

Global Analysis of Transverse-Spin Observables

Daniel Pitonyak

Lebanon Valley College, Annville, PA

Workshop on Transverse Spin and TMDs

Institute for Nuclear Theory

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Outline

- TMD and collinear twist-3 (CT3) functions
- Sivers and Collins effects & A_N in pp collisions
- Toward a global analysis of transverse spin observables
- Summary and outlook

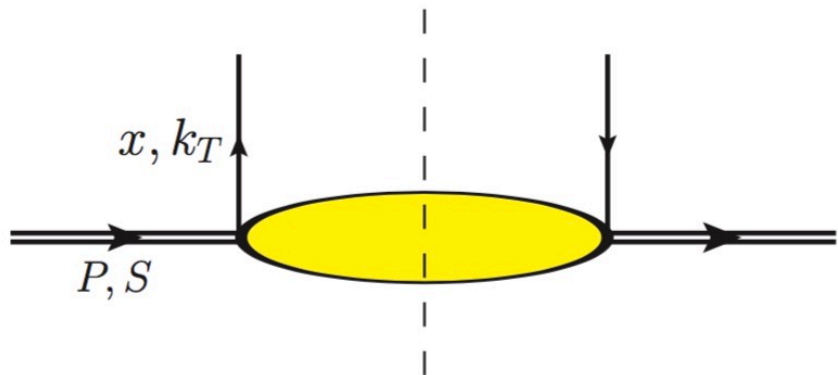


TMD and Collinear Twist-3 Functions

TMD PDFs (x, k_T)

q pol. \ H pol.	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_{1T} h_{1T}^\perp

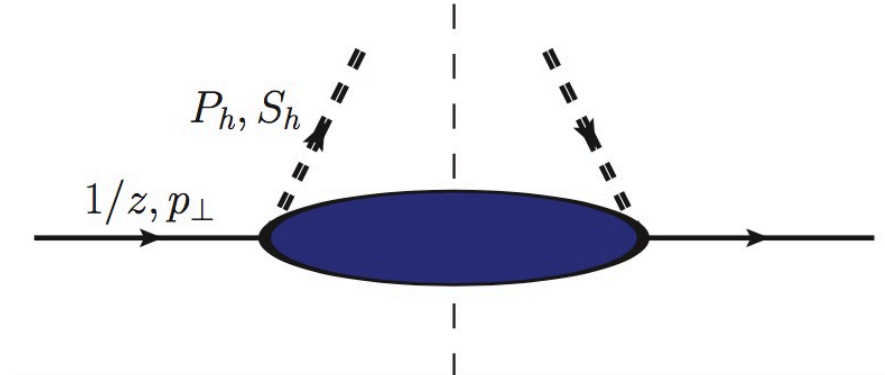
(Mulders, Tangerman (1996); Goeke, Metz, Schlegel (2005))



TMD FFs (z, p_\perp)

q pol. \ H pol.	U	L	T
U	D_1		H_1^\perp
L		G_{1L}	H_{1L}^\perp
T	D_{1T}^\perp	G_{1T}	H_{1T} H_{1T}^\perp

(Boer, Jakob, Mulders (1997))



Collinear PDFs (x)

q pol. \ H pol.	U	L	T
U	f_1 unpolarized		
L		g_1 helicity	
T			h_1 transversity

 Collinear FFs (z)

q pol. \ H pol.	U	L	T
U	D_1		
L		G_1	
T			H_1

Integrate TMDs over k_T (or p_\perp) \rightarrow collinear PDFs and FFs

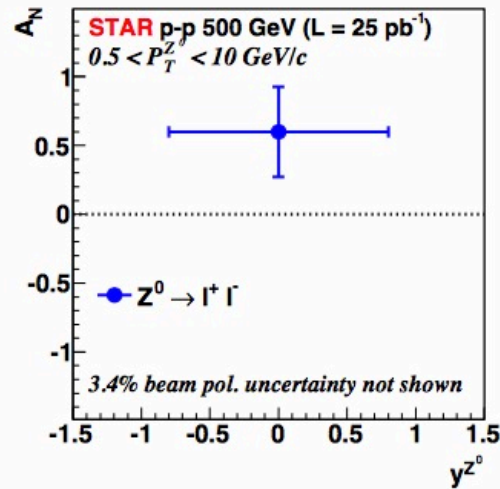
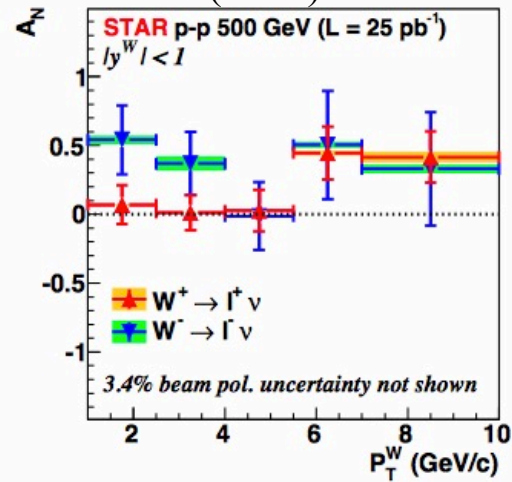
	CT3 PDF (x)		CT3 PDF (x, x_1)	CT3 FF (z)		CT3 FF (z, z_1)
Hadron Pol.						
U	<u>intrinsic</u> e	<u>kinematical</u> $h_1^{\perp(1)}$	<u>dynamical</u> H_{FU}	<u>intrinsic</u> E, H	<u>kinematical</u> $H_1^{\perp(1)}$	<u>dynamical</u> $\hat{H}_{FU}^{\mathcal{R}, \mathcal{S}}$
L	h_L	$h_{1L}^{\perp(1)}$	H_{FL}	H_L, E_L	$H_{1L}^{\perp(1)}$	$\hat{H}_{FL}^{\mathcal{R}, \mathcal{S}}$
T	g_T	$f_{1T}^{\perp(1)},$ $g_{1T}^{\perp(1)}$	F_{FT}, G_{FT}	D_T, G_T	$D_{1T}^{\perp(1)},$ $G_{1T}^{\perp(1)}$	$\hat{D}_{FT}^{\mathcal{R}, \mathcal{S}}, \hat{G}_{FT}^{\mathcal{R}, \mathcal{S}}$



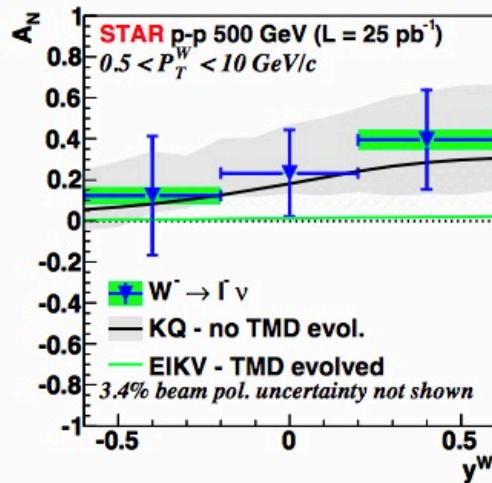
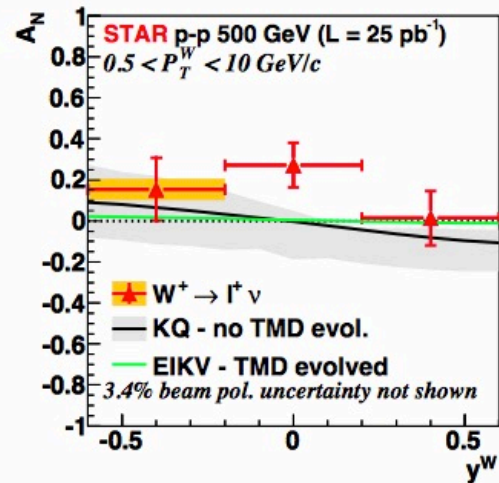
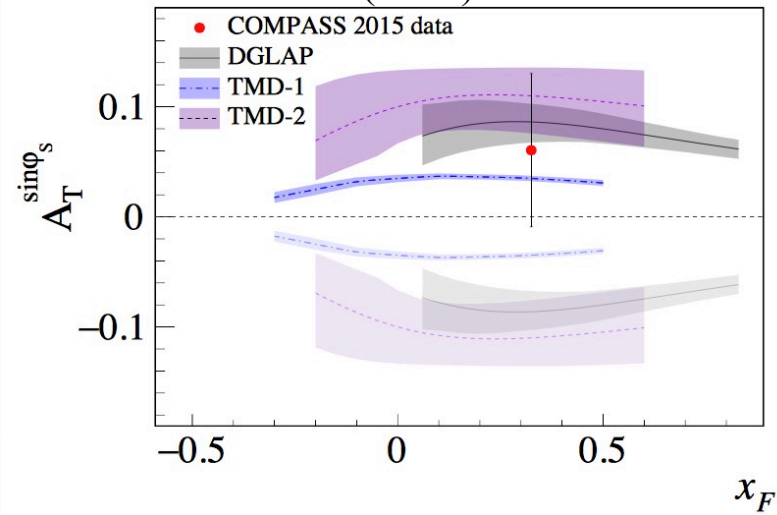
Sivers and Collins Effects & A_N in pp Collisions

Drell-Yan Sivers effect

STAR (2016)

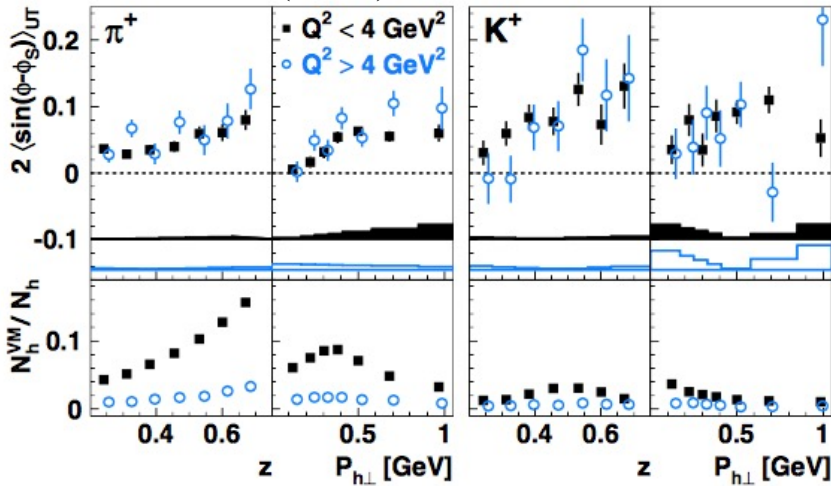


COMPASS (2017)

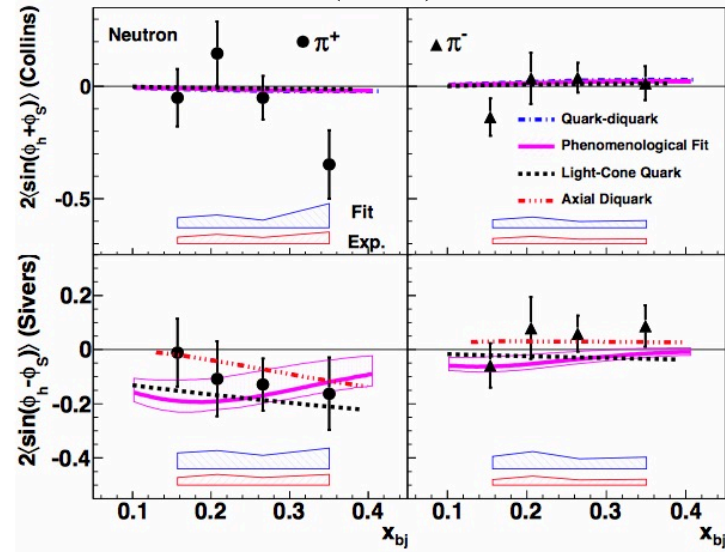


SIDIS Siverts effect ($\sin(\phi_h - \phi_s)$)

