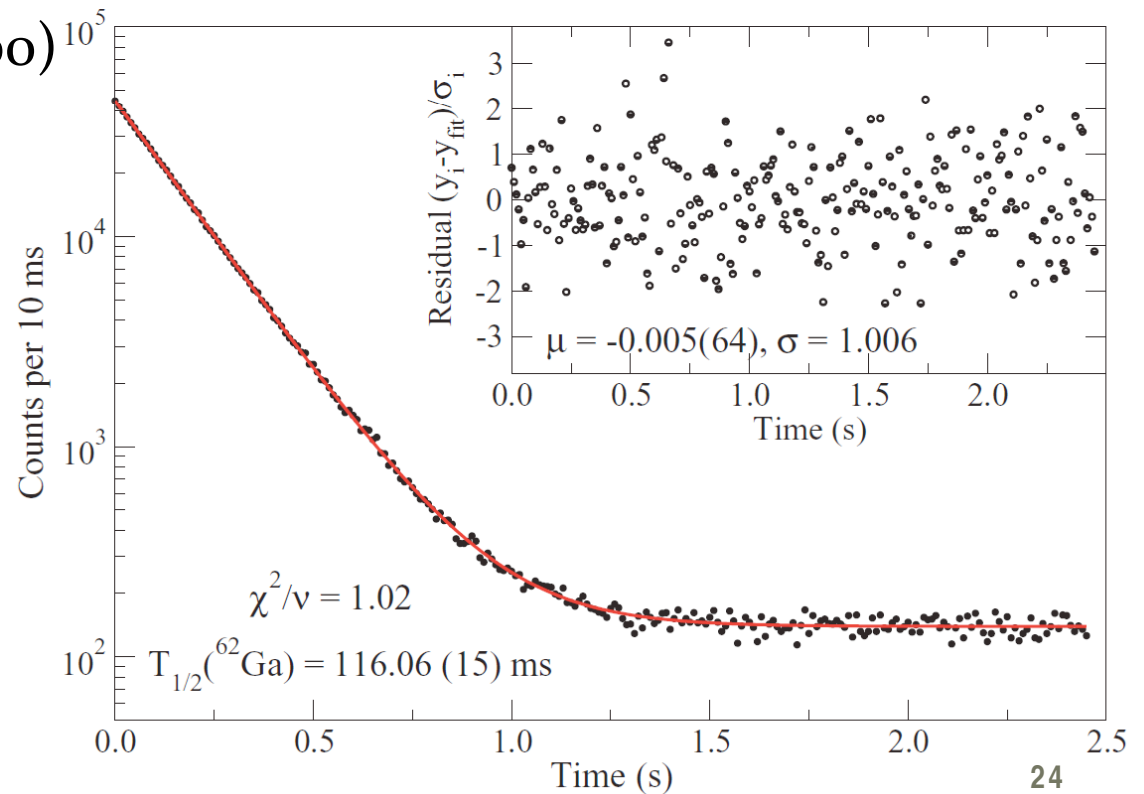
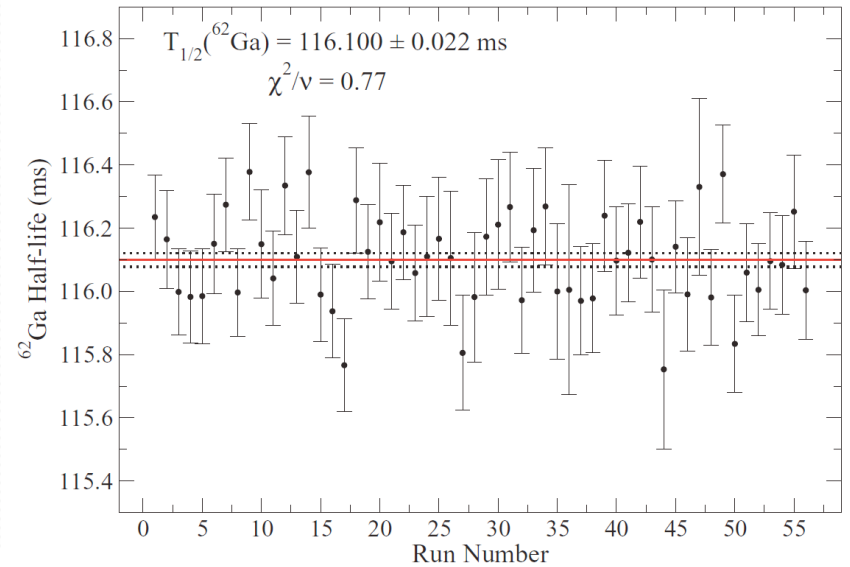
Half-life of ^{62}Ga

G.F. Grinyer et al. Phys. Rev. C 77 015501 (2008)

- ^{62}Ga beam of 8000 ions/s
 - Beam on 0.5 s, move tape
 - Count decay 2.5 s
 - Repeat 1200 times/hour
 - Run 65 hours (80,000)

- Half-life

- Deduced from $Ae^{-\lambda t}$
- 80,000 fits!
- Sum & fit = 56 fits
- Average = $T_{1/2}$



