

*“Probing Nucleons and Nuclei in High Energy Collisions”, Oct. 11, 2018, INT, Seattle, WA*

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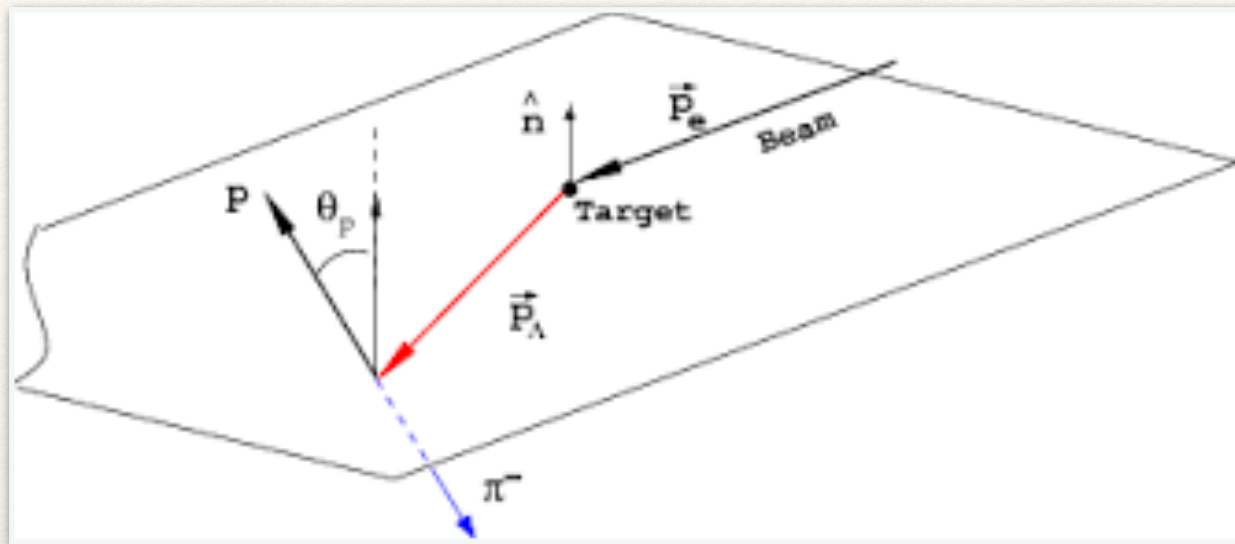
# Collinear twist-3 formalism for $\Lambda$ polarization

Marc Schlegel  
Department of Physics  
New Mexico State University

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# Measurement of $\Lambda$ -spin through decay $\Lambda^0 \rightarrow p\pi^-$



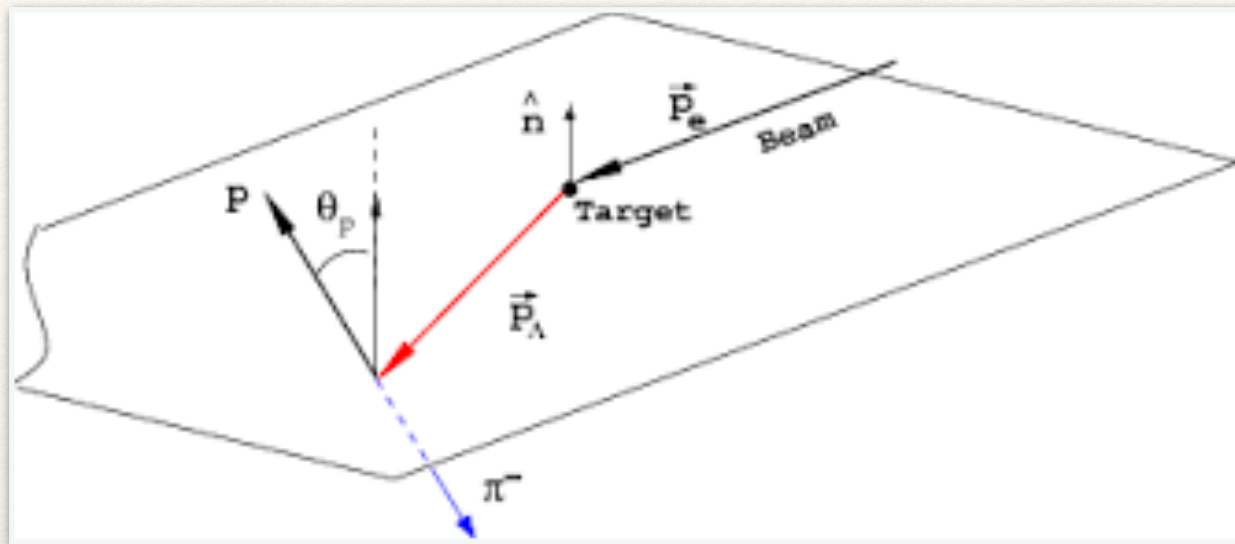
- Proton preferentially emitted along  $\Lambda$ -spin
- In  $\Lambda$  rest frame: pol. decay distribution

$$\left(\frac{dN}{d\Omega_p}\right)_{\text{pol}} = \left(\frac{dN}{d\Omega_p}\right)_{\text{unpol}} (1 + \alpha P_n^\Lambda \cos(\theta_p))$$

$P^\Lambda$ : Transverse Lambda Polarization



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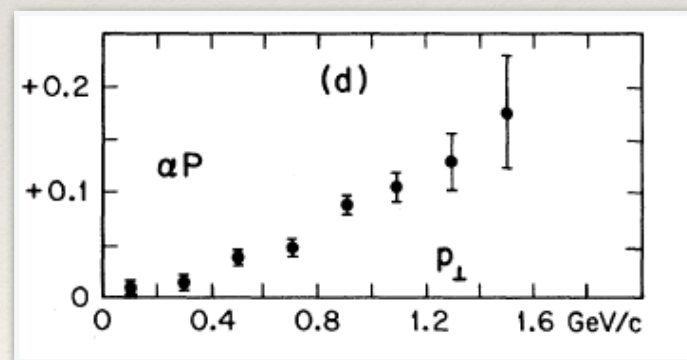
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## Transverse $\Lambda$ polarization in pA: long history...

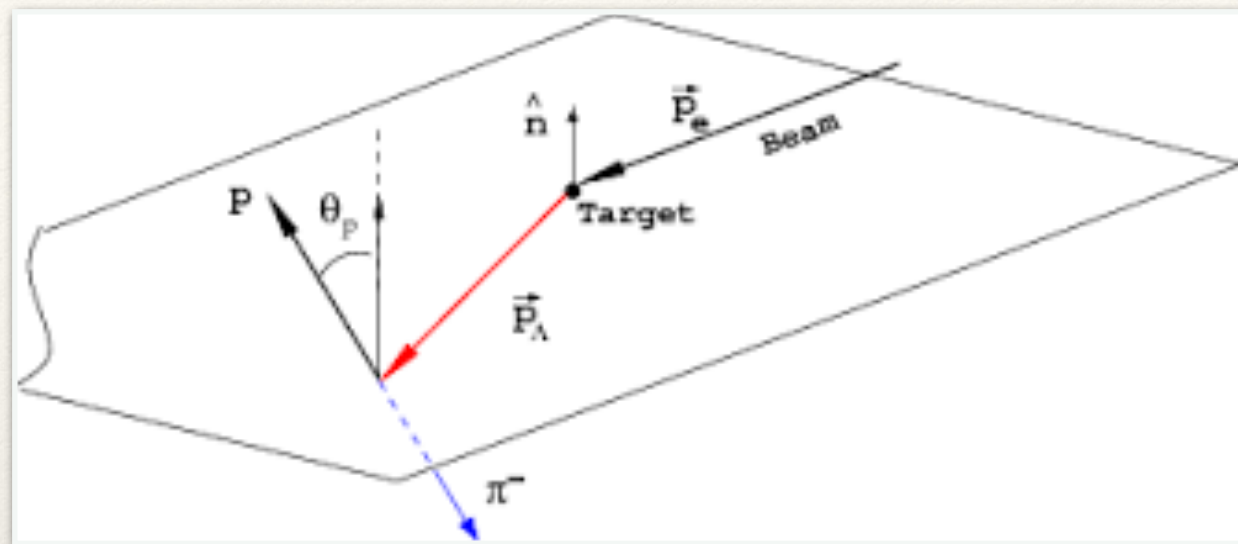
One of the first transverse spin effects at Fermilab (1976):  $p+\text{Be} \rightarrow \Lambda^0 + X$   
and many more follow-up measurements, also at CERN SPS (NA48), HERA-B



$\Lambda$  polarization was found to be sizeable!



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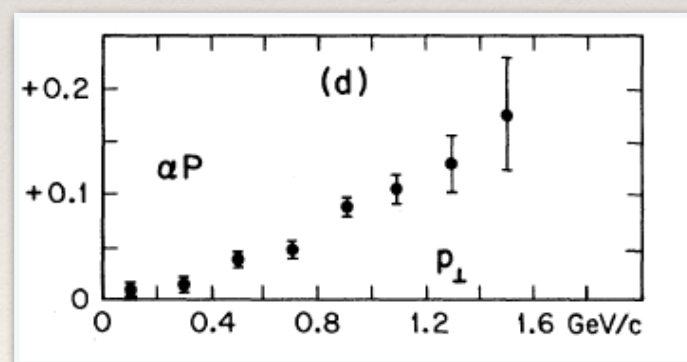
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HERMES (2007, 2014): [0704.3133, 1406.3236]

non-zero polarization in quasi-real photoproduction of  $\Lambda^\uparrow$  off nuclei!



What about LHC? Is it feasible at a high energy collider?



