

# Hadron Physics

Lecture #1: The quark model, QCD, and hadron spectroscopy

Lecture #2: Internal structure of hadrons: momentum and spin

Lecture #3: Internal structure of hadrons: charge, magnetism, polarizability

Lecture #4: Hadrons as laboratories (weak interactions, fundamental symmetries, and other miscellaneous topics)

*(much of this material is from my seminar at NNPSS 2008)*

# Hadrons as Laboratories

Weak interactions inside the nucleon

Weak NN interactions

Fundamental symmetry tests with  $A=1$  ( $Z=1$  or 0)

→ see also seminars by *Tim Chupp, John Hardy*

# references for this section

- General:
  - “Subatomic Physics 3<sup>rd</sup> edition”, Henley & Garcia (2007)
  - “Quarks & Leptons”, Halzen & Martin
  - “Experimental Foundations of Particle Physics”, Cahn & Goldhaber
- PV electron scattering:
  - M.J. Musolf et al., Phys. Rep. 239 (1994) 1.
  - E. Beise, M.L. Pitt and D.T. Spayde, Prog. Part. Nucl. Phys. 54 (2005) 289.
  - M. Pitt, NNPSS 2004 Lecture notes (Bar Harbor, ME)
- Neutrons:
  - J. Nico and W.M. Snow, Ann. Rev. Nucl. Part. Sci. 55 (2005) 27.
  - M.J. Ramsey-Musolf & S. Page, Ann. Rev. of Nucl. Part. Sci. 56 (2006) 1.
  - NNPSS 2007 Lecture notes, G. Greene.

Thanks for slides, photos, etc.: B. Filippone, P. Mumm, J. Nico, K. Pashke, M. Pitt, X. Zheng, P. Reimer, H. Gao, J. Martin, and many others

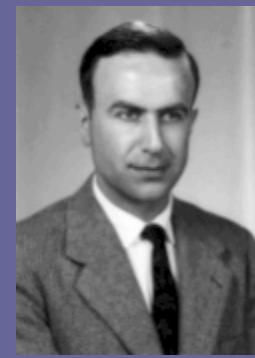
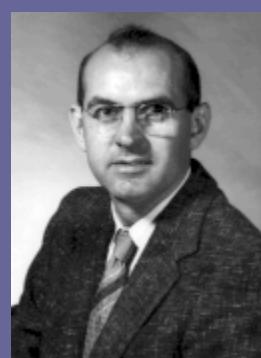
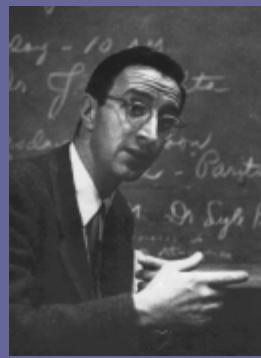
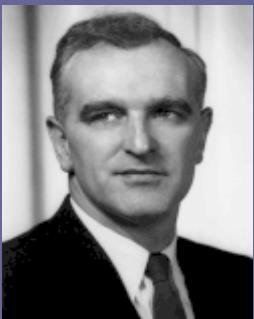
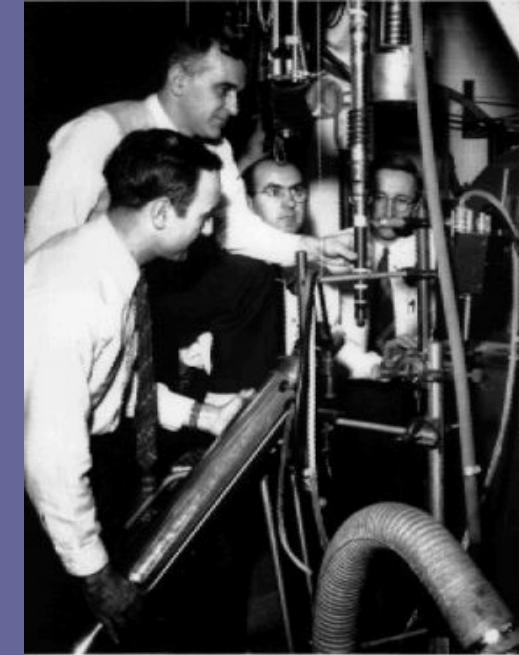
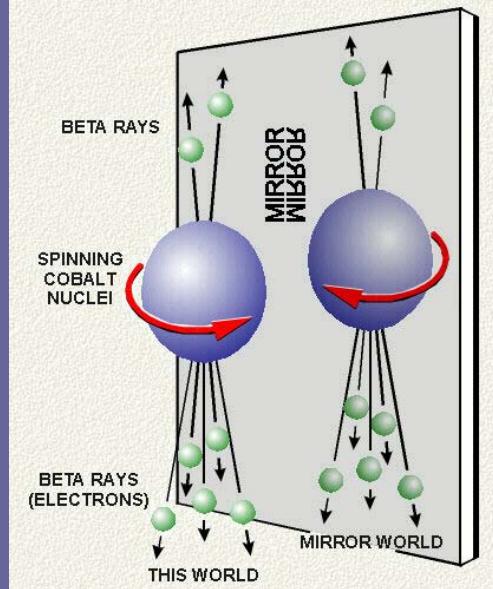
# “The Fall of Parity”

from the NIST virtual museum

<http://physics.nist.gov/GenInt/Parity/cover.html>



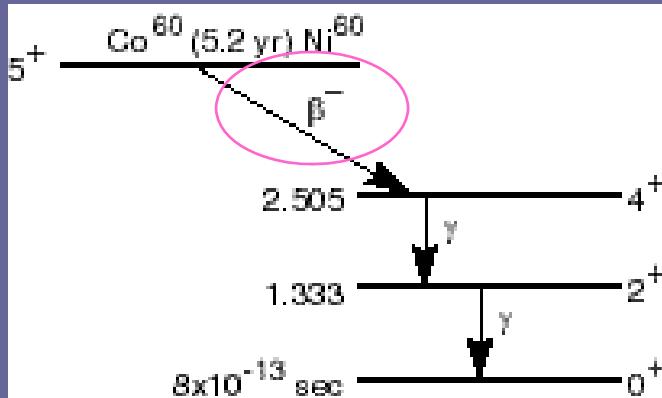
T.D. Lee & C.N. Yang,  
Phys. Rev. 104 (1956) 254.  
(October 1956)



C.-S. Wu, E. Ambler, R. W. Hayward, D. D. Hoppes, and R. P. Hudson,  
Phys Rev 105 (1957) 1413. (January 1957)

# The $^{60}\text{Co}$ experiment

(see the NIST museum!)

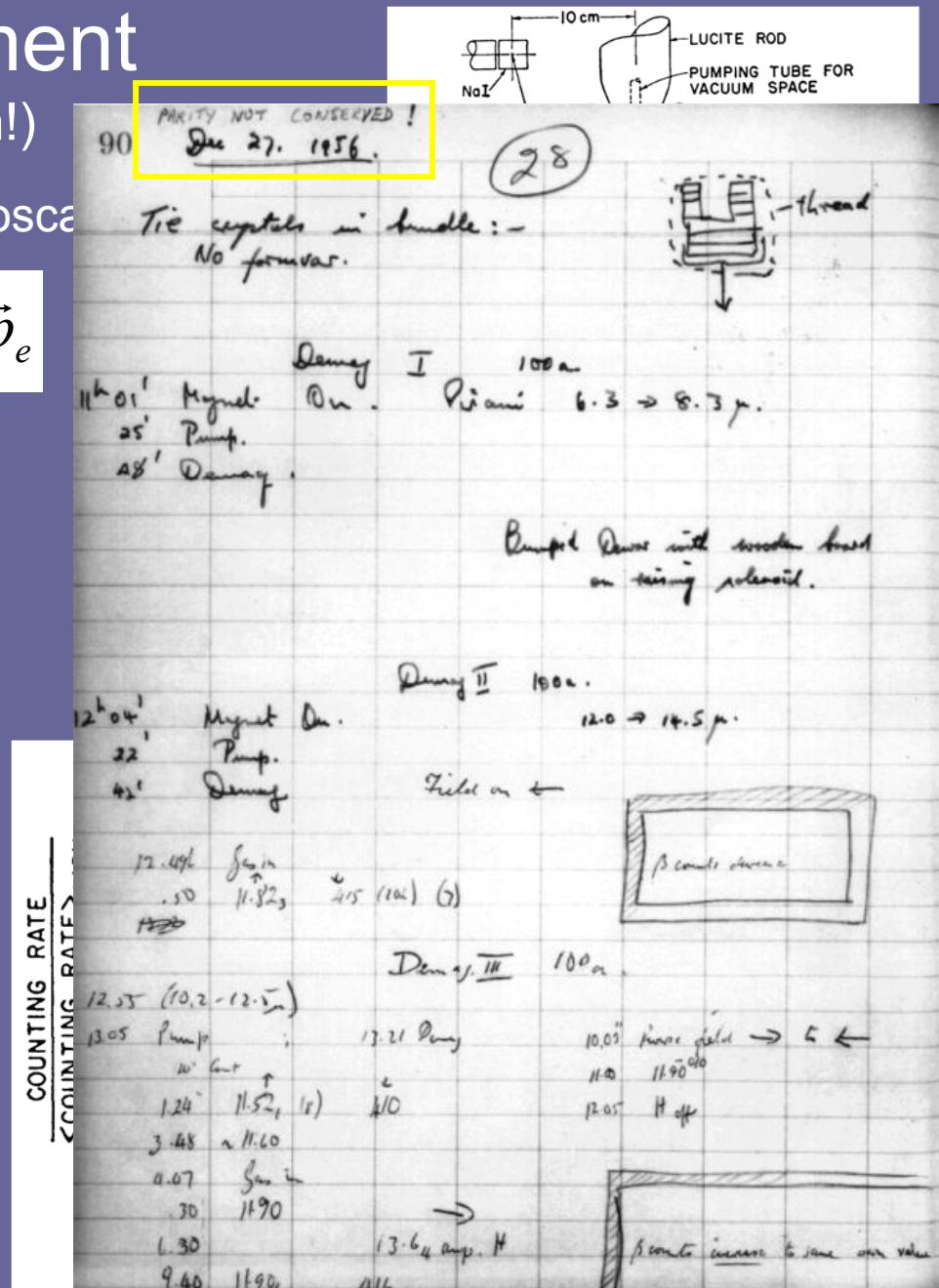


$\langle J \rangle$  = average nuclear polarization of  $^{60}\text{Co}$   
(in B-field and low T (0.01K))

anisotropy of  $\gamma$ 's measures the degree of pol'n of  $^{60}\text{Ni}$  (now used for thermometry – see F. Pobell, "Matter and Methods at Low Temperature")

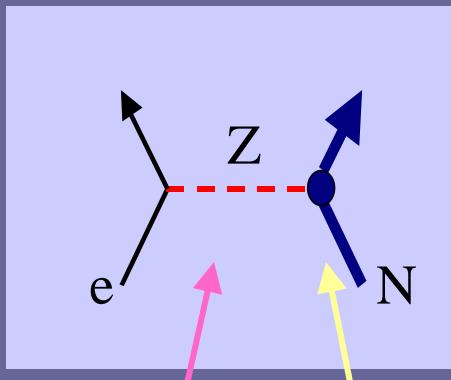
pseudoscalar

$$\langle \vec{J} \rangle \cdot \vec{p}_e$$



# Weak interactions involving nucleons (and electrons)

- electron scattering

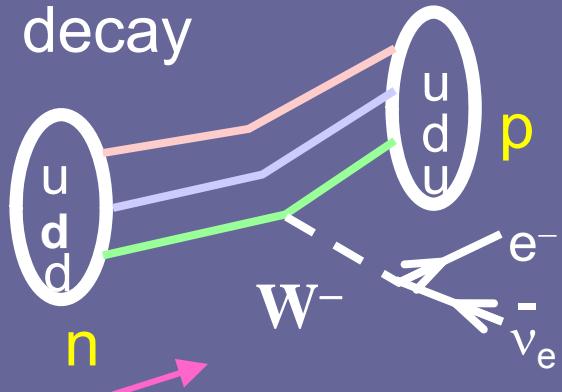


use well-understood  
target (nucleon) to  
perform precision  
test of EW physics

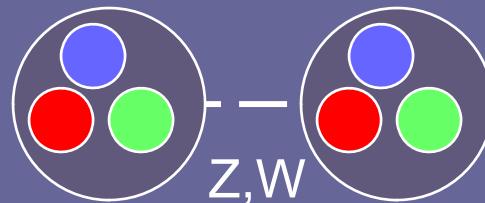
OR

use weak force to probe new aspects  
of strong interaction

- neutron decay



- weak NN interactions



**EXPERIMENTAL SIGNATURE IS  
VIOLATION OF PARITY**

# Neutral current e-quark interactions

$$L_{SM}^{PV} = -\frac{G_F}{\sqrt{2}} \sum_q [C_{1q} \bar{e} \gamma_\mu \gamma_5 e [\bar{q} \gamma^\mu q] + C_{2q} \bar{e} \gamma_\mu e [\bar{q} \gamma^\mu \gamma^5 q]]$$

A x V

V x A

product of V and A violates parity

$C_{iq}$ 's are “weak charges” of quarks (1=“vector”, 2=“axial”)  
If scattering directly from quarks, can measure these, e.g.

$$C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W$$

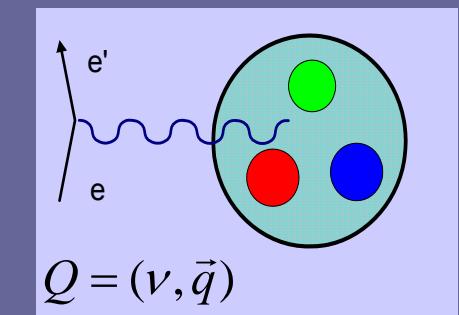
If scattering from a nucleon, have to take into account that quarks are bound inside nucleon via strong interaction to disentangle the  $C_{iq}$ 's:  
these are typically embedded in nucleon “form factors”

# Elastic electron (electromagnetic) scattering from nucleons

quark “charges” are buried in nucleon form factors

$$\frac{d\sigma}{d\Omega} = \sigma_{Mott} \frac{E'}{E} \frac{1}{\varepsilon(1+\tau)} [\varepsilon G_E^2(Q^2) + \tau G_M^2(Q^2)]$$

