

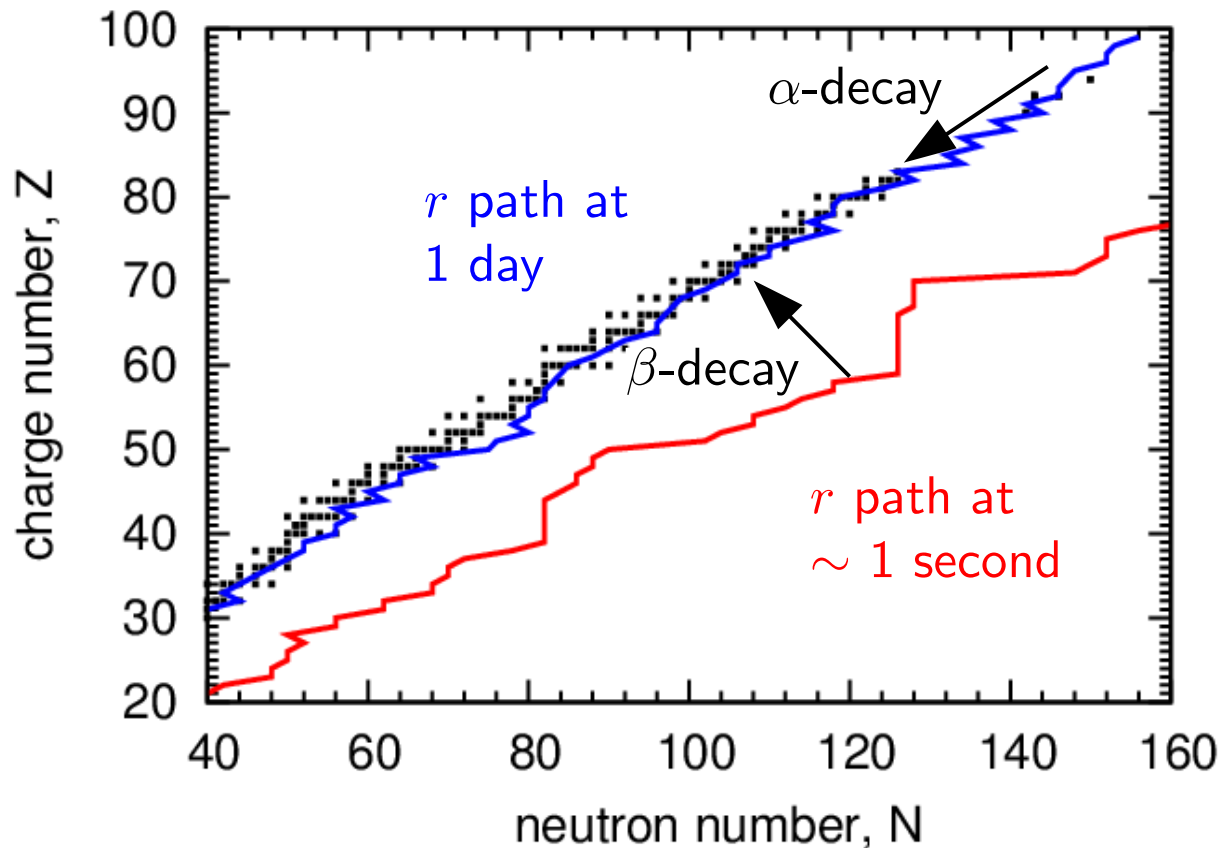
r-process nucleosynthesis yields and their heating rates

Meng-Ru Wu (Institute of Physics, Academia Sinica)

INT-JINA Symposium: First multi-messenger observations of a neutron star merger and its implications for nuclear physics
March 12 – March 14, 2017, INT, Seattle

Radioactive decay of the r -process nuclei

The nuclear energy release from the radioactive decay of the r -process nuclei is the source powering the kilonovae [e.g., Arnett's law: $L(t_{\text{peak}}) = \dot{\epsilon}(t_{\text{peak}})$]



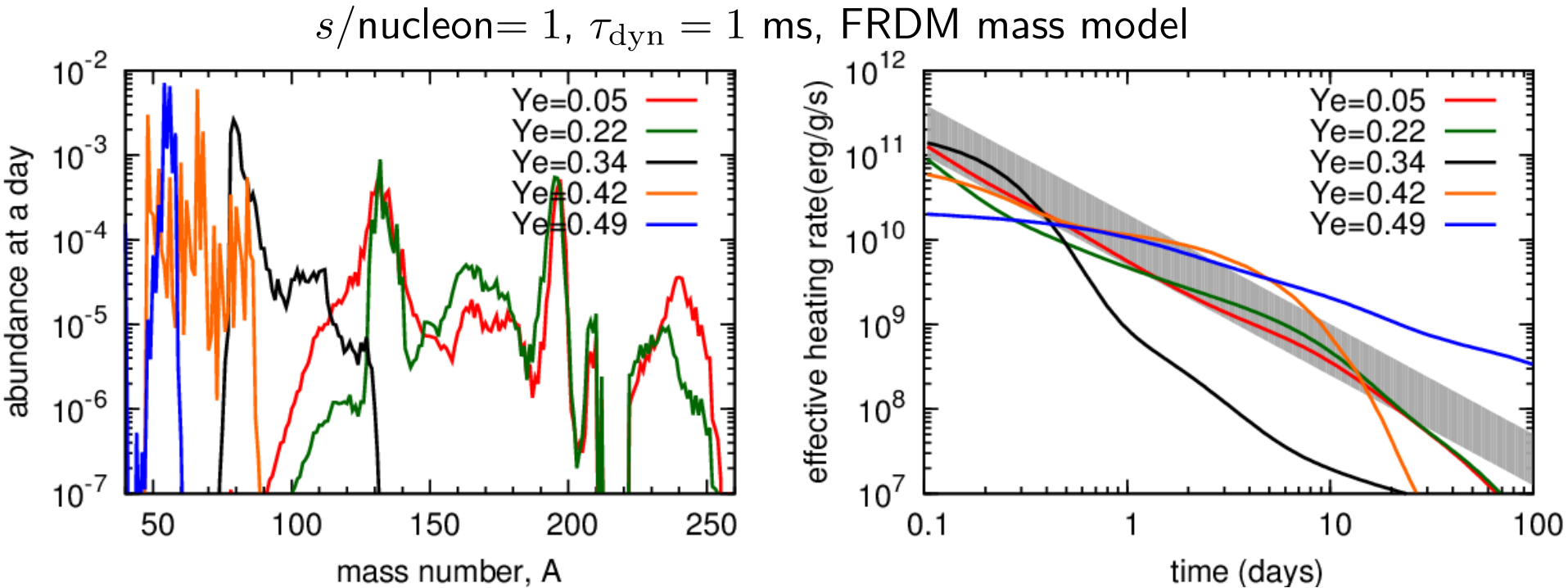
At \sim a day, the nuclear properties, e.g., the half-lives, the Q -value, are known. However, their abundances are determined by the initial astrophysical conditions and the unknown properties of the neutron-rich nuclei.

Abundances and heating rates

- Parametrized trajectories characterized by the initial Y_e , entropy s , and expansion timescale τ_{dyn} at 6 GK (see e.g., Lippuner & Roberts 2016).

Abundances and heating rates

- Parametrized trajectories characterized by the initial Y_e , entropy s , and expansion timescale τ_{dyn} at 6 GK (see e.g., Lippuner & Roberts 2016).



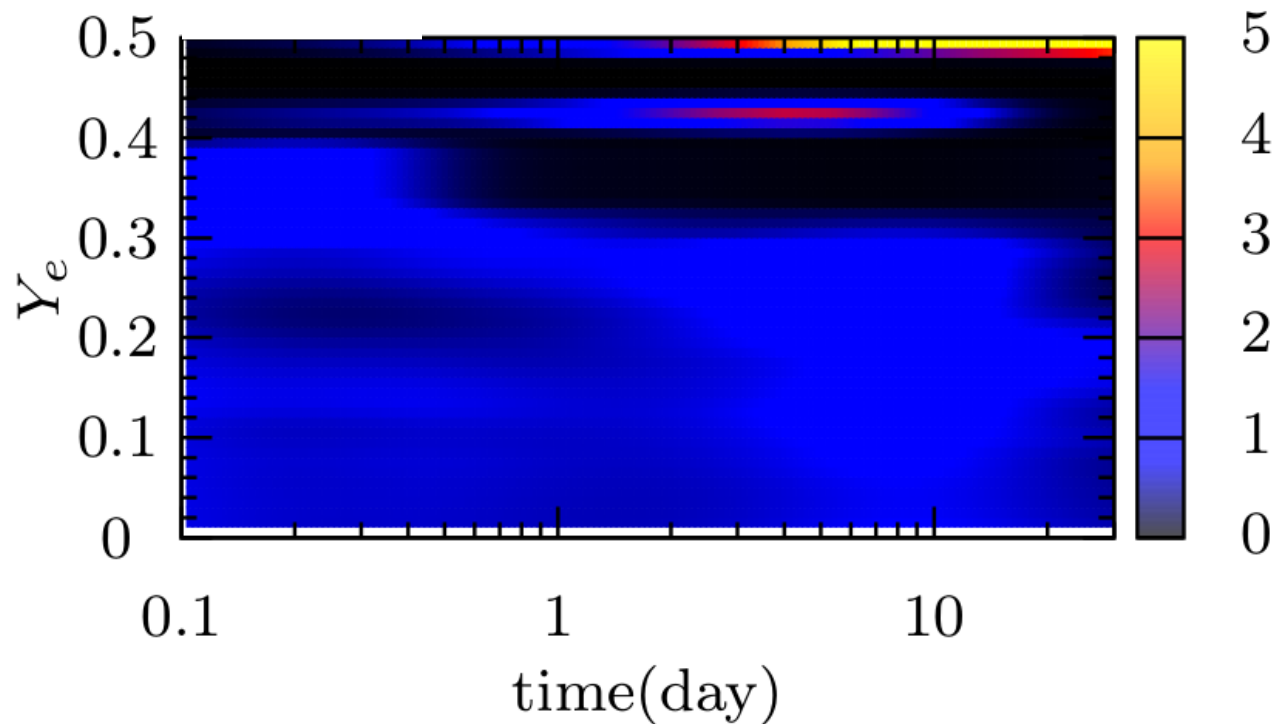
$$\text{Effective heating rate} = 0.25 \times \dot{Q}_\beta + \dot{Q}_\alpha + \dot{Q}_{\text{fis}}$$

