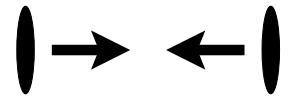


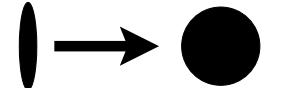
QCD at High Temperature (Experiment)

Kinematics

CMS: $s = (p_1 + p_2)^2 = 4E_{CM}^2$



Lab: $p_1 = (m, 0)$ $p_2 = (E_L, p_z) = (E_L, \sqrt{E_L^2 - m^2})$



$$s = (m + E_L)^2 - (E_L^2 - m^2) = 2m(E_L + m) \quad E_{CM} = \sqrt{mE_L/2}$$

SPS : 200 GeV (LAB)

$E_{CM} = 10$ GeV $\gamma = 10$

RHIC : 100 GeV (CMS)

$E_{CM} = 100$ GeV $\gamma = 100$

LHC : 2.75 TeV (CMS)

$E_{CM} = 2.75$ TeV $\gamma = 2750$

Rapidity:

$$y = \frac{1}{2} \log \left(\frac{E + p_z}{E - p_z} \right)$$

SPS : $\Delta y = 6$ RHIC : $\Delta y = 10.6$ LHC : $\Delta y = 17.3$

Bjorken Expansion

Experimental observation: At high energy ($\Delta y \rightarrow \infty$) rapidity distributions of produced particles (in both pp and AA) are “flat”

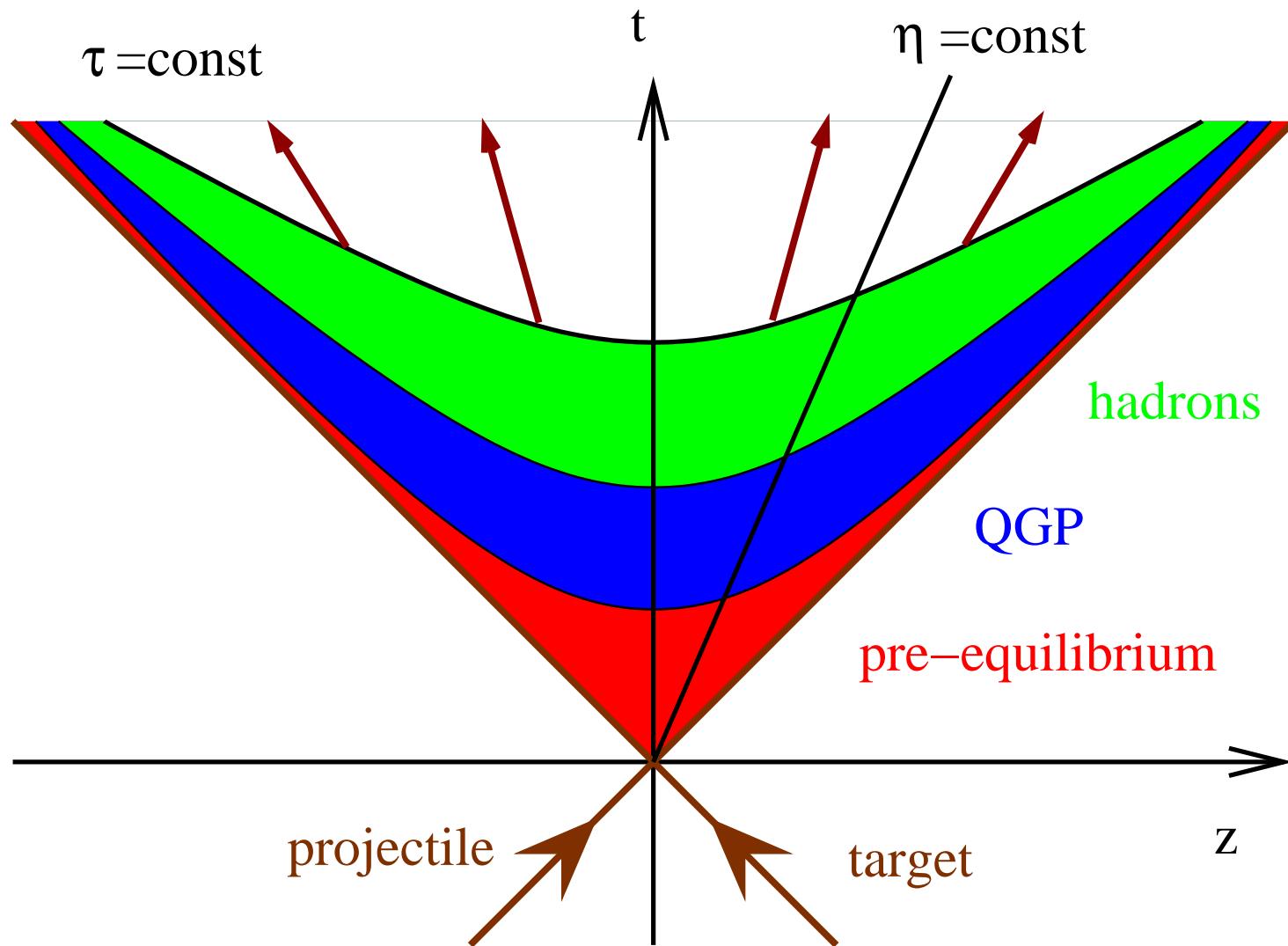
$$\frac{dN}{dy} \simeq \text{const}$$

Physics depends on proper time $\tau = \sqrt{t^2 - z^2}$, not on y

All comoving ($v = z/t$) observers are equivalent

Analogous to Hubble expansion

Bjorken Expansion



Bjorken Expansion: Hydrodynamics

Consider perfect relativistic fluid; 4-velocity $u_\mu = (1, \vec{v})\gamma$

$$T_{\mu\nu} = (\epsilon + P)u_\mu u_\nu - Pg_{\mu\nu}$$

Hydro = Conservation Laws ($\partial^\mu T_{\mu\nu} = 0$) + Equ. of State ($P = P(\epsilon)$)

$$\partial^\mu T_{\mu\nu} = (\partial^\mu \epsilon + \partial^\mu P)u_\mu u_\nu + (\epsilon + P)((\partial^\mu u_\mu)u_\nu + u_\mu \partial^\mu u_\nu) - \partial_\nu P = 0$$

Contract with u_ν , use $u^2 = 1$

$$(\partial^\mu \epsilon + \partial^\mu P)u_\mu + (\epsilon + P)\partial^\mu u_\mu - u^\nu \partial_\nu P = 0$$

$$u_\mu \partial^\mu \epsilon + (\epsilon + P)\partial^\mu u_\mu = 0$$

Thermodynamic relations

$$d\epsilon = Tds \quad \epsilon + P = Ts$$

Hydrodynamic equations

$$u^\mu(T\partial_\mu s) + (Ts)\partial^\mu u_\mu = 0$$

$$\boxed{\partial_\mu (su^\mu) = 0} \quad \text{isentropic expansion}$$

Variables: $t = \tau \cosh \alpha$, $z = \tau \sinh \alpha$. $\Rightarrow u_\mu = (\cosh \alpha, 0, 0, \sinh \alpha)$

$$\partial^\mu (su_\mu) = 0 \quad \Rightarrow \quad \frac{d}{d\tau} [\tau s(\tau)] = 0$$

Solution for ideal Bj hydrodynamics

$$\boxed{s(\tau) = \frac{s_0 \tau_0}{\tau}} \quad T = \frac{const}{\tau^{1/3}}$$

Exact boost invariance, no transverse expansion, no dissipation, . . .