$$G \rightarrow \langle N | \sum e_q \bar{q} \gamma^\mu q | N \rangle$$



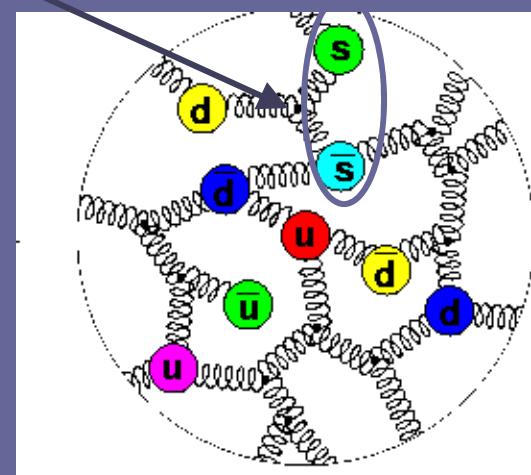
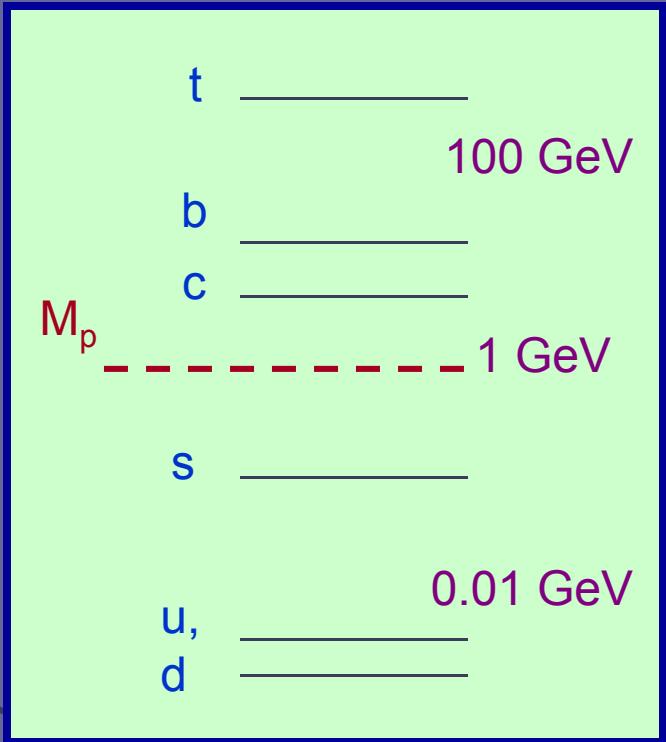
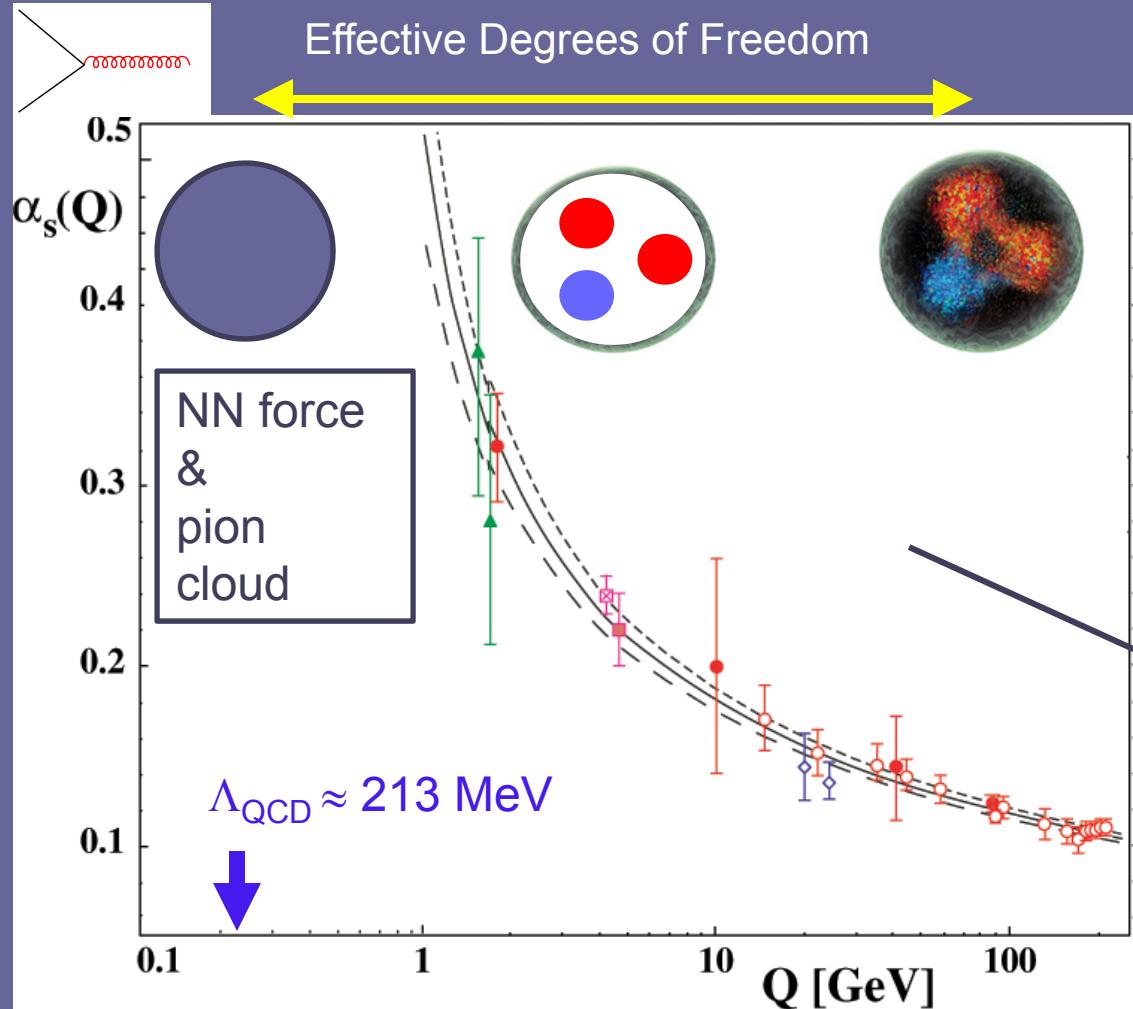
$$G_E^p = \frac{4}{3} G_E^u - \frac{1}{3} G_E^d - \frac{1}{3} G_E^s$$

$$G_E^n = \frac{2}{3} G_E^u - \frac{2}{3} G_E^d - \frac{1}{3} G_E^s$$

contributions to charge,  
similar for magnetism

$G_E$  and  $G_M$  for both proton and neutron are now very well known via precision electron scattering exps. But not enough information to disentangle individual quark pieces from these alone.

# Inside the Nucleon



# Weak nucleon form factors

point-like fermions:

$$\text{EM: } ieQ_f\gamma_\mu$$

$$\text{Weak: } i\frac{gM_Z}{4M_W}\gamma_\mu(g_V^f + g_A^f\gamma_5)$$

	$Q_f$	$g_V^f$	$g_A^f$
v	0	1	-1
e,μ <sup>-</sup>	-1	$-1 + 4 \sin^2 \theta_W$	+1
u,c,t	+2/3	$1 - 8/3 \sin^2 \theta_W$	-1
d,s,b	-1/3	$-1 + 4/3 \sin^2 \theta_W$	+1

nucleons:

$$\begin{aligned} G_{E,M}^{u,p} &= (3 - 4 \sin^2 \theta_W) G_{E,M}^{\gamma,p} - G_{E,M}^{Z,p} \\ G_{E,M}^{d,p} &= (2 - 4 \sin^2 \theta_W) G_{E,M}^{\gamma,p} - G_{E,M}^{\gamma,n} - G_{E,M}^{Z,p} \\ G_{E,M}^{s,p} &= (1 - 4 \sin^2 \theta_W) G_{E,M}^{\gamma,p} - G_{E,M}^{\gamma,n} - G_{E,M}^{Z,p} \end{aligned}$$

}

combine weak and EM information to separate (u,d,s) contributions to proton's charge and magnetism

$$J_{\mu 5}^Z = \bar{u}_N \left[ G_A^Z(q^2)\gamma_\mu + \frac{1}{M_N} G_P(q^2)q_\mu \right] \gamma_5 u_N$$

ignore

axial part is related to proton's spin

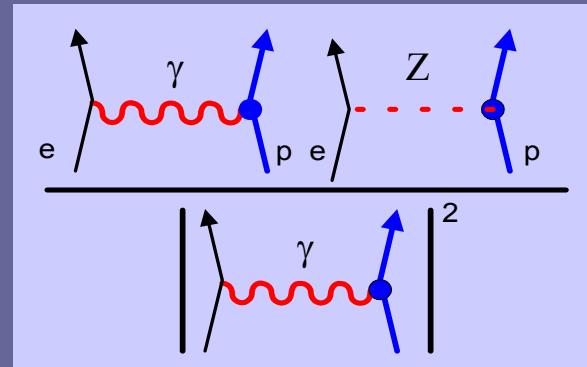
# Parity Violating Electron-Proton Scattering

polarized electrons, unpolarized target

$$A = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}$$

~ few parts per million

and want to measure to ~ 5%



Statistics :  $10^{13}$ - $10^{14}$  counts for precision measurement

high beam intensity, very thick targets, large acceptance detector  
(sometimes can't count individually scattered particles....)

Systematic Effects:

High quality beam (intensity, position, energy) over a long time

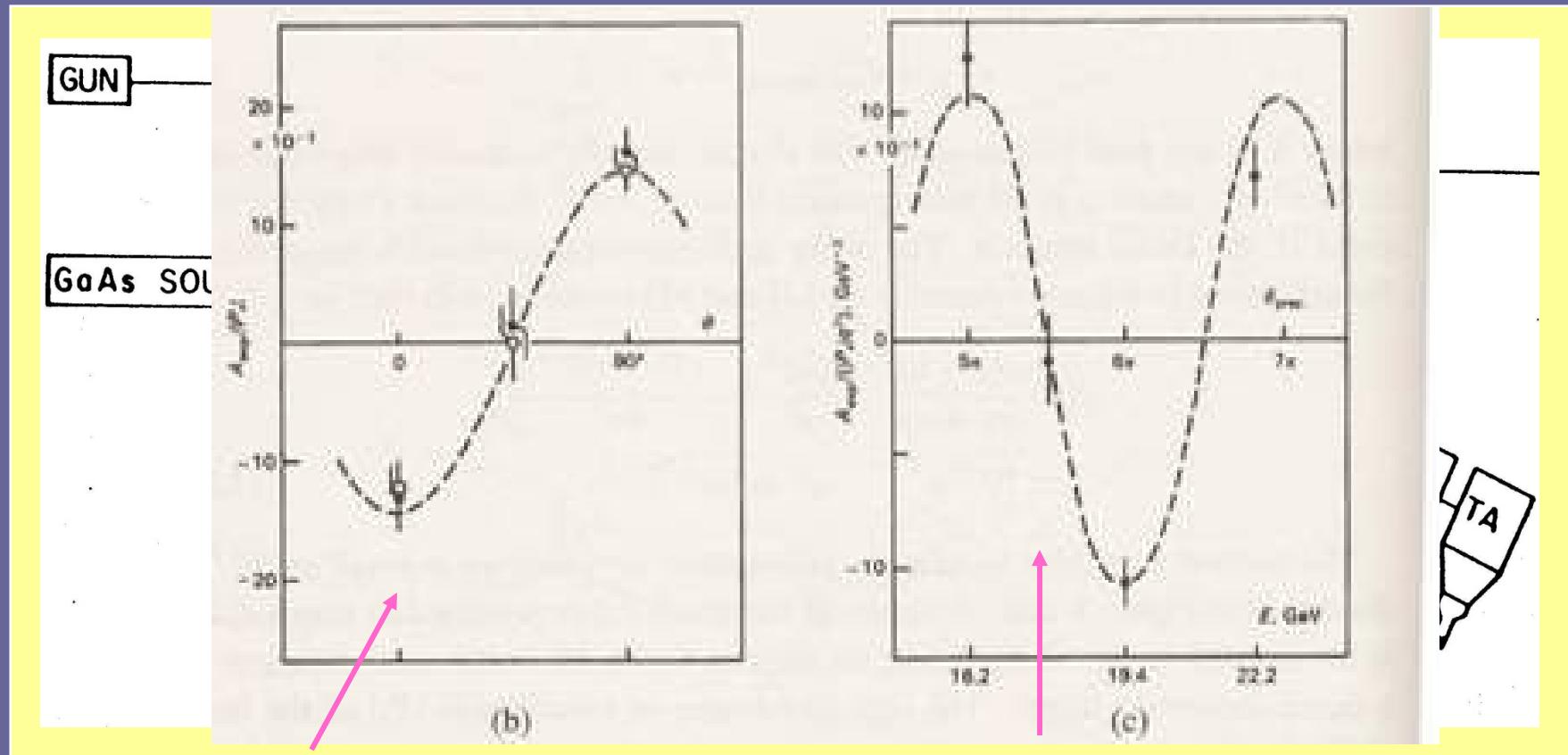
Minimize systematic effects in detection that depend on initial beam state  
(feedback on energy, position, angle and intensity of incident beam)

# Deep inelastic e-D scattering

C.Y. Prescott, et al, PL 77B (1978) 347

SLAC 1978:  $e + d$  (DIS):  $A \sim -60 \pm 10$  ppm

led to correct description of neutral weak force:  $\sin^2 \theta_W = 0.20 \pm 0.03$



reverse handedness of  
electron beam

reverse beam via  
g-2 precession

# SLAC 1978 and JLab 2006

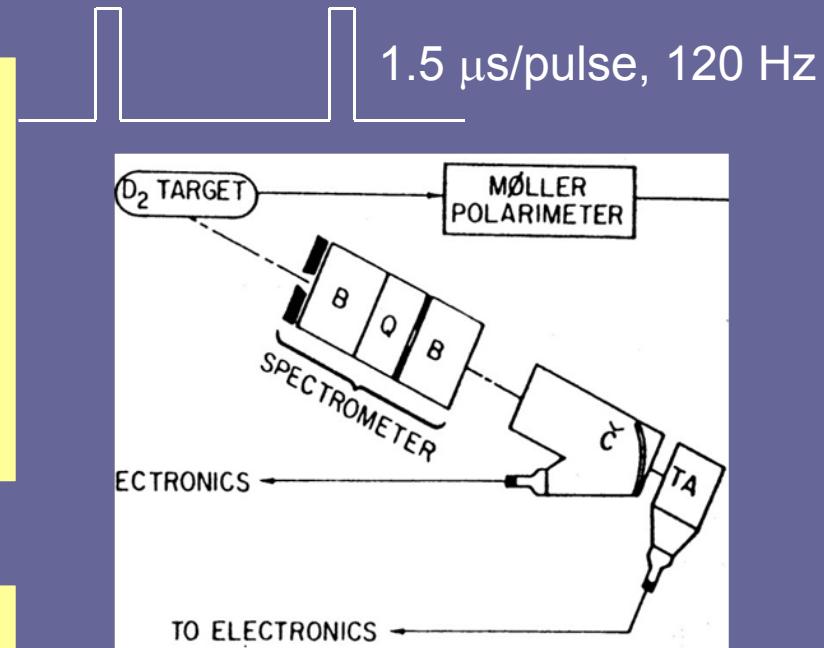
Prescott et al., SLAC DIS e-d:

$$P = 37\%$$

$$L = (7 \mu\text{A}, \text{ pulsed}) \times (30 \text{ cm target})$$

$$\Delta\Omega = 5 \times 10^{-4} \text{ sr}$$

1000 scattered electrons/pulse  
→ integrate



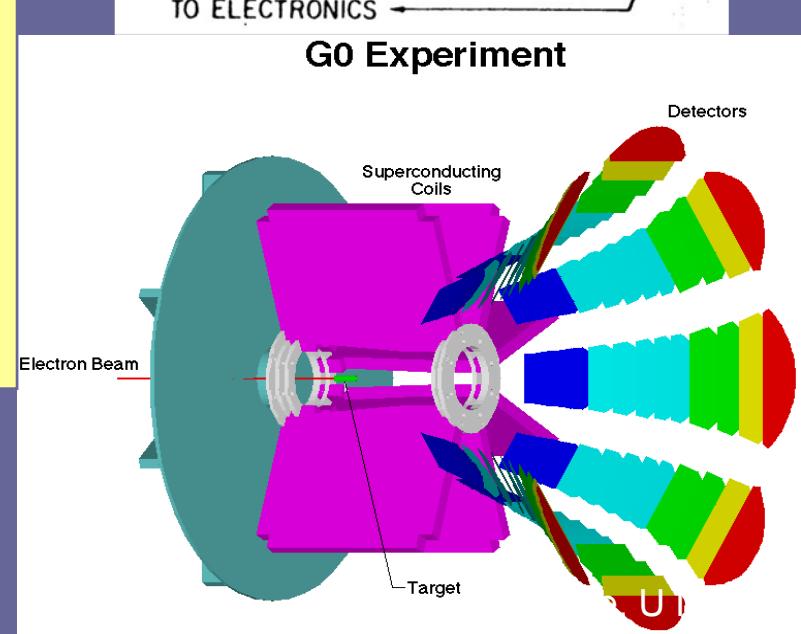
JLAB G0 Experiment (representative):

$$P = 83\%$$

$$L = (60 \mu\text{A}, \text{ CW}) \times (20 \text{ cm target})$$

$$\Delta\Omega = 2 \text{ sr}$$

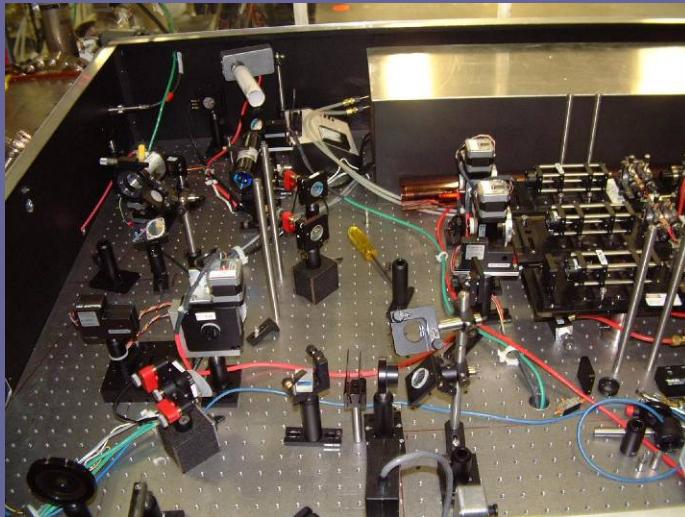
rate ~ maximum few MHz per detector  
→ count



$$P^2 (L) (\Delta\Omega) = (\text{JLAB}/\text{SLAC}) = 10^5$$

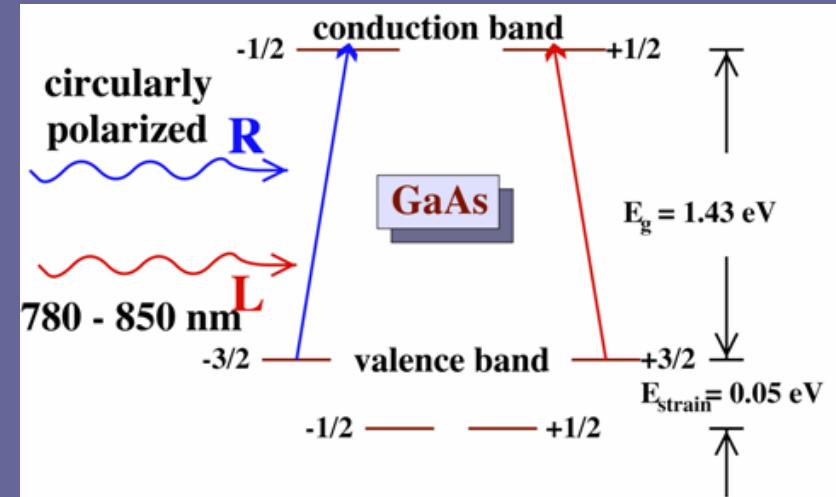
# Polarized Electrons

D.T. Pierce et al., Phys. Lett. 51A (1975) 465.

