

IV. Low Energy Probes of SUSY

- How is SUSY broken? $M_{SUSY} \neq M_{SM}$
- How viable is SUSY dark matter ?
- Is there enough SUSY CP violation to account for the matter-antimatter asymmetry?

High precision, low energy measurements can probe new physics in the “desert”

Complementary to Z^0 pole

- LEP & SLD programs are complete
- Sensitivity to off Z^0 -resonance physics
- Precision is improving

Precision ~ Mass Scale

$$\delta_{SUSY} = \frac{\Delta O^{SUSY}}{O^{SM}} \approx \frac{\alpha}{\pi} \left(\frac{M}{\tilde{M}} \right)^2$$

$$M=m_\mu \quad \delta \sim 2 \times 10^{-9}$$

$$\delta^{\text{exp}} \sim 1 \times 10^{-9}$$

$$M=M_W \quad \delta \sim 10^{-3}$$

Weak decays

$$d \rightarrow u e^- \bar{\nu}_e$$

$$s \rightarrow u e^- \bar{\nu}_e$$

$$b \rightarrow u e^- \bar{\nu}_e$$

$$(u \quad c \quad t) \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$$

SM

0.996 ± 0.001 Expt

0.948 ± 0.001

0.048 ± 0.0000

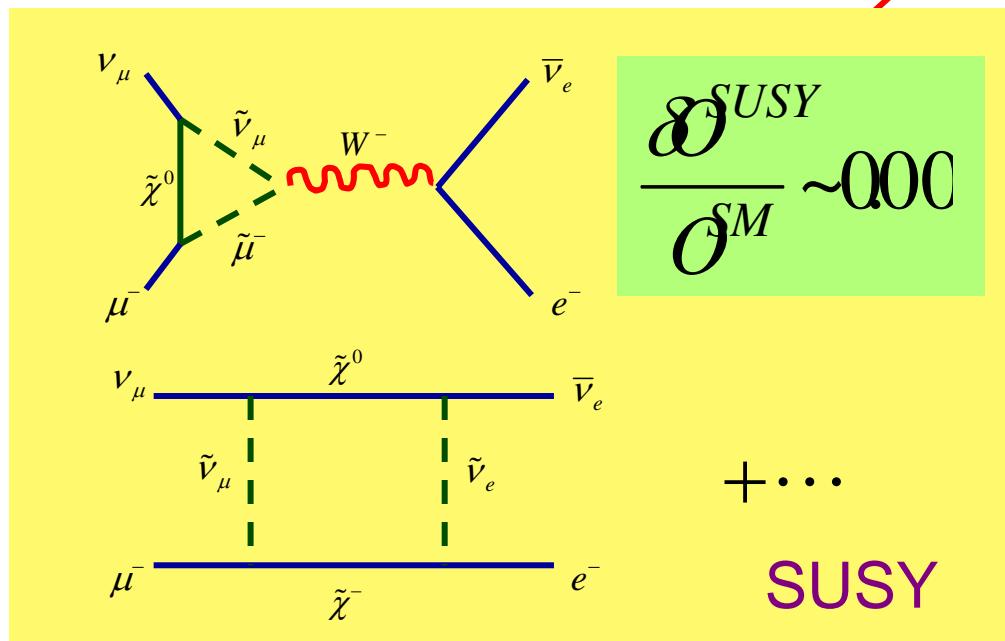
0.000 ± 0.0000

Weak decays

$$\begin{aligned} d &\rightarrow u \ e^- \ \bar{\nu}_e \\ s &\rightarrow u \ e^- \ \bar{\nu}_e \\ b &\rightarrow u \ e^- \ \bar{\nu}_e \end{aligned}$$

$$(u \ c \ t) \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

A red circle highlights the V_{ud} element of the CKM matrix.



$$\frac{G_F^\beta}{G_F^\mu} = |V_{ud}| \left(1 + \Delta r_\beta - \Delta r_\mu \right)$$

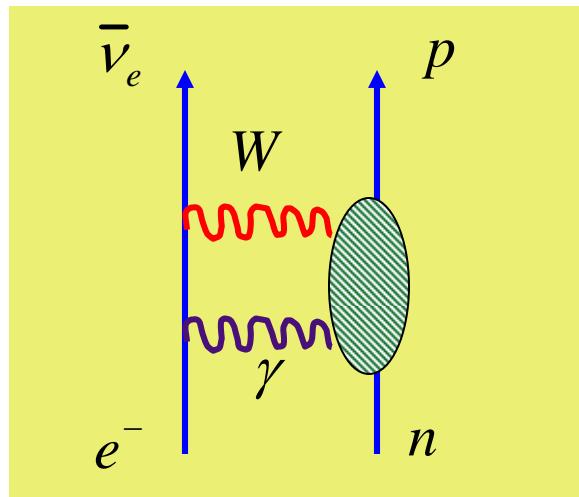
A red oval highlights the term $\Delta r_\beta - \Delta r_\mu$.

New physics

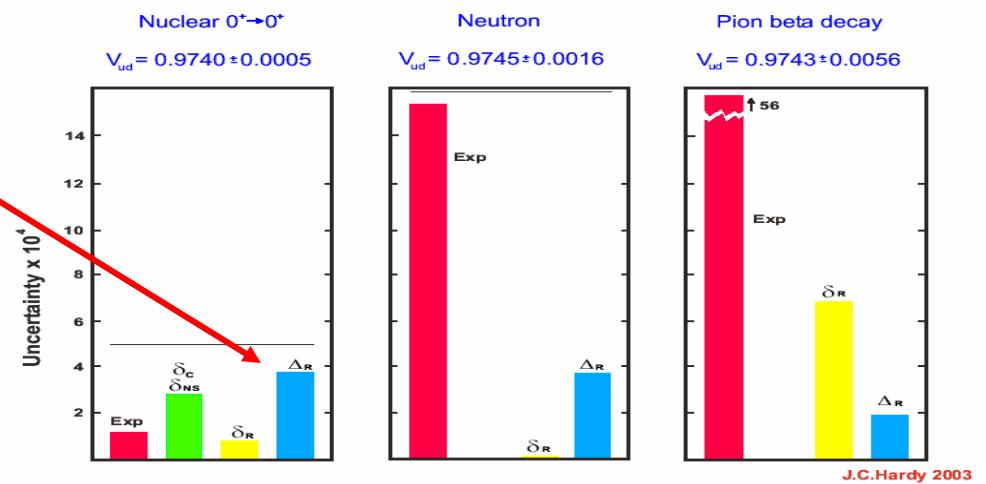
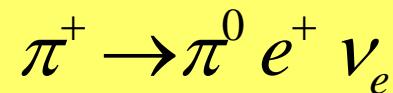
Weak decays

$$\frac{G_F^\beta}{G_F^\mu} = |V_{ud}| \left(1 + \Delta r_\beta - \Delta r_\mu \right)$$

SM theory input



β -decay



$$M_{W\gamma} = \frac{G_F}{\sqrt{2}} \frac{\hat{\alpha}}{2\pi} \left[\ln\left(\frac{M_Z^2}{\Lambda^2}\right) + C_{\gamma W}(\Lambda) \right]$$

Weak decays

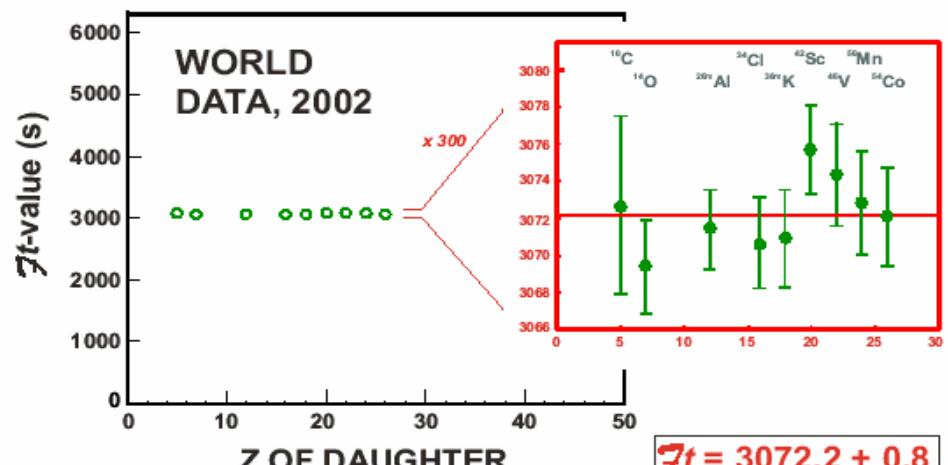
$$\frac{G_F^\beta}{G_F^\mu} = |V_{ud}| \left(1 + \Delta r_\beta - \Delta r_\mu \right)$$

