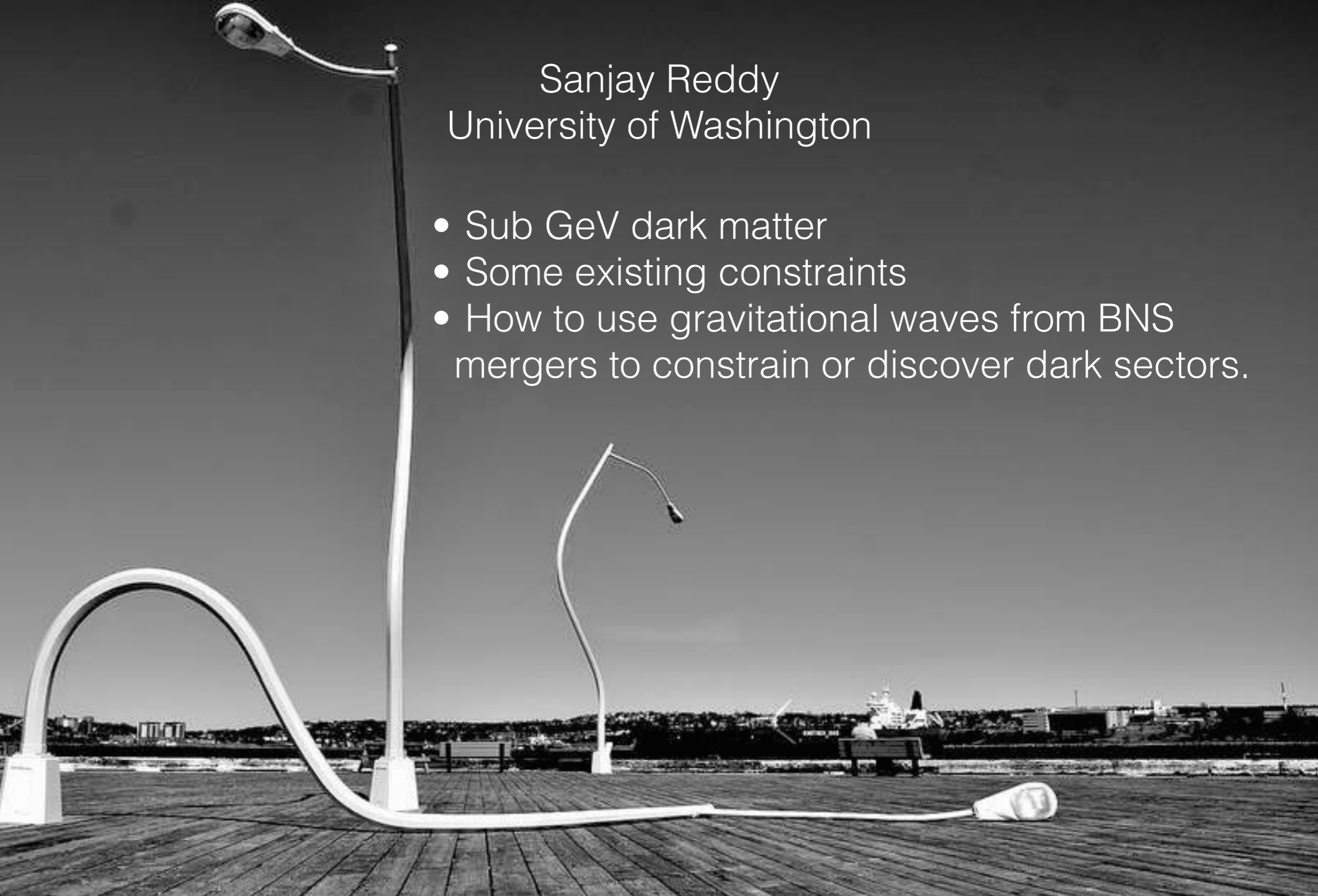


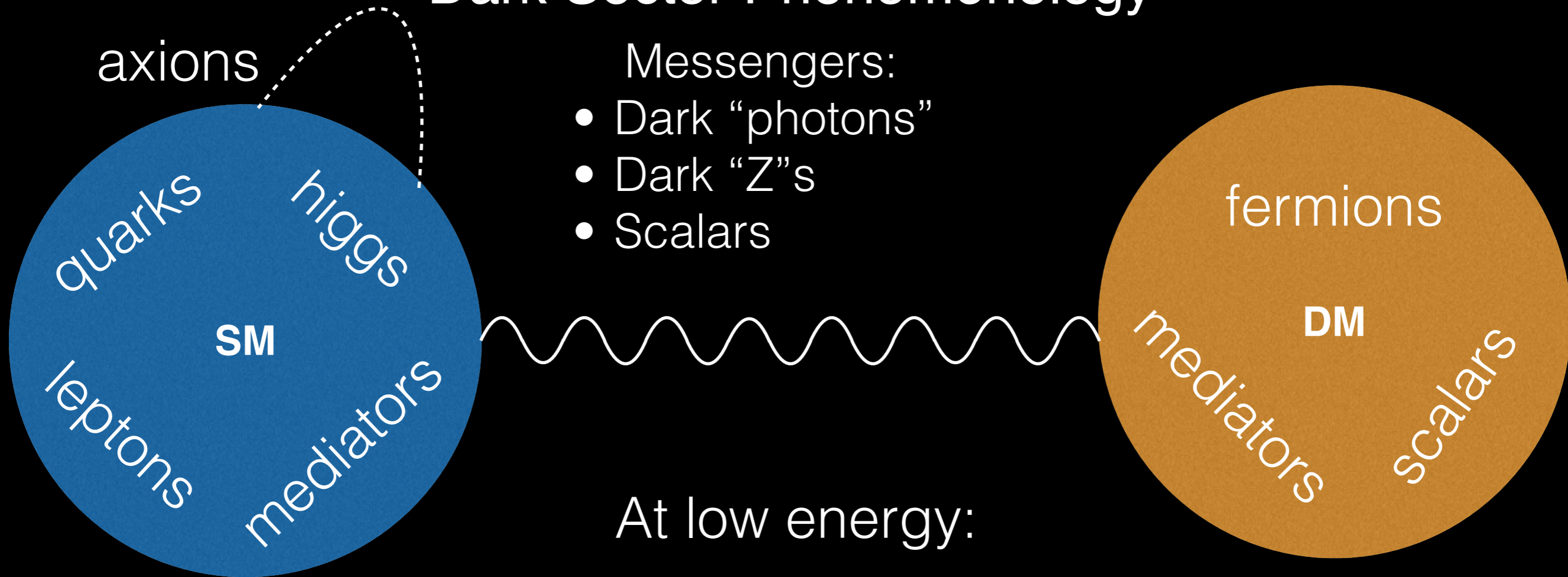
Dark Matter in Neutron Star Mergers

Sanjay Reddy
University of Washington

- Sub GeV dark matter
- Some existing constraints
- How to use gravitational waves from BNS mergers to constrain or discover dark sectors.

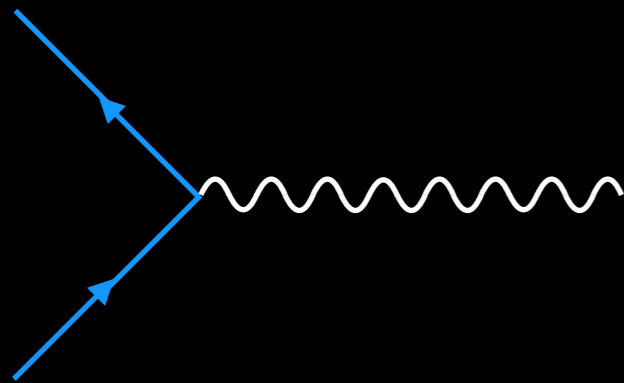


Dark Sector Phenomenology

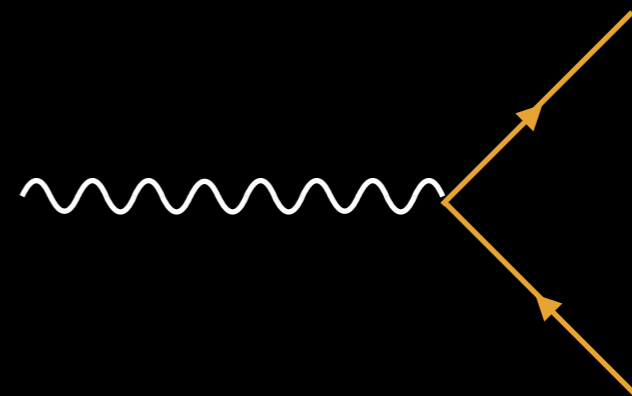


At low energy:

$$\mathcal{L}_{A'f} = g_f A'_\mu \bar{\psi}_f \gamma^\mu \psi_f$$



$$V_{\chi B} = \frac{g_\chi g_B}{q^2 + m_\phi^2} = g_\chi g_B \frac{e^{-m_\phi r}}{r}$$



$$V_{\chi\chi} = \frac{g_\chi^2}{q^2 + m_\phi^2} = g_\chi^2 \frac{e^{-m_\phi r}}{r}$$