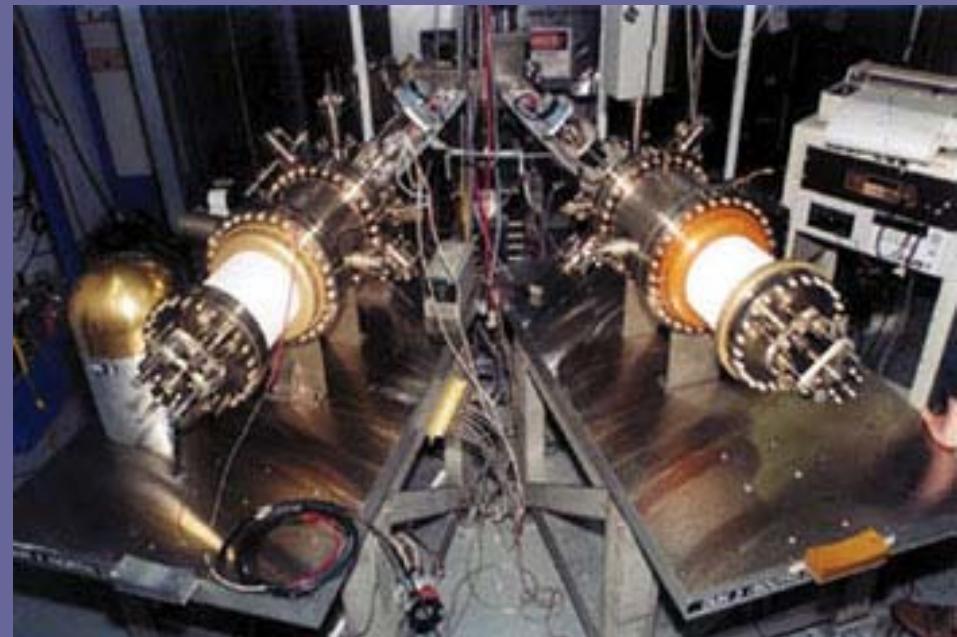


Reverse pol'n of beam  
at rate of 30 Hz

Feedback on laser intensity  
and position at high rate

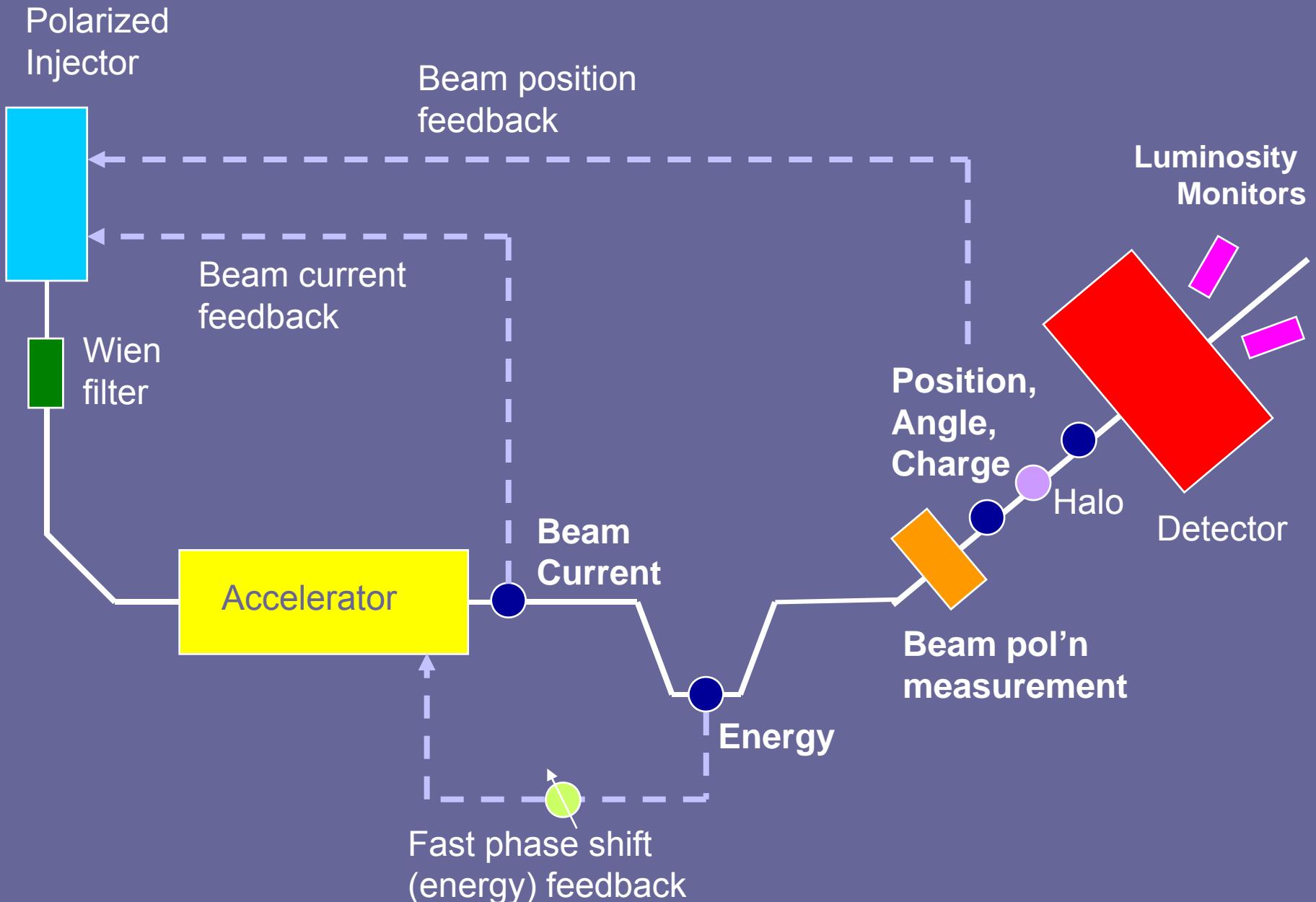


Electron retains circular polarization of  
laser beam:  $P_e \sim 85\%$



See also Physics Today, Dec 2007

# A Typical beam line for PV experiments



# False asymmetries from the beam

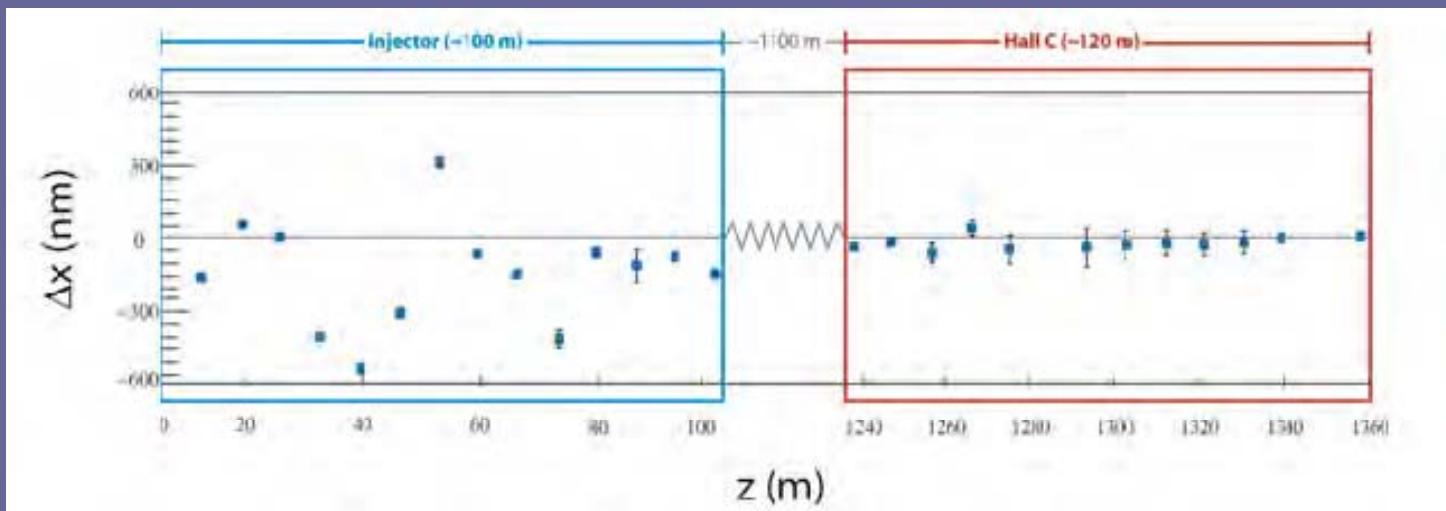
$$A_{meas} = A_{phys} + \sum_{i=1}^N \frac{1}{2Y} \left( \frac{\partial Y}{\partial P_i} \right) \Delta P_i$$

detector sensitivity  
beam property  
(x,θ,E,intensity)



At the injector:  
 $\Delta x < 1$  micron

Measurement at > 1 km:  
 $\Delta x \sim 5 - 10$  nm



# JLab polarized beam : 2005

G0 beam:

- strained GaAs ( $P_B \sim 73\%$ )
- 32 ns pulse spacing
- 40  $\mu\text{A}$  beam current

HAPPEX-II beam:

- superlattice ( $P_B > 85\%$ )
- 2 ns pulse spacing
- 35  $\mu\text{A}$  beam current

Beam Parameter	G0 beam	HAPPEX beam
Charge asymmetry	$-0.14 \pm 0.32 \text{ ppm}$	$-2.6 \pm 0.15 \text{ ppm}$
Position difference	$4 \pm 4 \text{ nm}$	$-8 \pm 3 \text{ nm}$
Angle difference	$1.5 \pm 1 \text{ nrad}$	$4 \pm 2 \text{ nrad}$
Energy difference	$29 \pm 4 \text{ eV}$	$66 \pm 3 \text{ eV}$
Total correction to Asymmetry	$-0.02 \pm 0.01 \text{ ppm}$	$0.08 \pm 0.03 \text{ ppm}$

# Electron accelerators (for nuclear physics)

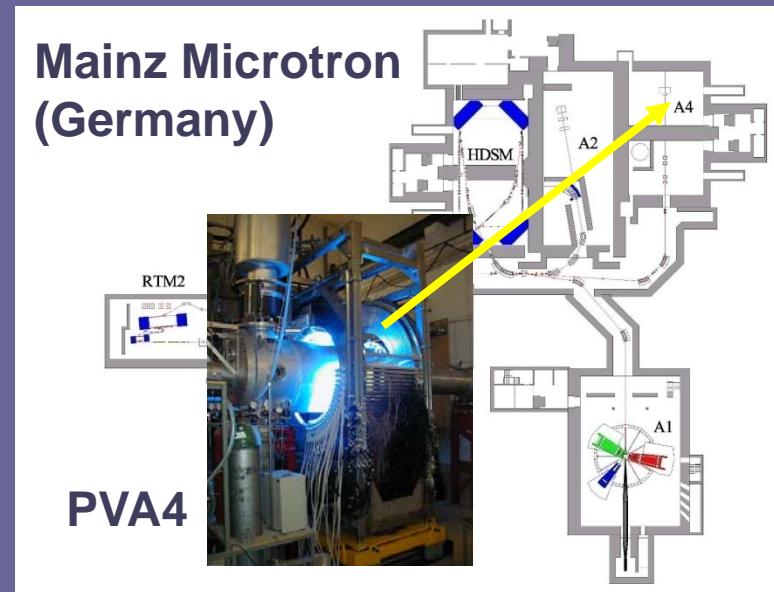
Jefferson Lab,  
Newport News, VA



JLab: 6 GeV beam, getting ready  
for upgrade to 12 GeV

Also, parity-violating Moller scattering  
(E158) carried out at SLAC....

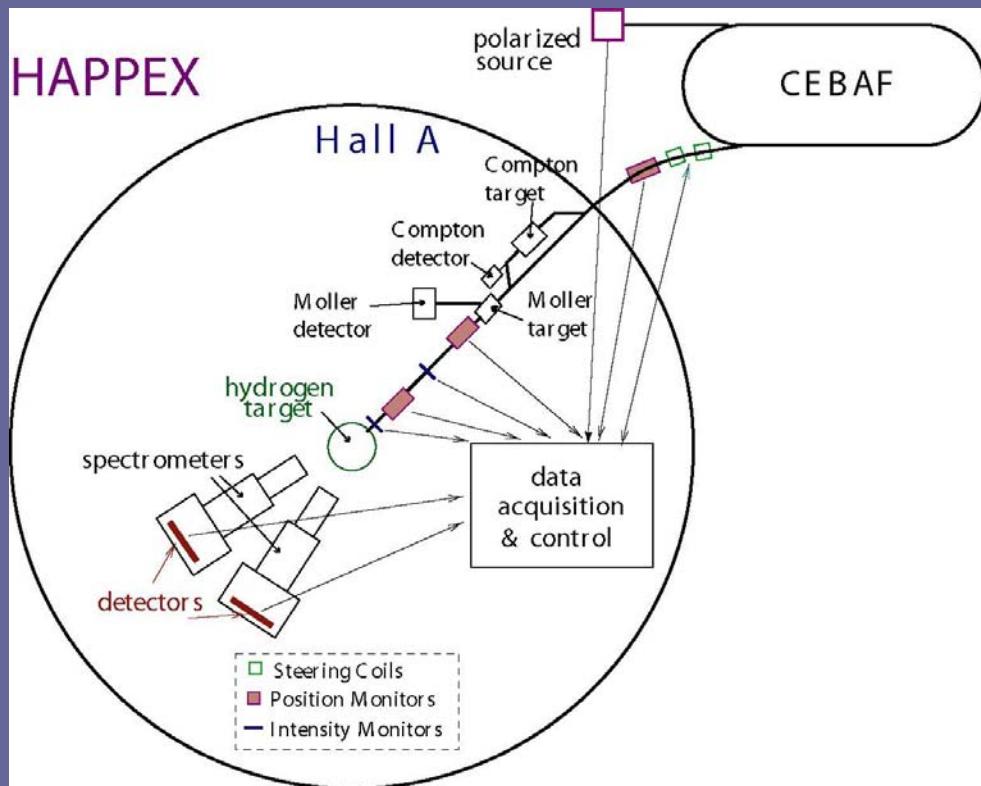
Mainz Microtron  
(Germany)



now the MIT R&E center

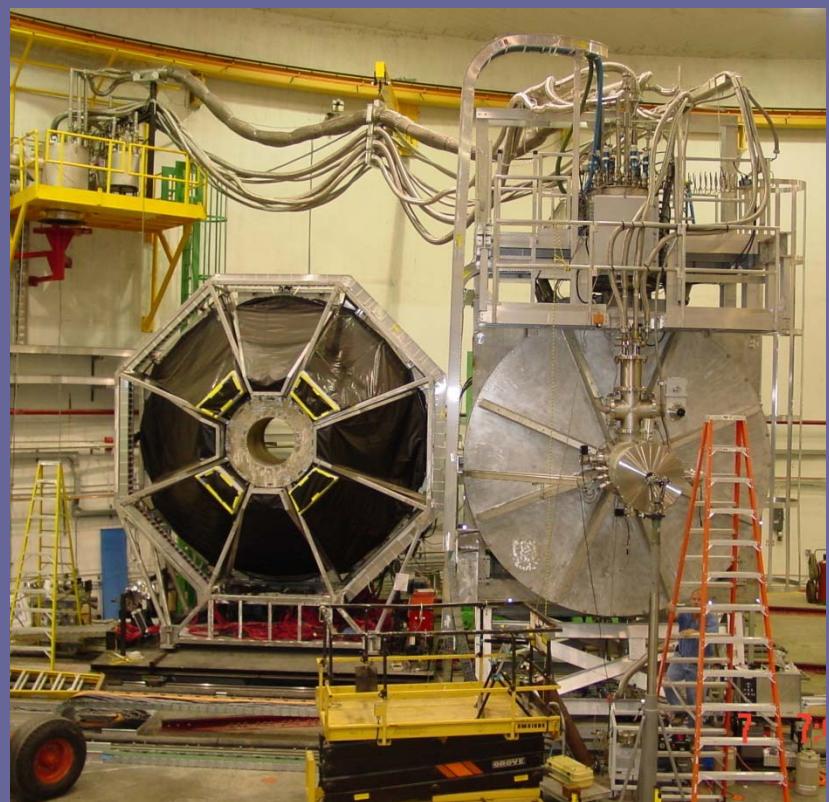
# Weak form factors at JLAB

G-Zero in Hall C



High precision,  
forward angles,  
selected  $Q^2$   
H and  ${}^4\text{He}$  targets