$$\text{Analytical heating rate} = 10^{10} \times t_d^{-1.3} \text{ erg/g/s}$$

Y_e dependence

Colors: ratio between the effective heating rate and the analytical heating rate

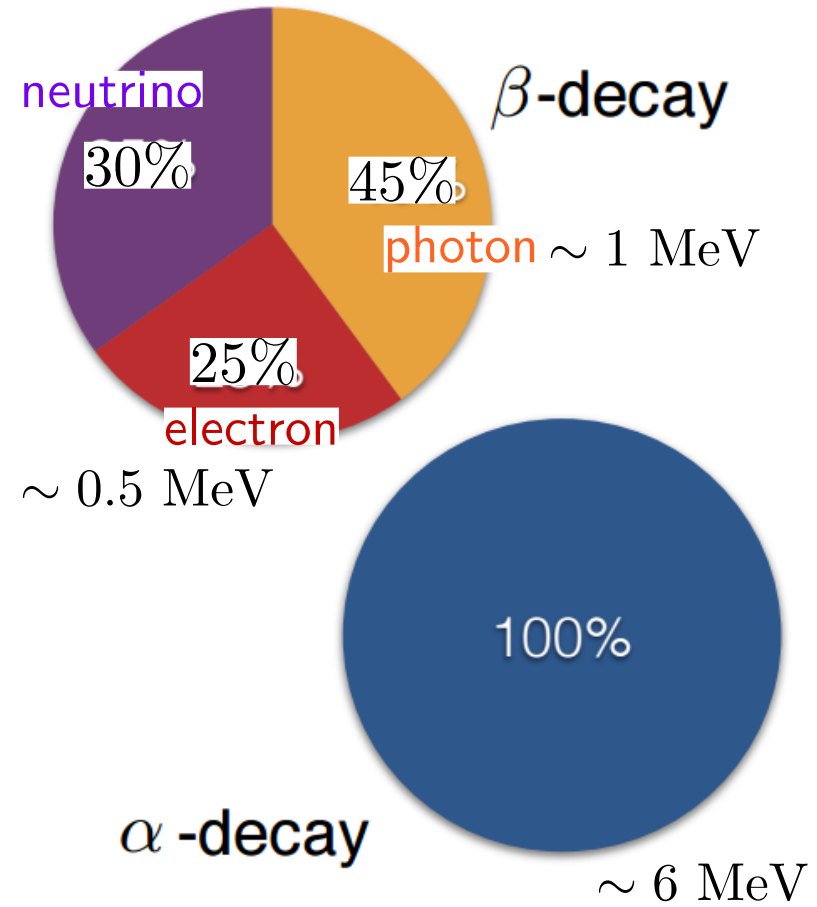
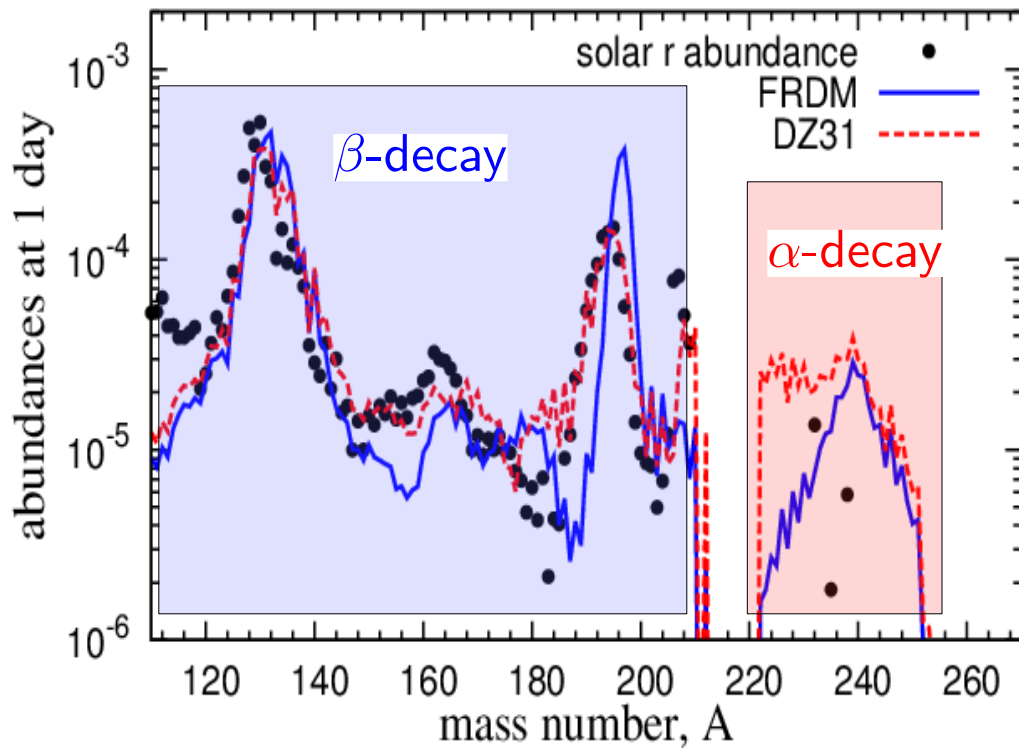
$s/\text{nucleon} = 1$, $\tau_{\text{dyn}} = 1 \text{ ms}$, FRDM mass model



- relative little heating power for $Y_e \gtrsim 0.3$
- more or less agree with the analytical heating rate for $Y_e \lesssim 0.3$

Nuclear physics impact on kilonova heating

For ejecta with very low $Y_e \lesssim 0.1$, initially cold ($s \sim 1$ per nucleon), different nuclear mass inputs can give rise to large abundance difference for $220 \lesssim A \lesssim 240$

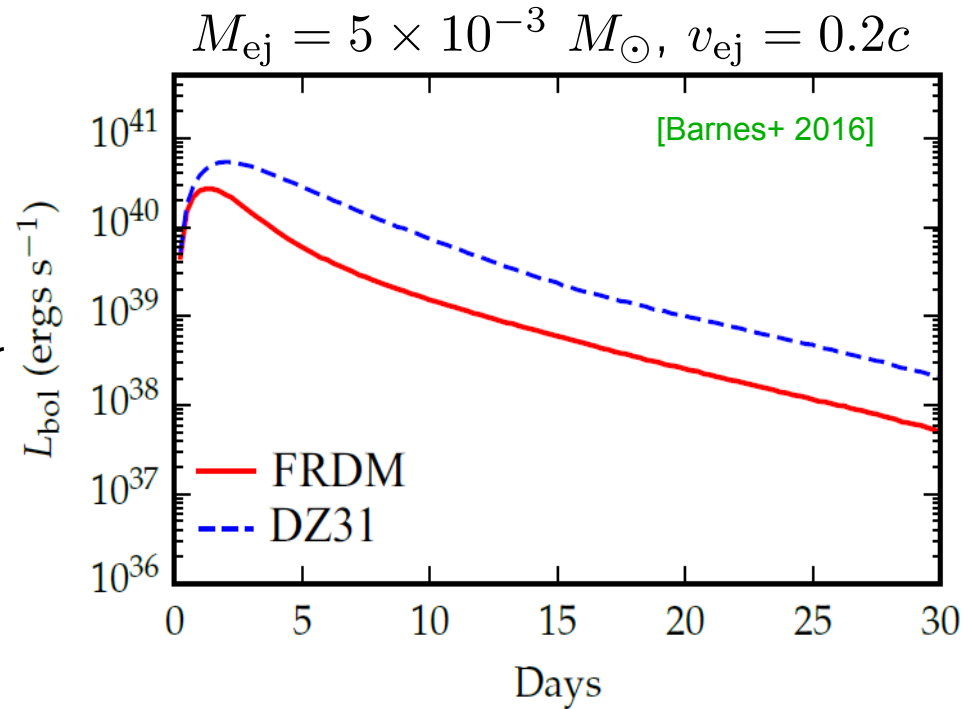


[Barnes+ 2016]

Nuclear physics impact on kilonova heating

- α decays may release similar amount of energy as β decays per second, sensitively depending on the adopted nuclear mass model
- α & β particles thermalize in a similar way while γ -ray thermalization quickly become inefficient

[Hotokezaka+2016, Barnes+2016]

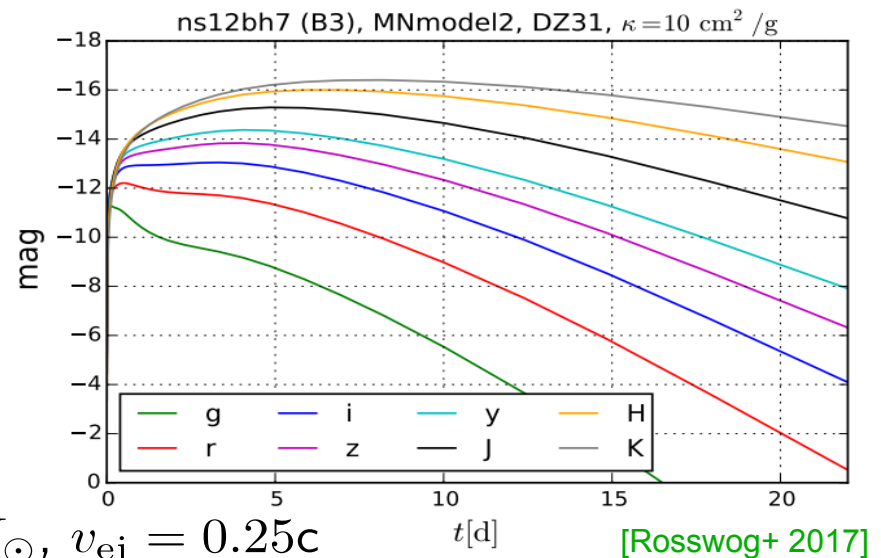
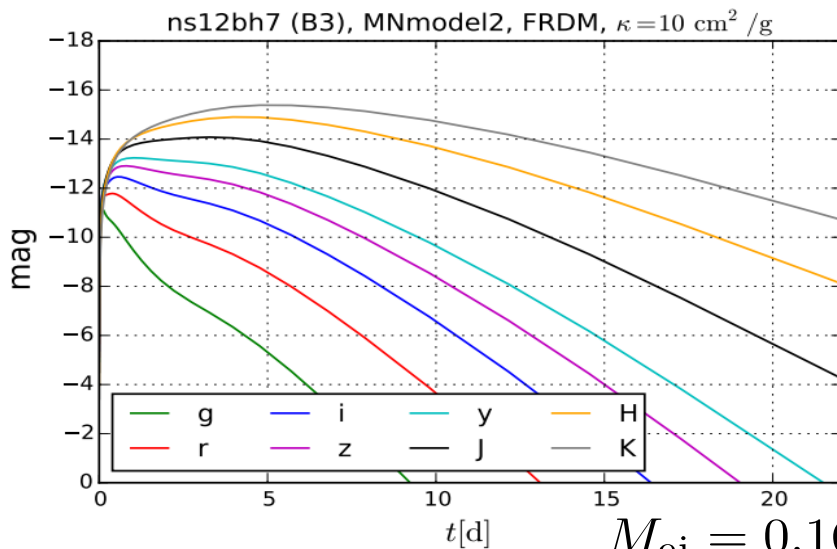
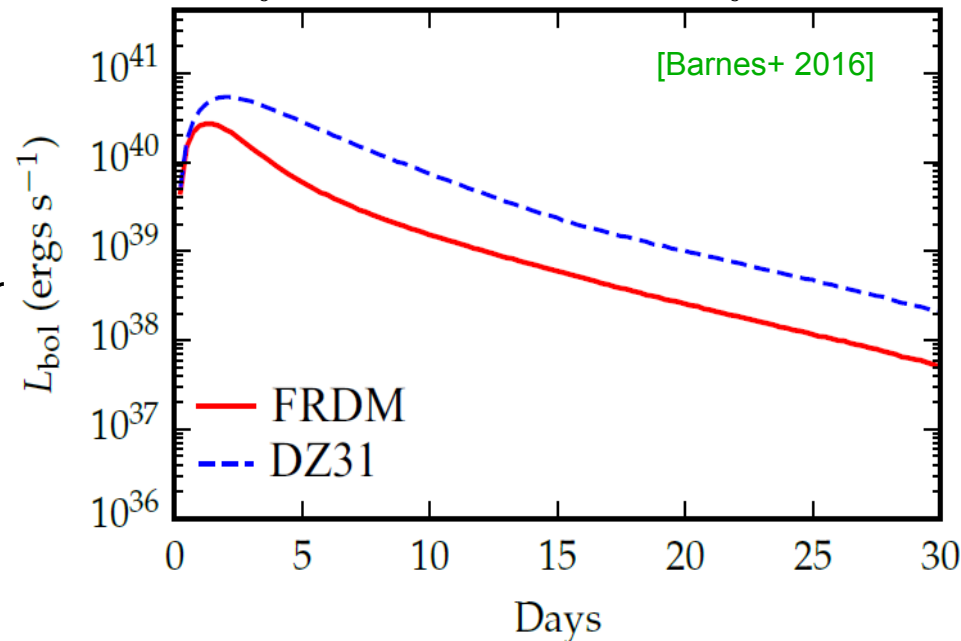