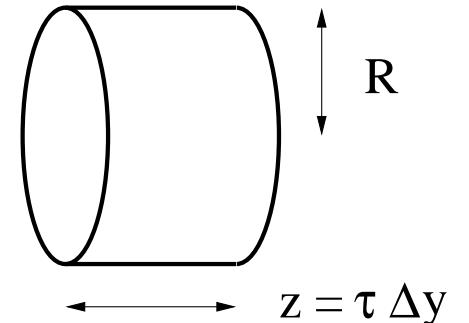
Numerical Estimates

Total entropy in rapidity interval $[y, y + \Delta y]$

$$S = s\pi R^2 z = s\pi R^2 \tau \Delta y = (s_0 \tau_0) \pi R^2 \Delta y$$

$$s_0 \tau_0 = \frac{1}{\pi R^2} \frac{S}{\Delta y}$$

Use $S/N \simeq 3.6$



$$s_0 = \frac{3.6}{\pi R^2 \tau_0} \left(\frac{dN}{dy} \right) \quad \text{Bj estimate}$$

$$\epsilon_0 = \frac{1}{\pi R^2 \tau_0} \left(\frac{dE_T}{dy} \right)$$

Depends on initial time τ_0

RHIC: Au-Au collisions ($\sqrt{s} = 200$ GeV)

$$\frac{dN}{dy} \simeq 998 \quad \tau_0 = 1 \text{ fm} \quad s_0 \simeq 33 \text{ fm}^{-3}$$