$G_E^s$ ,  $G_M^s$  and  $G_A^e$  separated  
over range  $Q^2 \sim 0.1 - 1.0 (\text{GeV}/c)^2$



H and D targets,  
forward/backward angles

# PV-A4 at Mainz

PbF<sub>2</sub> Calorimeter

phases I &II:  
forward angles, H target

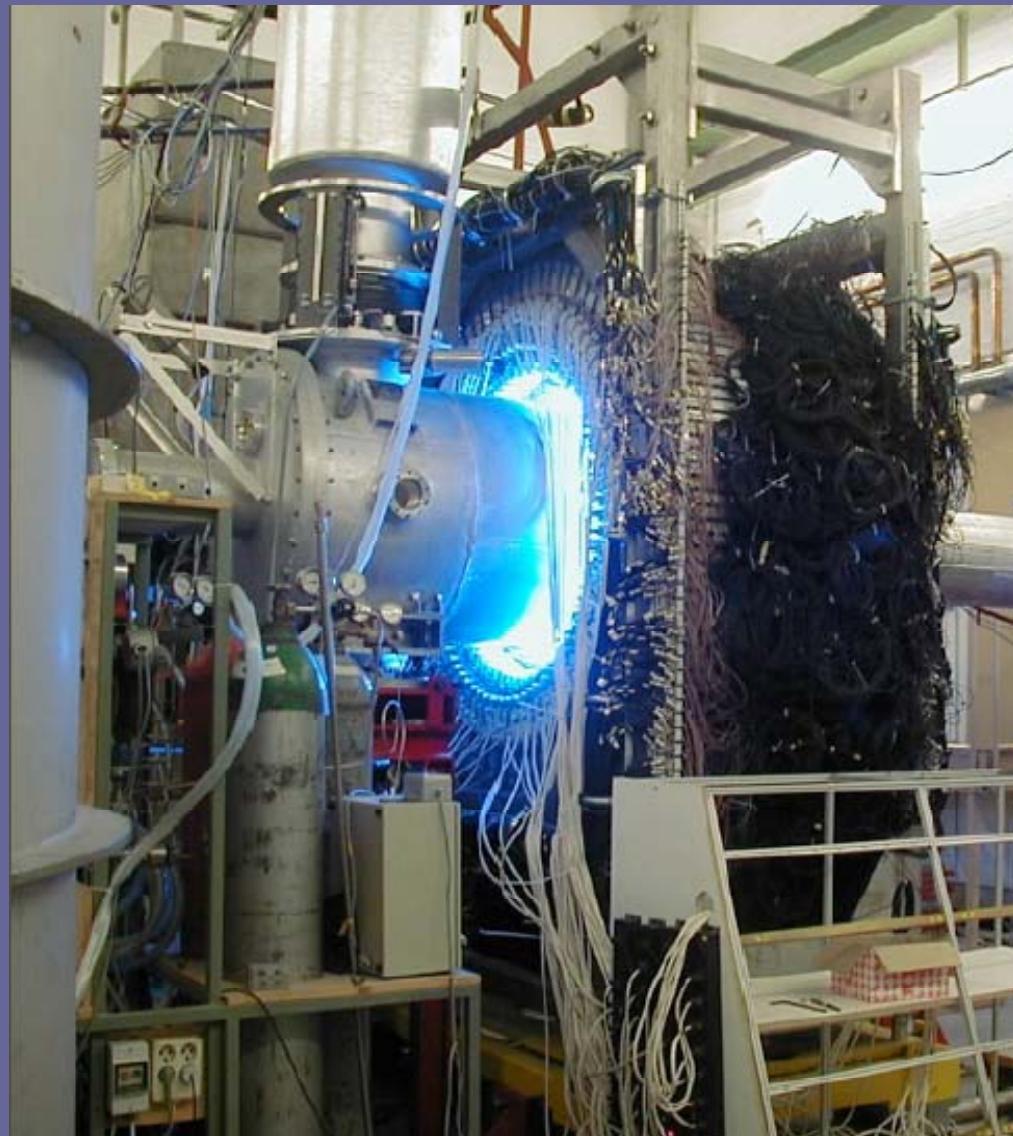
$$\text{I: } A_{\text{exp}} = -5.6 \pm 0.6 \pm 0.2 \text{ ppm}$$

$$\text{II: } A_{\text{exp}} = -1.36 \pm 0.29 \pm 0.13 \text{ ppm}$$

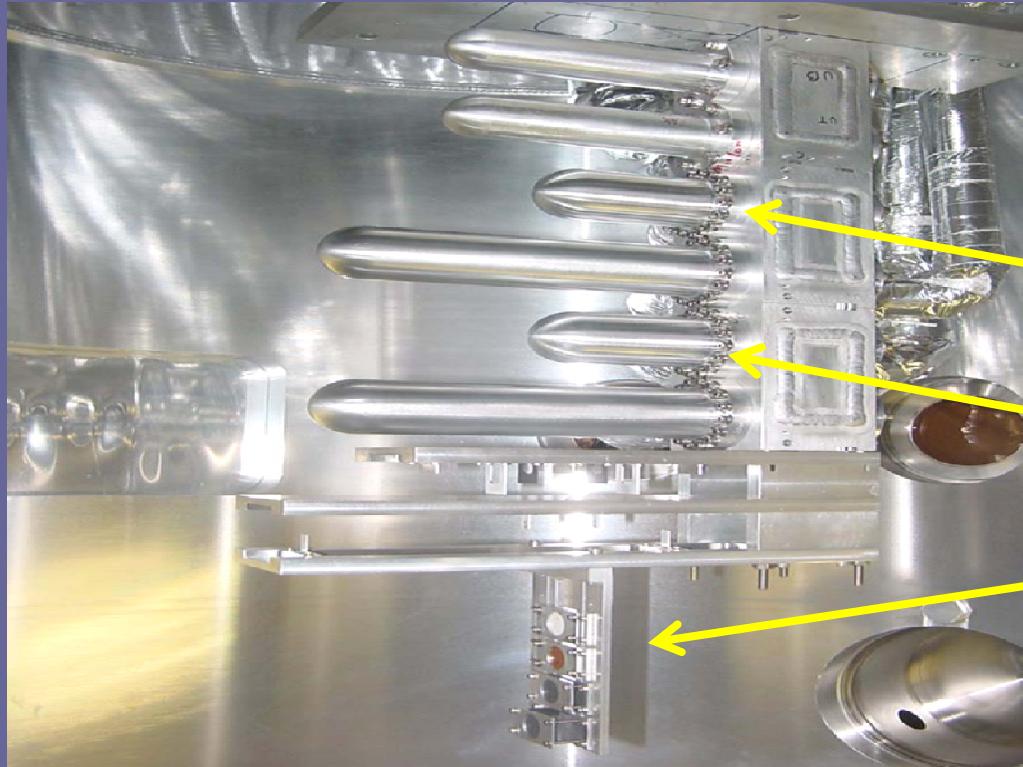
F. Maas et al.,  
PRL 93 (2004) 022002,  
PRL 94 (2005) 152001

phase III:  
backward angles, higher E

$$\text{III: } A_{\text{exp}} = -17.23 \pm 0.82 \pm 0.89 \text{ ppm}$$



# Cryogenic Targets



Hall C targets

liquid Hydrogen

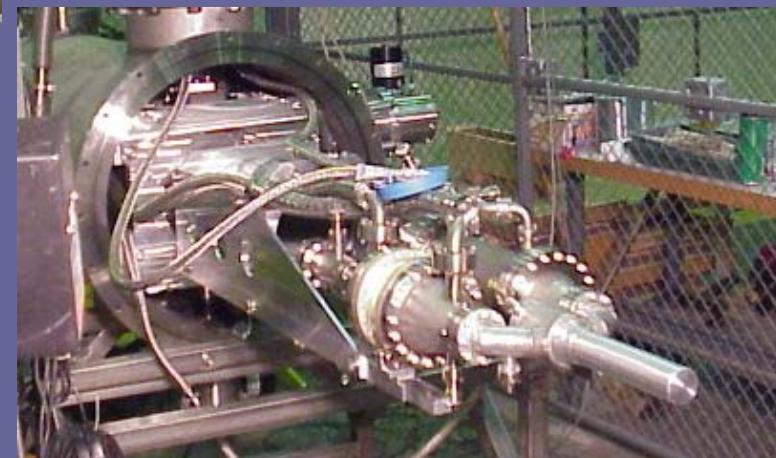
liquid deuterium

solid targets

G0 target

$$\begin{aligned}\Delta E &= (4.5 \text{ MeV} \cdot \text{cm}^2/\text{g}) \\ &\times (0.07 \text{ g/cc}) \times (20 \text{ cm}) \\ &= 6 \text{ MeV}\end{aligned}$$

$$\Delta P = (10 \text{ MeV}) \times (100 \mu\text{A}) = 1 \text{ kW}$$



# Summary of data at $Q^2 = 0.1 \text{ GeV}^2$

Solid ellipse:

K. Pashke, private comm, [same as J. Liu, et al PRC 76, 025202 (2007)], uses theoretical constraints on the axial form factor

Dashed ellipse:

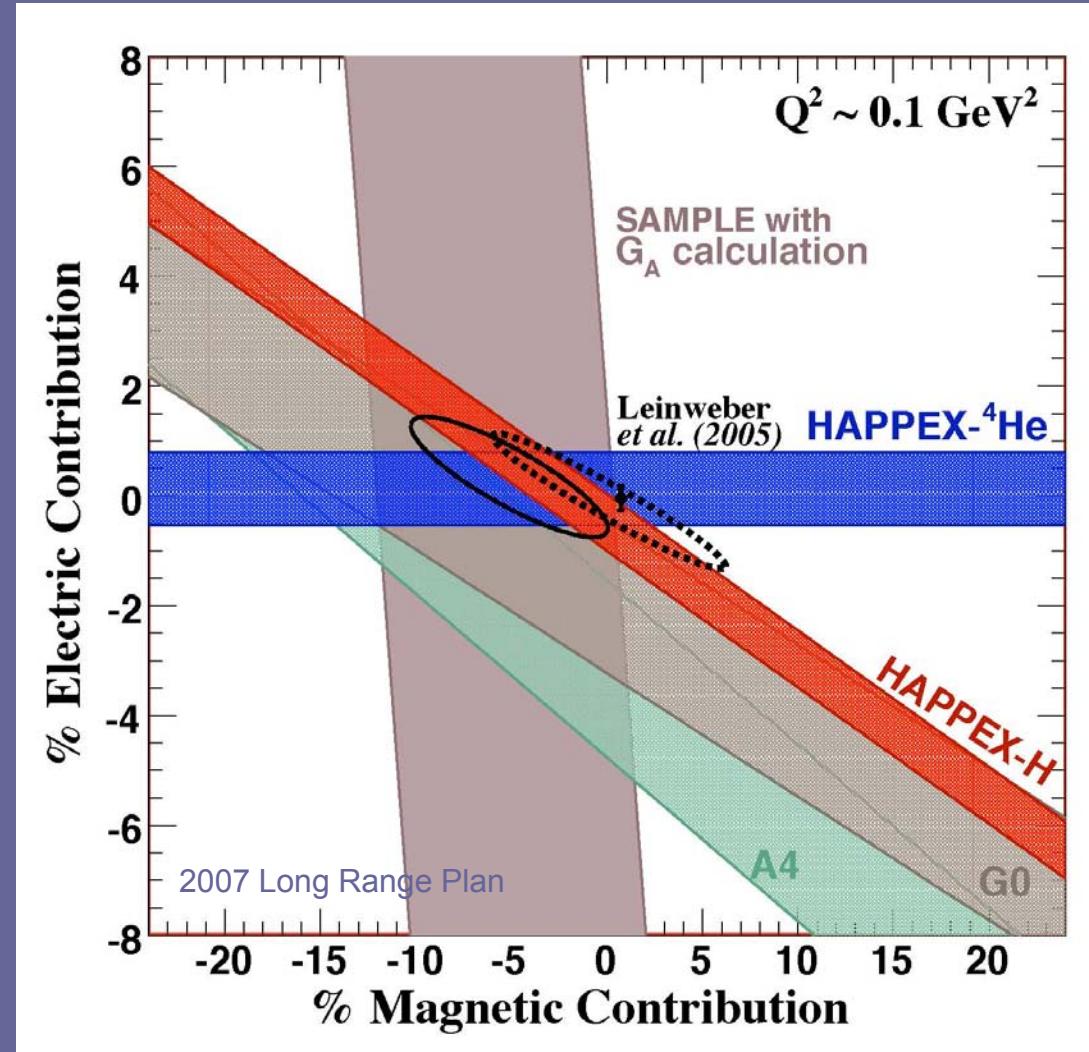
R. Young ,et al. PRL 97 (2006) 102002, does not constrain  $G_A$

Placement of SAMPLE band on the graph depends on choice for  $G_A$

at  $Q^2=0.23 \text{ GeV}^2$

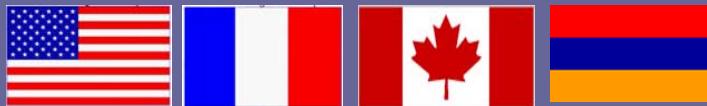
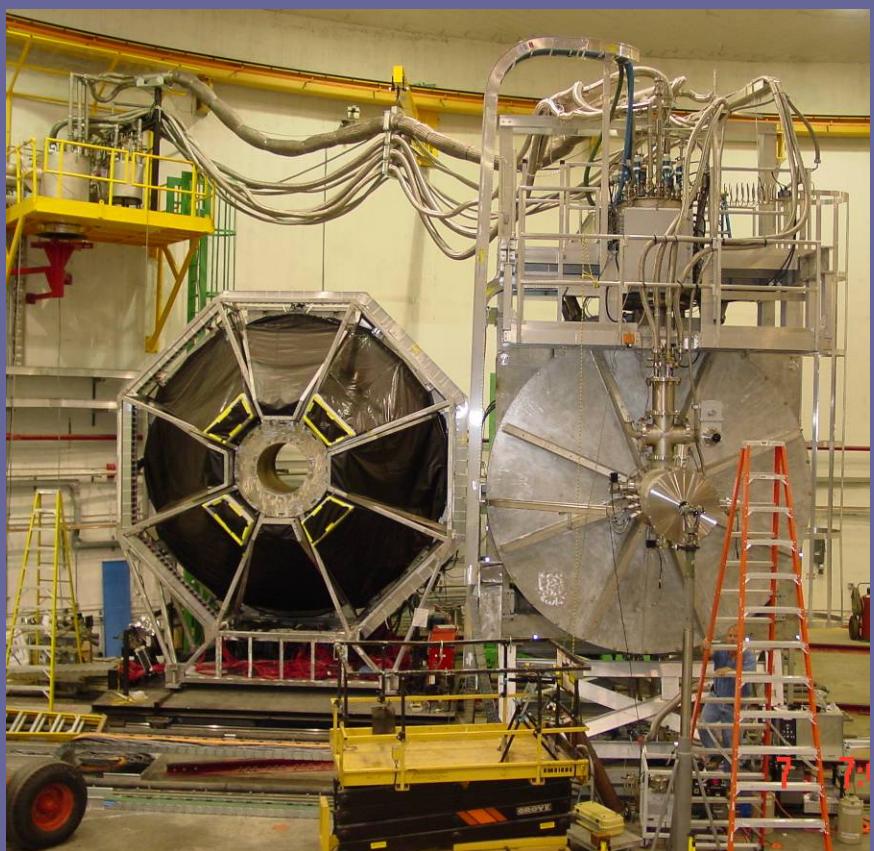
% contribution to proton:  
electric:  $-3.0 \pm 2.5 \%$   
magnetic:  $+2.9 \pm 3.2 \%$

