

#### TRANSVERSITY DISTRIBUTION AND ITS EXTRACTION



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## the "silver" measurement



Deliverables	Observables	What we learn
Sivers &	SIDIS with	Quantum Interference & Spin-Orbital correlations
unpolarized	Transverse	3D Imaging of quark's motion: valence $+$ sea
TMD quarks	polarization;	3D Imaging of gluon's motion
and gluon	di-hadron (di-jet)	QCD dynamics in a unprecedented $Q^2$ ( $P_{hT}$ ) range
Chiral-odd	SIDIS with	$3^{\rm rd}$ basic quark PDF: valence + sea, tensor charge
functions:	Transverse	Novel spin-dependent hadronization effect
Transversity;	polarization	QCD dynamics in a chiral-odd sector
Boer-Mulders		with a wide $Q^2 (P_{hT})$ coverage

Table 2.2: Science Matrix for TMD: 3D structure in transverse momentum space: (upper) the golden measurements; (lower) the silver measurements.

Accardi et al., E.P.J. A52 (16) 268

#### why transversity ?

## the leading-twist PDF/TMD map



1- h1 needed as the 3rd basic quark PDF for spin-1/2 objects

2- address novel QCD dynamics in the chiral-odd sector, also as TMD

Moreover, tensor charge not associated to conserved current in  $\mathcal{L}_{QCD}$  $\delta q(Q^2) = \int_0^1 dx \left[ h_1^q(x,Q^2) - h_1^{\bar{q}}(x,Q^2) \right]$ 

# potential for BSM discovery ?



## **Examples of indirect access**

 nuclear β-decay: effective field theory including operators not in SM Lagrangian; for example, tensor operator



- **neutron EDM**: estimate CPV induced by quark chromo-EDM  $d_q$ 



# **Examples of direct access**

-  $\mathbf{p} \mathbf{p} \rightarrow \mathbf{e}^- \mathbf{v} + \mathbf{X}$  search for W'  $\rightarrow \mathbf{e}^- \mathbf{v}$  with W' heavy partner of W

 $M_{W'} > 5.1-5.2$  TeV at 95% C.L.

puts contraints on BSM operators including tensor operator

see Gupta et al. (PNDME), P.R. D98 (18) 034503



Aaboud et al. (ATLAS), E.P.J. C78 (18) 401



# extraction of transversity

transversity is chiral-odd  $\rightarrow$  need a chiral-odd partner



- hadron-in-jet mechanism : mixed framework h1 as PDF

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$$A_{\text{SIDIS}}^{\sin(\phi_R + \phi_S)}(x, z, M_h^2) \sim -\frac{\sum_q e_q^2 h_1^q(x) \frac{|\mathbf{R}_T|}{M_h} H_{1,q}^{\triangleleft}(z, M_h^2)}{\sum_q e_q^2 f_1^q(x) D_{1,q}(z, M_h^2)}$$

collinear framework  $\rightarrow$  - simple product of PDF and IFF

- x-dependence of AsiDis all in PDF
- flavor sum simplified by symmetries of IFF

 $\begin{array}{cccc} \pi^{+}\pi^{-} & H_{1}^{\triangleleft u} & = & -H_{1}^{\triangleleft d} & \text{isospin symmetry} \\ \text{tree level} & H_{1}^{\triangleleft q} & = & -H_{1}^{\triangleleft \overline{q}} \\ & & & & \\ & & & & \\ & & & & D_{1}^{q} & = & D_{1}^{\overline{q}} \end{array} \right\} \text{charge conjugation}$ 

$$A_{\text{SIDIS}}^{\sin(\phi_R + \phi_S)}(x, z, M_h^2) \sim -\frac{\sum_q e_q^2 h_1^q(x) \frac{|\mathbf{R}_T|}{M_h} H_{1,q}^{\triangleleft}(z, M_h^2)}{\sum_q e_q^2 f_1^q(x) D_{1,q}(z, M_h^2)}$$

