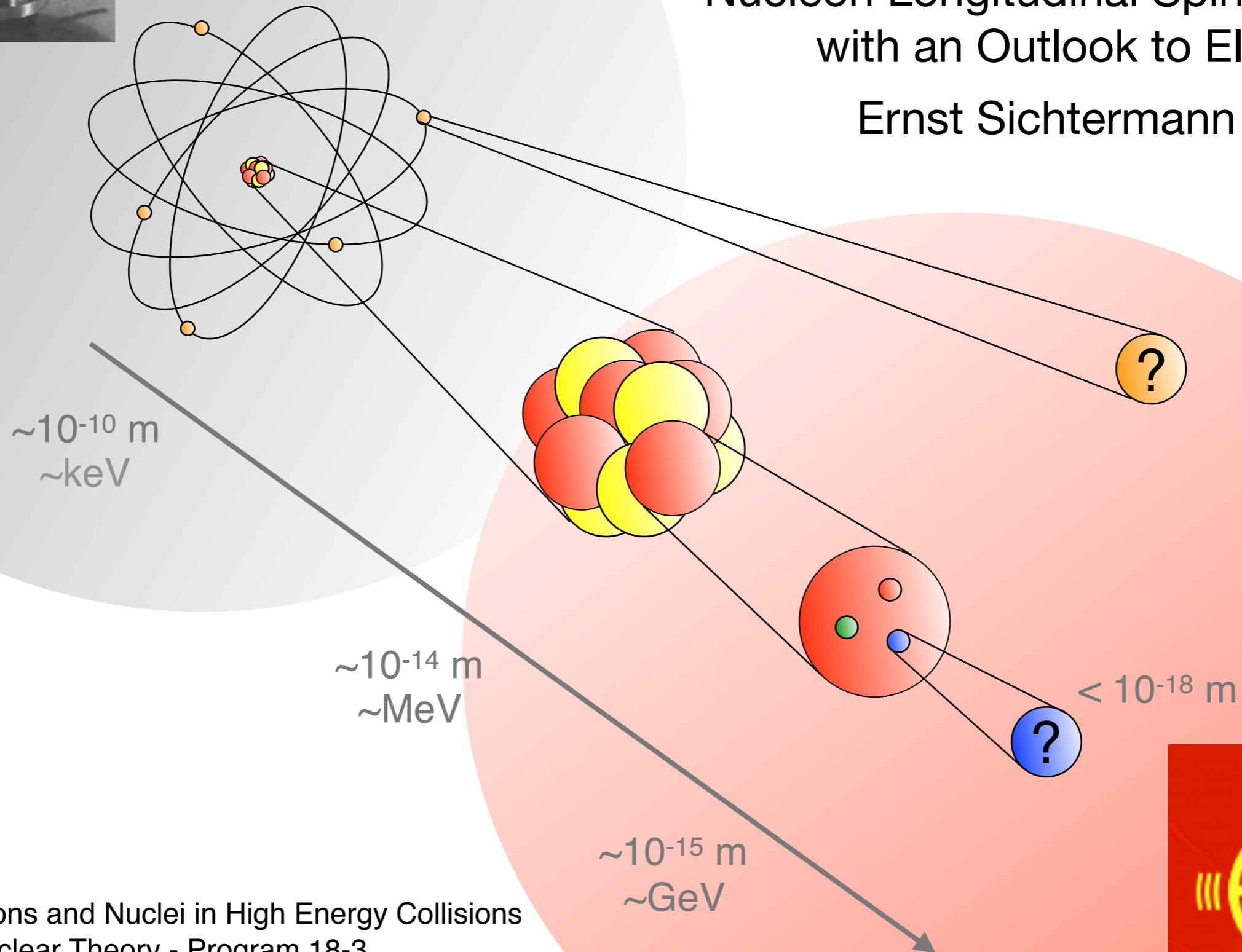


Nucleon Longitudinal Spin Status with an Outlook to EIC

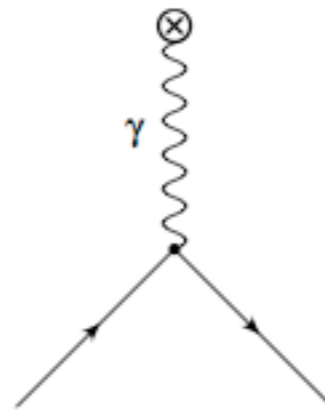
Ernst Sichteremann



Proton Magnetic Moment - circa 1935

- The magnetic moment $\vec{\mu}$ of a particle is related to its spin \vec{S} according to:

$$\vec{\mu} = g \frac{e}{2mc} \vec{S}$$



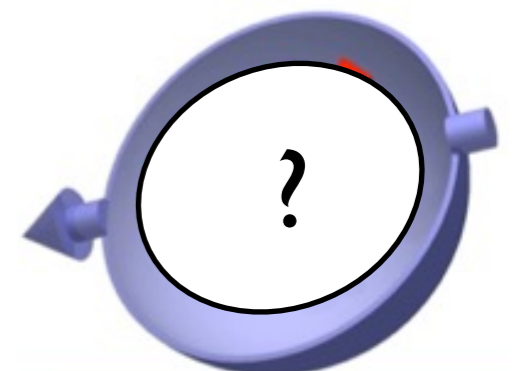
$$g = 2 + \mathcal{O}(\alpha_e)$$

for Dirac particles

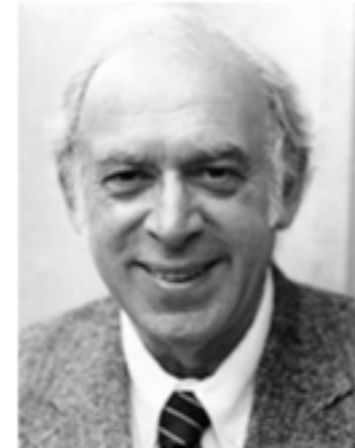
- 1933 - Frisch and Stern:

$$\vec{\mu}_p = 5.8 \frac{e}{2mc} \vec{S}_p - \textit{Proton has (spin-)substructure}$$

But, what is it?



Deep-Inelastic Electron Scattering



J.T. Friedman



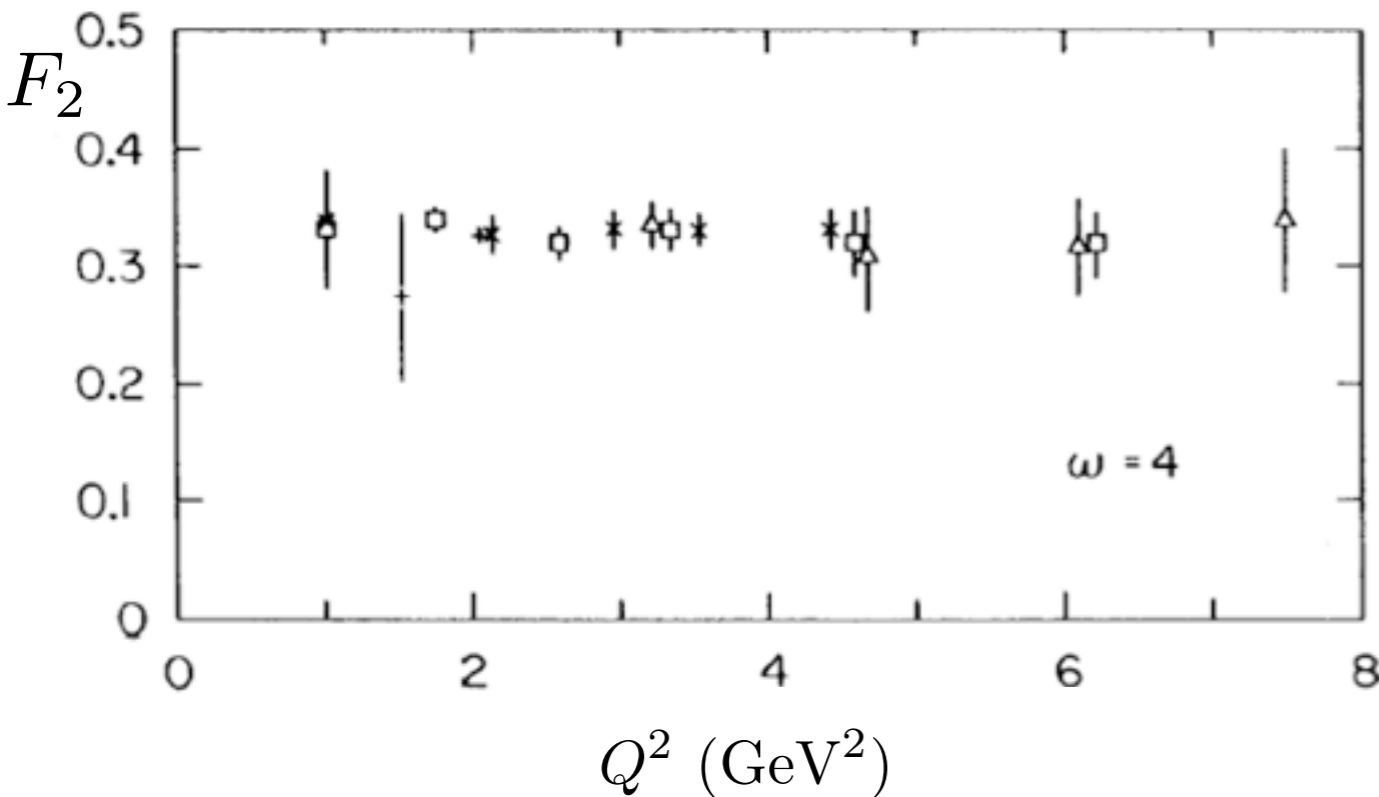
R. Taylor
Nobel Prize 1990



H.W. Kendall

Bjorken scaling:

+ 6° □ 18°
× 10° △ 26°

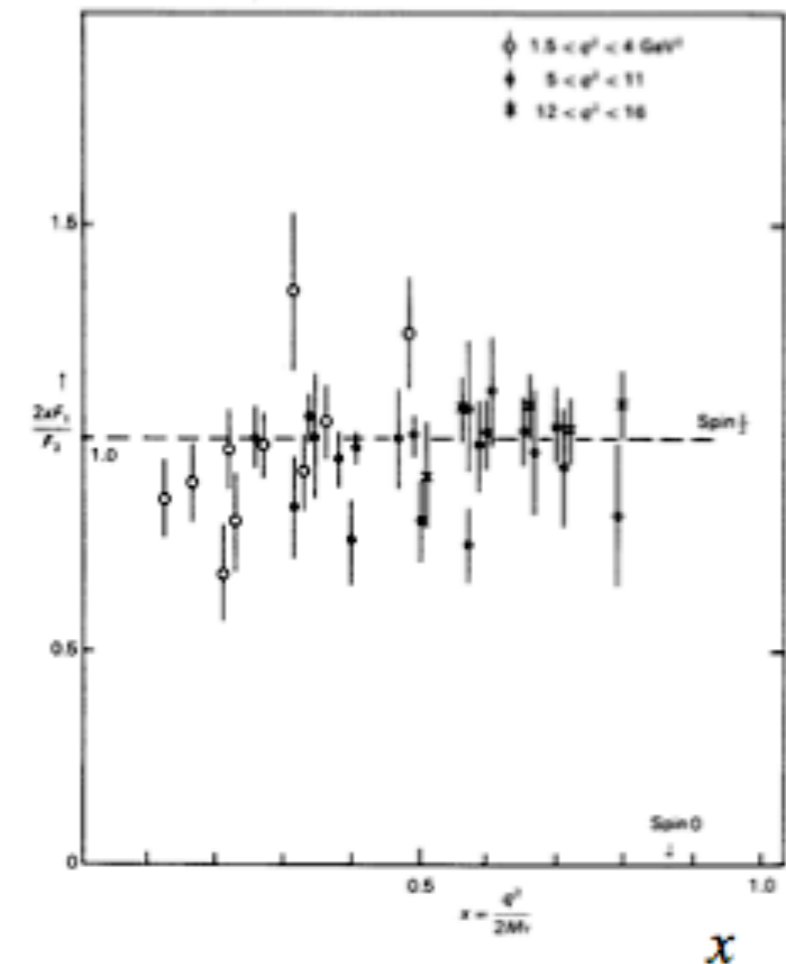


Point particles cannot be further resolved; their measurement does not depend on wavelength, hence Q^2 ,

Spin-1/2 quarks cannot absorb longitudinally polarized vector bosons and, conversely, spin-0 (scalar) quarks cannot absorb transversely polarized photons.

Callan-Gross relation:

$$\frac{2xF_1}{F_2}$$



← spin 1/2

← spin 0

Proton Structure

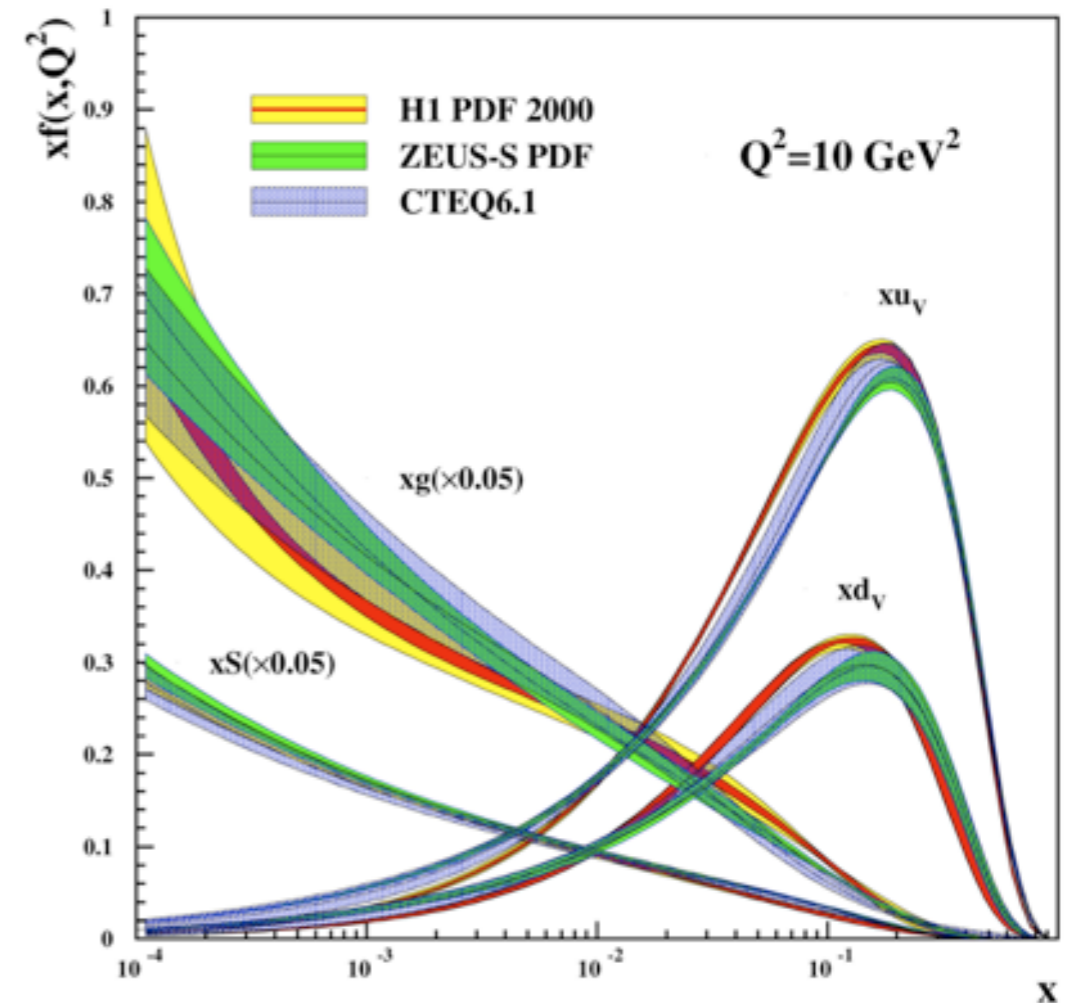
- Development of the (initial) quark-parton model,

$$F_1(x) = \frac{1}{2} \sum_q e_q^2 q(x)$$

$$F_2(x) = x \sum_q e_q^2 q(x) + \mathcal{O}(\alpha_s)$$

$$g_1(x) = \frac{1}{2} \sum_q e_q^2 \Delta q(x)$$

$$h_1(x) = \frac{1}{2} \sum_q e_q^2 \delta q(x)$$



improved by perturbative QCD radiation.

- Gluons *dominate*; responsible for $\sim 98\%$ of observed mass, carry $\sim 50\%$ of the proton momentum, *role(s) in nucleon spin?*

Proton Structure

- Development of the (initial) quark-parton model,

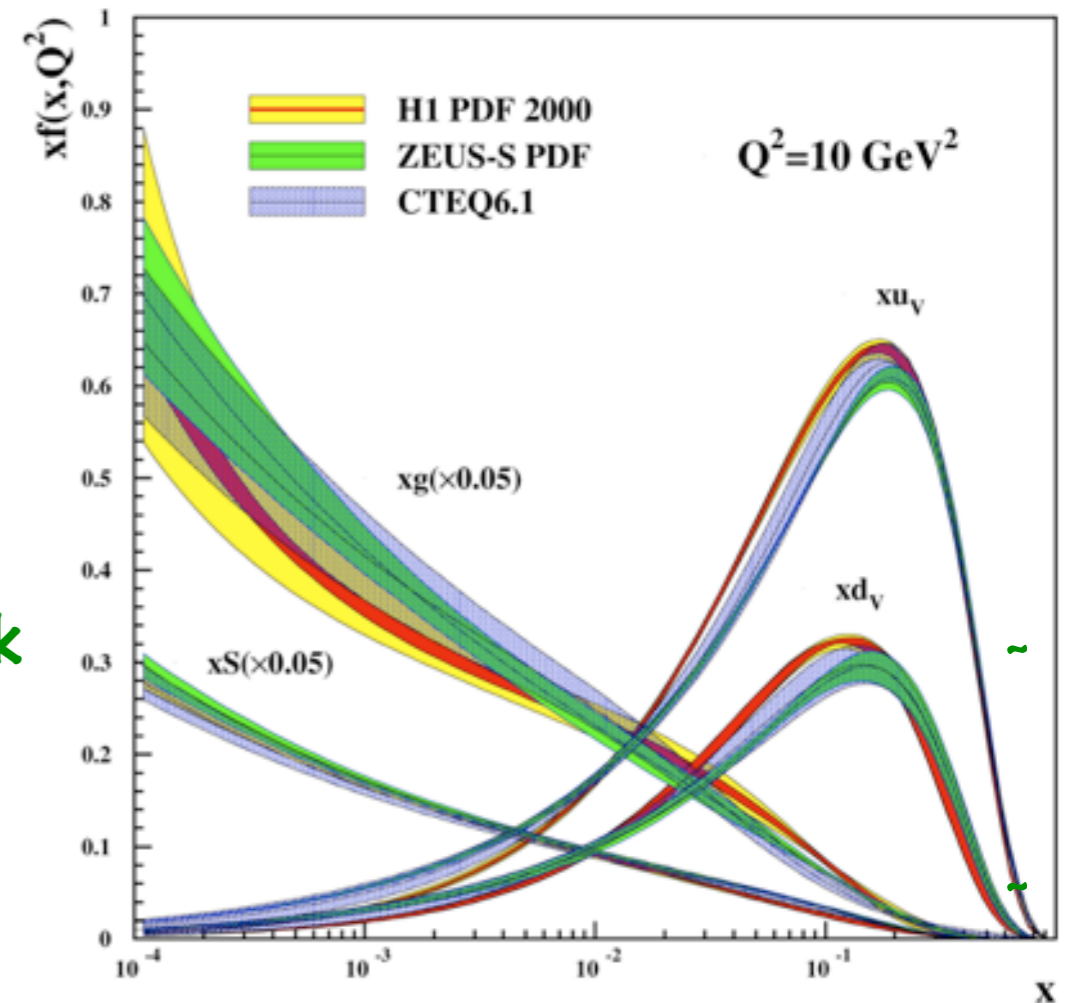
$$F_1(x) = \frac{1}{2} \sum_q e_q^2 q(x)$$

$$F_2(x) = x \sum_q e_q^2 q(x)$$

$$g_1(x) = \frac{1}{2} \sum_q e_q^2 \Delta q(x) \quad \text{much of this talk}$$

$$h_1(x) = \frac{1}{2} \sum_q e_q^2 \delta q(x) \quad \text{c.f. Marco's talk}$$

+ $\mathcal{O}(\alpha_s)$



improved by perturbative QCD radiation.

- Gluons *dominate*; responsible for $\sim 98\%$ of observed mass, carry $\sim 50\%$ of the proton momentum, *role(s) in nucleon spin?*

Proton Structure

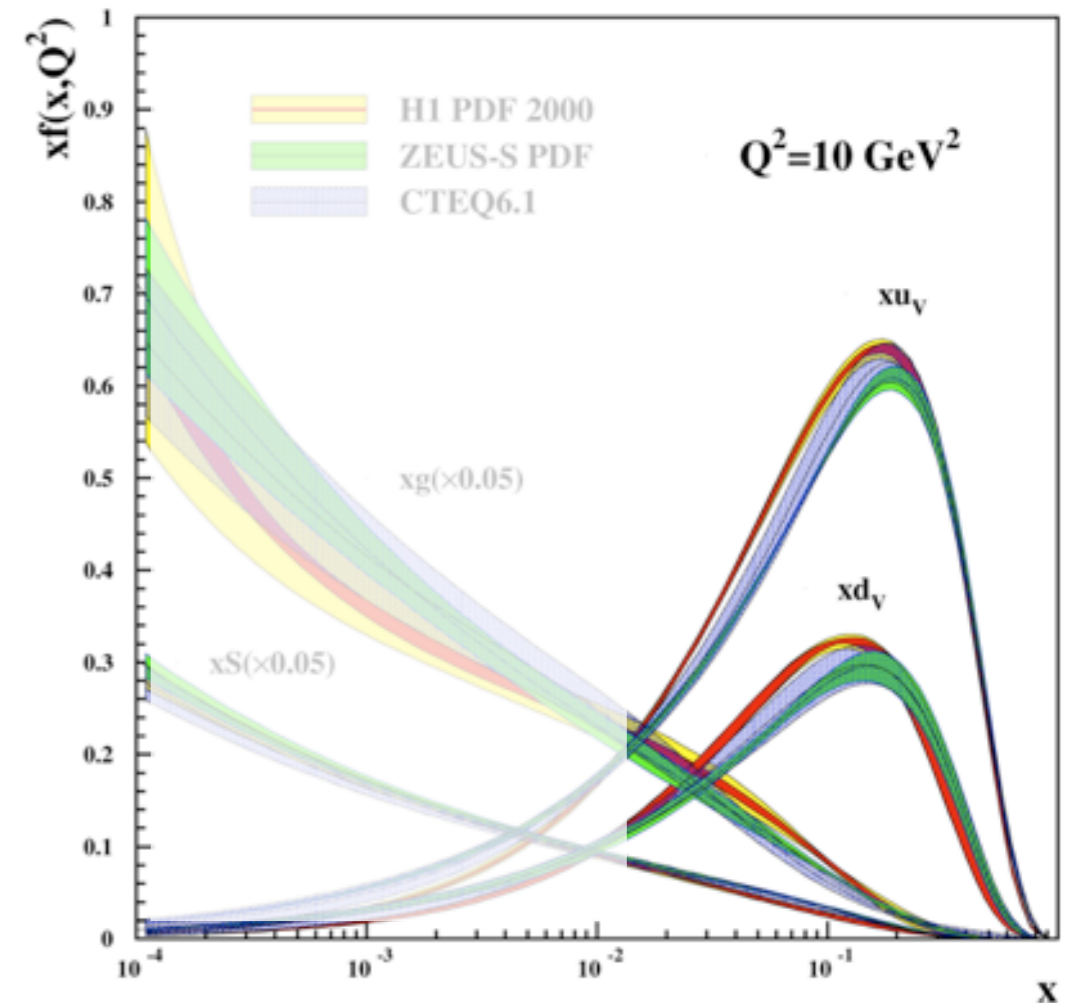
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$$h_1(x) = \frac{1}{2} \sum_q e_q^2 \delta q(x)$$



- Naively,

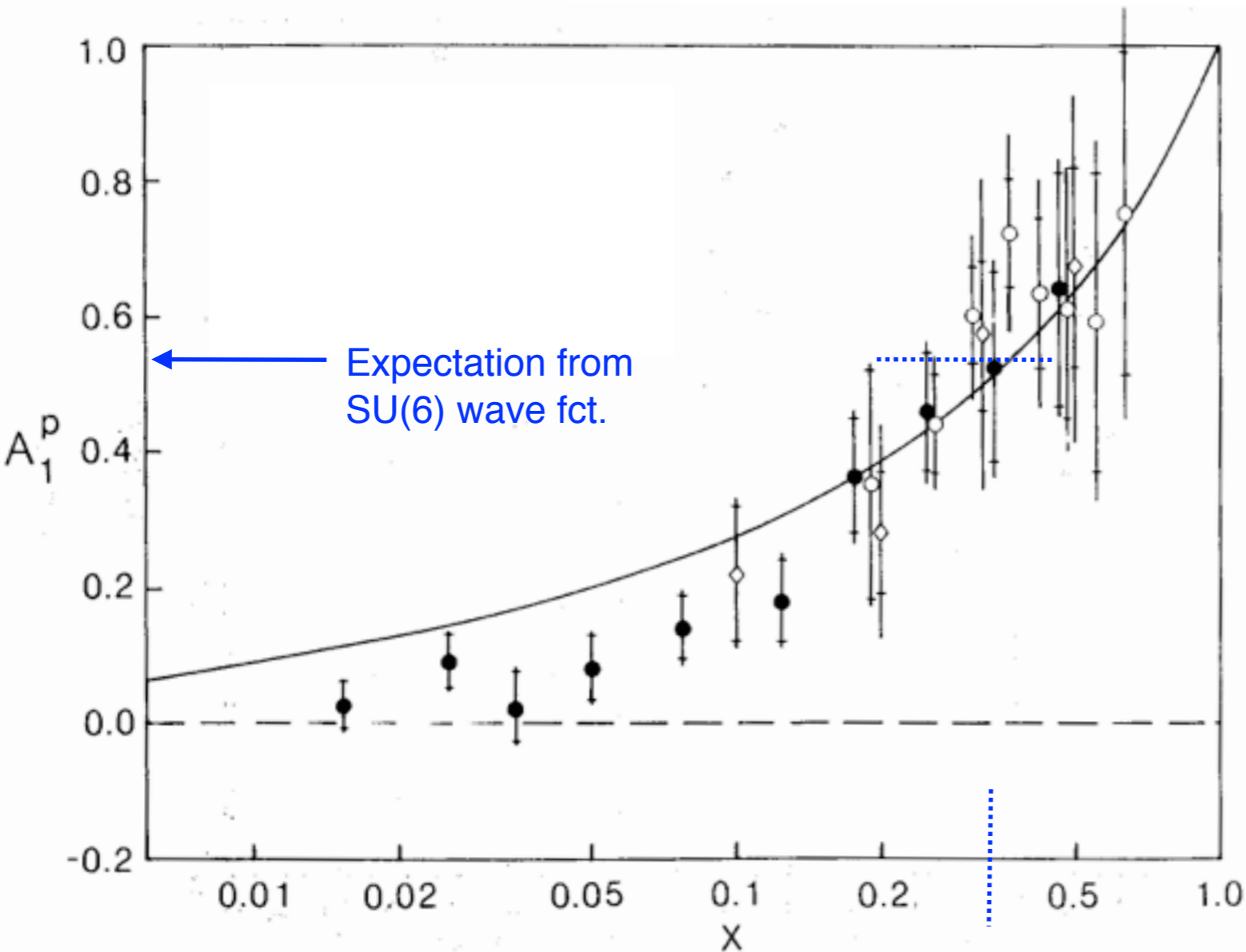
$$g_{\text{proton}}(\text{QPM}) \simeq 5.6 \quad \checkmark$$

and perhaps even more naive,

$$A_1^{\text{P}}(x) \equiv \frac{g_1(x)}{F_1(x)} \simeq \frac{5}{8}$$

Proton spin

Early measurements:



The spin asymmetry,

$$A_1^p \simeq \frac{\Delta\sigma}{\sigma} \simeq \frac{\sum e_q^2 (q^\uparrow - q^\downarrow)}{\sum e_q^2 (q^\uparrow + q^\downarrow)}$$

undershoots expectation.

