Free Fermi Gas

Pauli Potential

Conclusion

Semi Classical Simulation of the Pauli Potential

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Outline

- Introduction
- Pree Fermi gas
 - Statistical properties of a Fermi system
- Pauli Potential
 - The existing Phenomenological Pauli Potential
 - Our Model of Pauli Potential
- Conclusion



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Pasta Phase in Neutron Stars





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Equation of State for Pasta

 In earlier work (by Horowitz, Perez Garcia, and Piekarewicz) the structure and spin-independent neutrino response of the pasta was studied via purely classical simulation.

$$H = T + \sum_{i < j} V_{\rm nuc}(i, j)$$

• V_{nuc}(i,j) is the two body interaction term given as

 $V_{\text{nuc}}(i,j) = ae^{-r_{ij}^2/\Lambda} + \left[b + c\tau_z(i)\tau_z(j)\right]e^{-r_{ij}^2/2\Lambda} + V_{\text{coulomb}}(i,j)$

- I am interested in generalizing this work to the study of the spin-dependent response.
- To do so, one must (at least) understand fermionic correlations induced by the Pauli exclusion principle.



• We incorporate Pauli exclusion principle by adding a fictitious momentum-dependent *Pauli potential*

$$H_{\text{new}} = T + \sum_{i < j} V_{\text{nuc}}(i, j) + V_{\text{Pauli}}(i, j)$$

- We will focus our attention to this added term (*V*_{Pauli}).
 - We will discuss the existing Pauli potential. We will show that the existing model does not reproduce most of the properties of a quantum free Fermi gas.
 - Proposed a modified Pauli potential to overcome this shortcomings.

(This work has been submitted to Phys. Rev. C (archieve nucl.-th/0702086))



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Properties of a Free Fermi Gas

 A zero-temperature quantum free Fermi gas can be described by a Slater determinant

$$\Psi_{FG}(\mathbf{p}_1,\ldots,\mathbf{p}_N;\mathbf{r}_1,\ldots,\mathbf{r}_N) = \frac{1}{\sqrt{N!}} \begin{vmatrix} \varphi_{\mathbf{p}_1}(\mathbf{r}_1) & \cdots & \varphi_{\mathbf{p}_1}(\mathbf{r}_N) \\ \vdots & \cdots & \vdots \\ \varphi_{\mathbf{p}_N}(\mathbf{r}_1) & \cdots & \varphi_{\mathbf{p}_N}(\mathbf{r}_N) \end{vmatrix}$$

- Properties of the Slater determinant:
 - ▶ No two fermions can have the same momentum.
 - No two fermions can have the same spatial coordinates
 - No fermion can have a momentum larger than p_F
- The kinetic energy, correlation function, and momentum distribution of a Free Fermi gas



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1 $E/N \simeq 3/5\epsilon_F$ **2** $g_2(r) = 1 - \left(\frac{3j_1(p_F r)}{p_F r}\right)^2$; $j_1(p_F r) = \frac{sin(p_F r)}{(p_F r)^2} - \frac{cos(p_F r)}{(p_F r)}$ **3** $f(p) = \frac{3}{p_F^3}p^2\theta(p_F - p)$; $\left(\int f(p)dp = 1\right)$





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Semi Classical Simulation of the Pauli Potential

$$V_{ ext{Pauli}}(\mathbf{p}_1, \dots, \mathbf{p}_N; \mathbf{r}_1, \dots, \mathbf{r}_N) = \sum_{i < j=1}^N V_0 \exp(-s_{ij}^2/2) \delta_{\tau_i \tau_j} \delta_{\sigma_i \sigma_j}$$

 s_{ij} is the dimensionless phase-space distance between two particles

$$s_{ij}^2 \equiv rac{|\mathbf{p}_i - \mathbf{p}_j|^2}{p_0^2} + rac{|\mathbf{r}_i - \mathbf{r}_j|^2}{r_0^2}$$

- It vanishes for different fermionic species
- When *s_{ij}* is small, the Pauli potential is big so the particles are penalized for getting too close in phase space.
- When *s_{ij}* is large, Pauli potential is small so particles are no longer penalized
- No penalty, if either r_{ij} or p_{ij} is large

This is less restrictive than a Slater determinant.



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Model parameters (V_0 , p_0 , q_0) are determined by fitting to the kinetic energy of a quantum free Fermi gas at various densities.

Simulation results (red circles) are in good agreement with exact (black line).



	Pauli Potential	

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- The Pauli potential does not provide enough suppression between particles at short distances.
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Image: A matrix



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- ► However, if particles are close in space (r_{ij} ≪ 1), as long as momentum between the particles is large (p_{ij} ≫ 1) (or visa versa) the Pauli potential vanishes.
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New Model for Pauli Potential

$$V_{\text{Pauli}}(\mathbf{p}_1, \dots, \mathbf{p}_N; \mathbf{r}_1, \dots, \mathbf{r}_N) = \sum_{i < j}^N V_A \ e^{-r_{ij}/r_o} \delta_{\sigma_i \sigma_j} \delta_{\tau_i \tau_j}$$

$$+\sum_{i< j}^{N} \mathbf{V}_{\mathrm{B}} \mathbf{e}^{-\mathbf{p}_{ij}/\mathbf{p}_{\mathrm{o}}} \delta_{\sigma_{i}\sigma_{j}} \delta_{\tau_{i}\tau_{j}} + \sum_{i=1}^{N} \frac{\mathbf{V}_{\mathrm{C}}}{1+e^{-\eta(q_{i}^{2}-1)}}$$

- The first term penalizes particles if they get too close in space.
- The second term penalizes particles if they get too close in momentum space.
- The last term provides a cut-off for the $p > p_F$ region.

Model parameters are determined by fitting to the properties of a quantum free Fermi gas.



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Model parameters (V_A , V_B , V_C , r_o , p_o , η) are determined by fitting to the kinetic energy of a quantum free Fermi gas at various densities.

Our model is in good agreement with the analytic prediction.



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The new Pauli potential reproduces accurately the momentum distribution of a quantum free Fermi gas.



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- The new Pauli potential reproduces accurately the two-fermion. correlation function of a quantum free Fermi gas.
- The potential gives enough suppression at r=0.
- There is no crystallization at low (or any) densities.



Pauli Potential

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Conclusions & Future Study

- To include spin into our calculation we need to add a Pauli term into our classical Hamiltonian.
- The "old" Pauli potential fails to accurately describe all properties of a quantum free Fermi gas.
- We proposed a new Pauli potential that has successfully mimicked the Pauli exclusion principle
- We add this new spin dependent potential to our pasta Hamiltonian and proceed (after re-fitting) to calculate the spin dependence response of the nuclear pasta.

Work is still in progress ...