Nuclear physics impact on kilonova heating

- α decays may release similar amount of energy as β decays per second, sensitively depending on the adopted nuclear mass model
- α & β particles thermalize in a similar way while γ -ray thermalization quickly become inefficient

[Hotokezaka+2016, Barnes+2016]

$$M_{\text{ej}} = 5 \times 10^{-3} M_{\odot}, v_{\text{ej}} = 0.2c$$



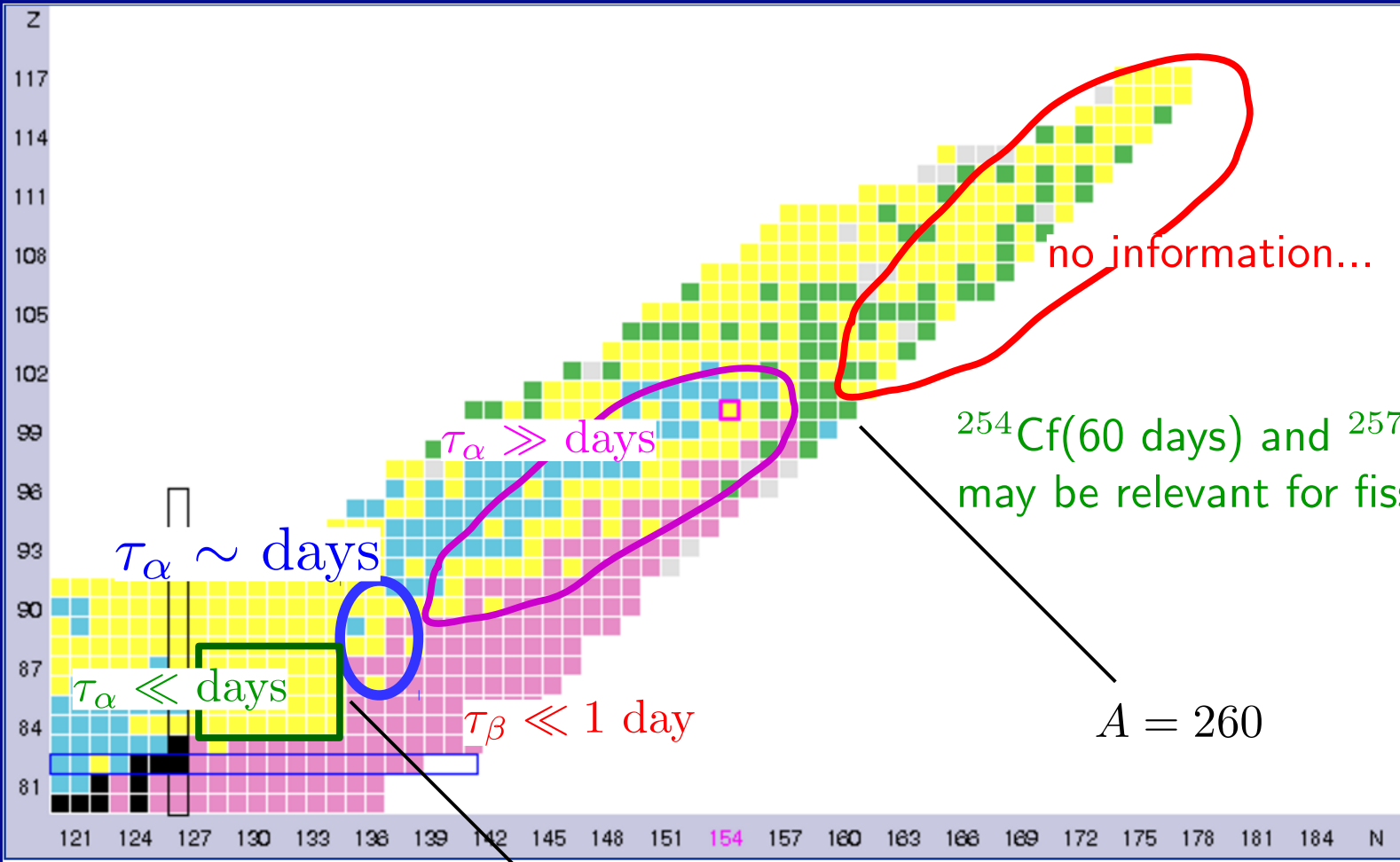
what heavy nuclei can be relevant for kilonova heating?



Chart of Nuclides

Click on a nucleus for information

Color code	Half-life	Decay Mode	Q_{β^-}	Q_{EC}	Q_{β^+}	S_n	S_p	Q_α	S_{2n}	S_{2p}	$Q_{2\beta^-}$	Q_{2EC}	Q_{ECp}
$Q_{\beta-n}$	BE/A	(BE-LDM Fit)/A	$E_{1st\ ex. st.}$	E_{2+}	E_{3-}	E_{4+}	E_{4+}/E_{2+}	β_2	$B(E2)_{42}/B(E2)_{20}$	$\sigma(n,\gamma)$	$\sigma(n,F)$	235U FY	239Pu FY



Tooltips
 On
 Off

Zoom: 1, 2, 3, 4, 5, 6, 7
 Uncertainty: NDS, Standard
 Screen Size: Narrow, Wide

^{254}Cf (60 days) and ^{257}Es (7.7 days) may be relevant for fission

EC+ β^+
 β^-
 α
 P
 N
 SF
 Unknown

Search options:
 Levels and Gammas
 Nuclear Wallet Cards

$A = 220$

$A = 260$

production of actinides

How neutron-rich may the α -decay be important?

Initially, $Y_n^0 = 1 - Y_e$, $Y_p^0 = Y_e$

Assuming all protons are locked in seed nuclei, right before n-captures

$$R_{n/s} = \frac{Y_n}{Y_{\text{seed}}} = \frac{Y_n^0 - N_{\text{seed}}(Y_p^0/Z_{\text{seed}})}{(Y_p^0/Z_{\text{seed}})}$$

$$\langle A \rangle_{\text{final}} \approx A_{\text{seed}} + R_{n/s}$$

To have nuclei with $A \sim 220 - 230$ that may alpha decay at relevant time,

$$\langle A \rangle_{\text{final}} \approx 190$$

With $N_{\text{seed}} \approx 50$, $Z_{\text{seed}} \approx 30$, it gives $Y_e = 0.16$.

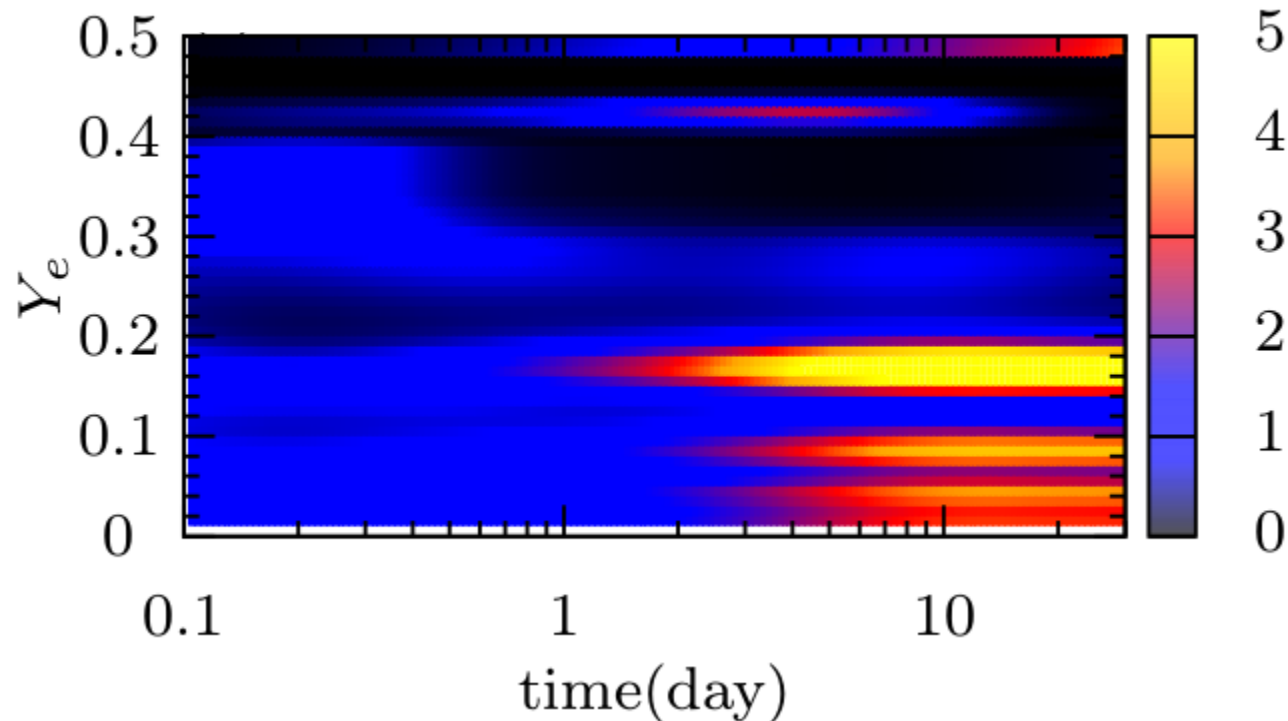
→ α -heating may become dominant for ejecta with $Y_e \lesssim 0.2$