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## Recent ATLAS measurement at $\sqrt{S} = 7$ TeV

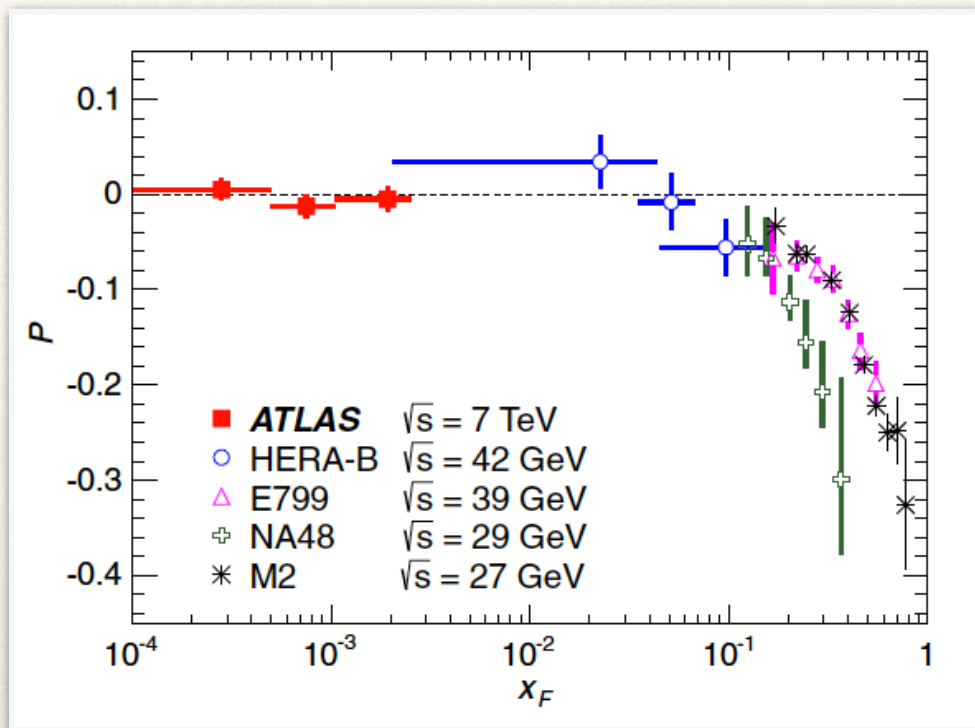
[ATLAS, PRD 91, 032004 (2015)]

Polarization small at mid-rapidity

$\Lambda$  polarization at LHC possible

Can  $\Lambda$  polarization be useful for LHC physics?

Tool in particle physics?

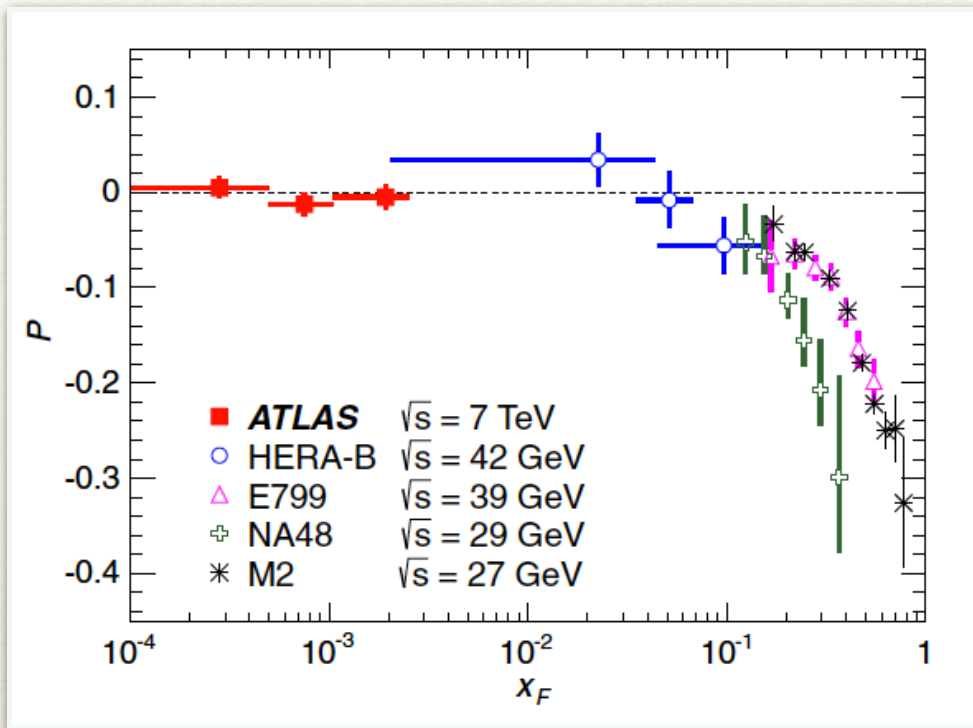




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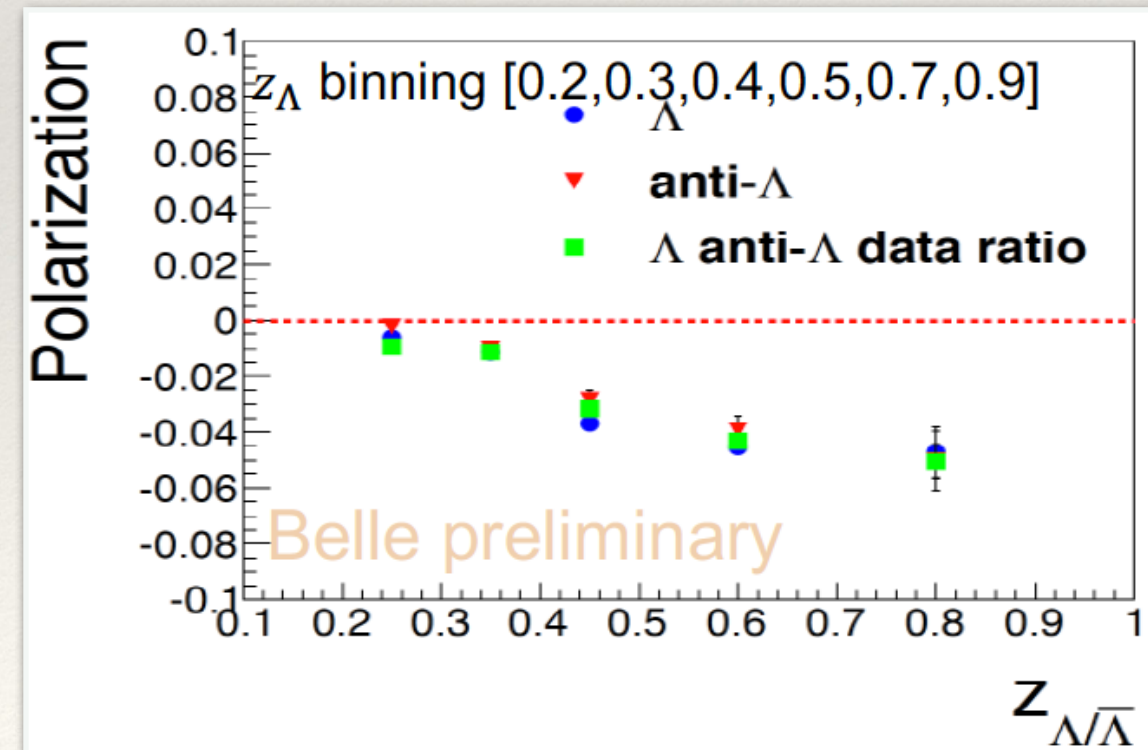


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Can  $\Lambda$  polarization be useful for LHC physics?  
 Tool in particle physics?

Simplest and cleanest process (like DIS):  $e^+ e^- \longrightarrow \Lambda^\uparrow X$

- ❖ [OPAL at LEP on Z-pole \[Eur.Phys.J C2, 49 \(1998\)\]](#):  
 Longitudinal Polarization,  
 no significant Transverse Polarization
- ❖ [Preliminary Belle data: Transverse Polarization \[Yinghui Guan, SPIN 2016\]](#)  
 → talk by A. Vossen  
 ⇒ significant transverse polarization  
 (measured w.r.t. thrust axis !)

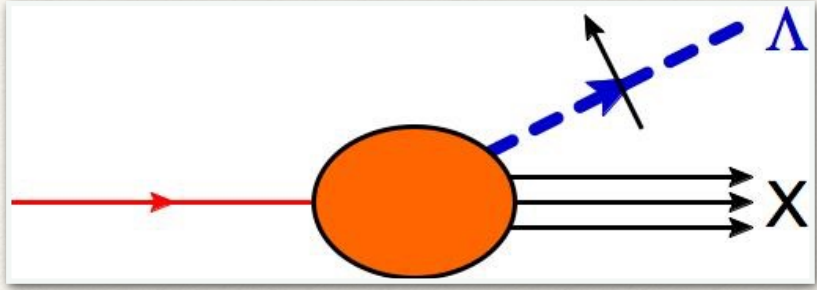




$\Lambda^\uparrow$  - production  
in pQCD



# Perturbative QCD at leading twist: $\Lambda$ fragmentation

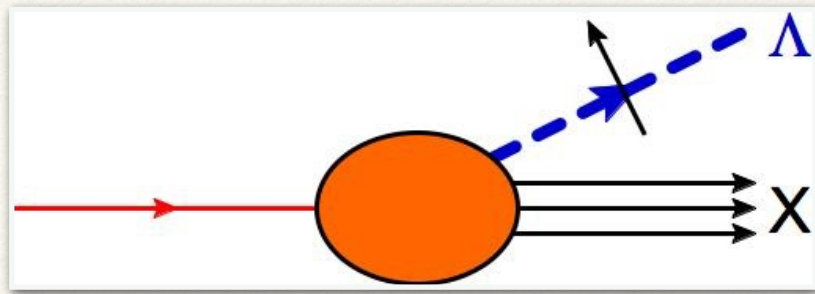


parton  $\longrightarrow$   $\Lambda$  +  $X$  transition:

$$\langle P_\Lambda, S_\Lambda; X | \bar{q}(0) | 0 \rangle$$



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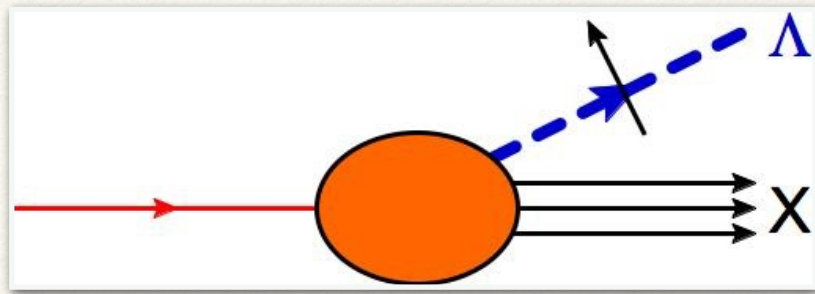
$$\langle P_\Lambda, S_\Lambda; X | \bar{q}(0) | 0 \rangle$$

‘square of the amplitude’

$$\Delta_{ij}(z) = \frac{1}{N_c} \sum_X \int \frac{d\lambda}{2\pi} e^{-i\frac{\lambda}{z}} \langle 0 | [\infty m, 0] q_i(0) | P_\Lambda, S_\Lambda; X \rangle \langle P_\Lambda, S_\Lambda; X | \bar{q}_j(\lambda m) [\lambda m, \infty m] | 0 \rangle$$