Similar to  $Q^2=0.1 \text{ GeV}^2$

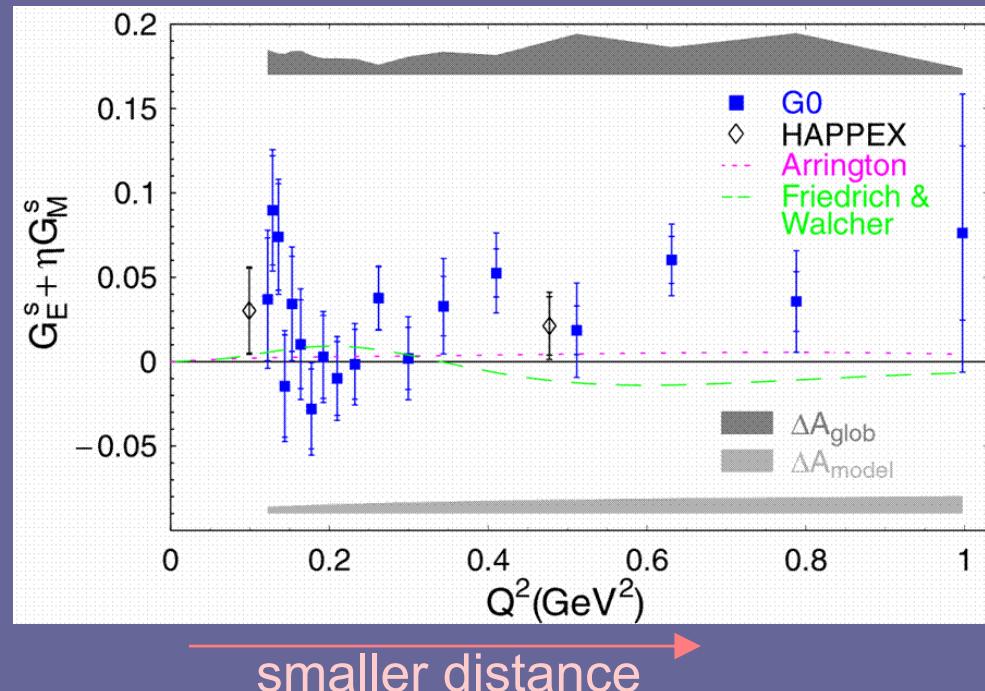


$$\% \text{ contrib} = \frac{G_{E,M}^s}{G_{E,M}^p} \times \left( -\frac{1}{3} \right) \times 100$$

# The “G0” experiment at JLab



D.S. Armstrong et al, PRL 95 (2005) 092001



- H and D targets, wide range of distance scales
- Independently determine neutral weak charge, magnetism and axial form factors

# G0 Backward Angle (at Jlab)

Electron detection:  $\theta = 108^\circ$ , H and D targets

Add Cryostat Exit Detectors (CED) to define electron trajectory

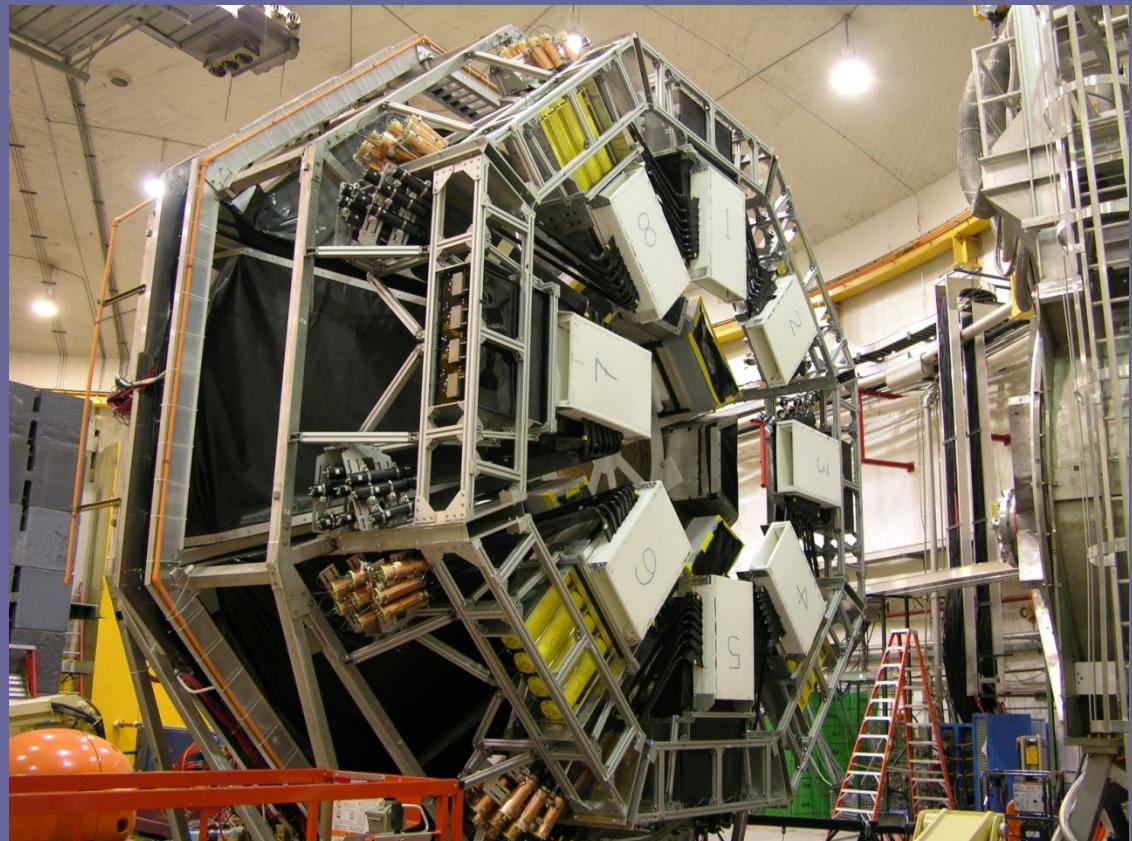
Aerogel Cerenkov detector for  $\pi/e$  separation ( $p_\pi < 380$  MeV/c)

1 scaler per channel FPD/CED pair (w/ and wo/ CER)

$E_e$ (MeV)	$Q^2$ (GeV $^2$ )
362	0.23
687	0.62

Both H and D  
at each kinematic setting

Common  $Q^2$  with  
HAPPEX-III and PVA4



# G0 fun facts

Run start to run end ~ 8940 hours, 330 C of beam, 3.5 Tb of data

$$A = \frac{Y_1 + Y_4 - Y_2 - Y_3}{Y_1 + Y_2 + Y_3 + Y_4}; \quad 1,4 = "+" \quad 2,3 = "-"$$



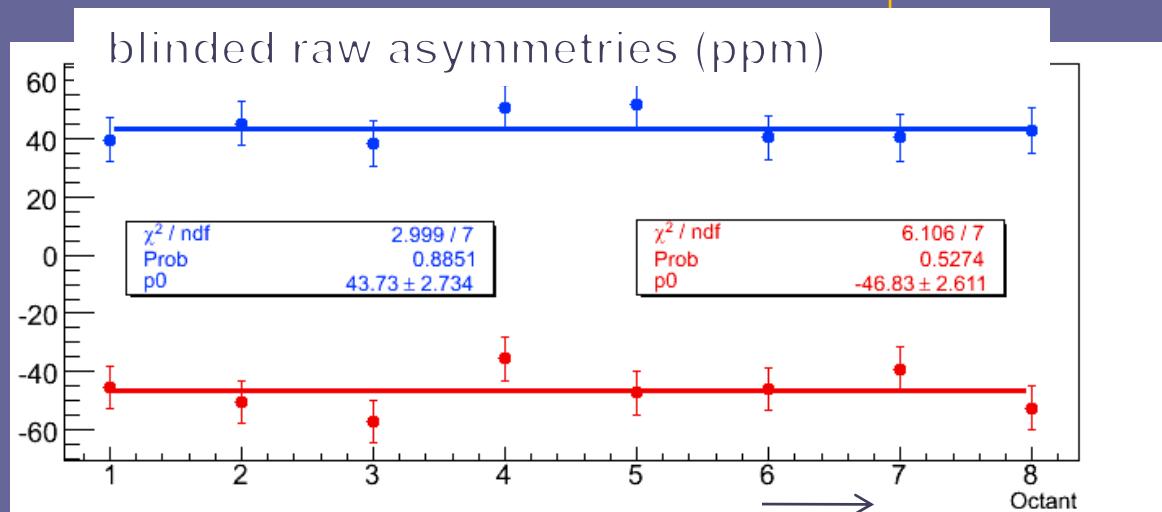
measured 2,000 asymmetries 15 times / second

In 2400 hours, we measured  
250 billion asymmetries  
(or, each asymmetry is measured 13  
million times)

3.5 Tb  $\sim 5.5 \times 10^8$  inches of 6250 bpi tape  
 $\sim 14,000$  km

distance from Newport News, VA to  
Katmandu, Nepal

$\sim 14,000$  data tapes



H-687 MeV data sample

# Precision tests of TeV-scale Electroweak Physics

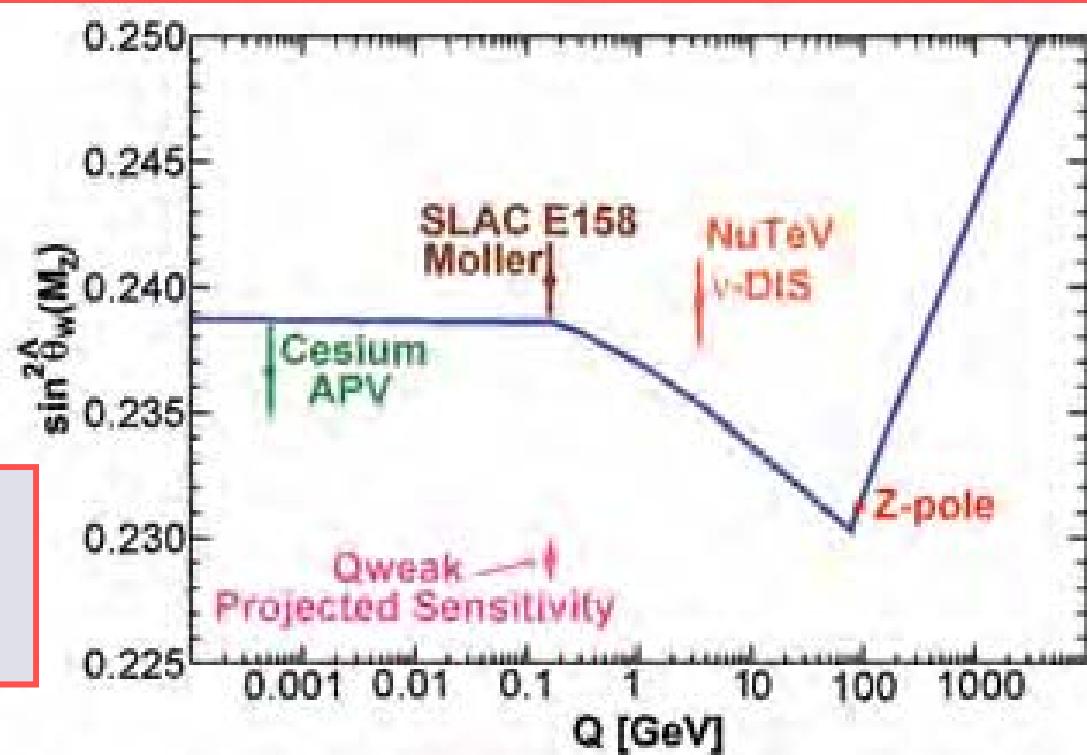
Indirect searches for extensions to Standard Model of particle physics using precision measurements of  $\sin^2\theta_W$  at several energy scales are complementary to experiments at the LHC.

Case 1: electron-electron scattering (Moller) → SLAC E158

$A_{\text{meas}} = -131 \pm 14 \pm 10$   
parts per billion

P. Anthony et al.,  
PRL 95 (2005) 081601

Limits existence of new  
e-e interactions on scale  
of 7-16 TeV

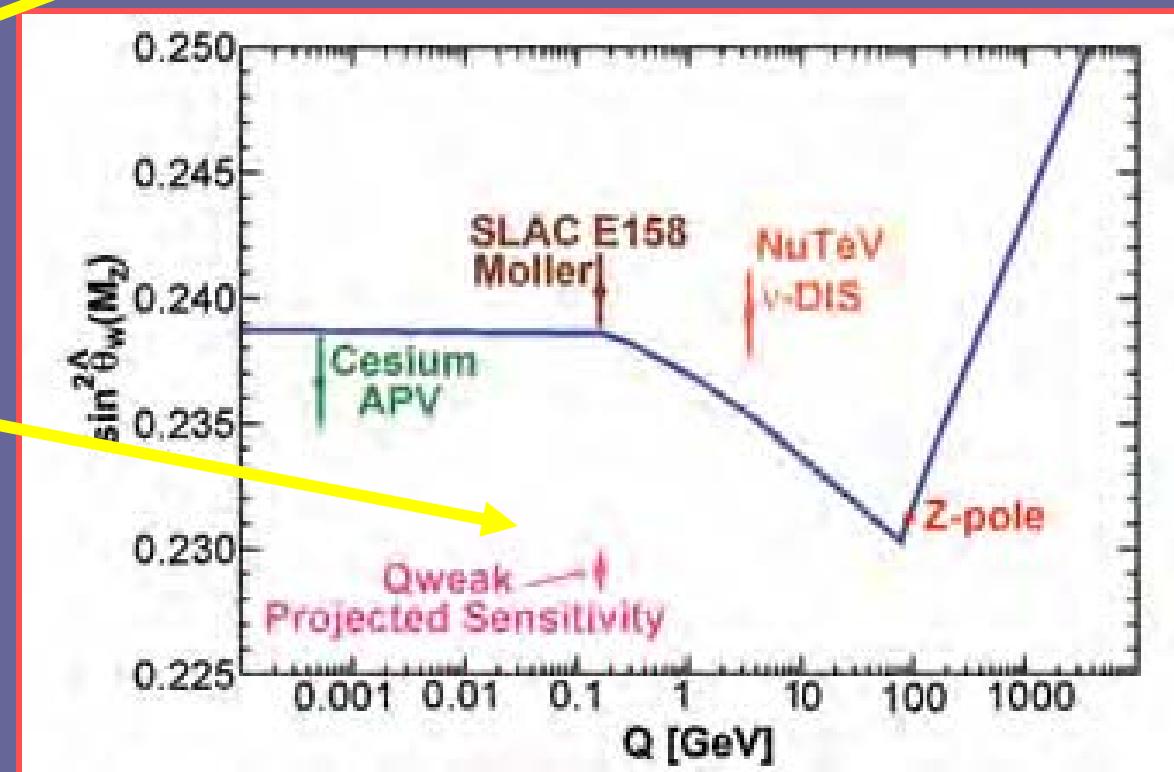


# Precision tests of TeV-scale Electroweak Physics

## Case 2: e-p scattering: Weak Charge of the Proton

$$L = L_{SM}^{PV} + L_{NEW}^{PV} = -\frac{G_F}{\sqrt{2}} \bar{e} \gamma_\mu \gamma_5 e \sum_q C_{1q} \bar{q} \gamma^\mu q + \boxed{\frac{g^2}{4\Lambda^2} \bar{e} \gamma_\mu \gamma_5 e \sum_q h_V^q \bar{q} \gamma^\mu q}$$

precise measurement  
of known e-q interactions  
( $C_{1u}$ ,  $C_{1d}$ ) can place limits  
on existence of new  
ones on scale of a few  
TeV