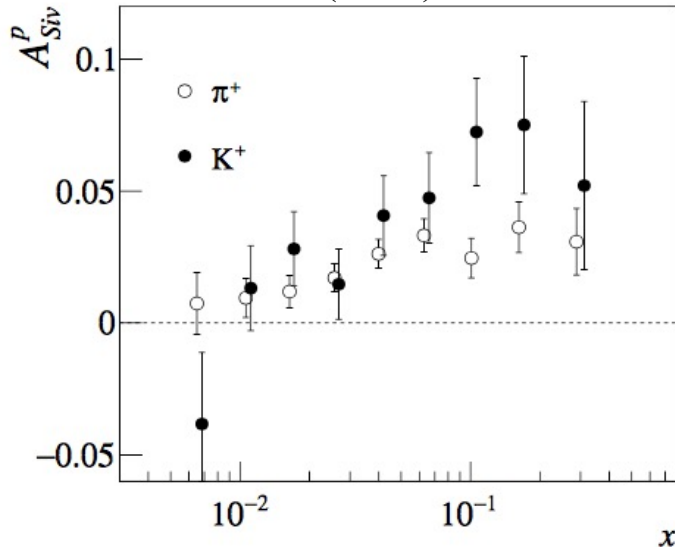
HERMES (2009)



JLab, Hall A (2011)

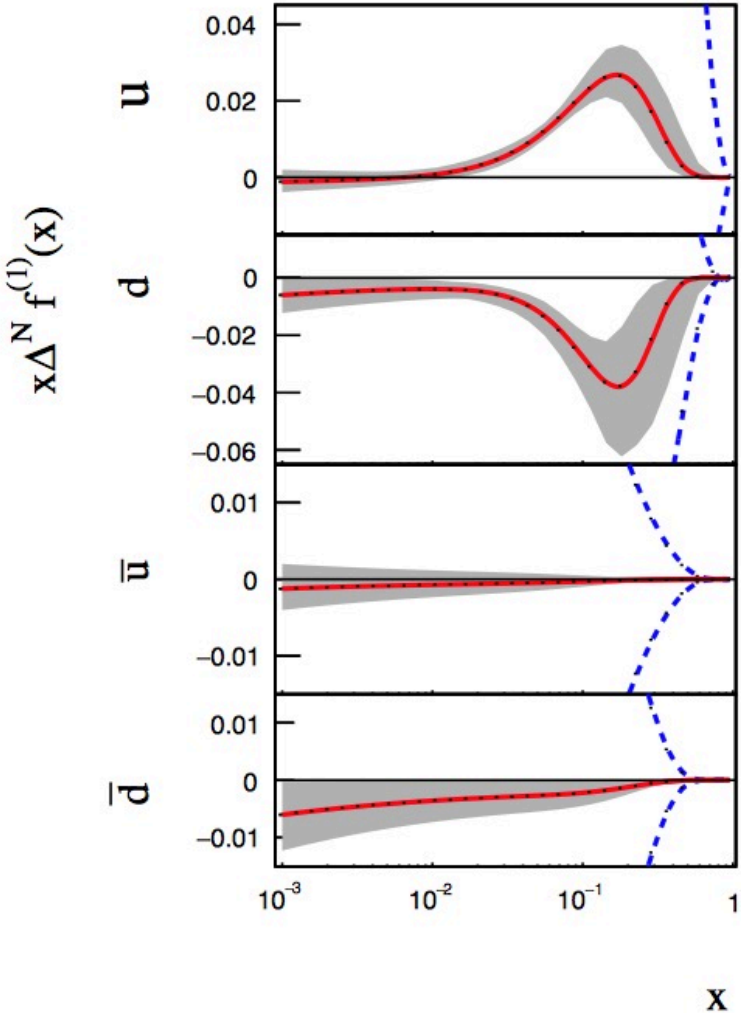


COMPASS (2015)

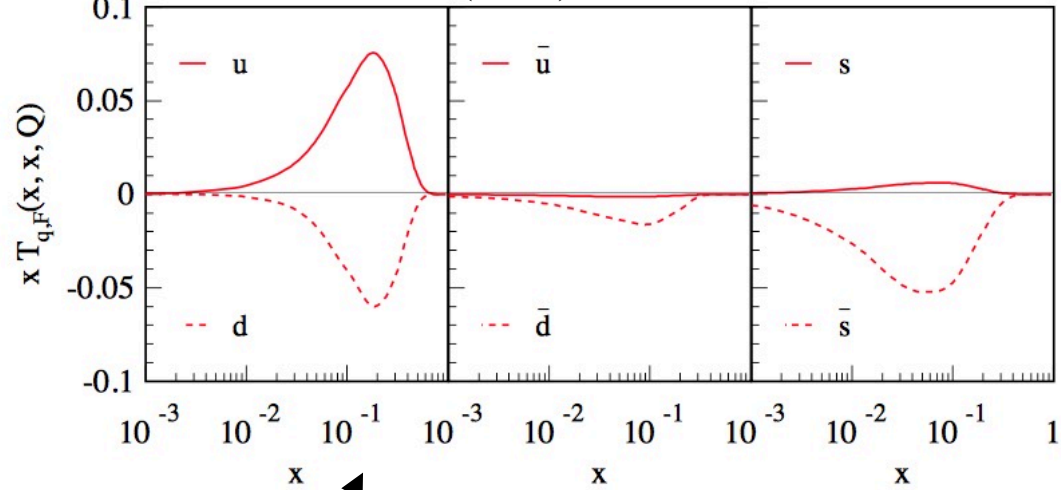


$$F_{UT}^{\sin(\phi_h - \phi_s)} = \mathcal{C} \left[-\frac{\hat{h} \cdot \vec{k}_T}{M} f_{1T}^\perp D_1 \right]$$

Anselmino, et al. (2017)

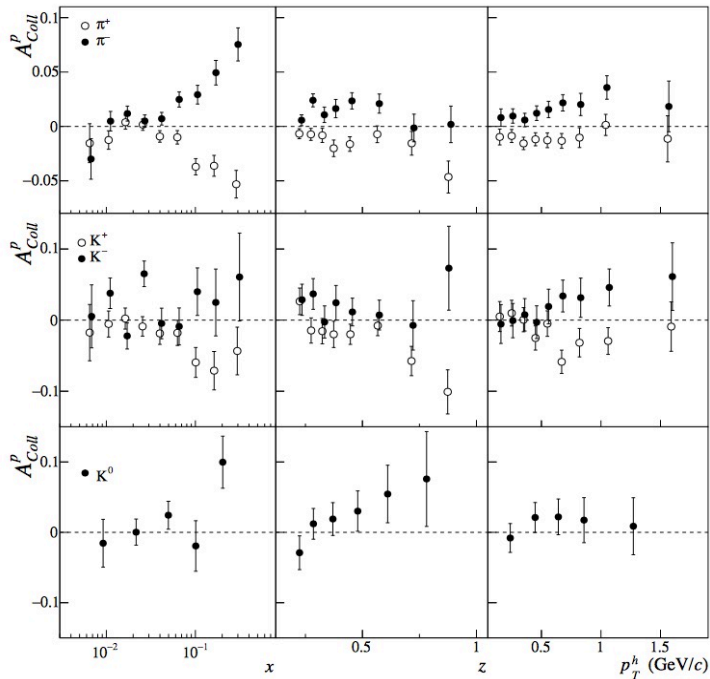


Echevarria, et al. (2014)



TMDs in Collins-Soper-Sterman (CSS) evolution formalism

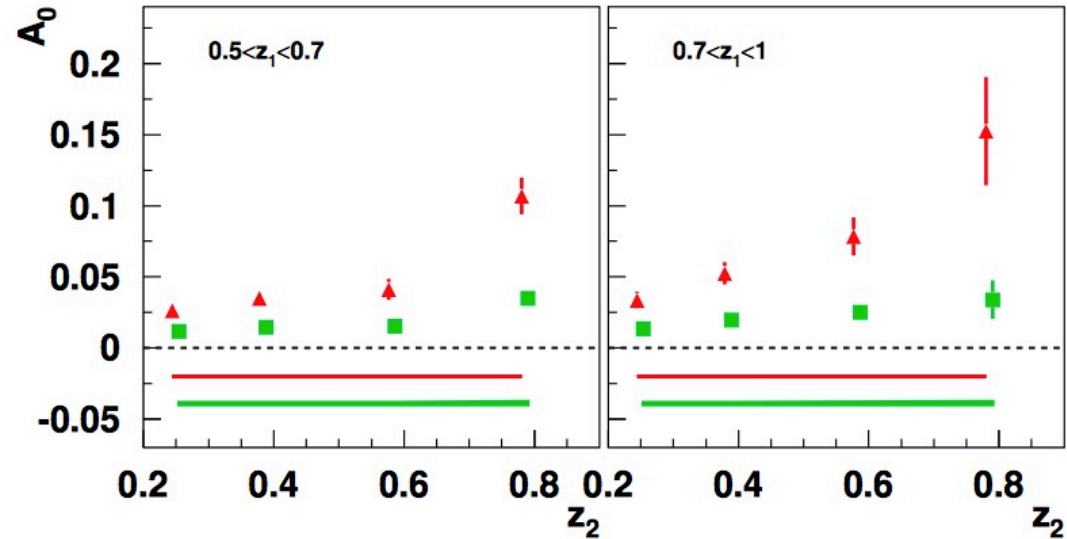
SIDIS Collins effect ($\sin(\phi_h + \phi_s)$)
COMPASS (2015)



Also data from JLab Hall A (2011, 2014) and HERMES

$$F_{UT}^{\sin(\phi_h + \phi_s)} = C \left[-\frac{\hat{h} \cdot \vec{p}_\perp}{M_h} h_1 H_1^\perp \right]$$

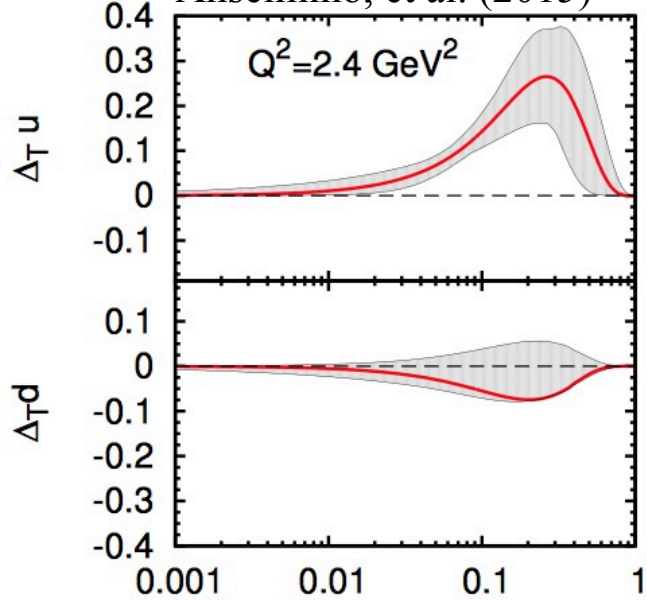
e^+e^- Collins effect ($\cos(2\phi_0)$)
Belle (2008)



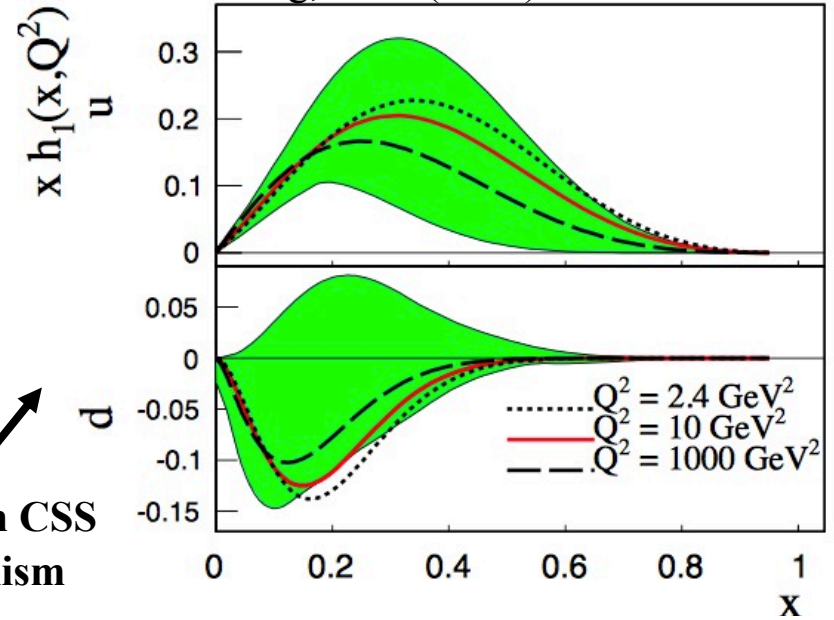
Also data from BaBar (2014) and BESIII (2016)

$$F_{UU}^{\cos(2\phi_0)} = C \left[\frac{2\hat{h} \cdot \vec{p}_{a\perp} \hat{h} \cdot \vec{p}_{b\perp} - \vec{p}_{a\perp} \cdot \vec{p}_{b\perp}}{M_a M_b} H_1^\perp \bar{H}_1^\perp \right]$$