Use QGP equation of state $s = 2g\pi^2T^3/45$

$$T_0 \simeq 240 \text{ MeV} \quad \epsilon_0 \simeq (5 - 6) \text{GeV/fm}^3$$

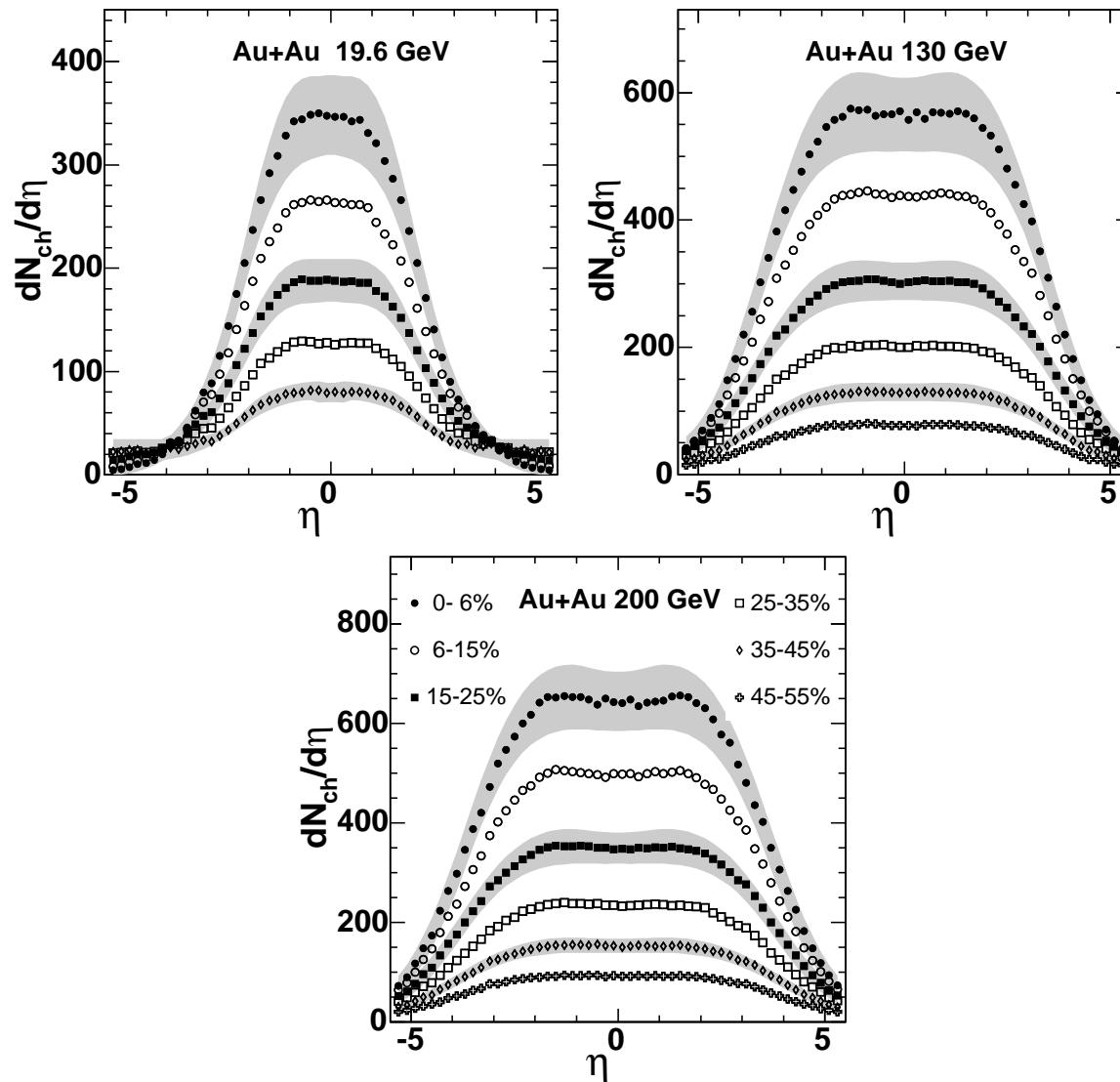
LHC: Factor ~ 2 in multiplicity

$$T_0 \simeq 300 \text{ MeV} \quad \epsilon_0 \simeq 15 \text{GeV/fm}^3$$

BNL and RHIC

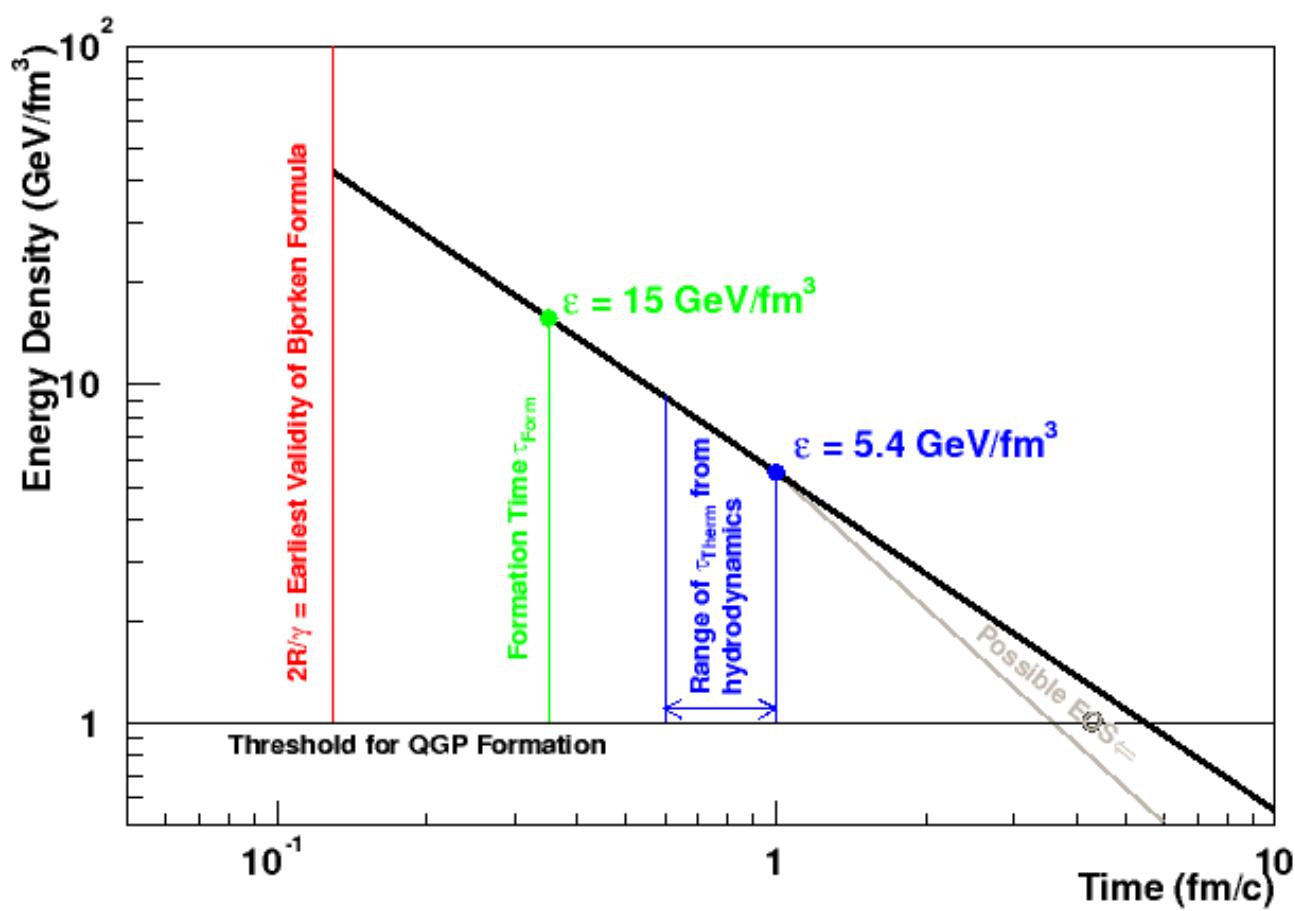


Multiplicities



Phobos White Paper (2005)