# Weak charge of the nucleon

As  $Q^2 \rightarrow 0$ , for far forward angles,

$$A_{LR} = \frac{-G_\mu}{4\pi\alpha\sqrt{2}} \left[ (1 - 4 \sin^2 \theta_W) + \frac{-\varepsilon G_E^{Y,p} (G_E^{Y,n} + G_E^s) - \tau G_M^{Y,p} (G_M^{Y,n} + G_M^s) + A_A}{\varepsilon (G_E^{Y,p})^2 + \tau (G_M^{Y,p})^2} \right]$$

$$A_{LR} = \frac{-G_\mu}{4\pi\alpha\sqrt{2}} \left[ Q_{weak}^p Q^2 + B_4 Q^4 + \dots \right]$$

$$Q_{weak}^p = 1 - 4 \sin^2 \theta_W$$

$$Q_{weak}^p = -2 \left( 2 C_{1u} + C_{1d} \right)$$

e axial-vector, q vector couplings

# QWeak: Parity violating e-p scattering

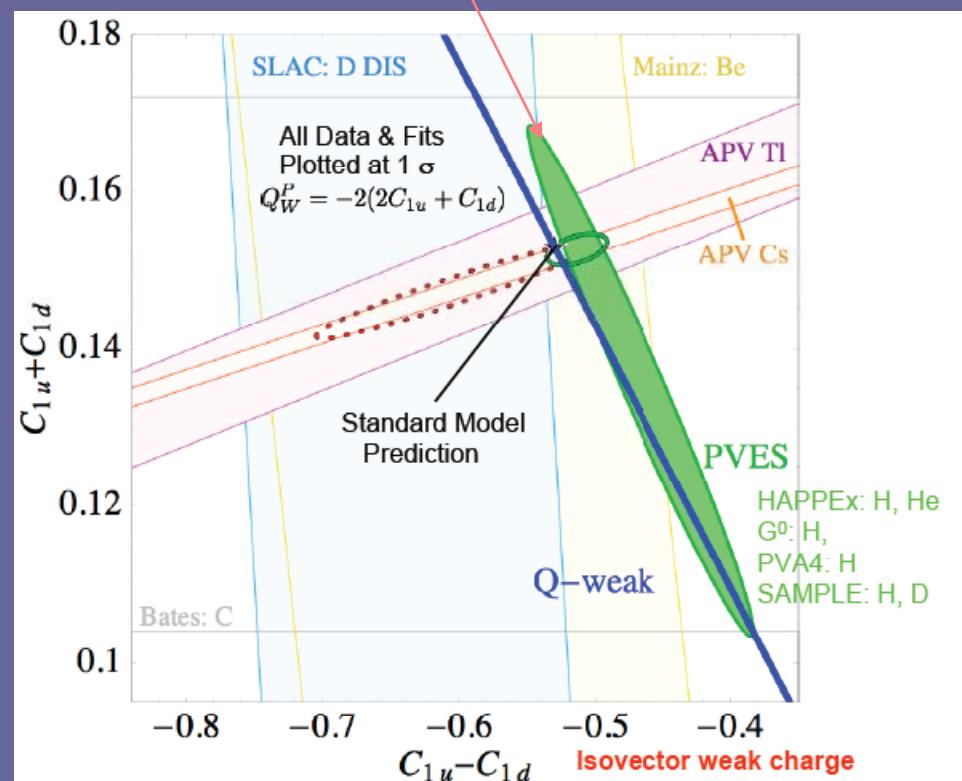
R. Carlini, et al

QWeak toroidal magnet at MIT-Bates lab



expected asymmetry  $\sim 0.3$  ppm  
measure to few %.

combined existing PV electron  
scattering data provide most precise  
experimental limit on proton's  
weak charges

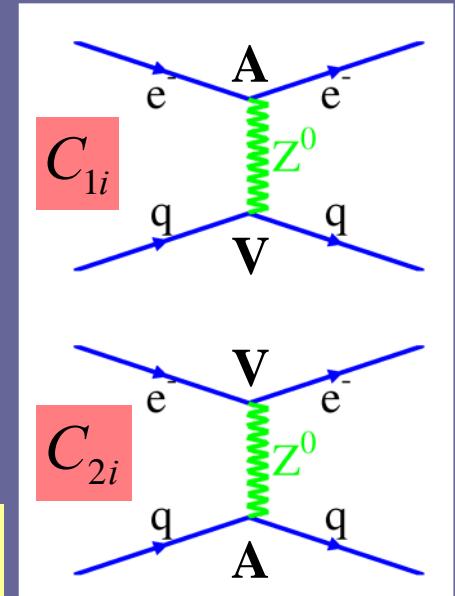
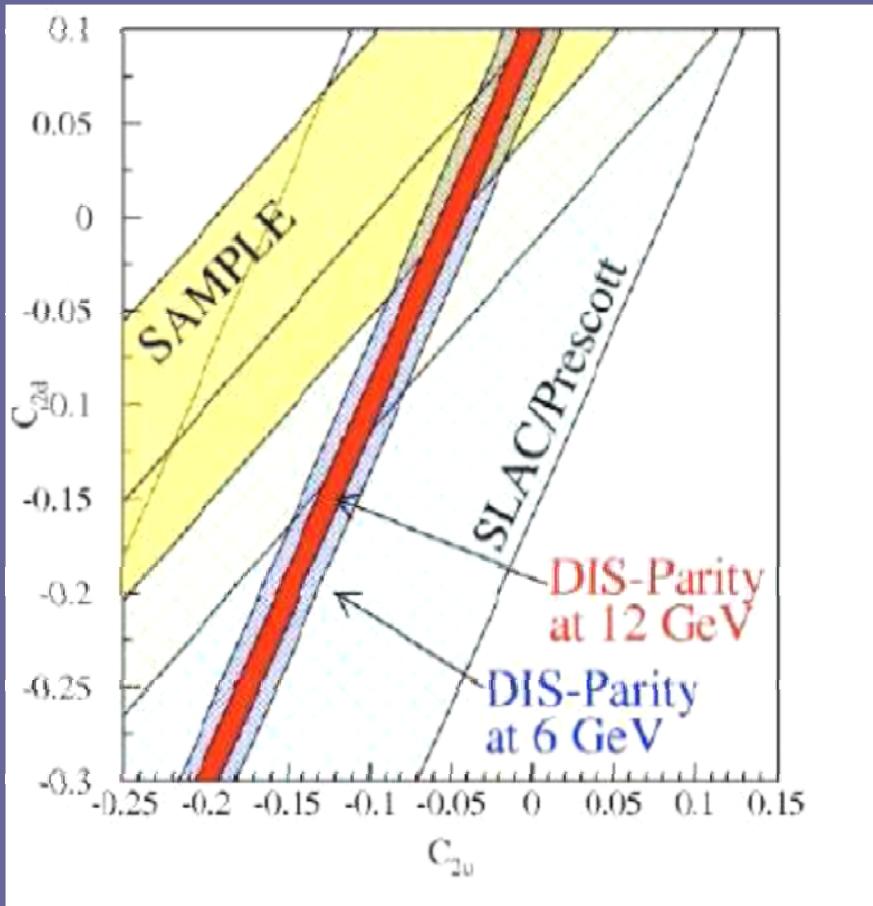


R. Young et al., PRL 99 (2007) 122003

# Precision tests of TeV-scale Electroweak Physics

## Case 3: Direct electron-quark scattering

high precision return to SLAC 1978 measurement



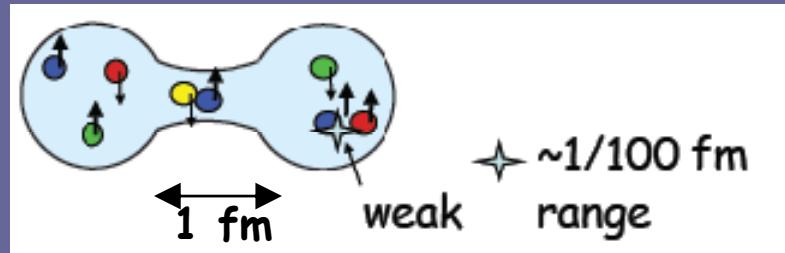
$$y = \frac{E - E'}{E}$$

$$A_{PV} = \frac{G_F Q^2}{\sqrt{2} \pi \alpha} [a(x) + f(y)b(x)]$$

$$b(x) \rightarrow [2C_{2u} - C_{2d}]$$

w/ 6 GeV, strong interaction important

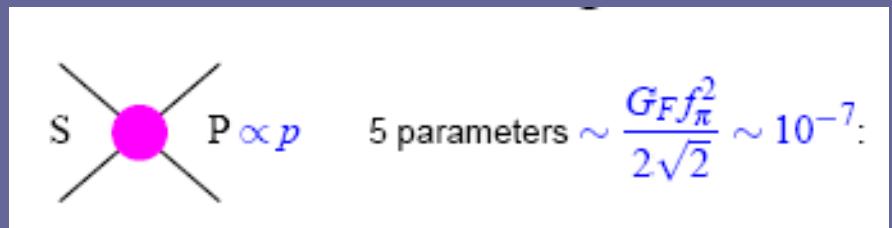
# Weak interactions between nucleons



from P. Mumm, NIST

Fundamental interaction is q-q weak interactions, but these are buried inside strongly bound systems

→ typical size of effect  $< 10^{-7}$   
5 leading contributions



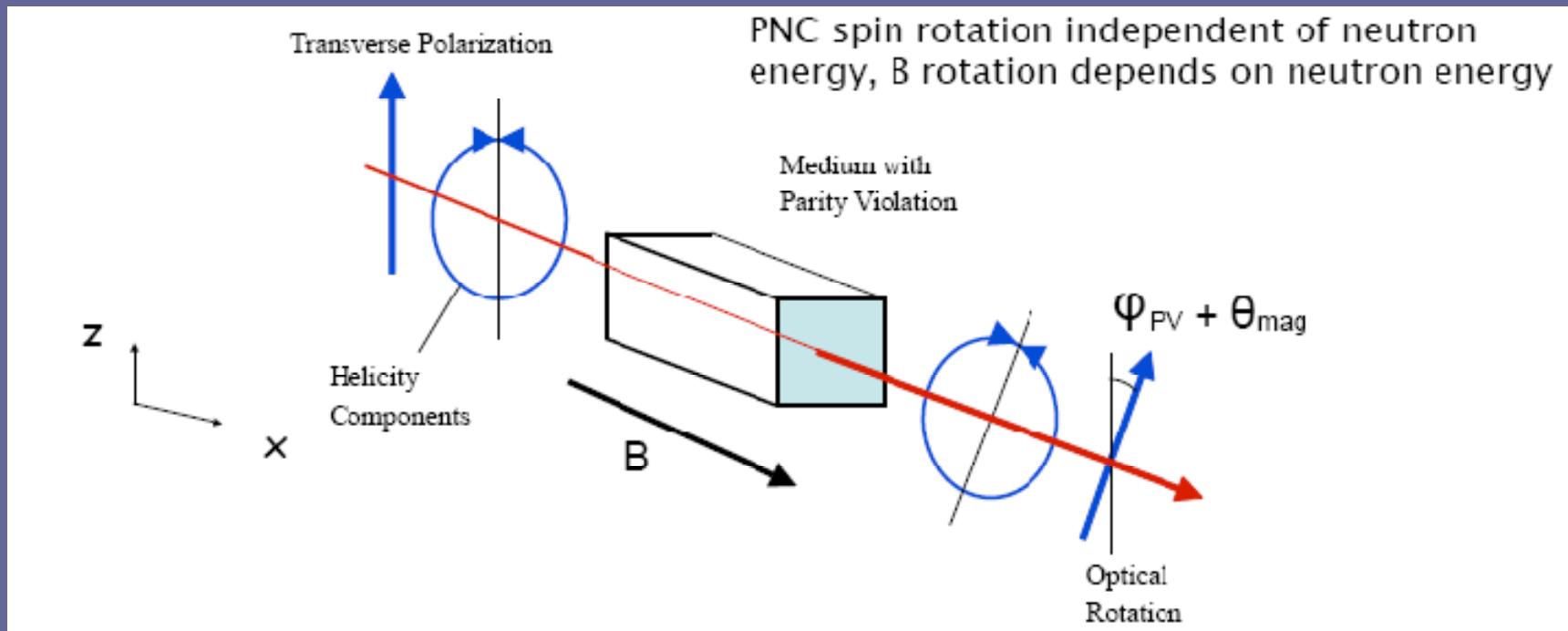
see H. Giesshammer, NNPDF 2008

- need a minimum of 5 experiments (more is better, and preferably without too many nucleons)
- need a theoretical framework to interpret the results → chiral EFT

# Required Experimental program

- Asymmetry in  $\vec{p}$ -p and  $\vec{p}$ - ${}^4\text{He}$  scattering  
p-p measurement from TRIUMF
- spatial  $\gamma$  asymmetry in  $\vec{n} + p \rightarrow d + \gamma$  and  $\vec{n} + d \rightarrow t + \gamma$   
 $n + p \rightarrow d + \gamma$  underway, moved from LANSCE to SNS
- circular pol'n of  $\gamma$  in  $n + p \rightarrow d + \gamma$
- spin rotation of  $\vec{n}$  in  ${}^4\text{He}$  and  $\vec{n}$  in p  
recently completed at NIST
- Also: laser spectroscopy on heavy rare isotopes (e.g., Francium)  
installation in progress at TRIUMF (L. Orozco, et al.)

# neutron spin rotation in ${}^4\text{He}$ (at NIST, Gaithersburg, MD)



$\phi_{PV}$  from  $\sigma \cdot p$  term in forward scattering amplitude (P odd)

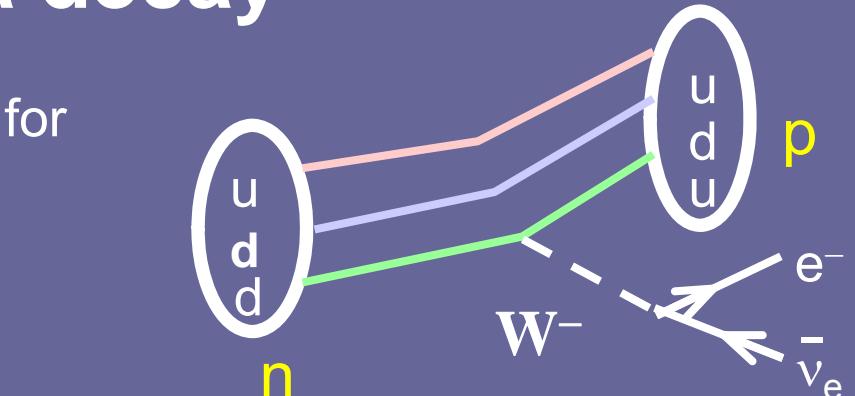
$$\frac{d\phi}{dz} \sim 10^{-7} \text{ rad/m}$$

# Polarized neutron beta decay

$G_F$  from  $\mu$ -decay, apparently the same for quarks and leptons

$$g_V^2 \propto (G_F V_{ud})^2$$

define  $\lambda = \frac{g_A}{g_V} = 1.2695 \pm 0.0029$



$$n \rightarrow p + e^- + \bar{\nu}_e + 782 \text{ keV}$$

decay rate

$$dW \propto (g_V^2 + 3g_A^2) F(E_e) \left[ 1 + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + b \frac{m}{E_e} + \vec{\sigma}_n \cdot \left( A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} - D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right) \right]$$



neutron lifetime  
measures this piece

measurable and non-zero

T odd, P even

a, b, B, and A each depend differently on  $\lambda$

# “cold” and “ultracold” neutrons

“cold”:

$$\text{temperature} = \frac{K.E.}{\frac{1}{40} eV} \times 300K = 36K$$

$$\beta = \sqrt{\frac{2 \times K.E.}{m_n}} \rightarrow v = 760 \text{ m/s}$$

$$\lambda = \frac{h}{p} = 0.5 \text{ nm}$$

“ultra-cold”:

$$v = 5 \text{ m/s}$$

$$T \sim \text{few mK}$$

$$\lambda \sim 50 \text{ nm}$$

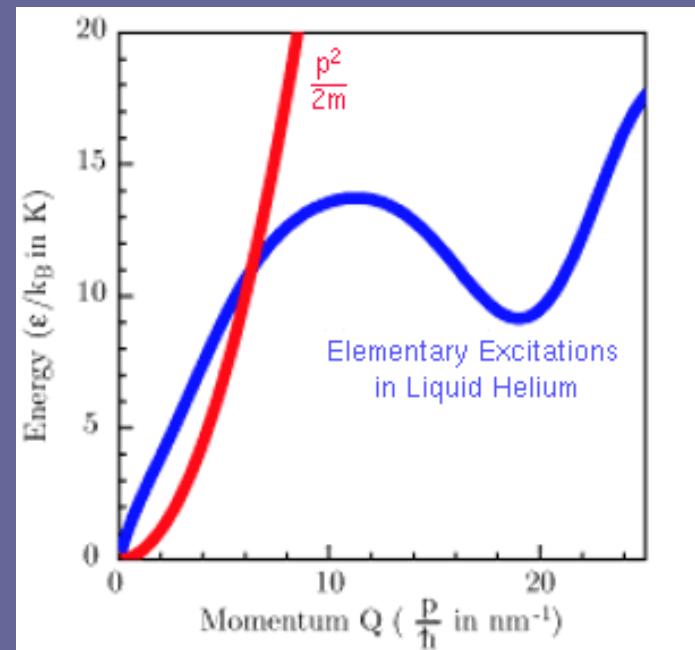
Scatter cold neutrons from solids, liquids, etc., that can absorb bulk of neutron KE into a “phonon” resonance

Can be trapped:

- material walls
- magnetic fields
- gravity!