Y_e dependence with DZ31 nuclear mass model

- vary the (n, γ) and (γ, n) rates for $35 \leq Z \leq 83$ [Mendoza-Temis+ 2015]

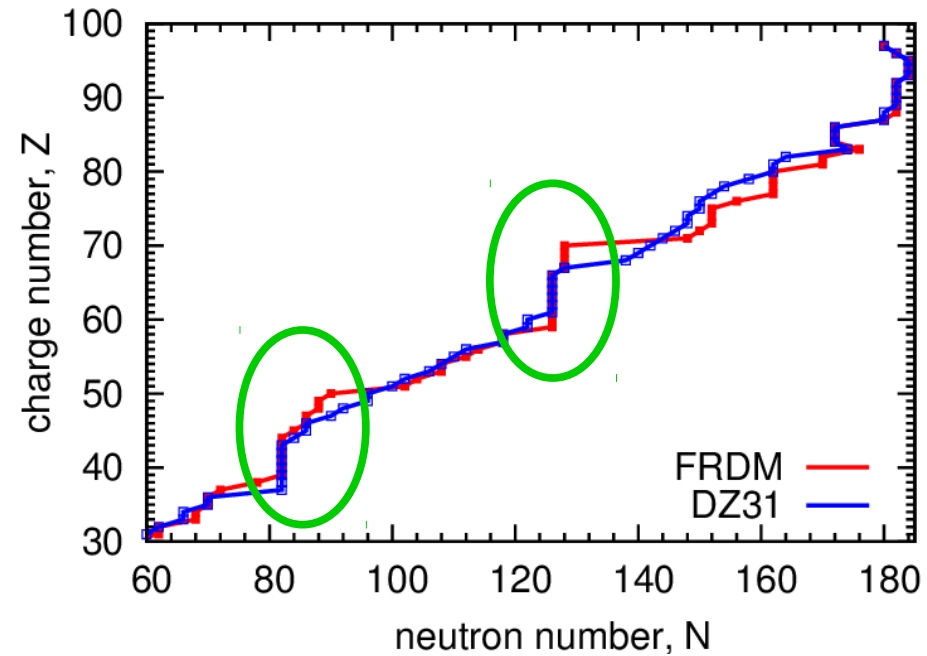
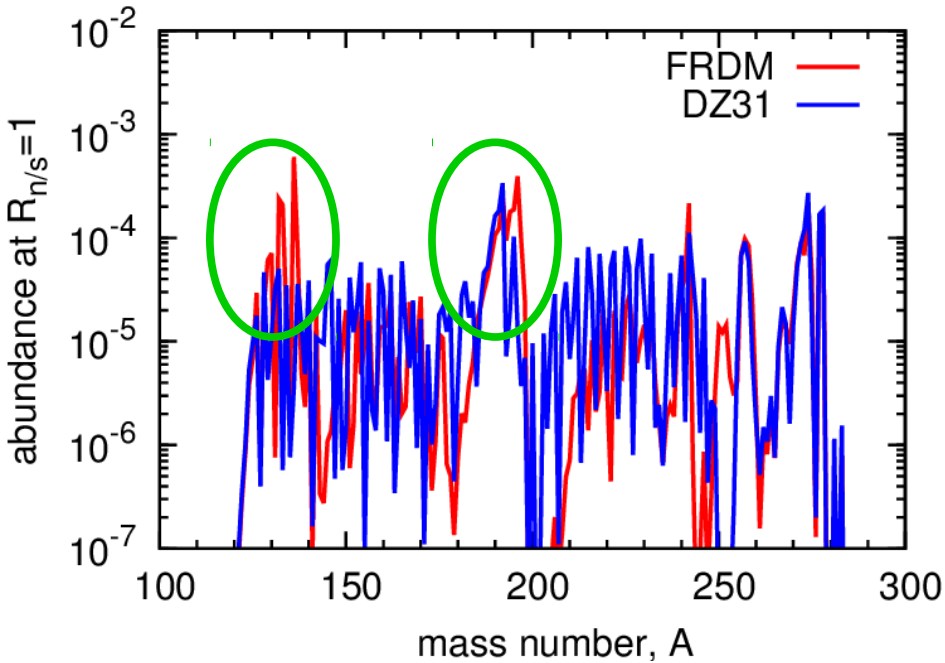
Colors: ratio between the effective heating rate and the analytical heating rate

$s/\text{nucleon} = 1$, $\tau_{\text{dyn}} = 1$ ms, DZ31 mass model



- the heating rate nearly unaffected for $Y_e \gtrsim 0.25$
- large enhancement due to α -decay up to $Y_e \lesssim 0.2$ for $t \gtrsim 3$ days
- the heating rate enhancement oscillates as Y_e changes

FRDM v.s. DZ31

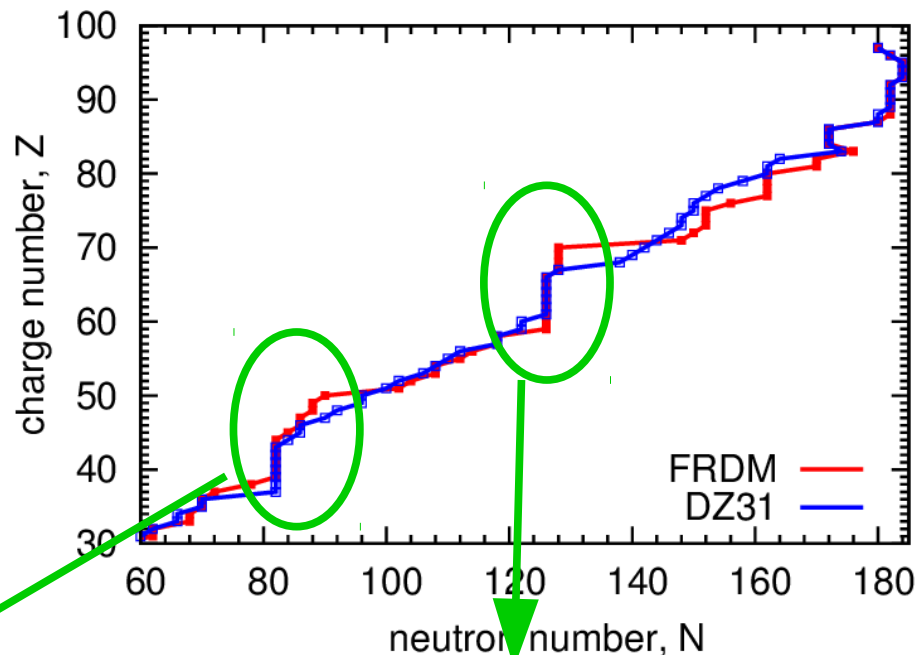
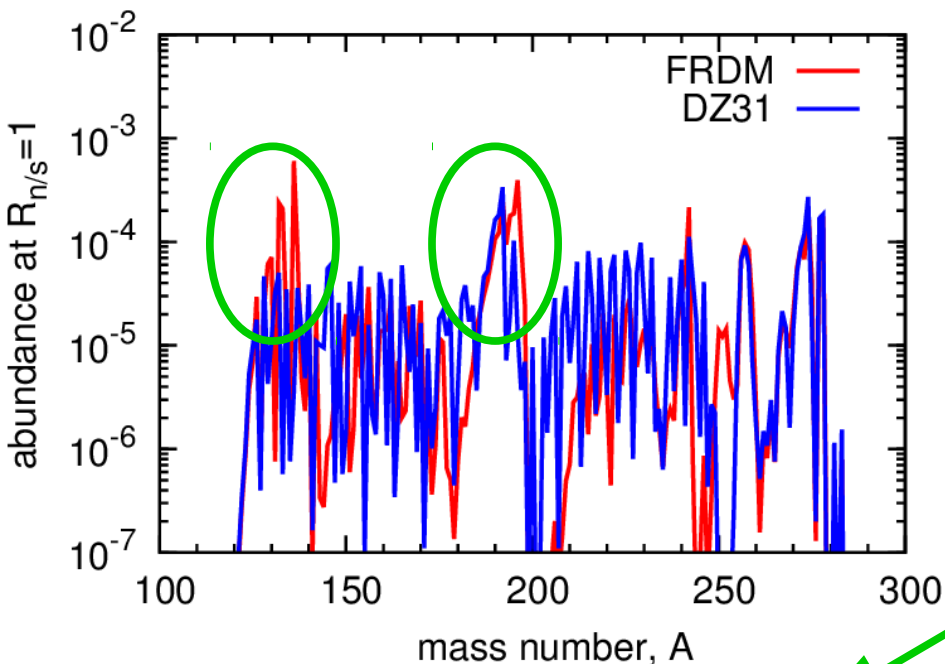


the r -process path with FRDM mass model extends to larger Z for regions slightly above the shell closure