0^+ ! 0^+ “Superallowed”

$$Ft = ft(1 + \delta_R + \delta_{NS})(1 - \delta_C) \\ = K / 2(G_F^\beta)^2$$

Nuclear structure-dependent corrections

β -decay



Weak decays

$$\frac{G_F^\beta}{G_F^\mu} = |V_{ud}| \left(1 + \Delta r_\beta - \Delta r_\mu \right)$$

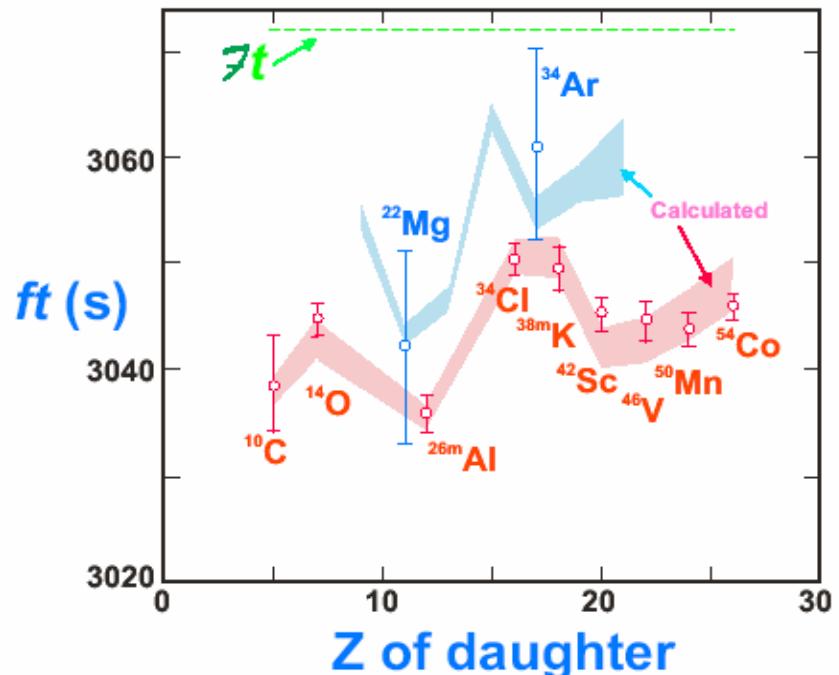
β -decay



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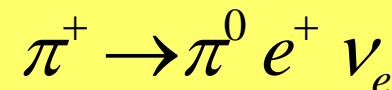
Nuclear structure-dependent corrections



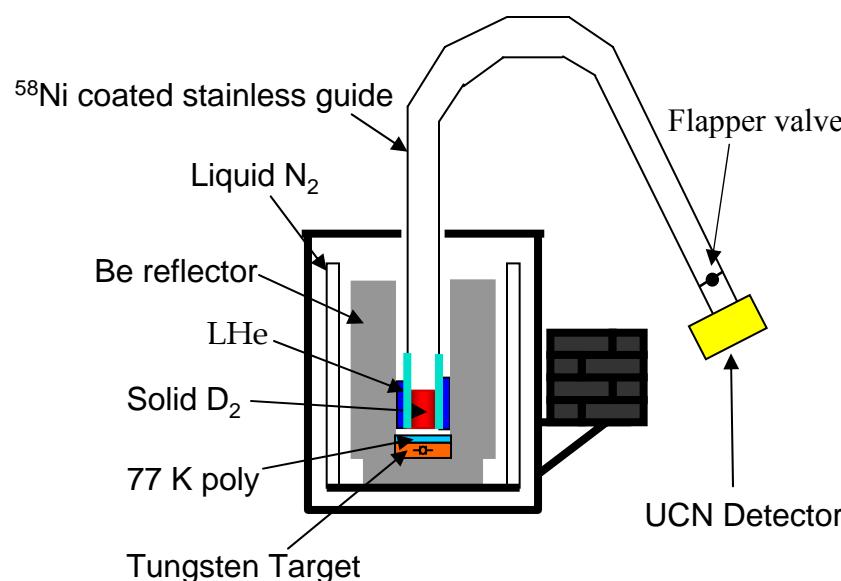
Weak decays

$$\frac{G_F^\beta}{G_F^\mu} = |V_{ud}| \left(1 + \Delta r_\beta - \Delta r_\mu \right)$$

β -decay



Ultra cold neutrons



LANSCE: "UCN A"

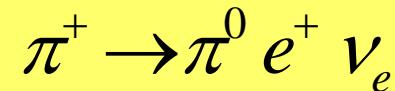
$$dW \propto 1 + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + A \vec{\sigma}_n \cdot \frac{\vec{p}_e}{E_e} + \dots$$

Future SNS: Pulsed
Cold Neutrons: "abBA"

Weak decays

$$\frac{G_F^\beta}{G_F^\mu} = |V_{ud}| \left(1 + \Delta r_\beta - \Delta r_\mu \right)$$

β -decay



$$dW \propto \frac{1}{\tau_n} F(E_e) \left[1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_v}{E_e \cdot E_v} + b \frac{m_e}{E_e} + A \frac{\sigma_n \cdot \mathbf{p}_e}{E_e} + B \frac{\sigma_n \cdot \mathbf{p}_v}{E_v} \right]$$

$\tau_n \propto 1/(g_V^2 + 3g_A^2)$
 $a = \frac{1 - \left(\frac{g_A}{g_V}\right)^2}{1 - 3\left(\frac{g_A}{g_V}\right)}$
 $b = 0$
 $A = -2 \frac{\left(\frac{g_A}{g_V}\right)^2 + \left(\frac{g_A}{g_V}\right)}{1 - 3\left(\frac{g_A}{g_V}\right)^2}$
 $B = 2 \frac{\left(\frac{g_A}{g_V}\right)^2 - \left(\frac{g_A}{g_V}\right)}{1 - \left(\frac{g_A}{g_V}\right)^2}$

Future SNS: Pulsed
Cold Neutrons: “abBA”

G. Greene

Decay correlations

Weak decays

$$\frac{G_F^\beta}{G_F^\mu} = |V_{ud}| \left(1 + \Delta r_\beta - \Delta r_\mu \right)$$

Los Alamos Neutron Science Center (LANSCE)



