# Searches for gravitational waves from known pulsars: constraining the pulsar ellipticity distribution

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# Image Image produced with produced  $\frac{1}{2}$ [psrqpy](http://psrqpy.readthedocs.io), Pitkin, JOSS *JOSS* [\(2018\)](http://joss.theoj.org/papers/10.21105/joss.00538)  $(2018)$

# **Background**

Known pulsars are a great target in searches for continuous gravitational wave signals.

Known frequency evolution and sky location allow coherent searches over long data stretches (~years) using (*almost*) a single phase template

> Pulsars with periods accessible to the LIGO/Virgo gravitational wave detectors



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#### Searches

Previous known pulsar searches carried out by the LIGO Scientific Collaboration & Virgo Collaboration (LVC) have made the following assumption:

- signals are emitted from a triaxial star ( $l=m=2$  mass quadrupole mode) rotating about its principal moment of inertia  $I_{zz}$  (no precession)
- GW signals are phase locked with the electromagnetic emission (which is itself locked to the star's rotation)<sup>†</sup> giving emission at twice the rotation rate  $f_{\rm rot}$
- each star's amplitude is independent (i.e. each target is independent of all others)

<sup>†</sup>Some targeted searches have been performed relaxing the very strong assumption about GW emission being phase locked to the rotation, e.g., Abbott *et al*, *ApJL*[, 683 \(2008\)](https://arxiv.org/abs/0805.4758) & [Abbott](https://arxiv.org/abs/1710.02327) *et al*, *PRD* [96, 122006 \(2017\)](https://arxiv.org/abs/1710.02327)

#### Searches

For each pulsar, searches attempt to evaluate the probability distribution of:

- *h*<sup>0</sup> : the gravitational wave strain detected at the Earth
- $\bullet$  cos  $\iota$ : the cosine of the inclination of the rotation axis to the line-of-site
- $\bullet$   $\phi_0$ : the phase of the signal at some epoch
- $\bullet$   $\psi$ : the polarisation angle

When no signal is found an upper limit on  $h_0$  can be set (often at 95% credible level). This can be compared to the *spin-down limit* set by assuming all rotational kinetic energy is dissipated through *l*=*m*=2 mass quadrupole GW emission:

$$
h_0^{\rm sd} = 2.55 \times 10^{-25} \left( \frac{I_{zz}}{10^{38} \,\text{kg} \,\text{m}^2} \right)^{1/2} \left( \frac{1 \,\text{kpc}}{d} \right) \left( \frac{100 \,\text{Hz}}{f_{\rm rot}} \right)^{1/2} \left( \frac{| \dot{f}_{\rm rot} |}{10^{-11} \,\text{Hz} \, s^{-1}} \right)^{1/2}
$$

## LVC Searches



LIGO S2: 28 pulsars. [Abbott](https://arxiv.org/abs/gr-qc/0410007) *et al*, *PRL* [94, 181103 \(2005\)](https://arxiv.org/abs/gr-qc/0410007)

LIGO S3+S4: 78 pulsars. [Abbott](https://arxiv.org/abs/gr-qc/0702039)  *et al*, *PRD* [76, 042001 \(2007\)](https://arxiv.org/abs/gr-qc/0702039)

LIGO S5: 116 pulsars. [Abbott](https://arxiv.org/abs/0909.3583) *et al*, *ApJ*[, 713 \(2010\)](https://arxiv.org/abs/0909.3583)

Rely on up-to-date ephemerides from EM pulsar observations (radio, X-ray,  $\gamma$ -ray) preferably overlapping GW observing runs.



*et al*, *ApJ*[, 839 \(2017\)](https://arxiv.org/abs/1701.07709)



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#### Searches

The probability distribution of  $h_0^{\parallel}$  can be converted into a distribution on the mass quadrupole moment  $\mathcal{Q}_{22}$ , or fiducial ellipticity<sup>†</sup>  $\varepsilon$  assuming a known distance (often known to  $\sim$  20%):

$$
\varepsilon = 10^{-6} \left( \frac{h_0}{4.2 \times 10^{-26}} \right) \left( \frac{100 \,\text{Hz}}{f_{\text{rot}}} \right)^2 \left( \frac{10^{38} \,\text{kg m}^2}{I_{zz}} \right) \left( \frac{d}{1 \,\text{kpc}} \right)
$$

and (following [Ushomirsky, Cutler & Bildsten,](https://arxiv.org/abs/astro-ph/0001136) *MNRAS* 319 ( 2000))

$$
Q_{22} = \sqrt{\frac{15}{8\pi}} I_{zz} \varepsilon
$$

This can in-turn be converted to a limit on the model-dependent internal *B*-field strength (e.g. Cutler, *PRD* [66, 084205 \(2002\)](https://arxiv.org/abs/gr-qc/0206051) for toroidal field with  $B < 10^{15}$  G):

$$
\varepsilon_B \approx -1.6\!\times\!10^{-6}\left(\frac{\langle B_t\rangle}{10^{15}\,\mathrm{G}}\right)
$$