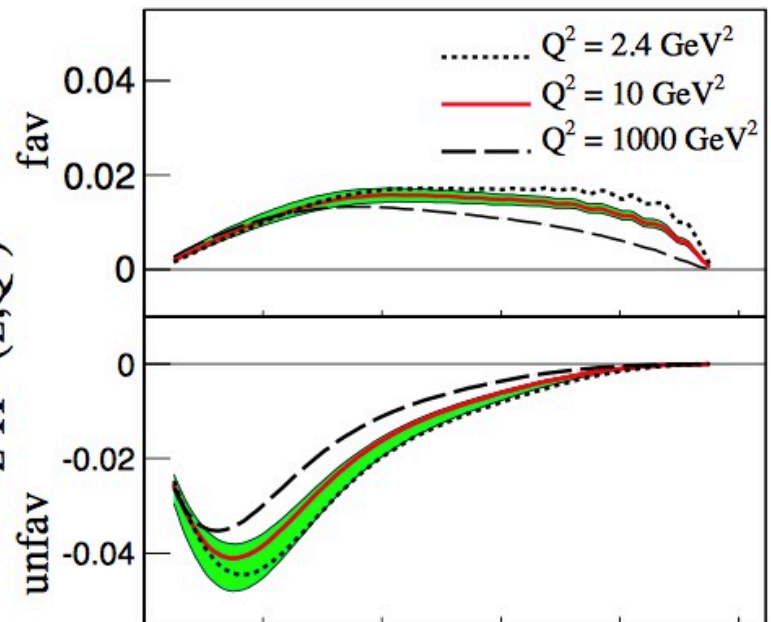
Anselmino, et al. (2015)



Kang, et al. (2016)

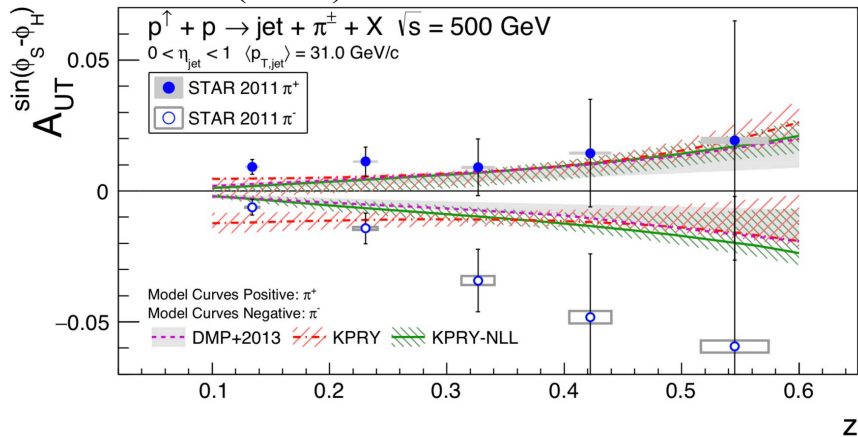


TMDs in CSS formalism



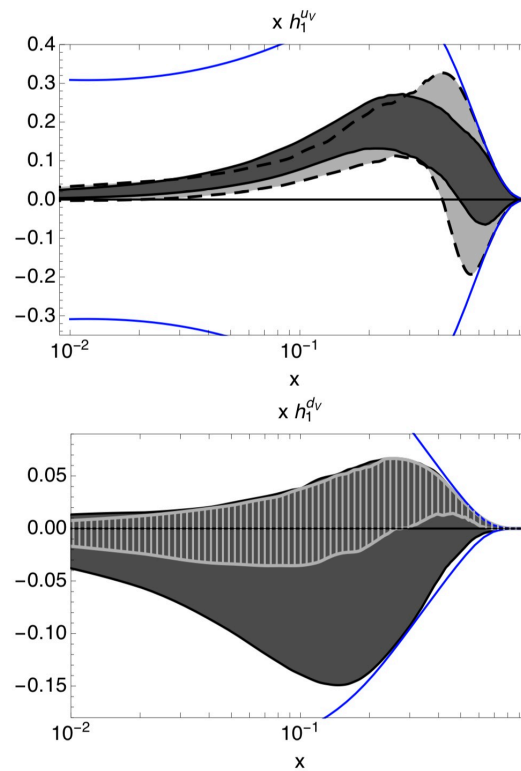
Hadron in a jet Collins effect

STAR (2017)



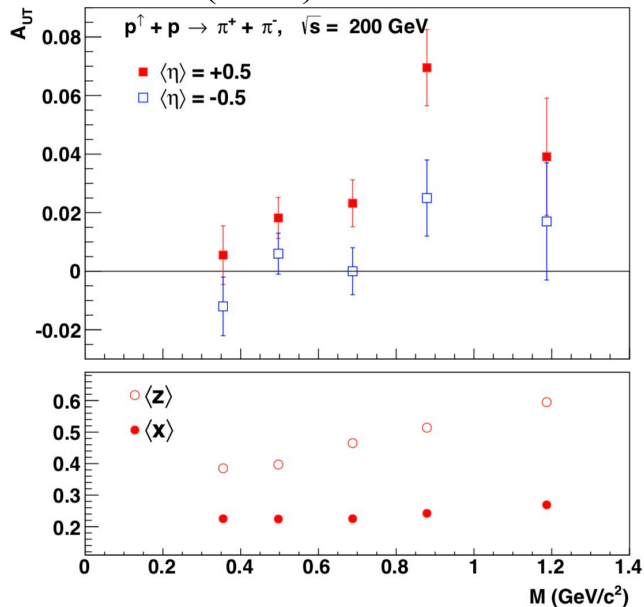
Theory curves from
 D'Alesio, et al. (2017)
 & Kang, et al. (2017)

Transversity from dihadron FF

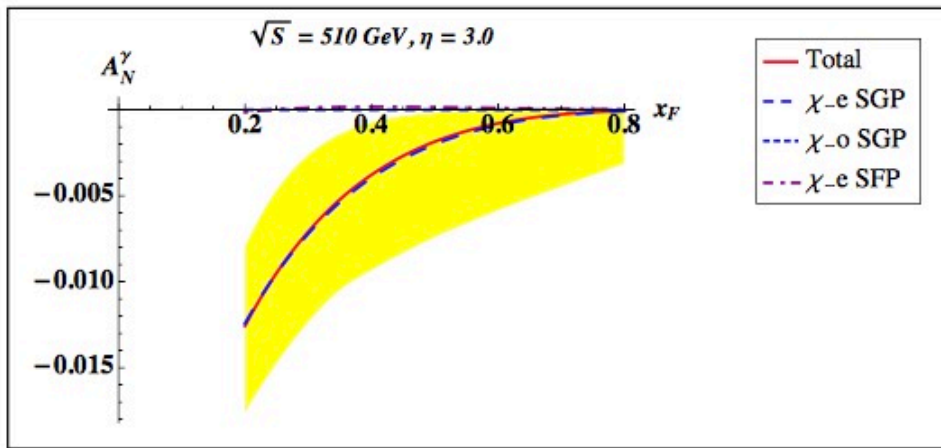
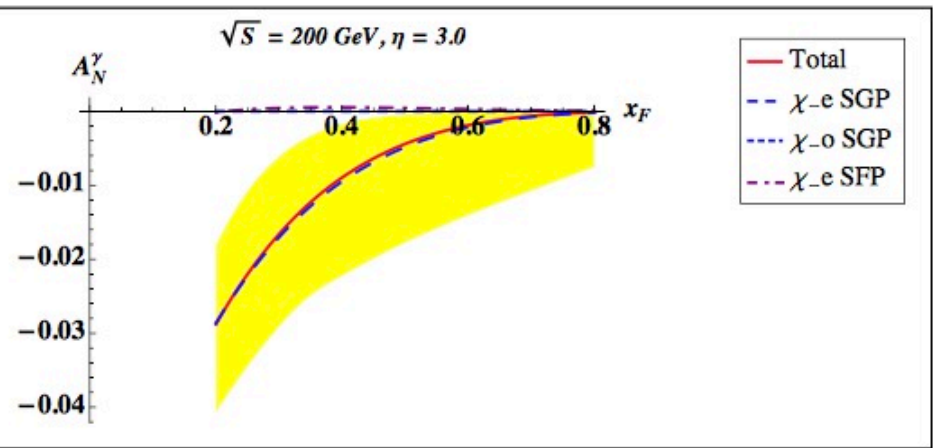


Bacchetta and
 Radici (2017)

STAR (2015)



A_N in $pp \rightarrow \gamma X$



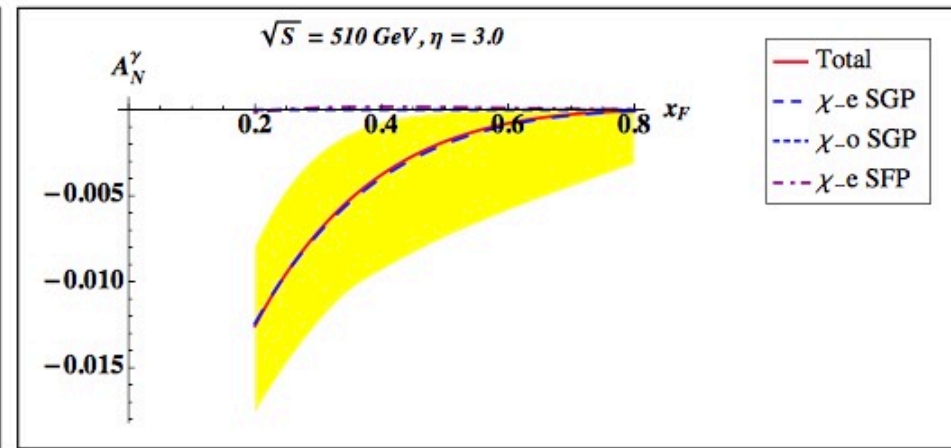
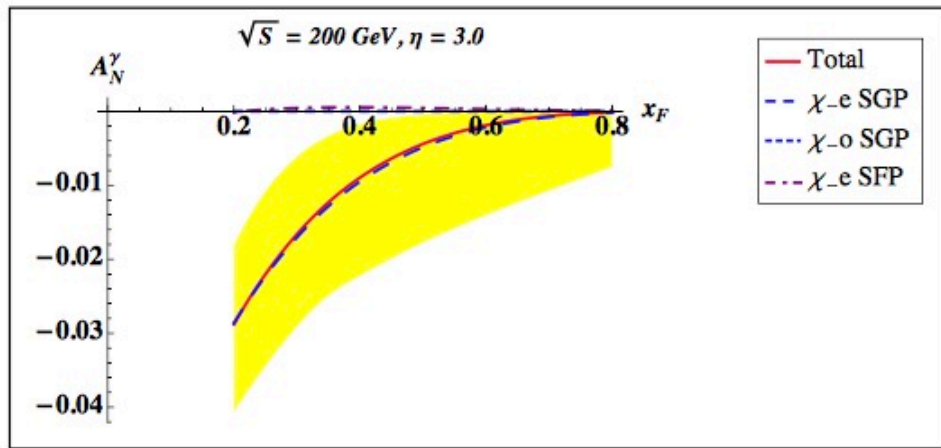
(Kanazawa, Koike, Metz, DP – PRD **91** (2015))
 (See also Gamberg, Kang, Prokudin (2013))

Qiu-Sterman term is the main cause of A_N in $pp \rightarrow \gamma X$ \rightarrow Test of the process dependence of the Sivers function

$$d\Delta\sigma^\pi \sim H \otimes f_1 \otimes F_{FT}(x, x)$$

\swarrow
 Qiu-Sterman function

A_N in $pp \rightarrow \gamma X$



(Kanazawa, Koike, Metz, DP – PRD **91** (2015))