LANSCE

β -decay



The Spallation Neutron Source



SNS

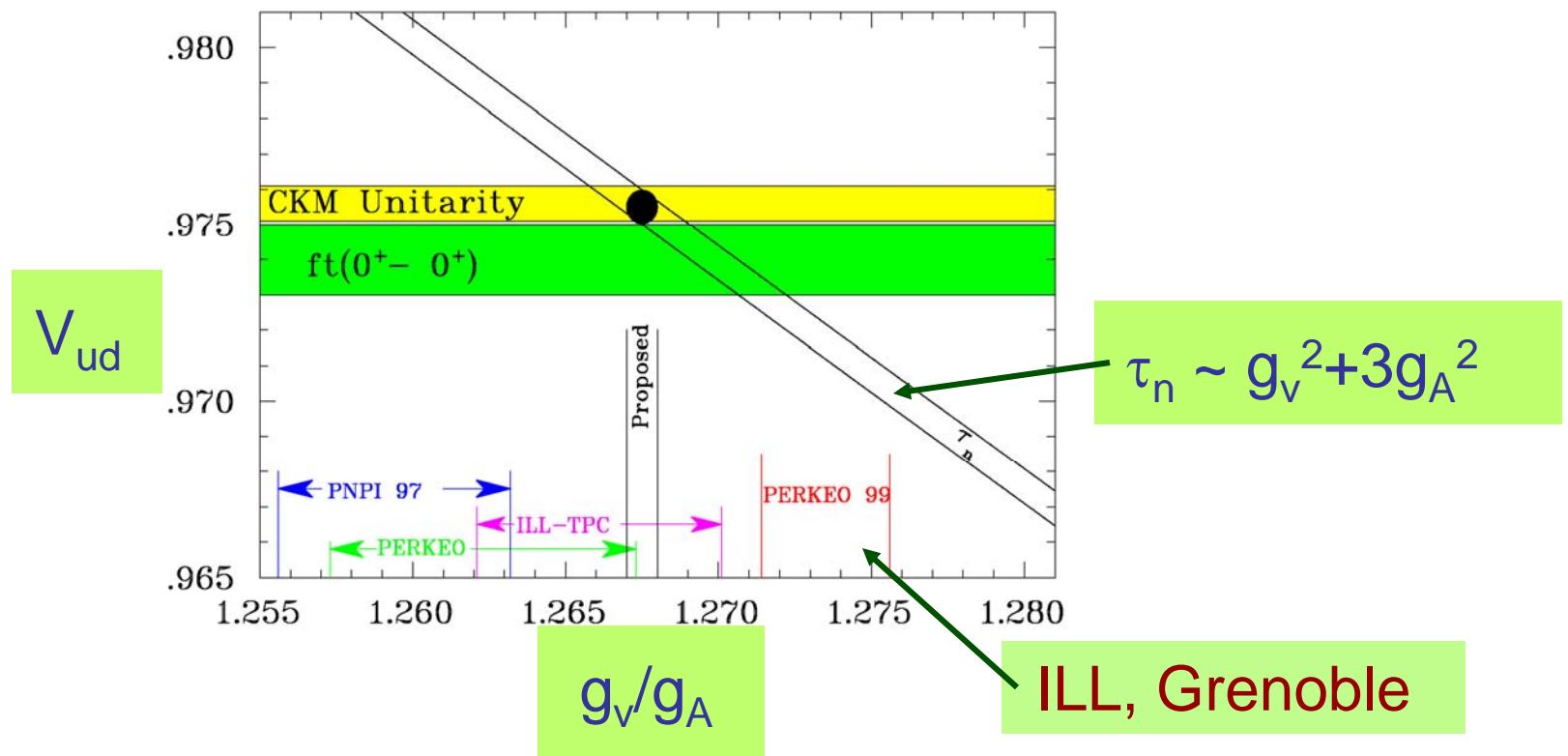


NIST

Weak decays

$$\frac{G_F^\beta}{G_F^\mu} = |V_{ud}| \left(1 + \Delta r_\beta - \Delta r_\mu \right)$$

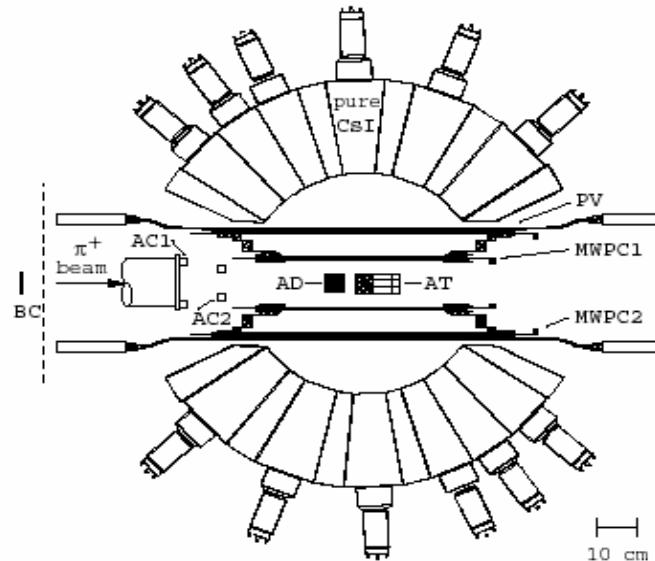
β -decay



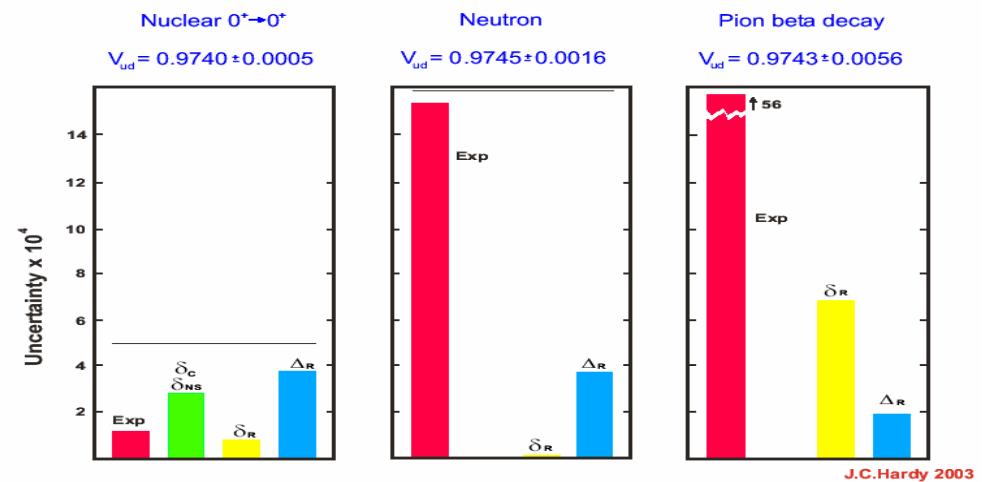
Weak decays

$$\frac{G_F^\beta}{G_F^\mu} = |V_{ud}| \left(1 + \Delta r_\beta - \Delta r_\mu \right)$$

PSI: “Pi-Beta”



β -decay



$$\Gamma(\pi^+ \rightarrow \pi^0 e^+ \nu_e) / \Gamma(\pi^+ \rightarrow \mu^+ \nu_\mu) \sim 1 \times 10^8$$

Weak decays

$$d \rightarrow u e^- \bar{\nu}_e$$

$$s \rightarrow u e^- \bar{\nu}_e$$

$$b \rightarrow u e^- \bar{\nu}_e$$

$$(u \quad c \quad t) \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} (d \quad s \quad b)$$

kaon decay

$$K^+ \rightarrow \pi^0 e^+ \bar{\nu}_e$$

$$\frac{G_F^K}{G_F^\mu} = |V_{us}| (1 + \Delta r_K - \Delta r_\mu)$$

Value of V_{us} important

SUSY Loops: Too small

Weak decays

$$d \rightarrow u e^- \bar{\nu}_e$$

$$s \rightarrow u e^- \bar{\nu}_e$$

$$b \rightarrow u e^- \bar{\nu}_e$$

$$(u \quad c \quad t) \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} (d \quad s \quad b)$$

kaon decay

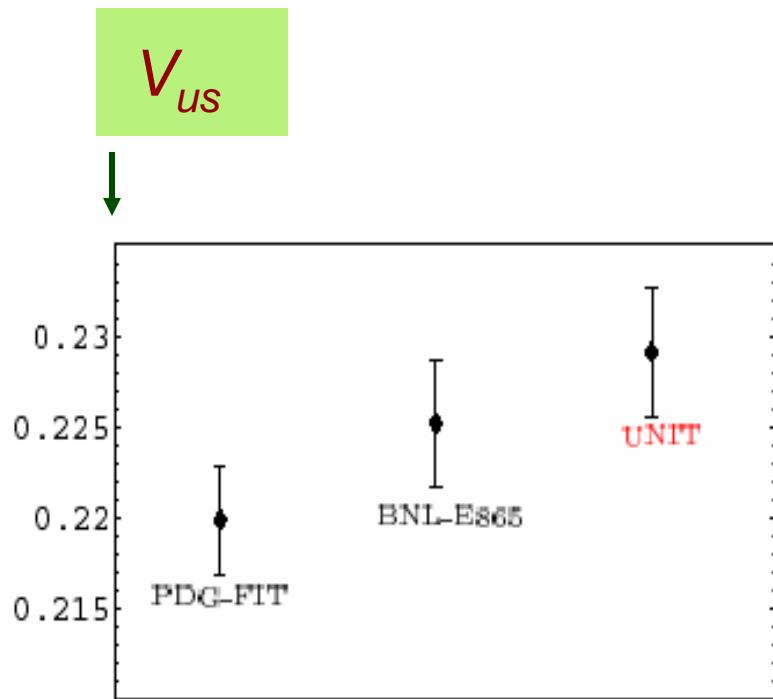
$$K^+ \rightarrow \pi^0 e^+ \bar{\nu}_e$$

$$\Gamma = |V_{us}|^2 \frac{G_F^2 S_{ew} M_K^5}{192\pi^3} \cdot |f_+(0)|^2 \cdot (1 + \Delta_{SU(2)} + \Delta_{EM}) \cdot I_{PS}(m_i, \lambda)$$