DM Accretion onto Neutron Stars

For a concise recent review see Kouvaris (2013)

Mass accretion rate:

$$M_{\text{acc}} = 1.3 \times 10^{43} \left(\frac{\rho_{\text{dm}}}{0.3 \text{ GeV/cm}^3} \right) \left(\frac{t}{\text{Gyr}} \right) f \text{ GeV}$$

where $f = \text{Min} \left[1, \frac{\sigma}{10^{-45} \text{ cm}^2} \right]$

Thermalization:

$$\frac{GM(r_{\text{th}})m_{\chi}}{r_{\text{th}}} \approx \frac{3}{2}T \Rightarrow r_{\text{th}} \approx 2.2 \text{ m} \left(\frac{T}{10^5 \text{ K}} \right)^{1/2} \left(\frac{\text{GeV}}{m_{\chi}} \right)^{1/2}$$

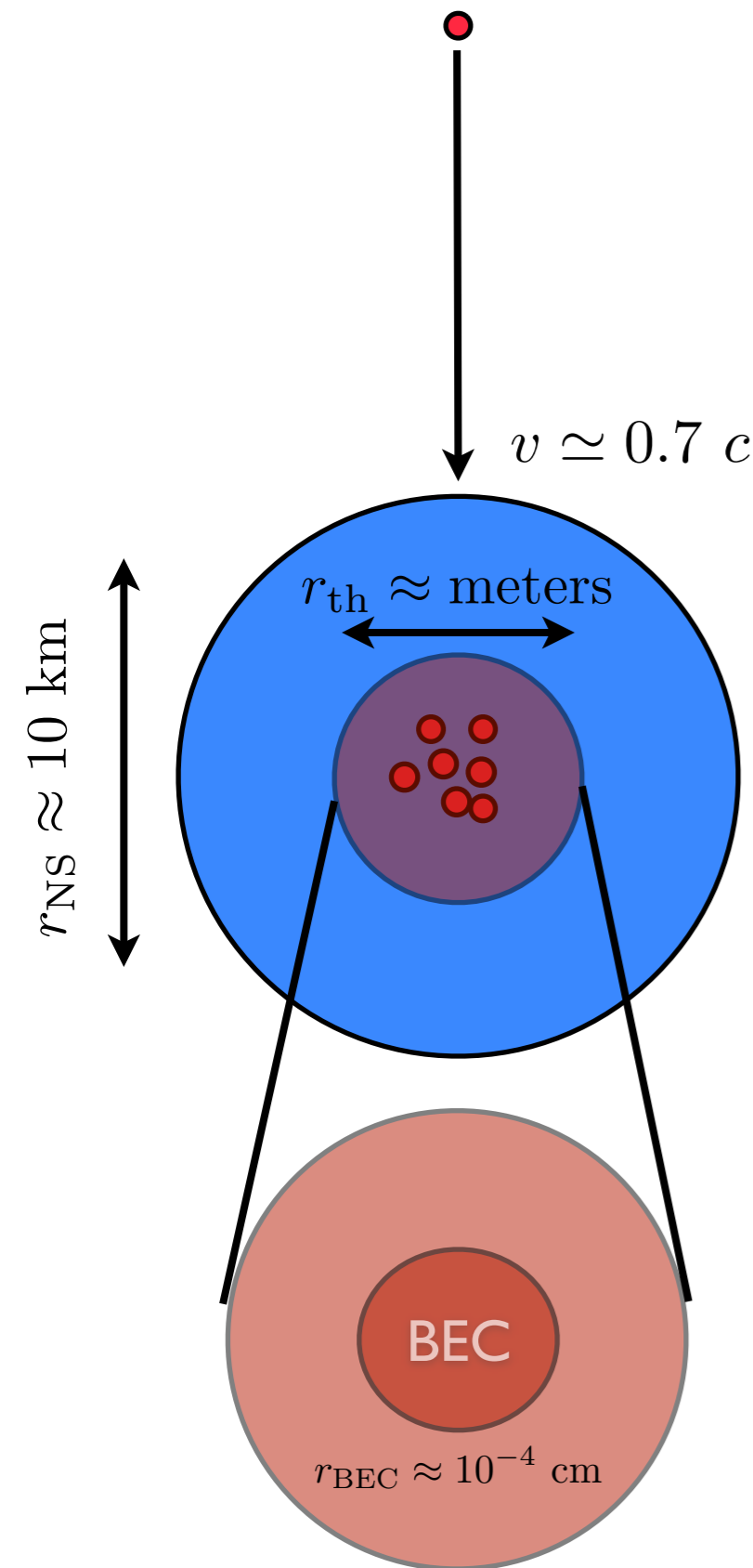
Self-Gravitation:

$$M_{\text{sg}} > \frac{4}{3} \pi \rho_c r_{\text{th}}^3 = 2.2 \times 10^{46} \text{ GeV} \left(\frac{m}{\text{GeV}} \right)^{-3/2}$$

Bose Einstein Condensation:

$$M_{\text{BEC}} > 8 \times 10^{27} \left(\frac{\text{GeV}}{m} \right)^{1.5} \text{ GeV}$$

Formation of BEC triggers collapse.



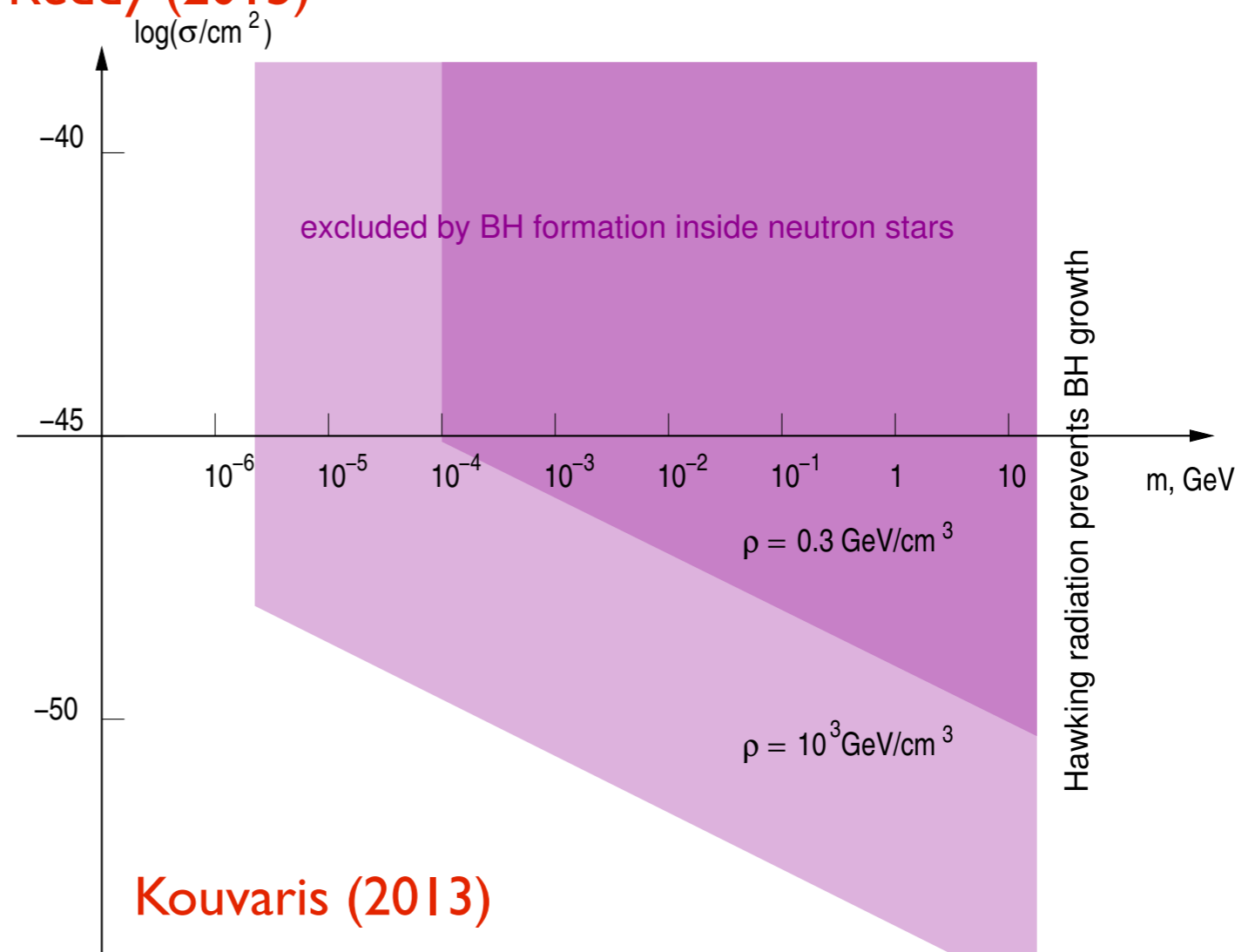
Black-hole Formation

Idea: Asymmetric bosonic dark matter can induce the collapse of the NS to a black hole. Goldman & Nussinov (1989)