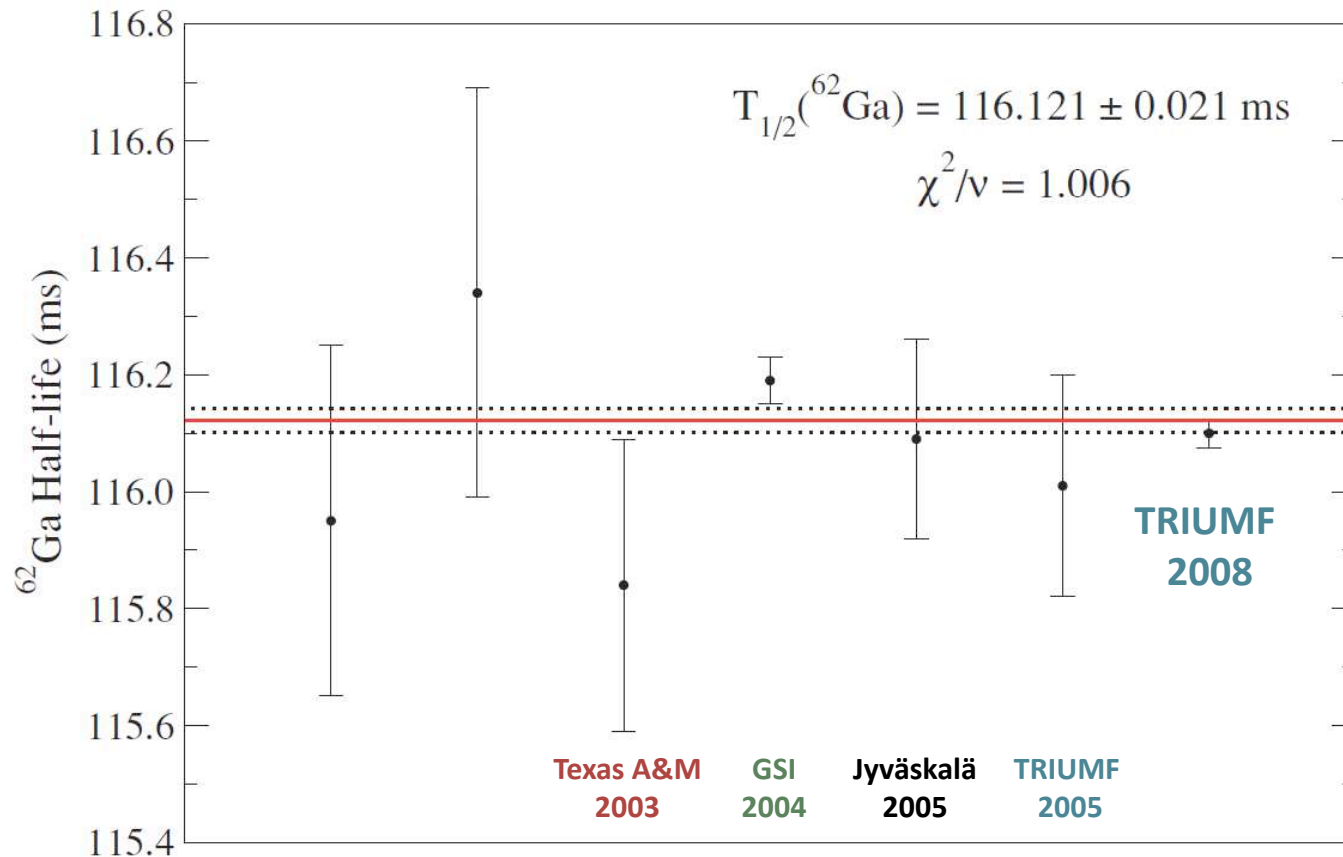
Half-life of ^{62}Ga

G.F. Grinyer et al. Phys. Rev. C 77 015501 (2008)

Most precise $T_{1/2}$ measurement
for any superallowed decay!

- ^{62}Ga Half-life

- Average TRIUMF result (2008): 116.100 ± 0.025 s (0.02%)
- Previous world average (2005): 116.175 ± 0.038 s (0.04%)



New Result: Half-life of ^{26m}Al

- ^{26m}Al half-life
 - New result at TRIUMF: Half-life measured to 0.01%!
 - See Poster by Paul Finlay (University of Guelph)

Ultra-high precision half-life measurement for the superallowed β^+ emitter $^{26}\text{Al}^m$

UNIVERSITY OF GUELPH

The Cabibbo-Kobayashi-Maskawa (CKM) quark mixing matrix

The CKM matrix plays a central role in the Standard Model by describing the mixing of the different quark generations.

$$|V\rangle = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} |u\rangle \\ |s\rangle \\ |b\rangle \end{pmatrix}$$

This matrix represents a unitary transformation between the quark mass eigenstates and their weak-interaction eigenstates.

The first row of the CKM matrix provides the most demanding experimental test of the unitarity condition:

$$V_{ud}^2 + V_{us}^2 + V_{ub}^2 = 1$$

Measuring V_{ud} via superallowed Fermi β decay $f t$ and $F t$ values

Hadronic correction, estimated at order $\sim 2\%$ (10% (model dependent))

Hadronic structure-dependent radiative corrections (2%–10%)

$$F t = f t \left(1 + \delta_{NS} + \delta_{SU(2)} - (\delta_{\text{RC}} + \delta_{\text{RC}}) \right) \text{ constant}$$

Dependent on hadronic, hadronic, isospin symmetry-breaking (strong and electromagnetic interactions)

CKM dependent and experimental branching ratio

The conserved vector current (CVC) hypothesis states that the vector coupling constant for the weak interaction is not renormalized in the nuclear medium. For all superallowed β values should be the same after applying small $\sim 1\%$ theoretical corrections. We can then average the experimental $F t$ values to determine a precise value for V_{ud} .

$$V_{ud}^2 = \frac{\sum_i F t_i}{2.3056 \times 10^4}$$

Using radiative corrections (model dependent)

The uncertainty in $F t$ is dominated by the uncertainty in the theoretical corrections, including a contribution accounting for a systematic difference between two different calculations of δ_{NS} using Woods-Saxon and Hartree-Fock radial wave functions.

Since $^{26}\text{Al}^m$ has the smallest nuclear structure dependent correction of all superallowed transitions, a higher precision $^{26}\text{Al}^m$ value will decrease these calculations for the other transitions as well.

Experiment and Data Analysis

The experiment was performed at TRIUMF, Canada's National Laboratory for Particle and Nuclear Physics Research, where 40 nA of 200 MeV protons impinged upon a ^{26}Al target and the resulting radioactive reaction products were created and sent through a mass separator.

A 30 nA $^{26}\text{Al}^m$ beam was implanted under various sizes (22 mm wide Mylar tape for 6–34). Due to contamination from $^{26}\text{Al}^g$ ($T_{1/2} = 1.0728(25)$ s) the beam spot was allowed to cool for 20–25 s, reducing the $^{26}\text{Al}^g$ contamination to a negligible level before moving the $^{26}\text{Al}^m$ sample to the centre of a 40-cm continuous flow gas proportional counter, where the beta decays were counted.

The β particles from the decay of the sample were registered independently using two multichannel scaler modules (MCS) with 250 time bins of variable width. The bin times were varied on a run-by-run basis.

The decay data were analyzed using a maximum likelihood method to extract the $^{26}\text{Al}^m$ half-life. The $^{26}\text{Al}^m$ half-life was determined to be $6.344(4)$ s.