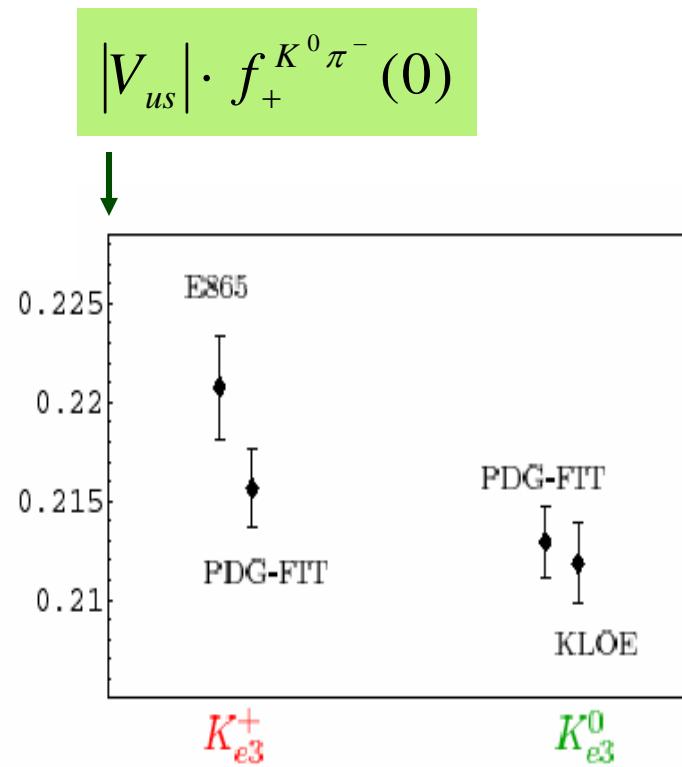
SM Theory

Weak decays

$$\frac{G_F^K}{G_F^\mu} = |V_{us}| \left(1 + \Delta r_K - \Delta r_\mu \right)$$



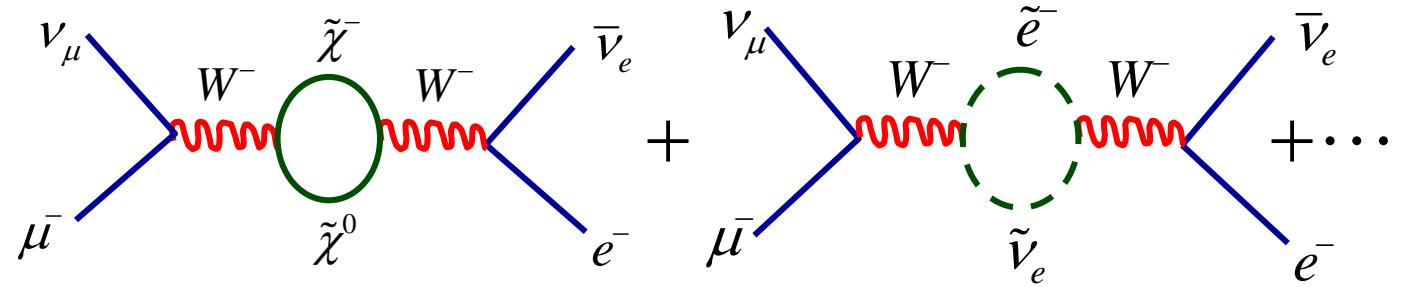
kaon decay



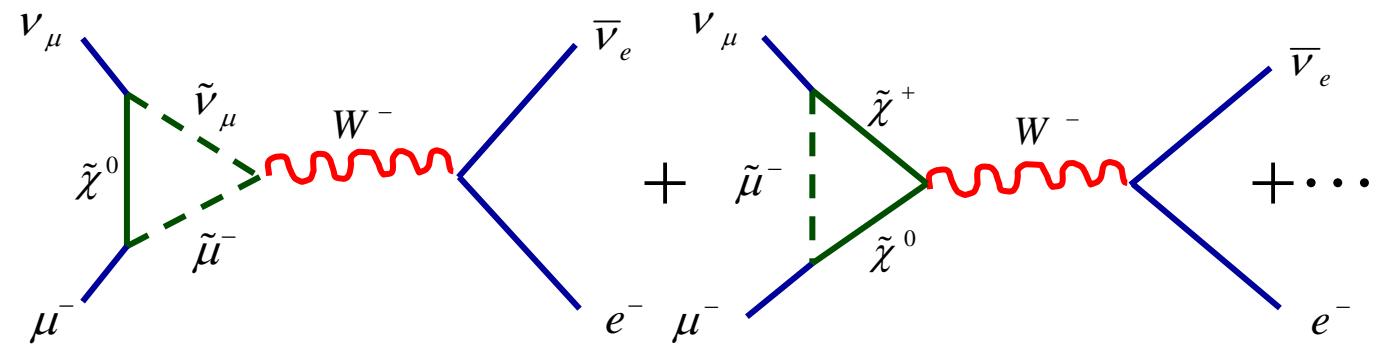
SUSY Radiative Corrections

Δr_μ

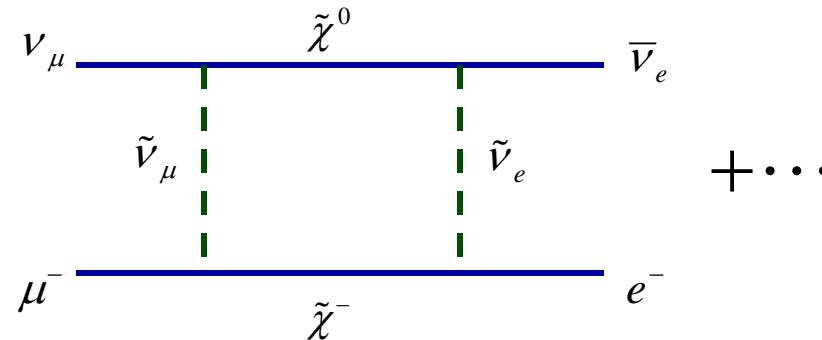
Propagator



Vertex & External leg



Box



Other Inputs

1. Muon (g-2):

Size of error bar crucial

2. W Mass:

$$G_F^\mu = \frac{\pi \alpha M_Z^2}{\sqrt{2} M_W^2 (M_Z^2 - M_W^2) (1 - \Delta r_\mu)}$$

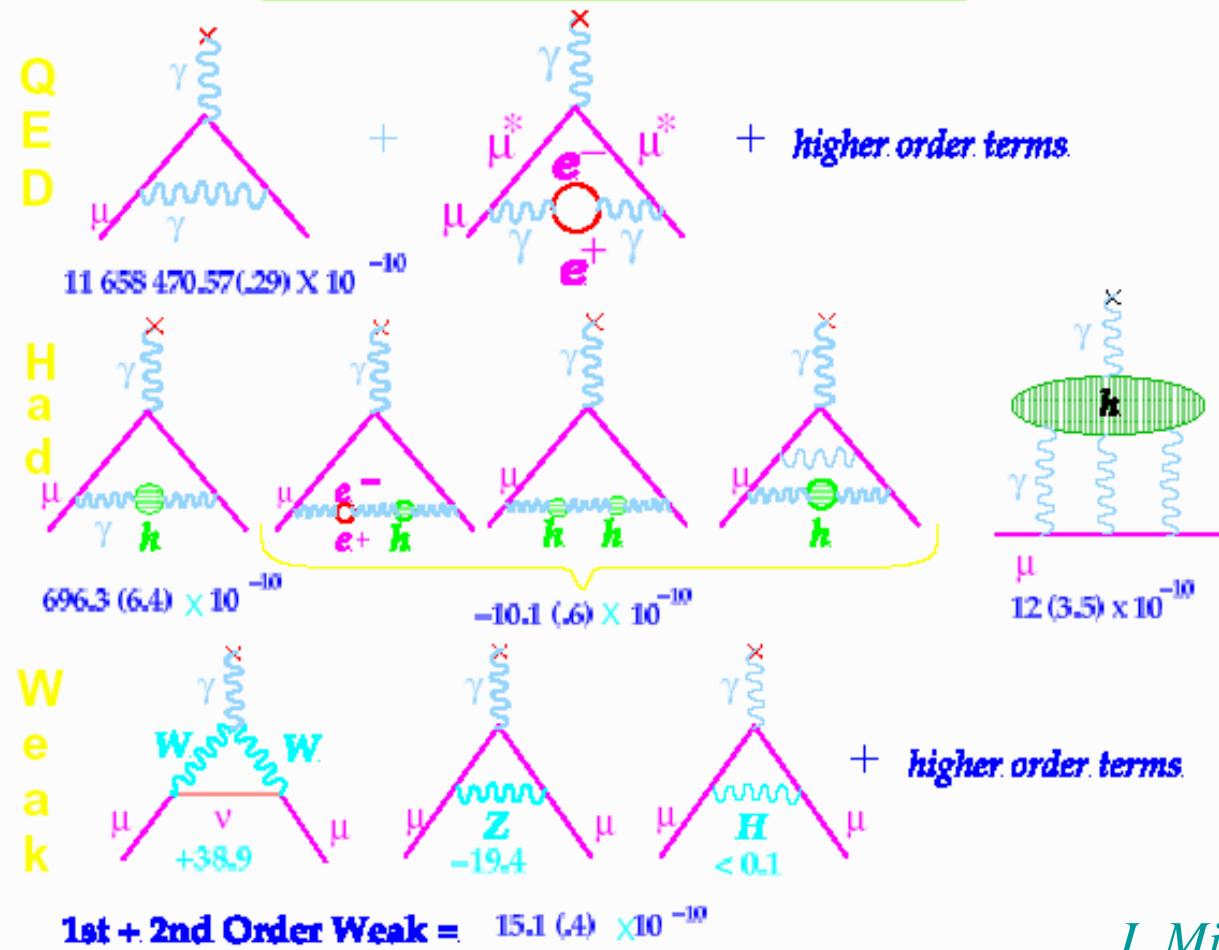
$$-0.009 \leq \Delta r_\mu^{SUSY} \leq 0.003$$