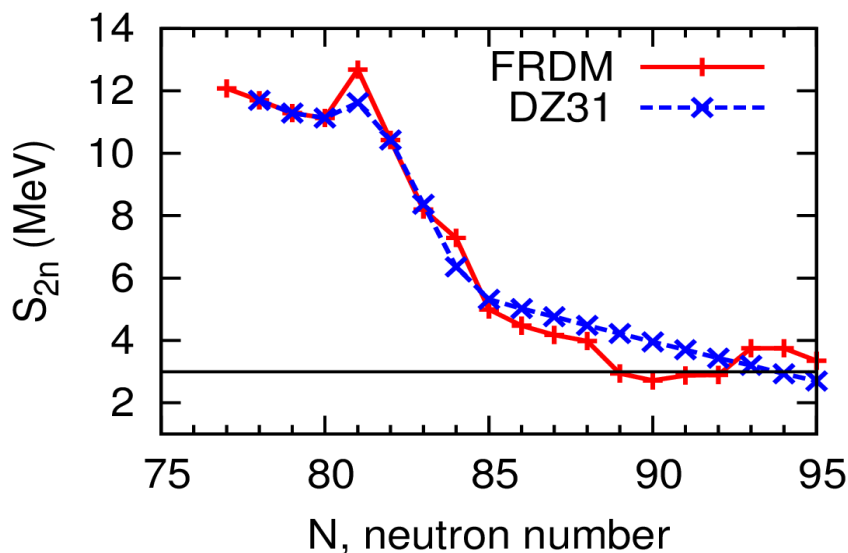
→ longer β -decay half-lives

→ larger abundances at peaks and smaller abundances in-between the peaks

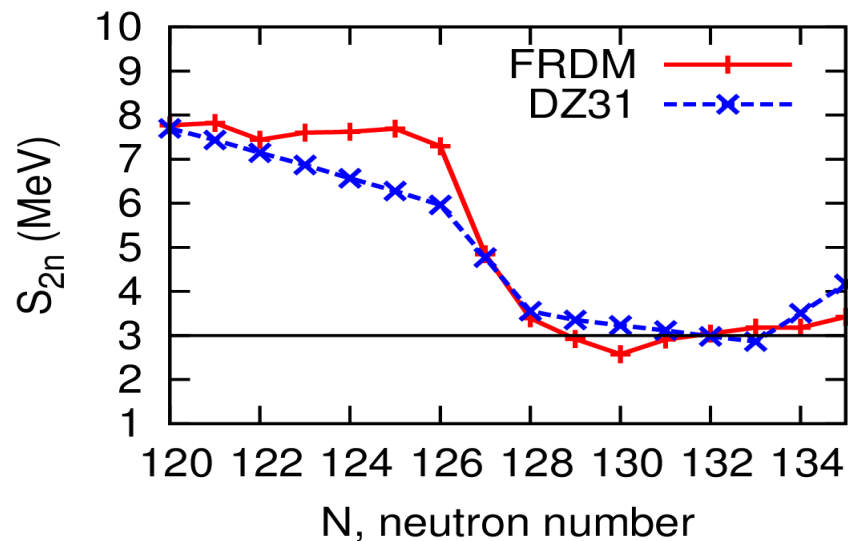
FRDM v.s. DZ31



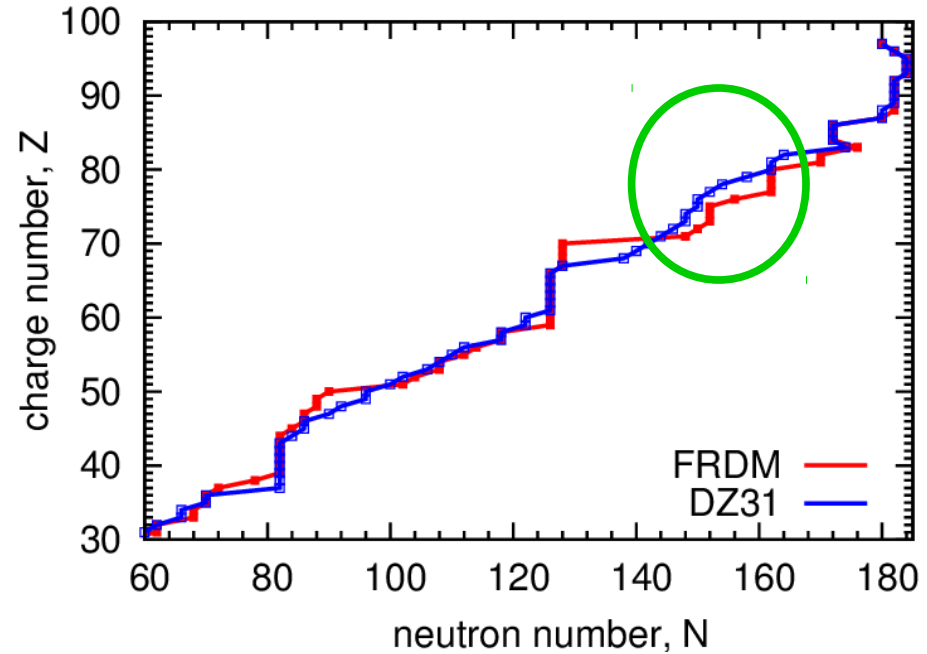
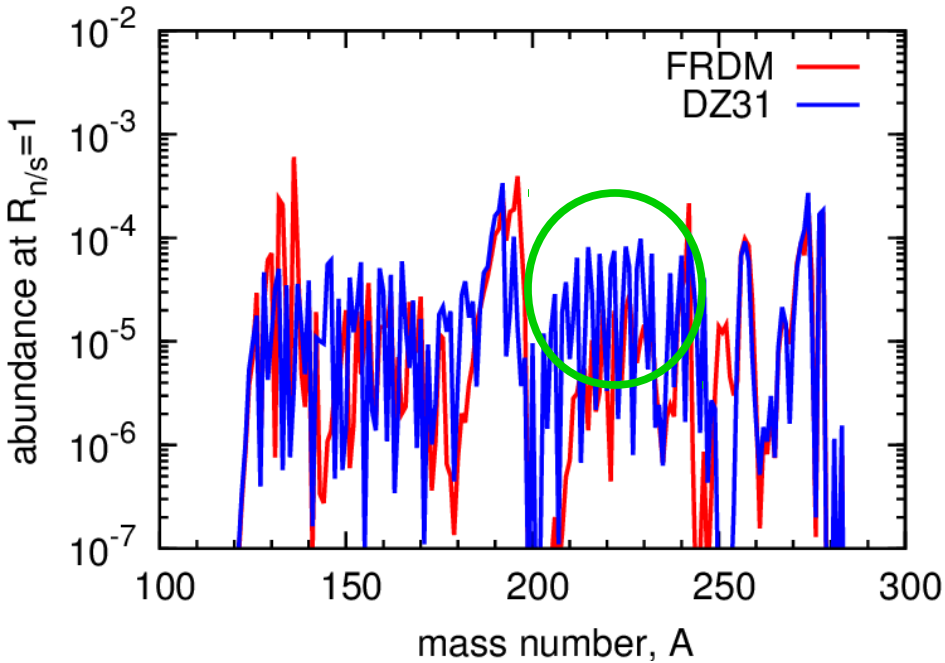
Cd ($Z = 48$) isotopic chain



Er ($Z = 68$) isotopic chain



FRDM v.s. DZ31

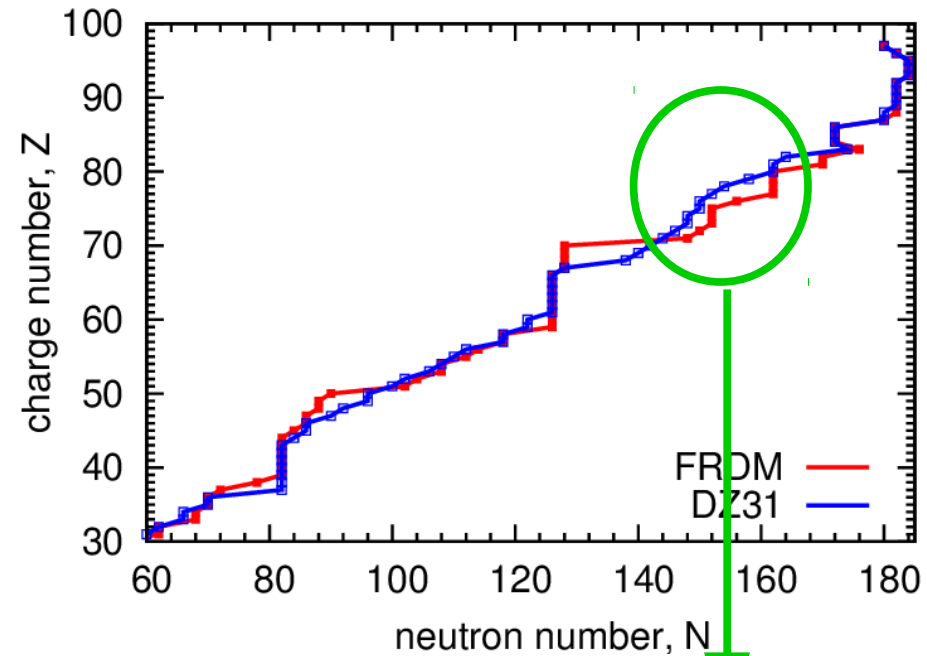
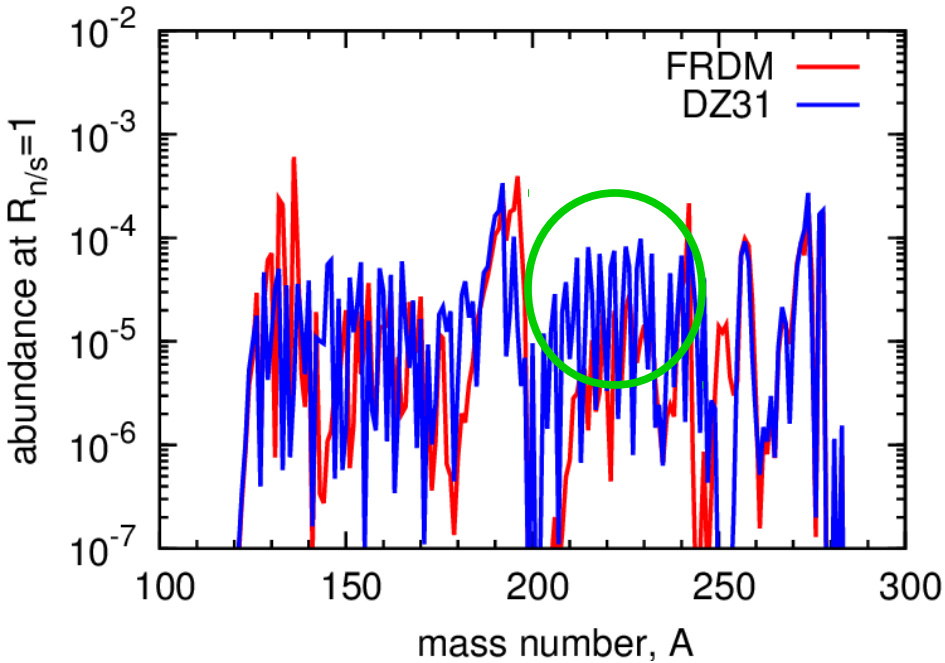


on the other hand, at the region around $A \sim 230$, the path of FRDM extends to larger N