We can convert to surface deformation, maximised over EoS, using $^\dagger$ :

$$
R\varepsilon_{\text{surf},22} \approx 25 \left(\frac{\varepsilon}{10^{-4}}\right) \text{ cm}
$$

8 † [Johnson-McDaniel,](https://arxiv.org/abs/1303.3259) *PRD* 88, [044016 \(2013\)](https://arxiv.org/abs/1303.3259)



Smallest spin-down ratio: Crab pulsar,  $\varepsilon$  <  $3.6 \times 10^{-5}$ , which is 20 times below the spin-down limit

MSP closest to spin-down limit: J0437-4715 (GW frequency 347 Hz, at 0.16 kpc)  $\varepsilon$  < 2.8  $\times$  10<sup>-8</sup>, which is only 1.4 times spin-down limit. Limits toroidal *B* field to  $\leq 10^{13}$  G

# Narrow-band search

LIGO O1 data: 11 pulsars. [Abbott](https://arxiv.org/abs/1710.02327) *et al*, *PRD*[, 96 122006 \(2017\)](https://arxiv.org/abs/1710.02327)



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#### Ellipticity distribution (with C. Messenger & X. Fan)

Known pulsar searches assume each pulsar's amplitude is independent. However, we could instead assume that there is an underlying distribution from which the ellipticities are drawn, e.g., an **exponential** distribution

$$
p(\varepsilon | \mu_{\varepsilon}, I) = \frac{1}{\mu_{\varepsilon}} e^{\varepsilon / \mu_{\varepsilon}}
$$

defined by the *hyperparameter*  $\mu_{\varepsilon}$ , the mean of the distribution.

We can combined data from all pulsars to estimate the probability distribution of  $\mu_{\varepsilon}^{}$ : *hierarchical Bayesian inference* (already used in GW field for black holes mass and spin distributions, e.g., Abbott *et al.*, *PRX* [6, 041015 \(2016\)](https://arxiv.org/abs/1606.04856) & [Stevenson,](https://arxiv.org/abs/1703.06873)  [Berry & Mandel,](https://arxiv.org/abs/1703.06873) *MNRAS*, 471 (2017)).

# Ellipticity distribution

We form a **joint likelihood** of the data from all pulsars marginalised over independent  $\cos\iota$ ,  $\psi$ , and  $\phi_{_0}$  values for each pulsar, and also marginalising over the uncertainty on distance (assuming a 20% Gaussian error):

The same kind of analysis could be used on results from a blind all-sky search by changing the prior on the distance

 $p(\mu_{\varepsilon}|\{\mathbf{D}\},I) \propto \int_{\varepsilon_{\min}}^{\varepsilon_{\max}} \left( \prod_{i=1}^{N_{\mathrm{pulsars}}} \int^{\theta_{i}} p(\mathbf{D}_{i}|\varepsilon,\theta_{i},d_{i},I) p(\theta_{i}|I) p(d_{i}|I) d\theta_{i} \mathrm{d}d_{i} \right) p(\varepsilon|\mu_{\varepsilon},I) p(\mu_{\varepsilon}|I) d\varepsilon$ 

where  $\mathbf{D}_i$  is the data for each pulsar,  $\theta_i^2 = {\phi_{0i^*}} \psi_i^2 \cos \iota_i^3$ , and *I* assumes knowledge such as pulsar frequencies.

Calculating this for all pulsars at once would be a difficult problem (100s of parameters), but instead we can make use of posteriors distributions on  $\varepsilon$ (or equivalently  $Q_{22}$ ) individually marginalised over other parameters.

The probability distribution for the ellipticity distribution (and prior on hyperparameters) can be changed to your favourite function

# **Simulations**

We generated independent realisations of signals from a population of 200 known pulsars (those searched for in O1 data) with  $\varepsilon$  drawn from exponential distributions with a range of  $\mu_{_\mathcal{E}}$  values, injected into Gaussian noise based on the aLIGO design sensitivity.

We also generated a set of "background" realisations, in which all pulsars had zero ellipticity (i.e. detector noise only).