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A final systematic uncertainty estimated by propagating the data fits by the statistical and experimental uncertainties using the method outlined in the text. The final value for the $^{26}\text{Al}^m$ half-life is:

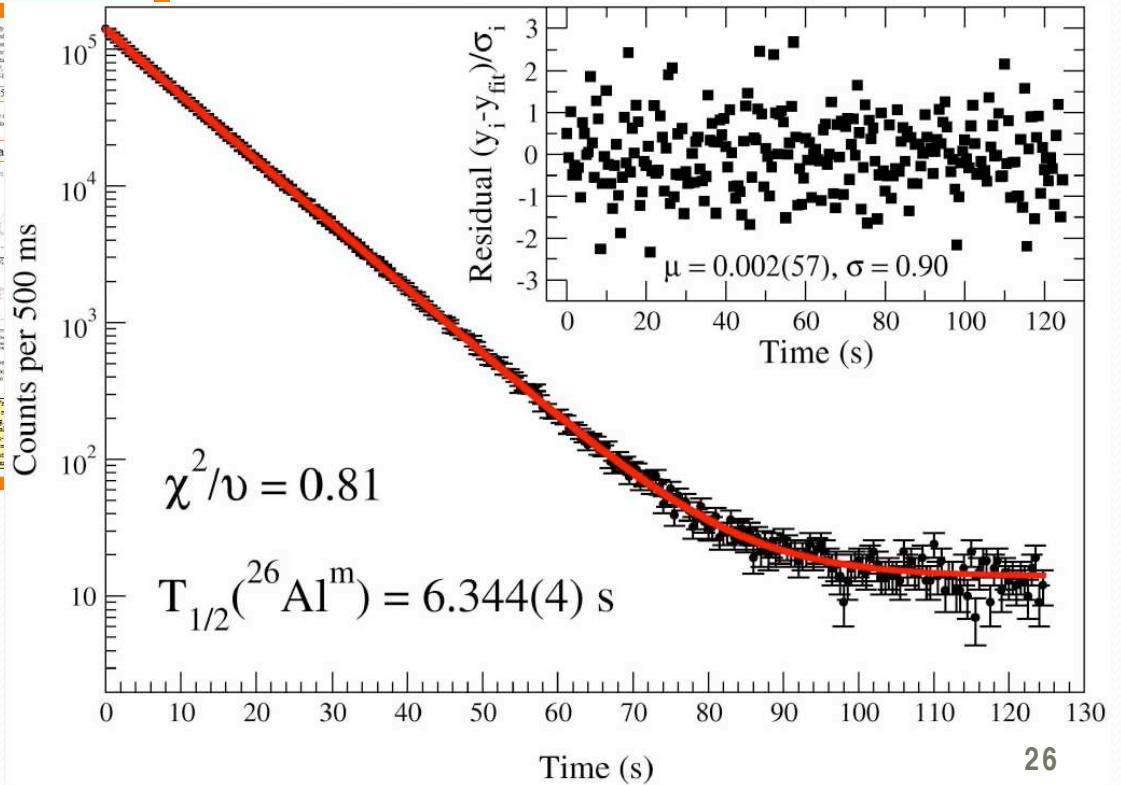
$$T_{1/2}(^{26}\text{Al}^m) = 6.3445$$

At 0.01%, this represents the measurement of any superallowed β transition.

Compa

AP- $F t$ value agrees with the value determined by the $^{26}\text{Al}^m$ value also determined by this analysis.

The high precision AP- $F t$ values for the other superallowed β decays of the first row are: ^{10}C , ^{14}O , ^{18}F , ^{22}F , ^{26}Mg , ^{34}S , ^{38}S , ^{42}Ca , ^{46}Ca , ^{50}Ca , ^{54}Fe , ^{58}Fe , ^{62}Ni , ^{66}Ni , ^{70}Zn , ^{74}Zn , ^{78}Zn , ^{82}Ge , ^{86}Ge , ^{90}Se , ^{94}Se , ^{98}Kr , ^{102}Kr , ^{106}Xe , ^{110}Xe , ^{114}Ba , ^{118}Ba , ^{122}Ra , ^{126}Ra .



^{62}Ga Branching Ratio

P.Finlay et al. Phys. Rev. C 78 025502 (2008), B.Hyland et al. Phys. Rev. Lett. 97 102501 (2006)

- To measure a B.R. to states in the β decay daughter we need
 - Total number of β decays (count β particles)
 - What fraction feeds each state (count γ rays)

$$BR(j) = \frac{N_{\gamma}(j)}{N_{\beta}}$$

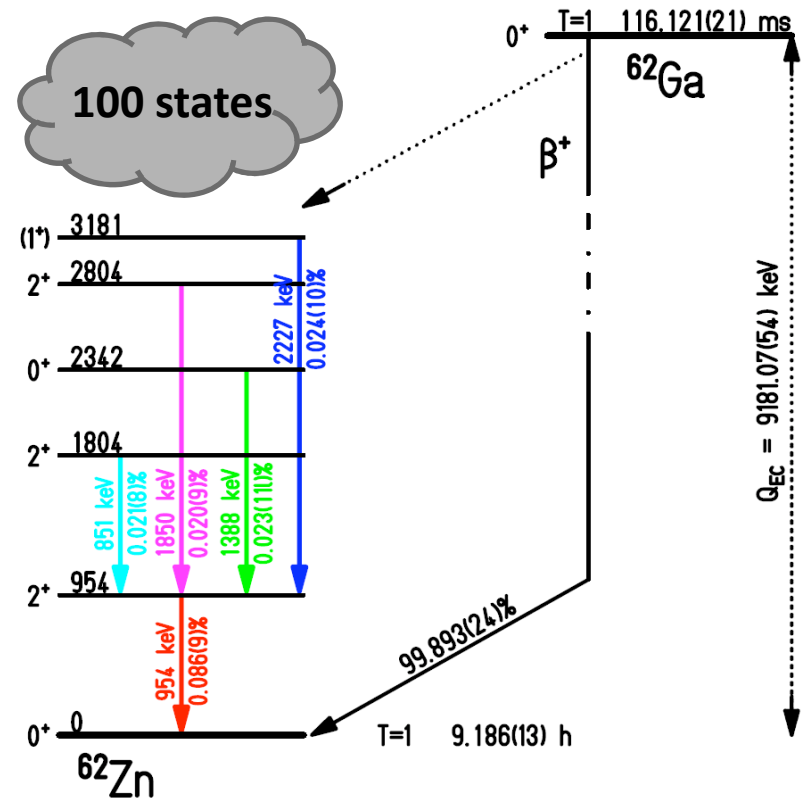
- Ground state – no γ rays

$$BR(\text{g.s.}) = 1 - \sum_{j=0}^n BR(j)$$

- Obstacles for high-precision
 - Efficiency and “Pandemonium”