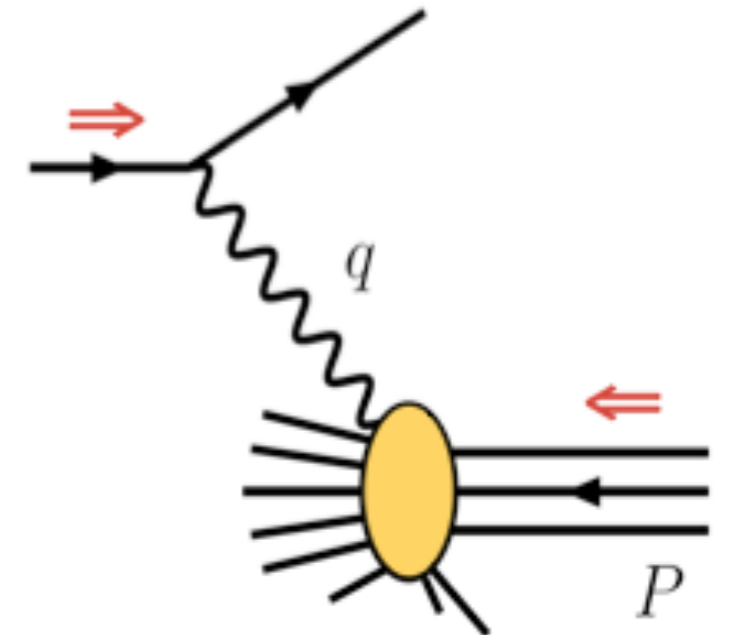
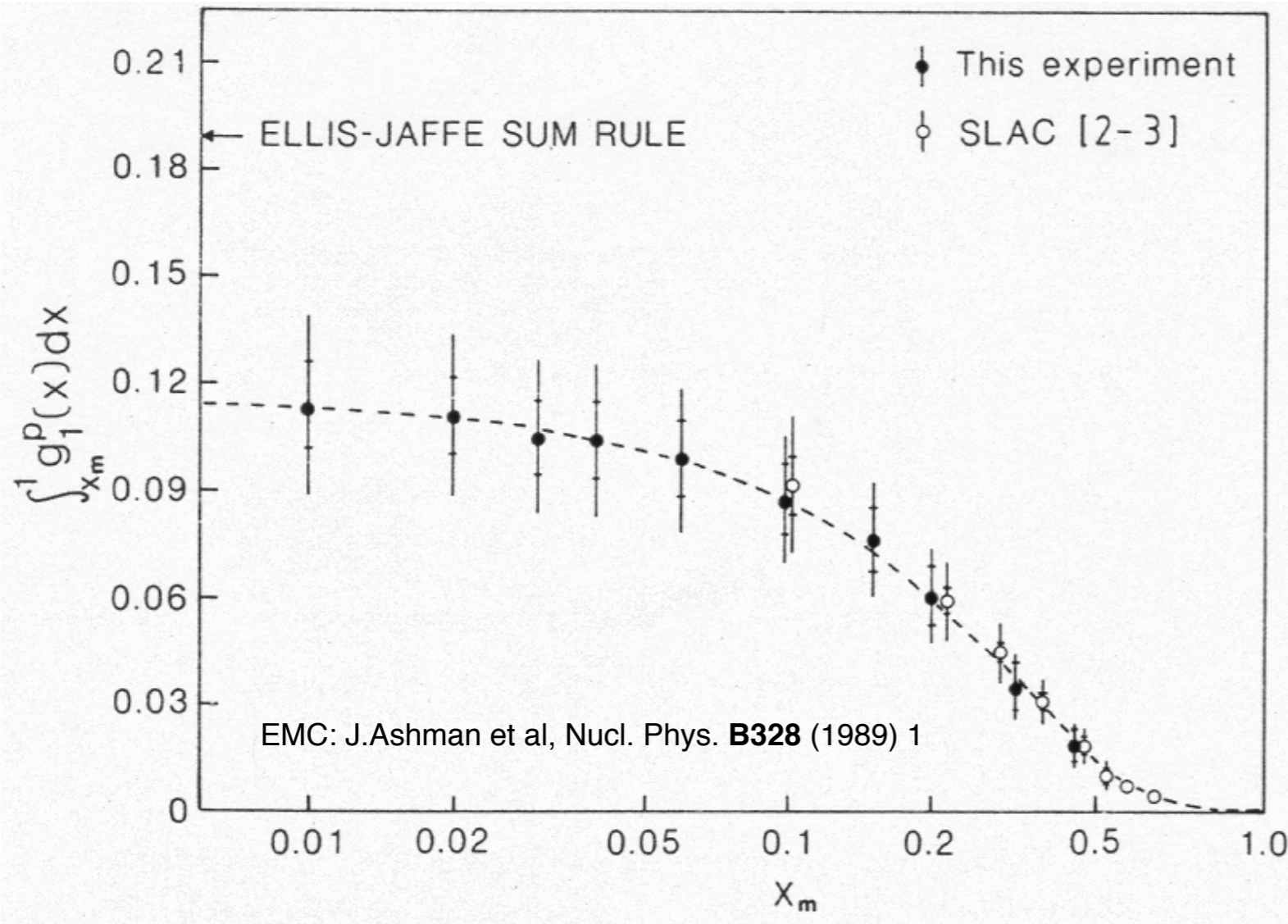
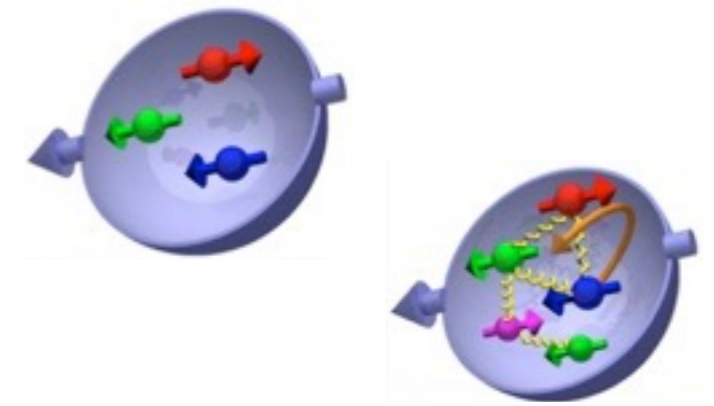
European Muon Collaboration (1989)

Naively,

$$A_1^p(x) \equiv \frac{g_1(x)}{F_1(x)} \simeq \frac{5}{8}$$

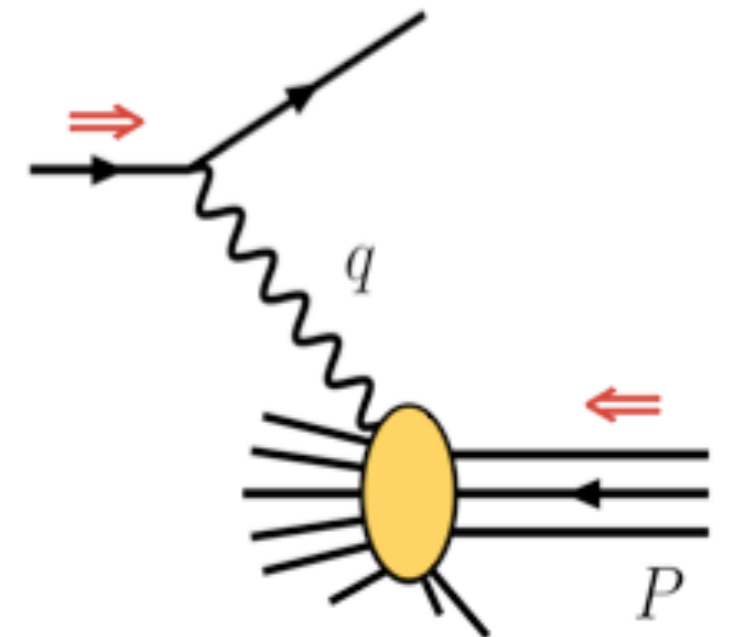
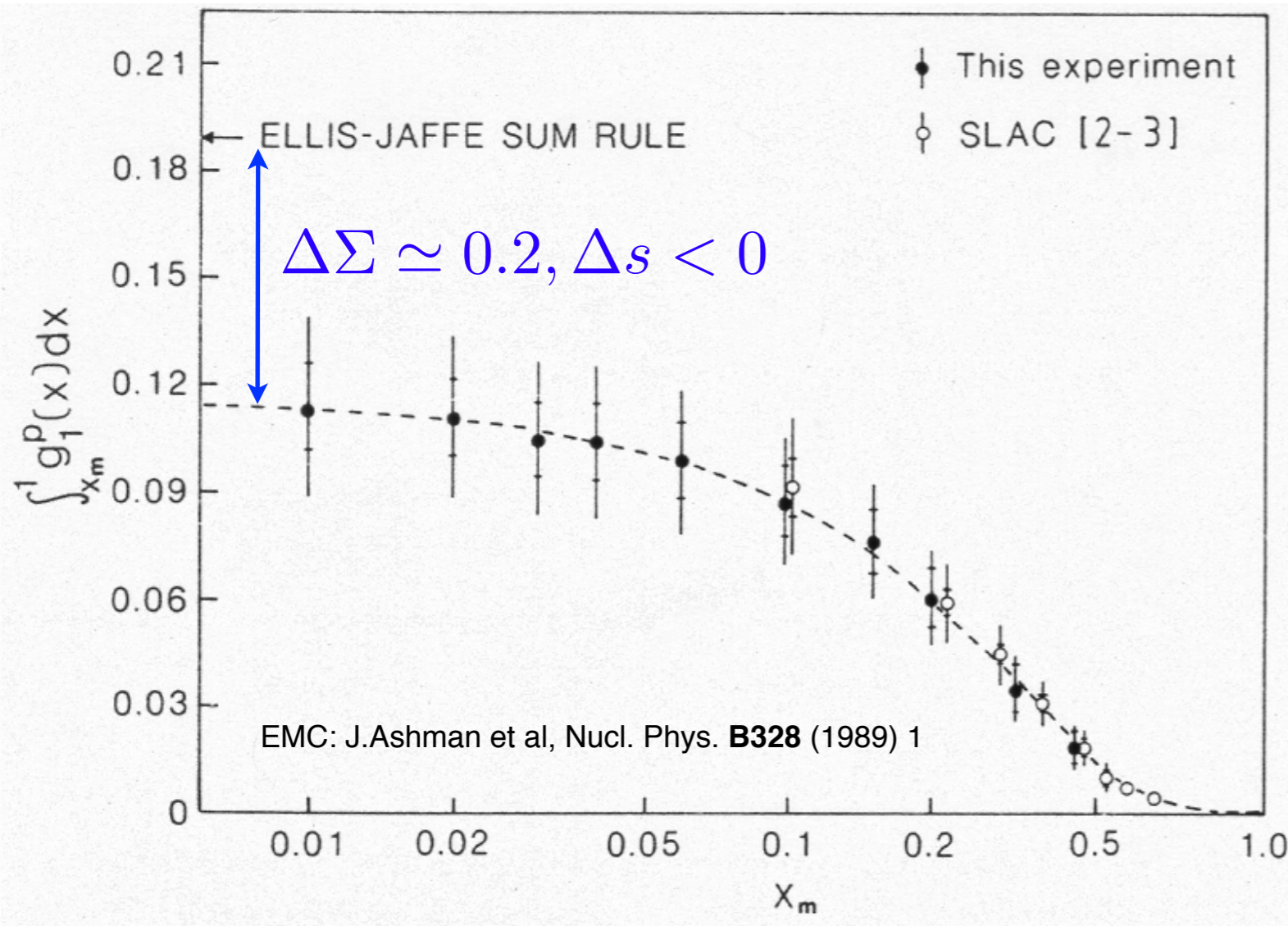
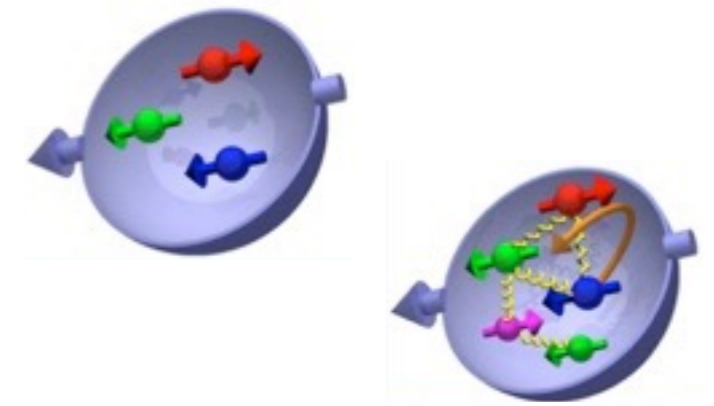
Better, relate $\int g_1(x) dx$ to the couplings in weak n and Σ^- decays.

Proton spin



$$\sigma(\Rightarrow, \Leftarrow) - \sigma(\Rightarrow, \Rightarrow) \sim g_1(x, Q^2)$$

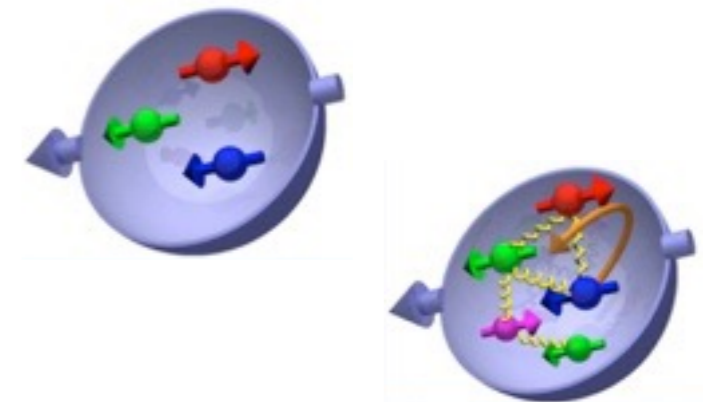
Proton spin



$$\sigma(\Rightarrow, \Leftarrow) - \sigma(\Rightarrow, \Rightarrow) \sim g_1(x, Q^2)$$

“ The sum of quark and anti-quark spins contribute little to the proton spin, and strange quarks are negatively polarized. ”

Proton spin



For the proton,

$$\Gamma_1 = \int_0^1 g_1(x) dx = \int_0^1 \left(\frac{1}{2} \sum e_q^2 \Delta q(x) \right) dx = \frac{1}{2} \left(\frac{4}{9} \Delta_1 u + \frac{1}{9} \Delta_1 d + \frac{1}{9} \Delta_1 s \right)$$

$$= \frac{1}{12} (\Delta_1 u - \Delta_1 d) + \frac{1}{36} (\Delta_1 u + \Delta_1 d - 2\Delta_1 s) + \frac{1}{9} (\Delta_1 u + \Delta_1 d + \Delta_1 s)$$

Known from weak neutron to proton decay

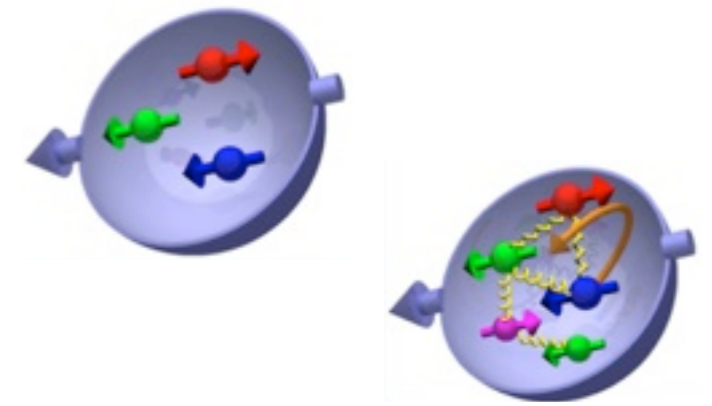
Known from weak neutron to proton decay,
combined with weak Σ to neutron decay

Unique to DIS, $\Delta\Sigma$

which becomes a prediction if $\Delta_1 s = 0$

No (reliable) substitute for energy; $x \propto 1/\sqrt{s}$


Proton spin



For the proton,

$$\Gamma_1 = \int_0^1 g_1(x) dx = \int_0^1 \left(\frac{1}{2} \sum e_q^2 \Delta q(x) \right) dx = \frac{1}{2} \left(\frac{4}{9} \Delta_1 u + \frac{1}{9} \Delta_1 d + \frac{1}{9} \Delta_1 s \right)$$


$$= \frac{1}{12} (\Delta_1 u - \Delta_1 d) + \frac{1}{36} \underbrace{(\Delta_1 u + \Delta_1 d - 2\Delta_1 s)}_{a_8 = 3F - D = 0.59 \pm 0.03} + \frac{1}{9} (\Delta_1 u + \Delta_1 d + \Delta_1 s)$$



 $< 10\%$



 Unique to DIS, $\Delta\Sigma$



 Known from weak neutron to proton decay, combined with weak Σ to neutron decay

Since,

$$\left. \begin{array}{l} \frac{\partial \Gamma_1}{\partial a_8} \Big|_{\text{Ellis-Jaffe}} \simeq \frac{5}{36} \\ \frac{\partial \Gamma_1}{\partial a_8} \Big|_{\text{experiment}} \simeq 0 \end{array} \right\} \text{one can recover the E-J expectation with a sizable shift of } a_8 = 3F - D, \quad a_8 \simeq 0.2 \pm 0.1$$

Proton spin - SU(3)

Such a *sizable* shift, however, is hard to support from (new) data:

Table 1. Present world HSD rate and angular-correlation data [14]. Numerical values marked g_1/f_1 are as extracted from angular and spin correlations.

Decay $A \rightarrow B\ell\nu$	Rate(10^6 s^{-1})		g_1/f_1		g_1/f_1	
	$\ell = e^\pm$	$\ell = \mu^-$	$\ell = e^-$		SU(3)	
$n \rightarrow p$	1.1291 ± 0.0010		1.2670 ± 0.0030		$F + D$	} Close & Roberts
$\Lambda^0 \rightarrow p$	3.161 ± 0.058	0.60 ± 0.13	0.718 ± 0.015		$F + \frac{1}{3}D$	
$\Sigma^- \rightarrow n$	6.88 ± 0.23	3.04 ± 0.27	-0.340 ± 0.017		$F - D$	
$\Sigma^- \rightarrow \Lambda^0$	0.387 ± 0.018				$-\sqrt{\frac{2}{3}} D^\dagger$	
$\Sigma^+ \rightarrow \Lambda^0$	0.250 ± 0.063				$-\sqrt{\frac{2}{3}} D^\dagger$	
$\Xi^- \rightarrow \Lambda^0$	3.35 ± 0.37	2.1 ± 2.1	0.25 ± 0.05		$F - \frac{1}{3}D$	
$\Xi^- \rightarrow \Sigma^0$	0.53 ± 0.10				$F + D$	
$\Xi^0 \rightarrow \Sigma^+$	0.876 ± 0.071	$0.012 \pm 0.007^*$	1.32 ± 0.21		$F + D$	KTeV

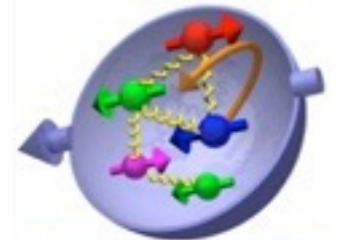
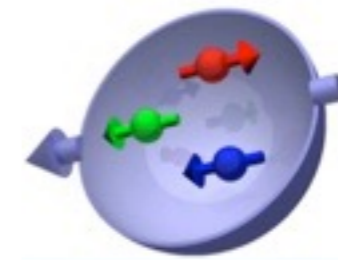
* KTeV data [5]—not included in the fits presented here.

† The absolute expression for g_1 is given, not g_1/f_1 (as $f_1 = 0$).

P. Ratcliffe, Czech J. Phys. 54, A21 (2004).

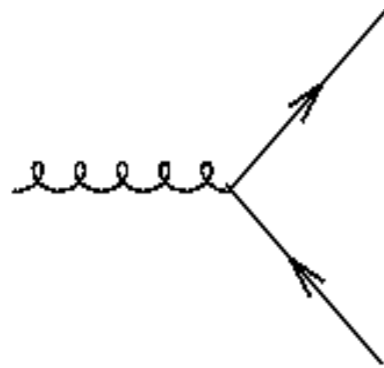
Smaller shifts have been reported, e.g. by Ehrnsperger (1995), Song (1997), Flores-Mendieta (1998), Yamanashi (2007), Sasaki (2009), ...

Proton spin



Numerous follow-up questions and experiment programs,

Among the early attempts at a resolution,



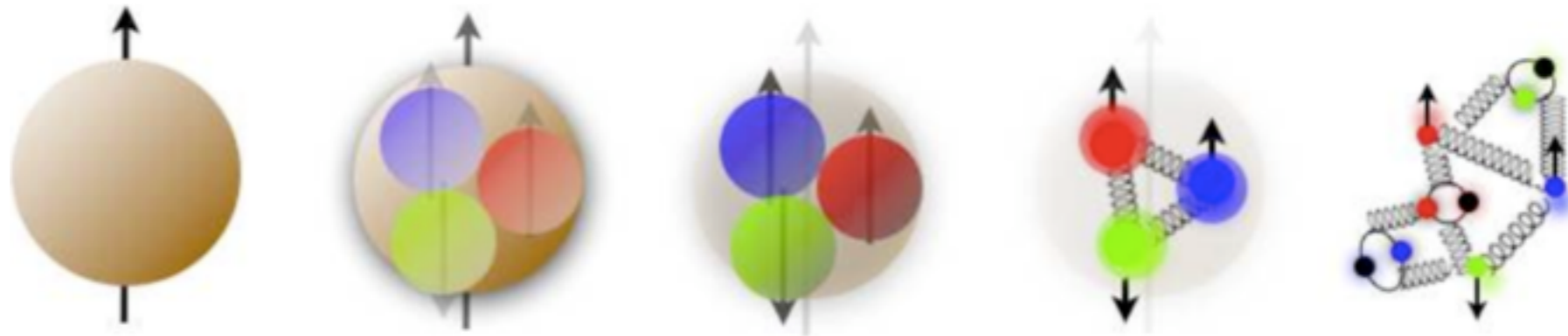
with the gluons *polarized*.

G. Altarelli and G.G. Ross Phys. Lett. **B212** (1998) 391

Note: this attempt requires *very* large polarization, *factors* larger than the nucleon spin itself, and by inference, *huge* compensating *orbital momenta*. Quite the proton, a ground-state object and all.

Other attempts include e.g extrapolation over unmeasured low-x.

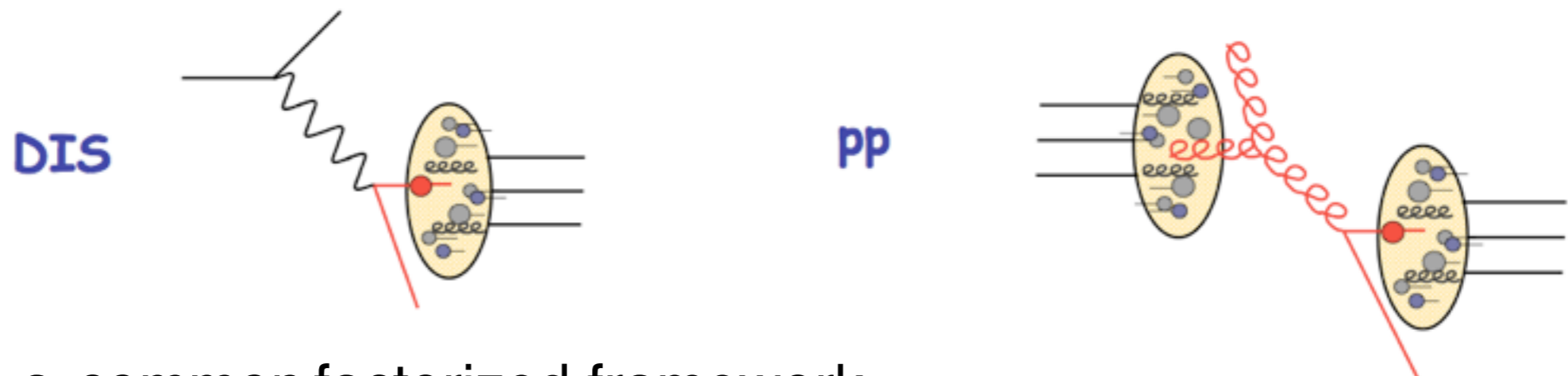
Renewed Interest in Nucleon Spin



What is the nucleon spin decomposition?

What is the role of spin in QCD?

Complementary ways to address several of the open questions,



within a common factorized framework.

Nucleon Spin - A World-Wide Quest



SLAC

E142, E143,
E154, E155



CERN

EMC, SMC,
COMPASS



DESY

HERMES



JLab

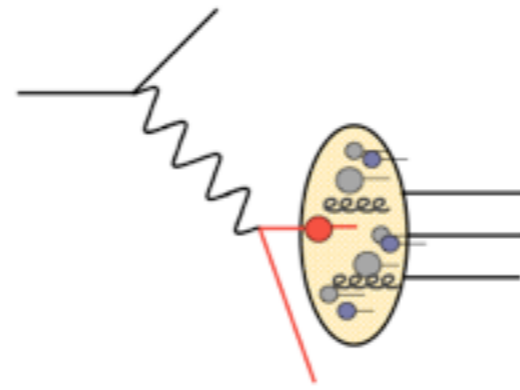
Hall A, CLAS



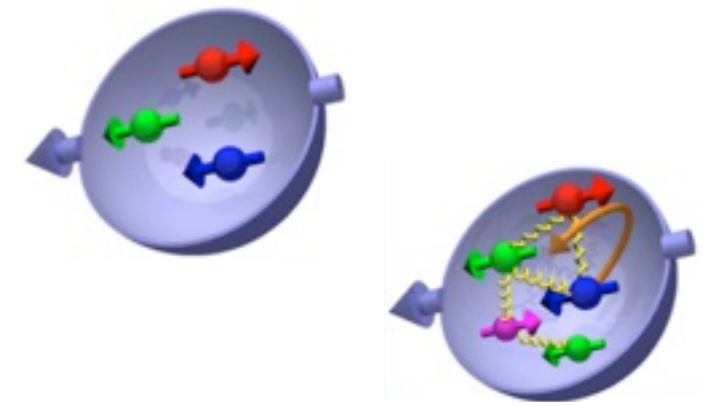
BNL

PHENIX, STAR

DIS



DIS - Nucleon spin



For the proton,

$$\Gamma_1 = \int_0^1 g_1(x) dx = \int_0^1 \left(\frac{1}{2} \sum e_q^2 \Delta q(x) \right) dx = \frac{1}{2} \left(\frac{4}{9} \Delta_1 u + \frac{1}{9} \Delta_1 d + \frac{1}{9} \Delta_1 s \right)$$

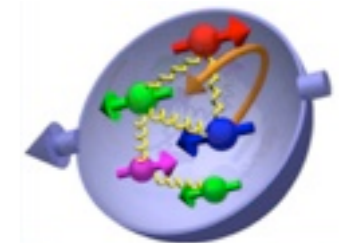
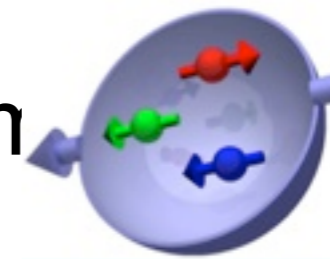
$$= \frac{1}{12} (\Delta_1 u - \Delta_1 d) + \frac{1}{36} (\Delta_1 u + \Delta_1 d - 2\Delta_1 s) + \frac{1}{9} (\Delta_1 u + \Delta_1 d + \Delta_1 s)$$