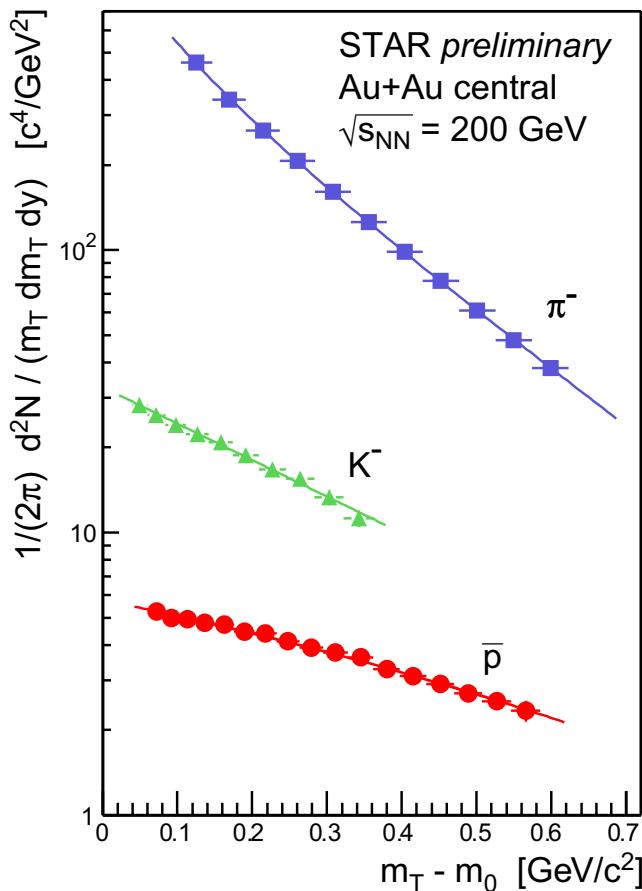
Bjorken Expansion



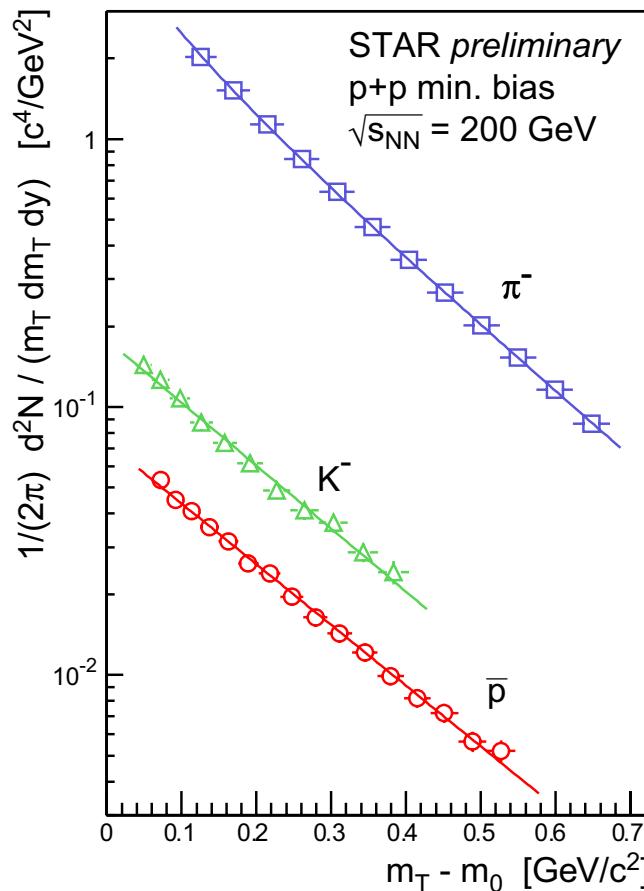
Collective Behavior: Radial Flow

Radial expansion leads to blue-shifted spectra in Au+Au

Au+Au



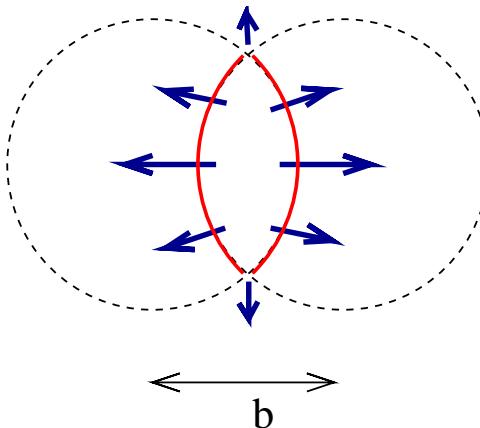
p+p



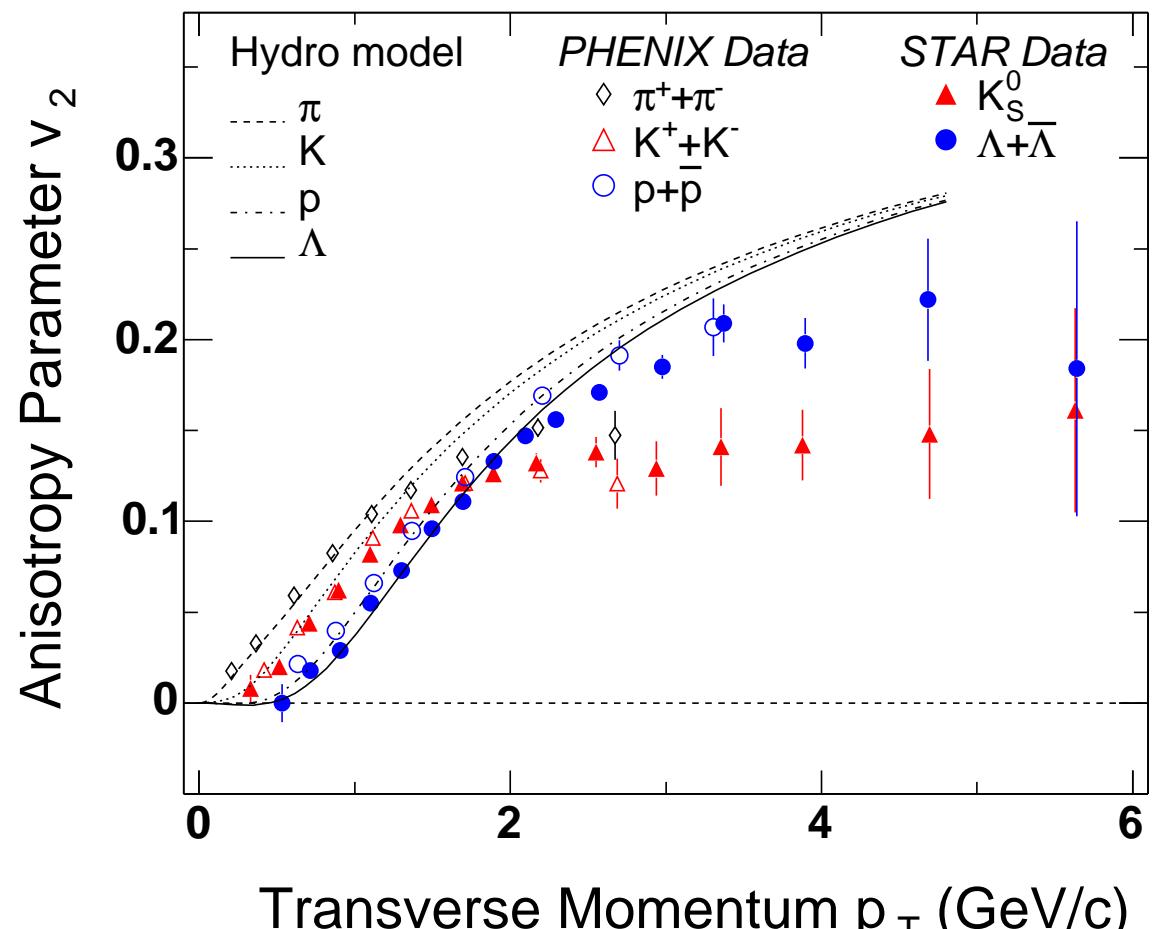
$$v_T \sim 0.6c!$$

$$m_T = \sqrt{p_T^2 + m^2}$$

Hydrodynamic
expansion converts
coordinate space
anisotropy
to momentum space
anisotropy

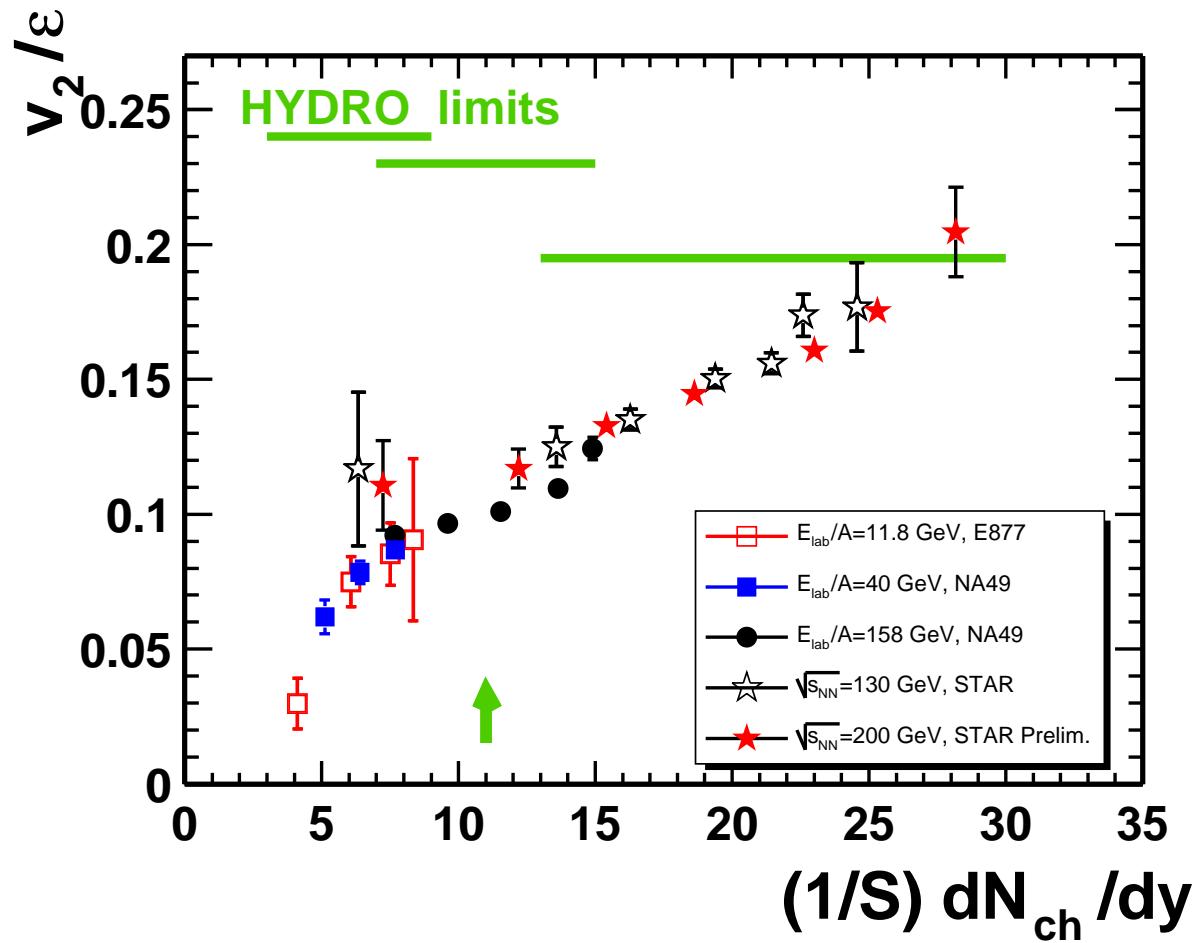


Elliptic Flow



source: U. Heinz (2005)

Elliptic Flow II



source: U. Heinz (2005)