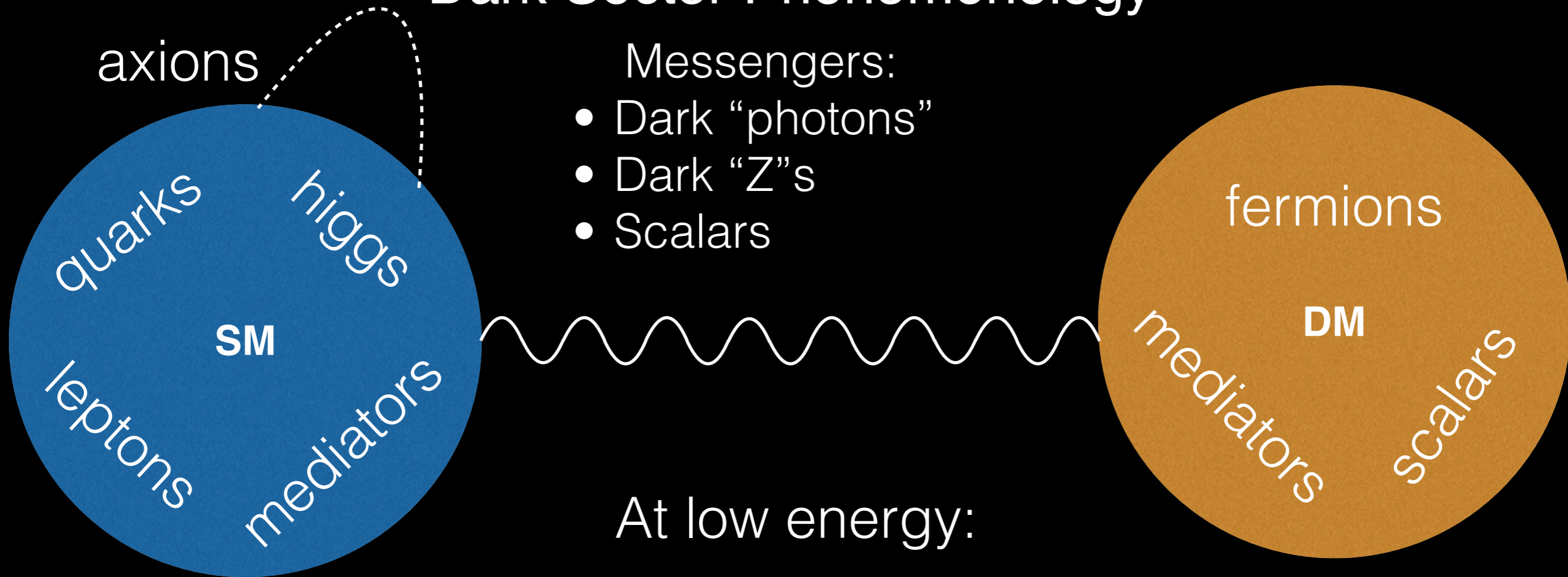
- Kouvaris and Tinyakov (2011)
- McDermott, Yu and Zurek (2012)
- Kouvaris (2012) & (2013)
- Guver, Erkoca, Reno, Sarcevic (2012)
- Fan, Yang, Chang (2012)
- Bell, Melatos and Petraki (2013)
- Jamison (2013)
- Bertoni, Nelson, Reddy (2015)

This idea has been explored in more detail by:

Existence of old neutron stars with estimated ages $\sim 10^{10}$ years provide strong constraints on asymmetric DM.

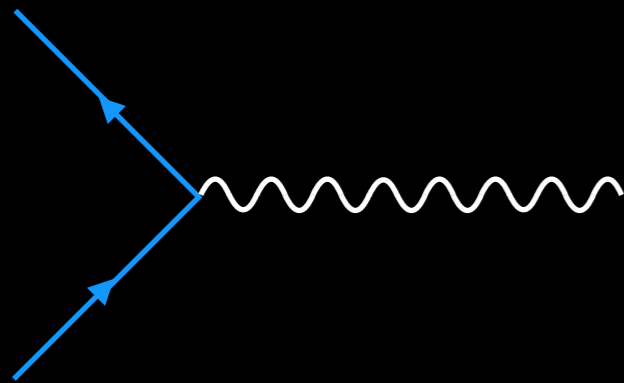


Dark Sector Phenomenology

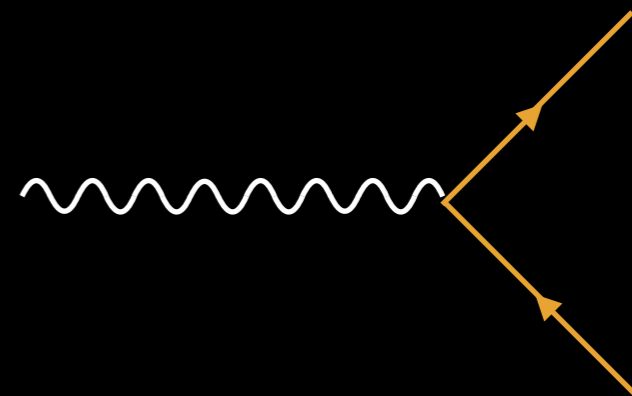


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Dark Photons

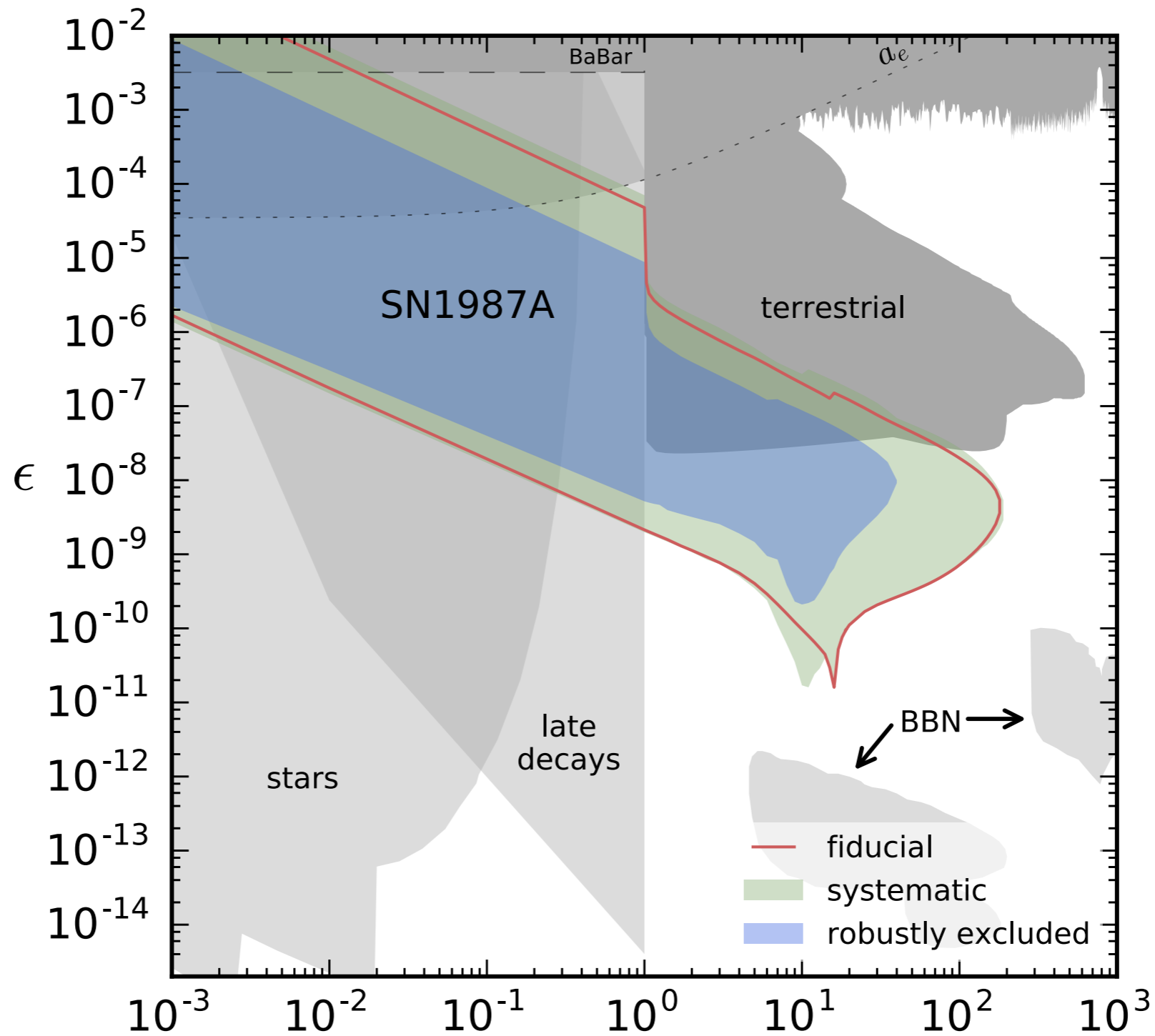
$$\mathcal{L} \supset g_Q A'_\mu J_\mu^{\text{EM}} + g_B B_\mu J_\mu^{\text{B}} - \frac{1}{2} m_{\gamma_Q}^2 A'_\mu A'^\mu - \frac{1}{2} m_{\gamma_B}^2 B_\mu B^\mu$$