m_t, M_H

3. Superpartner Masses: Start analysis with collider lower bounds and vary later

Muon g-2

SM Theory

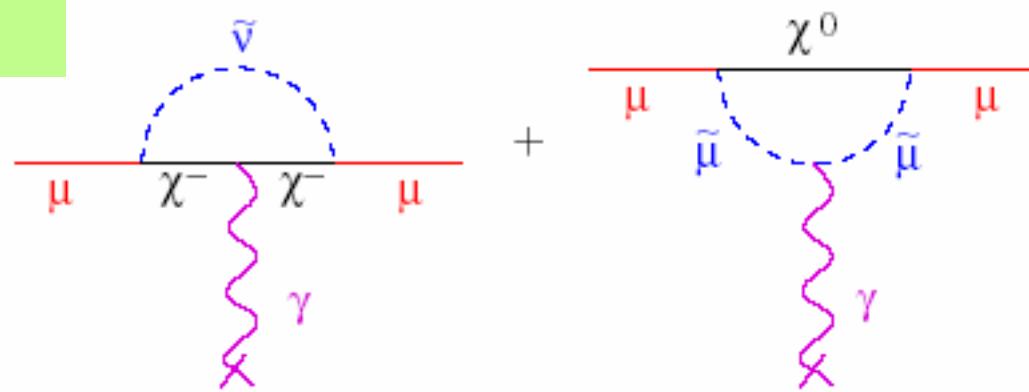


J. Miller

FNAL James Miller - The Muon Magnetic Moment Anomaly: Experiment

Muon g-2

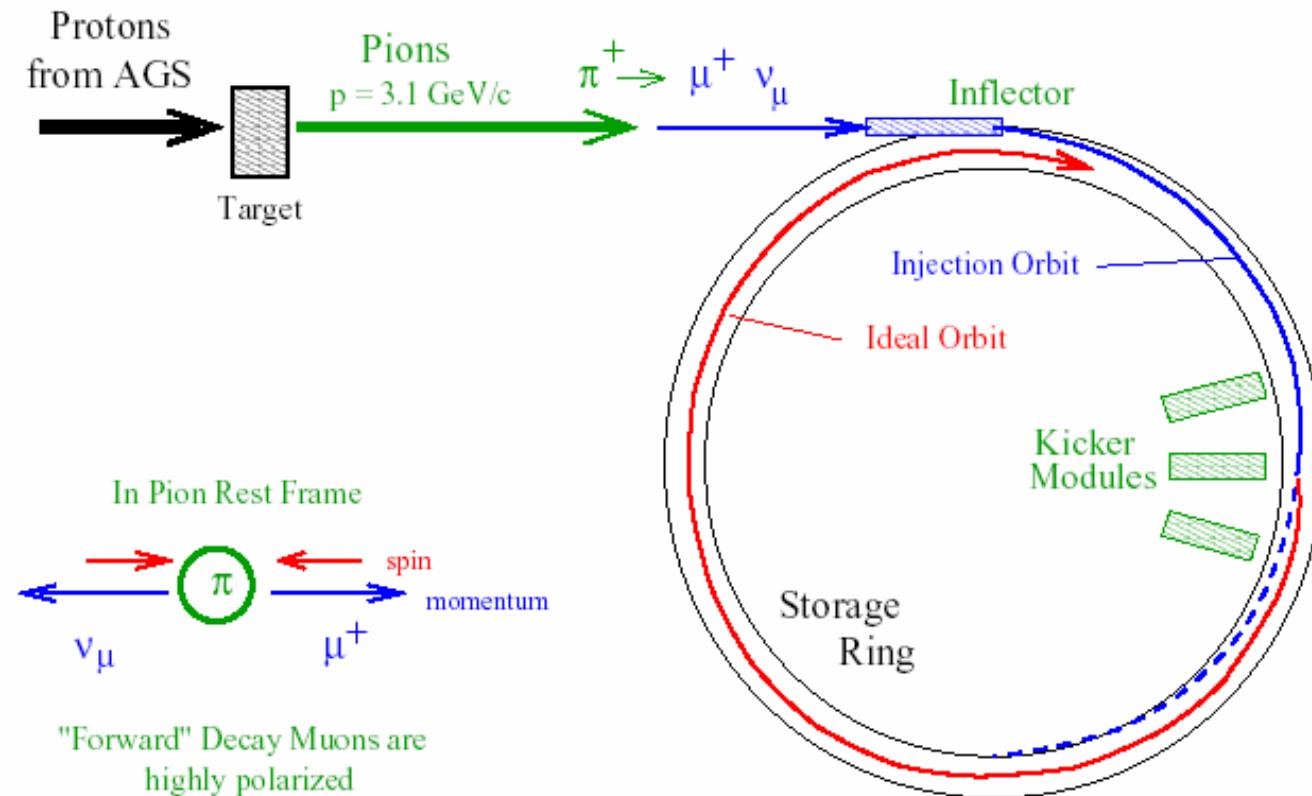
SUSY



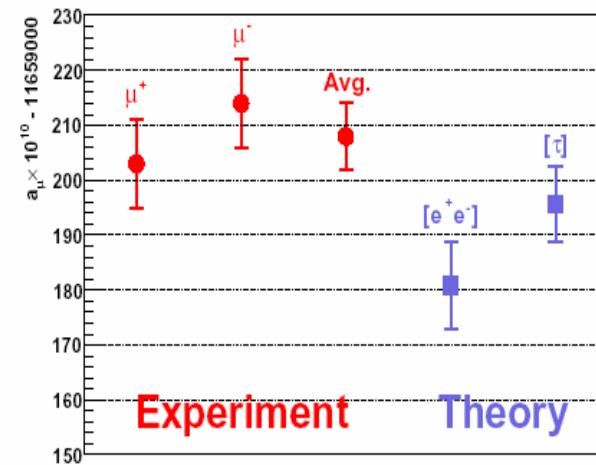
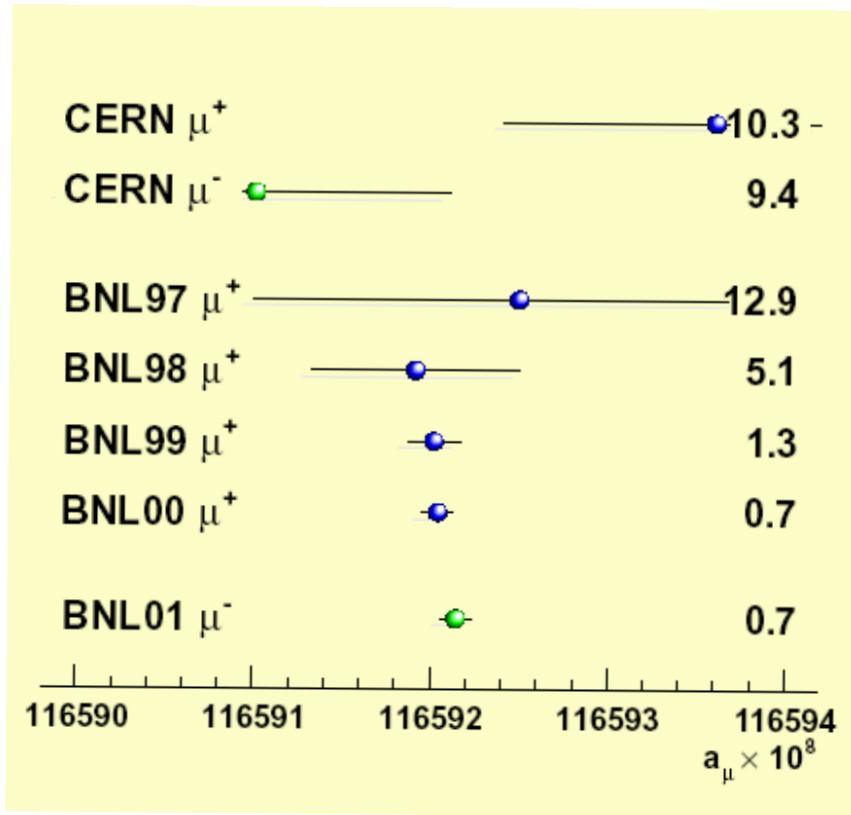
SUSY