→ shorter β -decay half-lives

→ smaller abundances

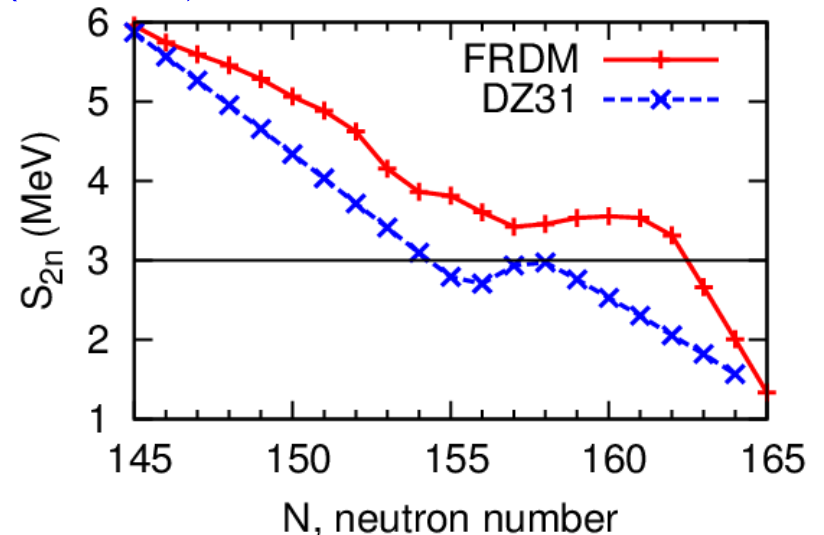
FRDM v.s. DZ31



on the other hand, at the region around $A \sim 230$, the path of FRDM extends to larger N

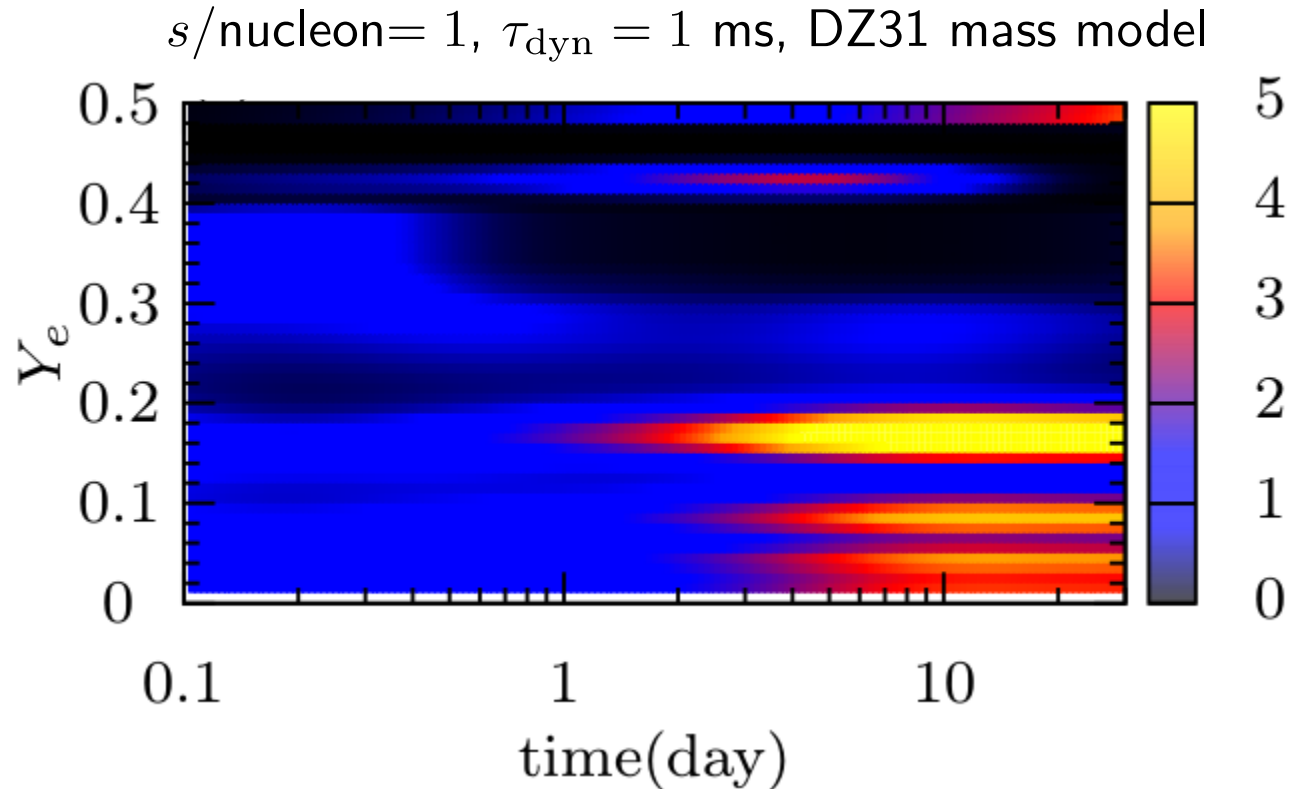
- shorter β -decay half-lives
- smaller abundances

Pt ($Z = 78$) isotopic chain



Y_e dependence with DZ31 nuclear mass model

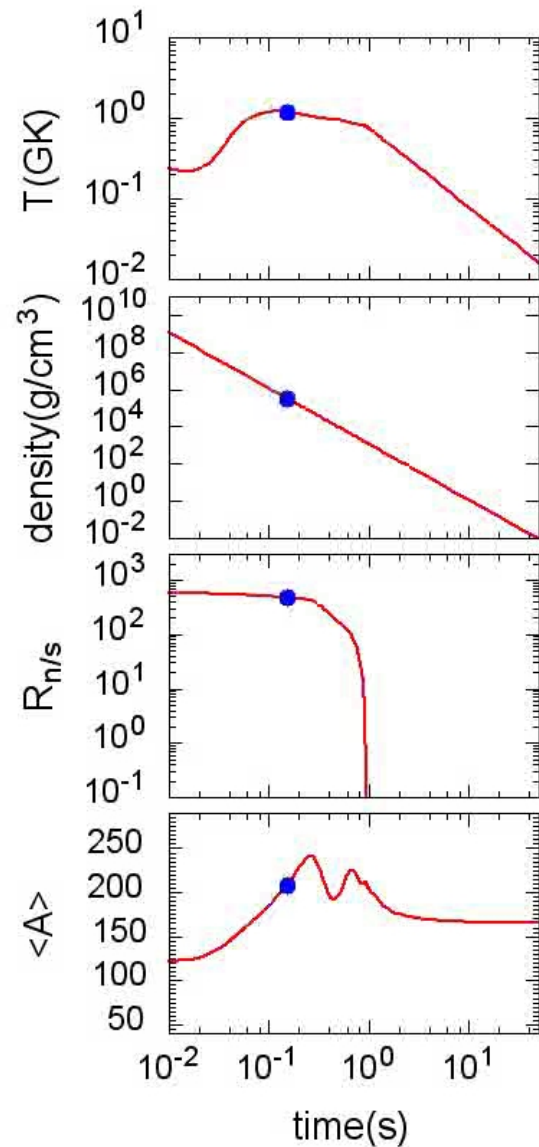
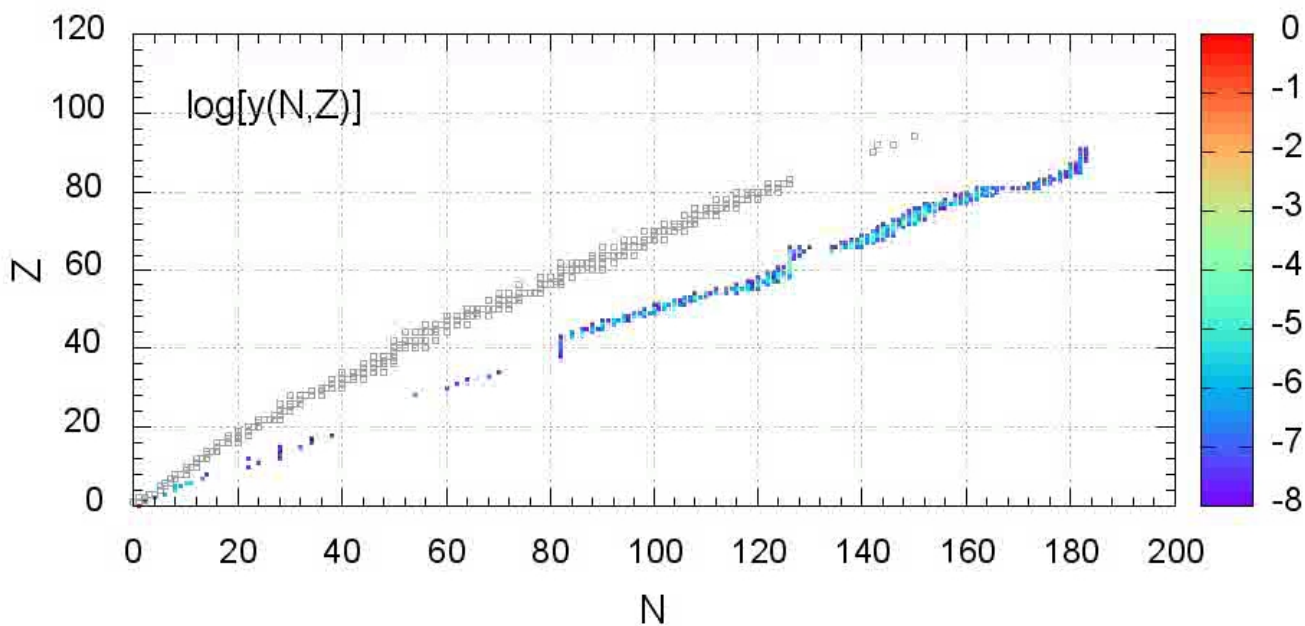
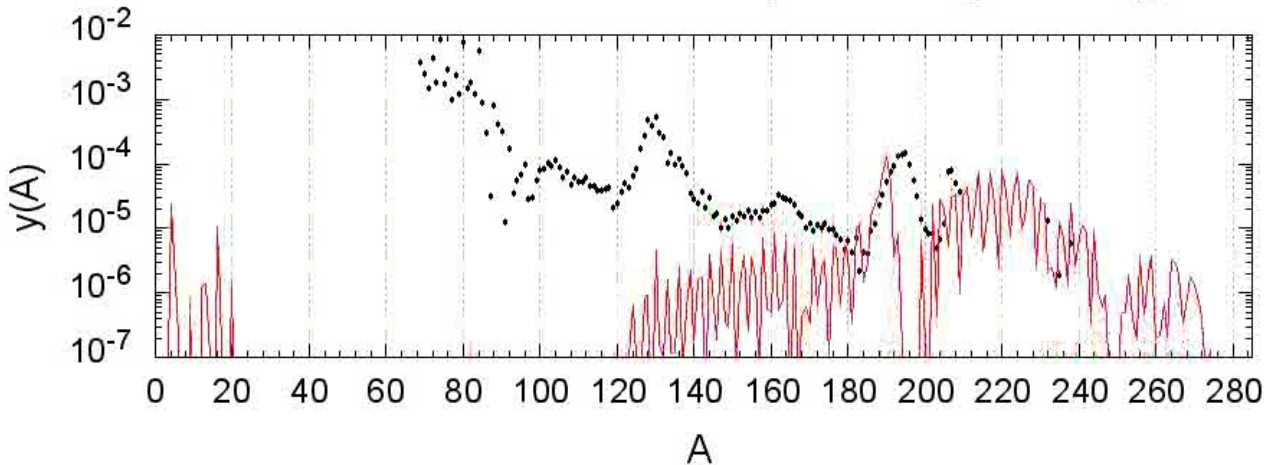
- vary the (n, γ) and (γ, n) rates for $35 \leq Z \leq 83$ [Mendoza-Temis+ 2015]