$$g_Q = \sqrt{4\pi\alpha} \epsilon$$

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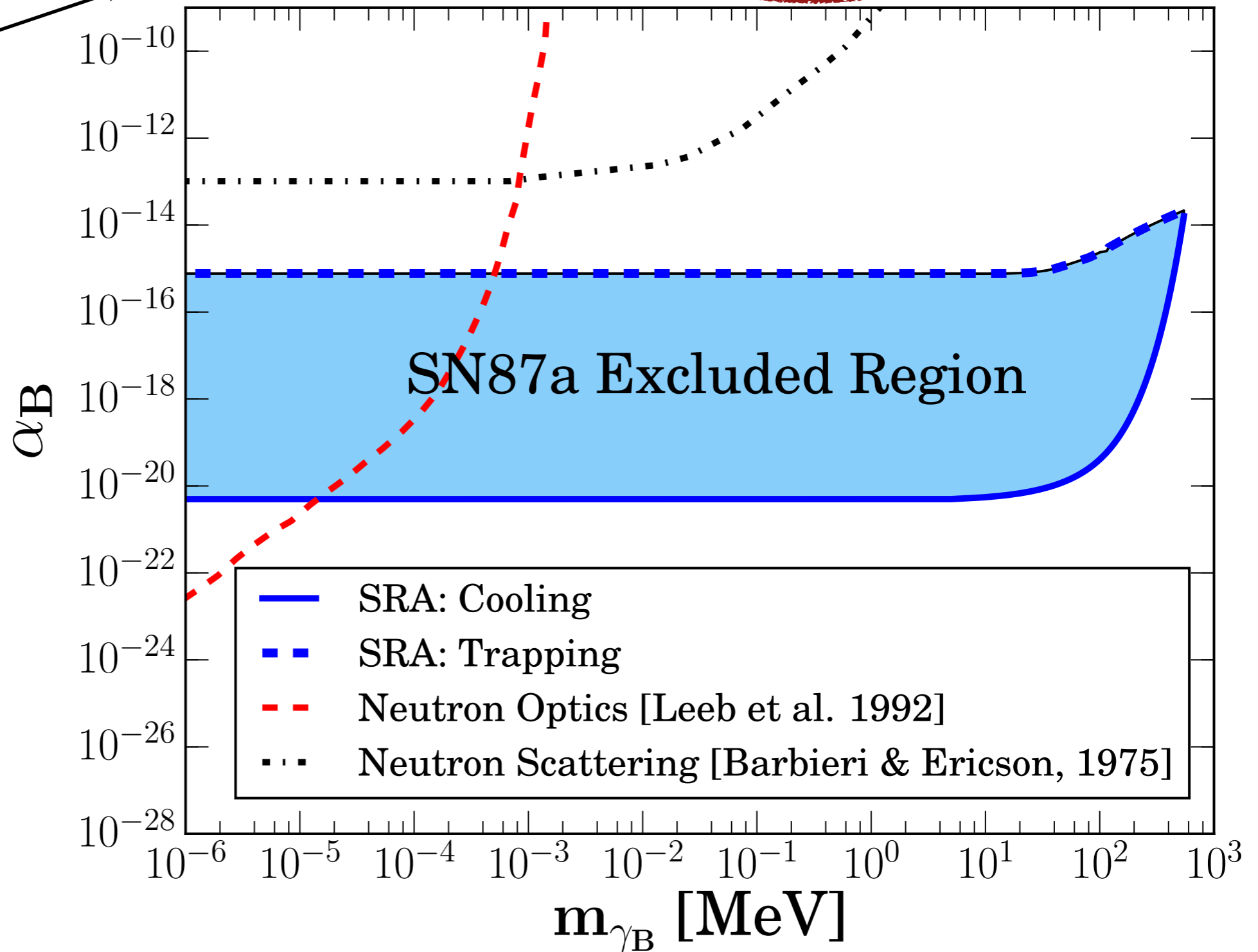
m' [MeV] Chang, Essig, McDermott (2017)

Dark Baryon Number Gauge Boson

$$\mathcal{L} \supset g_Q A'_\mu J_\mu^{\text{EM}} + g_B B_\mu J_\mu^{\text{B}} - \frac{1}{2} m_{\gamma_Q}^2 A'_\mu A'^\mu - \frac{1}{2} m_{\gamma_B}^2 B_\mu B^\mu$$

$$\alpha_B = \frac{g_B^2}{4\pi}$$

- Neutrino burst from SN87a provides stringent constraints.
- Constraints from RG and HB stars are relatively weaker.



SN87a bound on energy loss to exotic particles

Raffelt's "local" bound:

$$\mathcal{E}(\rho = 3 \times 10^{14} \text{ g/cm}^3, T = 30 \text{ MeV}) < \mathcal{E}_{\text{Raffelt}} = 10^{19} \frac{\text{ergs}}{\text{g s}}$$

This bound was found empirically by comparing to a suite of proto-neutron star simulations.

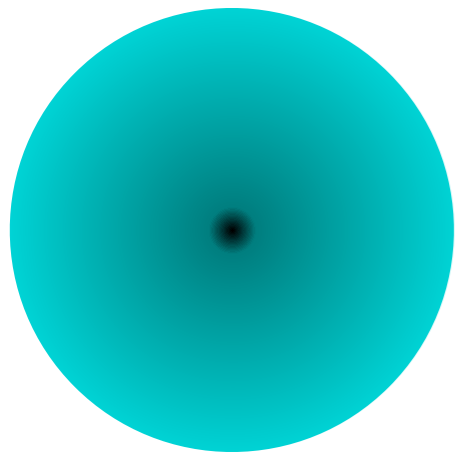
The corresponding bound on the luminosity is

$$L_{\text{exotic}} < \mathcal{E}_{\text{Raffelt}} \times M_{NS} \simeq 2 \times 10^{52} \frac{M}{M_{\odot}} \frac{\text{ergs}}{\text{s}}$$