(See also Gamberg, Kang, Prokudin (2013))

Qiu-Sterman term is the main cause of A_N in $pp \rightarrow \gamma X$

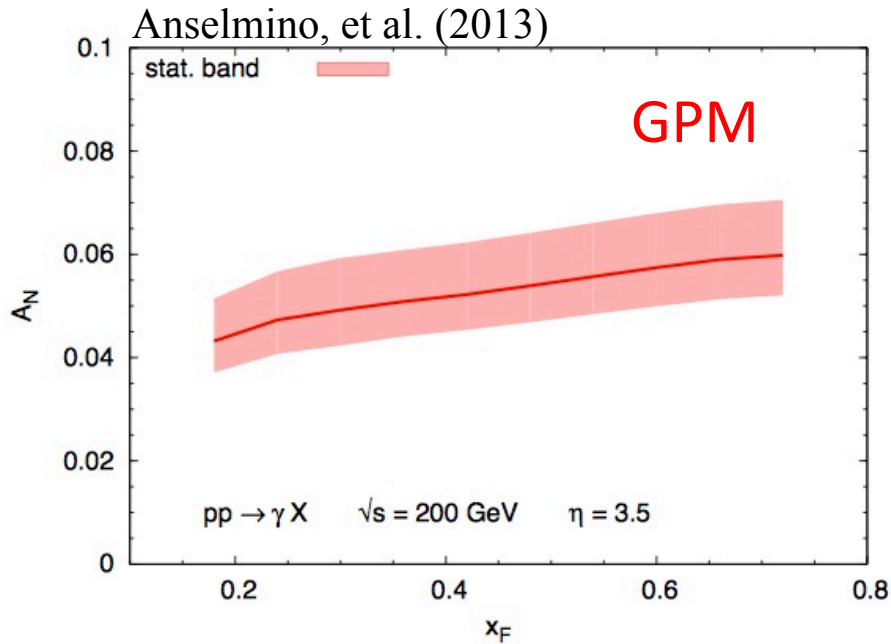


Test of the process dependence of the Sivers function

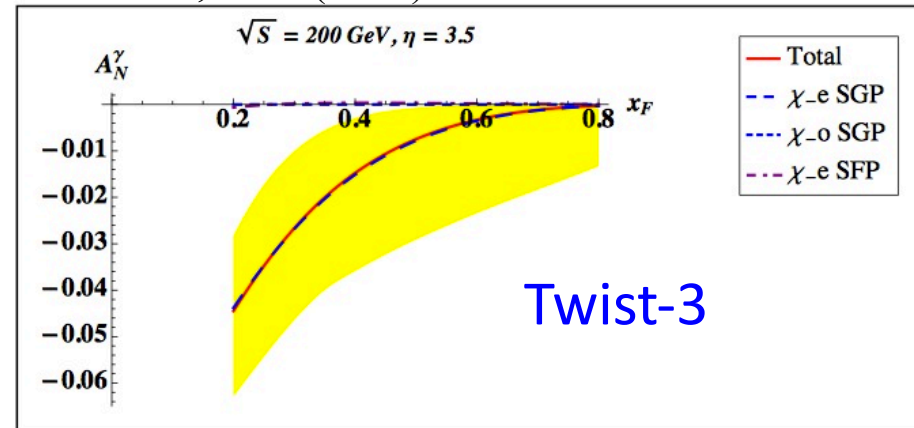
$$d\Delta\sigma^\pi \sim H \otimes f_1 \otimes F_{FT}(x, x)$$

Can direct photon observables like A_N in $ep \rightarrow \gamma X$ be measured at an EIC?

A_N in $pp \rightarrow \gamma X$

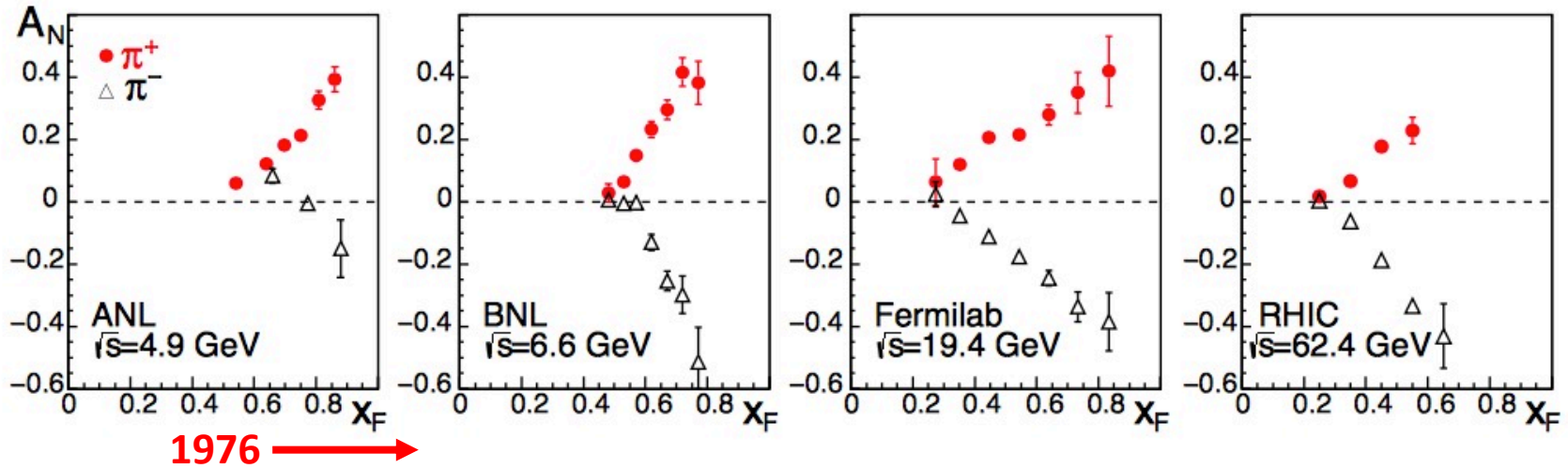


Kanazawa, et al. (2015)



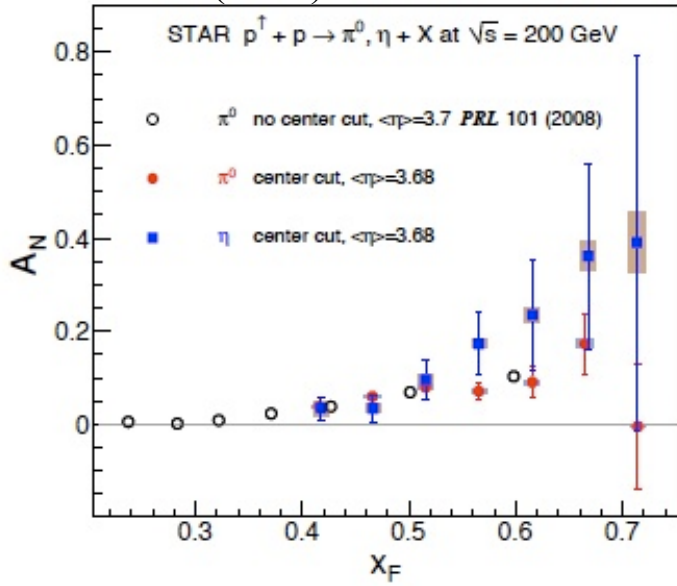
GPM predicts a **positive** asymmetry while **twist-3** predicts a **negative** one

A_N in $pp \rightarrow \pi X$ – PUZZLE FOR 40+ YEARS!

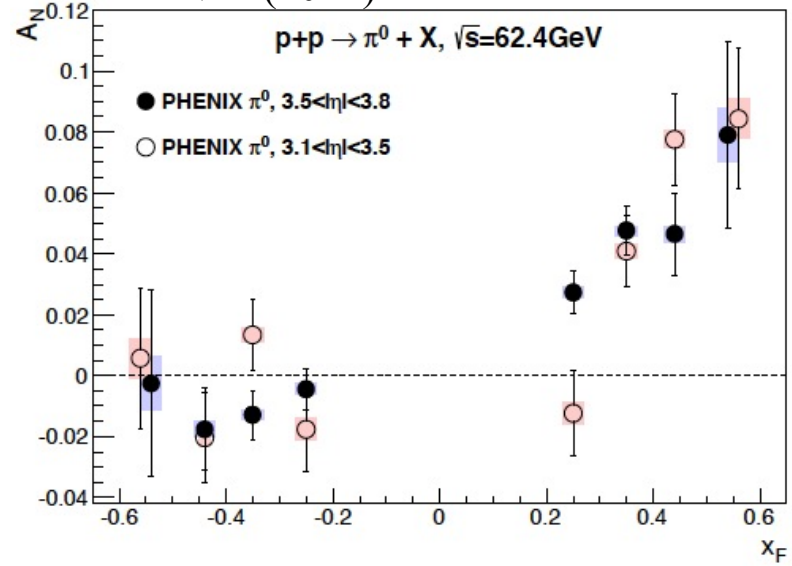


A_N in $pp \rightarrow \pi X$ – PUZZLE FOR 40+ YEARS!

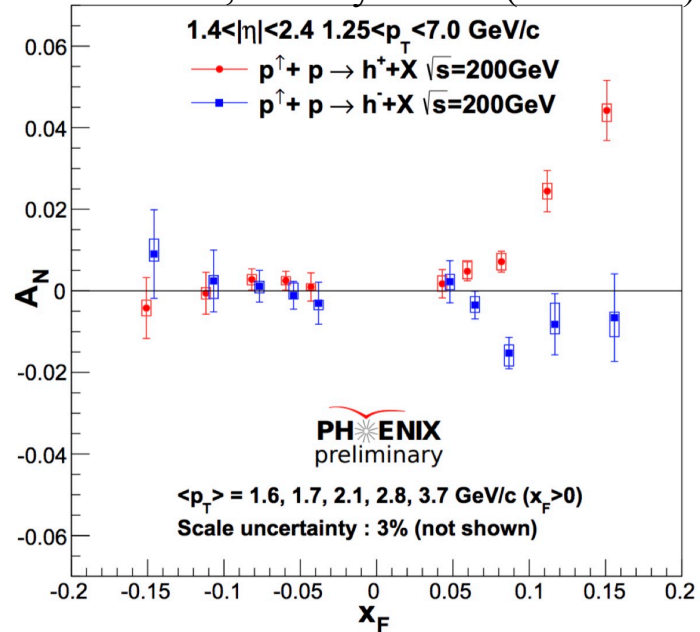
STAR (2012)



PHENIX (2014)



PHENIX, Talk by J. Bok (DIS 2018)



$$d\Delta\sigma^\pi \sim H \otimes f_1 \otimes \mathbf{F}_{FT}(\mathbf{x}, \mathbf{x})$$

$$E_\ell \frac{d^3\Delta\sigma(\vec{s}_T)}{d^3\ell} = \frac{\alpha_s^2}{S} \sum_{a,b,c} \int_{z_{\min}}^1 \frac{dz}{z^2} D_{c \rightarrow h}(z) \int_{x'_{\min}}^1 \frac{dx'}{x'} \frac{1}{x'S + T/z} \phi_{b/B}(x')$$

$$\times \sqrt{4\pi\alpha_s} \left(\frac{\epsilon^{\ell s_T n \bar{n}}}{z \hat{u}} \right) \frac{1}{x} \left[T_{a,F}(x, x) - x \left(\frac{d}{dx} T_{a,F}(x, x) \right) \right] H_{ab \rightarrow c}(\hat{s}, \hat{t}, \hat{u})$$

$$\boxed{F_{FT} \sim T_F}$$

(Qiu and Sterman (1999), Kouvaris, et al. (2006))

For many years the Qiu-Sterman/Sivers-type contribution was thought to be the dominant source of TSSAs in $p^\uparrow p \rightarrow \pi X$

~~$$d\Delta\sigma^\pi \sim H \otimes f_1 \otimes F_{FT}(x, x)$$~~

(Kang, Qiu, Vogelsang, Yuan (2011); Kang and Prokudin (2012);
 Metz, DP, Schäfer, Schlegel, Vogelsang, Zhou (2012))