Produce by moderating thermal neutrons from reactor or spallation source

Can be transported long distances by total internal reflection



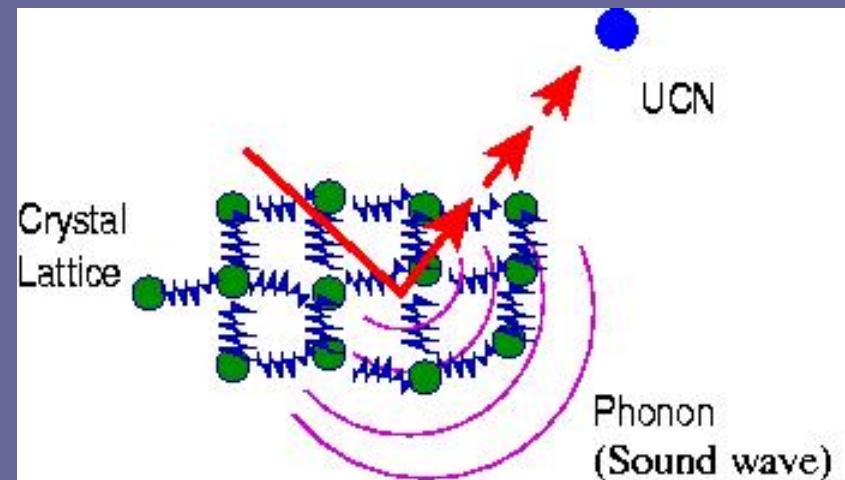
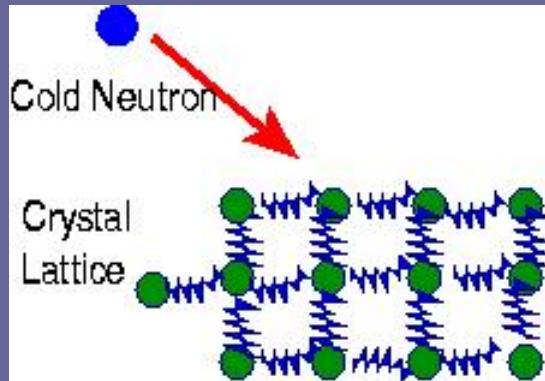
Golub & Pendelbury, 1977

# Ultracold Neutrons: Superthermal Process

slide from C.-Y. Liu, Indiana Univ.

R. Golub and J. M. Pendlebury, Phys. Lett, A53, 133 (1975)

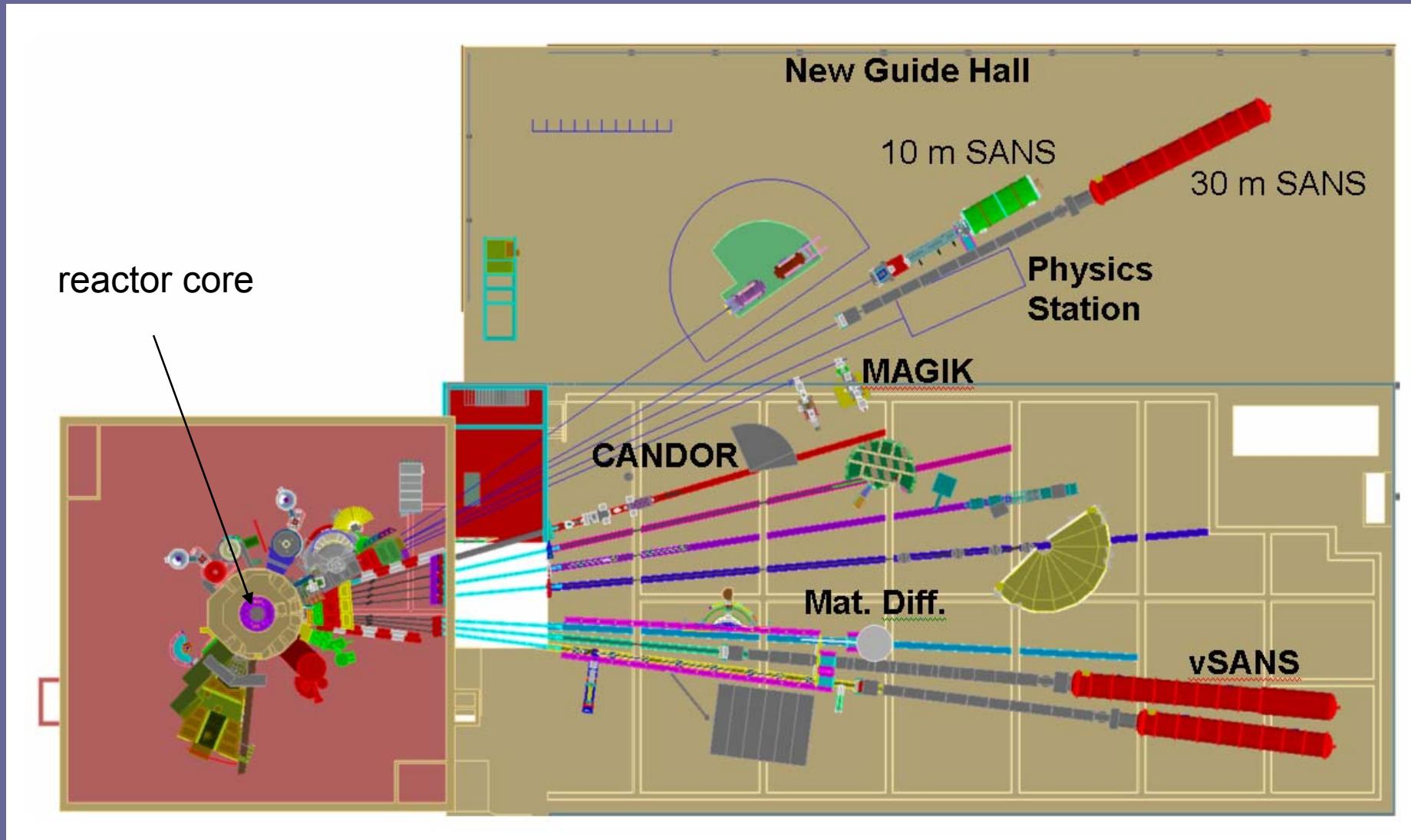
- Cold neutrons downscatter in the solid, giving up almost all their energy, becoming UCN.



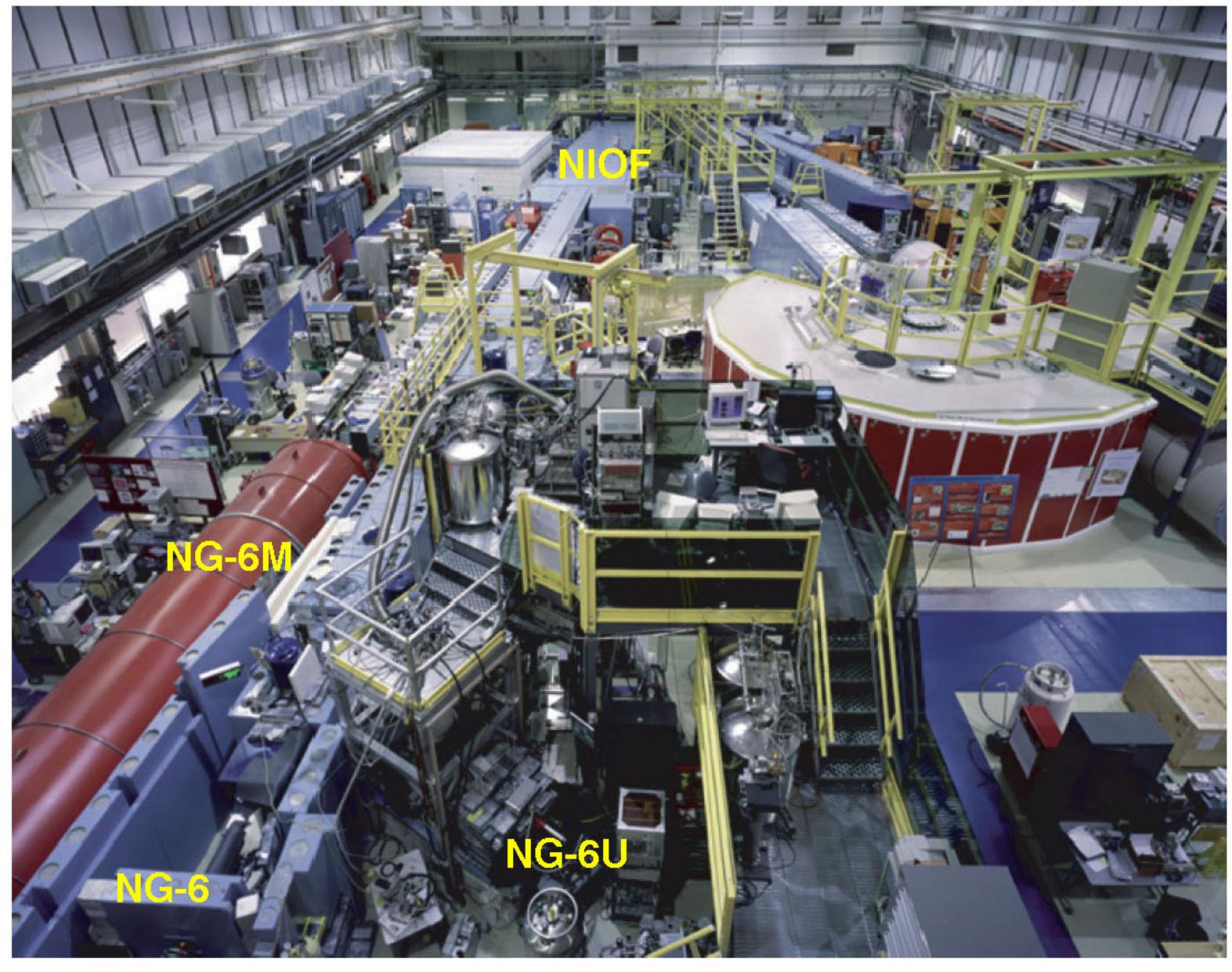
- UCN upscattering (the reverse process) is suppressed by cooling the moderator to low temperatures.

# NIST Center for Neutron Research

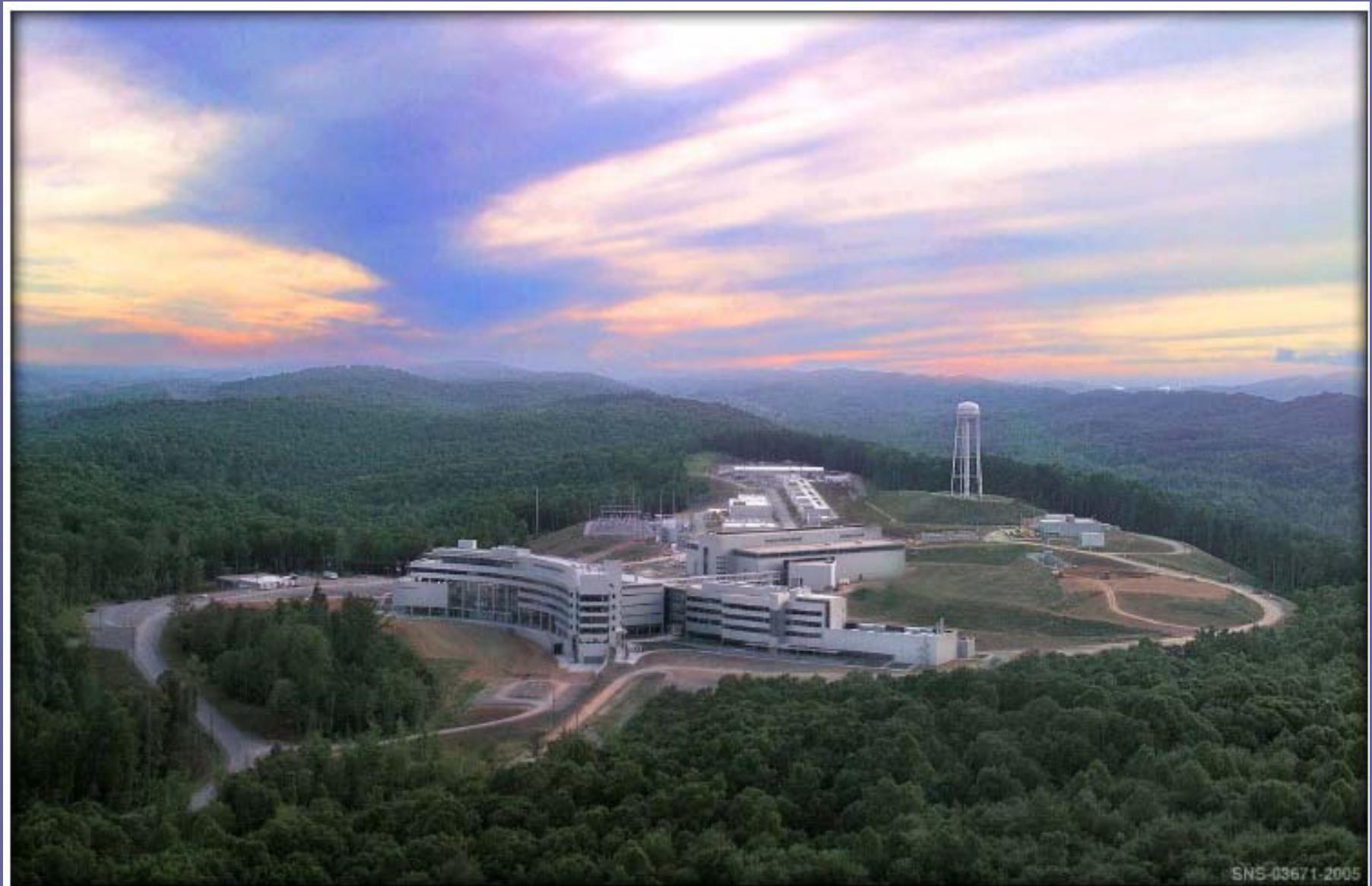
20 MW reactor is source of “cold” neutrons:  $T_n \sim 10^{-5}$ - $10^{-2}$  eV



# NIST Center for Cold Neutron Research



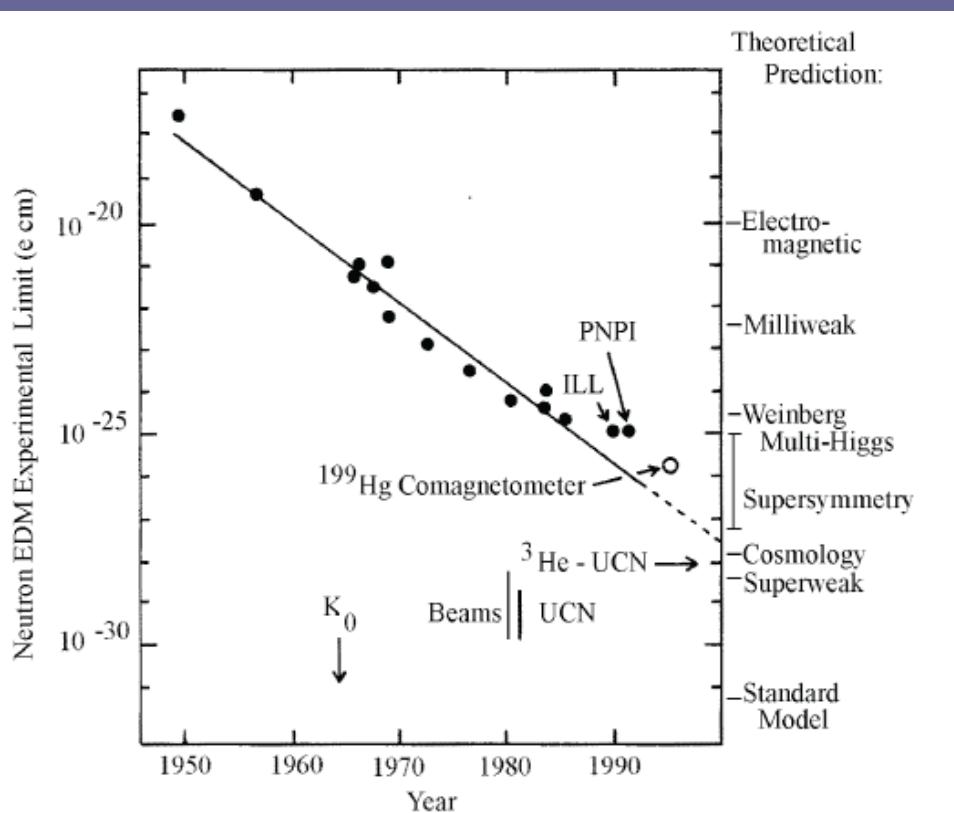
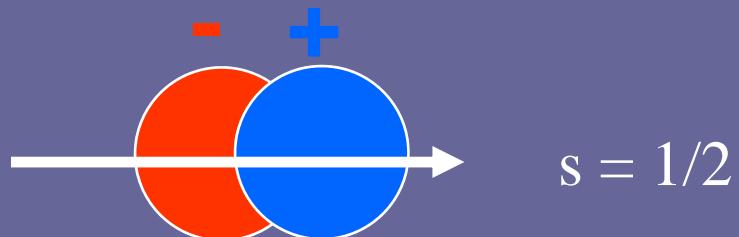
# The Spallation Neutron Source



# T-odd $\rightarrow$ CP odd: neutron EDM

CP-violation is important to understanding of origin of matter (vs antimatter)

$$\vec{d} \cdot \vec{E}$$



If neutron possesses an electric dipole moment, then, in an electric field:

$$H_E = -d_n \vec{\sigma} \cdot \vec{E}$$

will reverse sign under reversal of E.  
This is CP-odd observable.

many EDM experiments underway or proposed, in many systems.