Muon g-2

$(g-2)_\mu$ Experiment at BNL



Muon g-2



$$\Delta a_\mu(ee) = (23.9 \pm 9.9) \times 10^{-10} \quad 2.4 \text{ s.d.}$$

$$\Delta a_\mu(\tau) = (7.6 \pm 8.9) \times 10^{-10} \quad 0.9 \text{ s.d.}$$

W Boson & Top Quark Masses

$M_Z = 91.1875 \pm 0.0021$

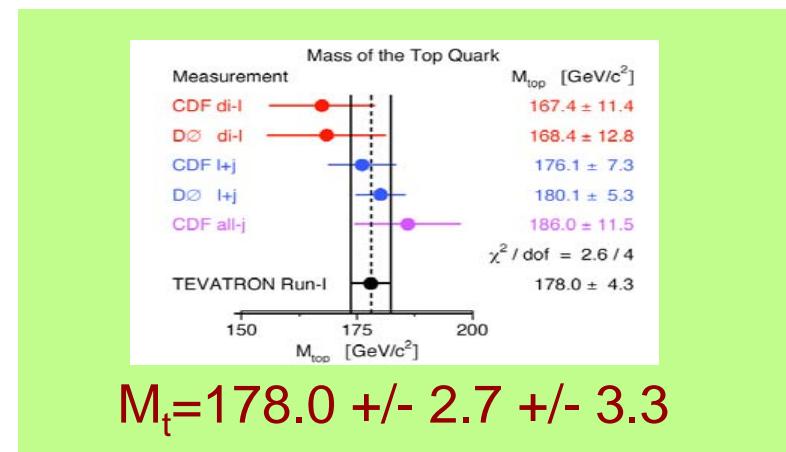
$$G_F^\mu = \frac{\pi \alpha M_Z^2}{\sqrt{2} M_W^2 (M_Z^2 - M_W^2) (1 - \Delta r_\mu)}$$

Most precisely measured

Dominant uncertainties

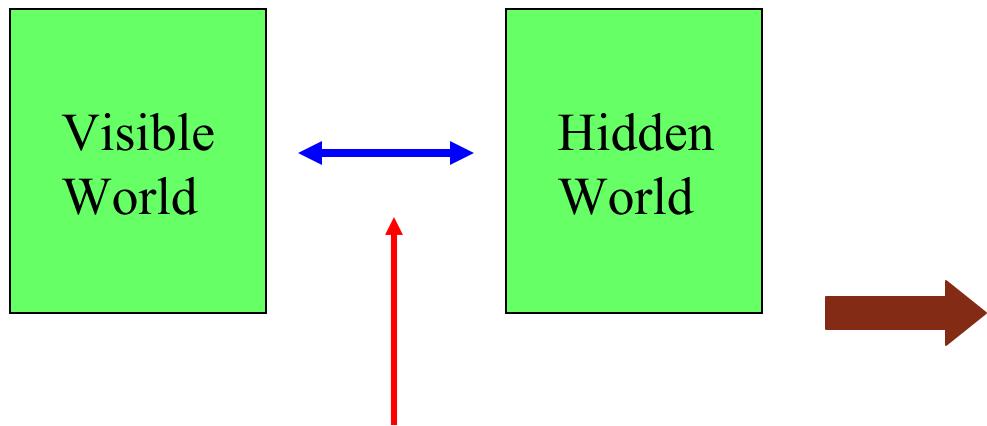
$$-0.009 \pm \Delta r_\mu^{SUSY} \leq 0.003$$

$m_t \dots$



SUSY theory must explain why no SUSY particles yet seen

SUSY Breaking



Flavor-blind mediation

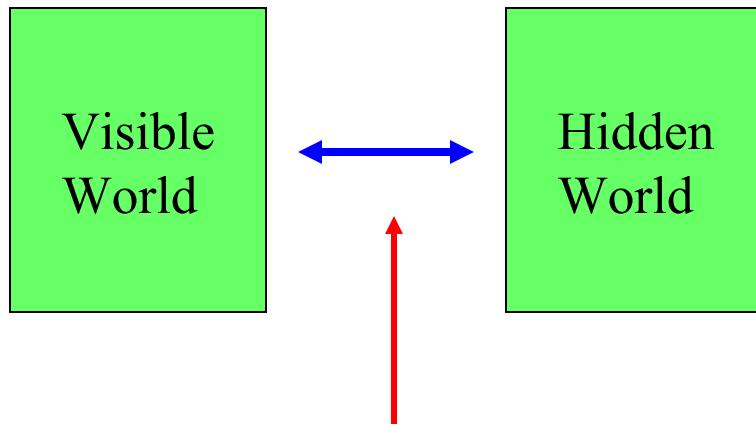
$$M_{\tilde{q}} > M_{\tilde{\ell}}$$

~ 5 parameters

$$\begin{aligned} M_{\tilde{e}} &>> m_e \\ M_{\tilde{q}} &>> m_q \\ M_{\tilde{\chi}} &>> M_{W,Z,\gamma} \end{aligned}$$

SUSY Theory clashes with Experiment

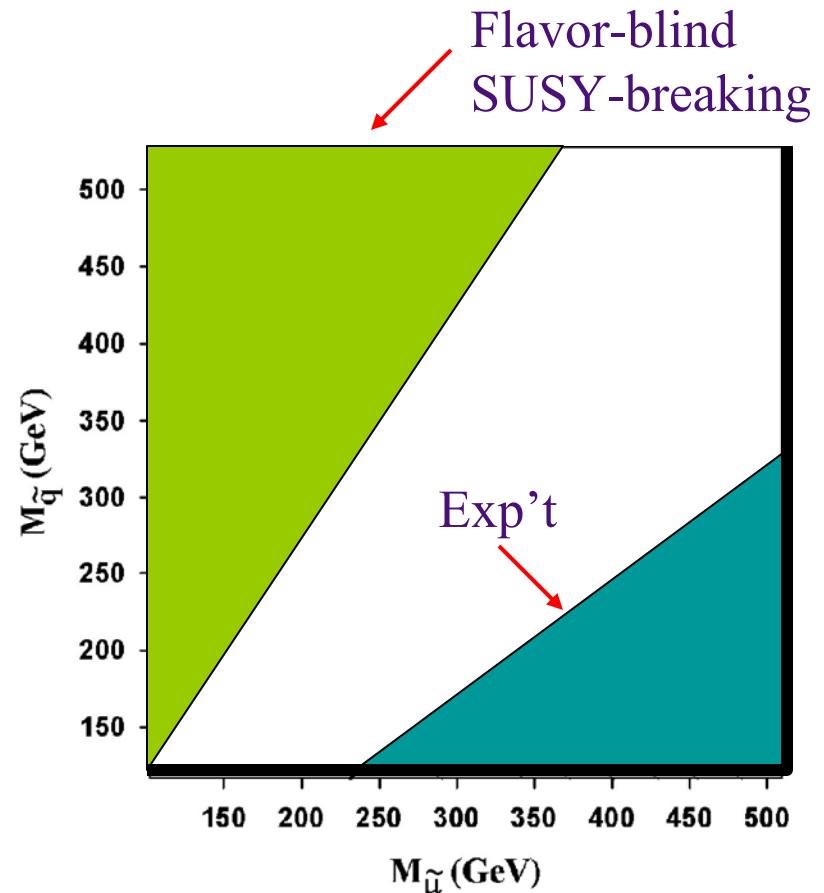
SUSY Breaking



Flavor-blind mediation

$$M_{\tilde{q}} > M_{\tilde{\ell}}$$

Kurylov & R-M



$$M_{\tilde{\mu}L} > M_{\tilde{q}L}$$

CC (non) universality & the MSSM

Data appear to conflict with MSSM & models having flavor-blind SUSY-breaking mediation

Possible resolutions:

1. $M_{\text{SUSY}} > \text{TeV}$
2. Hadronic effects in SM predictions
3. Something is wrong with exp't
4. New models of SUSY-breaking
5. Go beyond the MSSM

New ideas needed !