~~$$d\Delta\sigma^\pi \sim H \otimes f_1 \otimes F_{FT}(x, x)$$~~

$$d\Delta\sigma^\pi \sim h_1 \otimes S \otimes \left(H_1^{\perp(1)}, H, \int \frac{dz_1}{z_1^2} \frac{\hat{H}_{FU}^{\mathfrak{S}}}{(1/z - 1/z_1)^2} \right)$$

$$E_h \frac{d\Delta\sigma^{Frag}(S_T)}{d^3\vec{P}_h} = - \frac{4\alpha_s^2 M_h}{S} \epsilon^{P' P P_h S_T} \sum_i \sum_{a,b,c} \int_0^1 \frac{dz}{z^3} \int_0^1 dx' \int_0^1 dx \delta(\hat{s} + \hat{t} + \hat{u}) \frac{1}{\hat{s}(-x'\hat{t} - x\hat{u})}$$

$$\times h_1^a(x) f_1^b(x') \left\{ \left[H_1^{\perp(1),c}(z) - z \frac{dH_1^{\perp(1),c}(z)}{dz} \right] S_{H_1^\perp}^i + \frac{1}{z} H^c(z) S_H^i \right.$$

$$\left. + \frac{2}{z} \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\left(\frac{1}{z} - \frac{1}{z_1}\right)^2} \hat{H}_{FU}^{c,\mathfrak{S}}(z, z_1) S_{\hat{H}_{FU}}^i \right\}$$

(Metz and DP - PLB 723 (2013))

We now believe the TSSAs in $p^\uparrow p \rightarrow \pi X$
 are due fragmentation effects as the partons
 form pions in the final state

$$d\Delta\sigma^\pi \sim h_1 \otimes S \otimes \left(H_1^{\perp(1)}, H, \int \frac{dz_1}{z_1^2} \frac{\hat{H}_{FU}^{\mathfrak{S}}}{(1/z - 1/z_1)^2} \right)$$

$$H^q(z) = -2z H_1^{\perp(1),q}(z) + 2z \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\frac{1}{z} - \frac{1}{z_1}} \hat{H}_{FU}^{q,\mathfrak{S}}(z, z_1)$$

QCD e.o.m.
relation
(EOMR)

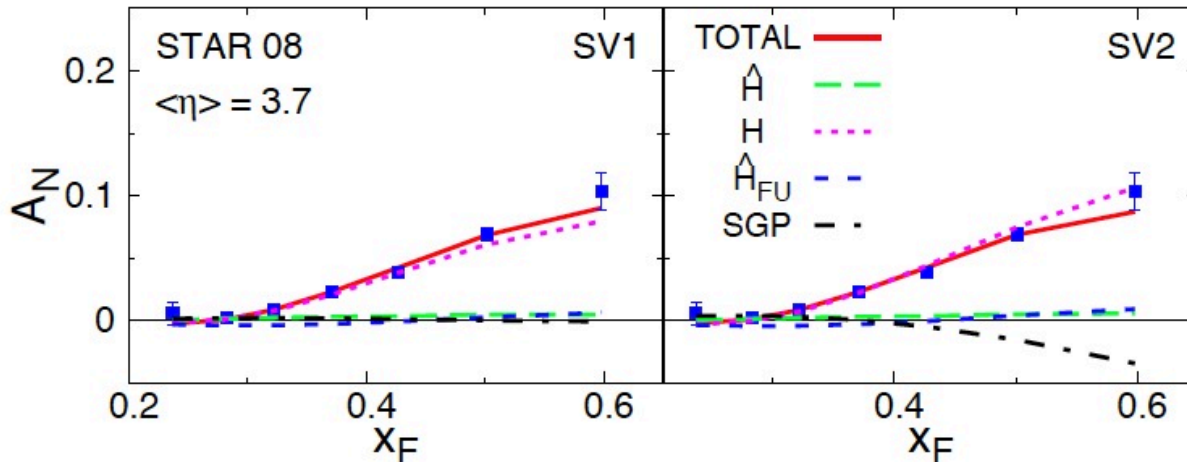
$\longrightarrow \equiv \tilde{H}^q(z)$

$$d\Delta\sigma^\pi \sim \mathbf{h}_1 \otimes S \otimes \left(\mathbf{H}_1^{\perp(1)}, \mathbf{H}, \int \frac{dz_1}{z_1^2} \frac{\hat{H}_{FU}^{\mathfrak{S}}}{(1/z - 1/z_1)^2} \right)$$



$$d\Delta\sigma^\pi \sim \mathbf{h}_1 \otimes \hat{S} \otimes \left(\mathbf{H}_1^{\perp(1)}, \tilde{\mathbf{H}}, \int \frac{dz_1}{z_1^2} \frac{\hat{H}_{FU}^{\mathfrak{S}}}{(1/z - 1/z_1)^2} \right)$$

Also included the Qiu-Sterman term $\pi F_{FT}(\mathbf{x}, \mathbf{x}) = f_{1T}^{\perp(1)}(\mathbf{x})$



Fragmentation term is the main cause of A_N in $pp \rightarrow \pi X$

(Kanazawa, Koike, Metz, DP, PRD 89(RC) (2014))

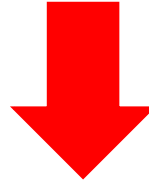
$$d\Delta\sigma^\pi \sim h_1 \otimes \hat{S} \otimes \left(H_1^{\perp(1)}, \tilde{H}, \int \frac{dz_1}{z_1^2} \frac{\hat{H}_{FU}^{\mathfrak{S}}}{(1/z - 1/z_1)^2} \right)$$

$$\frac{H^q(z)}{z} = - \left(1 - z \frac{d}{dz} \right) H_1^{\perp(1),q}(z) - \frac{2}{z} \int_z^\infty \frac{dz_1}{z_1^2} \frac{\hat{H}_{FU}^{q,\mathfrak{S}}(z, z_1)}{(1/z - 1/z_1)^2}$$

Lorentz
invariance
relation (LIR)

(Kanazawa, Koike, Metz, DP, Schlegel, PRD **93** (2016))

$$d\Delta\sigma^\pi \sim \mathbf{h}_1 \otimes \hat{S} \otimes \left(\mathbf{H}_1^{\perp(1)}, \tilde{H}, \int \frac{dz_1}{z_1^2} \frac{\hat{H}_{FU}^{\mathfrak{S}}}{(1/z - 1/z_1)^2} \right)$$



$$d\Delta\sigma^\pi \sim \mathbf{h}_1 \otimes \tilde{S} \otimes \left(\mathbf{H}_1^{\perp(1)}, \tilde{H} \right)$$

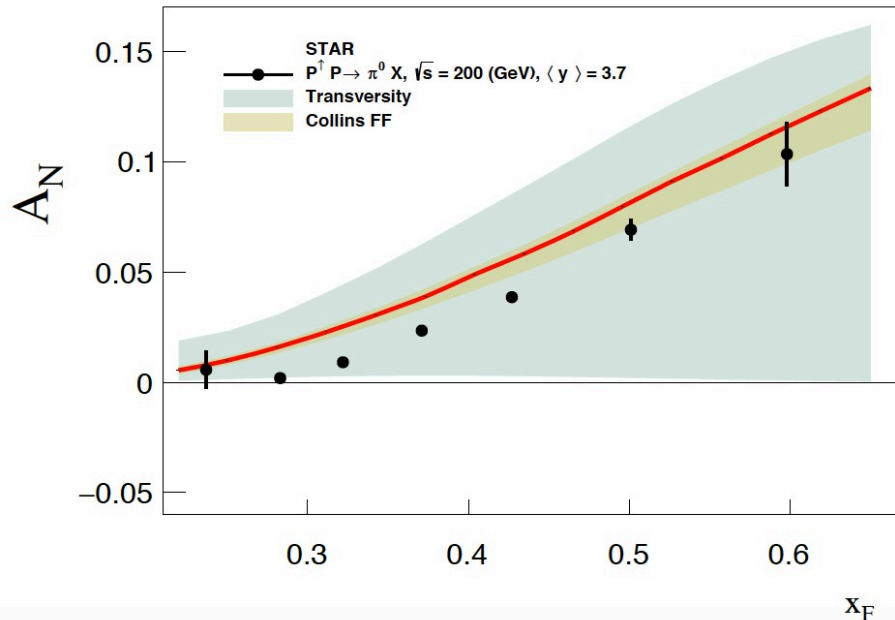
$$E_h \frac{d\Delta\sigma^{Frag}(S_T)}{d^3\vec{P}_h} = - \frac{4\alpha_s^2 M_h}{S} \epsilon^{P' P P_h S_T} \sum_i \sum_{a,b,c} \int_0^1 \frac{dz}{z^3} \int_0^1 dx' \int_0^1 dx \delta(\hat{s} + \hat{t} + \hat{u}) \frac{1}{\hat{s}} \\ \times h_1^a(x) f_1^b(x') \left\{ \left[H_1^{\perp(1),c}(z) - z \frac{dH_1^{\perp(1),c}(z)}{dz} \right] \tilde{S}_{H_1^\perp}^i + \left[-2H_1^{\perp(1),c}(z) + \frac{1}{z} \tilde{H}^c(z) \right] \tilde{S}_H^i \right\}$$

where $\tilde{S}_{H_1^\perp}^i \equiv \frac{S_{H_1^\perp}^i - S_{HFU}^i}{-x'\hat{t} - x\hat{u}}$ and $\tilde{S}_H^i \equiv \frac{S_H^i - S_{HFU}^i}{-x'\hat{t} - x\hat{u}}$

$$d\Delta\sigma^\pi \sim h_1 \otimes \hat{S} \otimes \left(H_1^{\perp(1)}, \tilde{H}, \int \frac{dz_1}{z_1^2} \frac{\hat{H}_{FU}^S}{(1/z - 1/z_1)^2} \right)$$



$$d\Delta\sigma^\pi \sim h_1 \otimes \tilde{S} \otimes \left(H_1^{\perp(1)}, \tilde{H} \right)$$

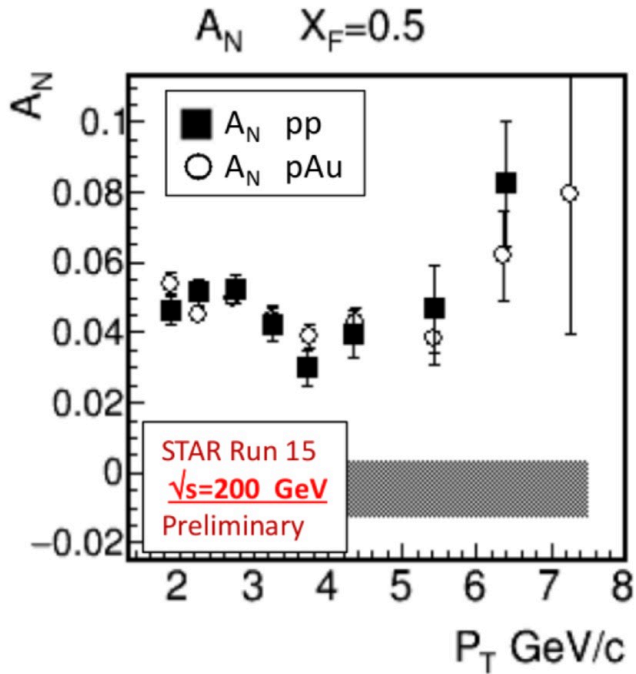


Fragmentation term is the main cause of A_N in $pp \rightarrow \pi X$

The A_N data from RHIC can be used along with measurements from SOLID at JLab to constrain transversity at large $x!$

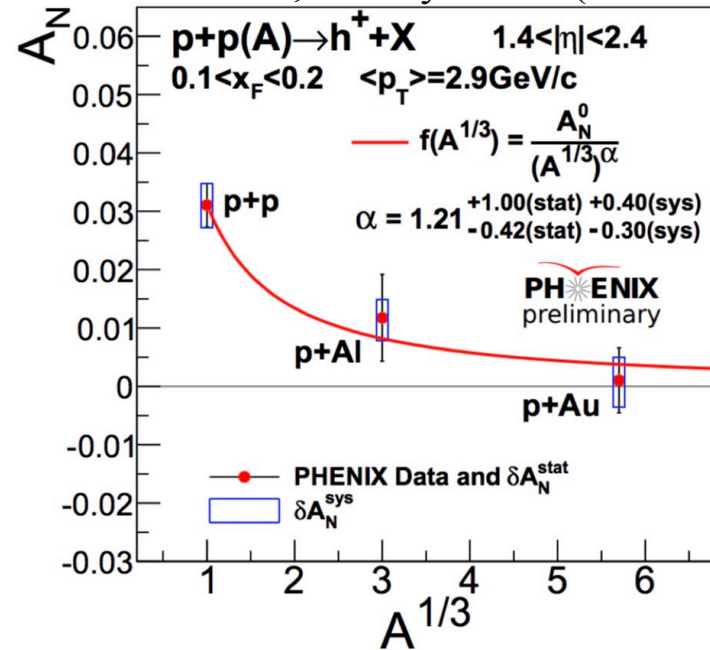
A_N in $pA \rightarrow \pi X$

STAR, Talk by C. Dilks (DIS 2016)