- collinear framework  $\rightarrow$  simple product of PDF and IFF - x-dependence of A<sub>SIDIS</sub> all in PDF
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- $\begin{array}{cccc} \pi^{+}\pi^{-} & H_{1}^{\triangleleft u} & = & -H_{1}^{\triangleleft d} & \text{isospin symmetry} \\ \text{tree level} & H_{1}^{\triangleleft q} & = & -H_{1}^{\triangleleft \overline{q}} \\ & & & \\ & & & \\ & & & \\ D_{1}^{q} & = & D_{1}^{\overline{q}} \end{array} \right\} \text{charge conjugation} + \begin{array}{c} \text{data on proton} \\ \text{and deuteron targets} \end{array}$

proton 
$$xh_1^{u-\bar{u}} - \frac{1}{4}xh_1^{d-\bar{d}} = F[A_{\text{SIDIS}}^p \text{ data}, H_1^{\triangleleft u}, f_1^q D_1^q]$$
 separate valence  
deuteron  $xh_1^{u-\bar{u}} + xh_1^{d-\bar{d}} = \tilde{F}[A_{\text{SIDIS}}^D \text{ data}, H_1^{\triangleleft u}, f_1^q D_1^q]$  up and down

collinear framework  $\rightarrow$  - factorization theorems for all hard processes - universality of h<sub>1</sub> H<sub>1</sub> $\stackrel{q}{}$  mechanism



collinear framework  $\rightarrow$  - factorization theorems for all hard processes - universality of  $h_1 H_1 \triangleleft$  mechanism



#### data used in the global fit



Airapetian et al., JHEP 0806 (08) 017

Adolph et al., P.L. **B713** (12) Braun et al., E.P.J. Web Conf. 85 (15)



Vossen et al., P.R.L. 107 (11) 072004



run 2006 (s=200)

Adamczyk et al. (STAR), P.R.L. 115 (2015) 242501

#### the phase space



- mostly high  $x \rightarrow$  not enough for sea quark explorations

- guess low-x behavior (relevant for calculation of tensor charge)

#### choice of functional form

functional form whose Mellin transform can be computed analytically and complying with Soffer Bound at any x and scale Q<sup>2</sup>

$$h_1^{q_v}(x;Q_0^2) = F^{q_v}(x) \begin{bmatrix} SB^q(x) + \overline{SB}^{\overline{q}}(x) \end{bmatrix}$$

$$\begin{array}{c} & & \\ & & \\ & & \\ & & \\ Soffer Bound \\ 2|h_1^q(x,Q^2)| \leq 2 SB^q(x,Q^2) = |f_1^q(x,Q^2) + g_1^q(x,Q^2)| \\ & & \\ & & \\ MSTW08 \quad DSSV \end{array}$$

$$(x) = \frac{N_{q_v}}{\max_x[|F^{q_v}(x)|]} x^{A_{q_v}} \left[1 + B_{q_v} \operatorname{Ceb}_1(x) + C_{q_v} \operatorname{Ceb}_2(x) + D_{q_v} \operatorname{Ceb}_3(x)\right] \\ & & \\ \operatorname{Ceb}_n(x) \text{ Cebyshev polynomial} \end{array}$$

10 fitting parameters

constrain parameters

 $F^{q_v}$ 

 $|N_{q_v}| \le 1 \Rightarrow |F^{q_v}(x)| \le 1$  Soffer Bound ok at any Q<sup>2</sup>

#### low-x behavior



#### low-x behavior



2) "massive" jet in DIS  $\rightarrow$  h<sub>1</sub> at twist 3 violation of Burkardt-Cottingham s.r.  $\int_{0}^{1} dx \, g_{2}(x) \propto \int_{0}^{1} dx \, \frac{h_{1}(x)}{x} \longrightarrow A_{q} + a_{q} > 1$ 

3) small-x dipole picture =>  $h_1^{q_v}(x) \stackrel{x \to 0}{\approx} x^{1-2\sqrt{\frac{\alpha_s(Q^2)N_c}{2\pi}}} \longrightarrow \text{at } Q_0 \quad A_q + a_q \sim 1$ *Kovchegov & Sievert, arXiv:1808.10354* 

#### low-x behavior

$$\lim_{x \to 0} x SB^{q}(x) \propto x^{a_{q}} \\ \lim_{x \to 0} F^{q_{v}}(x) \propto x^{A_{q}} \\ h_{1}^{q}(x) \stackrel{x \to 0}{\approx} x^{A_{q}} + a_{q} - 1 \\ \text{tensor charge} \quad \delta q(Q^{2}) = \int_{x_{\min}}^{1} dx h_{1}^{q-\bar{q}}(x, Q^{2}) \\ \text{constrain parameters} \\ \text{low-x behavior important} \\ \delta q \quad \text{finite} => A_{q} + a_{q} > 0 \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{violation of Burkardt-Cottingham s.r.} \int_{0}^{1} dx g_{2}(x) \propto \int_{0}^{1} dx \frac{h_{1}(x)}{x} \longrightarrow A_{q} + a_{q} > 1 \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{ at twist 3} \\ \text{``massive'' jet in DIS} \rightarrow h_{1} \text{``massive'' jet in$$