For each population realisation we calculate the "evidence" that the data contains pulsars with  $\varepsilon$  drawn from an exponential distribution, and also that the population is consistent with noise.

Exponential mean,  $\mu_{\varepsilon}$ , prior (Jeffreys)  $p(\mu_{\varepsilon}|I) = \frac{(\ln \mu_{\varepsilon \max} - \ln \mu_{\varepsilon \min})^{-1}}{\mu_{\varepsilon}}$  if  $\mu_{\varepsilon \min} < \mu_{\varepsilon} < \mu_{\varepsilon \max}$ 

> Bayesian evidence for population with ellipticities drawn from an exponential distribution

 $p(\{\mathbf{D}\}|H_{\text{exp}}, I) = \int^{\mu_{\varepsilon}} p(\{\mathbf{D}\}| \mu_{\varepsilon}, H_{\text{exp}}, I) p(\mu_{\varepsilon}|H_{\text{exp}}, I) d\mu_{\varepsilon}$ 

Bayesian odds comparing evidence for exponential distribution to data consistent with noise

$$
\mathcal{O} = \frac{p(\{\mathbf{D}\} | H_{\text{exp}}, I)}{p(\{\mathbf{D}\} | H_{\text{noise}}, I)}
$$

# **Results**

90% credible intervals from the posterior distributions on  $\mu_{\varepsilon}$  for the simulated population distributions (orange/green intervals are for the ensemble with the largest/smallest odds)



## Results

Odds comparing a hypothesis that the data is consistent with signals having ellipticities drawn from an exponential distribution against all data being consistent with noise

Odds comparing a hypothesis that the data is consistent with *any* combination of pulsars containing a signal (signal amplitudes are independently marginalised over) against all data being consistent with noise





#### Efficiency curves for detection of ensemble of pulsars.

#### **Results**

Left: using a "false alarm probability" set by the number of background realisations. *Right*: extrapolating the background to an equivalent " $5\sigma$ " level (using a KDE of the background).



#### **Results**

What about a different distribution? We can assume the distribution is a half-Gaussian (using the same prior on  $\sigma_{_{\!\!E}}$  as we had for  $\mu_{_{\!\!E}}$ )and compare models

$$
p(\varepsilon|\sigma_{\varepsilon}, I) = \frac{2}{\sqrt{2\pi}\sigma_{\varepsilon}} e^{-\frac{1}{2}(\varepsilon/\sigma_{\varepsilon})^2}
$$



#### S6 results

Using 92 pulsars from the S6 known pulsar search ([Aasi](https://arxiv.org/abs/1309.4027) *et al*, *ApJ*, [785 \(2014\)](https://arxiv.org/abs/1309.4027)) we applied the analysis assuming a 20% Gaussian error on pulsar distances (and upper bound on  $\varepsilon$  of  $\sim 10^{-4}$ ). **No detection** of an ensemble of sources**.**



# S6 results

Set 95% credible upper limits of:

- $\bullet$   $\mu_{\varepsilon}$  < 5.4 × 10<sup>-8</sup> for an exponential distribution
- $\bullet$  $\sigma_{\varepsilon}$  < 6.5 **×** 10<sup>-8</sup> for a half-Gaussian distribution



# **Conclusions**

Searches for known pulsars provide tight constraints on neutron star ellipticity that are starting to compete with *spin-down limits*. Further searches are underway we're currently looking at O2 data.

We can apply the method to constrain (a parameterised model-dependent) ellipticity distribution using results from real searches.

This could give us a new detection statistic for the ensemble of pulsars rather than relying on detecting individual sources (see also, e.g., [Cutler & Schutz,](https://arxiv.org/abs/gr-qc/0504011) *PRD* 72, [063006 \(2005\)](https://arxiv.org/abs/gr-qc/0504011), [Fan, Chen & Messenger, PRD 94, 084029 \(2016\)](https://arxiv.org/abs/1607.06735) & [Smith &](https://arxiv.org/abs/1712.00688)  [Thrane, arXiv:1712.00688](https://arxiv.org/abs/1712.00688)).

# Questions

What might be a good model for the underlying distribution?

- Exponential?
- Power law?
- Gaussian mixture?
- Non-parametric (i.e. histogram, [IGMM](https://www.seas.harvard.edu/courses/cs281/papers/rasmussen-1999a.pdf))?

Should we fit separate distributions for MSPs & young pulsars / globular cluster & field pulsars / binary pulsars and isolated pulsars?

Should/could we also constrain model-dependent internal *B*-field distributions?

What distance priors are reasonable if using sources from blind searches?