- large enhancement due to α -decay up to $Y_e \lesssim 0.2$ for $t \gtrsim 3$ days
- the heating rate enhancement oscillates as Y_e changes

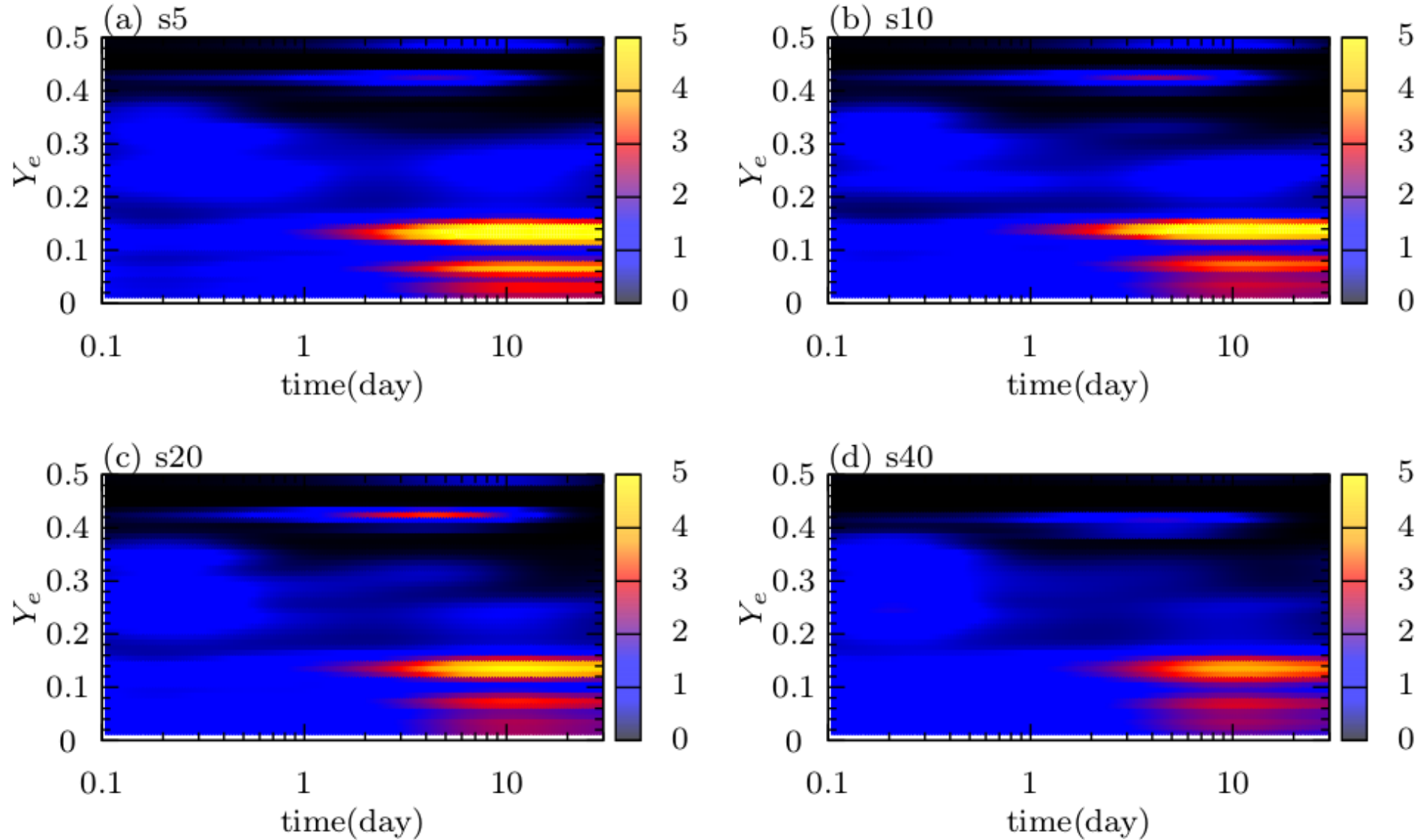
Y_e dependence with DZ31 nuclear mass model

time=1.46E-01 s, $T=1.20\text{E}+00$ GK, density= $3.63\text{E}+05$ g/cm³, $R_{n/s}=479.005$, $\langle A \rangle=206.1$, beta/(n,g)=0.00E+00



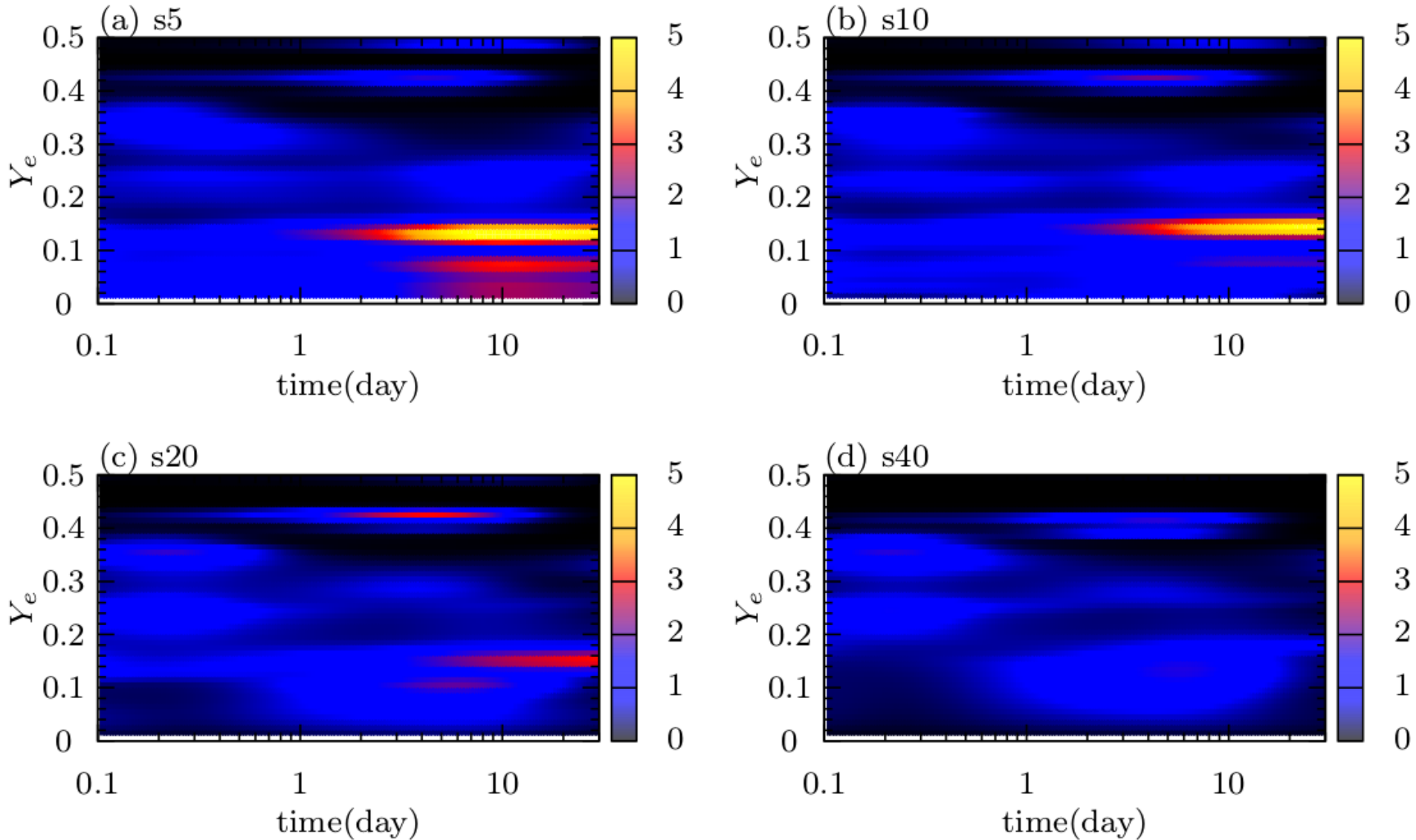
Y_e dependence with DZ31 nuclear mass model

- different entropy, for $\tau_{\text{dyn}} = 10$ ms



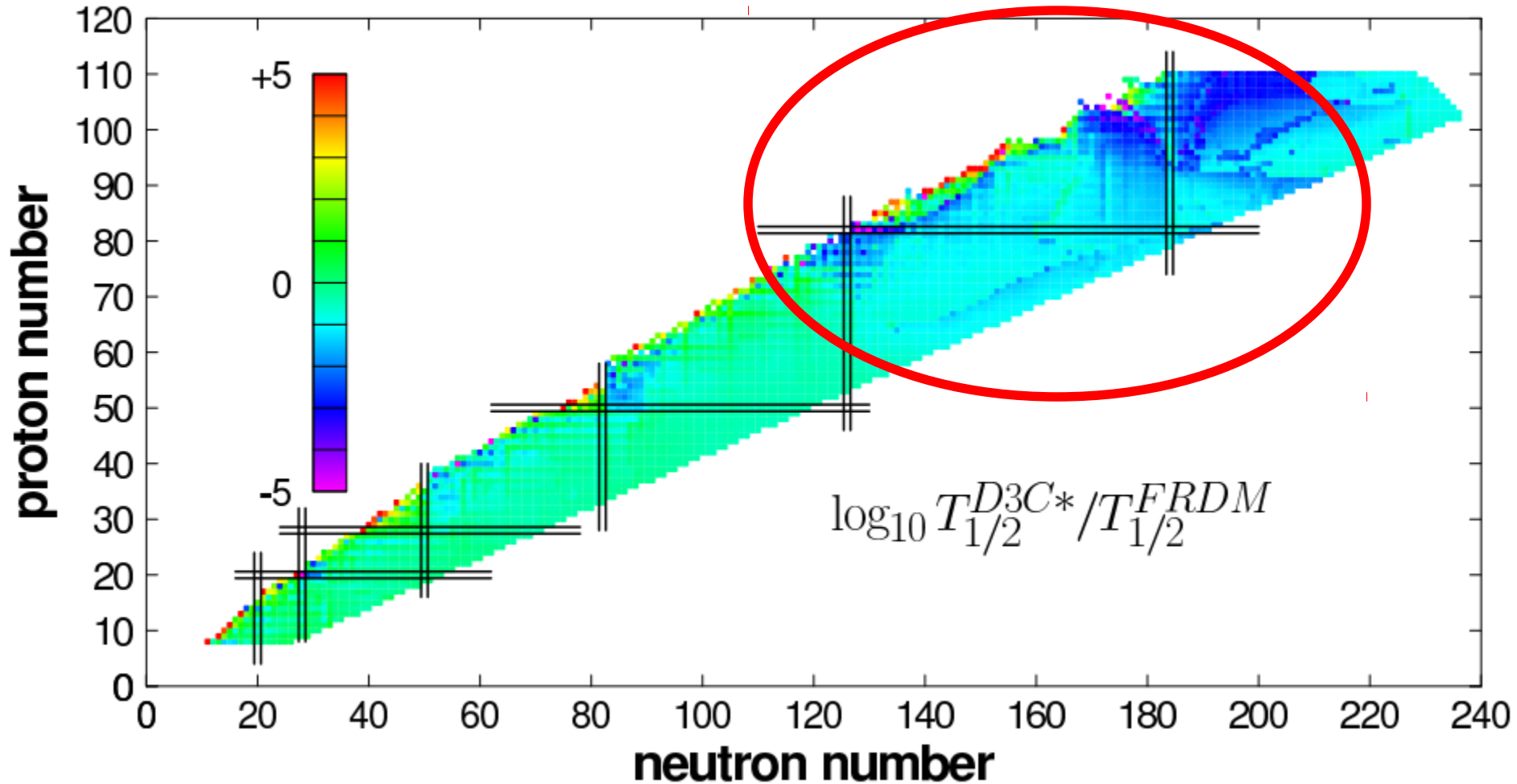
Y_e dependence with DZ31 nuclear mass model