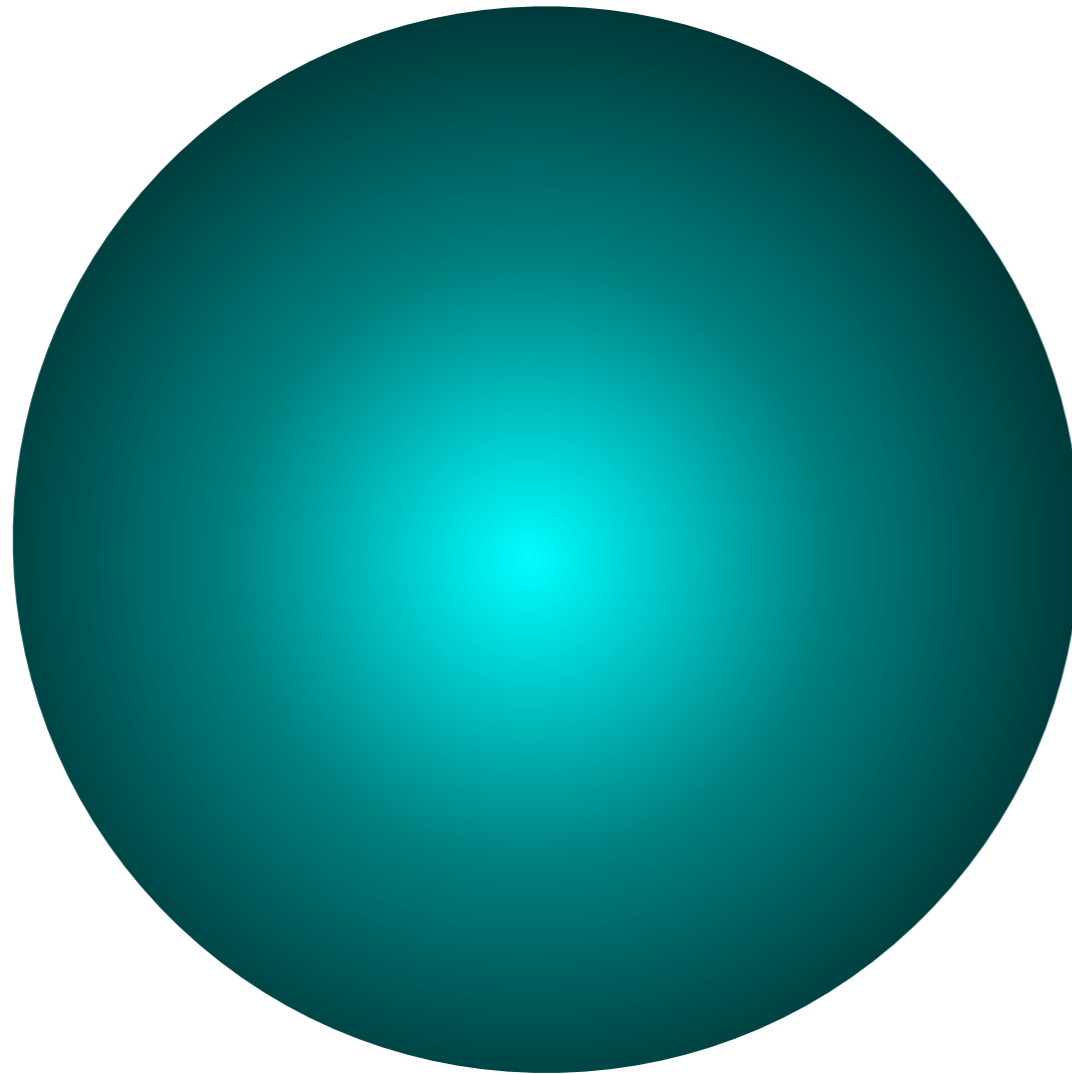
Requires $g_B < 10^{-10}$ for $m_{\phi} < 100 \text{ MeV}$

Dark Matter and Neutron Star Structure

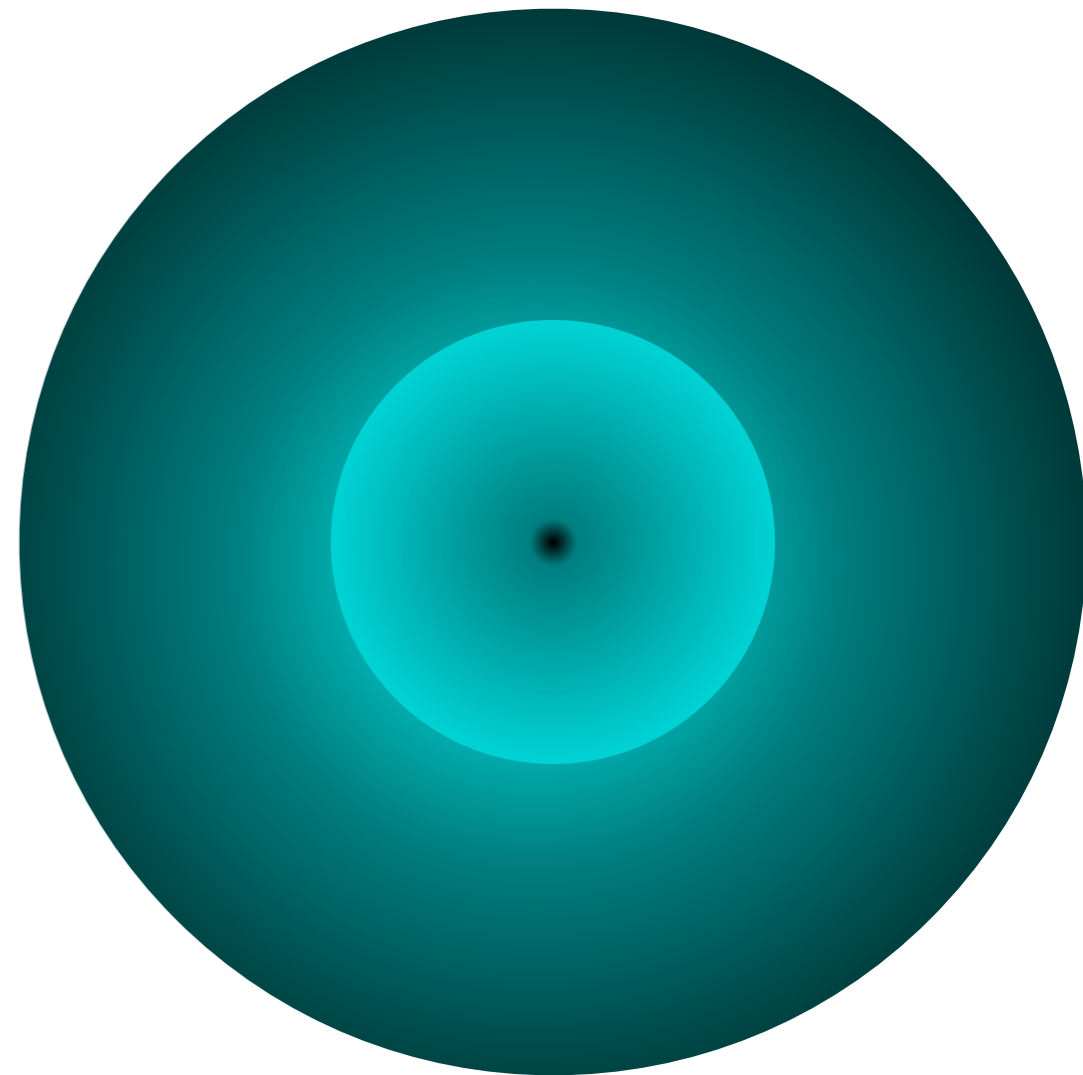
Trace amounts of dark matter can influence the structure of neutron stars.