Elliptic Flow III: Viscosity

$$T_{\mu\nu} \rightarrow T_{\mu\nu} + \eta(\nabla_\mu u_\nu + \nabla_\nu u_\mu - trace)$$

perturbative QCD

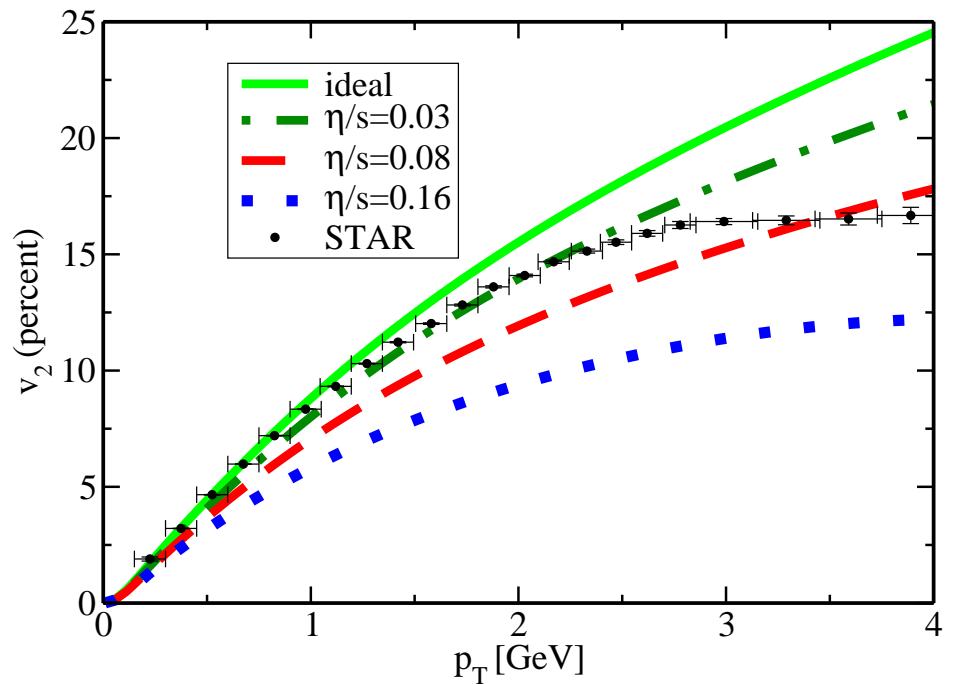
$$\frac{\eta}{s} = \frac{5.12}{g^4 \log(g^{-1})} \sim 1$$

Arnold, Moore, Yaffe

universal bound?

$$\frac{\eta}{s} \geq \frac{1}{4\pi}$$

Son, Starinets

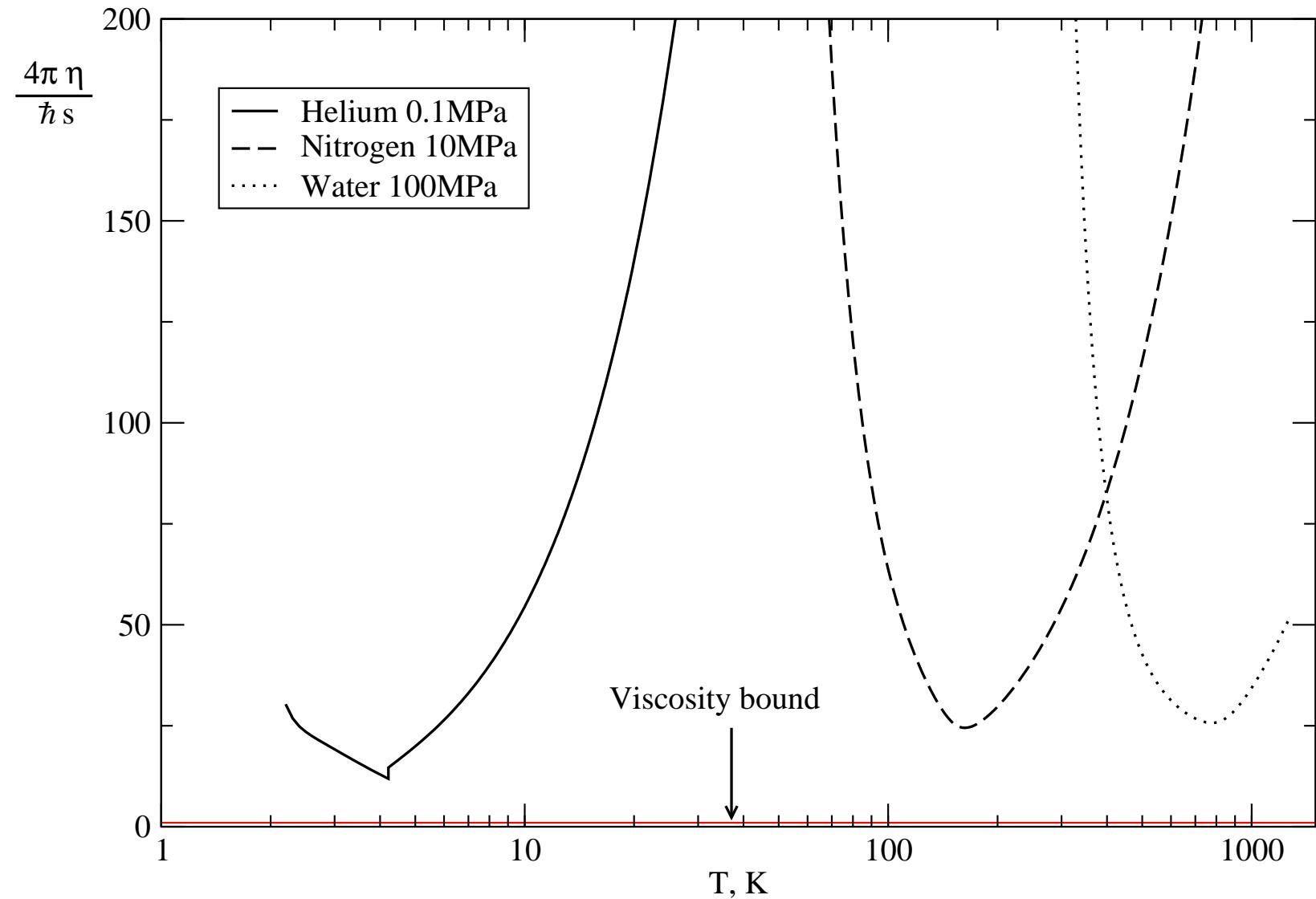


Romatschke (2007), Teaney (2003)

A (Most) Perfect Fluid?