\uparrow
 $\checkmark < 10\%$

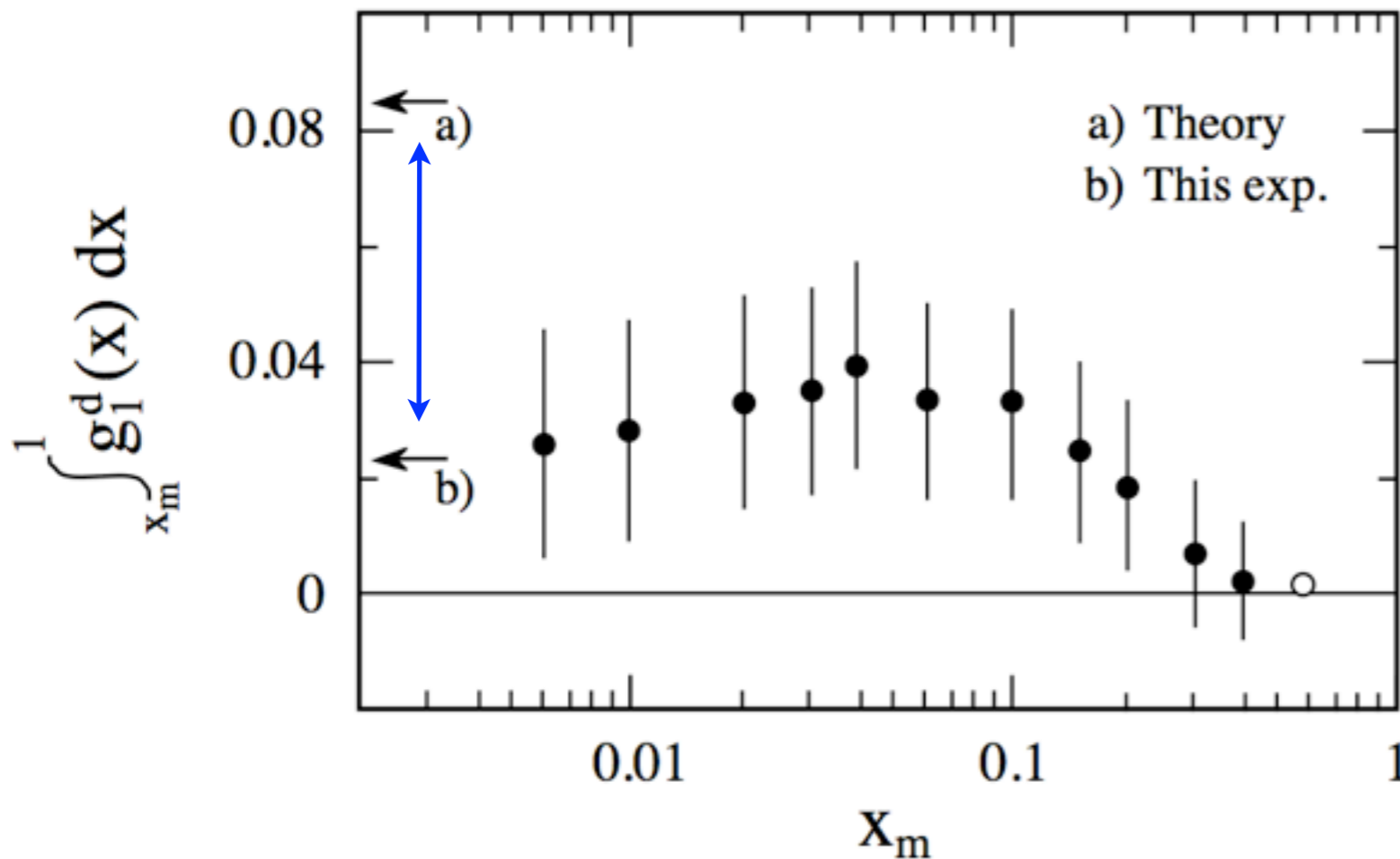
$+ \mathcal{O}(\alpha_s)$ now well known $+ \mathcal{O}(1/Q^2)$

Similar can be done for the neutron,
 experimentally via the deuteron or ^3He ,

Nucleon Spin - Bjorken Sun



$$\int_{0.006}^{0.6} g_1^d(x) dx = 0.024 \pm 0.020_{\text{stat}} \pm 0.014_{\text{sys}}$$



SMC: B.Adeva et al, Phys. Lett. B302, 533 (1993)

$$\Gamma_1^p = 0.126 \pm 0.010_{\text{stat}} \pm 0.015_{\text{sys}} \quad (\text{EMC})$$

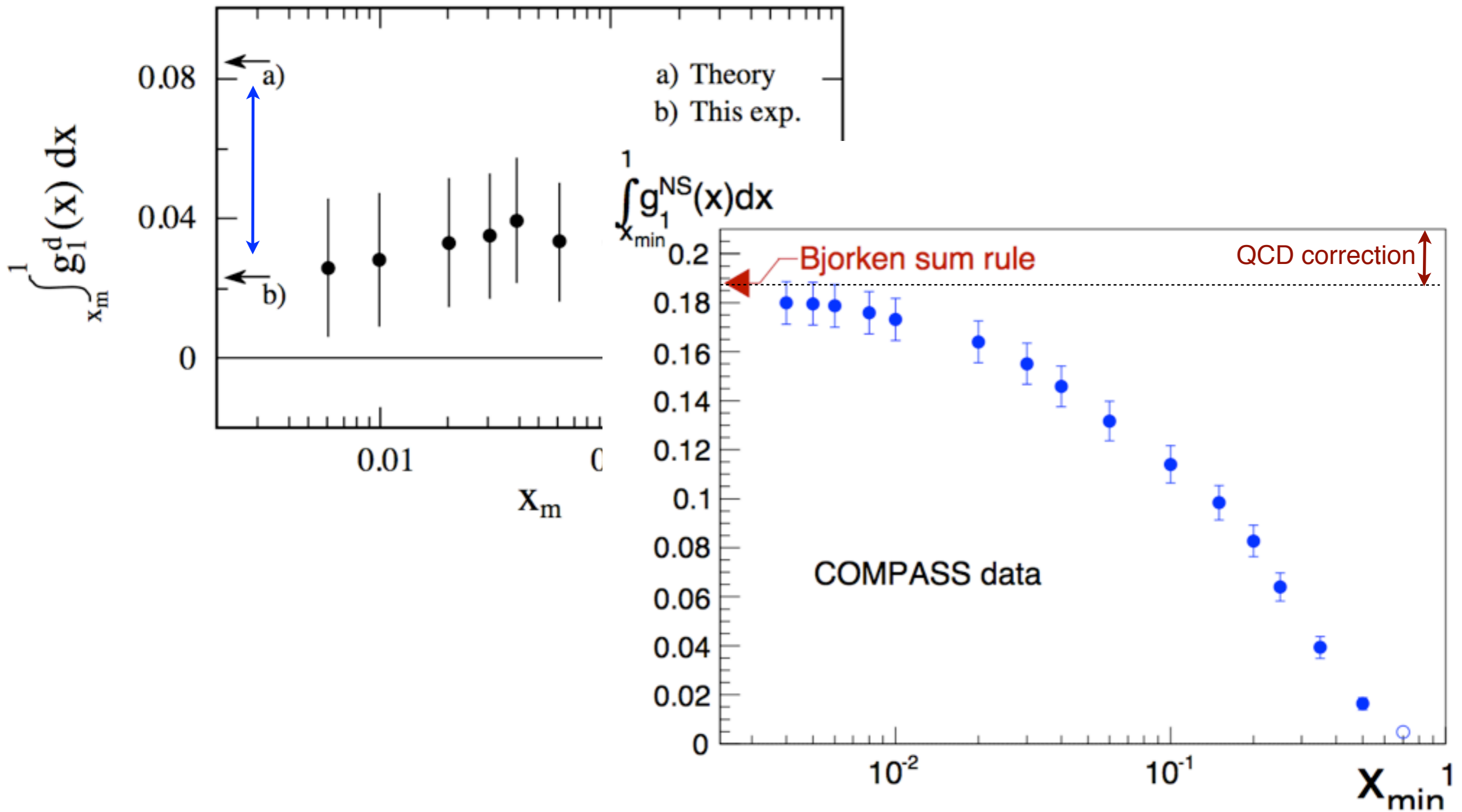
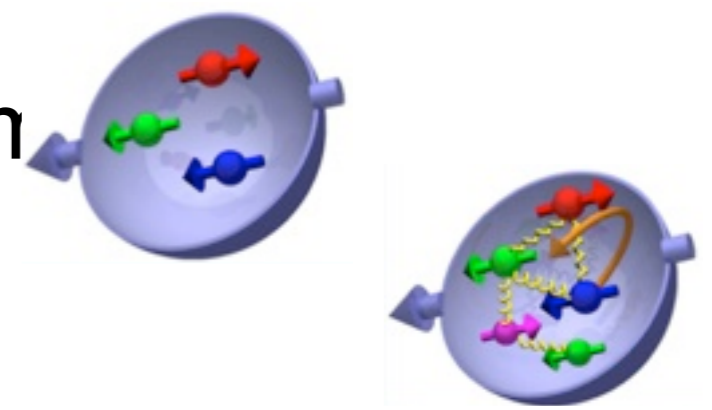
$$\Gamma_1^n = -0.08 \pm 0.04_{\text{stat}} \pm 0.04_{\text{sys}}$$

$$\Gamma_1^p - \Gamma_1^n = 0.20 \pm 0.05_{\text{stat}} \pm 0.04_{\text{sys}}$$

$$\Gamma_1^p - \Gamma_1^n = \frac{1}{6} \left| \frac{g_A}{g_V} \right| \left[1 - \frac{\alpha_s}{\pi} + \dots \right] \quad \checkmark \quad 25\%$$

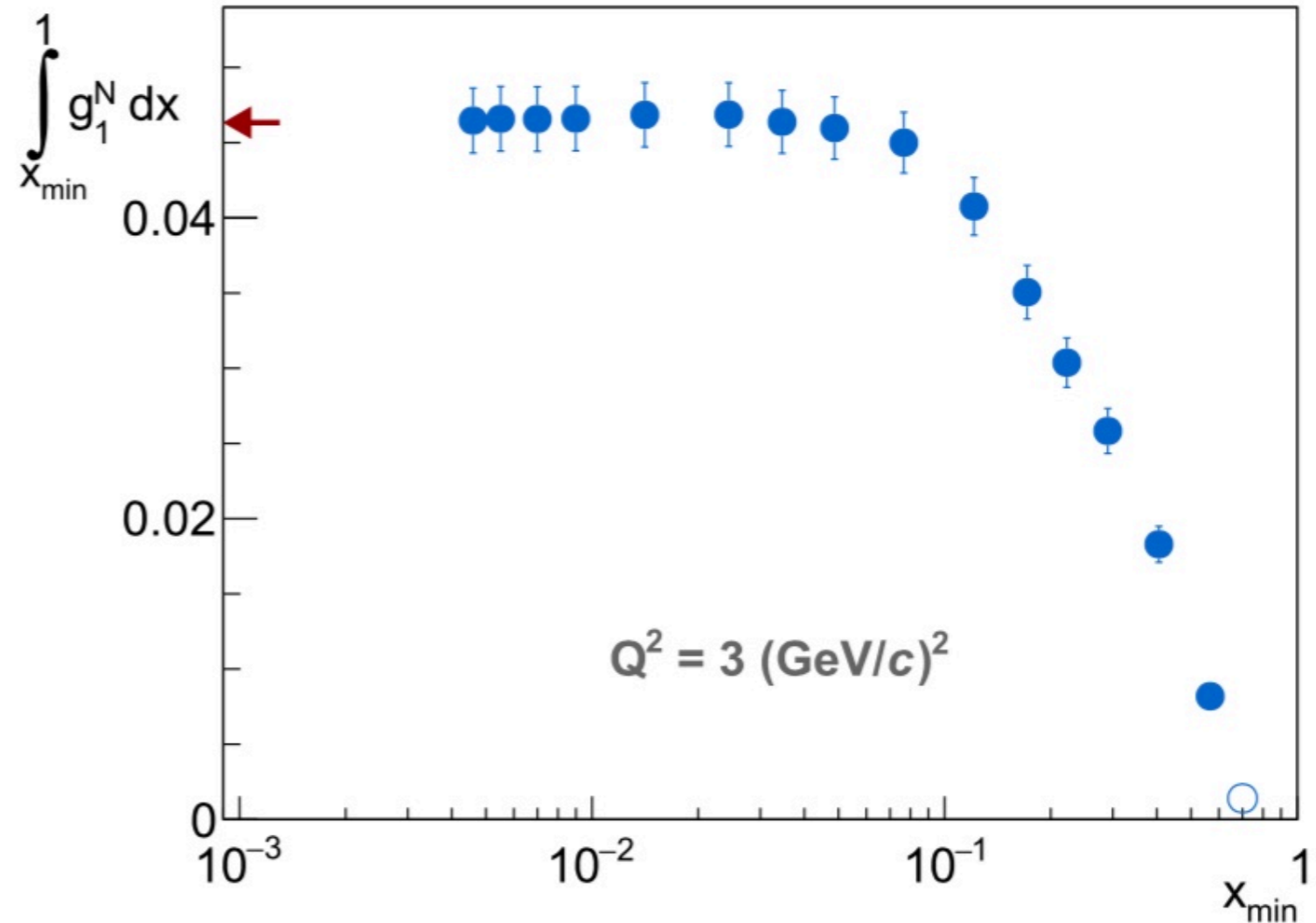
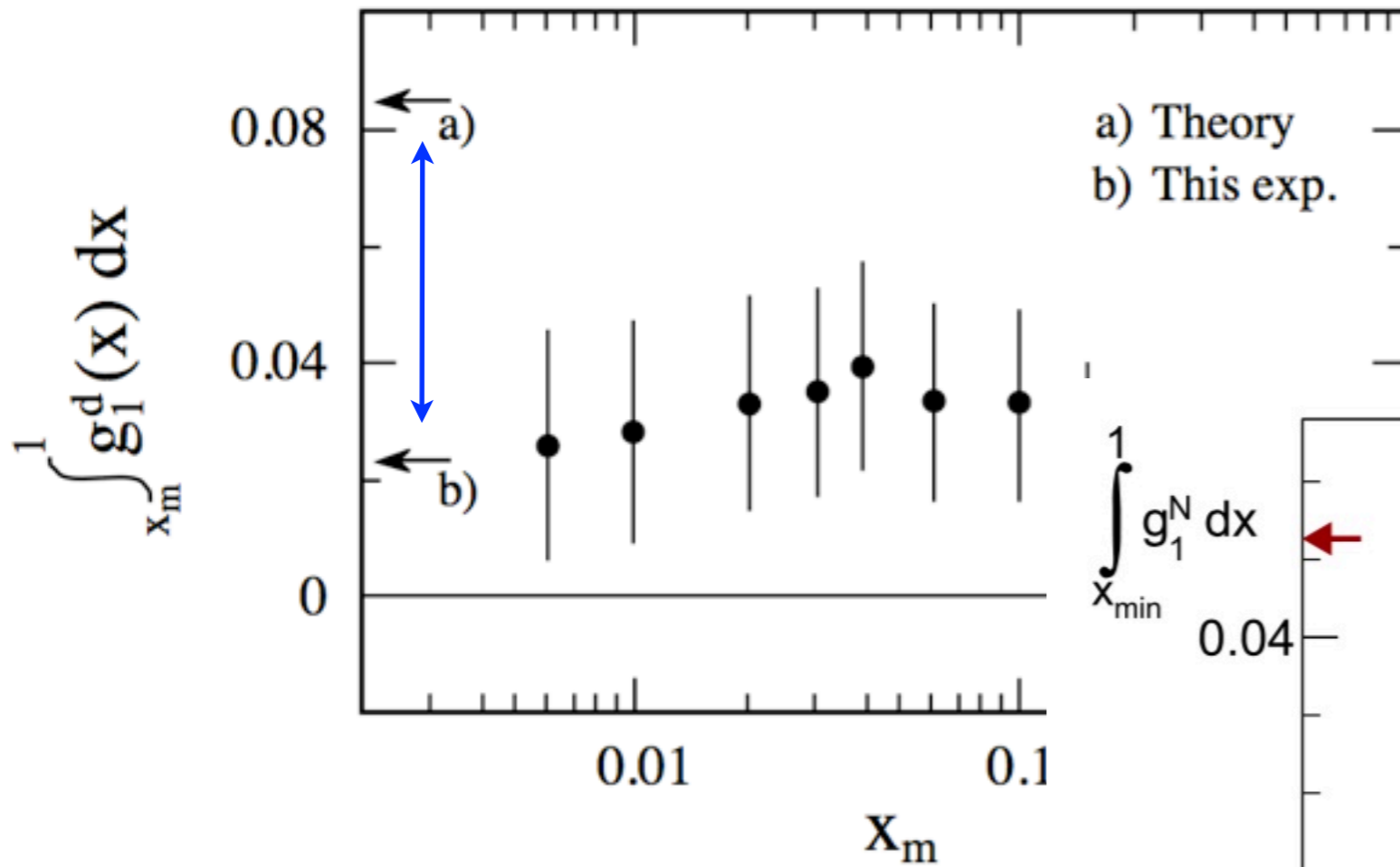
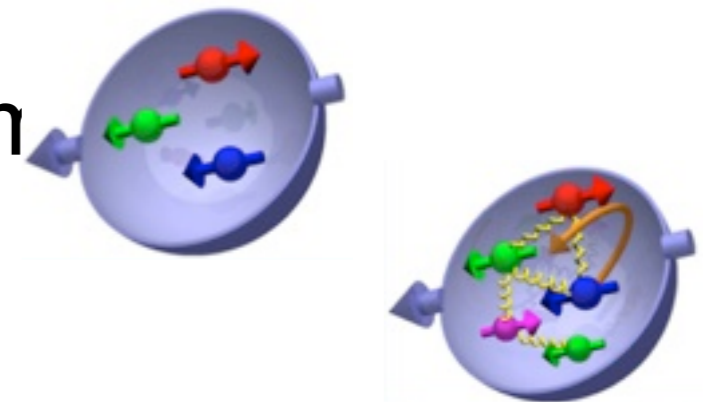
$$= 0.191 \pm 0.002 + \mathcal{O}(\alpha_s^2)$$

Nucleon Spin - Bjorken Sun



Nucleon Spin - Bjorken Sum

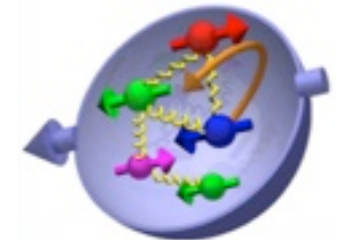
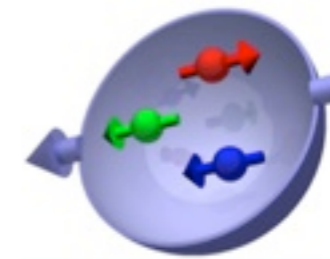
✓ < 10%



$$\Gamma_1^{NS} = 0.192 \pm 0.007_{\text{stat}} \pm 0.015_{\text{syst}}$$

$$|g_A/g_V| = 1.29 \pm 0.05_{\text{stat}} \pm 0.10_{\text{syst}}$$

Nucleon spin - $\Delta\Sigma$



$$a_8 = \underbrace{\Delta_1 u + \Delta_1 d - 2\Delta_1 s}_{3F-D} = 0.59 \pm 0.03$$

↑ Known from weak neutron to proton decay, combined with weak Σ to neutron decay

Table 1. Present world HSD rate and angular-correlation data [14]. Numerical values marked g_1/f_1 are as extracted from angular and spin correlations.

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$\Xi^0 \rightarrow \Sigma^+$	0.876 \pm 0.071	0.012 \pm 0.007 *	1.32 \pm 0.21		$F + D$

} Close & Roberts

KTeV

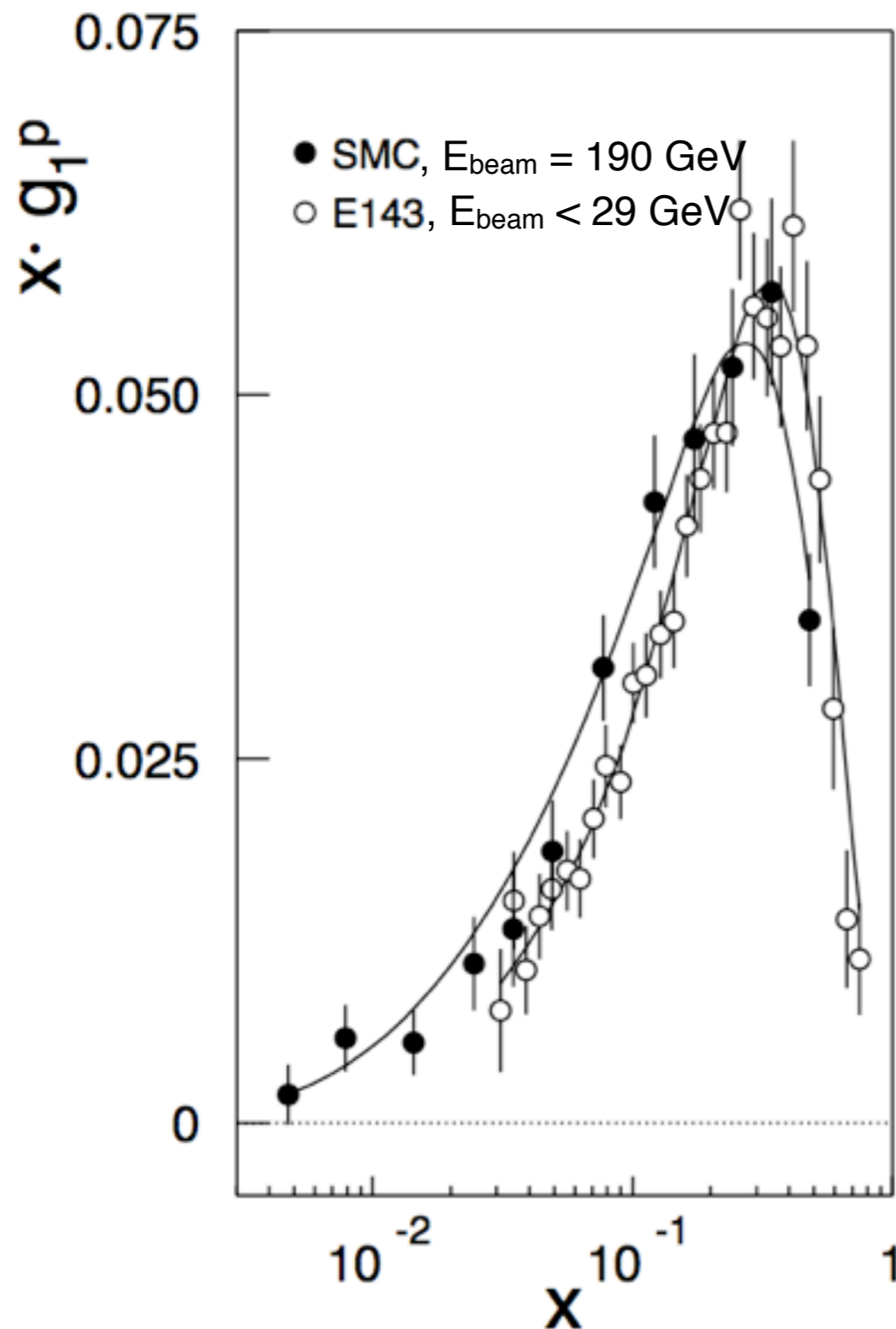
* KTeV data [5]—not included in the fits presented here.

† The absolute expression for g_1 is given, not g_1/f_1 (as $f_1 = 0$).



→ $a_0(Q^2 = 3 (\text{GeV}/c)^2) = 0.32 \pm 0.02_{\text{stat}} \pm 0.04_{\text{syst}} \pm 0.05_{\text{evol}}.$

Nucleon Spin - Scaling Violations, Global Analyses



Provide sensitivity to $\Delta G(x, Q^2)$

Closely related also to extrapolations over unmeasured small- x ,

$$g_1(x, Q^2) \propto \exp \sqrt{\ln \frac{1}{x} \ln \ln \frac{Q^2}{\Lambda^2}},$$

as $x \rightarrow 0$, $Q^2 \rightarrow \infty$,

R.D. Ball et al, Nucl. Phys. B444, 287 (1995),
B449, 680 (1995)

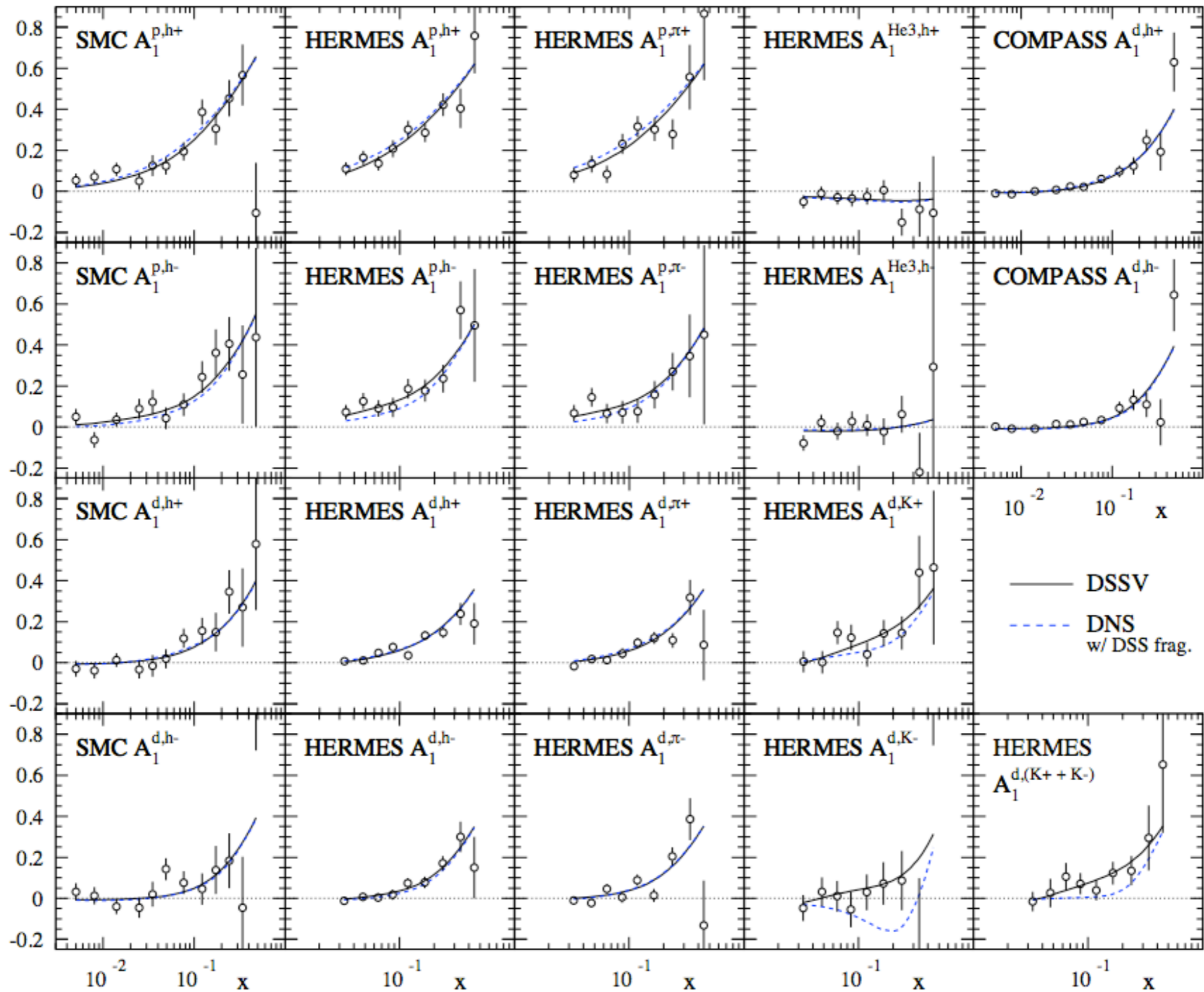
Start of polarized pQCD analyses.