dark-core



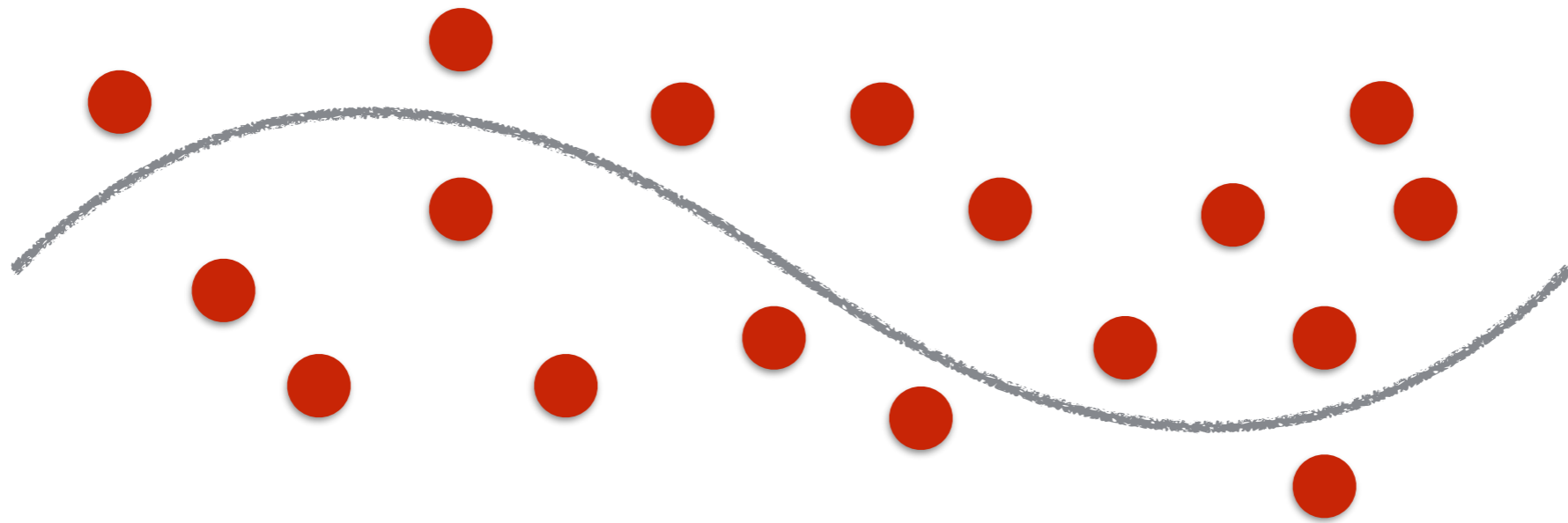
dark-halo



dark-halo +
anti-dark core

Equation of State of Dark Matter

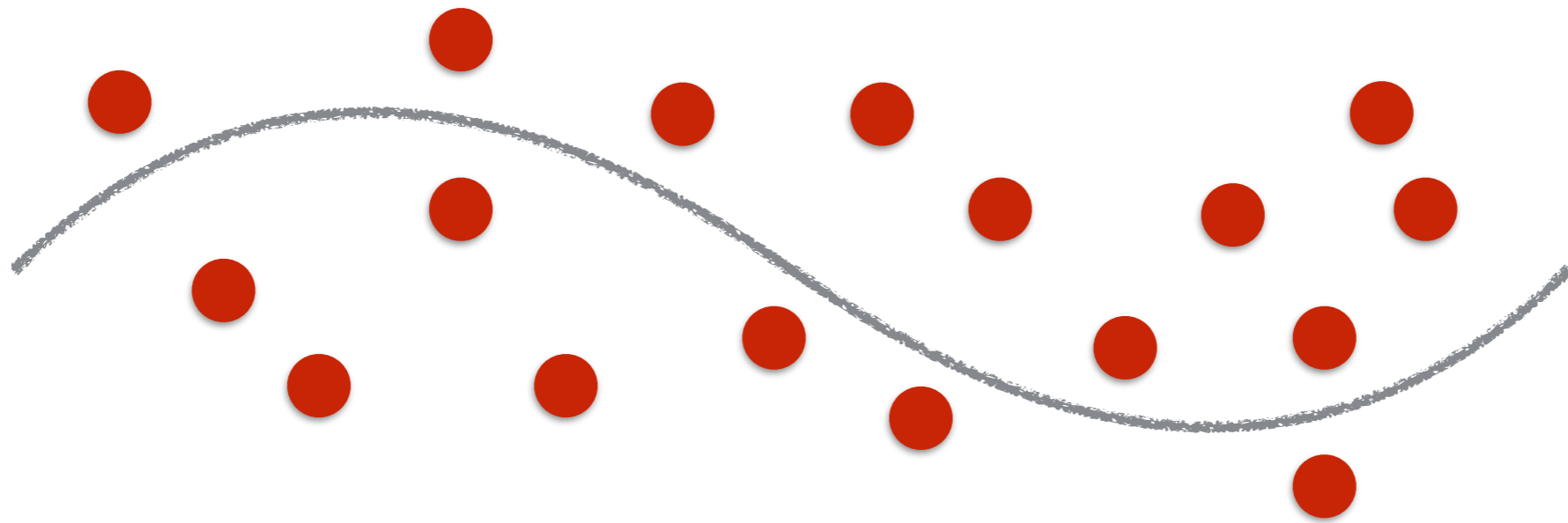
Energy density:
$$\epsilon_\chi = \epsilon_{\text{kin}} + m_\chi n_\chi + \frac{g_\chi^2}{2m_\phi^2} n_\chi^2$$



Large coherent enhancement of interactions when Compton wavelength of mediator is larger than the inter-particle separation.

Equation of State of Dark Matter

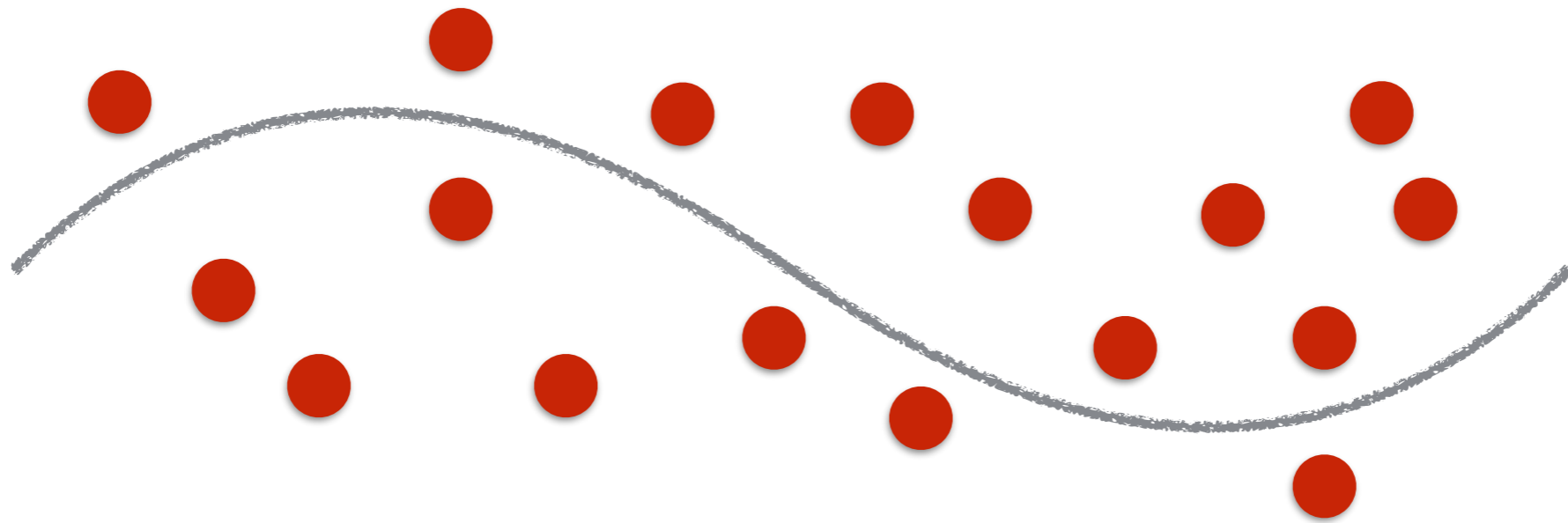
Energy density:
$$\epsilon_\chi = \epsilon_{\text{kin}} + m_\chi n_\chi + \frac{g_\chi^2}{2m_\phi^2} n_\chi^2 + \frac{g_\chi g_B}{m_\phi^2} n_B$$



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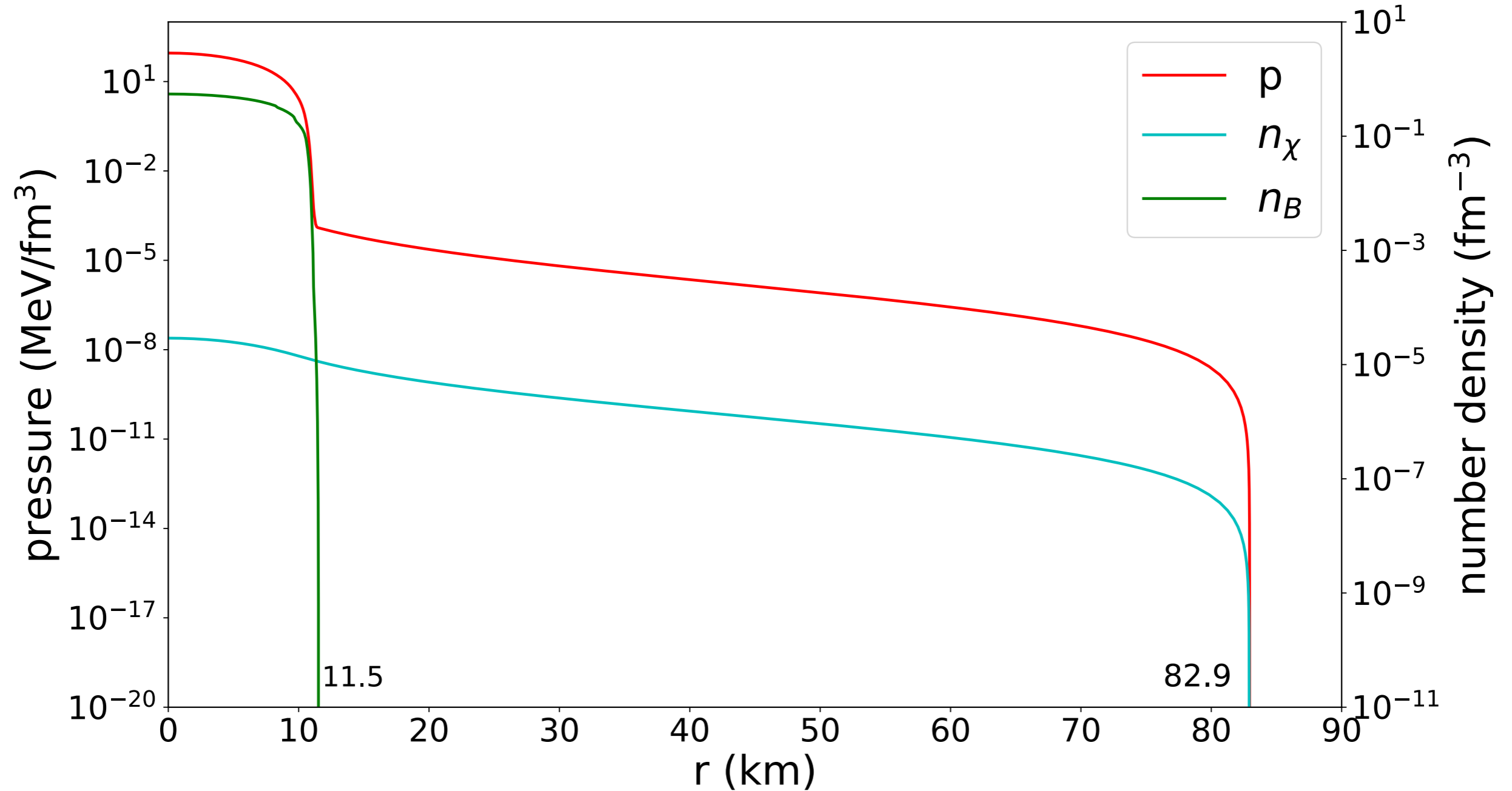
Equation of State of Dark Matter