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1

2

Kovchegov & Sievert, arXiv:1808.10354

**our choice** 
$$A_q + a_q > \frac{1}{3}$$
  $\left| \int_0^{x_{\min}} dx \right| \sim 1\% \text{ of } \left| \int_{x_{\min}}^1 dx \right|$ 

for  $x_{min}=10^{-6}$  from MSTW08

#### theoretical uncertainties

#### unpolarized Di-hadron Fragmentation Function D1

- quark D<sub>1</sub>q is well constrained by  $e^+e^- \rightarrow (\pi^+\pi^-) X$  (Montecarlo)
- **gluon**  $D_1^g$  is **not** constrained by  $e^+e^- \rightarrow (\pi^+\pi^-) X$  (currently, LO analysis)
- **no data** available yet for  $p p \rightarrow (\pi^+\pi^-) X$

#### we don't know anything about the gluon $D_1^g$

our choice: set 
$$D_{I^g}(Q_0) = \begin{cases} 0 \\ D_{I^u}(Q_0) / 4 \\ D_{I^u}(Q_0) \end{cases}$$

deteriorates our e<sup>+</sup>e<sup>-</sup> fit as  $\chi^2/dof =$ 

$$\begin{cases} 1.69 & 1.28 \\ 1.81 & 1.37 \\ 2.96 & 2.01 \end{cases}$$
background  $\rho$  channels

- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50



- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50, 100



- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50, 100, 200 sets...



- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50, 100, 200 sets... until average and standard deviation reproduce original exp. points (here, 200x3=600)



- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50, 100, 200 sets... until average and standard deviation reproduce original exp. points (here, 200x3=600)
- exclude largest and smallest 5% => 90% band



automatically accounts for correlations

#### results

#### global fit published in

Radici and Bacchetta, P.R.L. **120** (18) 192001



## $X^2$ of the fit



global fit **10** parameters











0.05

0.10

Х

0.50

0.01

#### tensor charge



8)	PNDME '18	Gupta et al., P.R. D98 (18) 034503
9)	FTMC '17	Alexandrou et al., P.R. D95 (17) 114514;
<i>J</i> ) <b>LIMC I</b> /		

- E P.R. D96 (17) 099906
- 10) RQCD '14 Bali et al., P.R. D91 (15)
- 11) LHPC '12 Green et al., P.R. D86 (12)

Radici & Bacchetta, P.R.L. <b>120</b> (18) 192001	4)	global fit '17
Kang et al., P.R. D <b>93</b> (16) 014009	5)	"TMD fit" * Q <sup>2</sup> =10
Anselmino et al., P.R. D87 (13) 094019	6)	Torino fit * Q <sup>2</sup> =1
Lin et al., P.R.L. <b>120</b> (18) 152502	7)	JAM fit '17 * Q <sub>0</sub> <sup>2</sup> =2

#### tensor charge



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- **PNDME** '18 Gupta et al., P.R. D98 (18) 034503 8)
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#### **Compass pseudo-data**

#### add to previous set of data a new set of SIDIS pseudo-data for deuteron target



statistical error  $\sim 0.6 \times [\text{error in } 2010 \text{ proton run }]$ <A> = average value of replicas in previous global fit

#### impact of pseudo-data



#### tensor charge



3) global fit + j	pseudoc	lata
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Radici & Bacchetta, P.R.L. <b>120</b> (18) 192001	4)	global fit '17
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K

Anseli

Lin et al., P.R.L. 120 (18) 152502 7) JAM fit '17 \* Q<sub>0</sub><sup>2</sup>=2

- 8) PNDME '18 Gupta et al., P.R. D98 (18) 034503
  9) ETMC '17 Alexandrou et al., P.R. D95 (17) 114514;
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- **11) LHPC '12** Green et al., P.R. D86 (12)

#### tensor charge



Lin et al., P.R.L. 120 (18) 152502 7) JAM fit '17 \*  $Q_0^2=2$ 

# better X<sup>2</sup>



 $\chi^2/dof = 1.32 \pm 0.09$ 

add to SIDIS+pp data + Compass SIDIS pseudo-data constraint to reproduce  $g_T$  from lattice



#### tensor charge



 $Q^2 = 4 \text{ GeV}^2 *$ 

#### 2) global fit + pseudodata + constrain g<sub>T</sub> 3) global fit + pseudodata

Radici & Bacchetta, P.R.L. <b>120</b> (18) 192001	4)	global fit '17
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Green et al., P.R. D86 (12)