LANSCE, NIST, SNS...

SUSY irrelevant to low energy observables

Better control of non-pQCD

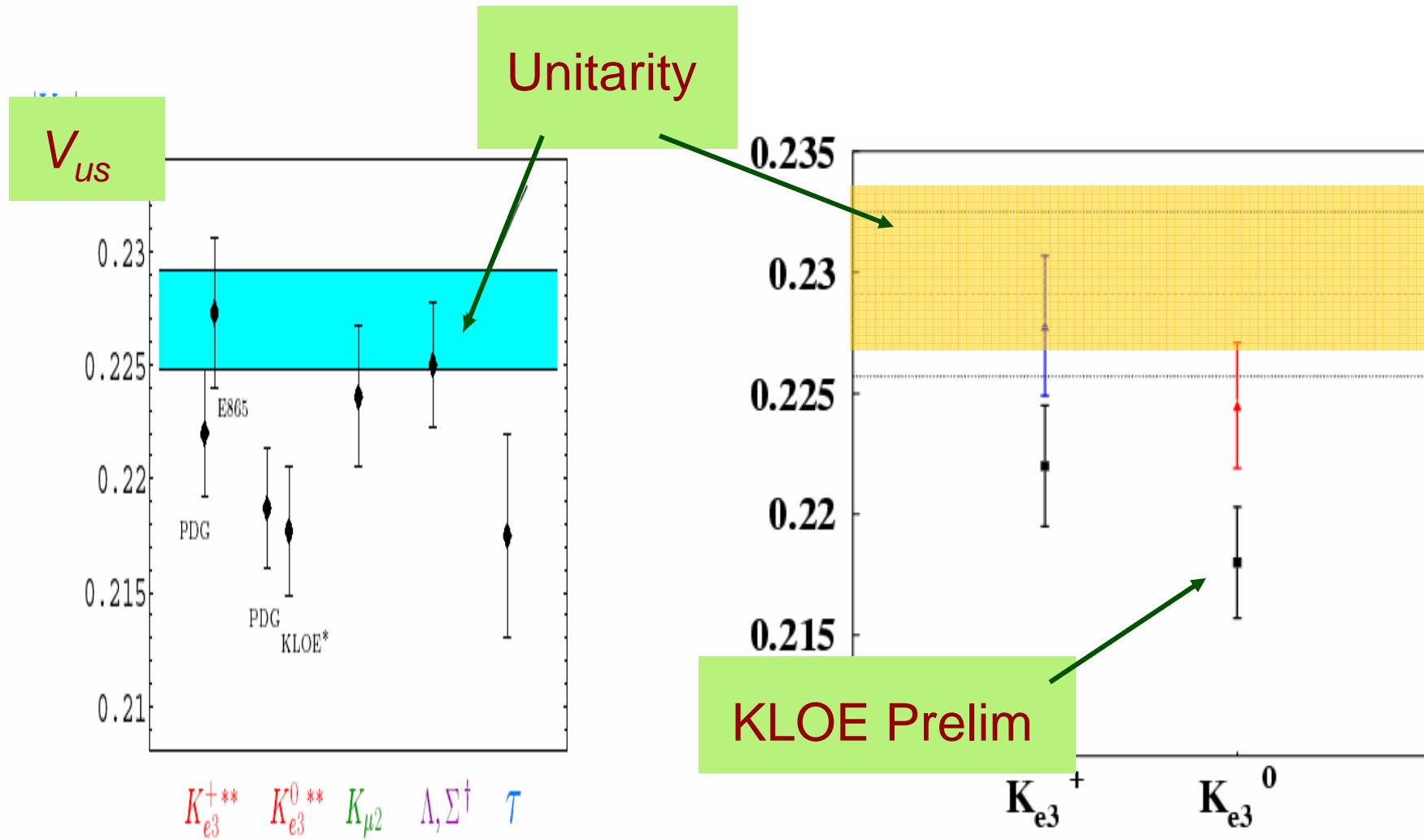
Exp't: $n \beta\text{-decay}$, K_{e3} decay

???

Break R-parity

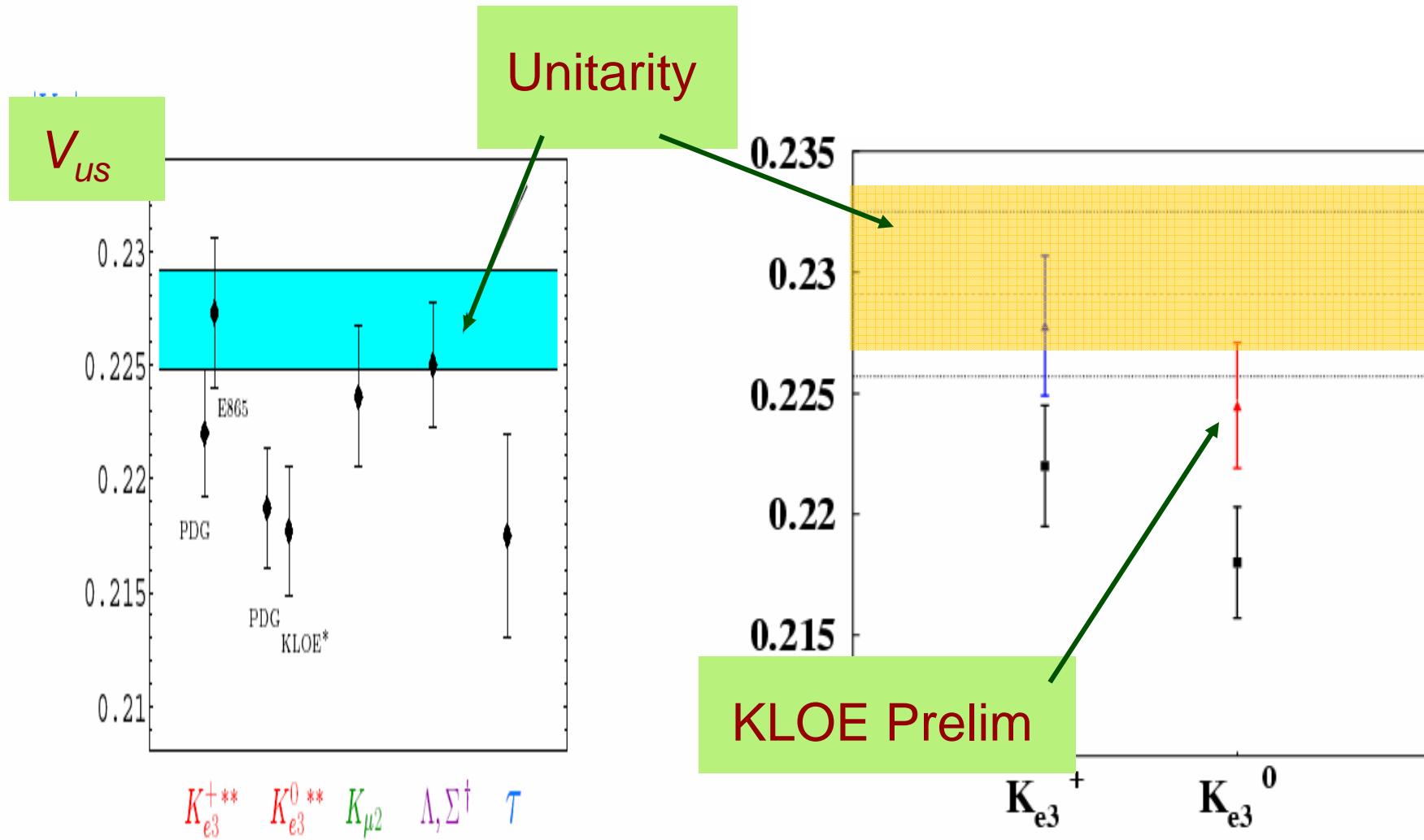
But NuTeV

Kaon Decays & V_{us}



Cirigliano & De Lucia

Kaon Decays & V_{us}



CC (non) universality & the MSSM

Data appear to conflict with MSSM & models having flavor-blind SUSY-breaking mediation