$$N_{\gamma}(j) = \frac{N_{\gamma}^{exp}(j)}{\epsilon(j)}$$

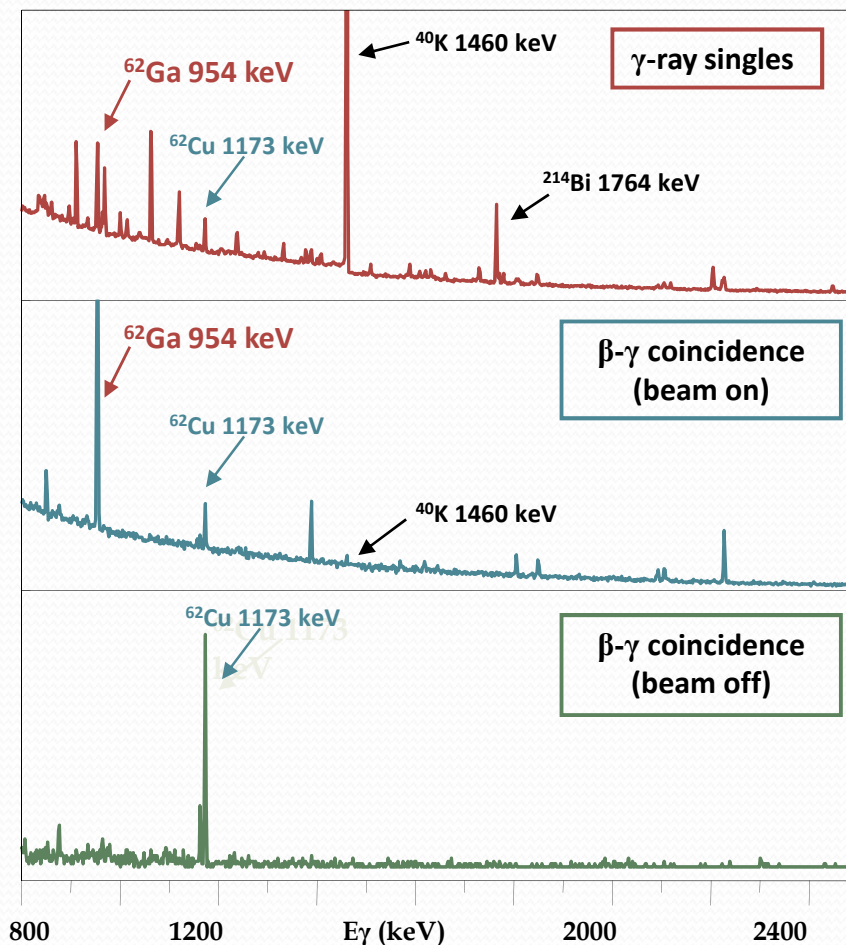
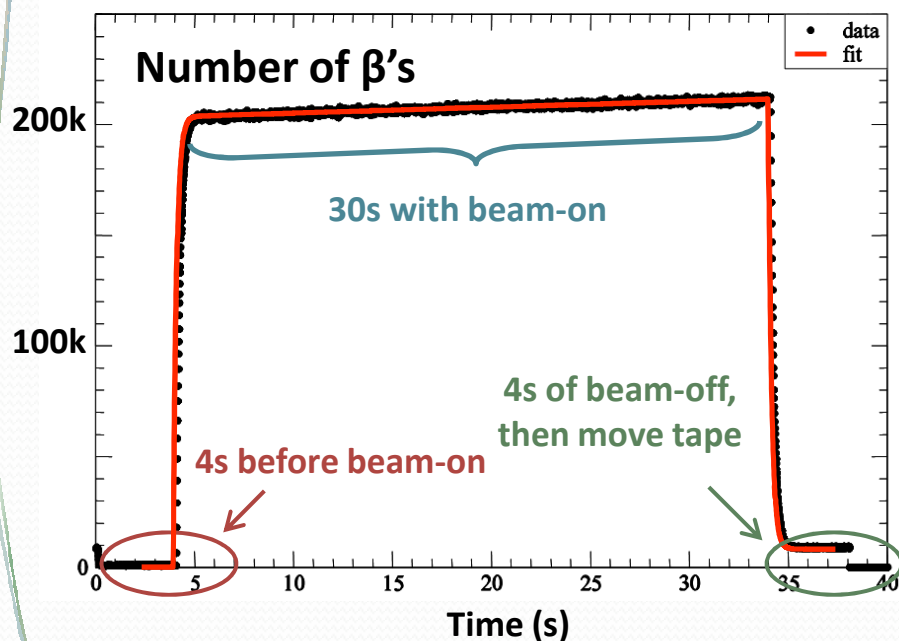
Uncertainty > 1%



^{62}Ga Branching Ratio

P.Finlay et al. Phys. Rev. C 78 025502 (2008), B.Hyland et al. Phys. Rev. Lett. 97 102501 (2006)

- Experiment at TRIUMF
 - 8π spectrometer
 - Inner array of β counters
 - $N_{\beta} = 6.3 \times 10^8$



^{62}Ga Branching Ratio

P.Finlay et al. Phys. Rev. C 78 025502 (2008)

- Observed + Pandemonium (Shell-model calculation):

$$\sum_{j=0}^n BR(j) = \mathbf{0.134(3)} + \mathbf{0.008(8)} = \mathbf{0.142(8)\%}$$