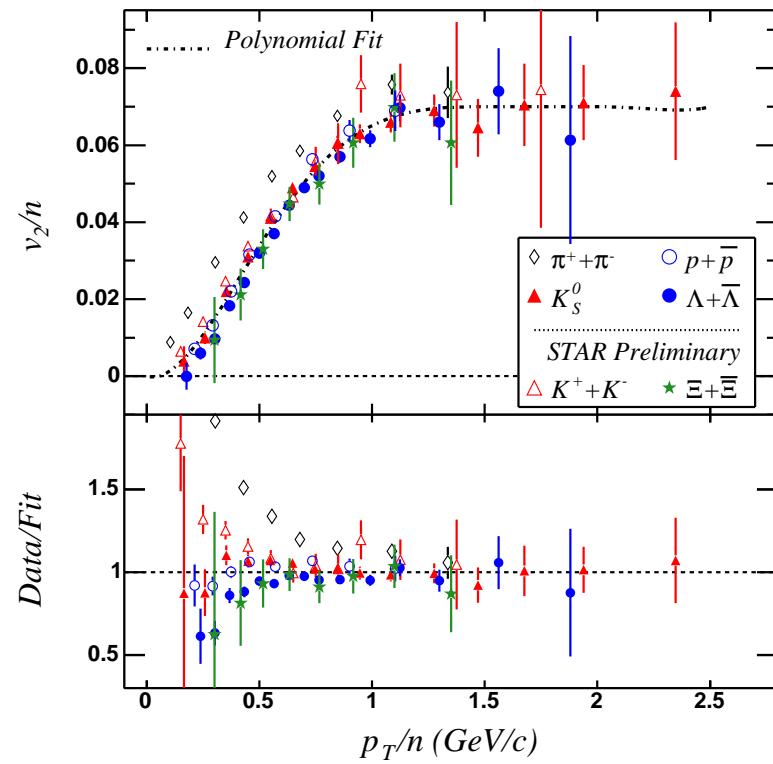
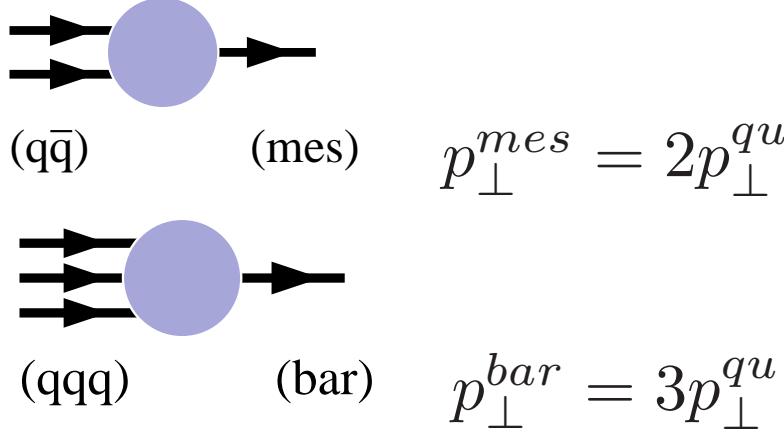
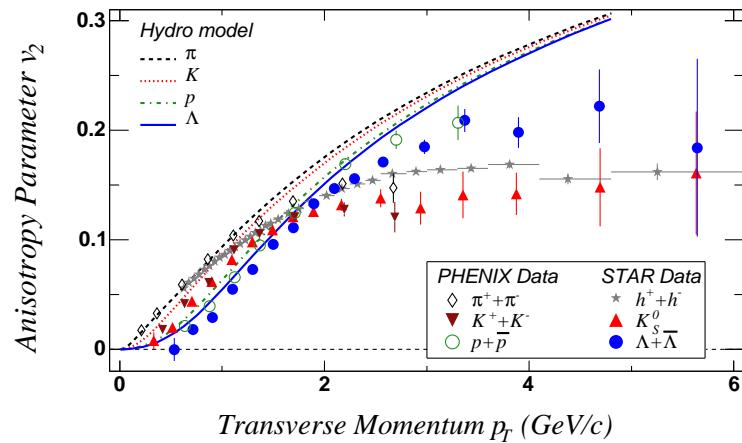


A (Most) Perfect Fluid?



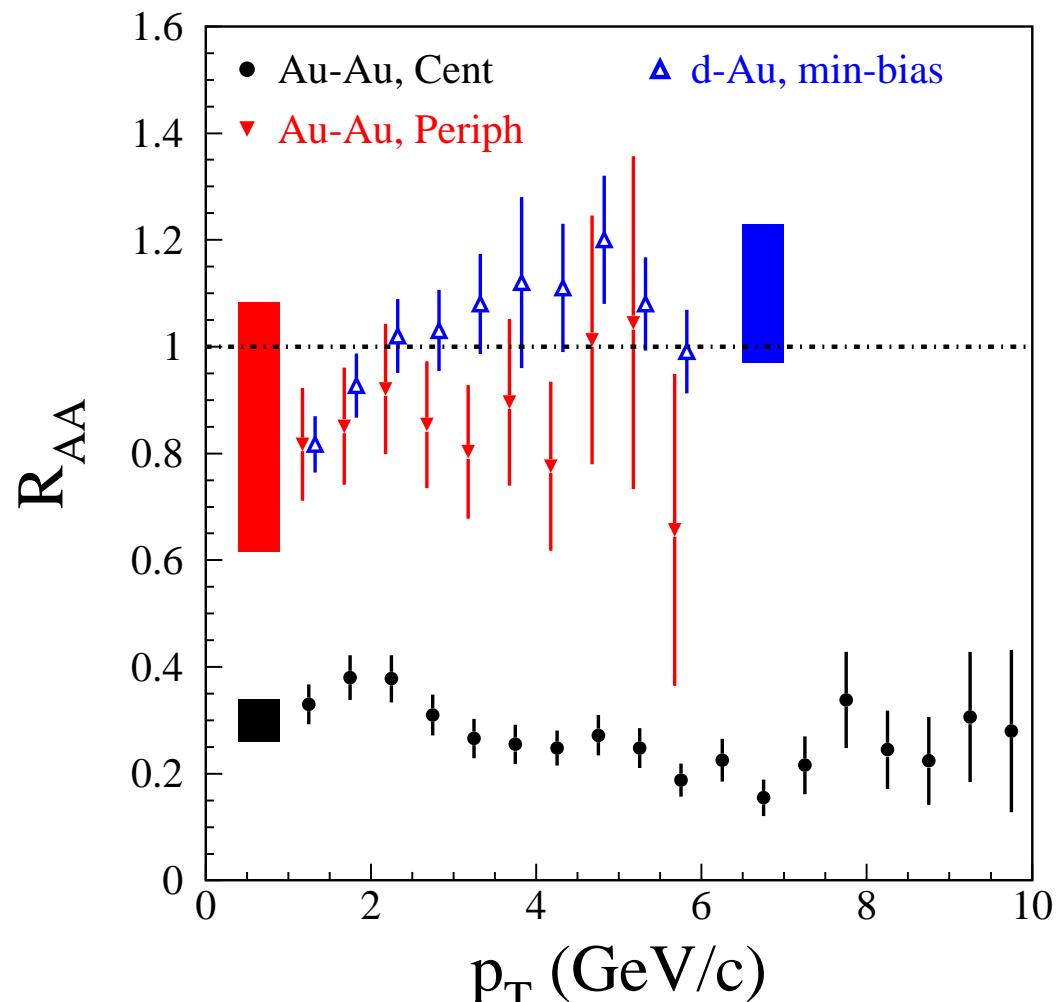
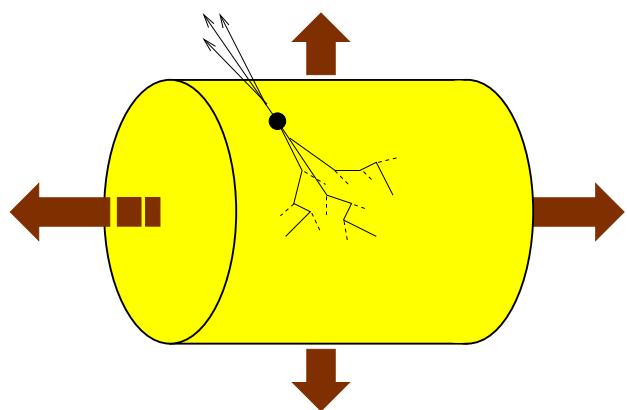
Elliptic Flow IV: Recombination