More recent work at NLO:

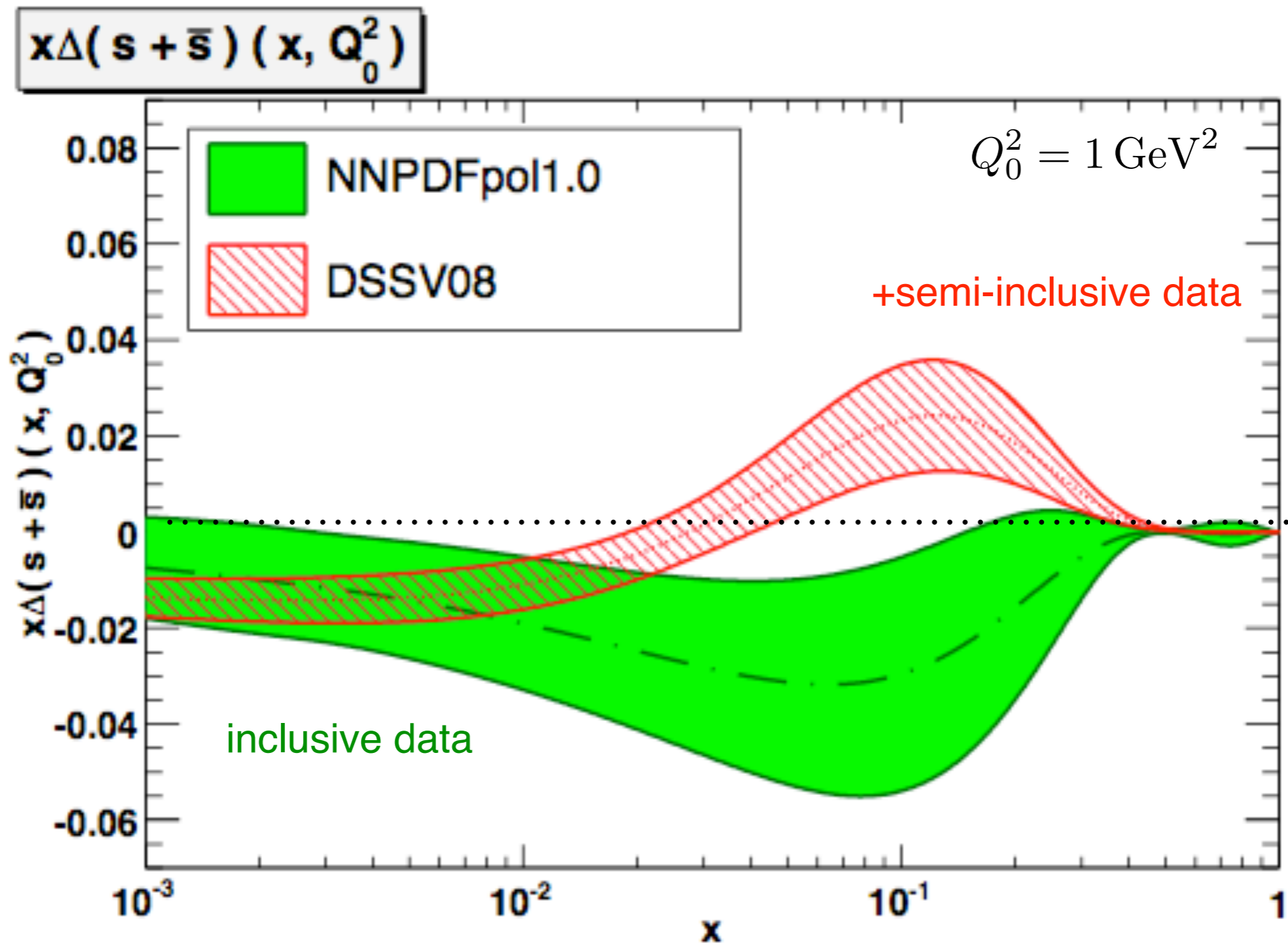
NNPDF: Ball et al,
initially inclusive DIS data only,
Nucl. Phys. B874, 36 (2013)

DSSV: De Florian et al,
inclusive DIS, SIDIS, RHIC
ArXiv 1304.0079 (2013)

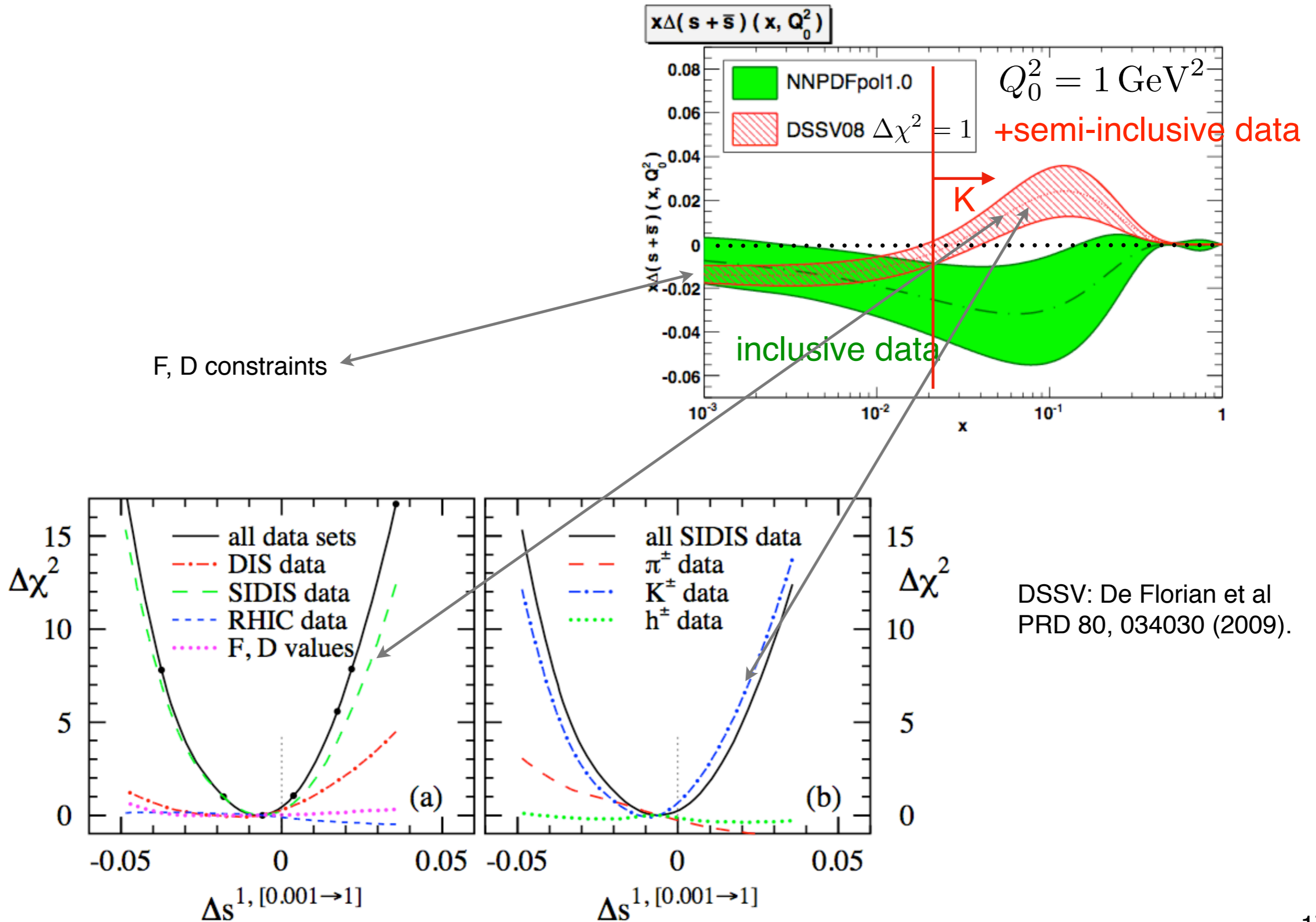
SIDIS - Quark Flavor Structure



Nucleon spin - strange quark polarization, Δs



Nucleon spin - strange quark polarization, Δ_s

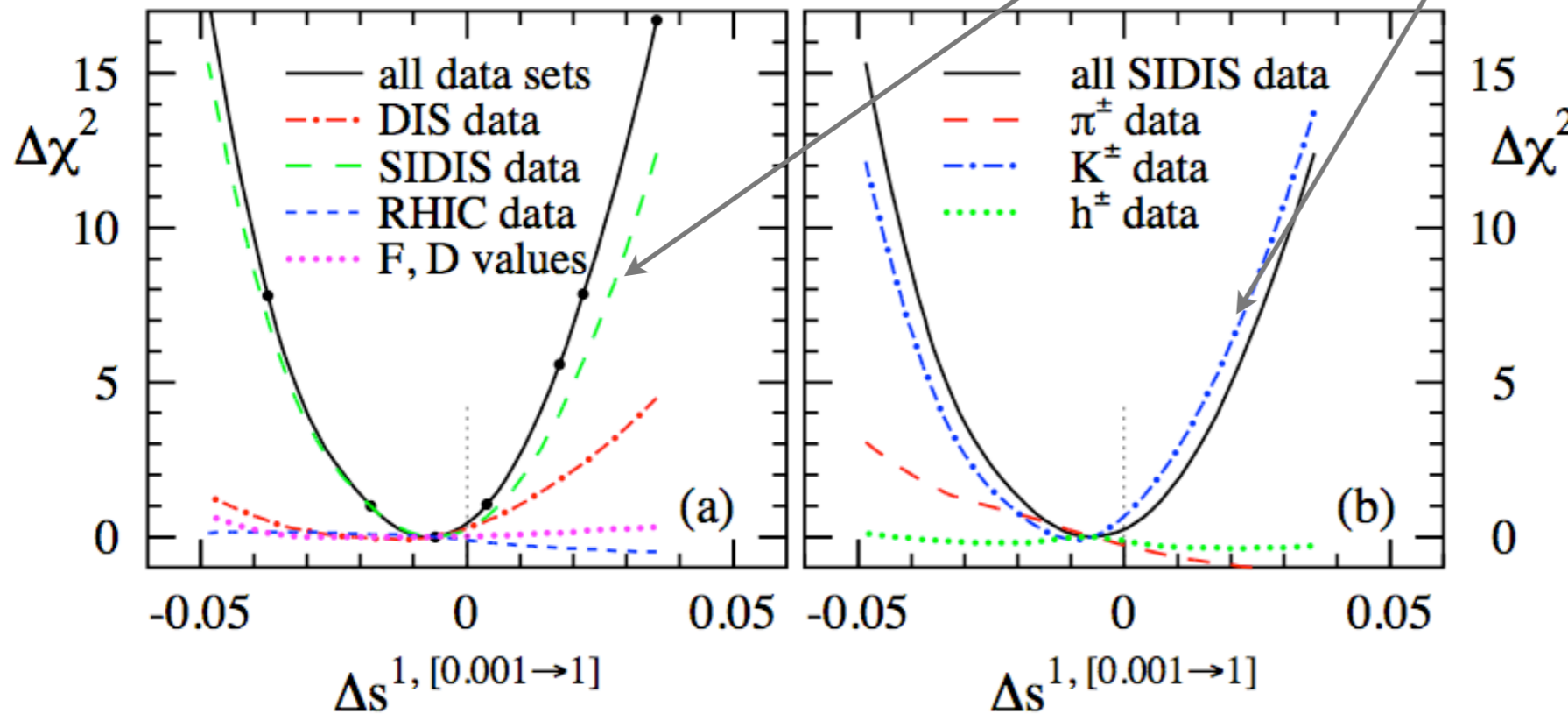
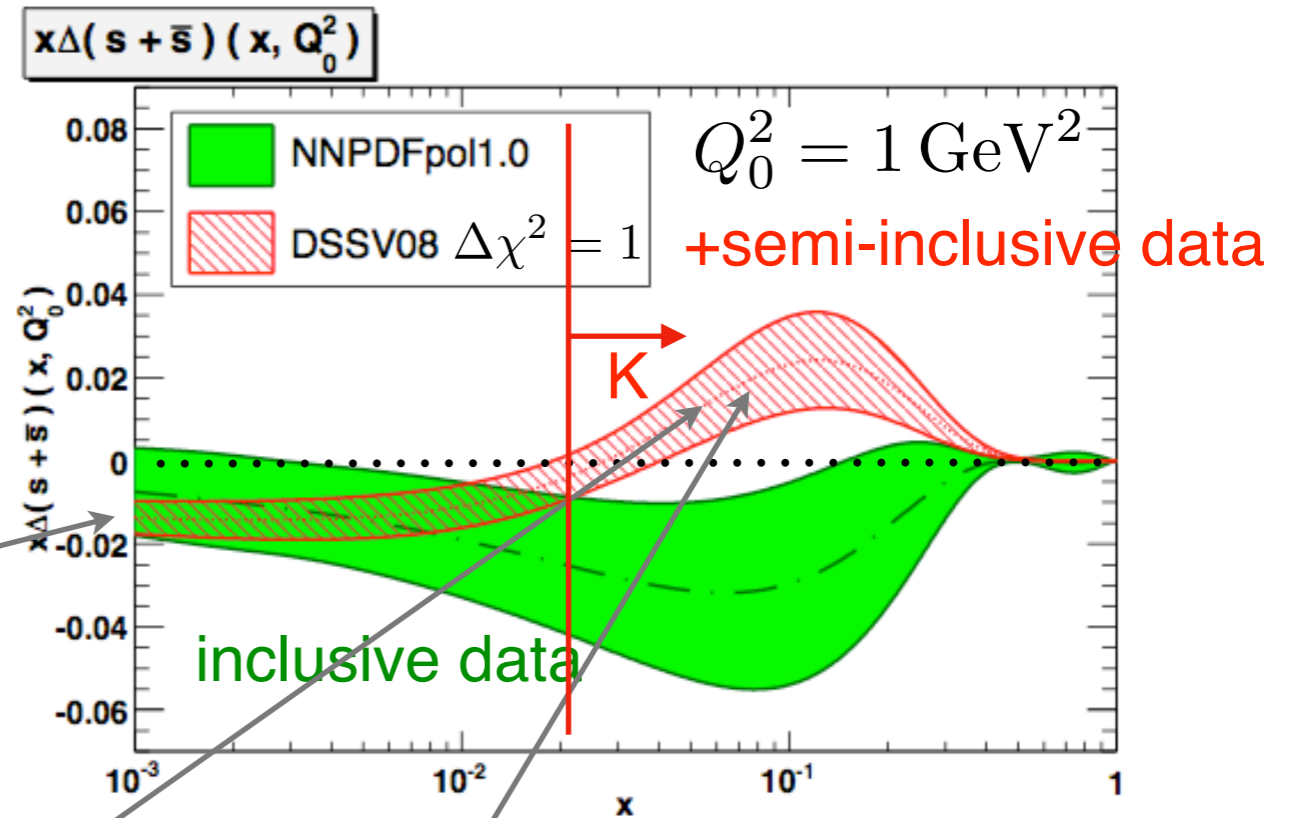


Nucleon spin - strange quark polarization, Δs

SIDIS relies on FF,
 in particular to Kaon FF,
 no "tension" in DSSV though,
 Larger F, D uncertainty in NNPDF
 Lattice indicative of small net Δs

Clear role for a future EIC

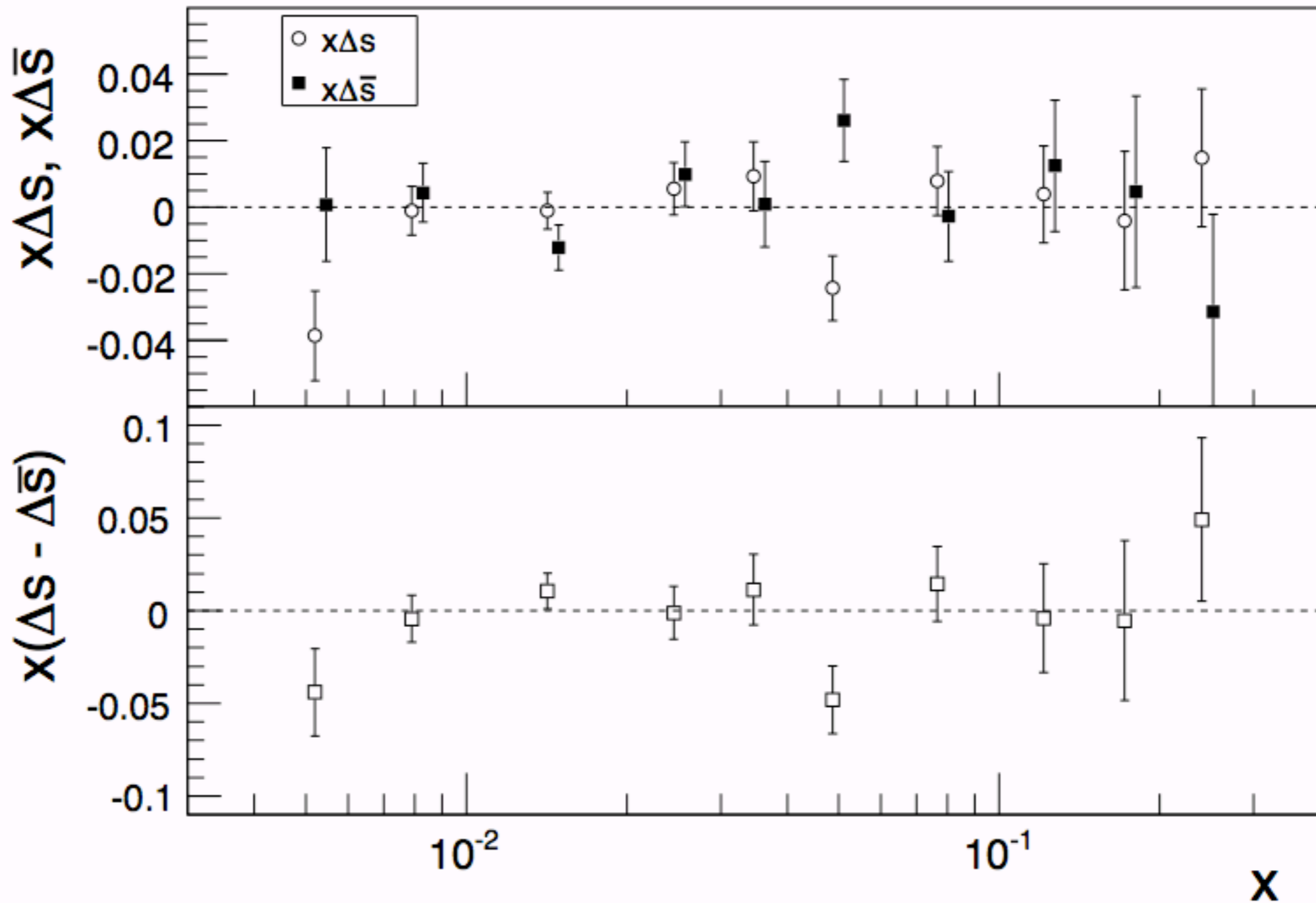
F, D constraints



DSSV: De Florian et al
 PRD 80, 034030 (2009).

Nucleon spin - COMPASS L.O. Δ_s

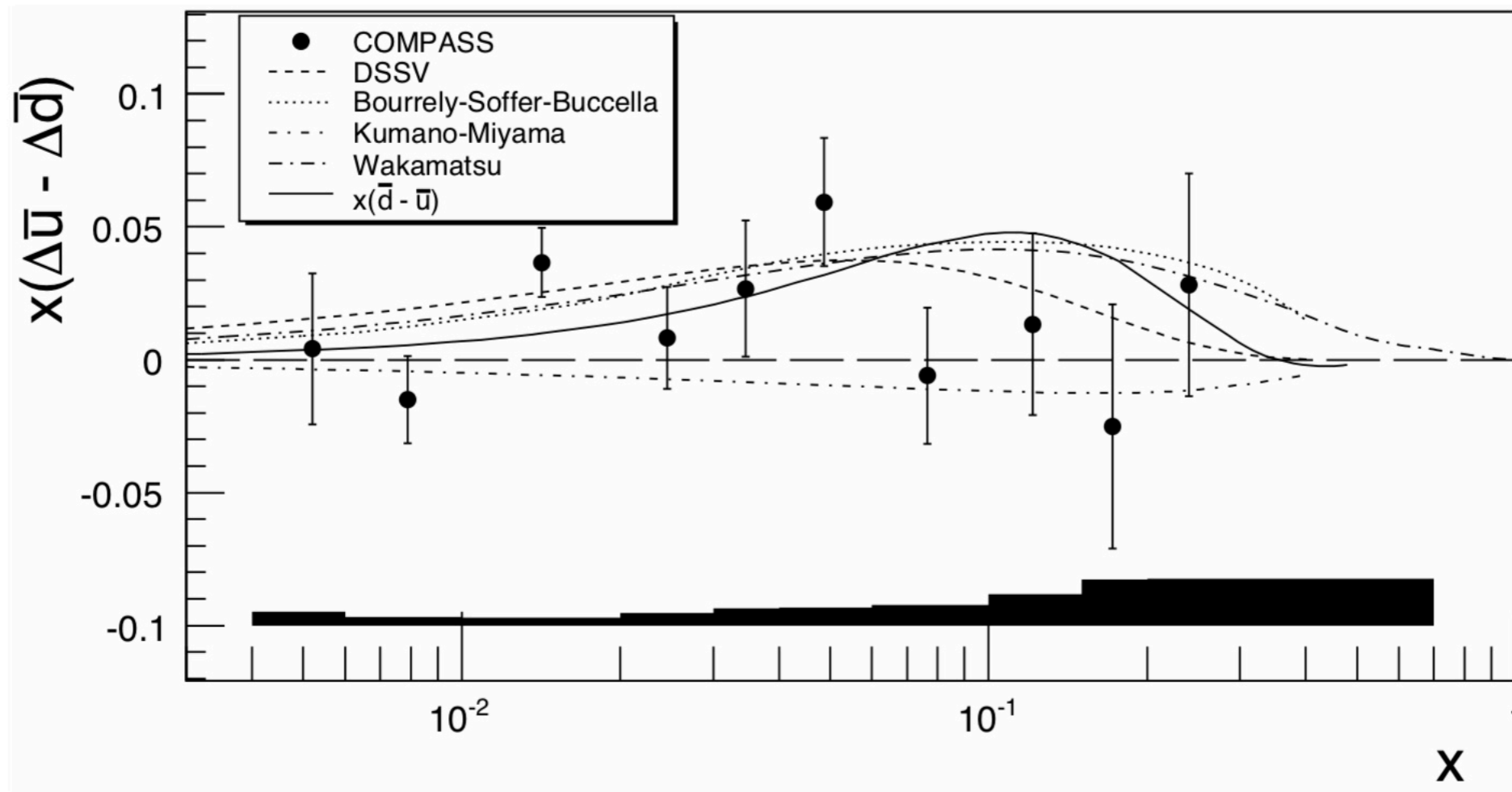
M.G. Alekseev, Phys.Lett. B693 (2010) 227-235



No evidence for negative net strange quark polarization,
not from $(K^+ + K^-)$ asymmetries of D either,
Renewed interest in FF, in-situ multiplicities, etc.

Nucleon spin - COMPASS L.O. $\Delta\bar{u} - \Delta\bar{d}$

M.G. Alekseev, Phys.Lett. B693 (2010) 227-235

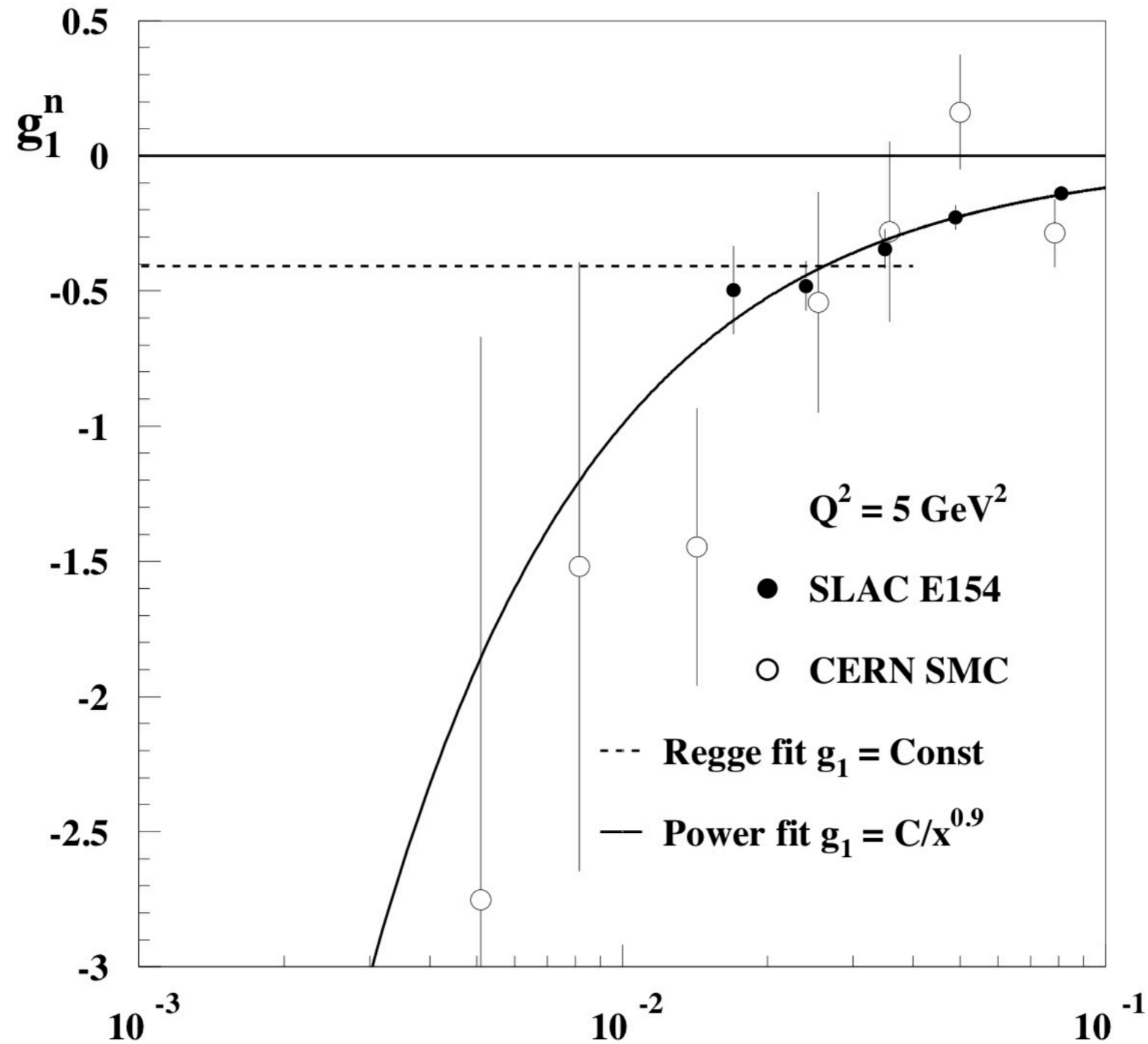


Well known flavor asymmetry in unpolarized sea,

COMPASS: $\Delta\bar{u} - \Delta\bar{d}$, is slightly positive, about 1.5 standard deviations from zero

Nucleon Spin - Scaling Violations, small-x

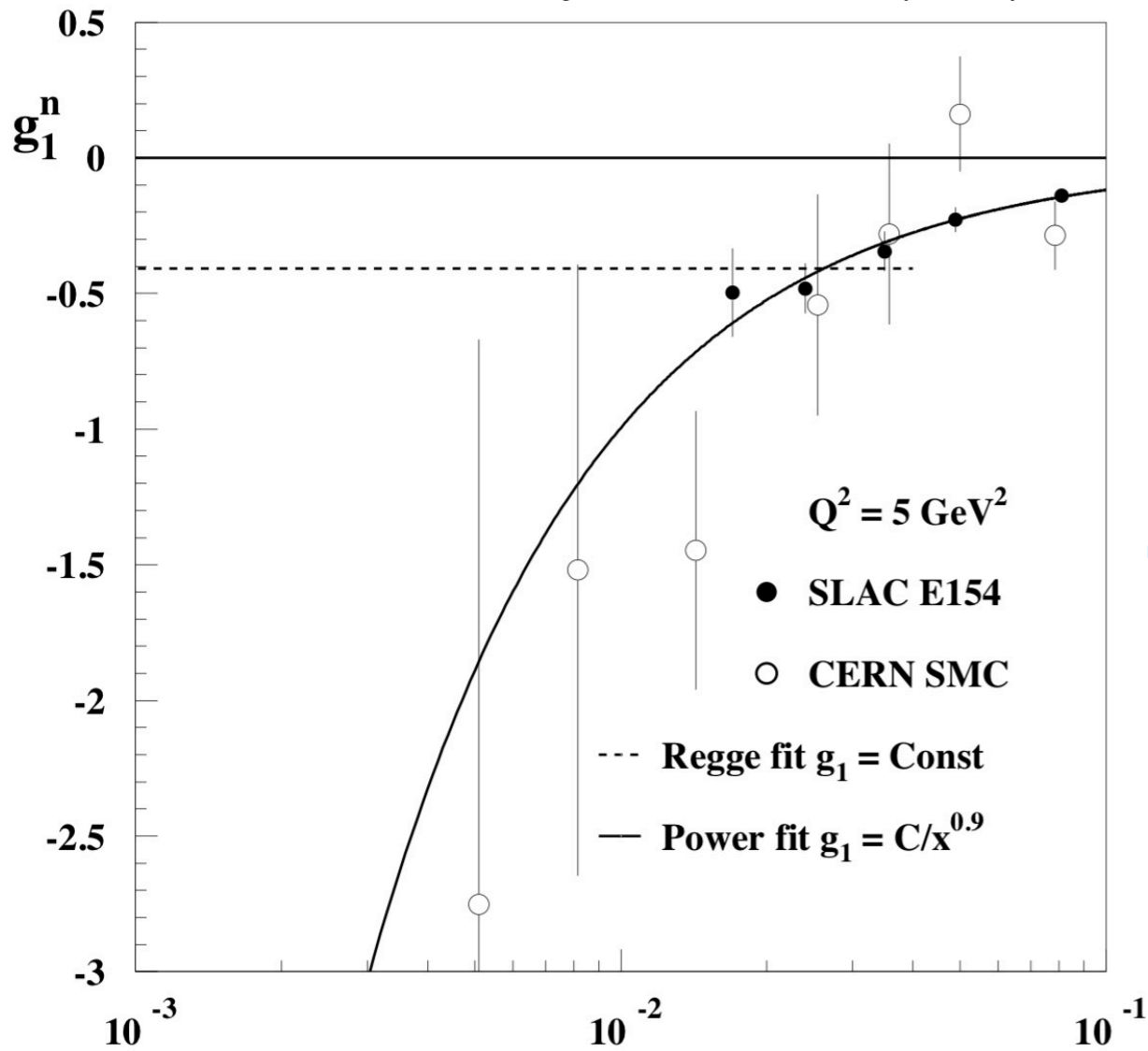
K. Abe, Phys. Rev. Lett. 79 (1997) 26



COMPASS has improved small-x uncertainties by a factor,
Nevertheless, the field has a ways to go...