- different entropy, for $\tau_{\text{dyn}} = 1$ ms



Impact of β -decay rates

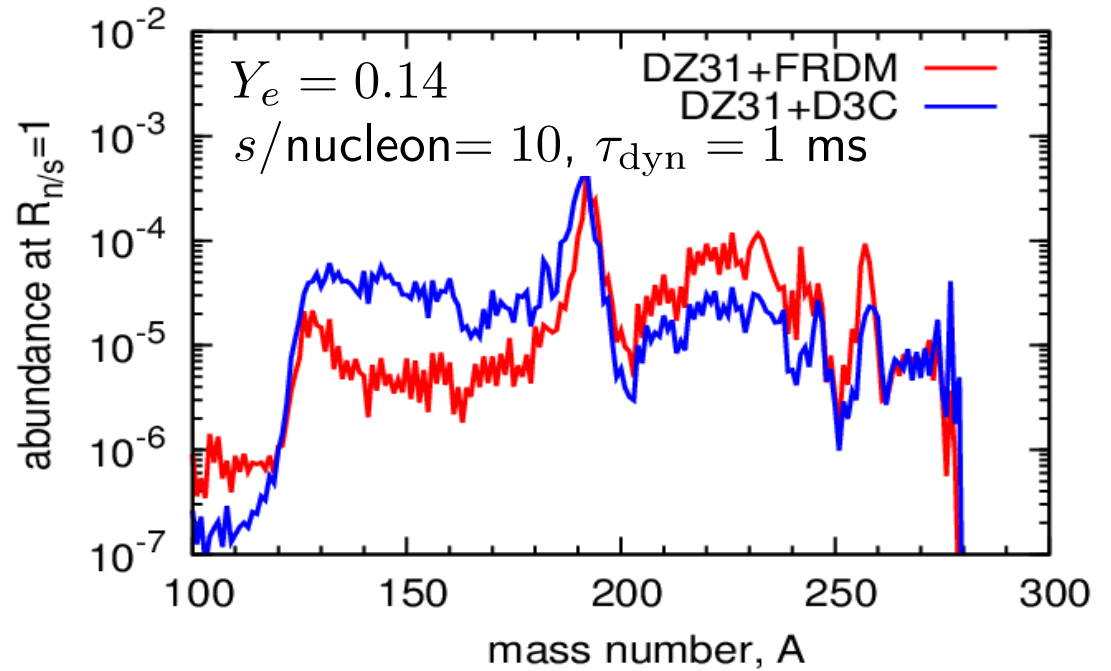
[Marketin+ 2015]



D3C* model predicts shorter τ_{β} for heavier nuclei

Impact of β -decay rates

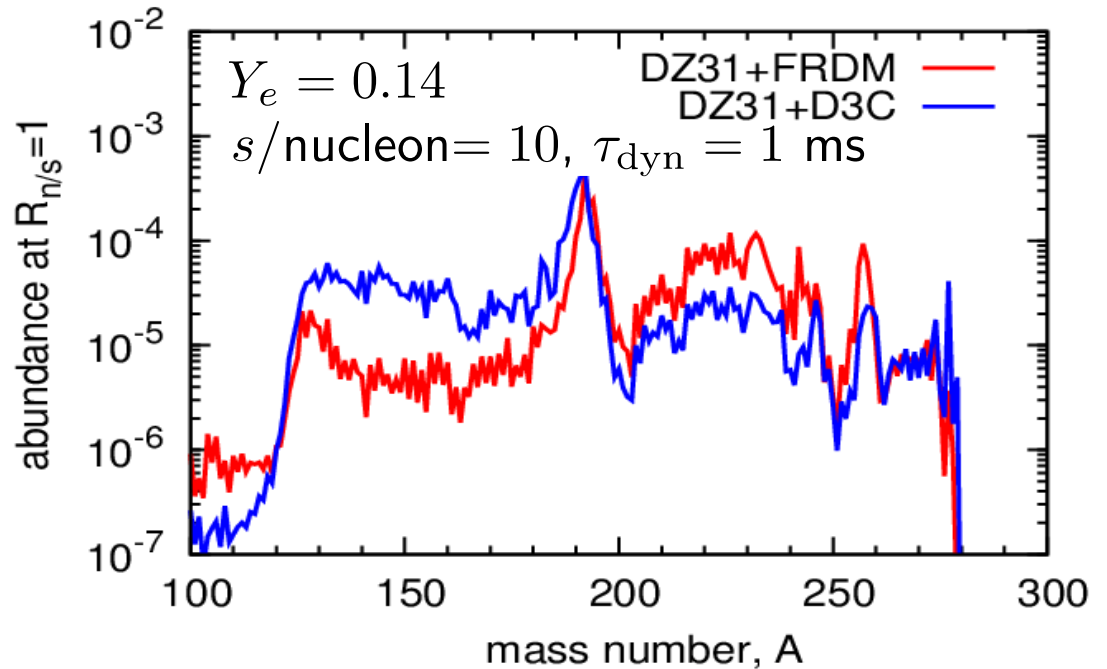
- more lanthanides and less actinides with D3C* β -decay half-lives



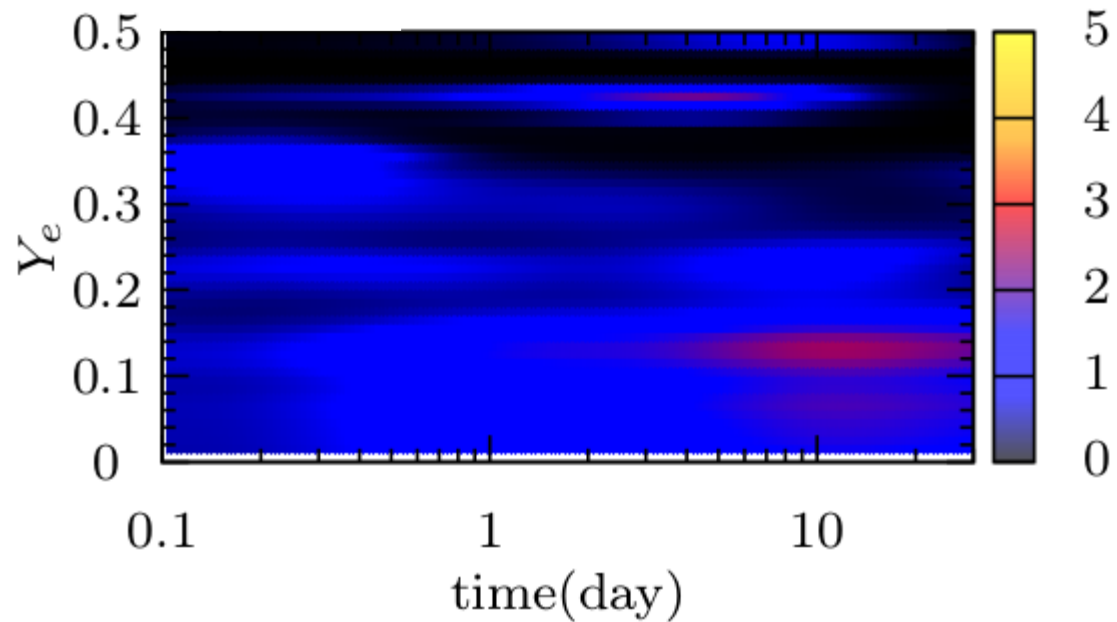
Impact of β -decay rates

- more lanthanides and less actinides with D3C* β -decay half-lives

- less energy release due to the α -decay



s/nucleon = 10, $\tau_{\text{dyn}} = 1$ ms, DZ31+D3C*



Summary

- The radioactive energy release during the kilonova time is insensitive to the nuclear physics inputs for $Y_e \gtrsim 0.25$.
- The nuclear heating rate can be sensitive to the composition, e.g., for ejecta with $Y_e \gtrsim 0.3$, the radioactive energy release is significantly smaller when compared to $Y_e \lesssim 0.3$.
- For $Y_e \lesssim 0.2$, the theoretical inputs of the radioactive heating rates are subject to the uncertain nuclear physics properties of neutron-rich nuclei and can vary by a factor of ~ 5 – 10 .

Summary

- The radioactive energy release during the kilonova time is insensitive to the nuclear physics inputs for $Y_e \gtrsim 0.25$.
 - The nuclear heating rate can be sensitive to the composition, e.g., for ejecta with $Y_e \gtrsim 0.3$, the radioactive energy release is significantly smaller when compared to $Y_e \lesssim 0.3$.
 - For $Y_e \lesssim 0.2$, the theoretical inputs of the radioactive heating rates are subject to the uncertain nuclear physics properties of neutron-rich nuclei and can vary by a factor of ~ 5 – 10 .
-
- more systematic sensitivity study is needed to better quantify the uncertainty
 - how to reduce the uncertainty?
 - how much does the composition-dependent heating rates affect the predicted kilonova lightcurves?