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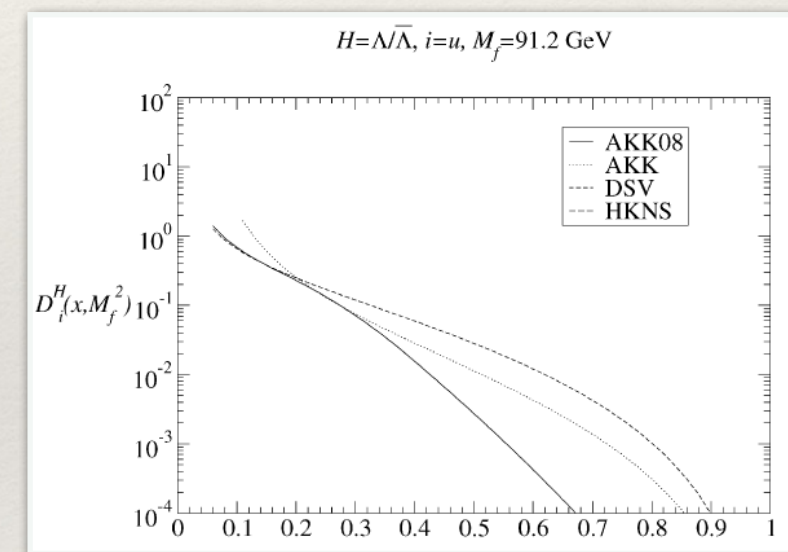
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## $\Lambda$ fragmentation functions at leading twist

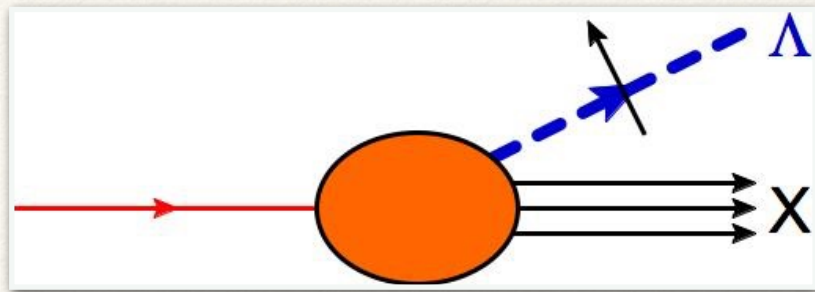
$$D_1^{\Lambda/q}(z)$$

FF of unpolarized  $q \rightarrow \Lambda$ :  
fairly known [fits by AKK08, DSV, ...]





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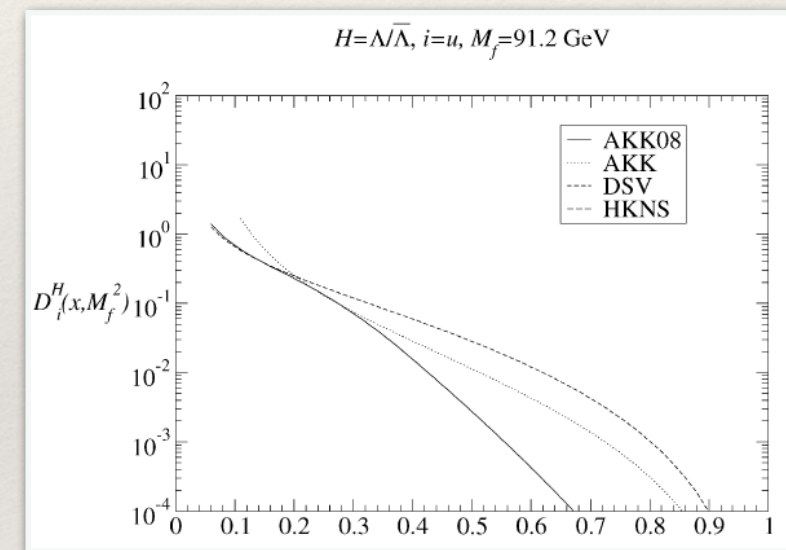
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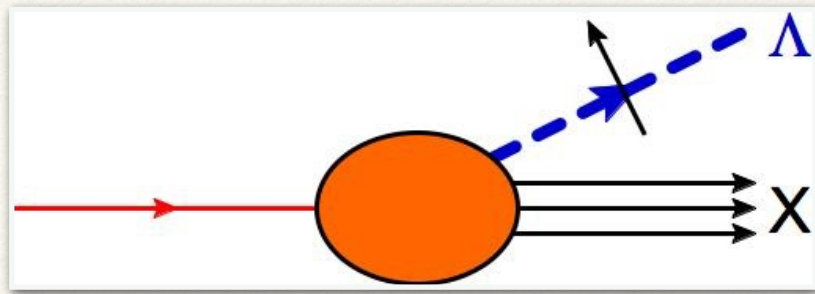


$$G_1^{\Lambda/q}(z)$$

FF of longitudinally pol.  $q \longrightarrow \Lambda$ :  
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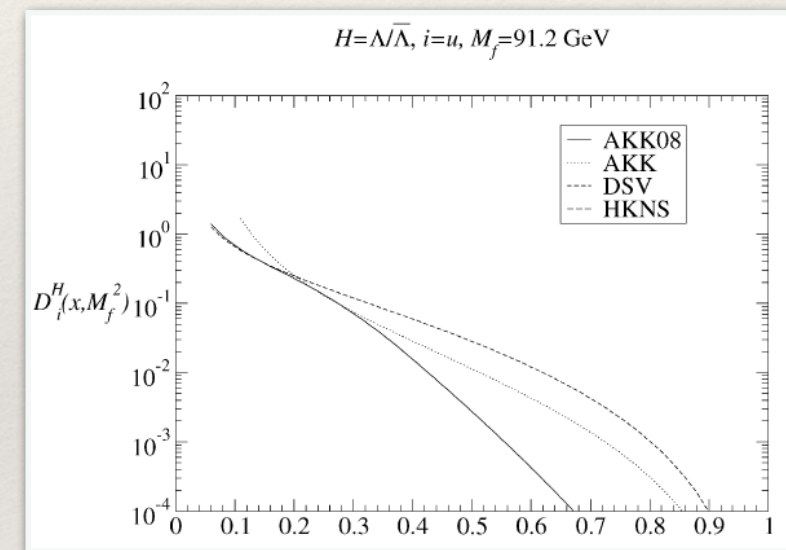
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poorly known [attempts by DSV to fit LEP data]

$$H_1^{\Lambda/q}(z)$$

FF of transversely pol.  $q \rightarrow \Lambda$ :  
unknown, chiral-odd, hard to extract from single-inclusive processes  
Candidate to explain large transverse  $\Lambda$  polarization?



# Collinear Twist-3 formalism

'intrinsic' twist-3 FF with transverse spin:

$$G_T^{\Lambda/q}(z)$$

$$D_T^{\Lambda/q}(z)$$



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'kinematic' twist-3 FF with transverse spin:

$$\Delta_{\partial}^{\alpha}(z) = \int d^2 p_T p_T^{\alpha} \Delta(z, z p_T)$$

→

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# Collinear Twist-3 formalism

'intrinsic' twist-3 FF with transverse spin:

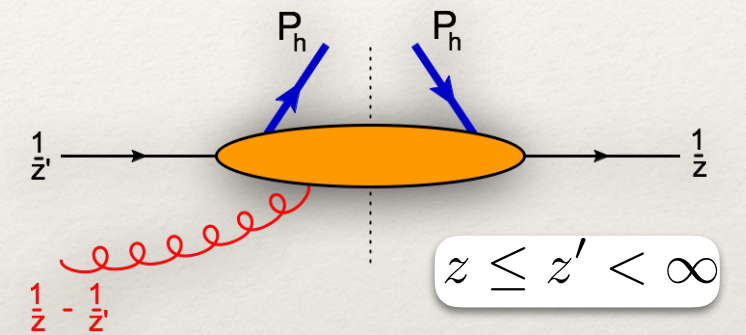
$$G_T^{\Lambda/q}(z) \quad D_T^{\Lambda/q}(z)$$

'kinematic' twist-3 FF with transverse spin:

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'dynamical' twist-3 FF with transverse spin:

$$\Delta_F^{\alpha}(z, z') \sim \langle 0 | q(\lambda m) g F^{m\alpha}(\mu m) | P_{\Lambda}, S_{\Lambda}; X \rangle \langle P_{\Lambda}, S_{\Lambda}; X | \bar{q}(0) | 0 \rangle \\ \implies \hat{D}_{FT}^{\Lambda/q}(z, z'), \hat{G}_{FT}^{\Lambda/q}(z, z')$$



complex functions:

$$FF(z, z) = 0$$

$$FF(z, 0) = 0$$

$$\left. \frac{\partial}{\partial z'} FF(z, z') \right|_{z'=z} = 0$$



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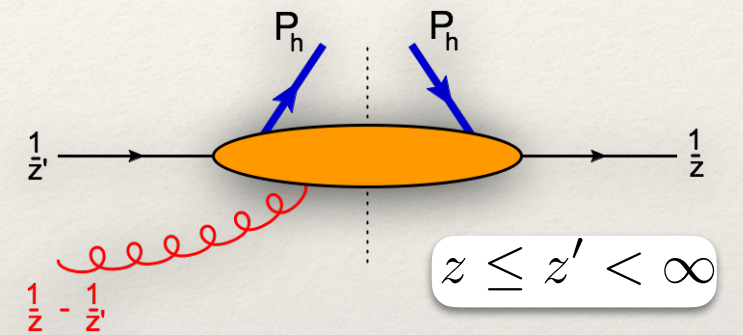
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## Relations: Equation of Motion & Lorentz-Invariance

[Kanazawa, Koike, Metz, Pitonyak, MS, PRD 93, 054024 (2016)]

$$D_{1T}^{\perp(1)}(z) + \frac{D_T(z)}{z} = \int_0^1 d\beta \frac{\Im[\hat{D}_{FT}(z, z/\beta)] - \Im[\hat{G}_{FT}(z, z/\beta)]}{1 - \beta}$$

$$G_{1T}^{\perp(1)}(z) - \frac{G_T(z)}{z} = \int_0^1 d\beta \frac{\Re[\hat{D}_{FT}(z, z/\beta)] - \Re[\hat{G}_{FT}(z, z/\beta)]}{1 - \beta}$$

$$\frac{D_T(z)}{z} = - \left( 1 - z \frac{d}{dz} \right) D_{1T}^{\perp(1)}(z) - 2 \int_0^1 d\beta \frac{\Im[\hat{D}_{FT}(z, z/\beta)]}{(1 - \beta)^2}$$

$$\frac{G_T(z)}{z} = \frac{G_1(z)}{z} + \left( 1 - z \frac{d}{dz} \right) G_{1T}^{\perp(1)}(z) - 2 \int_0^1 d\beta \frac{\Re[\hat{G}_{FT}(z, z/\beta)]}{(1 - \beta)^2}$$

Two equations, three functions → eliminate 'intrinsic & kinematical twist-3'



Single-inclusive  
Hard Processes suitable for  
 $\Lambda^\uparrow$  - production



$\Lambda^\uparrow$  in pp - collisions ( $p p \rightarrow \Lambda^\uparrow X$ ):

complete LO formulae not yet available, complicated!