11) LHPC '12

#### tensor charge



#### not yet full compatibility

$$Q^2 = 4 \text{ GeV}^2 *$$

8)

9)

#### 2) global fit + pseudodata + constrain g<sub>T</sub> 3) global fit + pseudodata

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## impact of lattice g<sub>T</sub> constraint



 $X^2$ 

$$\chi^2/dof = 1.32 \pm 0.09$$

 $\chi^2/dof = 1.77 \pm 0.19$ 



probability density function for a  $\chi^2$  distribution with 31 and 32 dof, respectively

#### compatibility with lattice

# add to SIDIS+pp data + Compass SIDIS pseudo-data constraint to reproduce from lattice g<sub>T</sub> , δu , δd

 $\overline{g_T}^{latt} = 1.004 \pm 0.057$ 



 $\overline{\delta d}^{\text{latt}} = -0.218 \pm 0.026$ 



#### tensor charge



 $Q^2 = 4 \text{ GeV}^2 *$ 

global fit + pseudodata + constrain g<sub>T</sub>, δu, δd
 global fit + pseudodata + constrain g<sub>T</sub>
 global fit + pseudodata

Radici & Bacchetta, P.R.L. <b>120</b> (18) 192001	4)	global fit '17
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11)	I HPC '12	Green et al PR $D86(12)$

#### tensor charge



 $Q^2 = 4 \text{ GeV}^2 *$ 

#### compatible, but...

global fit + pseudodata + constrain g<sub>T</sub>, δu, δd
 global fit + pseudodata + constrain g<sub>T</sub>
 global fit + pseudodata

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**11) LHPC '12** Green et al., P.R. D86 (12)



probability density function for a  $\chi^2$  distribution with 31 and 34 dof, respectively

#### impact of "full" lattice constraint



#### truncated tensor charge

truncated  $\delta q^{[0.0065, 0.35]}$  Q<sup>2</sup> = 10



1) global fit + pseudodata + constrain  $g_T$ ,  $\delta u$ ,  $\delta d$ 

- 2) global fit + pseudodata + constrain g<sub>T</sub>
  - 3) global fit + pseudodata

4) global fit '17

Radici & Bacchetta, P.R.L. **120** (18) 192001 5) **"TMD fit"** 



# Conclusions

- first global fit of di-hadron inclusive data leading to extraction of transversity as a PDF in collinear framework
- inclusion of STAR p-p<sup>†</sup> data increases precision of up channel; large uncertainty on down due to unconstrained gluon unpolarized di-hadron fragmentation function
- no apparent simultaneous compatibility with lattice for tensor charge in up, down, and isovector channels
- adding Compass SIDIS pseudo-data for deuteron increases precision, particularly for down, but seems to confirm this scenario
- forcing the fit to reproduce lattice isovector tensor charge is not enough to reach simultaneous compatibility;  $\chi^2$  worsens
- it is possible to reach simultaneous compatibility with lattice but  $\chi^2$  worsens even more and probabilistic distribution is very unlikely



# Back-up



# **2-hadron**-inclusive production

framework collinear factorization



# **IFF** symmetries



$$\begin{array}{rcl} H_1^{\triangleleft u} &=& -H_1^{\triangleleft d} & \text{isospin symmetry} \\ H_1^{\triangleleft q} &=& -H_1^{\triangleleft \overline{q}} \\ D_1^q &=& D_1^{\overline{q}} \end{array} \right\} \text{charge conjugation}$$

#### valid only for $(\pi^+\pi^-)$ pairs and at tree level





Pτ





# To do list

 use also other (multi-dimensional) data from STAR run 2011 (s=500) and (later) run 2012 (s=200)





Radici et al., P.R. D94 (16) 034012

- → need data on p+p →  $(\pi\pi) X$  constrains gluon D<sub>1</sub><sup>g</sup>
- refit di-hadron fragmentation functions using new data:
   e<sup>+</sup>e<sup>-</sup> → (ππ) X constrains D<sub>1</sub><sup>q</sup>
   (currently only by Montecarlo)
- Seidl et al., P.R. D**96** (17) 032005
- use COMPASS data on πK and KK channels, and from Λ<sup>†</sup> fragmentation: constrain strange contribution ?
- explore other channels, like inclusive DIS via Jet fragm. funct.'s

#### more constraints on extrapolation



- of course, need more data