STAR shows *no* suppression with A
 (neutral pions, larger x_F region)

PHENIX, Talk by J. Bok (DIS 2018)



PHENIX shows $A^{1/3}$ suppression
 (charged hadrons, smaller x_F region)

2013 expression from Metz and DP

$$\begin{aligned}
 E_h \frac{d\Delta\sigma^{Frag}(S_T)}{d^3\vec{P}_h} = & - \frac{4\alpha_s^2 M_h}{S} \epsilon^{P' P P_h S_T} \sum_i \sum_{a,b,c} \int_0^1 \frac{dz}{z^3} \int_0^1 dx' \int_0^1 dx \delta(\hat{s} + \hat{t} + \hat{u}) \frac{1}{\hat{s}(-x'\hat{t} - x\hat{u})} \\
 & \times h_1^a(x) f_1^b(x') \left\{ \left[H_1^{\perp(1),c}(z) - z \frac{dH_1^{\perp(1),c}(z)}{dz} \right] S_{H_1^\perp}^i + \frac{1}{z} H^c(z) S_H^i \right. \\
 & \left. + \frac{2}{z} \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\left(\frac{1}{z} - \frac{1}{z_1}\right)^2} \hat{H}_{FU}^{c,\mathfrak{S}}(z, z_1) S_{\hat{H}_{FU}}^i \right\}
 \end{aligned}$$

2013 expression from Metz and DP

$$E_h \frac{d\Delta\sigma^{Frag}(S_T)}{d^3\vec{P}_h} = - \frac{4\alpha_s^2 M_h}{S} \epsilon^{P'P P_h S_T} \sum_i \sum_{a,b,c} \int_0^1 \frac{dz}{z^3} \int_0^1 dx' \int_0^1 dx \delta(\hat{s} + \hat{t} + \hat{u}) \frac{1}{\hat{s}(-x'\hat{t} - x\hat{u})}$$

$$\times h_1^a(x) f_1^b(x') \left\{ \left[H_1^{\perp(1),c}(z) - z \frac{dH_1^{\perp(1),c}(z)}{dz} \right] S_{H_1^\perp}^i + \frac{1}{z} H^c(z) S_H^i \right.$$

$\sim A^{-1/3}$

$$\left. + \frac{2}{z} \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\left(\frac{1}{z} - \frac{1}{z_1}\right)^2} \hat{H}_{FU}^{c,\mathfrak{S}}(z, z_1) S_{\hat{H}_{FU}}^i \right\}$$

$\sim A^0$

Include small x multiple-rescatterings to calculate pA TSSA (Hatta, Xiao, Yoshida, Yuan (2017))

2013 expression from Metz and DP

$$E_h \frac{d\Delta\sigma^{Frag}(S_T)}{d^3\vec{P}_h} = - \frac{4\alpha_s^2 M_h}{S} \epsilon^{P' P P_h S_T} \sum_i \sum_{a,b,c} \int_0^1 \frac{dz}{z^3} \int_0^1 dx' \int_0^1 dx \delta(\hat{s} + \hat{t} + \hat{u}) \frac{1}{\hat{s}(-x'\hat{t} - x\hat{u})}$$

$$\times h_1^a(x) f_1^b(x') \left\{ \left[H_1^{\perp(1),c}(z) - z \frac{dH_1^{\perp(1),c}(z)}{dz} \right] S_{H_1^\perp}^i + \frac{1}{z} H^c(z) S_H^i \right.$$

$$\left. + \frac{2}{z} \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\left(\frac{1}{z} - \frac{1}{z_1}\right)^2} \hat{H}_{FU}^{c,\mathfrak{S}}(z, z_1) S_{\hat{H}_{FU}}^i \right\}$$

$\sim A^0$

EOMR + LIR →

$$\frac{2}{z} \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\left(\frac{1}{z} - \frac{1}{z_1}\right)^2} \hat{H}_{FU}^{\pi/c,\mathfrak{S}}(z, z_1) = H_1^{\perp(1),c}(z) + z \frac{dH_1^{\perp(1),c}(z)}{dz} - \frac{1}{z} \tilde{H}^c(z)$$

2013 expression from Metz and DP

$$E_h \frac{d\Delta\sigma^{Frag}(S_T)}{d^3\vec{P}_h} = - \frac{4\alpha_s^2 M_h}{S} \epsilon^{P' P P_h S_T} \sum_i \sum_{a,b,c} \int_0^1 \frac{dz}{z^3} \int_0^1 dx' \int_0^1 dx \delta(\hat{s} + \hat{t} + \hat{u}) \frac{1}{\hat{s}(-x'\hat{t} - x\hat{u})}$$

$$\times h_1^a(x) f_1^b(x') \left\{ \left[H_1^{\perp(1),c}(z) - z \frac{dH_1^{\perp(1),c}(z)}{dz} \right] S_{H_1^\perp}^i + \frac{1}{z} H^c(z) S_H^i \right.$$

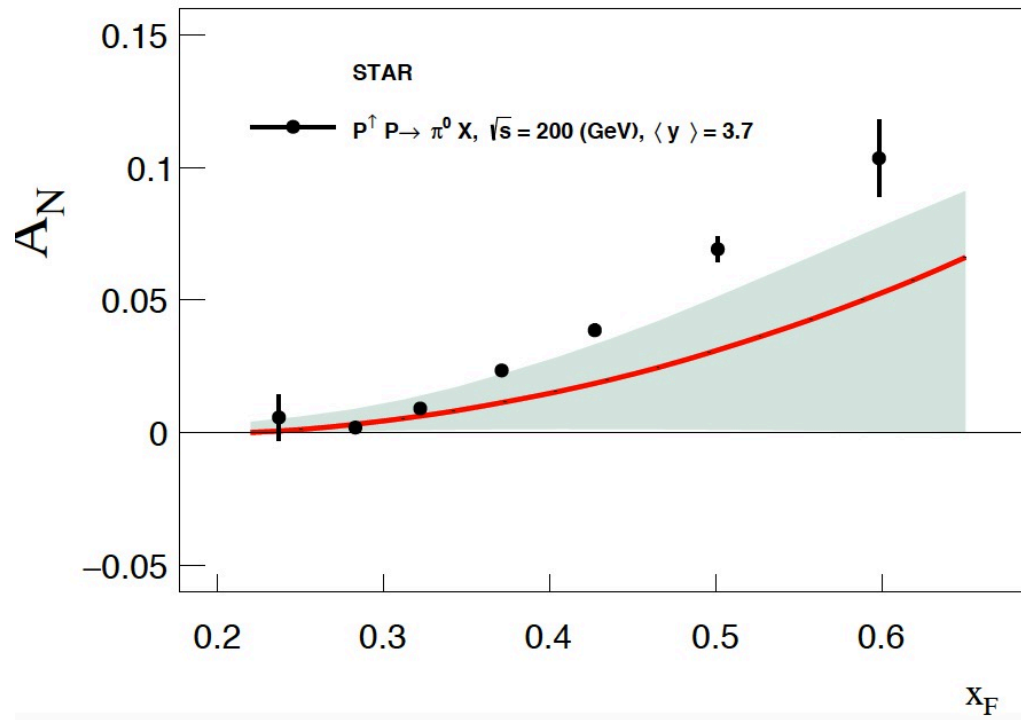
$$\left. + \frac{2}{z} \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\left(\frac{1}{z} - \frac{1}{z_1}\right)^2} \hat{H}_{FU}^{c,\mathfrak{S}}(z, z_1) S_{\hat{H}_{FU}}^i \right\}$$

$\sim A^0$

EOMR + LIR →

$$\frac{2}{z} \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\left(\frac{1}{z} - \frac{1}{z_1}\right)^2} \hat{H}_{FU}^{\pi/c,\mathfrak{S}}(z, z_1) = H_1^{\perp(1),c}(z) + z \frac{dH_1^{\perp(1),c}(z)}{dz} - \frac{1}{z} \tilde{H}^c(z)$$

Calculate pieces involving the (first k_T -moment of the) Collins function to get an updated estimate for the term in blue

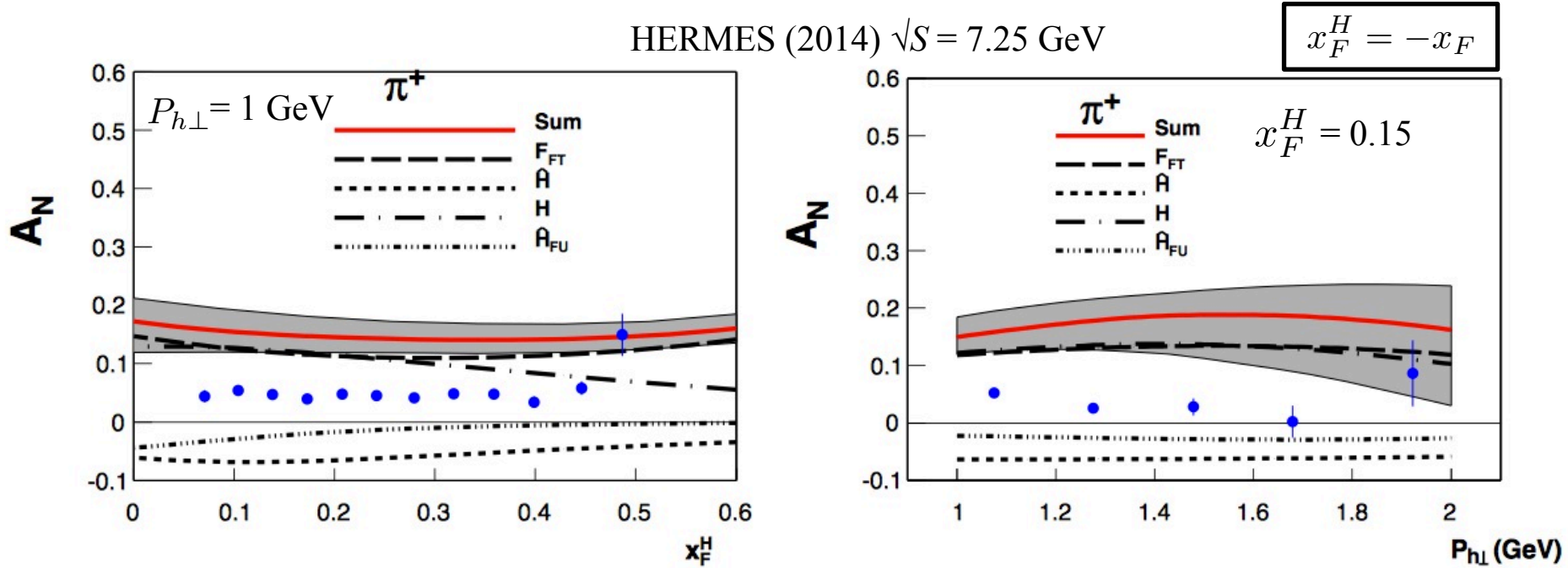


(Gamberg, Kang, DP, Prokudin, PLB 770 (2017))

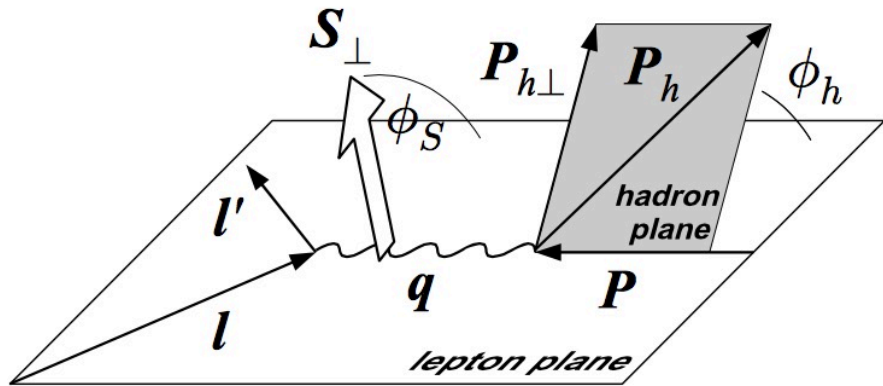
Fragmentation term as the cause of A_N in pp collisions
 is *not* inconsistent with the RHIC pA A_N data.