Energy density:
$$\epsilon_\chi = \epsilon_{\text{kin}} + m_\chi n_\chi + \frac{g_\chi^2}{2m_\phi^2} n_\chi^2 + \frac{g_\chi g_B}{m_\phi^2} n_B - \frac{g_{\bar{\chi}} g_B}{m_\phi^2} n_B$$



Large coherent enhancement of interactions when Compton wavelength of mediator is larger than the inter-particle separation.

Profile of a Dark Neutron Star



1.4 M_{solar} Neutron star with $10^{-4} M_{\text{solar}}$ of dark matter.

Dark matter: $m_\chi = 100 \text{ MeV}$

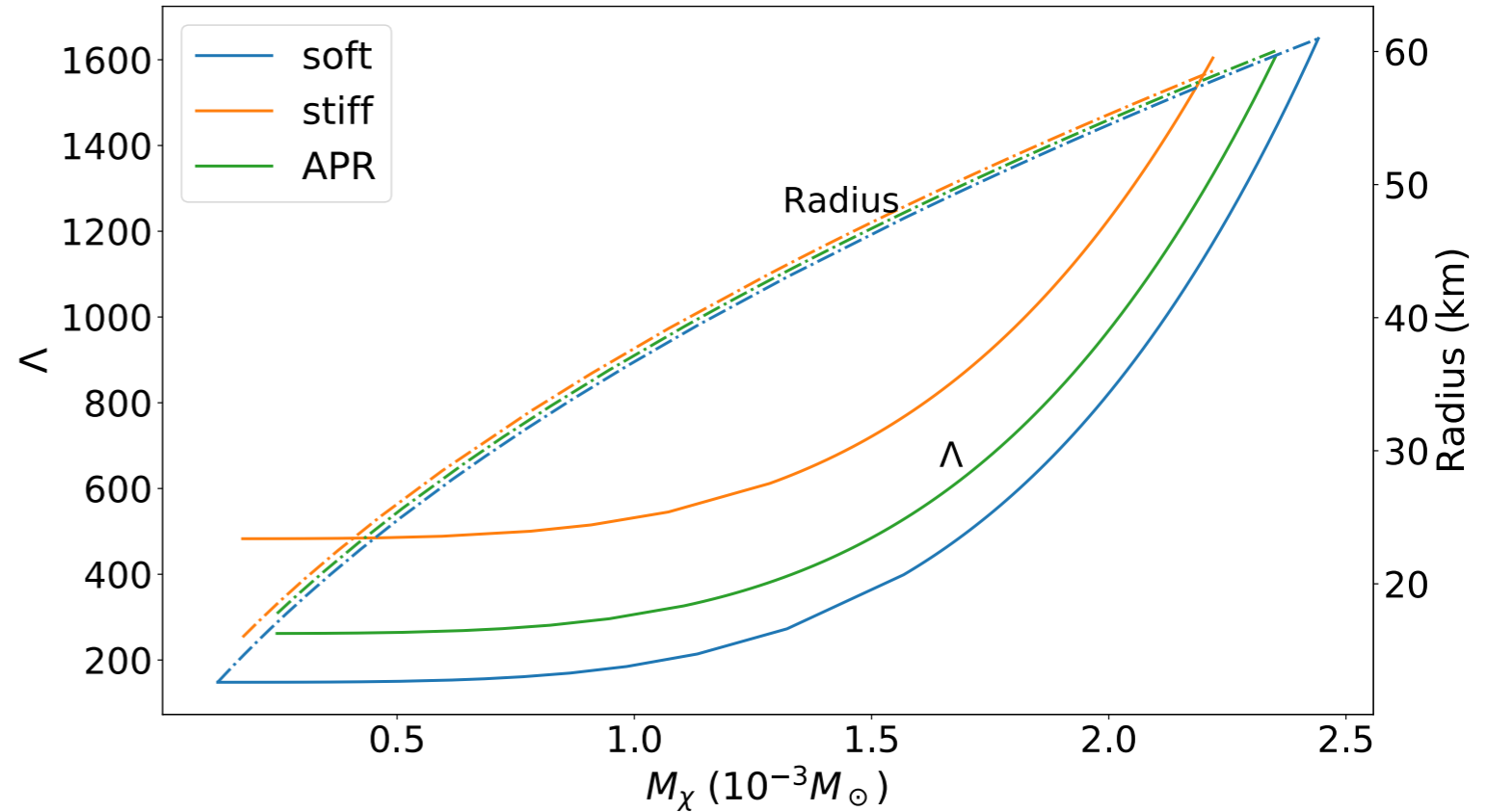
Interactions: $g_\chi/m_\phi = (0.5/\text{MeV})$ or $(0.5 \times 10^{-6}/\text{eV})$

For light mediators, only trace amounts are needed

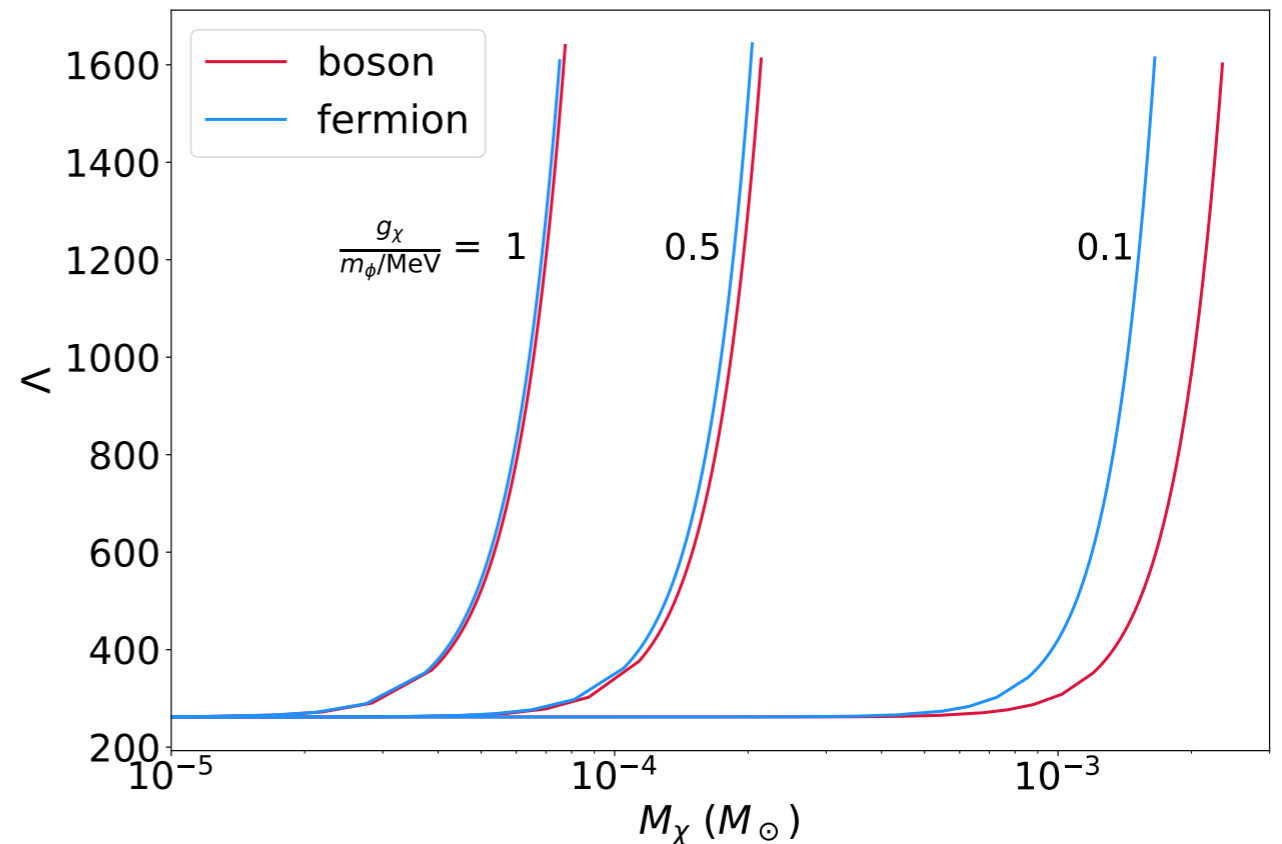
10^{-4} - $10^{-2} M_{\text{solar}}$ is adequate to enhance $\Lambda > 800$!