New U.S. experiment planned using ultracold neutrons at SNS –Oak Ridge

<http://p25ext.lanl.gov/edm/edm.html>

# Origin of Matter

from P. Huffman, NSCU

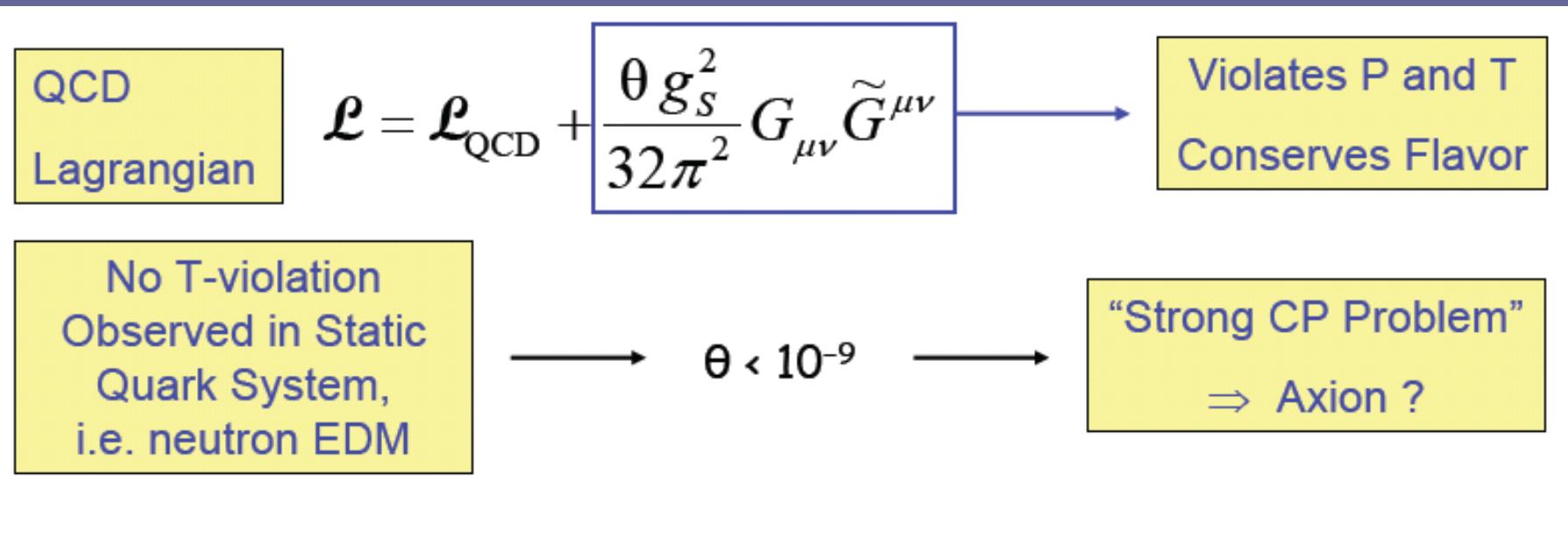


- **Sakharov conditions:** [JETP 5 (1967) 24]
  - A violation of baryon number
  - A violation of both C and CP
  - A departure from thermal equilibrium
- Could occur in several points during BBN
  - $T \sim 10^{29}$  K (GUT scale)
  - $10^{29}$  K >> T >>  $10^{15}$  K (leptogenesis)
  - $T \sim 10^{15}$  K (Electroweak)
- One can test Electroweak Baryogenesis (EWB) and parts of leptogenesis experimentally



A. Sakharov, 1943  
(from Wikipedia)

# neutron EDM and CP in the strong interaction



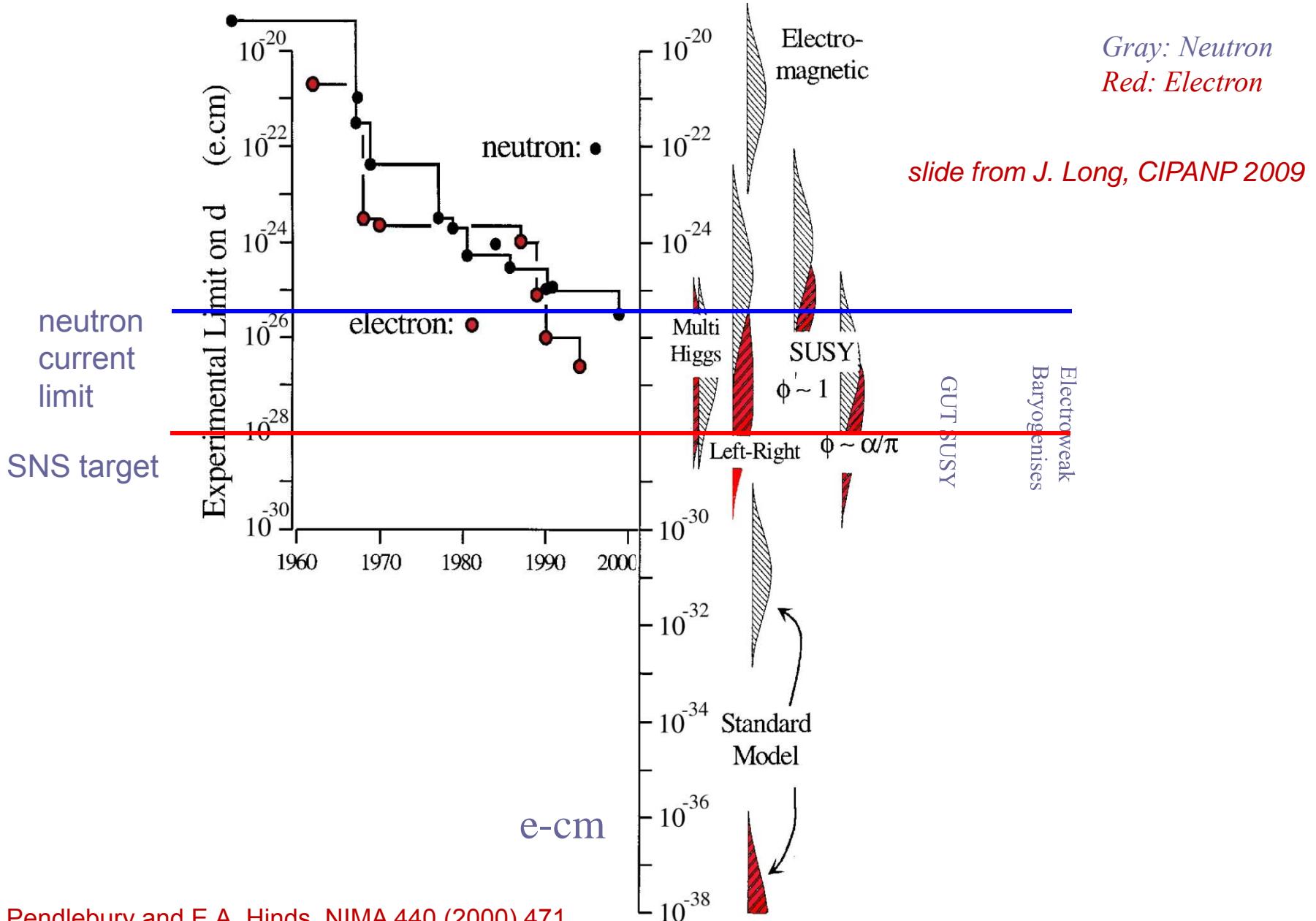
## CP Conservation in the Presence of Pseudoparticles\*

R. D. Peccei and Helen R. Quinn†

*Institute of Theoretical Physics, Department of Physics, Stanford University, Stanford, California 94305  
(Received 31 March 1977)*

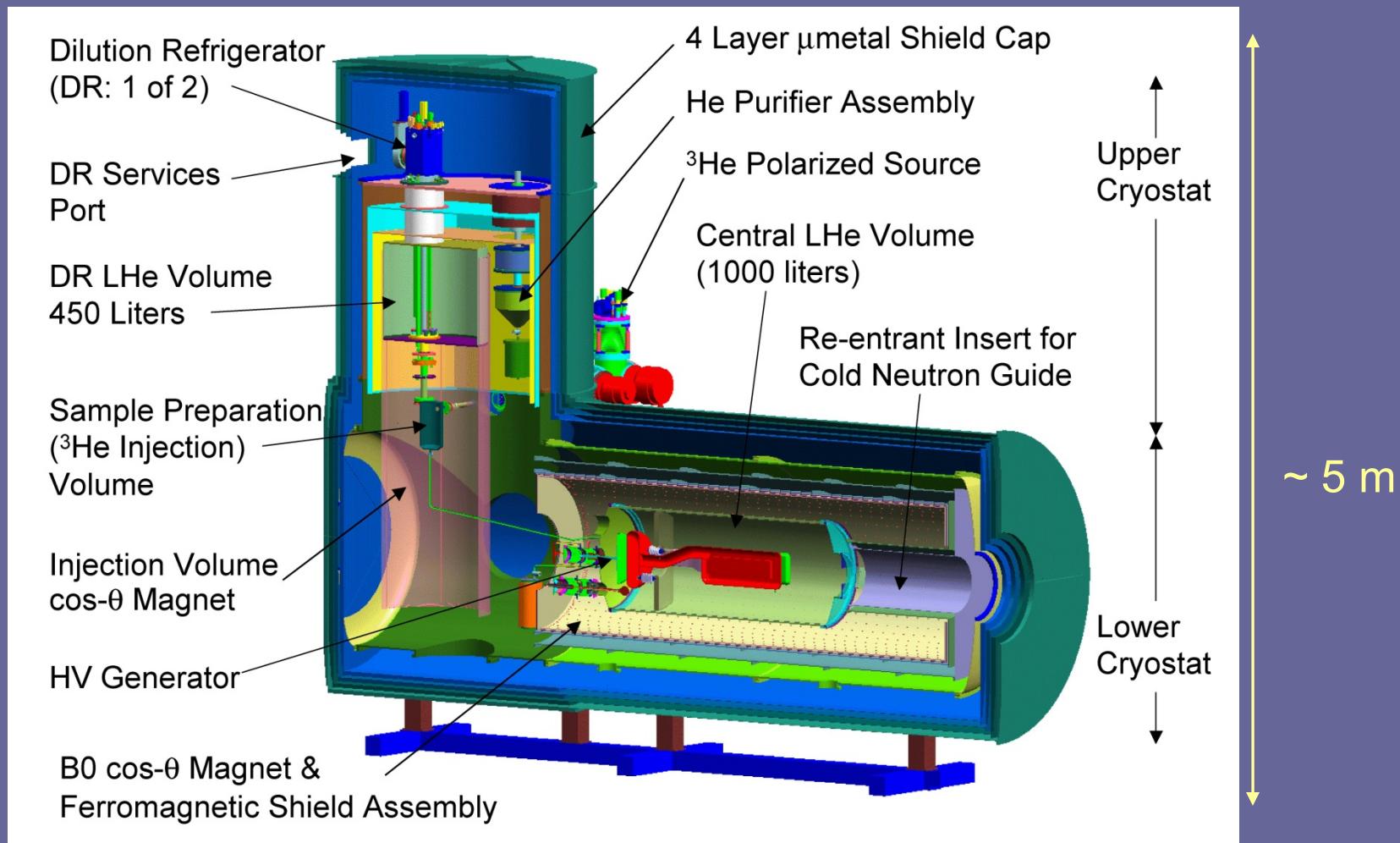
We give an explanation of the *CP* conservation of strong interactions which includes the effects of pseudoparticles. We find it is a natural result for any theory where at least one flavor of fermion acquires its mass through a Yukawa coupling to a scalar field which has nonvanishing vacuum expectation value.

# Model Sensitivity to EDM



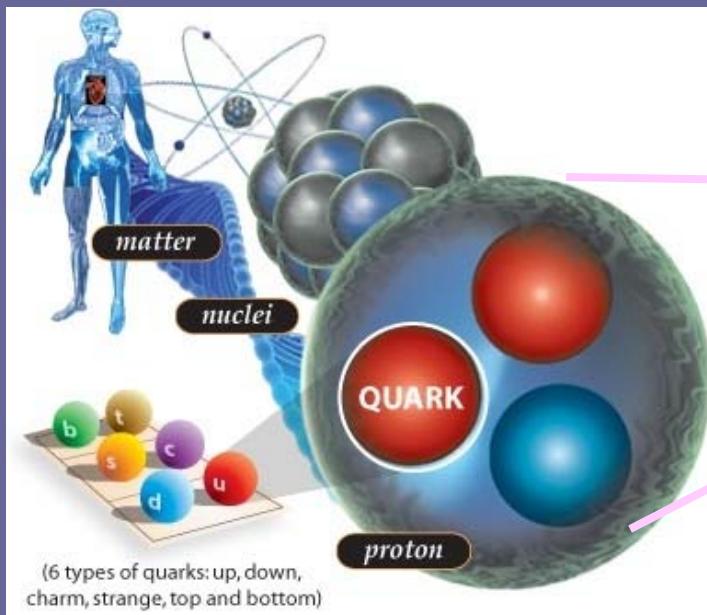
# proposed n-EDM experiment at the SNS

slide from J. Long, CIPANP 2009

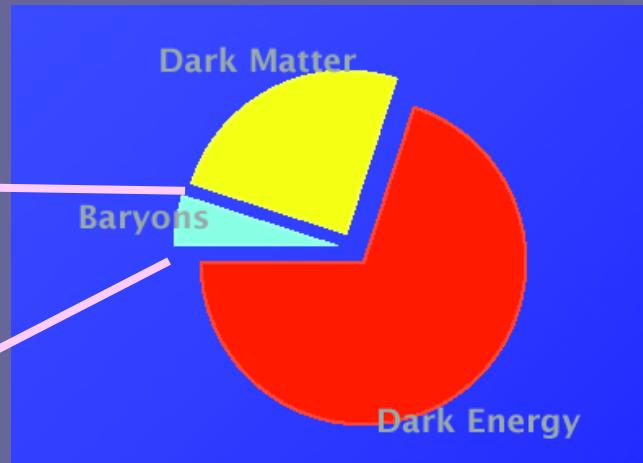


*stay tuned → 2016....*

# Thanks for your attention



picture courtesy of Jefferson Lab



*and enjoy the NSCL tour....*