Reduced theoretical uncertainties are needed.

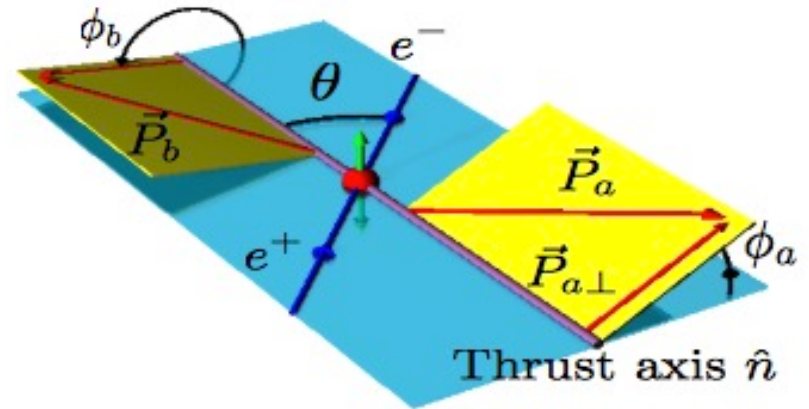
A_N in $ep \rightarrow \pi X$



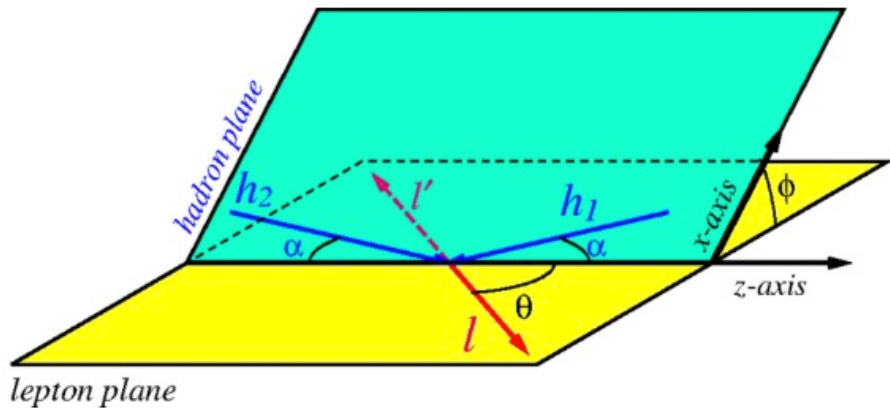
- JLab Hall A also has data for a neutron target, but P_T is too low
 - ➔ 12 GeV upgrade will give valuable data at higher P_T
- This process can help better constrain the 3-parton FFs that have been fit in pp
 - ➔ crucial to measure at at EIC – data in forward region!
 - ➔ need to update analysis to include constraints from LIRs



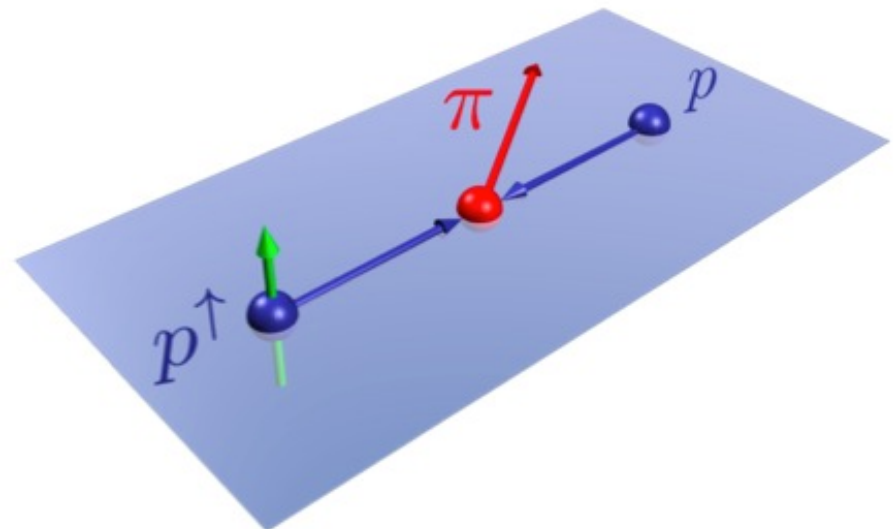
Sivers $\sim \sin(\phi_h - \phi_s)$, Collins $\sim \sin(\phi_h + \phi_s)$, ...



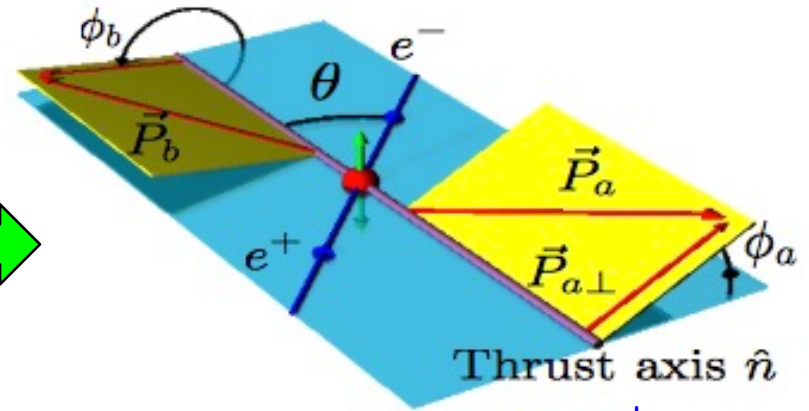
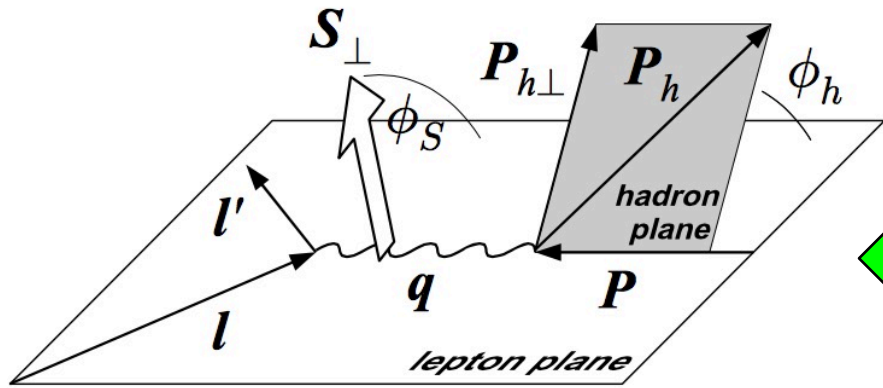
Collins $\sim \cos(\phi_a + \phi_b)$, ...



Sivers $\sim \sin(\phi_s)$ (lepton pair) / Sivers $\sim \cos(\phi_{W/Z})$ (boson)

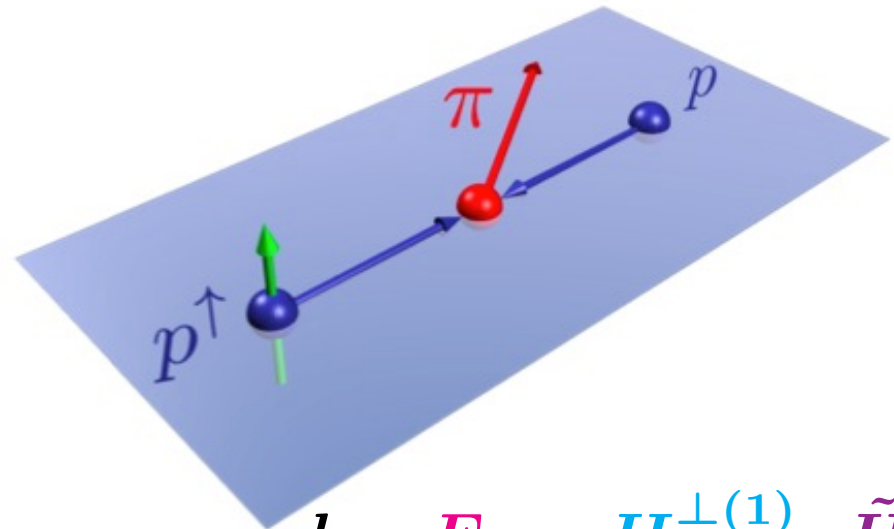
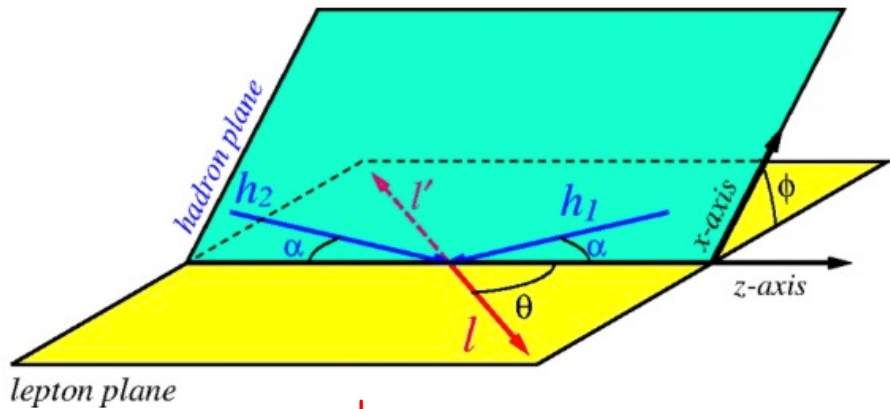


$A_N \sim d\sigma_L - d\sigma_R$



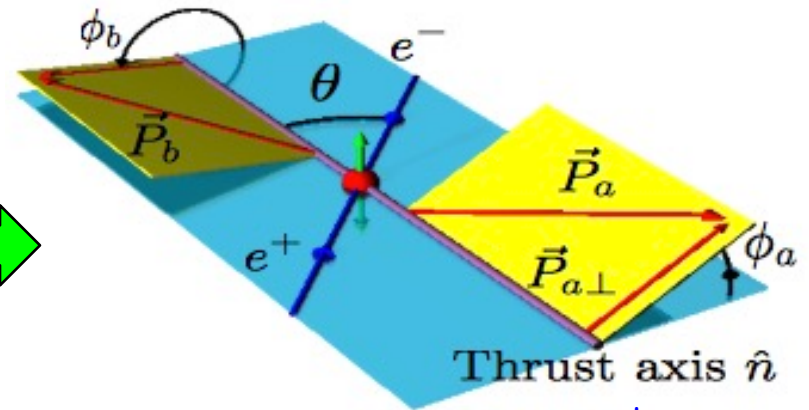
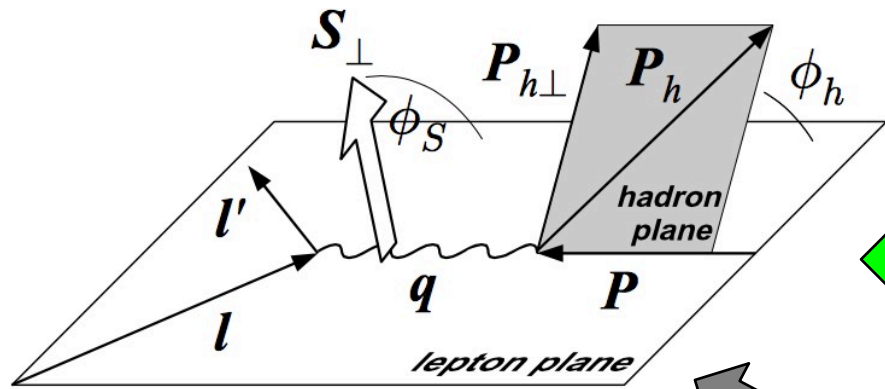
$$h_1, f_{1T}^\perp, H_1^\perp$$