$$m_{\chi} = 100 \text{ MeV}$$

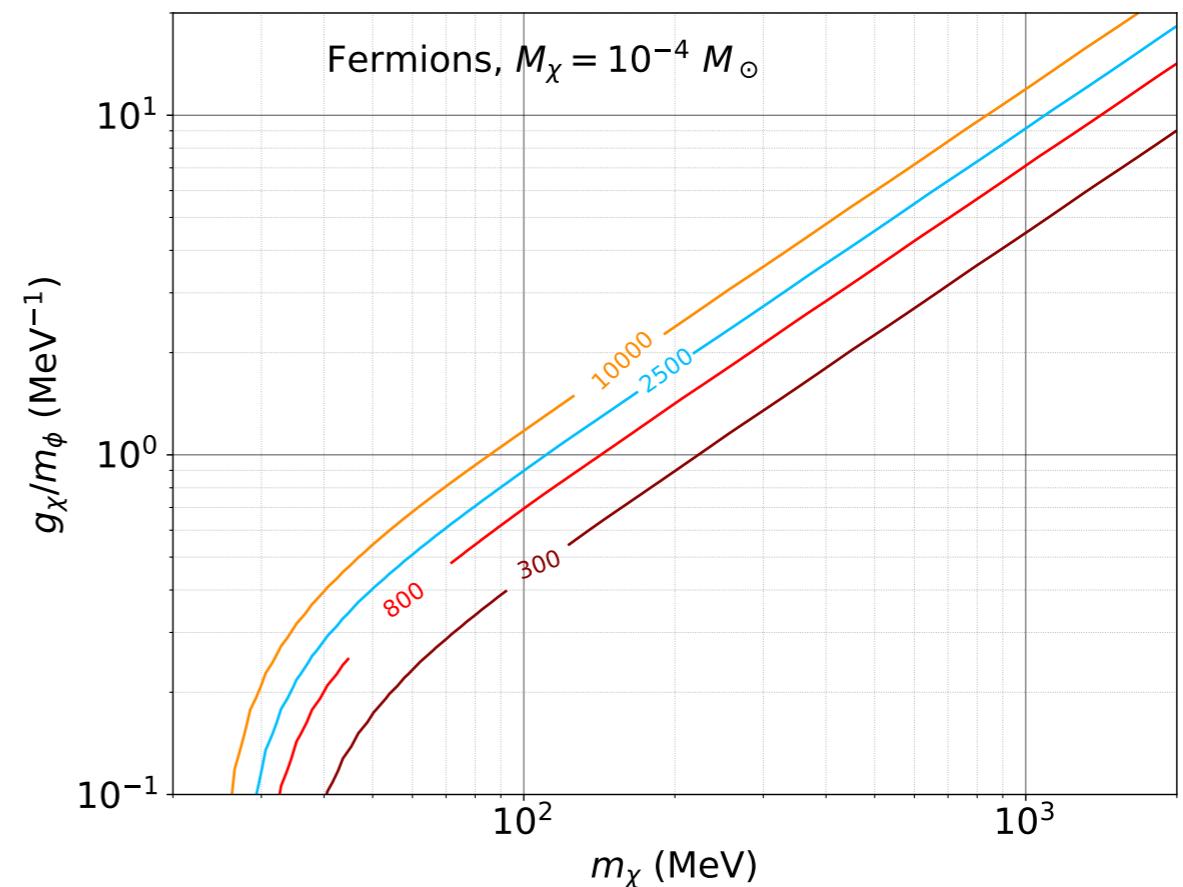
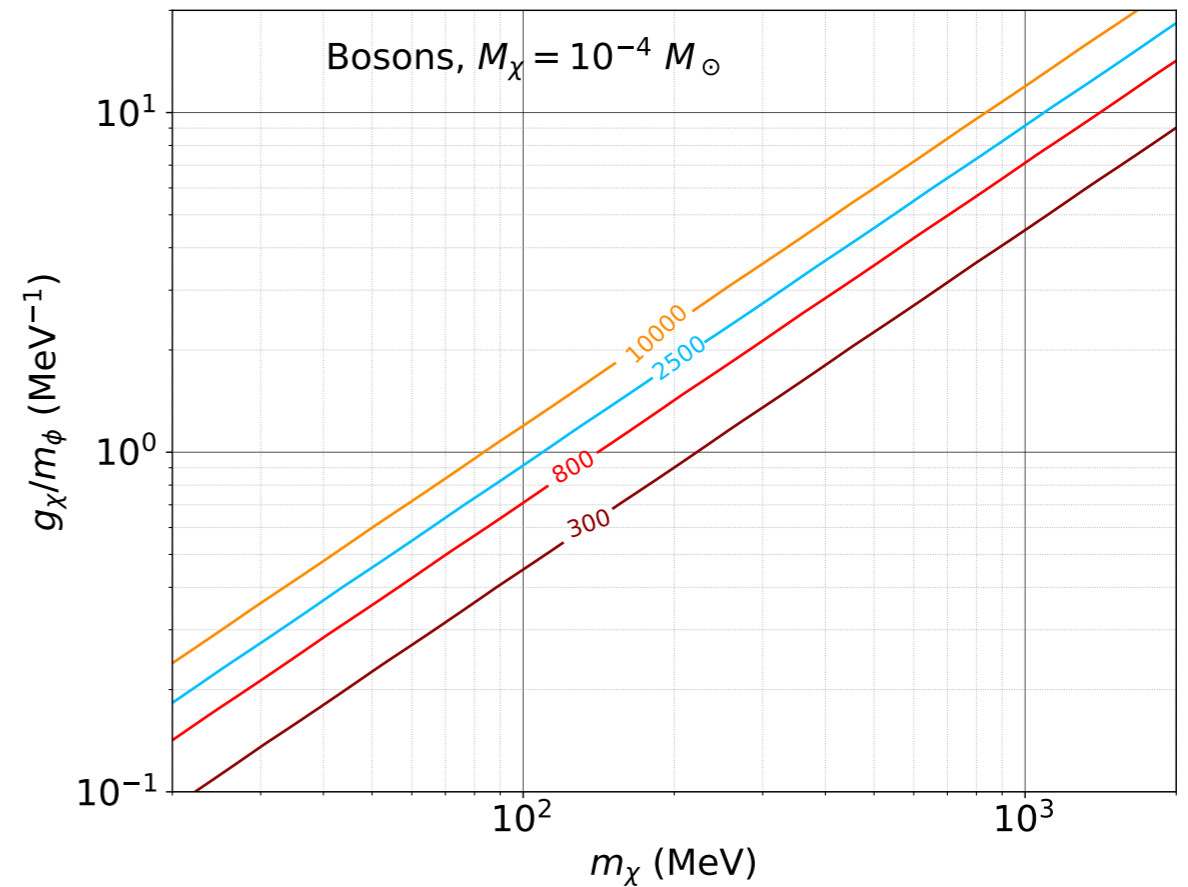
$$g_{\chi}/m_{\phi} = (0.1/\text{MeV}) \text{ or } (10^{-6}/\text{eV})$$



Interactions of “natural” size produce large Λ



GW170817 rules out extensive regions of the dark matter model space if NS contain dark matter.



Could/should NS contain dark matter ?

- Supernova can easily produce $10^{-4} M_{\text{solar}}$ of < 100 MeV dark matter.
- Coupling to baryons allows for dark charge separation.
- Dark matter might be clumpy.
- Dark clumps might seed star formation.
- Might be the best place to find them ?

Conclusions



MeV-GeV dark matter can play a role in mergers. Trace amounts of interacting dark matter in the neutron star can enhance their tidal polarizability (Λ) to discernible values.



Neutron stars can accrete, inherit, or create their own dark matter. Dark matter production during supernova and mergers can be significant even for very weak coupling.



If Adv. LIGO suggests either large Λ or a large variability in Λ , it may reveal the particle nature of dark matter - gravitationally !



If not, it provides useful constraints on generic dark matter models with light mediators.