[Koike, Metz, Pitonyak, Yabe, Yoshida, PRD 2017]

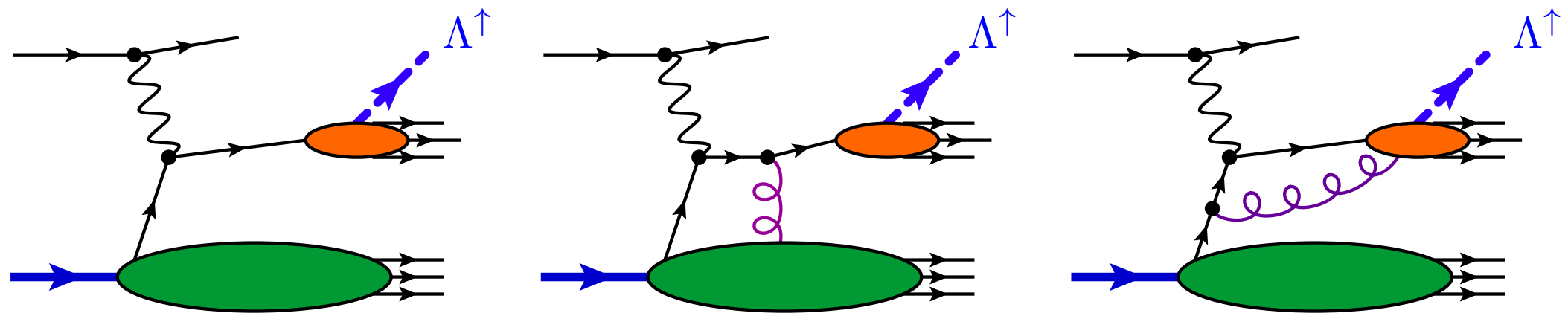


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## Single-inclusive $\Lambda$ - production ( $e p \rightarrow \Lambda^\uparrow X$ )



## LO-formula (including EoM & LIR)

[Kanazawa, Koike, Metz, Pitonyak, MS, PRD 2016]

$$\sigma(S_\Lambda) \propto \int \hat{\sigma}_1 \otimes f_1^q \otimes \Im[\hat{D}_{FT}^q] + \int \hat{\sigma}_5 \otimes f_1^q \otimes \Im[\hat{G}_{FT}^q] + \int \hat{\sigma} \otimes \frac{d\hat{H}_{FU}^q(x, x)}{dx} \otimes H_1^q(z)$$

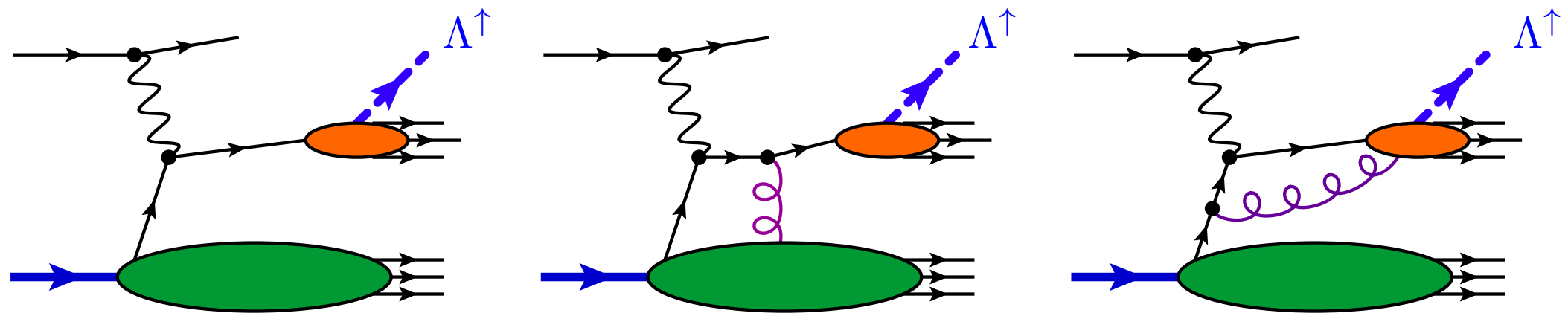


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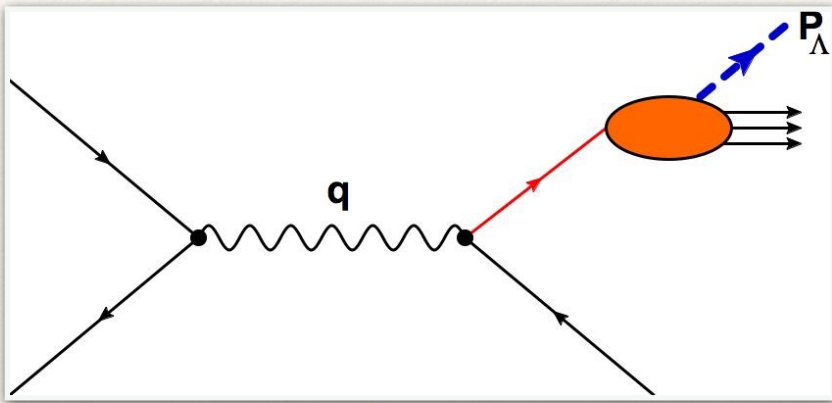
chiral-odd: may be small ?!

Physics opportunity at EIC with  $P_T >$  a few GeV!  
might help to solve a 40-year old puzzle...



# Unpolarized $e^+ e^- \rightarrow \Lambda X$ cross section

“Parton Model like” at LO



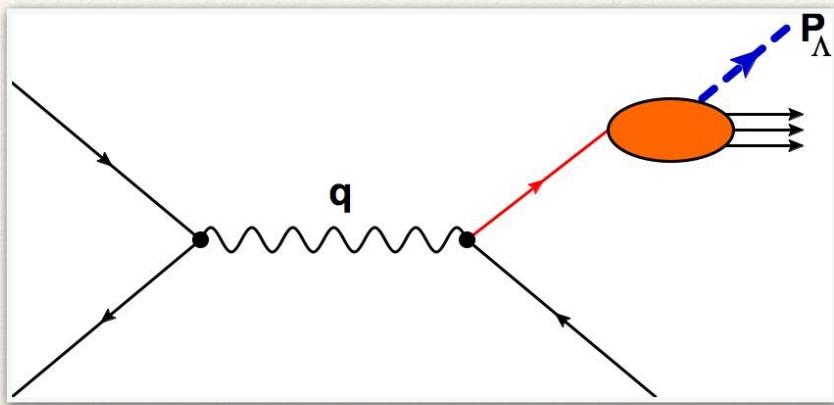
$$E_\Lambda \frac{d\sigma}{d^3 \vec{P}_\Lambda} \propto \sum_q e_q^2 D_1^{\Lambda/q}(z_h)$$

$$z_h = \frac{2P_\Lambda \cdot q}{q^2}$$



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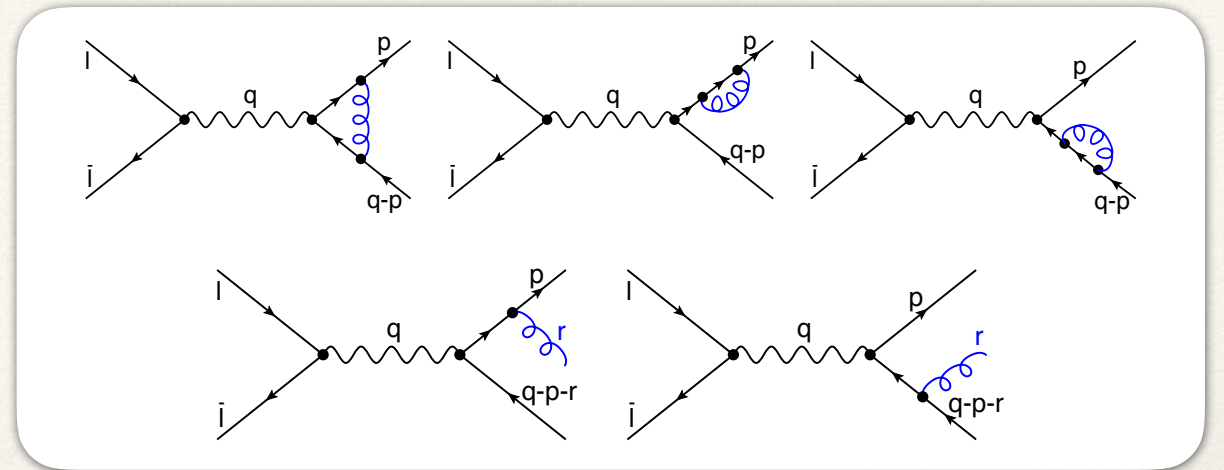
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NLO

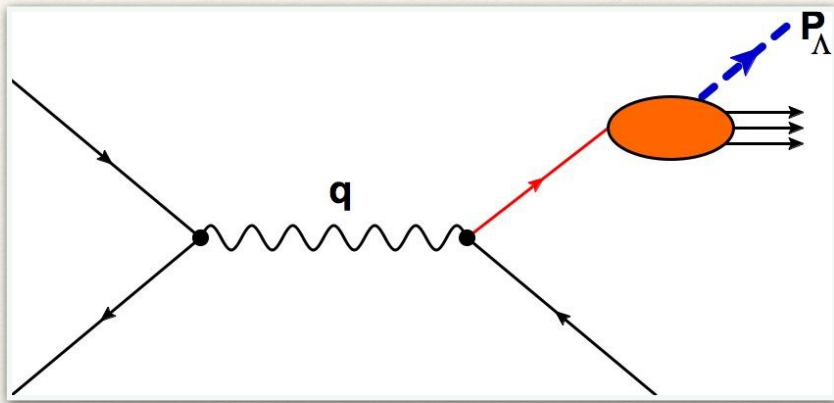


$$\left( E_\Lambda \frac{d\sigma}{d^3 \vec{P}_\Lambda} \right)_{\text{NLO}} \propto \sum_q e_q^2 \int_{z_h}^1 \frac{dw}{w} \left[ \hat{\sigma}^{\text{MS},q}(w, s/\mu^2) D_1^{\Lambda/q}(z_h/w, \mu) + \hat{\sigma}^{\text{MS},g}(w, s/\mu^2) D_1^{\Lambda/g}(z_h/w, \mu) \right]$$