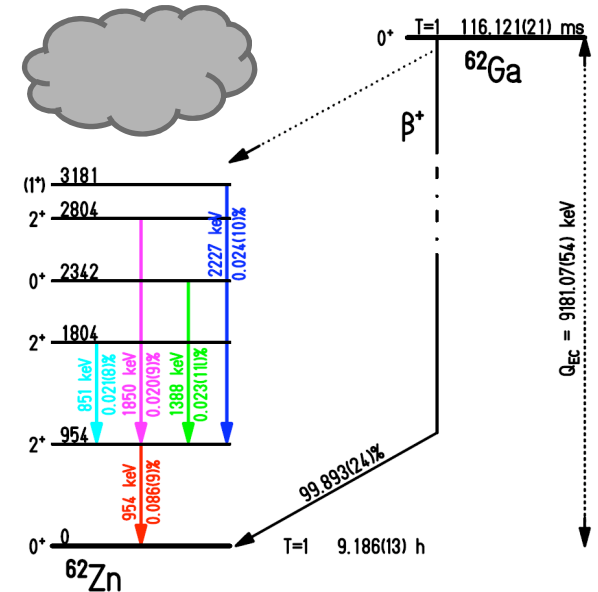
2% precision (pointing to 0.134(3))
6% precision (pointing to 0.008(8))

- Ground-state branching ratio

$$BR(\text{g.s.}) = 1 - \sum_{j=0}^n BR(j) = \mathbf{100 - 0.142(8)\%}$$

$$= \mathbf{99.858(8)\%}$$

0.008% precision !! (pointing to 0.008(8))



- As a general rule

- Works only if S.A.F. decay is to ground state and B.R. > 95%

$$\sum_{j=0}^n BR(j) = \mathbf{5.0(1)\%}$$

2% precision (pointing to 5.0(1))

$$BR(\text{g.s.}) = 1 - \sum_{j=0}^n BR(j) = \mathbf{95.0(1)\%}$$

0.1% precision (pointing to 95.0(1))

Experimental Impact for ^{62}Ga

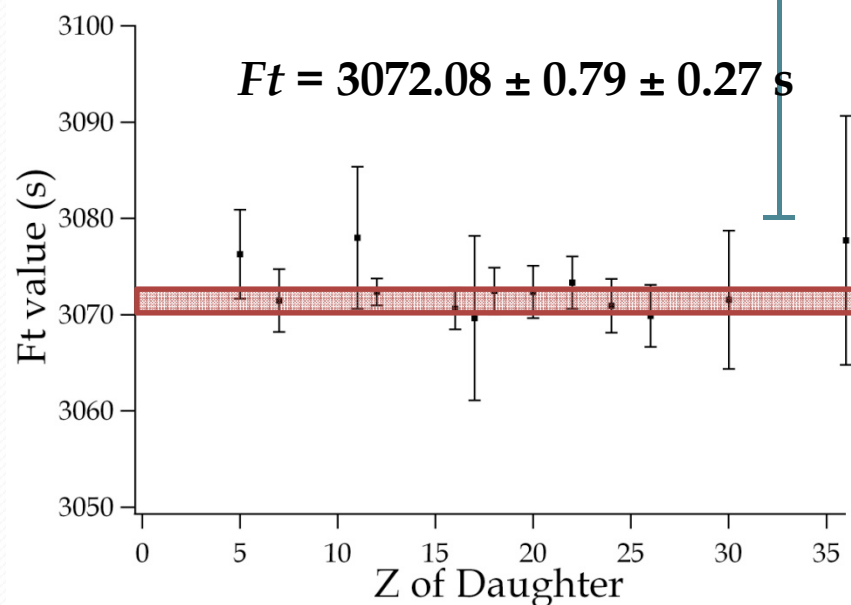
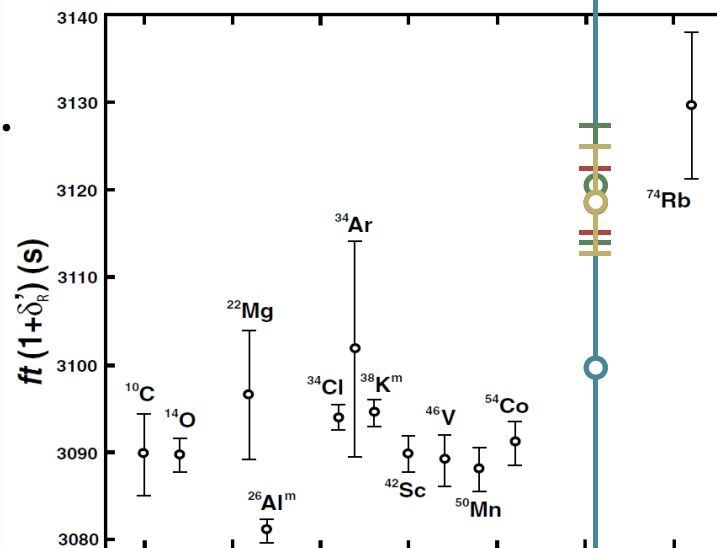
- Status of the ^{62}Ga ft value in 2005...

- $ft(1+\delta_R) = 3102(48) \text{ s}$
- $+Q\text{-value} = 3121(6) \text{ s}$
- $+T_{1/2} = 3119(5) \text{ s}$
- $+B.R. = 3119(3) \text{ s}$
- 16x improvement!

- Corrected Ft value for ^{62}Ga

- $Ft(2005) = 3058(47) \text{ s}$
- $Ft(2010) = 3072(7) \text{ s}$

Uncertainty dominated by theory
(experimental contribution negligible)



Can ft to Constrain Theory

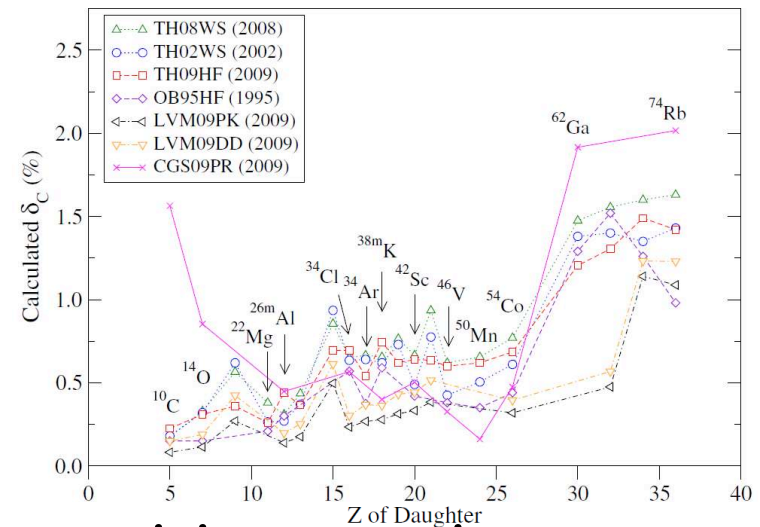
- ^{62}Ga Ft value limited by theory

Theoretical (TH08WS) $\delta_C = 1.5(2) \%$

Theoretical (TH02WS) $\delta_C = 1.4(2) \%$

Theoretical (TH09HF) $\delta_C = 1.2(7) \%$

Theoretical (OB95HF) $\delta_C = 1.3(2) \%$



- Constrain calculations using high-precision experiment
 - Assuming CVC we compute an “experimental” δ_C for ^{62}Ga

$$(1 - \delta_C) = \frac{(\overline{Ft})^*}{ft(1 + \delta'_R)}$$

“experimental” theoretical correction \rightarrow $(1 - \delta_C)$
 \leftarrow World average of 12 others $(\overline{Ft})^*$
 \leftarrow Rad. Correction (^{62}Ga) $(1 + \delta'_R)$
 \leftarrow Experimental ft (^{62}Ga) ft