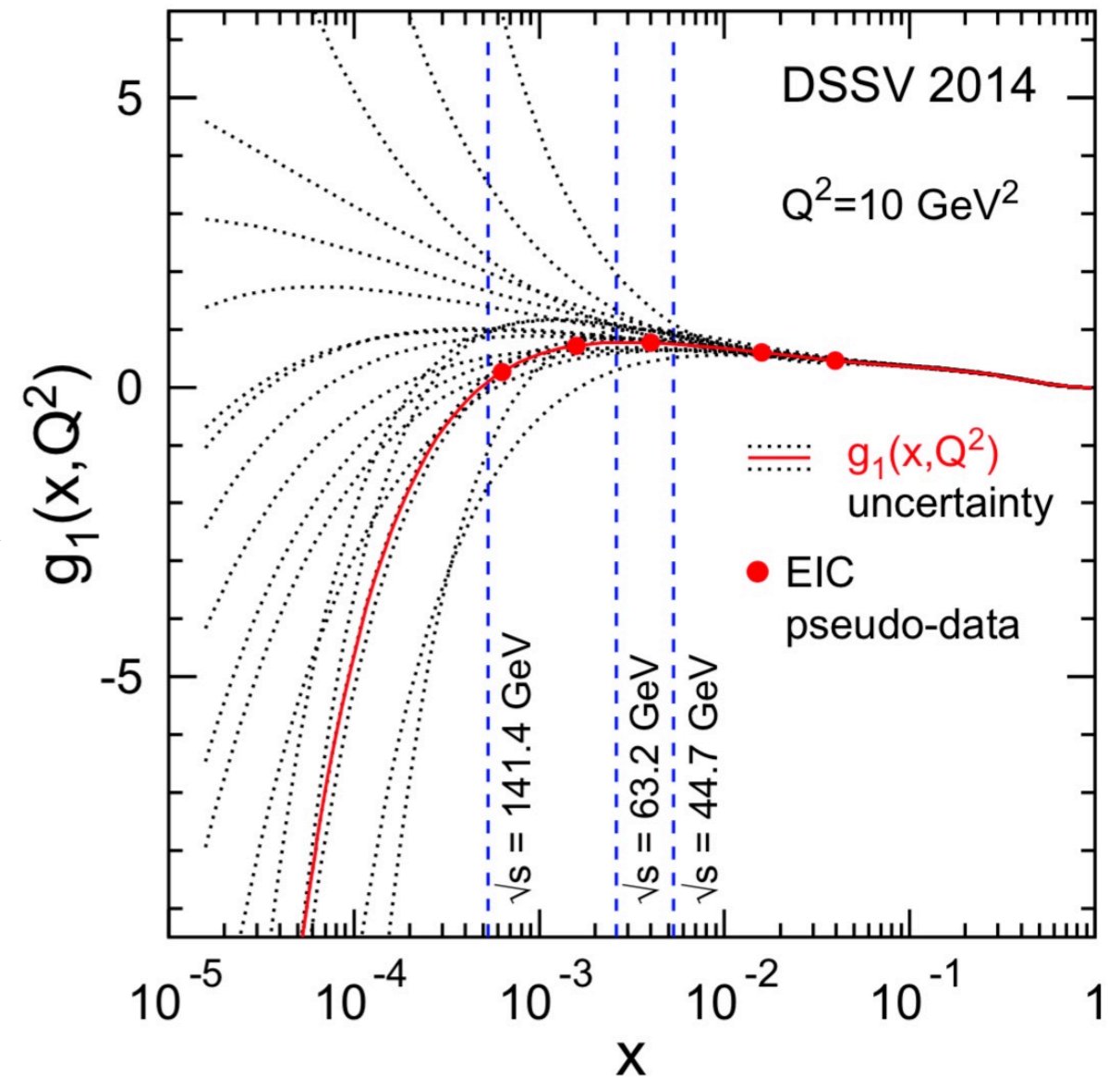
Nucleon Spin - Scaling Violations, small-x

K. Abe, Phys. Rev. Lett. 79 (1997) 26



neutron

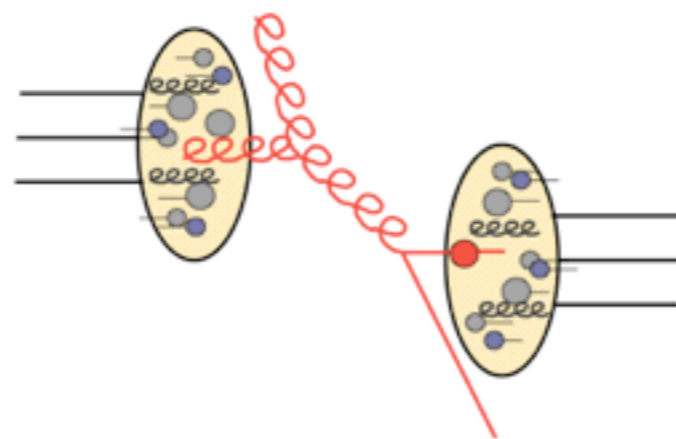
E.C. Aschenauer, ArXiv:1708.01527



proton

Clear role and impact for a future EIC

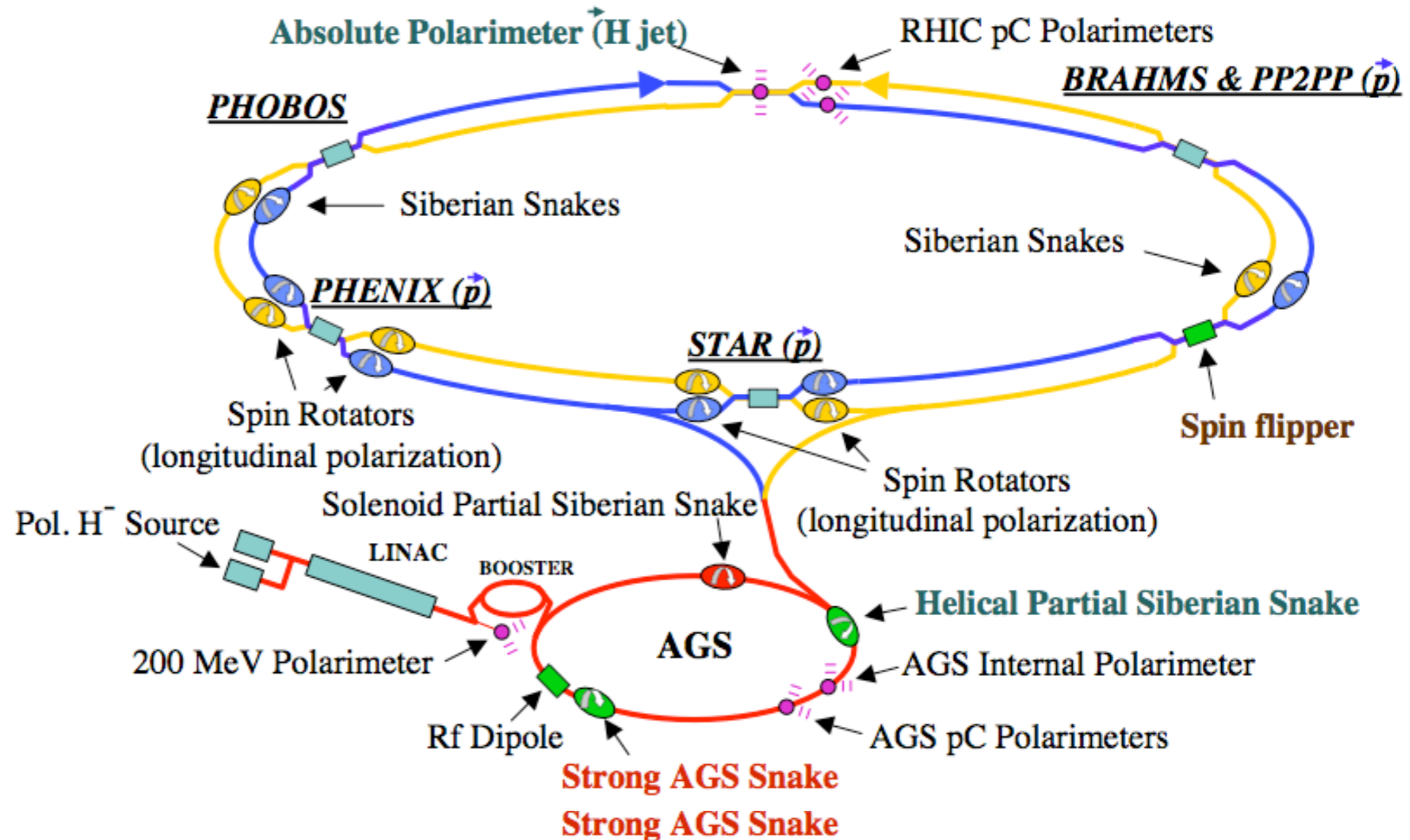
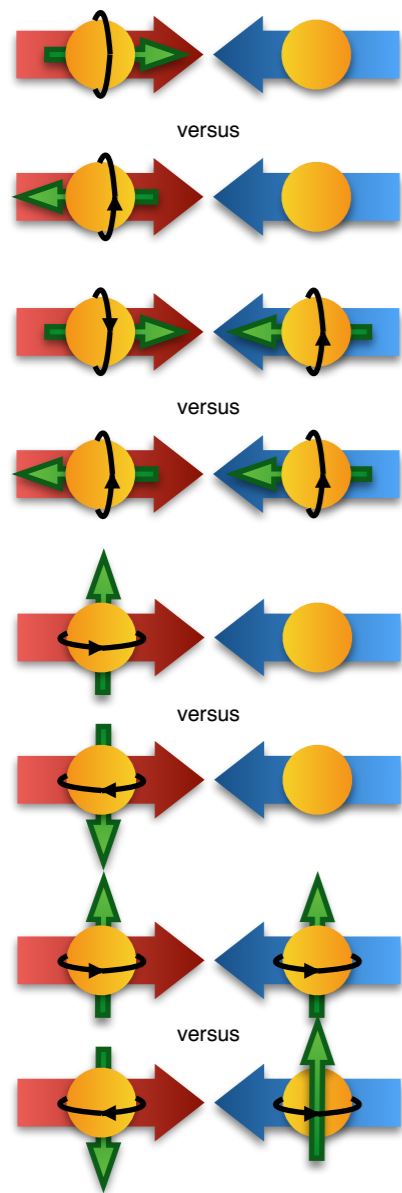
pp



RHIC - Polarized Proton-Proton Collider

Unique opportunities to study nucleon spin properties and spin in QCD,

$\sqrt{s} = 62, 200, \text{ and } 500 \text{ GeV}$



at hard (perturbative) scales with good systematic controls, e.g. from the $\sim 100\text{ns}$ succession of beam bunches with alternating beam spin configurations.

RHIC - Polarized Proton-Proton Collider

Unique opportunities to study nucleon spin properties and spin in QCD,

Longitudinal data

STAR

$\sqrt{s} = 200 \text{ GeV}$

2005

2006

2009

2015

35 pb⁻¹

50 pb⁻¹

$\sqrt{s} = 500 \text{ GeV}$

2009

2011

2012

2013

350 pb⁻¹

Transverse data

$\sqrt{s} = 200 \text{ GeV}$

2006

2008

2012

2015

38 pb⁻¹

50 pb⁻¹

$\sqrt{s} = 500 \text{ GeV}$

2011

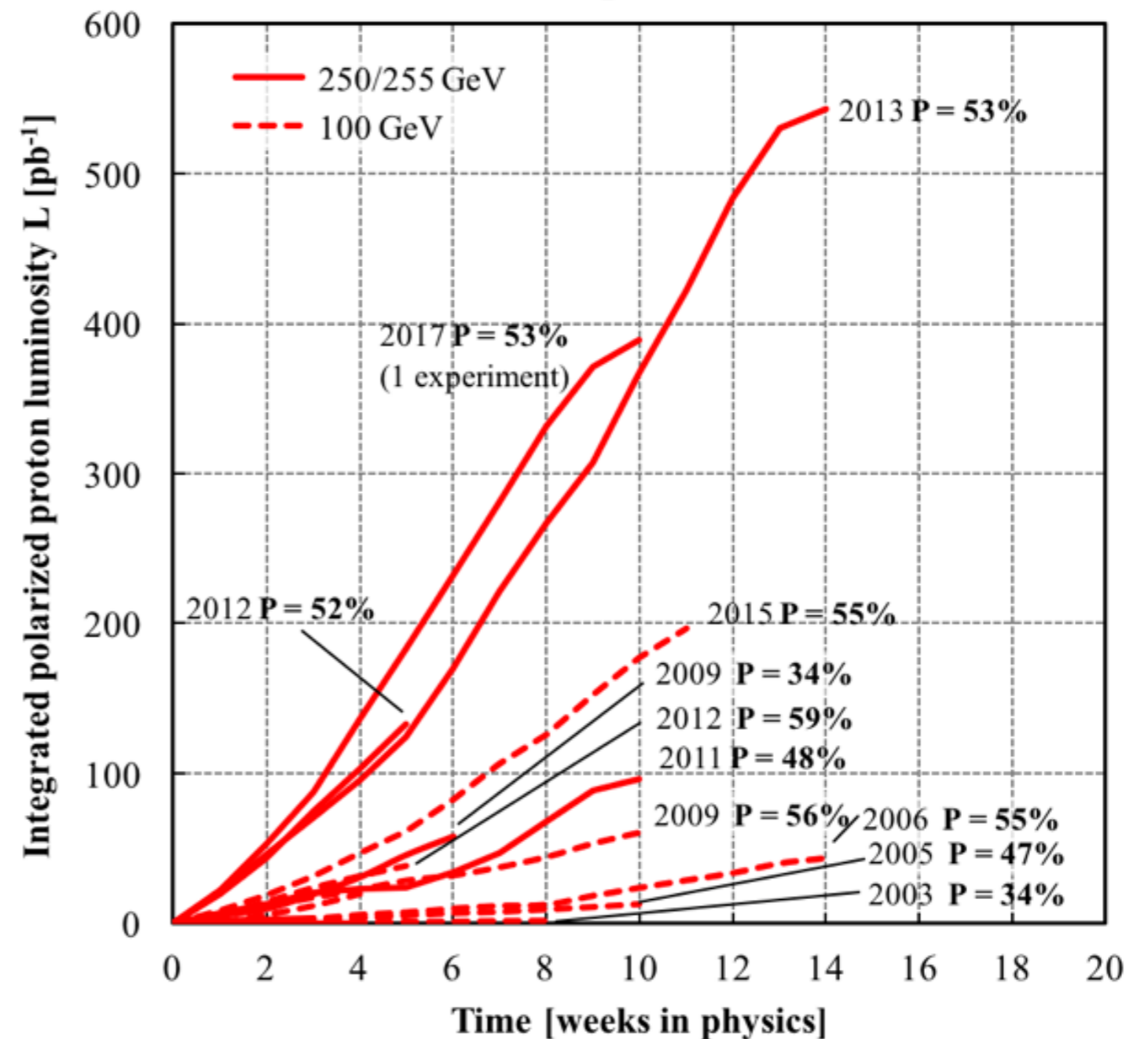
2017

25 pb⁻¹

350 pb⁻¹

50-60% polarization

Polarized protons



Gluon Polarization at RHIC

Measurement:
$$A_{LL} = \frac{\sigma^{\uparrow\uparrow} - \sigma^{\uparrow\downarrow}}{\sigma^{\uparrow\uparrow} + \sigma^{\uparrow\downarrow}} \stackrel{?}{=} \sum_{f=q,g} \frac{\Delta f_1}{f_1} \otimes \frac{\Delta f_2}{f_2} \otimes \hat{a}_{LL} \otimes (\text{fragmentation functions})$$

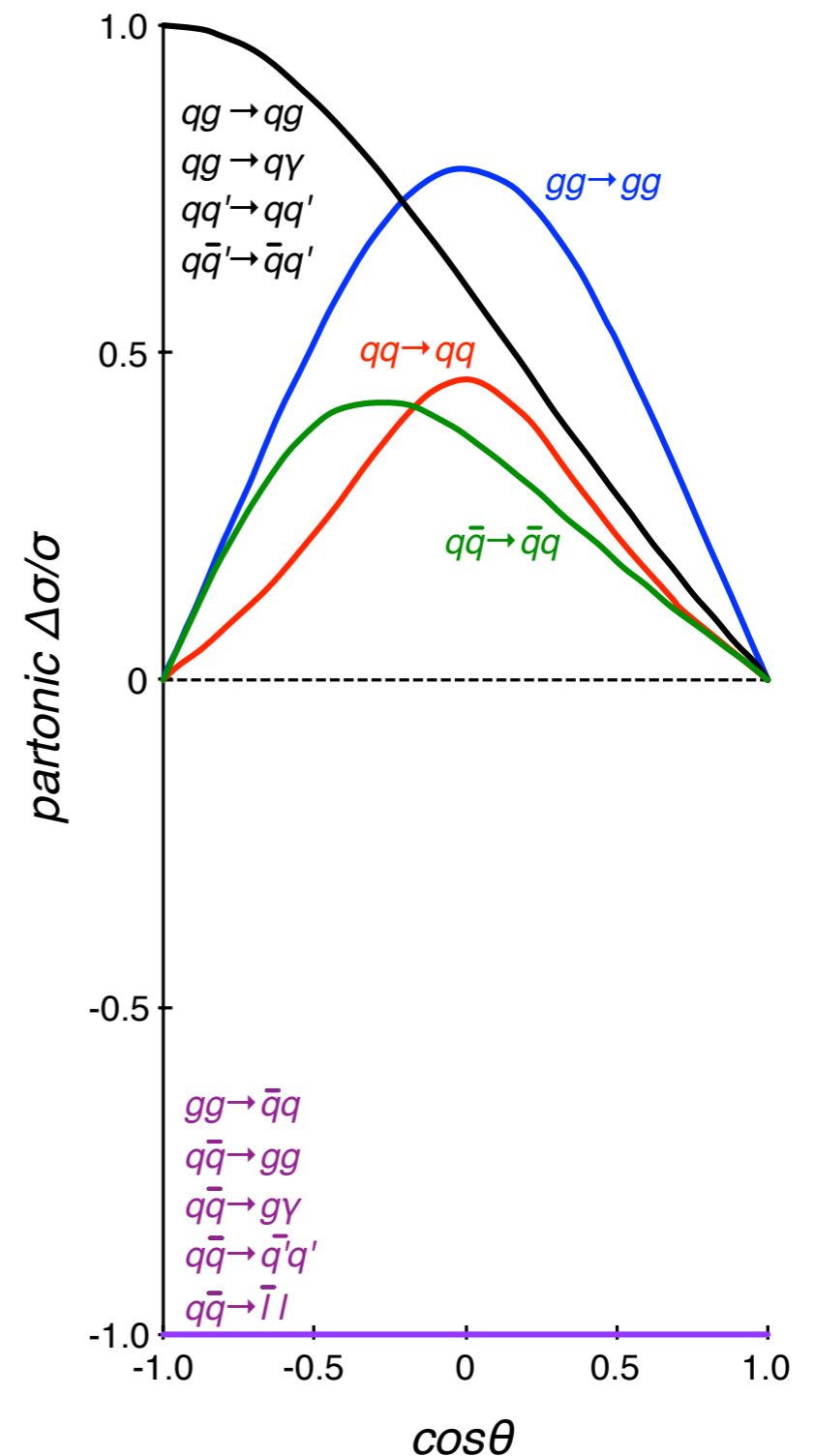
- Detect and reconstruct particle, jet,
- Extract beam-spin dependent yields,
- Measure relative luminosity, beam polarization
- Evaluate double beam-helicity asymmetry

Advantages:

- High yields of neutral pions, jets at RHIC,
- Relatively straightforward triggering,
- Understood reconstruction techniques,
- Sizable partonic asymmetries

Disadvantages:

- Contributions from several sub-processes,
- Wide x_g range sampled for each fixed p_T
- $x_g, x_q \sim p_T/\sqrt{s} \cdot \exp(-\eta)$



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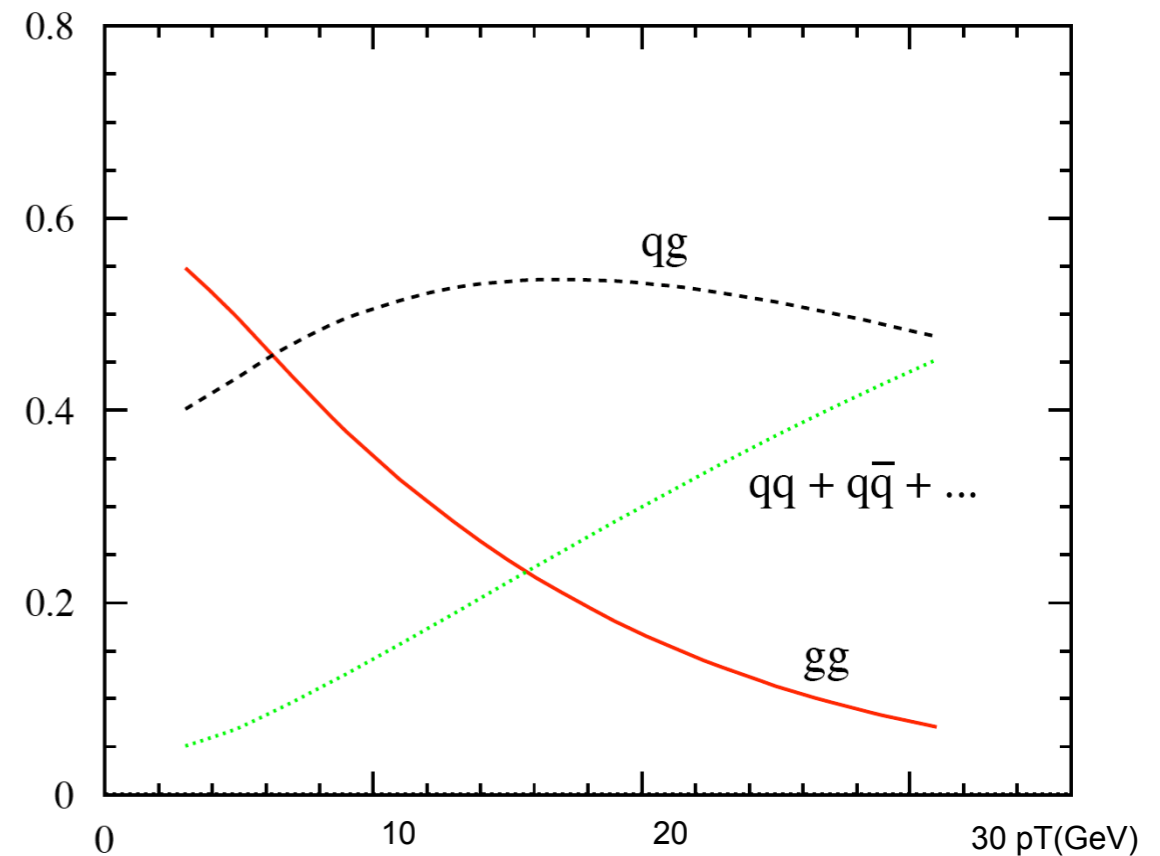
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gluon-gluon and *quark-gluon* scattering contributions dominate.