Possible resolutions:

1. $M_{\text{SUSY}} > \text{TeV}$
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SUSY irrelevant to low energy observables

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Exp't: $n \beta$ -decay, K_{e3} decay

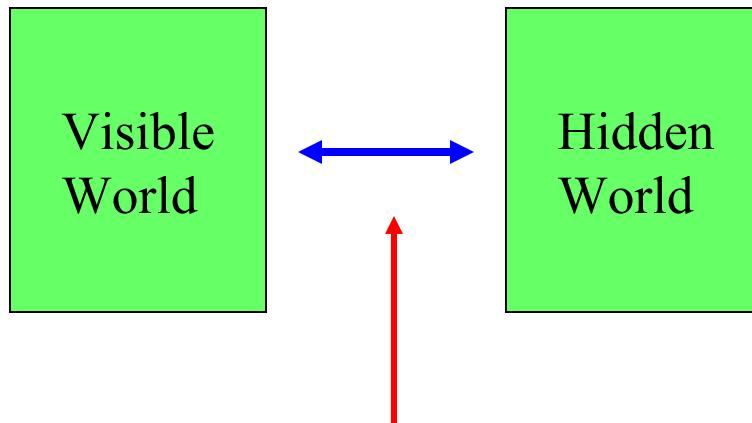
???

Break R-parity

But NuTeV

SUSY Theory OK if no SUSY dark matter

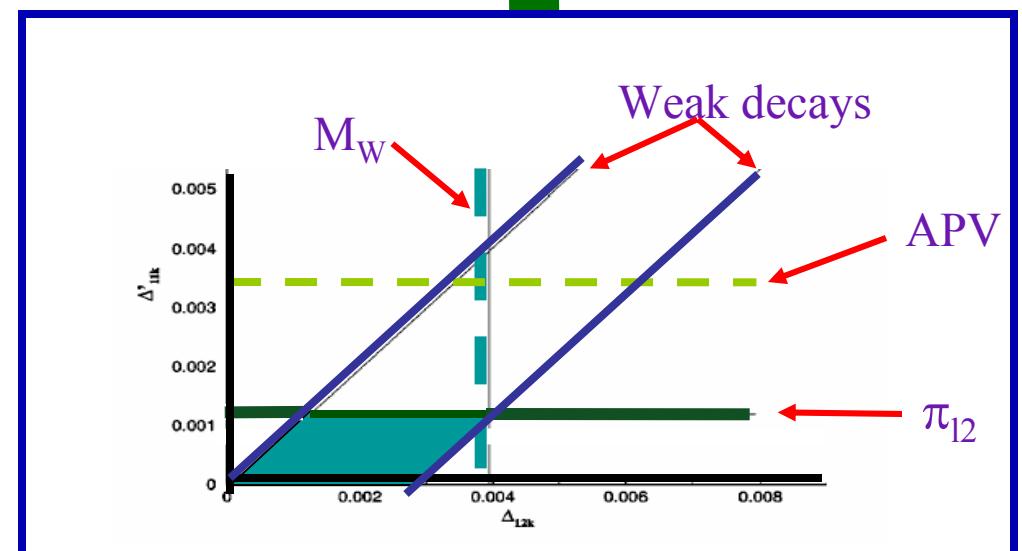
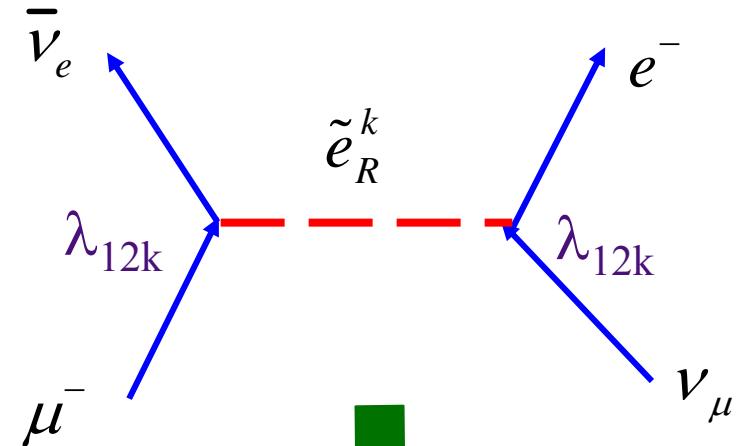
SUSY Breaking



Flavor-blind mediation

$$\chi^0 \rightarrow e^- \mu^+ \bar{\nu}_e$$

“R parity violation” (RPV)



R-Parity Violation (RPV)

$$W_{RPV} = \lambda_{ijk} L^i L^j E^k + \lambda'_{ijk} L^i Q^j D^k + \mu'_i L^i H_u$$

$\Delta L=1$

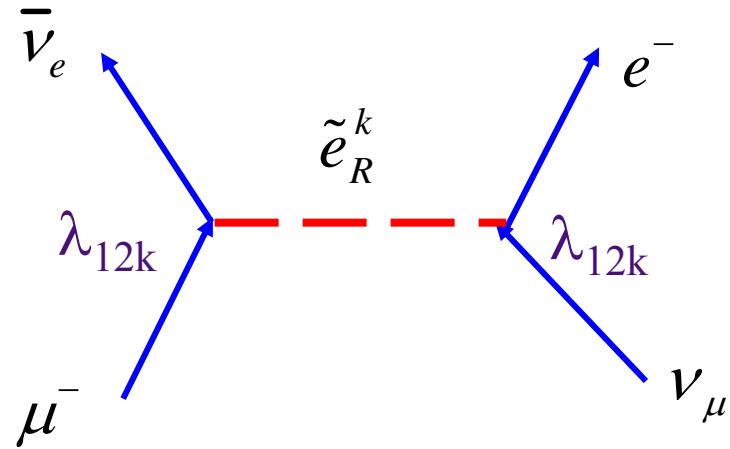
$$+ \lambda''_{ijk} U^i D^j D^k$$

$\Delta B=1$ proton decay:
Set $\lambda''_{ijk} = 0$

L^i, Q^i $SU(2)_L$ doublets

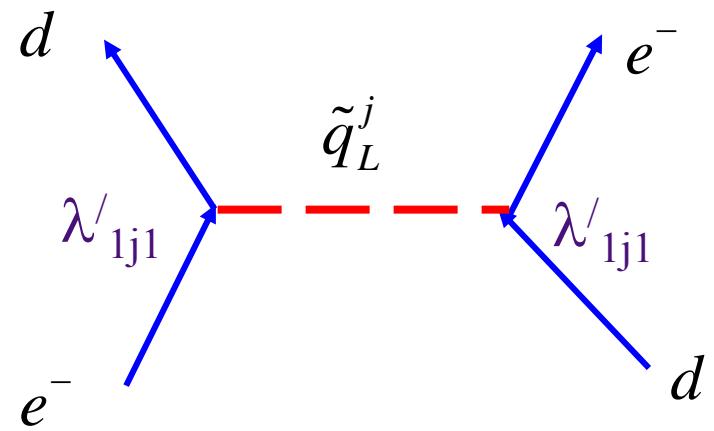
E^i, U^i, D^i $SU(2)_L$ singlets

Four-fermion Operators



$\Delta L=1$

$$\Delta_{12k} = \frac{|\lambda_{12k}|^2}{4\sqrt{2}G_F M_{\tilde{e}_R^k}^2}$$



$\Delta L=1$

$$\Delta'_{1j1} = \frac{|\lambda'_{ijj}|^2}{4\sqrt{2}G_F M_{\tilde{q}_L^j}^2}$$