- Using $ft = 3074.1(15) \text{ s}$, $\delta_R = 1.46(9) \%$, $Ft = 3072.1(8) \text{ s}$

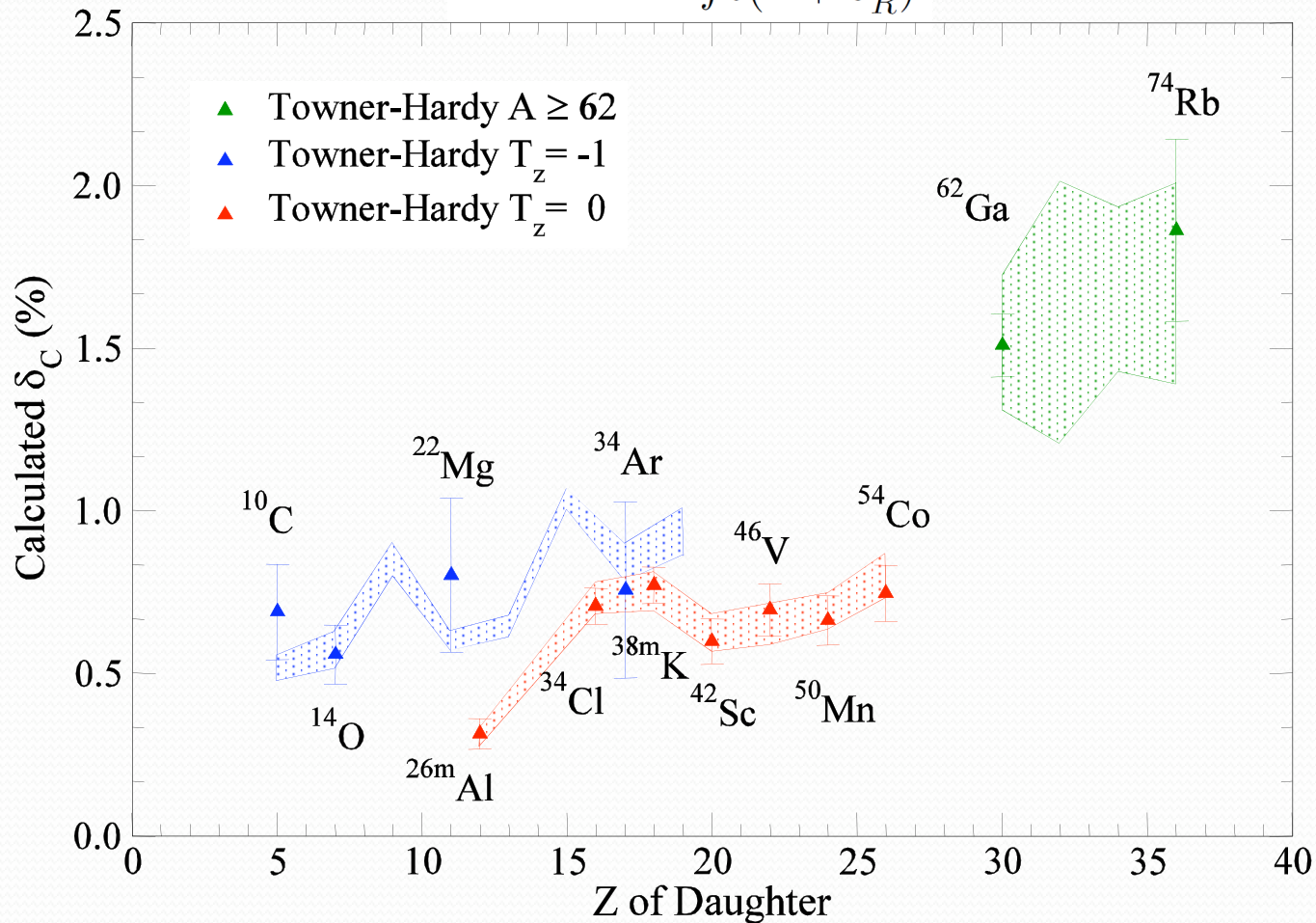
Experimental $\delta_C = 1.45(10) \%$

Not an ABSOLUTE test!