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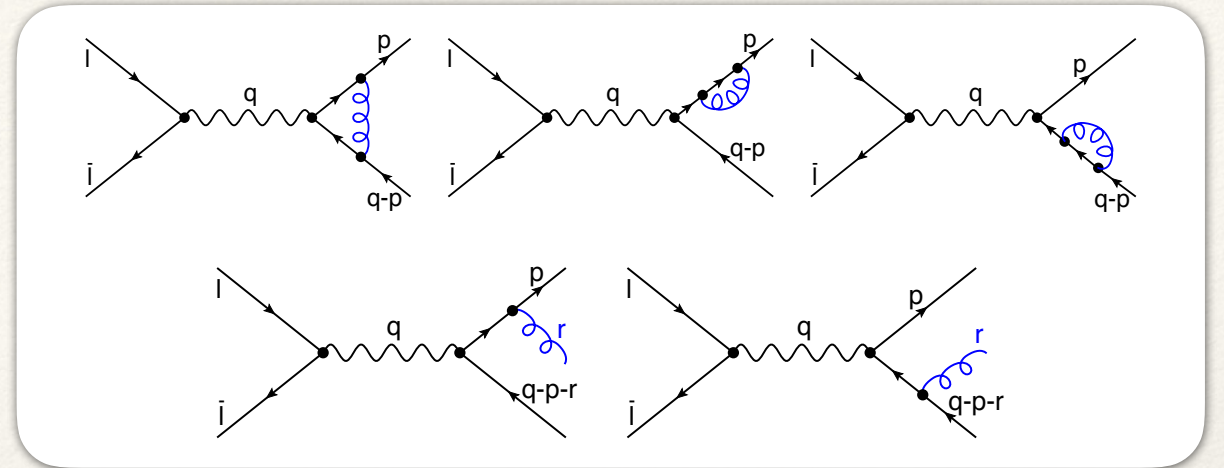
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Typical NLO features:

- ❖ infrared safe (cancellation of  $1/\epsilon^2$  - poles in dim. reg.)

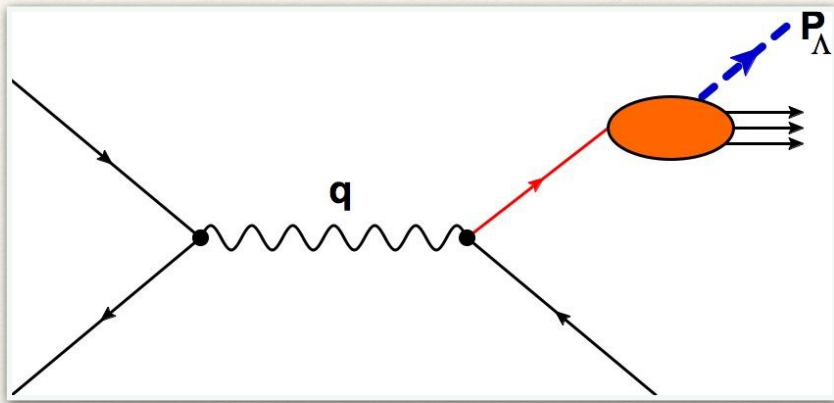
$$\hat{\sigma}_{\text{virt}} + \hat{\sigma}_{\text{real}} = \mathcal{O}(1/\epsilon)$$

$$\hat{\sigma}^{q/g} \propto -\frac{1}{\epsilon} P_{q/g} q(w) + \mathcal{O}(\epsilon^0)$$



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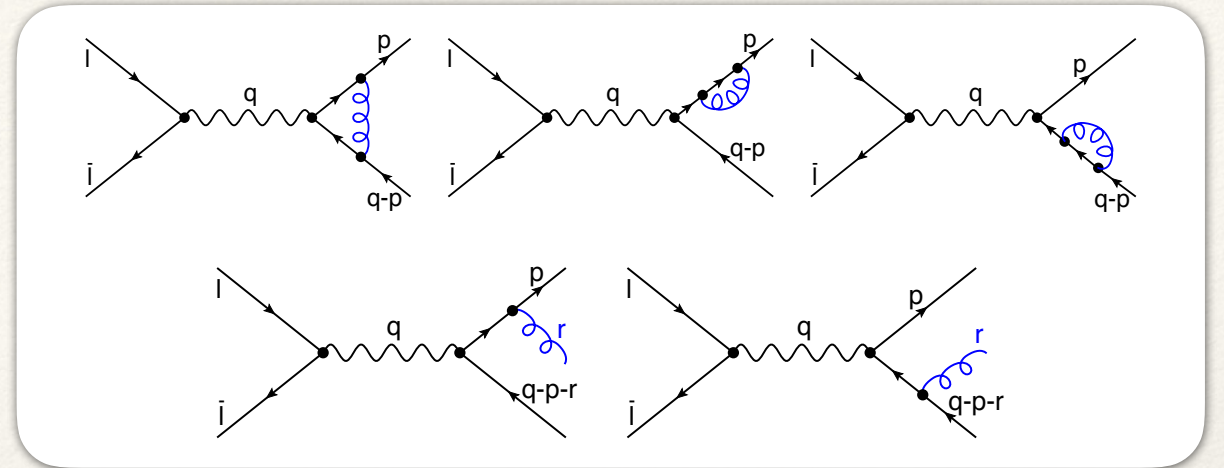
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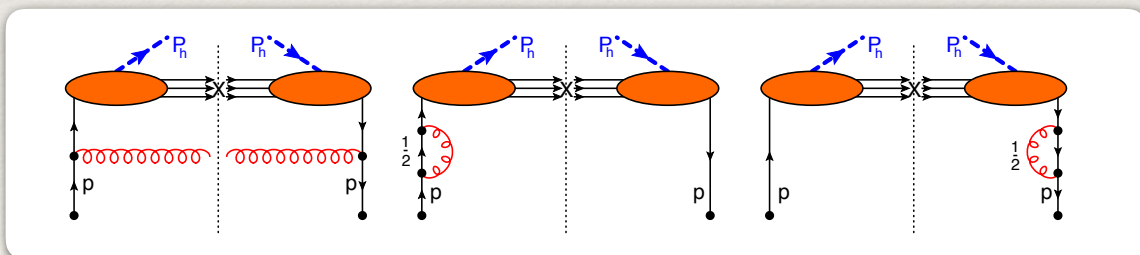
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- $\overline{\text{MS}}$  renormalization of fragmentation functions  $\rightarrow$  DGLAP evolution



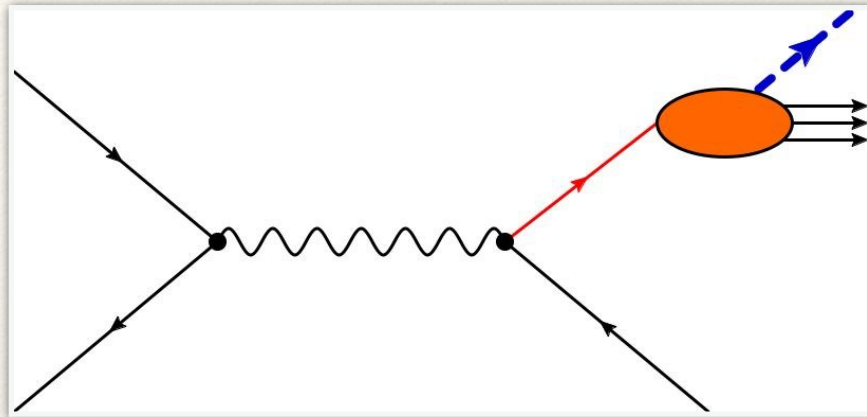
$$D_{1,\text{bare}}^{\Lambda/q}(z) = D_{1,\text{ren}}^{\Lambda/q}(z) + \frac{\alpha_s}{2\pi} \frac{S_\epsilon^{\overline{\text{MS}}}}{\epsilon} \sum_{i=q,g} \int_z^1 \frac{dw}{w} P_{iq}(w) D_{1,\text{ren}}^{\Lambda/i}\left(\frac{z}{w}\right) + \mathcal{O}(\alpha_s^2)$$

$\mathcal{O}(1/\epsilon)$  cancels,  
necessary condition for  
one-loop factorization!

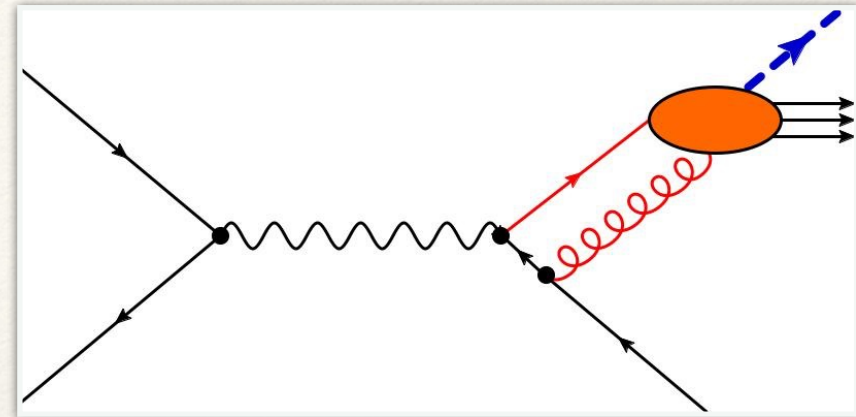


## Transverse $\Lambda$ polarization at LO

'intrinsic' & 'kinematical' twist-3 FF:



'dynamical' twist-3 FF:



$$\frac{d\sigma(S_{\Lambda T})}{dz_h d\phi} = C |S_{\Lambda T}| \sin(\phi_S) \sum_q e_q^2 \left[ \frac{D_T^{\Lambda/q}(z_h)}{z_h} - D_{1T}^{\perp(1)\Lambda/q}(z_h) + \int_0^1 d\beta \frac{\Im[\hat{D}_{FT} - \hat{G}_{FT}]^{\Lambda/q}(z_h, z_h/\beta)}{1-\beta} \right]$$



**Equation of Motion:**

$$\frac{d\sigma(S_{\Lambda T})}{dz_h d\phi} = C |S_{\Lambda T}| \sin(\phi_S) \sum_q e_q^2 \left[ 2 \frac{D_T^{\Lambda/q}(z_h)}{z_h} \right]$$

**or:**

$$\frac{d\sigma(S_{\Lambda T})}{dz_h d\phi} = C |S_{\Lambda T}| \sin(\phi_S) \sum_q e_q^2 \left[ -2 D_{1T}^{\perp(1)\Lambda/q}(z_h) + 2 \int_0^1 d\beta \frac{\Im[\hat{D}_{FT} - \hat{G}_{FT}]^{\Lambda/q}(z_h, z_h/\beta)}{1-\beta} \right]$$

### Single-Transverse $\Lambda$ spin asymmetry

- ❖ Unique effect driven by a single fragmentation function  $D_T \rightarrow$  absent in DIS ( $1\gamma$ )
- ❖ EoM needed at LO to preserve e.m. current conservation of hadronic tensor ( $q_\mu W^{\mu\nu} = 0$ ) (EoM not optional!)