“quark number” scaling of elliptic flow



Jet Quenching

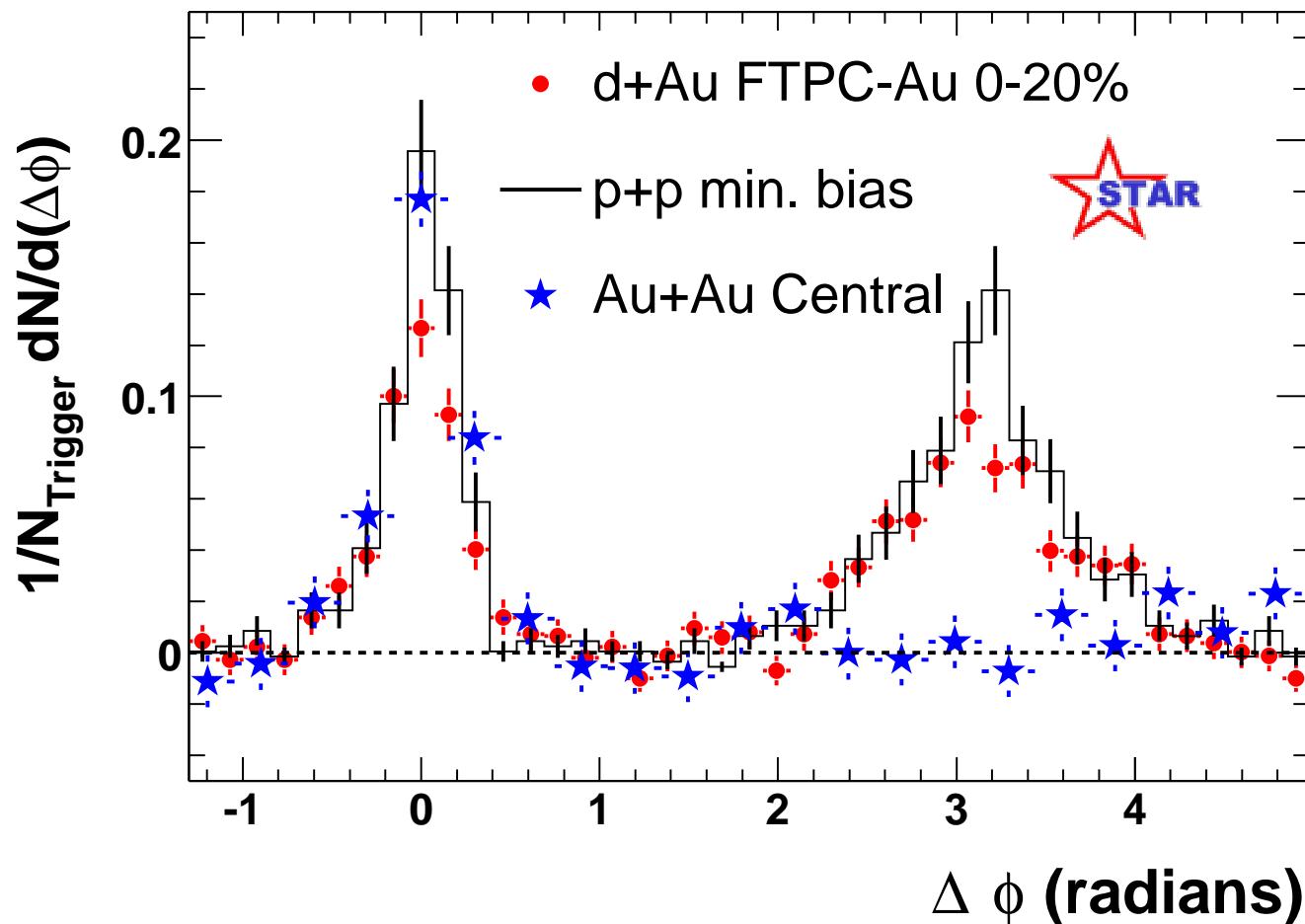
$$R_{AA} = \frac{n_{AA}}{N_{coll} n_{pp}}$$



source: Phenix White Paper (2005)

Jet Quenching II

Disappearance of away-side jet

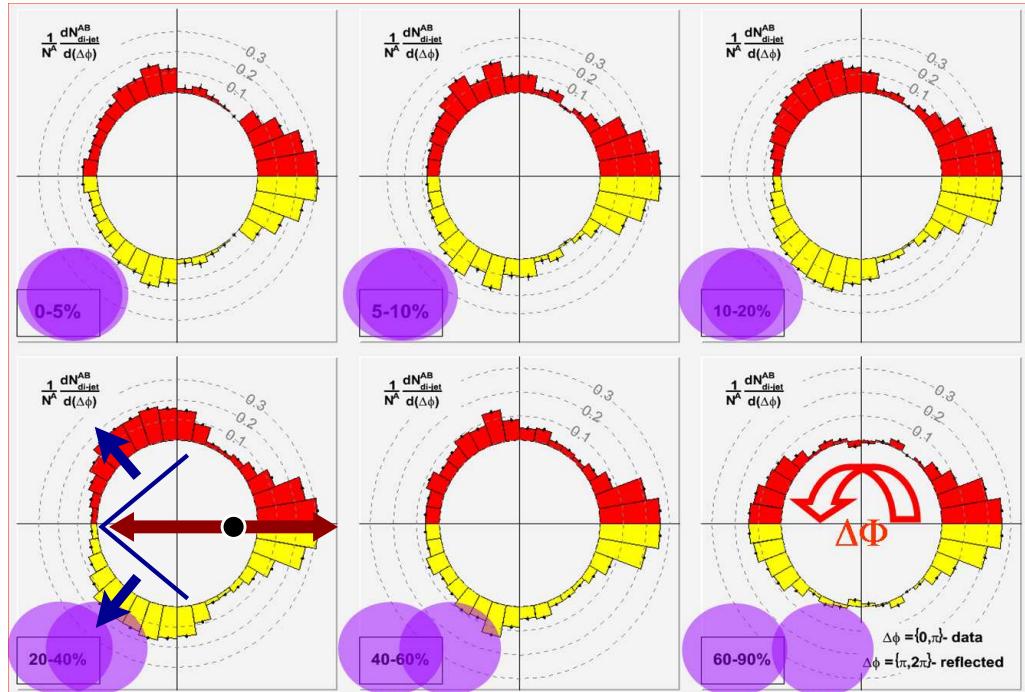


source: Star White Paper (2005)

Jet Quenching III: The Mach Cone

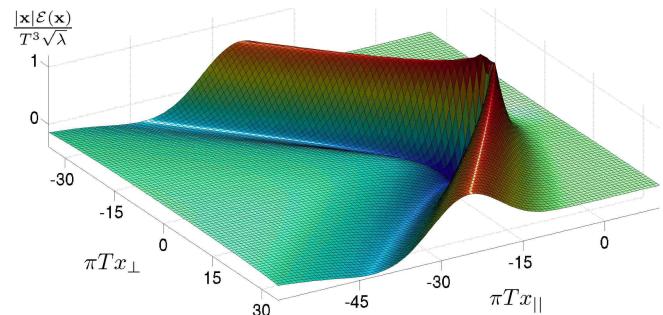
azimuthal multiplicity $dN/d\phi$

(high energy trigger particle at $\phi = 0$)