“Experimental” Corrections

- Using the THo8WS calculations, $Ft = 3072.1(8)s$

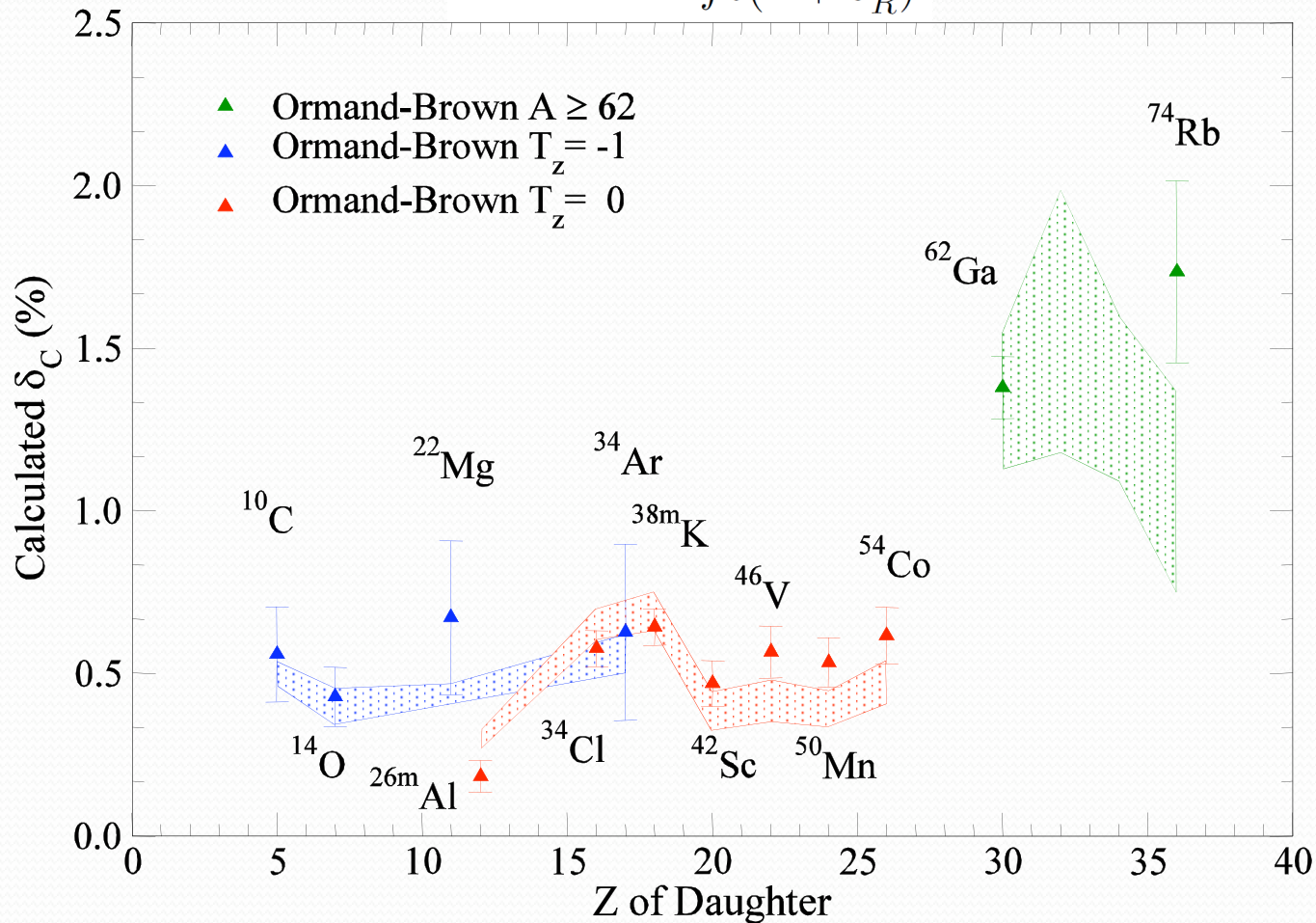
$$(1 - \delta_C) = \frac{(\overline{Ft})^*}{ft(1 + \delta'_R)}$$



“Experimental” Corrections

- Using the OB95HF calculations, $Ft = 3076.2 \text{ s}$

$$(1 - \delta_C) = \frac{(\overline{Ft})^*}{ft(1 + \delta'_R)}$$



Alternative Approaches

- Spectroscopic factors from low-energy transfer reactions
 - See poster by Kyle Leach (University of Guelph)

$$\delta_{C2} \approx \sum_{\pi, \alpha} \frac{T_f(T_f + 1) + \frac{3}{4} - T_\pi(T_\pi + 1)}{T_f(T_f + 1)} S_{\alpha, T_f}^{T_\pi} \Omega_\alpha^\pi$$

Experimental Guidance of ISB Corrections via Direct Nuclear Reactions
 K.G. Leach¹, P.E. Garrett², G.C. Ball³, J.C. Bangay⁴, L. Bianco⁵, G.A. Demand⁶, T. Faesslermann⁷, P. Finlay⁸, K.L. Green⁹, R. Harstenberger¹, R. Krücken¹, A.A. Phillips¹, E.T. Rand¹, C.S. Sumthararatchi¹, C.E. Svensson¹, I.S. Towne¹, S. Triambak¹, H.-F. Wirth¹, and J. Wong¹
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Motivation

Through precision measurements of f_1 values, studies of superallowed $0^+ \rightarrow 0^+$ nuclear β decays provide rigorous tests of the Standard Model description for electroweak interactions. The conserved vector current (CVC) hypothesis states that the weak vector coupling constant G_V is not renormalized within the nuclear medium. The corrected f_1 values are thus expected to be nucleus independent.

$\mathcal{F}t = f_1(1 + \delta_R)(1 - \delta_C) = \frac{K}{2G_V^2(1 + \Delta_R)} = \text{constant}$

The figure above shows the experimental f_1 values, as well as the corrected $\mathcal{F}t$ values for the 13 most precisely measured superallowed nuclei.

Isospin Symmetry Breaking

The ISB correction can be represented as a sum of two terms [1]. The first term, δ_C , accounts for the first-order effects of isospin mixing between parent and daughter nuclear states. The second term, δ_R , corrects for imperfect radial overlap between the initial and final spatial wavefunctions, which core orbitals are important to include in the radial overlap of the β transitions. Corrections are determined from an examination of experimental spectroscopic strengths in similar nuclear decays [1].

Where δ_C is the spectroscopic factor for pickup of a neutron in the final state α from an i particle state of spin T_i or $T_i \pm 1$ for a particular state of spin T_f .

Recent revisions to the CVC hypothesis in the $A \approx 20$ region [1].
 Accurate calculation of the spectroscopic factor requires a larger number of particles or holes are included in the calculation.
 Since the ^{68}Ga f_1 value is extremely precise, we need a precise and accurate f_1 value in order to provide a stringent test of CVC in the upper p -shell.

Neutron Particle-Hole Ground State Configurations

Current Tower and Hardy ^{68}Zn Calculation Model Space [2]	^{68}Zn Experimentally Accessible
$1g_{7/2}$	$1g_{7/2}$
$2p_{1/2}$	$2p_{1/2}$
$1f_{7/2}$	$1f_{7/2}$
$2p_{3/2}$	$2p_{3/2}$
$1f_{5/2}$	$1f_{5/2}$

⁶⁸Ni Closed-Shell Core

$^{68}\text{Zn}(d,p)^{68}\text{Zn}$ Experiment

Observing the states populated in the single neutron pickup reaction $^{68}\text{Zn}(d,p)^{68}\text{Zn}$ will help determine the relative strengths of each particle-hole configuration. This experiment directly probes neutron hole states in ^{68}Zn .