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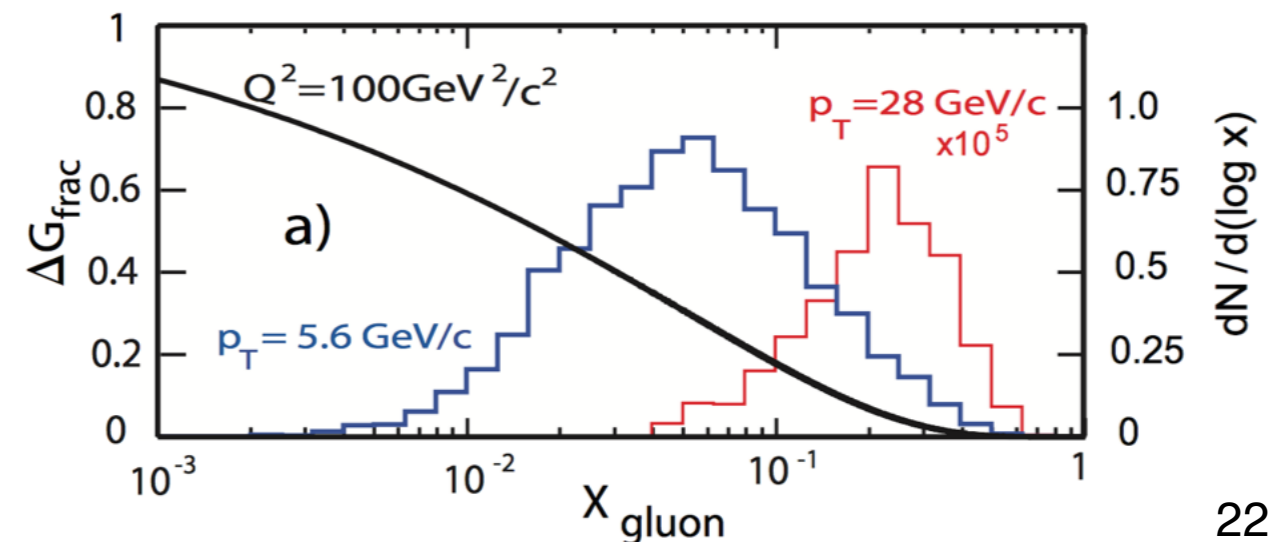
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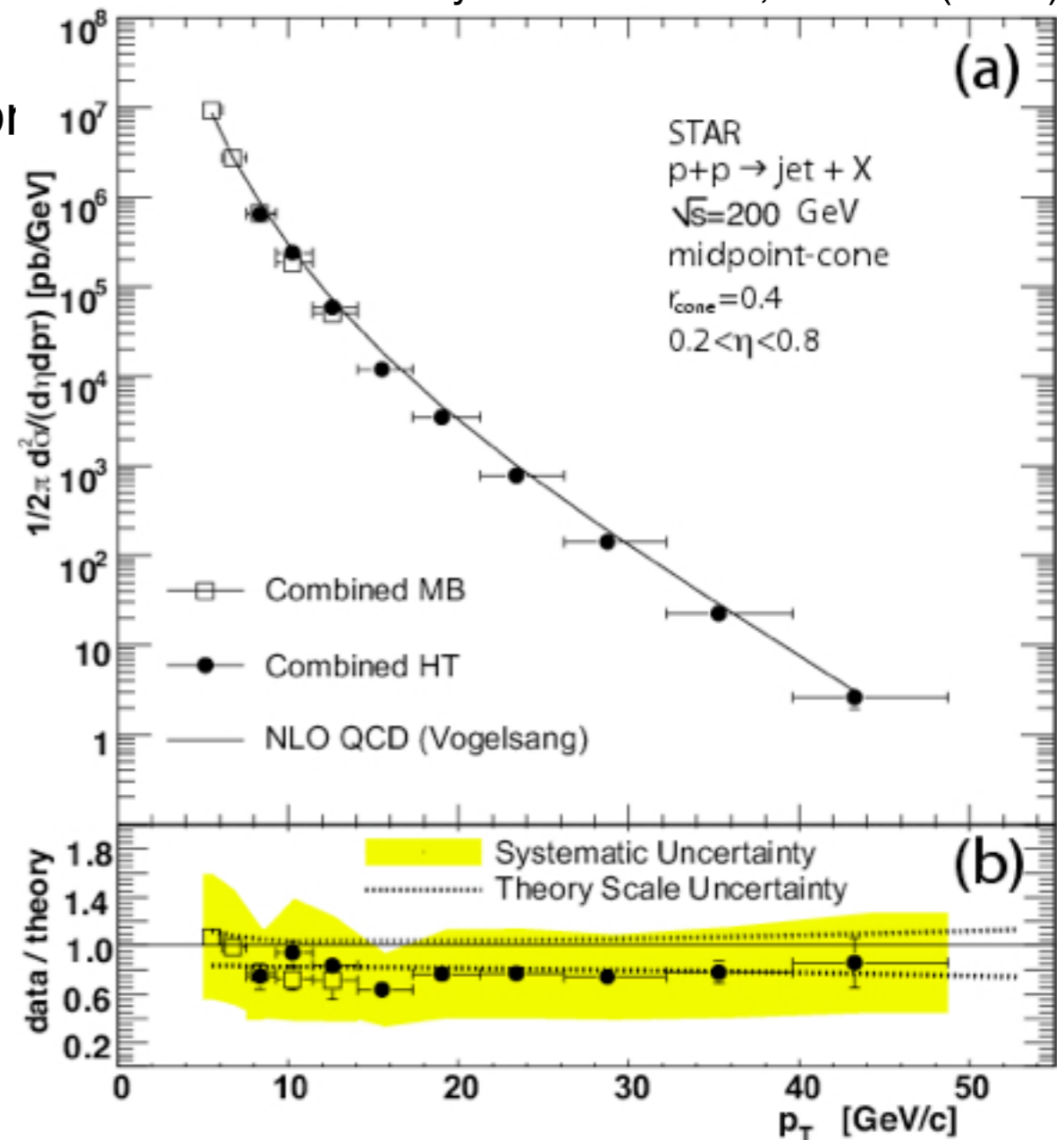
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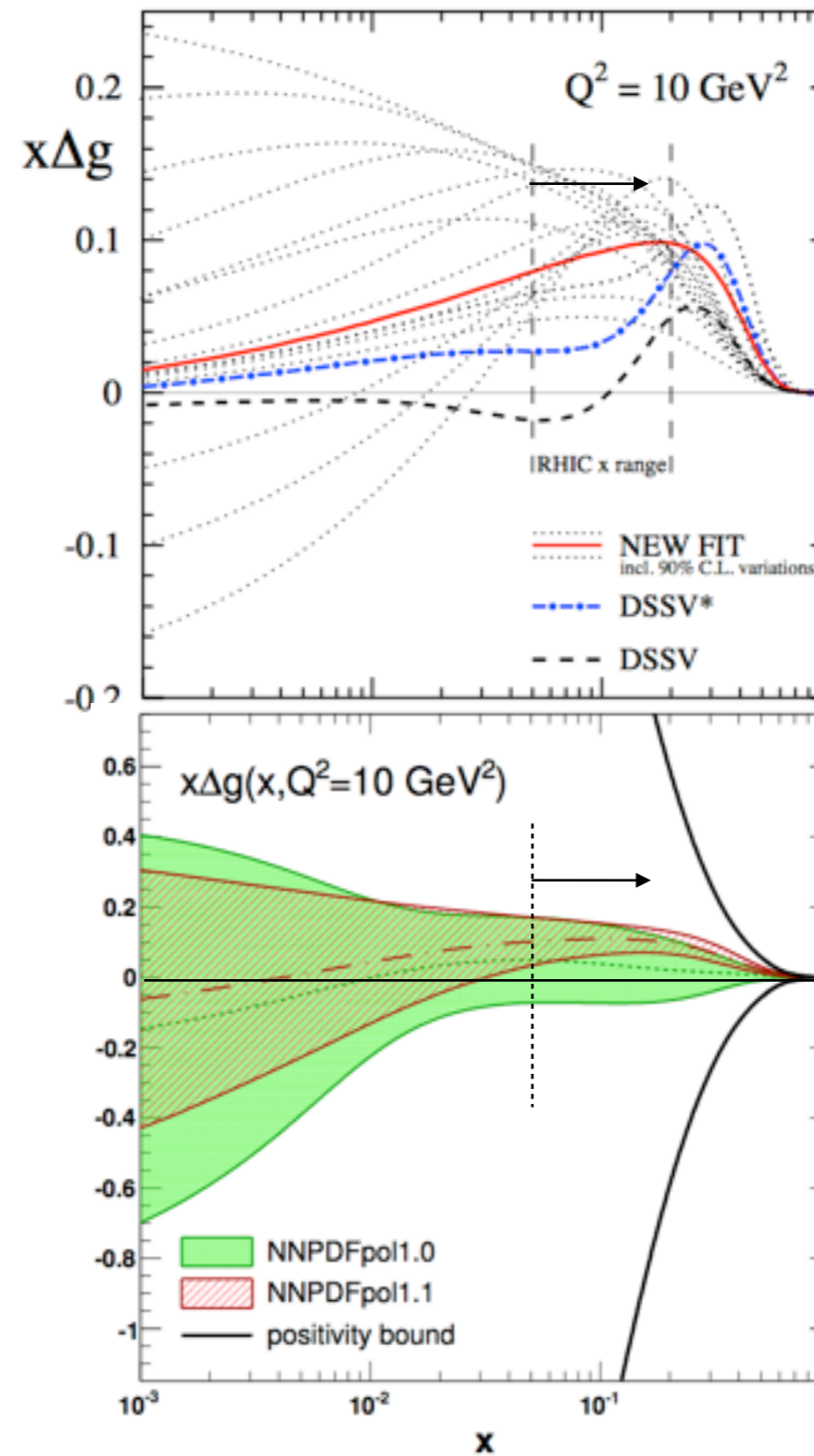
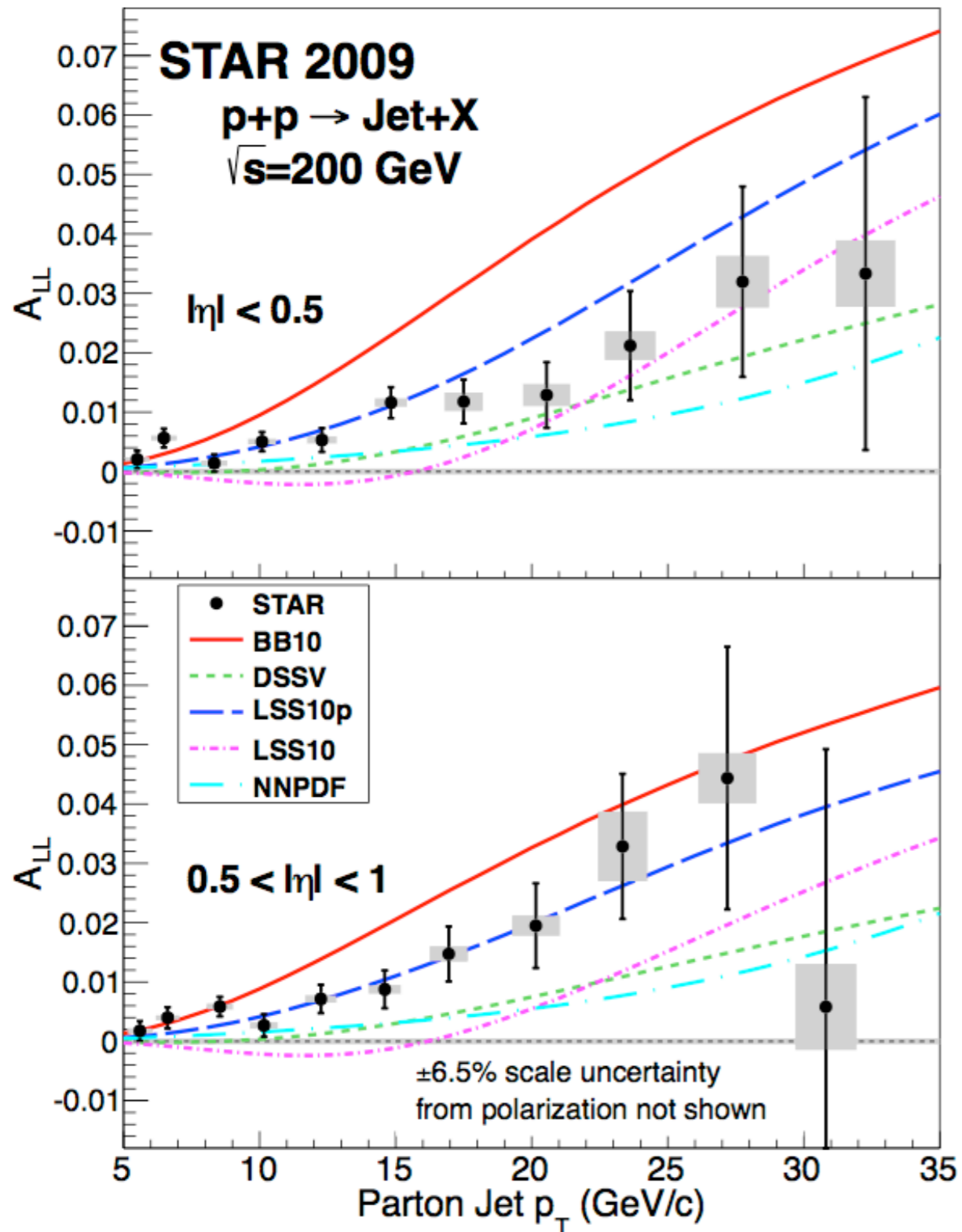
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Phys. Rev. Lett. 97, 252001 (2006)



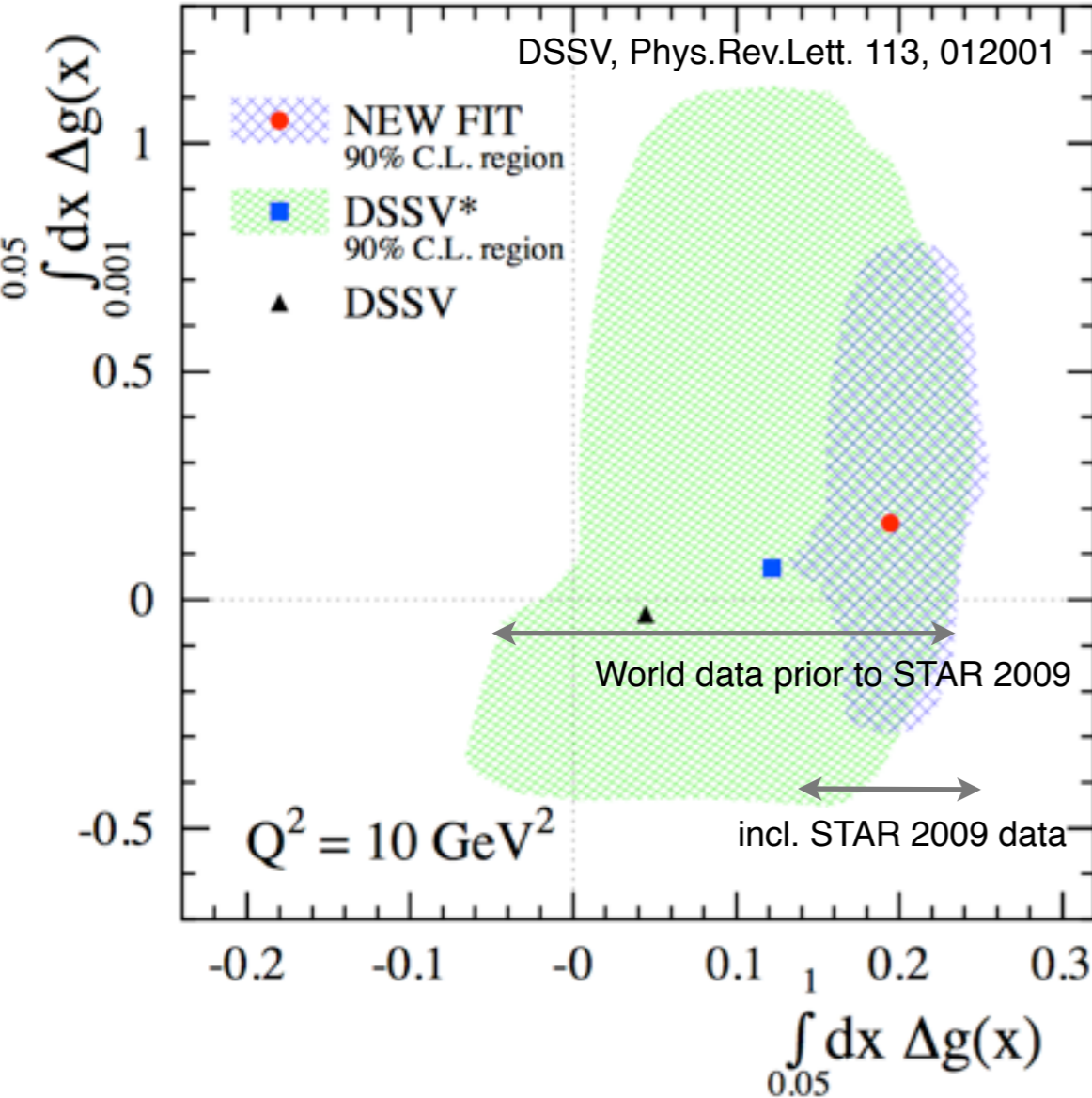
NLO pQCD agrees well with the “numerator”

Phys.Rev.Lett 115 (2015) 092002

 0.20 ± 0.07 DSSV, PRL 113,
012001(2014) 0.21 ± 0.10 NNPDF, Nucl. Phys. B
887, 276 (2014)Gluon polarization is positive in the region of the data; ~ 0.2

Gluon Polarization

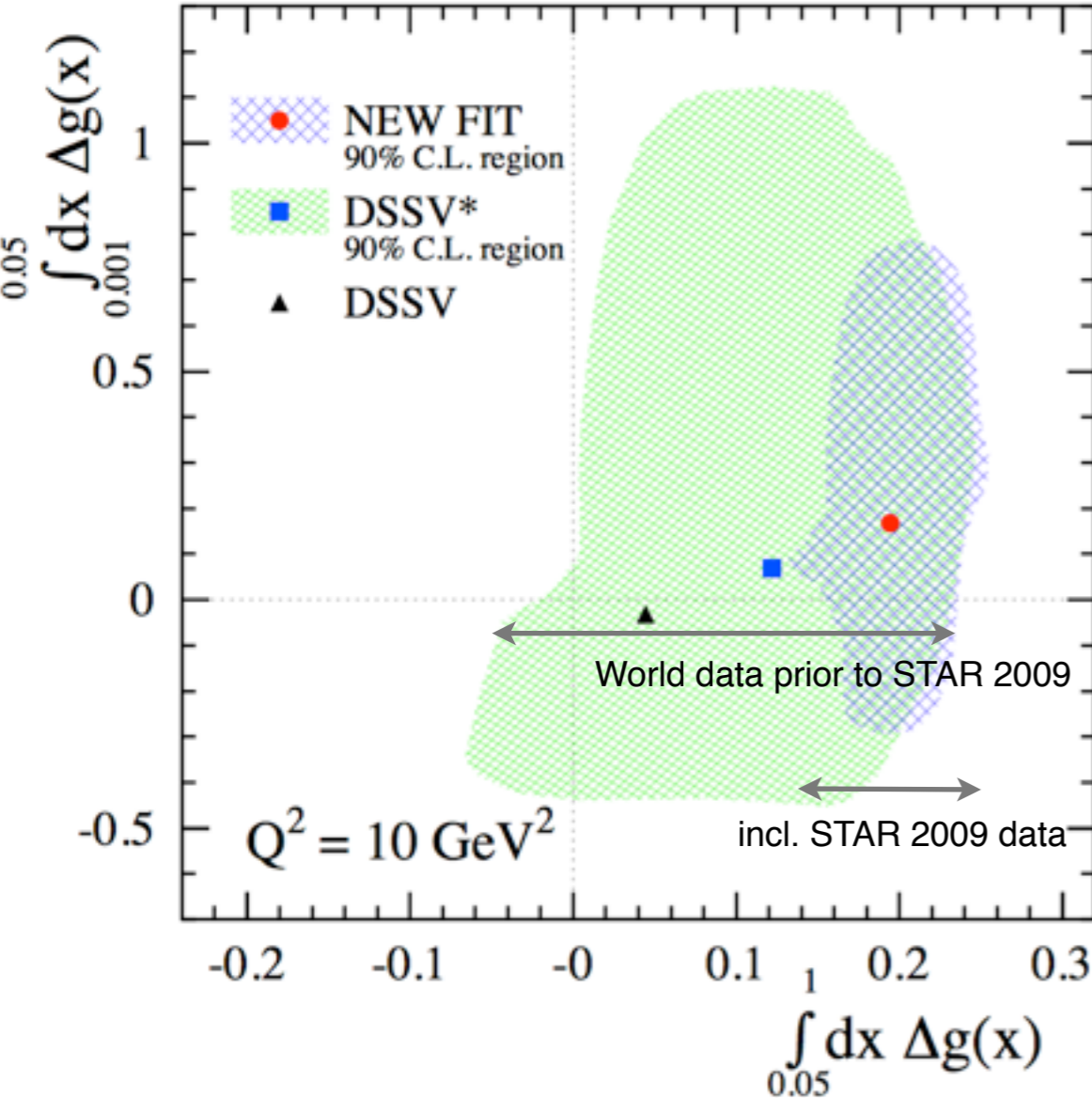
Evidence for positive gluon polarization



Easy to "hide" 1 h in the unmeasured region

Gluon Polarization

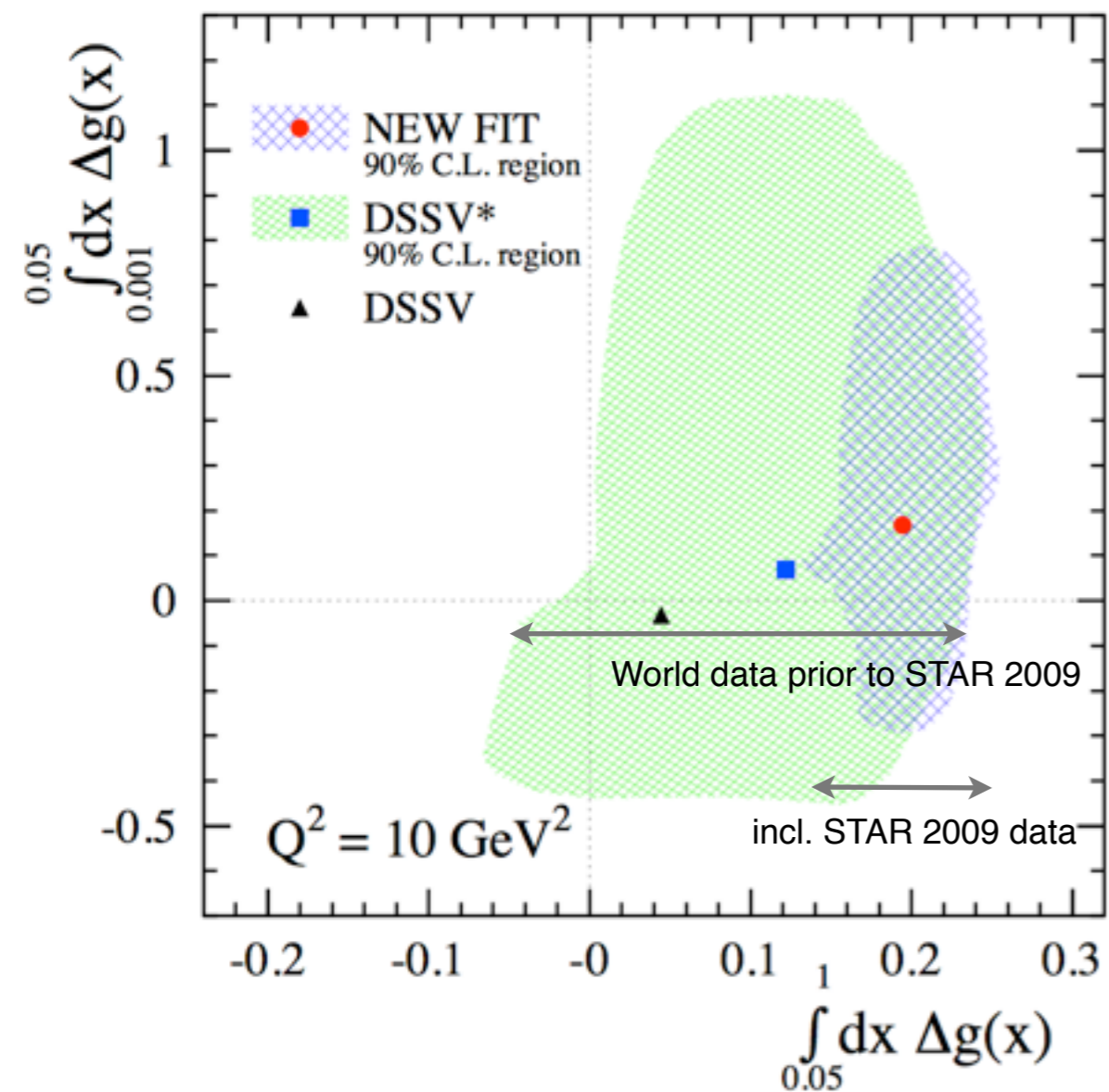
DSSV, Phys.Rev.Lett. 113, 012001



Extend sensitivity to *smaller* x_g
 $\sqrt{s} = 500 \text{ GeV}$ data, $x_g \sim 1/\sqrt{s}$,
 forward rapidity, $x_g \sim \exp(-\eta)$,

Gluon Polarization

DSSV, Phys.Rev.Lett. 113, 012001



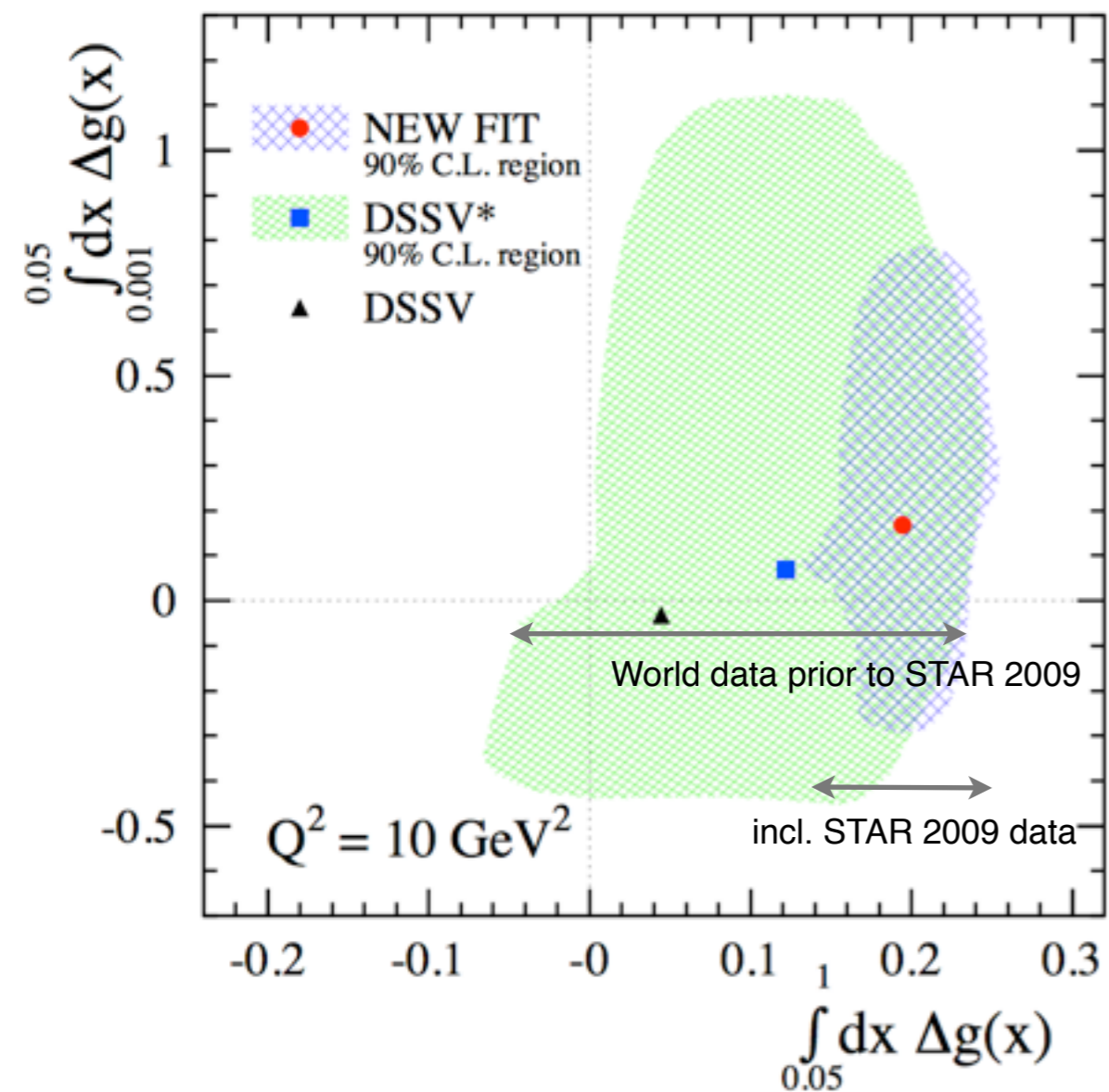
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Further *precision* from jet and neutral pion probes, and
 from *complementary* probes

Gain insight in x_g dependence.

Gluon Polarization

DSSV, Phys.Rev.Lett. 113, 012001



→

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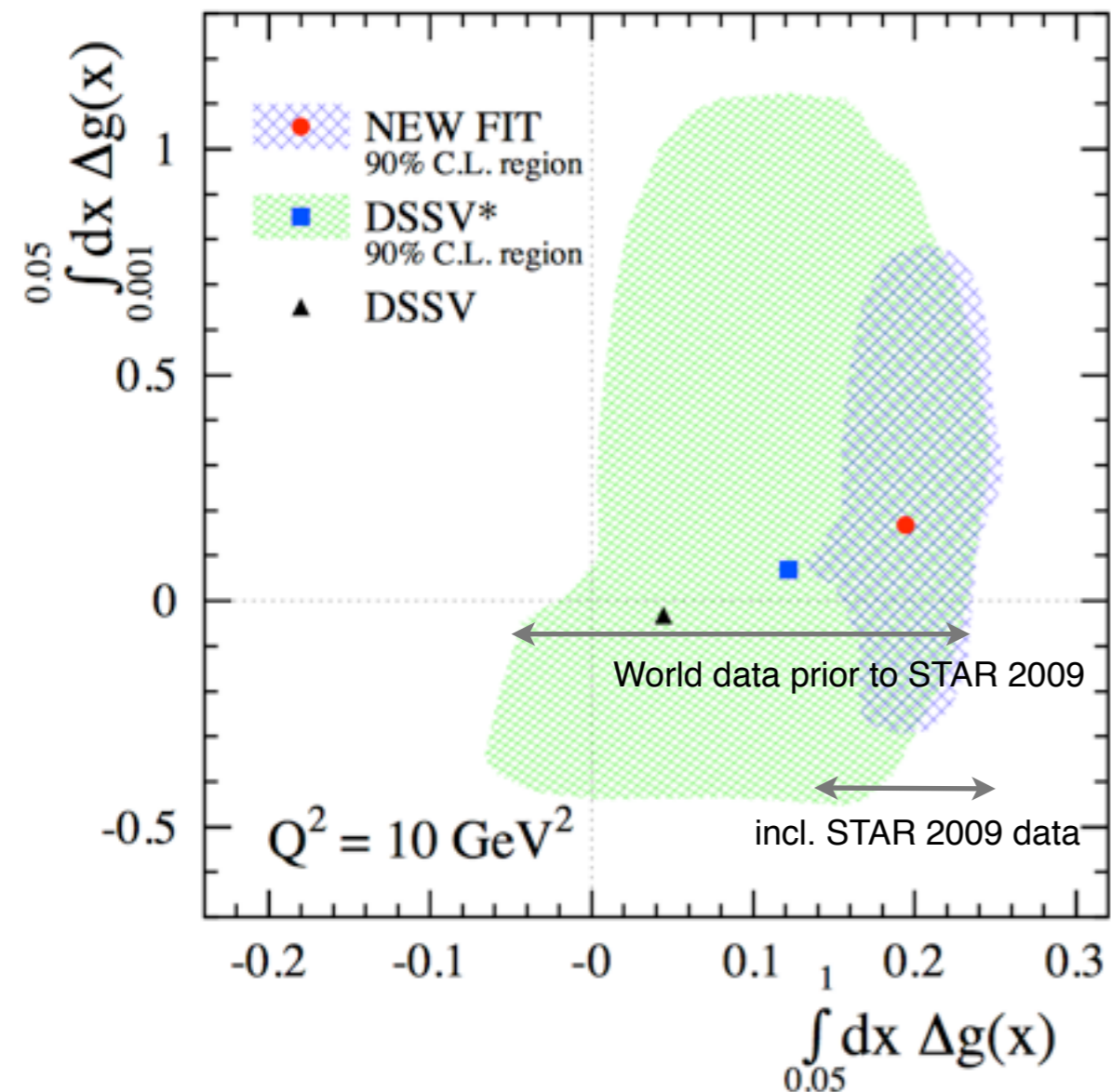
↓

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DSSV, Phys.Rev.Lett. 113, 012001



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 forward rapidity, $x_g \sim \exp(-\eta)$,

**RHIC continues to make progress and impact,
 see also Elke's talk yesterday,**

Clear role for a future EIC

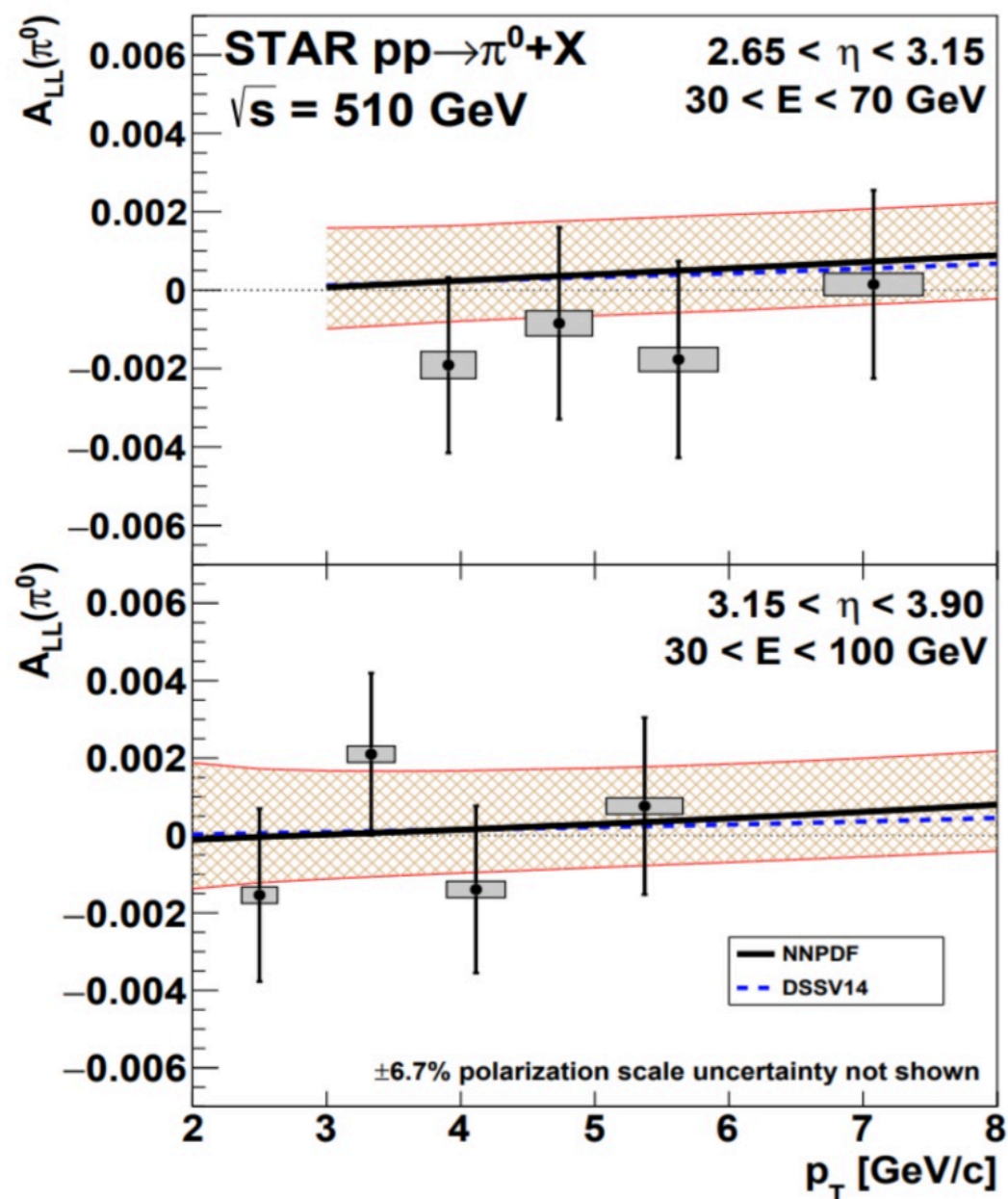
Further *precision* from jet and neutral pion probes, and
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Gain insight in x_g dependence.