wake of a fast quark

in $\mathcal{N} = 4$ plasma



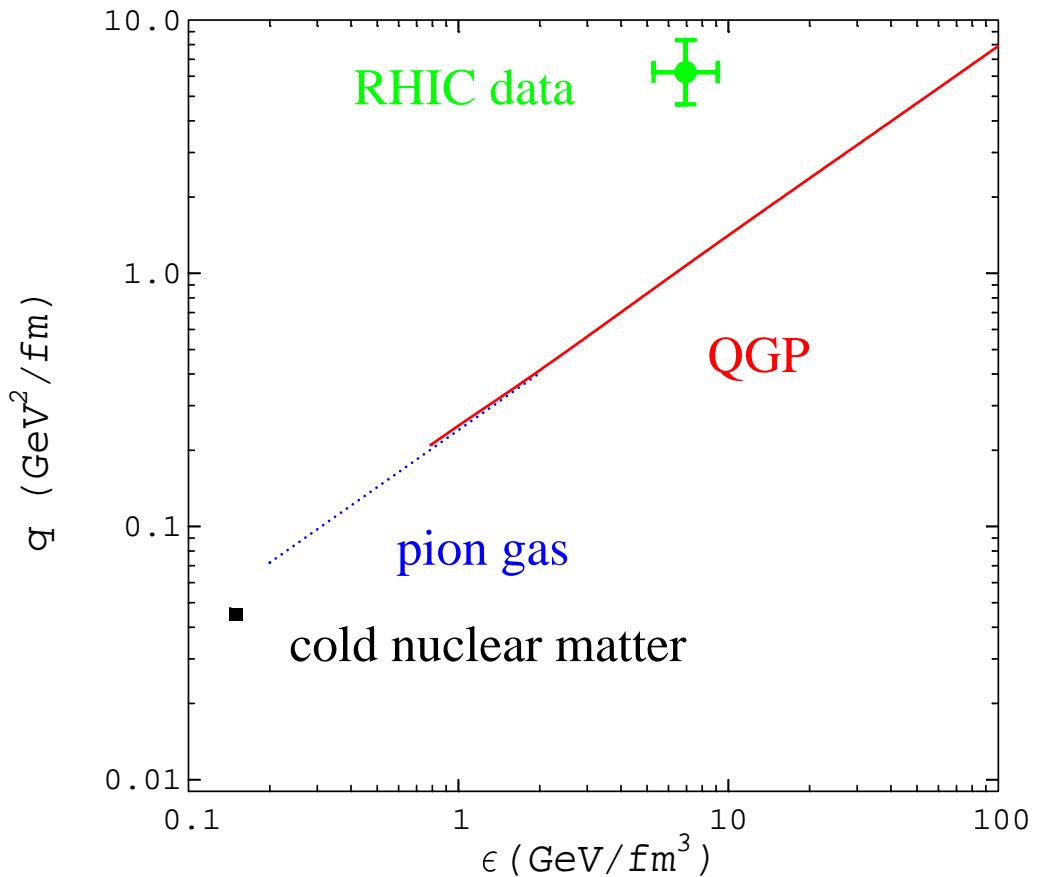
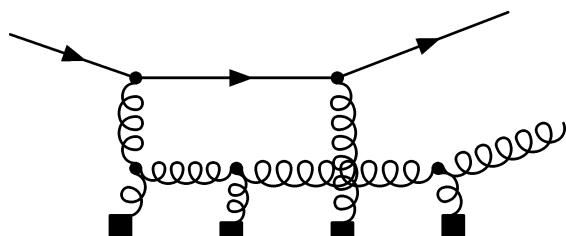
Chesler and Yaffe (2007)

source: Phenix (PRL, 2006), W. Zajc (2007)

Jet Quenching: Theory

energy loss governed by

$$\hat{q} = \rho \int q_\perp^2 dq_\perp^2 \frac{d\sigma}{dq_\perp^2}$$

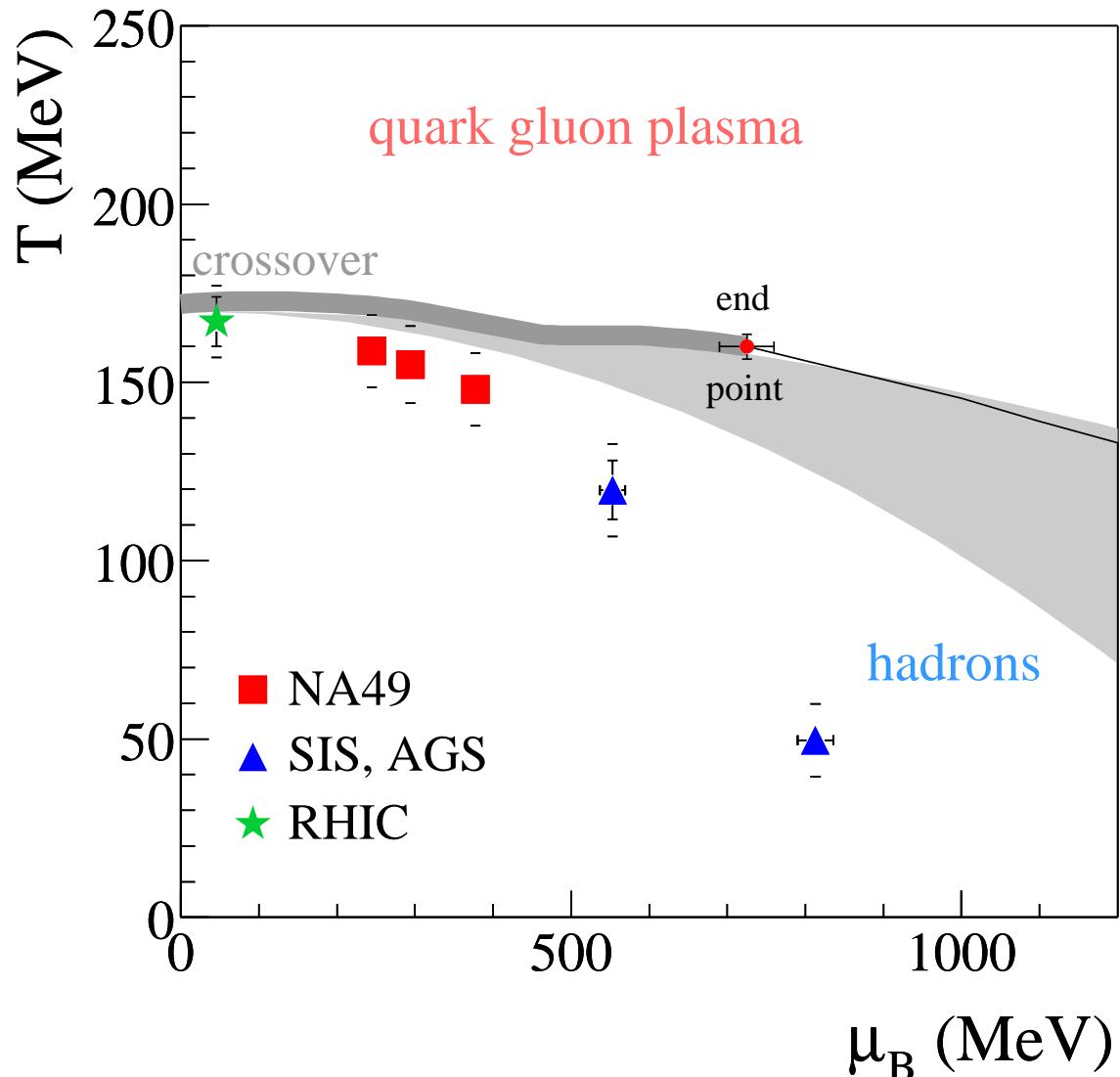


larger than pQCD predicts?

also: large energy loss of heavy quarks

[some recent doubts about \hat{q} , see P. Stankus seminar], source: R. Baier (2004)

Phase Diagram: Freezeout



Summary (Experiment)

Matter equilibrates quickly and behaves collectively

Little Bang, not little fizzle

Initial energy density in excess of 10 GeV/fm^3

Conditions for Plasma achieved

Evidence for strongly interacting Plasma (“sQGP”)

Fast equilibration $\tau_0 \ll 1 \text{ fm}$

Large elliptic flow, “perfect fluid”

Strong energy loss of leading partons

The Future: LHC