Experimental Details

- MP tandem Van de Graaff and Stern-Gerlach polarization source provided a 22 MeV deuteron beam with 60-64% polarization
- Beam was incident on 126 $\mu\text{g}/\text{cm}^2$ of ^{68}Zn with a 13 $\mu\text{g}/\text{cm}^2$ carbon backing

MLL-LMU OSD Magnetic Spectrograph

^{68}Zn Elastic Scattering and Global OMP Comparison

- a) The ratio to Rutherford scattering cross-sections
- b) The scattering cross-section polarization asymmetries

Comparison of four global optical-model parameter (OMP) sets with the experimental data, for 22 MeV deuteron-elastic scattering from ^{68}Zn

The set of An and Cai (2006) [Ref. [5]] best reproduce the observed elastic scattering data

$^{68}\text{Zn}(d,p)^{68}\text{Zn}$ Experimental Data

Outgoing protons were detected at 9 angles between 10° and 60° , covering excitation energies in ^{68}Zn of up to ~ 4 MeV, with both polarizations.

- ^{68}Zn energy spectrum above for 22 MeV ^{68}Zn polarized deuterons at $\theta_{lab} = 30^\circ$
- Angular distributions and asymmetries have currently been constructed for more than 100 peaks

Preliminary Comparison with Theory for $^{68}\text{Zn}(d,p)^{68}\text{Zn}$

2002 Calculations [3]

- Modified Surface-Orbit Interaction (mSOI) with a closed ^{68}Ni core

2006 Calculations [1]

- mSOI and OMP1 with a ^{68}Ni core, opening up to one $f_{7/2}$ hole

Experimental f_1 value is shown for comparison.

Summary

- Experimental f_1 values are needed to help guide superallowed ISB calculation model space truncations
- The need for guidance is most evident in the ^{68}Ga β decay daughter nucleus ^{68}Zn , since the experimental f_1 value of ^{68}Ga is extremely precise
- Ultimately, these reactions will determine which core orbitals are important to include in the radial overlap calculation model space for heavier superallowed nuclei, thus providing a more stringent test of the standard model

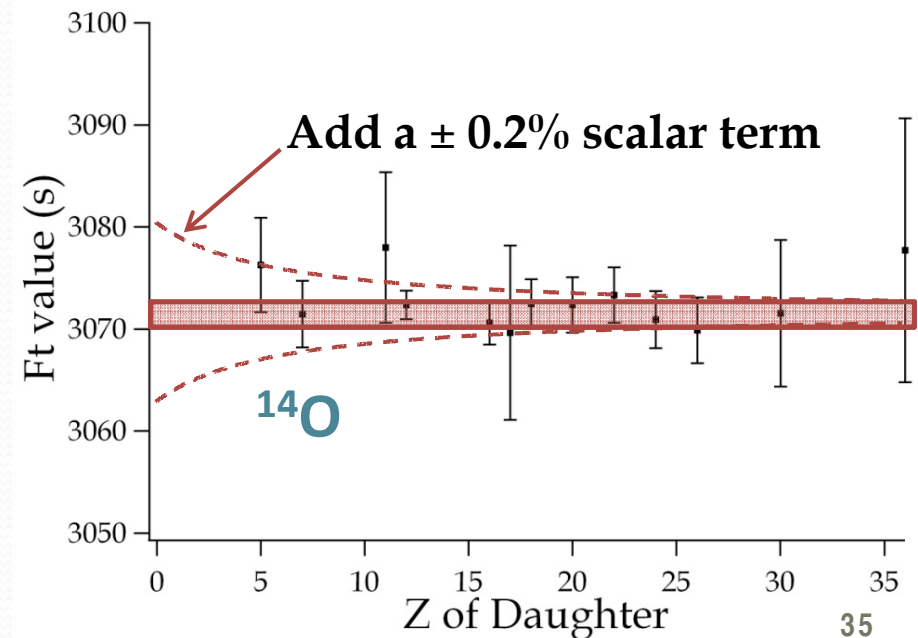
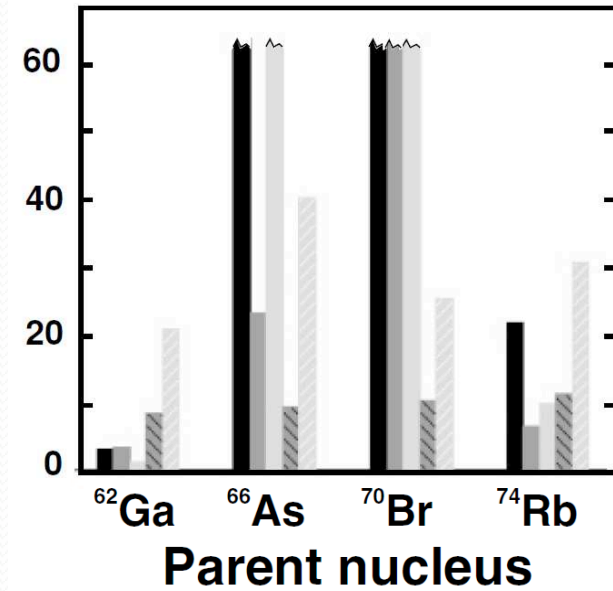
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- [2] J.C. Hardy and I.S. Towne, Phys. Rev. C 79, 055502 (2009)
- [3] I.S. Towne, Private Communication (2009)
- [4] H.-F. Wirth et al., MLL-LMU Jahresbericht, 71 (2000)
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On The Horizon

- Other heavy cases: ^{66}As , ^{70}Br , ^{74}Rb
 - All similar in principle to ^{62}Ga
 - Beam production is a challenge
- Light decays ^{10}C , ^{14}O
 - Search for scalars in β decay
 - *ab initio* calculations of δ_C
- Alternative techniques
 - $T=1/2$, $T=2$ decay
 - Transfer reactions



Summary

- Most precise way to extract V_{ud} (by 6x) is through ft values of S.A.F. decay
 - Unlikely to change in the next 10 yrs
- Uncertainty dominated by theory
 - I.S.B corrections are an essential ingredient for CVC and V_{ud}
 - Range of δ_C values unsatisfactory
- New facilities with state-of-the-art experiment and theory
 - Leading to an improved understanding of these effects