Gluon Polarization

An early glimpse in the forward acceptance region:

arXiv: 1805.09745, PRD 98, 032013 (2018)



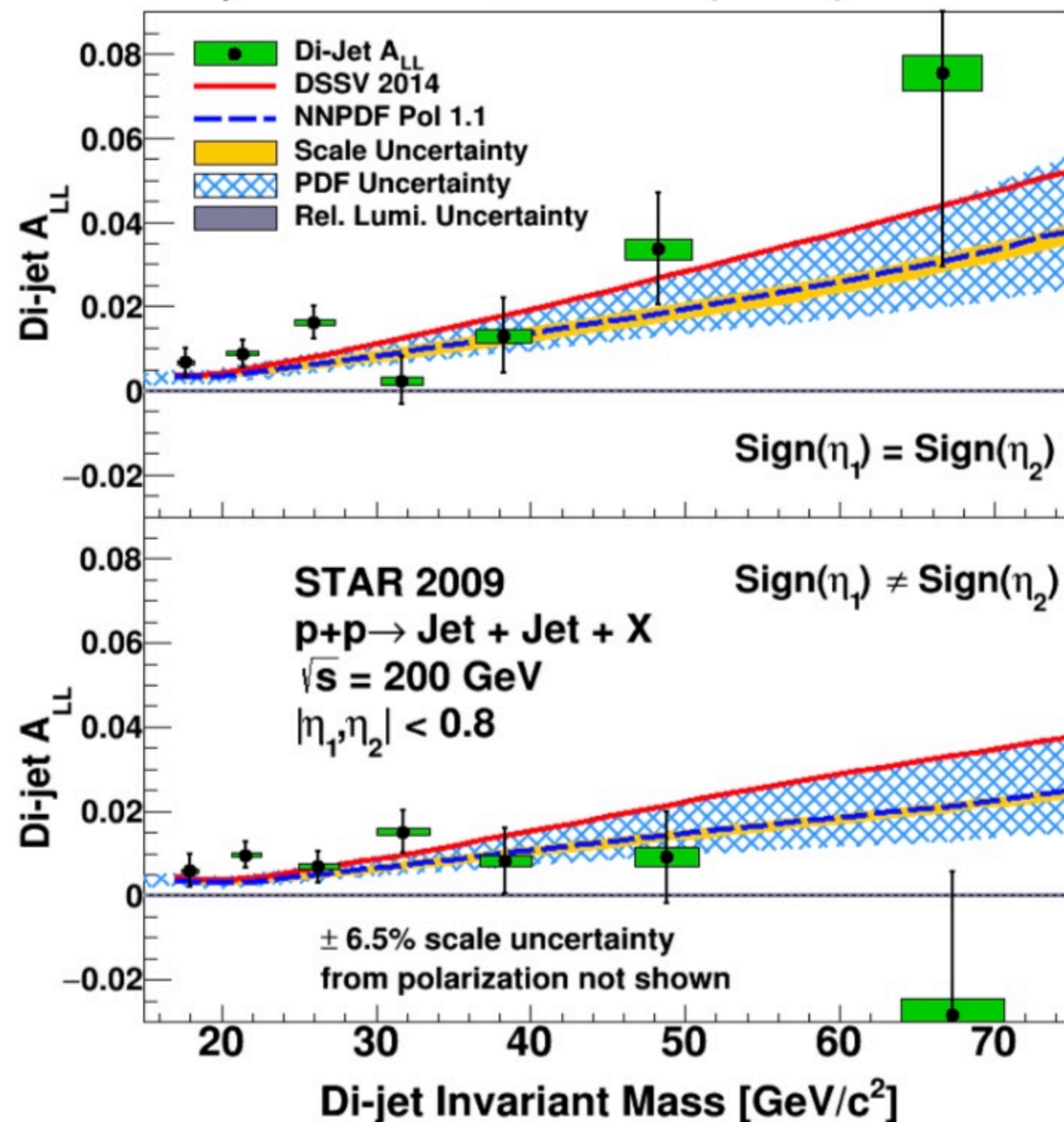
- Results are given for transverse momenta in the range $2 < p_T < 10 \text{ GeV/c}$ within two regions of pseudorapidity that span $2.65 < \eta < 3.9$
- These results are sensitive to the polarized gluon parton distribution function, $\Delta g(x)$, down to the region of parton momentum fraction $x \sim 0.001$
- These results will provide the first direct experimental constraints in $x \ll 0.01$

Correlation measurements will access larger (average) partonic asymmetries.

Gluon Polarization

Mid-central *di-jet* asymmetries:

Phys. Rev. D 95, 071103 (2017)

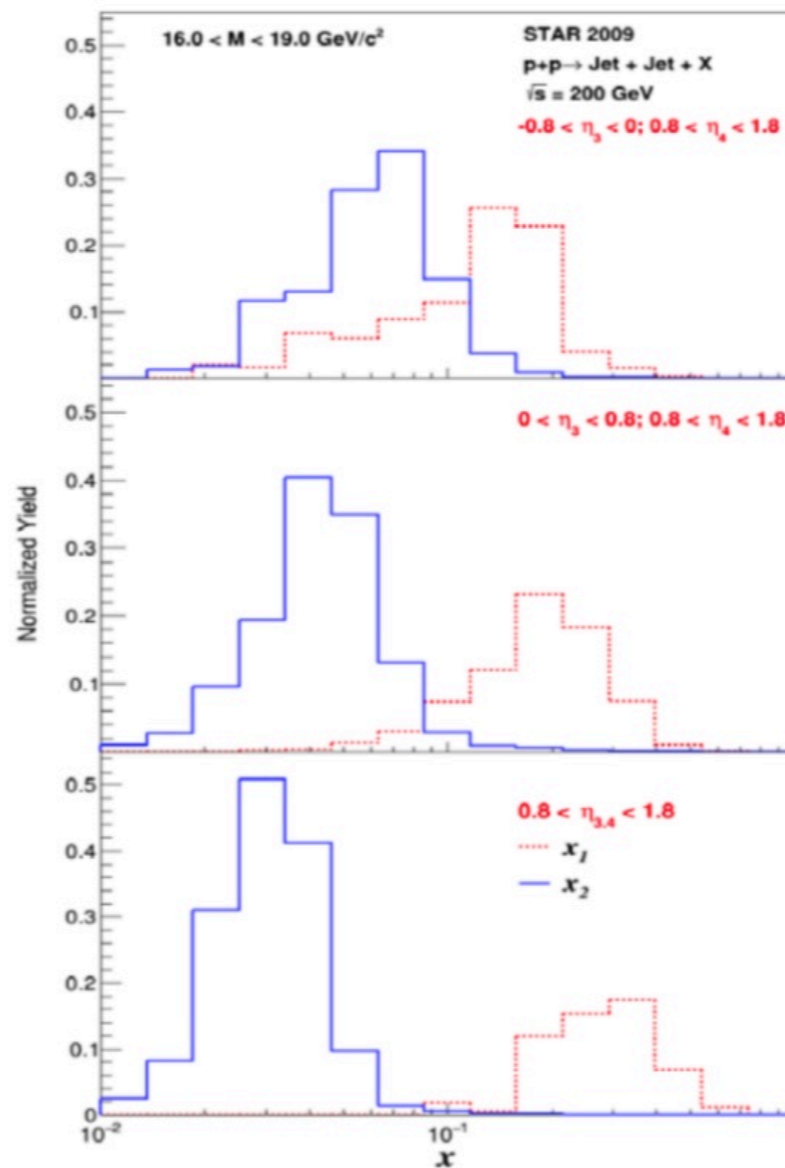
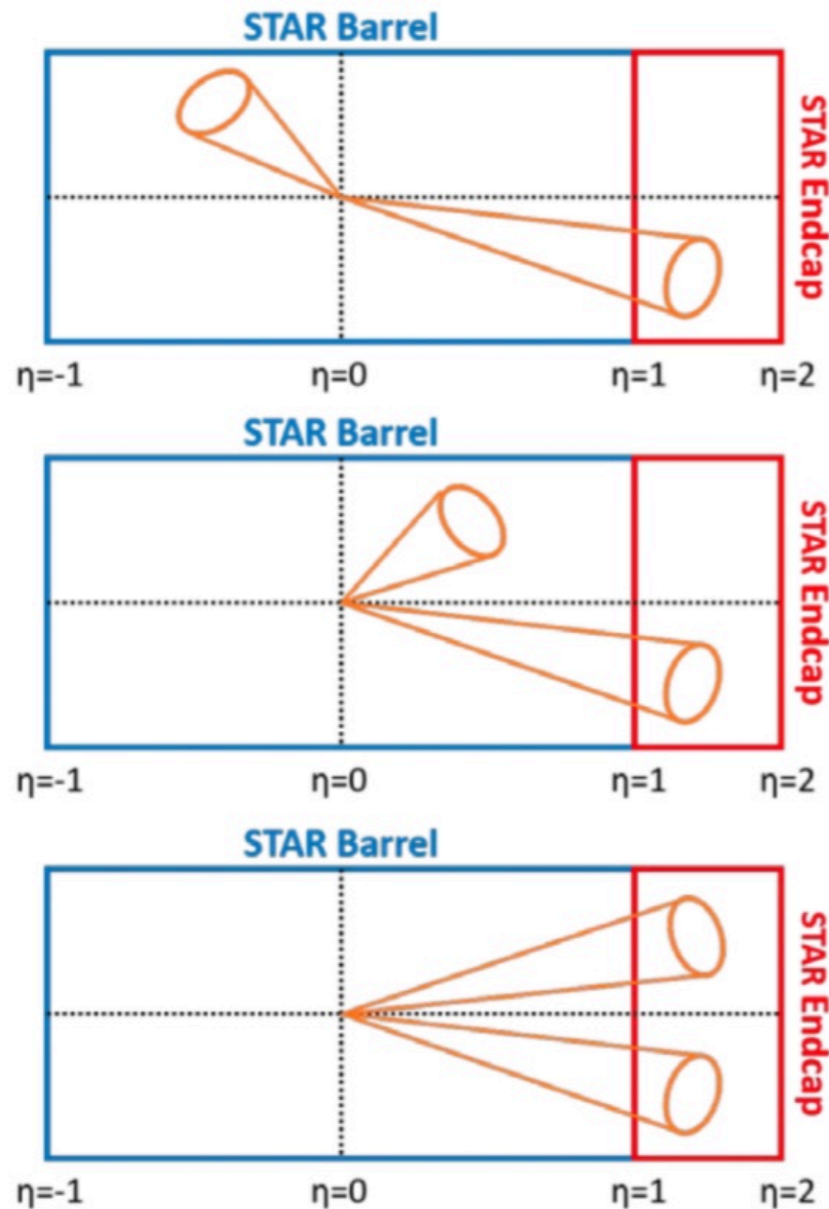


Towards sensitivity to Bjorken- x .

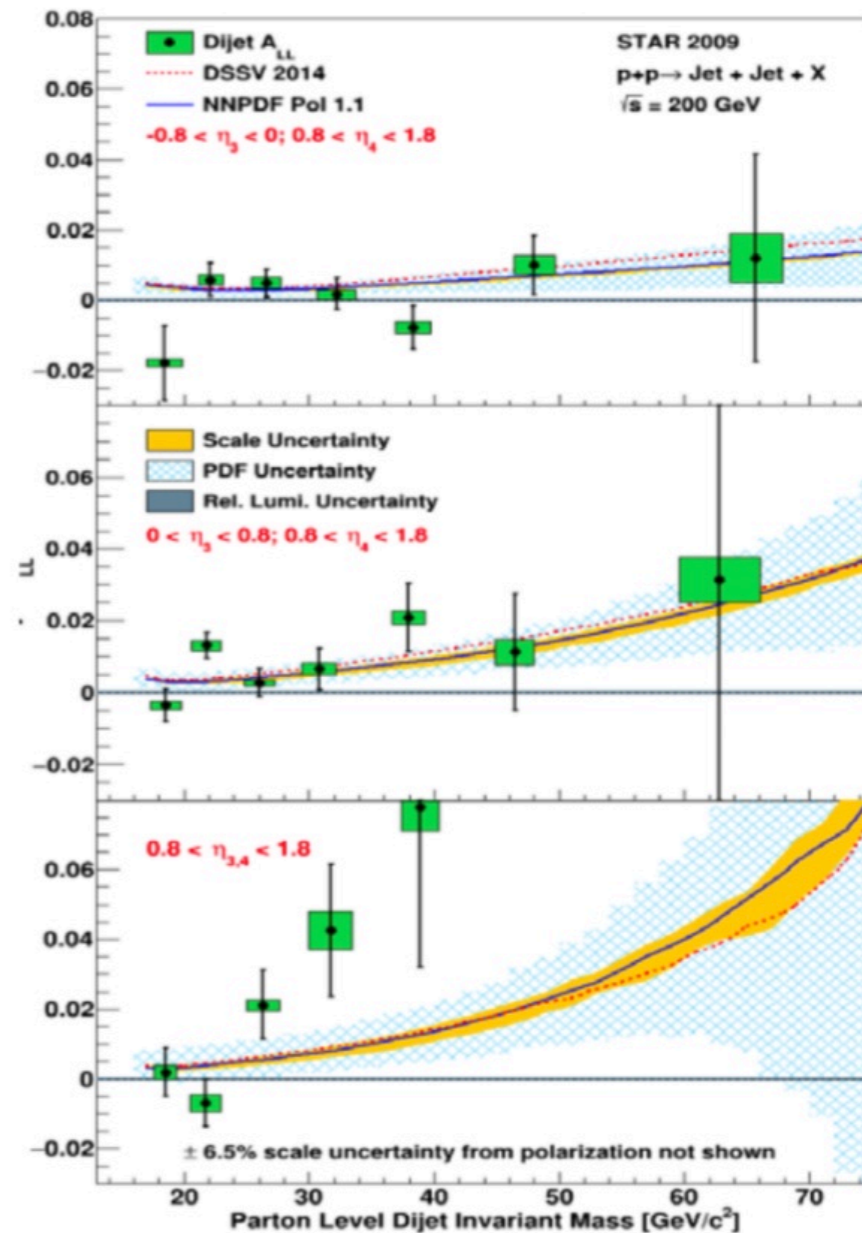
Preliminary results at 500 GeV have come out as well, paper in preparation

Gluon Polarization

di-jet asymmetries in a more forward region:



PRD 98 (2018) 032011



Impact clearly exists; quantifying it will require renewed global analysis (and/or reweighting)

Quark Polarization at RHIC

$\sqrt{s} = 500$ GeV above W production threshold,

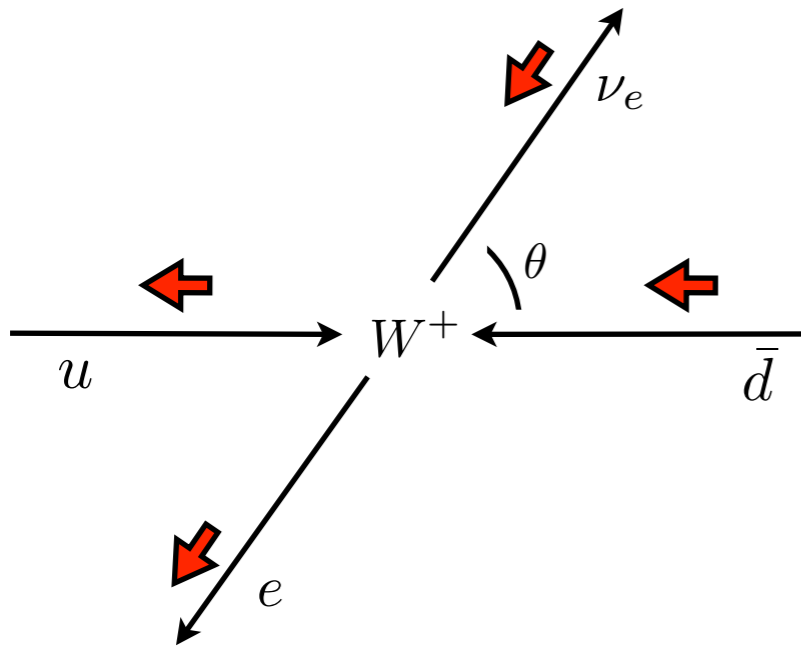
Experiment Signature:

large p_T lepton, missing E_T

Experiment Challenges:

charge-ID at large $|\eta|$
electron/hadron discrimination
luminosity hungry

Free of fragmentation (!)



$$\Delta\sigma^{\text{Born}}(\vec{p}p \rightarrow W^+ \rightarrow e^+ \nu_e) \propto -\Delta u(x_a)\bar{d}(x_b)(1+\cos\theta)^2 + \Delta\bar{d}(x_a)u(x_b)(1-\cos\theta)^2$$

Spin Measurements:

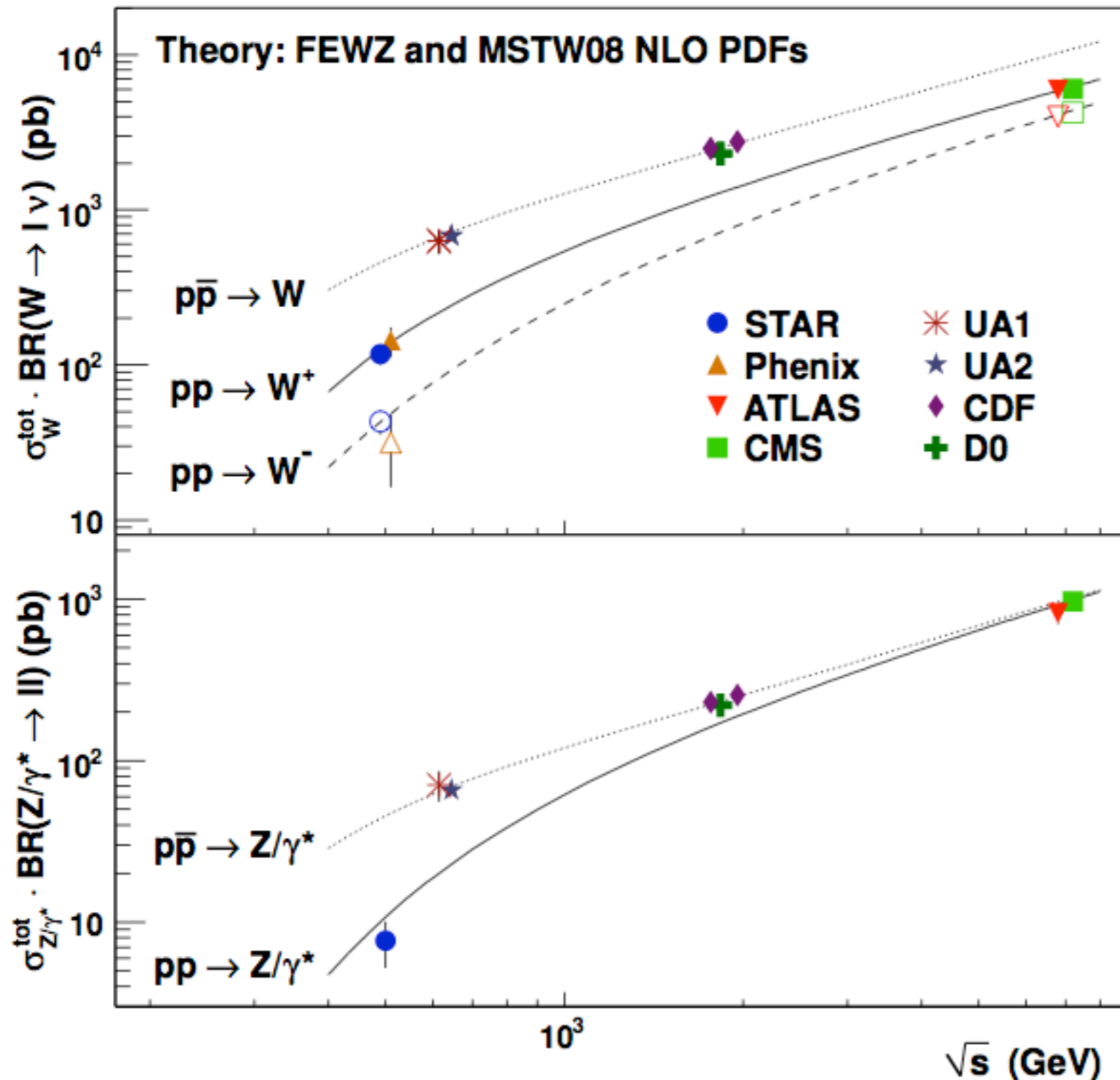
$$A_L(W^+) = \frac{-\Delta u(x_a)\bar{d}(x_b) + \Delta\bar{d}(x_a)u(x_b)}{u(x_a)\bar{d}(x_b) + \bar{d}(x_a)u(x_b)} = \begin{cases} -\frac{\Delta u(x_a)}{u(x_a)}, & x_a \rightarrow 1 \\ \frac{\Delta\bar{d}(x_a)}{\bar{d}(x_a)}, & x_b \rightarrow 1 \end{cases}$$

LO expressions to illustrate overall behavior,

$$A_L(W^-) = \begin{cases} -\frac{\Delta d(x_a)}{d(x_a)}, & x_a \rightarrow 1 \\ \frac{\Delta\bar{u}(x_a)}{\bar{u}(x_a)}, & x_b \rightarrow 1 \end{cases}$$

NLO known and used in extracting pPDFs.

W and Z Production Cross Sections



PHENIX: first W^+ and W^- production cross sections in proton-proton collisions, Phys.Rev.Lett. **106** (2011) 062001,

STAR: Initial NC cross section at RHIC, confirmation of PHENIX CC cross section measurements, Phys. Rev. **D85** (2012).

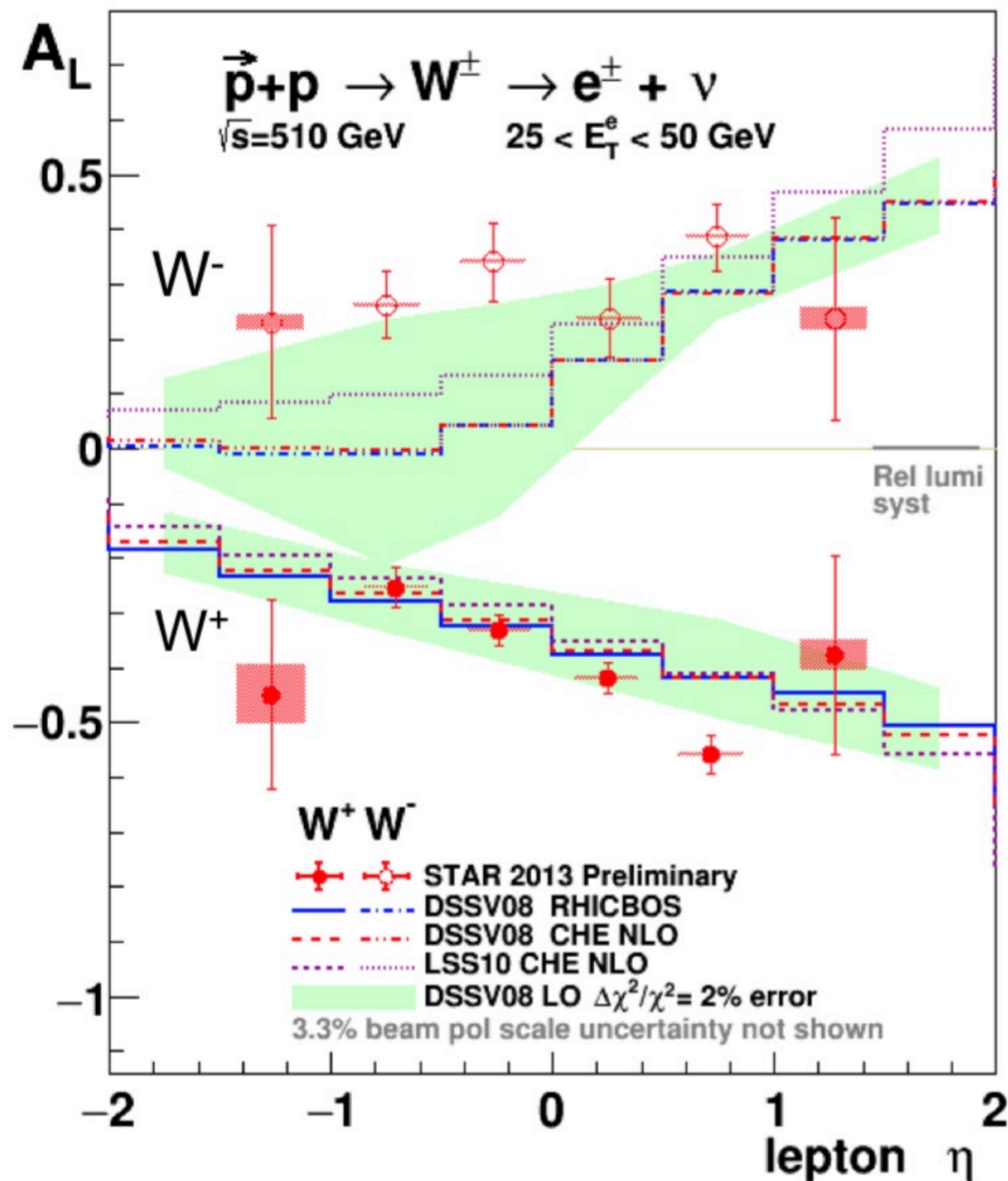
Data are well-described by NLO pQCD theory (FEWZ + MSTW08),

Necessary condition to interpret asymmetry measurements,

Future ratio measurements may provide insights in unpolarized light quark distributions

The “numerator” is again well-described; differential c.s. ratios complementary to E866, ...