# Transverse $\Lambda$ polarization at NLO

[Gamberg, Kang, Pitonyak, M.S., Yoshida, released soon]

- ❖ Study the NLO dynamics for twist-3 fragmentation in the simplest process
- ❖ Different compared to twist-3 distributions (no pole contributions)



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[Gamberg, Kang, Pitonyak, M.S., Yoshida, released soon]

- ❖ Study the NLO dynamics for twist-3 fragmentation in the simplest process
- ❖ Different compared to twist-3 distributions (no pole contributions)

Virtual & Real diagrams (qg/q - channel here, gg/g, qqb/g not shown)

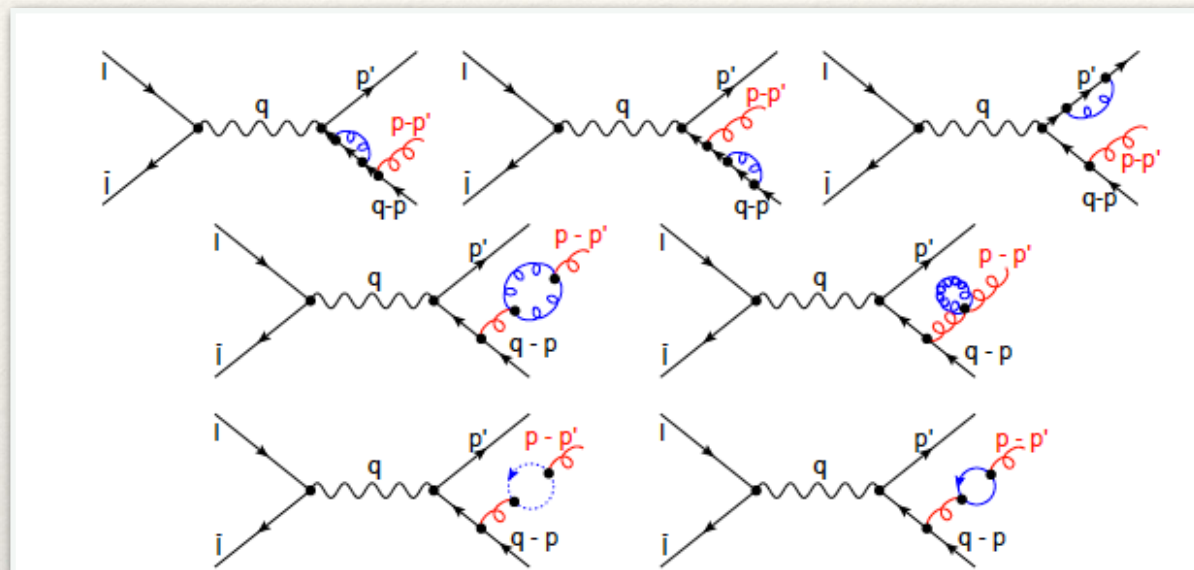


Figure 6: Self-energy corrections

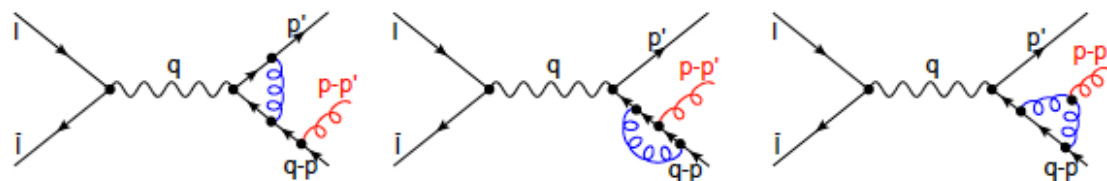


Figure 7: Vertex corrections

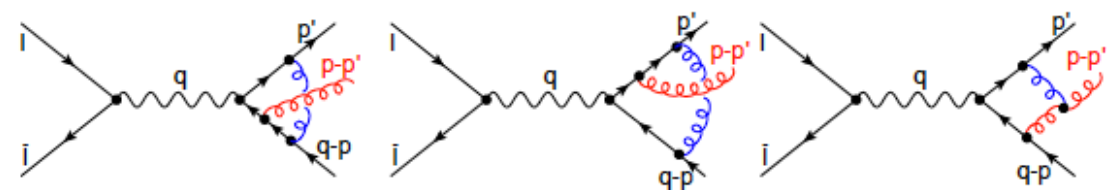


Figure 8: Box corrections

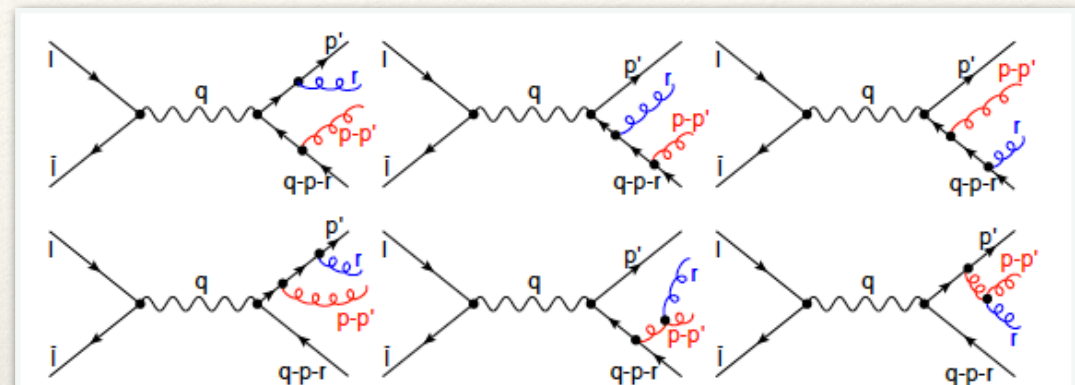


Figure 9: Real corrections

- ❖ E.o.M. - relations are crucial:  
Eliminate 'intrinsic' twist-3 contributions: only then color gauge invariance at NLO! ✓
- ❖ Imaginary parts: In the dynamical fragmentation process & loop diagrams  
Infrared  $1/\epsilon^2$  - poles cancel ✓
- ❖  $1/\epsilon$  - poles of imaginary parts of loops cancel through E.o.M. ✓
- ❖  $1/\epsilon$  - collinear poles of real parts of loops through MSbar - renormalization (?)



## Complete structure of the NLO result w/o intrinsic twist-3

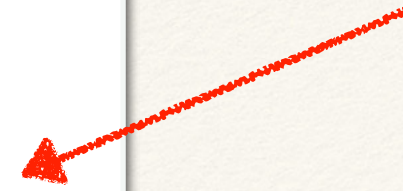
$$E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) = (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l S_h}}{s} (2v-1) \times$$
$$\sum_{q=u,\bar{u},\dots} e_q^2 \left[ -2 D_{1T}^{\perp(1),q}(z_h) + 2 \int_0^1 d\beta \frac{\Im[\hat{D}_{FT}^q - \hat{G}_{FT}^q](z_h, \frac{z_h}{\beta})}{1-\beta} \right]$$



## Complete structure of the NLO result w/o intrinsic twist-3

$$E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) = (4\pi^2 z_h^2)^\varepsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l S_h}}{s} (2v-1) \times$$
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LO

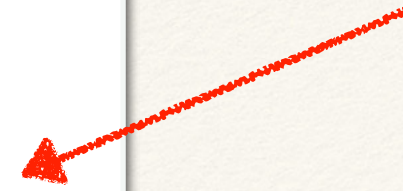




## Complete structure of the NLO result w/o intrinsic twist-3

$$\begin{aligned}
 E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) &= (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l S_h}}{s} (2v-1) \times \\
 &\sum_{q=u,\bar{u},\dots} e_q^2 \left[ -2 D_{1T}^{\perp(1),q}(z_h) + 2 \int_0^1 d\beta \frac{\Im[\hat{D}_{FT}^q - \hat{G}_{FT}^q](z_h, \frac{z_h}{\beta})}{1-\beta} \right. \\
 &\left. + \frac{\alpha_s}{2\pi} S_\epsilon \int_{z_h}^1 \frac{dw}{w^2} \int_0^1 d\beta \left\{ \hat{\sigma}_{D_{1T}^{\perp(1)}}^{q;\text{EoM}}(w) D_{1T}^{\perp(1),q}\left(\frac{z_h}{w}\right) \right. \right.
 \end{aligned}$$

LO





# Complete structure of the NLO result w/o intrinsic twist-3

$$\begin{aligned}
 E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) &= (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l S_h}}{s} (2v-1) \times \\
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 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM



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 E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) &= (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l S_h}}{s} (2v-1) \times \\
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 &\left. \left. + \hat{\sigma}_{D_{1T}^{\perp(1)}}^{g;\text{EoM}}(w) D_{1T}^{\perp(1),g}\left(\frac{z_h}{w}\right) + \hat{\sigma}_{H_1^{(1)}}^{g;\text{EoM}}(w) H_1^{(1)g}\left(\frac{z_h}{w}\right) \right\}
 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM



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 E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) &= (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l S_h}}{s} (2v-1) \times \\
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 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM

NLO

2-gluon correlation w/ EoM



# Complete structure of the NLO result w/o intrinsic twist-3

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 &+ \frac{\alpha_s}{2\pi} S_\epsilon \int_{z_h}^1 \frac{dw}{w^2} \int_0^1 d\beta \left\{ \hat{\sigma}_{D_{1T}^{\perp(1)}}^{q;\text{EoM}}(w) D_{1T}^{\perp(1),q}\left(\frac{z_h}{w}\right) \right. \\
 &+ \hat{\sigma}_{D_{1T}^{\perp(1)}}^{g;\text{EoM}}(w) D_{1T}^{\perp(1),g}\left(\frac{z_h}{w}\right) + \hat{\sigma}_{H_1^{(1)}}^{g;\text{EoM}}(w) H_1^{(1)g}\left(\frac{z_h}{w}\right) \\
 &\left. \left. + \hat{\sigma}_{\Im D_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{D}_{FT}^q](\frac{z_h}{w}, \frac{z_h}{w\beta})}{1-\beta} + \hat{\sigma}_{\Im G_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{G}_{FT}^q](\frac{z_h}{w}, \frac{z_h}{w\beta})}{1-\beta} \right\}
 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM

NLO

2-gluon correlation w/ EoM



# Complete structure of the NLO result w/o intrinsic twist-3

$$\begin{aligned}
 E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) &= (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l S_h}}{s} (2v-1) \times \\
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 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM

NLO

2-gluon correlation w/ EoM

NLO

q-gluon-q correlation w/ EoM



# Complete structure of the NLO result w/o intrinsic twist-3

$$\begin{aligned}
 E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) &= (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l S_h}}{s} (2v-1) \times \\
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 &+ \hat{\sigma}_1^{gg;\text{EoM}}(w, \beta) \Im[\hat{D}_{FT}^{gg} - \hat{G}_{FT}^{gg} + (1-\epsilon)\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 &+ \hat{\sigma}_3^{gg;\text{EoM}}(w, \beta) \Im[(1-\epsilon)\hat{D}_{FT}^{gg} + \hat{G}_{FT}^{gg} + \frac{\epsilon}{2}\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 &\left. \right]
 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM

NLO

2-gluon correlation w/ EoM

NLO

q-gluon-q correlation w/ EoM



# Complete structure of the NLO result w/o intrinsic twist-3

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 &+ \hat{\sigma}_1^{gg;\text{EoM}}(w, \beta) \Im[\hat{D}_{FT}^{gg} - \hat{G}_{FT}^{gg} + (1-\epsilon)\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 &+ \hat{\sigma}_3^{gg;\text{EoM}}(w, \beta) \Im[(1-\epsilon)\hat{D}_{FT}^{gg} + \hat{G}_{FT}^{gg} + \frac{\epsilon}{2}\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 &\left. \right]
 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM

NLO

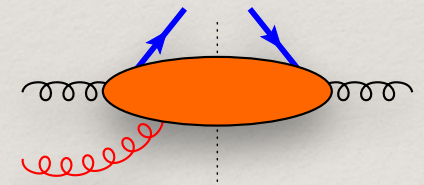
2-gluon correlation w/ EoM

NLO

q-gluon-q correlation w/ EoM

NLO

triple-gluon correlation w/ EoM





# Complete structure of the NLO result w/o intrinsic twist-3

$$\begin{aligned}
 E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) &= (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l S_h}}{s} (2v-1) \times \\
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 &+ \hat{\sigma}_{\Im D_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{D}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} + \hat{\sigma}_{\Im G_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{G}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} \\
 &+ \hat{\sigma}_1^{gg;\text{EoM}}(w, \beta) \Im[\hat{D}_{FT}^{gg} - \hat{G}_{FT}^{gg} + (1-\epsilon)\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 &+ \hat{\sigma}_3^{gg;\text{EoM}}(w, \beta) \Im[(1-\epsilon)\hat{D}_{FT}^{gg} + \hat{G}_{FT}^{gg} + \frac{\epsilon}{2}\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 &+ \hat{\sigma}_{D_{FT}}^{q\bar{q};\text{EoM}}(w) \left( \sum_{q=u,d,\dots} \Im[\hat{D}_{FT}^{q\bar{q}}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \right)
 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM

NLO

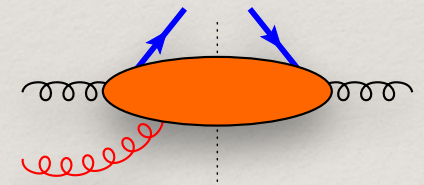
2-gluon correlation w/ EoM

NLO

q-gluon-q correlation w/ EoM

NLO

triple-gluon correlation w/ EoM





# Complete structure of the NLO result w/o intrinsic twist-3

$$\begin{aligned}
 E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) &= (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l} S_h}{s} (2v-1) \times \\
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 &+ \hat{\sigma}_{D_{1T}^{\perp(1)}}^{g;\text{EoM}}(w) D_{1T}^{\perp(1),g}\left(\frac{z_h}{w}\right) + \hat{\sigma}_{H_1^{(1)g}}^{g;\text{EoM}}(w) H_1^{(1)g}\left(\frac{z_h}{w}\right) \\
 &+ \hat{\sigma}_{\Im D_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{D}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} + \hat{\sigma}_{\Im G_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{G}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} \\
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 &+ \hat{\sigma}_3^{gg;\text{EoM}}(w, \beta) \Im[(1-\epsilon)\hat{D}_{FT}^{gg} + \hat{G}_{FT}^{gg} + \frac{\epsilon}{2}\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 &+ \hat{\sigma}_{D_{FT}}^{q\bar{q};\text{EoM}}(w) \left( \sum_{q=u,d,\dots} \Im[\hat{D}_{FT}^{q\bar{q}}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \right)
 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM

NLO

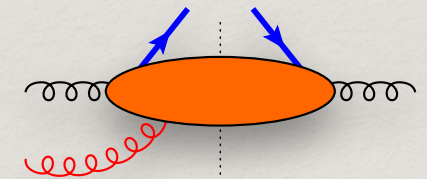
2-gluon correlation w/ EoM

NLO

q-gluon-q correlation w/ EoM

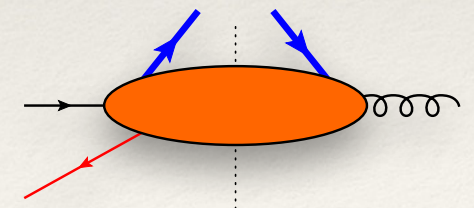
NLO

triple-gluon correlation w/ EoM



NLO

qq-gluon correlation w/ EoM





# Complete structure of the NLO result w/o intrinsic twist-3

$$\begin{aligned}
 E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) = & (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l} S_h}{s} (2v-1) \times \\
 & \sum_{q=u,\bar{u},\dots} e_q^2 \left[ -2 D_{1T}^{\perp(1),q}(z_h) + 2 \int_0^1 d\beta \frac{\Im[\hat{D}_{FT}^q - \hat{G}_{FT}^q](z_h, \frac{z_h}{\beta})}{1-\beta} \right. \\
 & + \frac{\alpha_s}{2\pi} S_\epsilon \int_{z_h}^1 \frac{dw}{w^2} \int_0^1 d\beta \left\{ \hat{\sigma}_{D_{1T}^{\perp(1)}}^{q;\text{EoM}}(w) D_{1T}^{\perp(1),q}\left(\frac{z_h}{w}\right) \right. \\
 & + \hat{\sigma}_{D_{1T}^{\perp(1)}}^{g;\text{EoM}}(w) D_{1T}^{\perp(1),g}\left(\frac{z_h}{w}\right) + \hat{\sigma}_{H_1^{(1)g}}^{g;\text{EoM}}(w) H_1^{(1)g}\left(\frac{z_h}{w}\right) \\
 & + \hat{\sigma}_{\Im D_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{D}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} + \hat{\sigma}_{\Im G_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{G}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} \\
 & + \hat{\sigma}_1^{gg;\text{EoM}}(w, \beta) \Im[\hat{D}_{FT}^{gg} - \hat{G}_{FT}^{gg} + (1-\epsilon)\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 & + \hat{\sigma}_3^{gg;\text{EoM}}(w, \beta) \Im[(1-\epsilon)\hat{D}_{FT}^{gg} + \hat{G}_{FT}^{gg} + \frac{\epsilon}{2}\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 & + \hat{\sigma}_{D_{FT}}^{q\bar{q};\text{EoM}}(w) \left( \sum_{q=u,d,\dots} \Im[\hat{D}_{FT}^{q\bar{q}}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \right) \\
 & \left. + \hat{\sigma}_{\Re}^q(w, \beta) \frac{\Re[\hat{D}_{FT}^q - \hat{G}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} \right\} + \mathcal{O}(\Lambda^2/s),
 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM

NLO

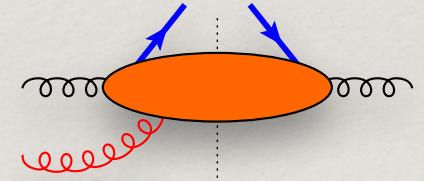
2-gluon correlation w/ EoM

NLO

q-gluon-q correlation w/ EoM

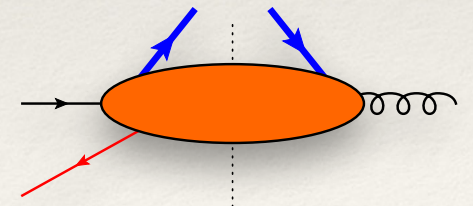
NLO

triple-gluon correlation w/ EoM



NLO

qq-gluon correlation w/ EoM





# Complete structure of the NLO result w/o intrinsic twist-3

$$\begin{aligned}
 E_h \frac{d\sigma_U^{\text{EoM},2}}{d^{d-1}\vec{P}_h}(S_h) &= (4\pi^2 z_h^2)^\epsilon \frac{2\alpha_{\text{em}}^2 N_c}{z_h s^2} \frac{2M_h \epsilon^{P_h m l} S_h}{s} (2v-1) \times \\
 &\sum_{q=u,\bar{u},\dots} e_q^2 \left[ -2 D_{1T}^{\perp(1),q}(z_h) + 2 \int_0^1 d\beta \frac{\Im[\hat{D}_{FT}^q - \hat{G}_{FT}^q](z_h, \frac{z_h}{\beta})}{1-\beta} \right. \\
 &+ \frac{\alpha_s}{2\pi} S_\epsilon \int_{z_h}^1 \frac{dw}{w^2} \int_0^1 d\beta \left\{ \hat{\sigma}_{D_{1T}^{\perp(1)}}^{q;\text{EoM}}(w) D_{1T}^{\perp(1),q}\left(\frac{z_h}{w}\right) \right. \\
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 &+ \hat{\sigma}_{\Im D_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{D}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} + \hat{\sigma}_{\Im G_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{G}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} \\
 &+ \hat{\sigma}_1^{gg;\text{EoM}}(w, \beta) \Im[\hat{D}_{FT}^{gg} - \hat{G}_{FT}^{gg} + (1-\epsilon)\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 &+ \hat{\sigma}_3^{gg;\text{EoM}}(w, \beta) \Im[(1-\epsilon)\hat{D}_{FT}^{gg} + \hat{G}_{FT}^{gg} + \frac{\epsilon}{2}\hat{H}_{FT}^{gg}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \\
 &+ \hat{\sigma}_{D_{FT}}^{q\bar{q};\text{EoM}}(w) \left( \sum_{q=u,d,\dots} \Im[\hat{D}_{FT}^{q\bar{q}}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \right) \\
 &\left. + \hat{\sigma}_{\Re}^q(w, \beta) \frac{\Re[\hat{D}_{FT}^q - \hat{G}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} \right\} + \mathcal{O}(\Lambda^2/s),
 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM

NLO

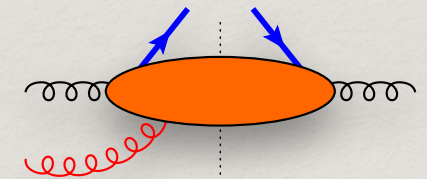
2-gluon correlation w/ EoM

NLO

q-gluon-q correlation w/ EoM

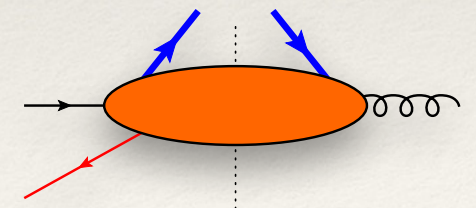
NLO

triple-gluon correlation w/ EoM



NLO

qq-gluon correlation w/ EoM



NLO

imaginary parts of loops



# Complete structure of the NLO result w/o intrinsic twist-3

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 & + \hat{\sigma}_{D_{1T}^{\perp(1)}}^{g;\text{EoM}}(w) D_{1T}^{\perp(1),g}\left(\frac{z_h}{w}\right) + \hat{\sigma}_{H_1^{(1)g}}^{g;\text{EoM}}(w) H_1^{(1)g}\left(\frac{z_h}{w}\right) \\
 & + \hat{\sigma}_{\Im D_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{D}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} + \hat{\sigma}_{\Im G_{FT}}^{qg;\text{EoM}}(w, \beta) \frac{\Im[\hat{G}_{FT}^q]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right)}{1-\beta} \\
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 & + \hat{\sigma}_{D_{FT}}^{q\bar{q};\text{EoM}}(w) \left( \sum_{q=u,d,\dots} \Im[\hat{D}_{FT}^{q\bar{q}}]\left(\frac{z_h}{w}, \frac{z_h}{w\beta}\right) \right) \\
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 \end{aligned}$$

LO

NLO

2-quark correlation w/ EoM

NLO

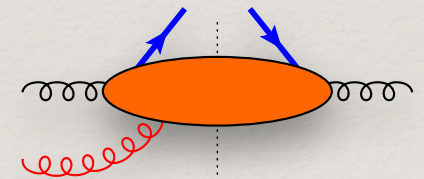
2-gluon correlation w/ EoM

NLO

q-gluon-q correlation w/ EoM

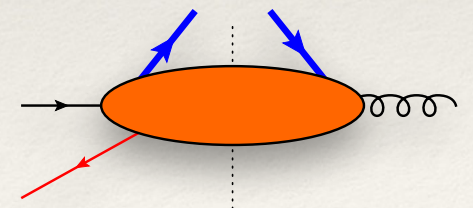
NLO

triple-gluon correlation w/ EoM



NLO

qq-gluon correlation w/ EoM



NLO

imaginary parts of loops

All partonic factors calculated in  
Feynman gauge & Light-cone gauge,  
both calculations *agree!*



*If we assume that twist-3 factorization holds...*



*If we assume that twist-3 factorization holds...*

read off evolution equations from collinear divergences for quark twist-3 FF  $D_T(z)$

$$\begin{aligned} \frac{\partial}{\partial \ln \mu^2} \left( D_T^f(z; \mu) \right) = & \frac{z}{2} \int_z^1 \frac{dw}{w^2} \int_0^1 d\beta \left[ P_{1,f \rightarrow f}^{[1]}(w) D_{1T}^{\perp(1),f} \left( \frac{z}{w}; \mu \right) + P_{1,f \rightarrow g}^{[1]}(w) D_{1T}^{\perp(1),g} \left( \frac{z}{w}; \mu \right) \right. \\ & + P_{2,f \rightarrow fg}^{[1]}(w, \beta) \frac{\Im[\hat{D}_{FT}^{fg} - \hat{G}_{FT}^{fg}](\frac{z}{w}, \beta; \mu)}{1 - \beta} + P_{3,f \rightarrow fg}^{[1]}(w, \beta) \frac{2 \Im[\hat{D}_{FT}^{fg}](\frac{z}{w}, \beta; \mu)}{(1 - \beta)^2} \\ & + \sum_{f'=q', \bar{q}'} P_{4,f \rightarrow f' \bar{f}'}^{[1]}(w, \beta) \Im[\hat{D}_{FT}^{f' \bar{f}'}](\frac{z}{w}, \beta; \mu) + \sum_{f'=q', \bar{q}'} P_{5,f \rightarrow f' \bar{f}'}^{[1]}(w, \beta) \Im[\hat{G}_{FT}^{f' \bar{f}'}](\frac{z}{w}, \beta; \mu) \\ & \left. + P_{6,f \rightarrow gg}^{[1]}(w, \beta) \frac{\Im[\hat{N}_2^s(\frac{z}{w}, \beta; \mu)]}{\beta^2(1 - \beta)^2} + P_{7,f \rightarrow gg}^{[1]}(w, \beta) \frac{\Im[\hat{N}_2^a(\frac{z}{w}, \beta; \mu)]}{\beta^2(1 - \beta)^2} + P_{8,f \rightarrow gg}^{[1]}(w, \beta) \frac{\Im[\hat{N}_1(\frac{z}{w}, \beta; \mu)]}{\beta^2(1 - \beta)^2} \right] \end{aligned}$$

ordinary DGLAP splitting functions

$$P_{1,f \rightarrow f}^{[1]}(w) = -2 \frac{C_F \alpha_s}{2\pi} \left( \frac{1+w^2}{(1-w)_+} + \frac{3}{2} \delta(1-w) \right)$$

$$P_{1,f \rightarrow g}^{[1]}(w) = 4 \frac{C_F \alpha_s}{2\pi} \left( \frac{1+(1-w)^2}{w} \right)$$

Others: more complicated



*If we assume that twist-3 factorization holds...*

read off evolution equations from collinear divergences for quark twist-3 FF  $D_T(z)$

$$\begin{aligned} \frac{\partial}{\partial \ln \mu^2} \left( D_T^f(z; \mu) \right) = & \frac{z}{2} \int_z^1 \frac{dw}{w^2} \int_0^1 d\beta \left[ P_{1,f \rightarrow f}^{[1]}(w) D_{1T}^{\perp(1),f} \left( \frac{z}{w}; \mu \right) + P_{1,f \rightarrow g}^{[1]}(w) D_{1T}^{\perp(1),g} \left( \frac{z}{w}; \mu \right) \right. \\ & + P_{2,f \rightarrow fg}^{[1]}(w, \beta) \frac{\Im[\hat{D}_{FT}^{fg} - \hat{G}_{FT}^{fg}](\frac{z}{w}, \beta; \mu)}{1 - \beta} + P_{3,f \rightarrow fg}^{[1]}(w, \beta) \frac{2 \Im[\hat{D}_{FT}^{fg}](\frac{z}{w}, \beta; \mu)}{(1 - \beta)^2} \\ & + \sum_{f'=q', \bar{q}'} P_{4,f \rightarrow f' \bar{f}'}^{[1]}(w, \beta) \Im[\hat{D}_{FT}^{f' \bar{f}'}](\frac{z}{w}, \beta; \mu) + \sum_{f'=q', \bar{q}'} P_{5,f \rightarrow f' \bar{f}'}^{[1]}(w, \beta) \Im[\hat{G}_{FT}^{f' \bar{f}'}](\frac{z}{w}, \beta; \mu) \\ & \left. + P_{6,f \rightarrow gg}^{[1]}(w, \beta) \frac{\Im[\hat{N}_2^s(\frac{z}{w}, \beta; \mu)]}{\beta^2(1 - \beta)^2} + P_{7,f \rightarrow gg}^{[1]}(w, \beta) \frac{\Im[\hat{N}_2^a(\frac{z}{w}, \beta; \mu)]}{\beta^2(1 - \beta)^2} + P_{8,f \rightarrow gg}^{[1]}(w, \beta) \frac{\Im[\hat{N}_1(\frac{z}{w}, \beta; \mu)]}{\beta^2(1 - \beta)^2} \right] \end{aligned}$$

ordinary DGLAP splitting functions

$$P_{1,f \rightarrow f}^{[1]}(w) = -2 \frac{C_F \alpha_s}{2\pi} \left( \frac{1 + w^2}{(1 - w)_+} + \frac{3}{2} \delta(1 - w) \right)$$

$$P_{1,f \rightarrow g}^{[1]}(w) = 4 \frac{C_F \alpha_s}{2\pi} \left( \frac{1 + (1 - w)^2}{w} \right)$$

Others: more complicated

Final proof of one-loop factorization:

Need to derive evolution equation directly from correlator!

Previous work on *unpolarized* chiral-odd twist-3 fragmentation:

[Belitsky, Kuraev, NPB 1996; Ma, Zhang, PLB 2017]

“The Gribov-Lipatov reciprocity fulfilled for two-particle cut-vertices only!”



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# Summary & Outlook

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- ❖  $\Lambda$  Polarization: Long history, measured in pp-collisions, recently at ATLAS  $\rightarrow$  feasible at a high-energy collider
- ❖ Recent measurement at Belle in  $e^+e^-$ : clean processes to determine polarized  $\Lambda$  fragmentation functions
- ❖ Theory for  $e^+e^-$ : Transverse  $\Lambda$  single-spin asymmetry through (LO)  $D_T$ , consequence of missing T-reversal  $\rightarrow$  unique feature
- ❖ Outlook/Implication: NLO completed,
  - $\rightarrow$  calculate 'splitting functions' for polarized  $\Lambda$  fragmentation function
  - $\rightarrow$  Double Spin Asymmetry  $A_{LT}$  equally important...

Then: extension to other processes, e.g.  $e^+e^-$  to be studied ( $\Lambda+\pi$  - final state)