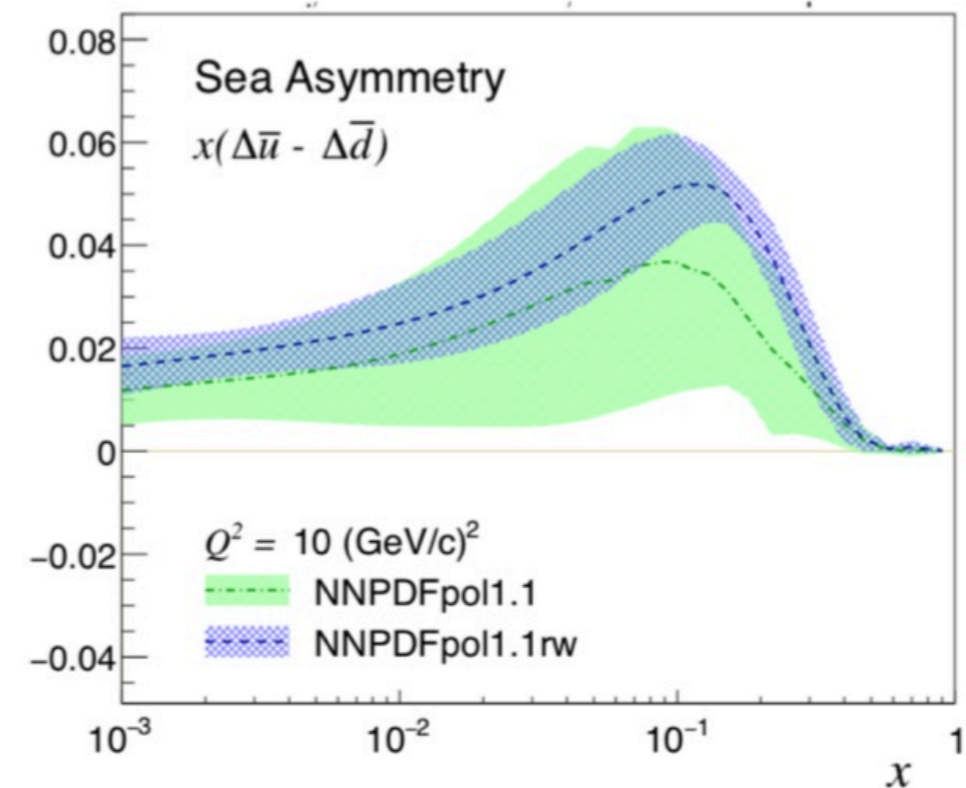
Quark Polarization at RHIC



See e.g. J. Zhang, INPC
 Q.H. Xu, DIS

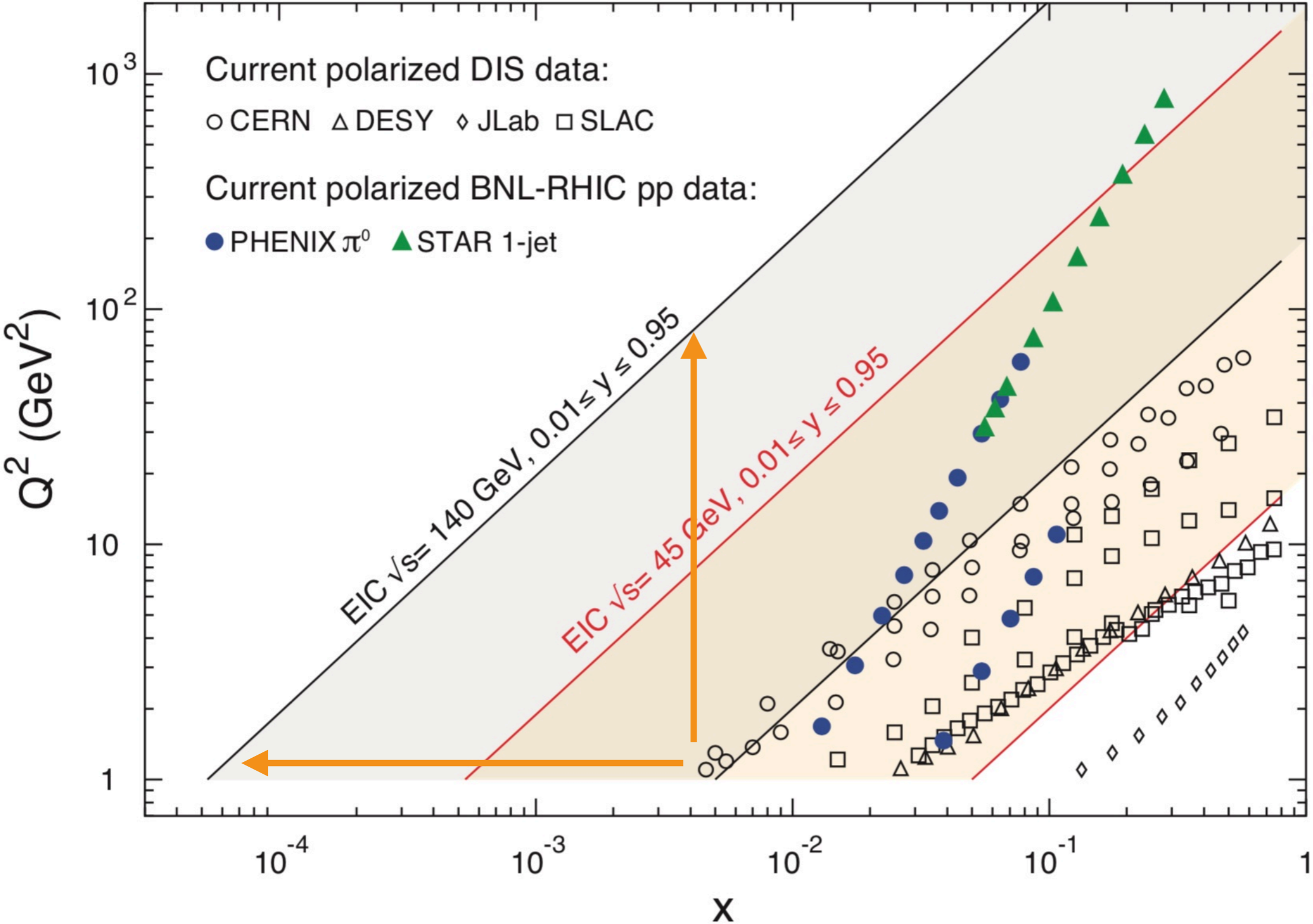
- Further confirmed the polarized sea asymmetry:

$$\Delta\bar{u} > \Delta\bar{d}$$



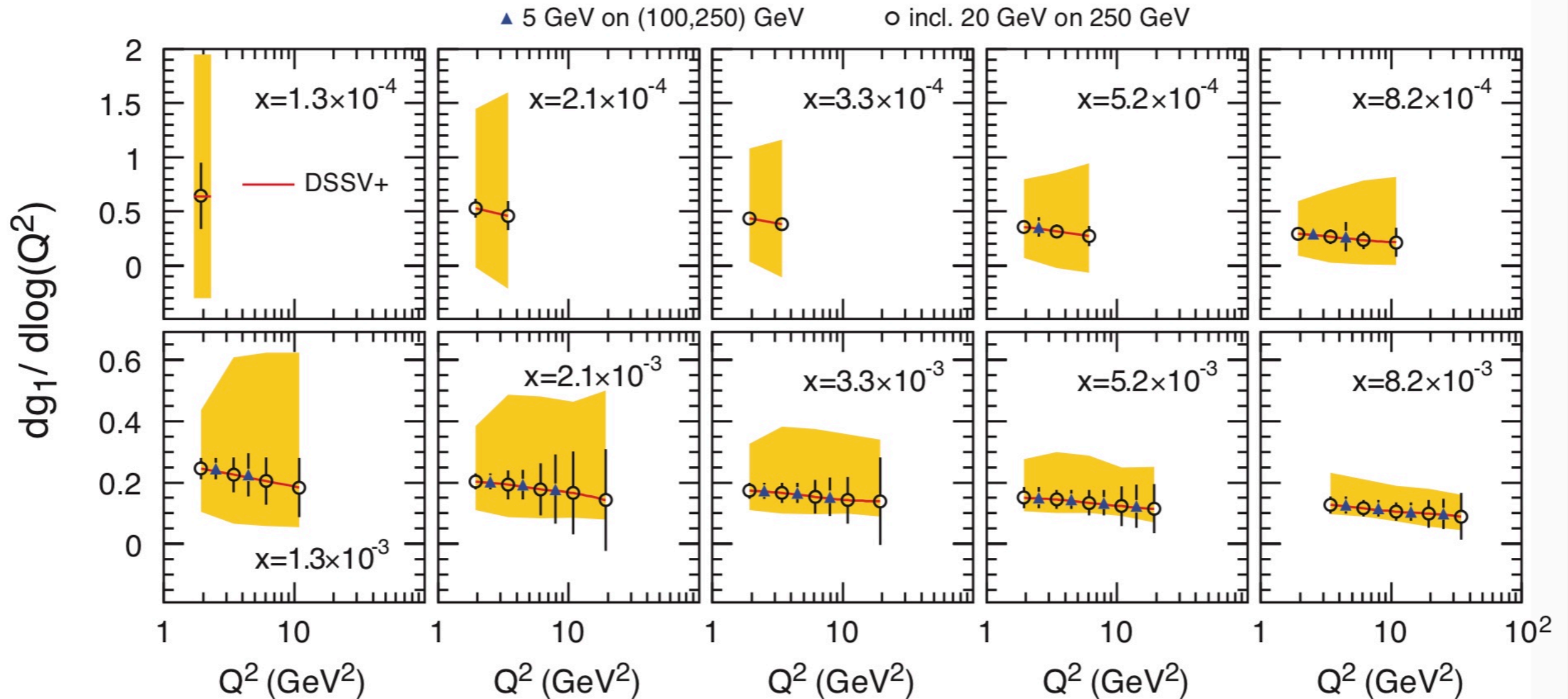
The light quark-sea is polarized and exhibits a flavor asymmetry.

Outlook to EIC



Two orders in x and Q^2 compared to existing data; few, if any, alternatives.

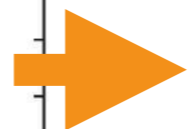
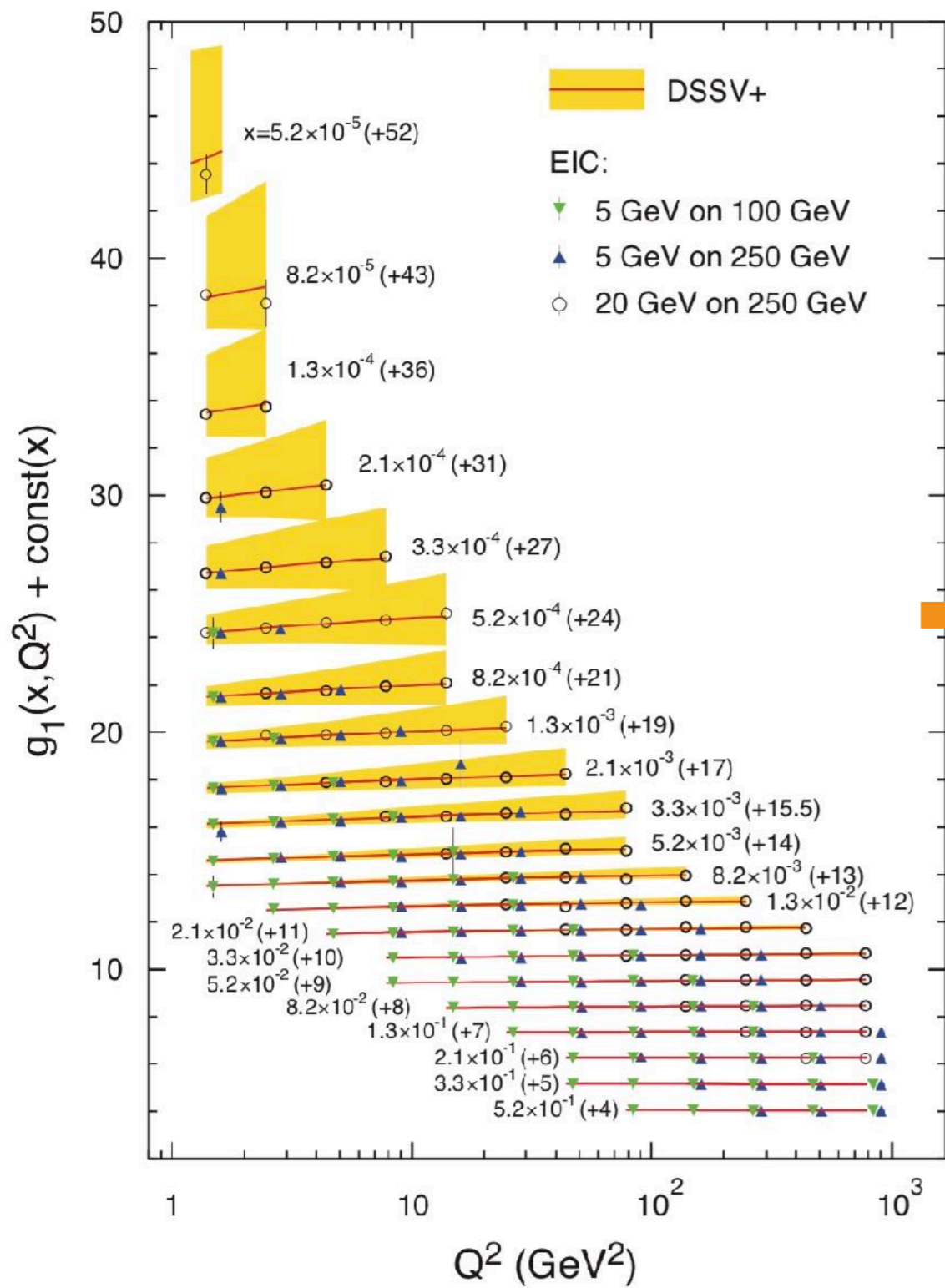
Outlook to EIC



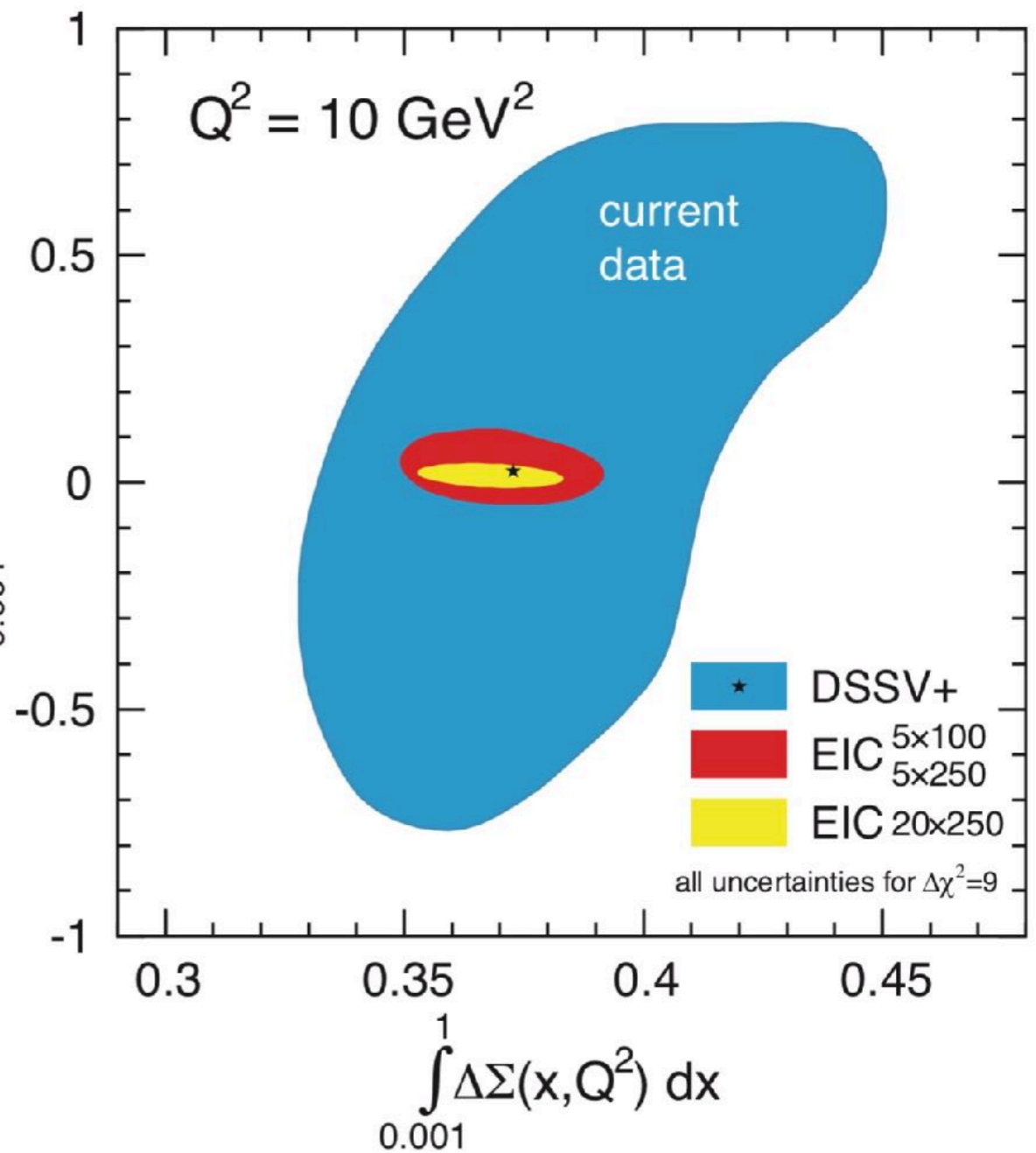
Direct sensitivity to scale-dependence, and hence gluon polarization, at least to $x \sim 10^{-3}$

Simultaneous access also to a host of complementary channels,
e.g. open charm production.

Outlook to EIC



$$\int_{0.001}^1 \Delta g(x, Q^2) dx$$



Conclusive insights in quark and gluon helicity from inclusive measurements, and orbital momentum by subtraction (!)

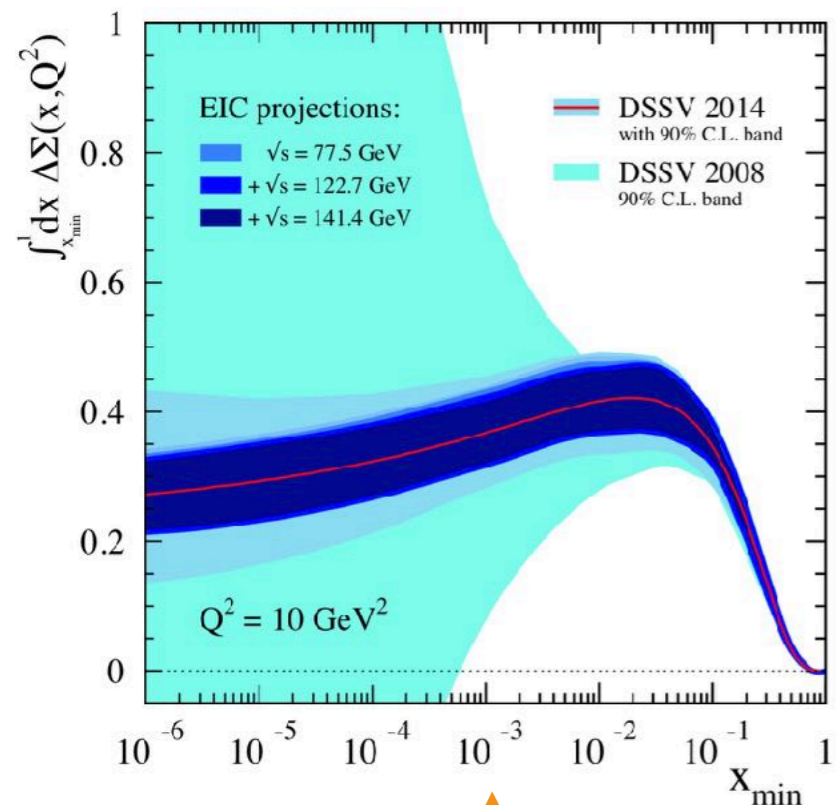
Outlook to EIC

A more up-to-date view; E. Aschenauer et al.
PRD 92 (2015) 094030

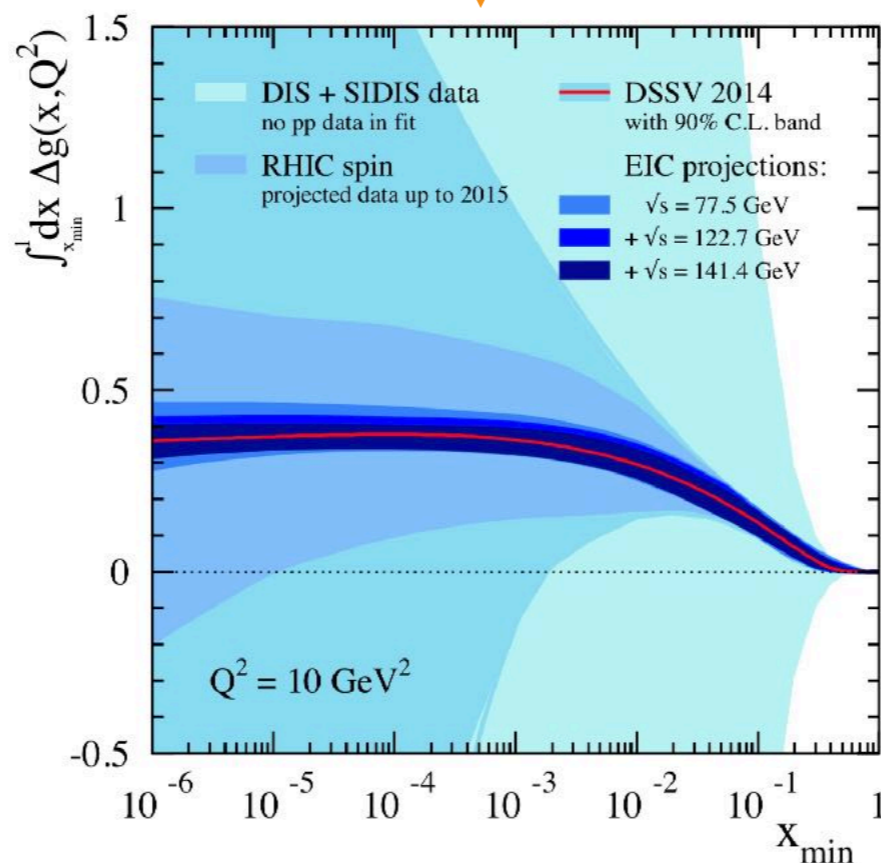
gluon helicity



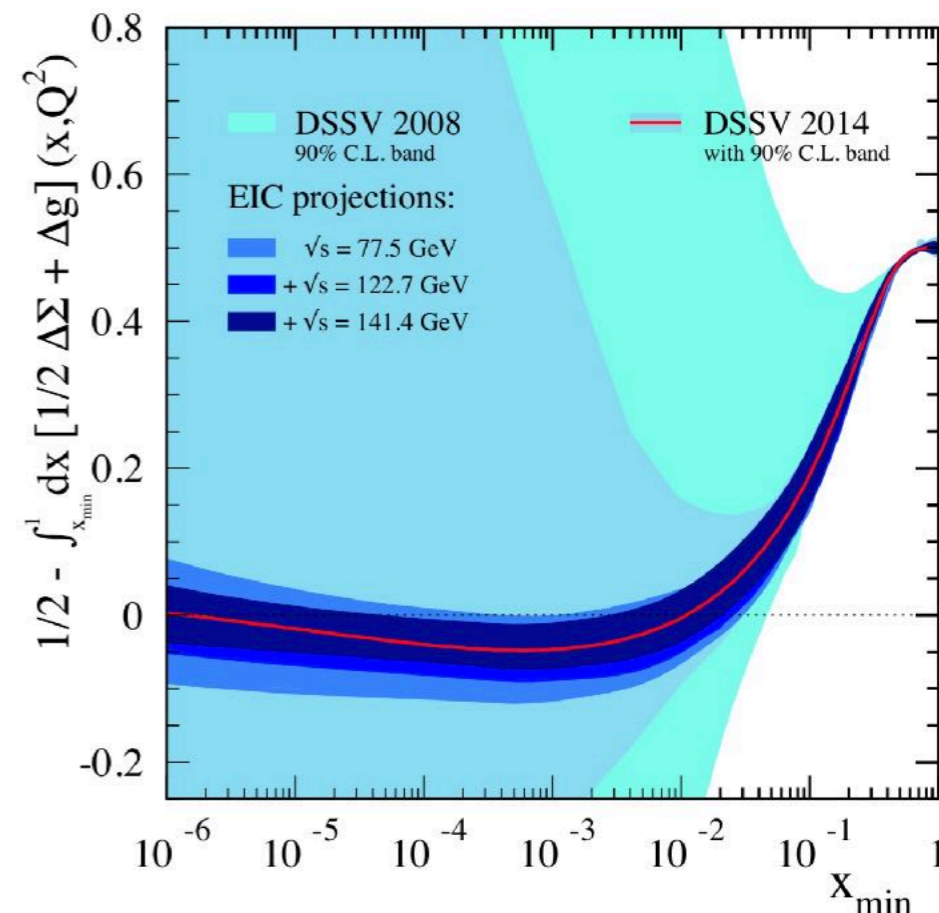
“orbital momenta”

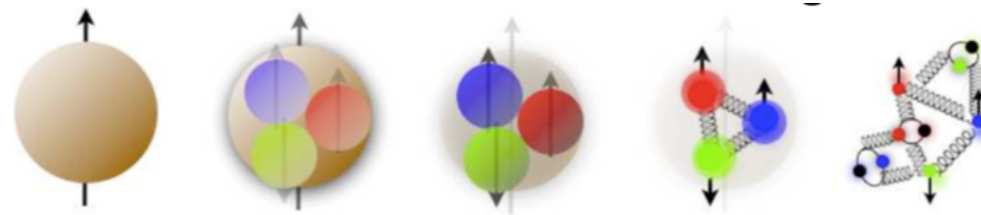


(anti-)quark helicity



Clearly requires EIC





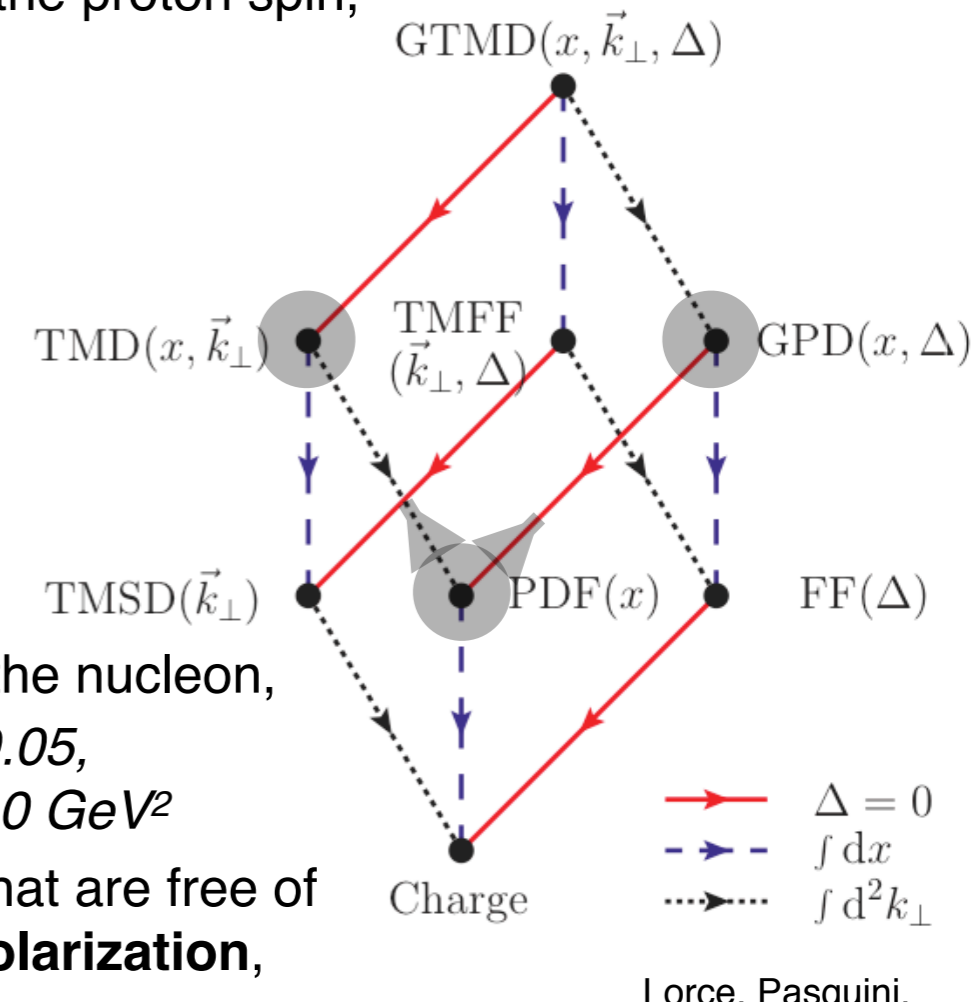
We are still far from fundamental knowledge and understanding of nucleon spin

DIS data:

- small- x measurements provided the impetus for renewed study of the proton spin,
- data on proton and neutron targets over a wide x -range, confirming the Bjorken Sum rule, decent insight in the sum of quark and anti-quark spins,
- initial sensitivities to scale dependence,
- best (lack of) insight in strangeness,
- going beyond collinear distributions,

RHIC spin program:

- has achieved the most sensitive insights in **gluon polarization** in the nucleon, *gluons are positively polarized for momentum fractions $x > 0.05$, at the level of $0.2 h$ for $Q^2 = 10 \text{ GeV}^2$*
- has provided evidence, with measurements at the W -mass scale that are free of fragmentation uncertainties, of non-perturbative **sea-quark polarization**,
- (quite promising TMD measurements; a talk by itself)



Lattice QCD:

- considerable progress, *also* on x -dependence,
- combination with data in the foreseeable future ?

EIC and theory will be essential to solve the “spin puzzle.”