$$H_1^\perp$$



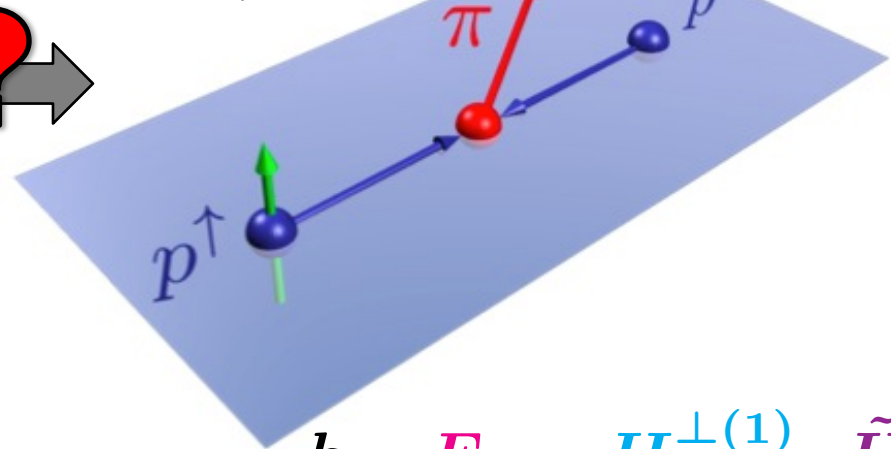
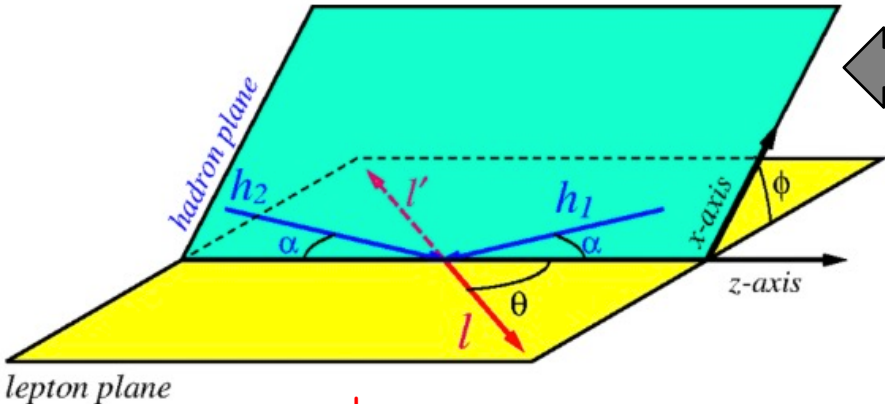
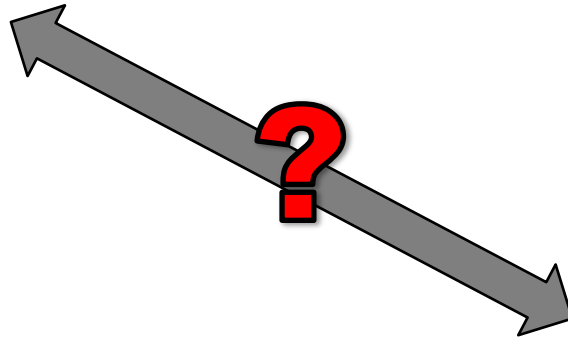
$$f_{1T}^\perp$$

$$h_1, F_{FT}, H_1^{\perp(1)}, \tilde{H}$$



$h_1, f_{1T}^{\perp}, H_1^{\perp}$

H_1^{\perp}



f_{1T}^{\perp}

$h_1, F_{FT}, H_1^{\perp(1)}, \tilde{H}$





Toward a Global Analysis of Transverse Spin Observables

Recall the current phenomenology of TMD observables...

$$\tilde{f}_{1T}^{\perp(1)}(x, b_T; Q^2, \mu_Q) \sim \boxed{F_{FT}(x, x; \mu_{b_*})} \exp \left[-S_{pert}(b_*(b_T); \mu_{b_*}, Q, \mu_Q) - S_{NP}^{f_{1T}^{\perp}}(b_T, Q) \right]$$

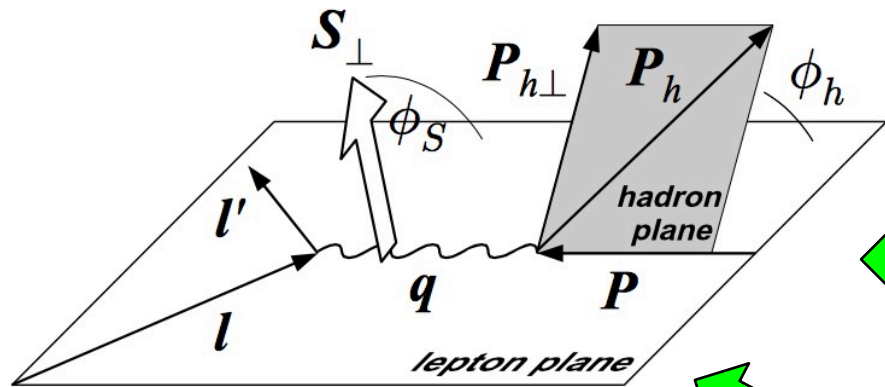
$$\boxed{g_{f_{1T}^{\perp}}(x, b_T)} + g_K(b_T) \ln(Q/Q_0)$$

$$\tilde{H}_1^{\perp(1)}(z, b_T; Q^2, \mu_Q) \sim \boxed{H_1^{\perp(1)}(z; \mu_{b_*})} \exp \left[-S_{pert}(b_*(b_T); \mu_{b_*}, Q, \mu_Q) - S_{NP}^{H_1^{\perp}}(b_T, Q) \right]$$

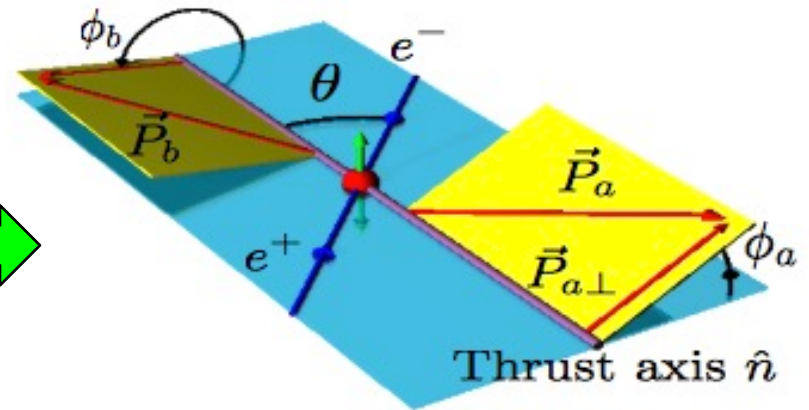
$$\boxed{g_{H_1^{\perp}}(z, b_T)} + g_K(b_T) \ln(Q/Q_0)$$

The **CT3 functions** (along with the NP g -functions) are what get extracted in analyses of TSSAs in **TMD processes** that use CSS evolution!

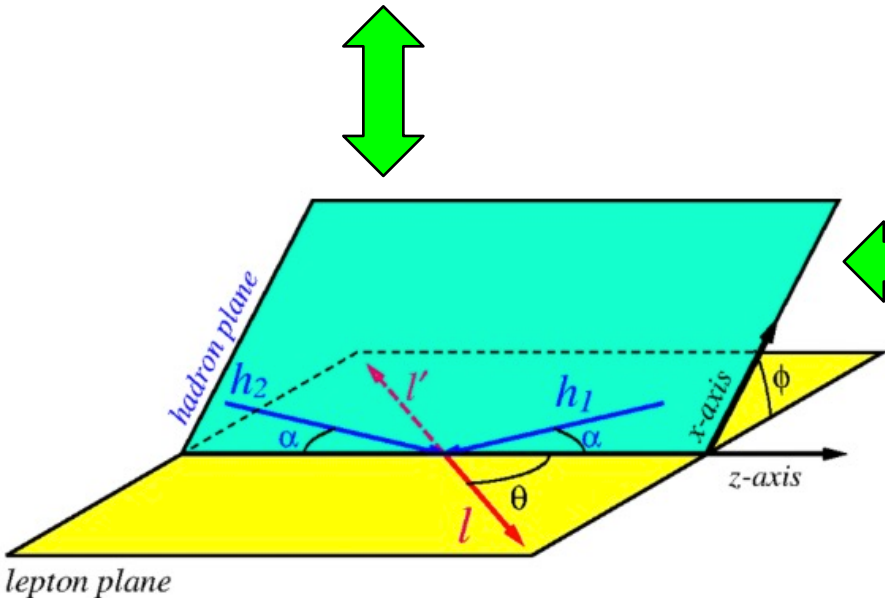
(Echevarria, Idilbi, Kang, Vitev (2014); Kang, Prokudin, Sun, Yuan (2016))



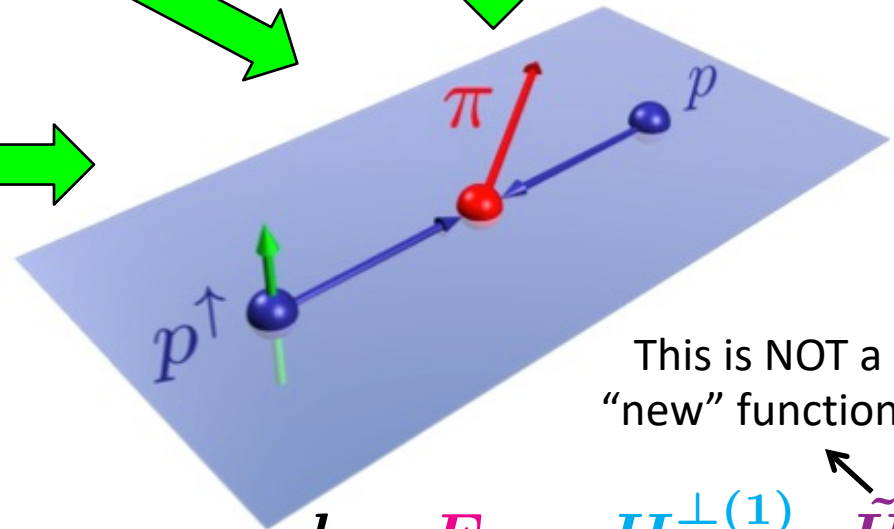
$h_1, F_{FT}, H_1^{\perp(1)}$



$H_1^{\perp(1)}$



F_{FT}

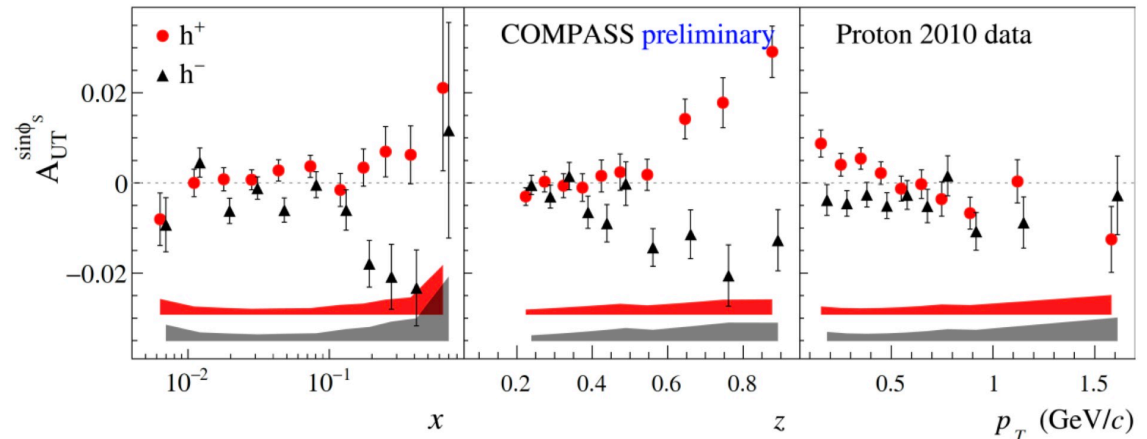


This is NOT a "new" function!

$h_1, F_{FT}, H_1^{\perp(1)}, \tilde{H}$

$A_{UT}^{\sin\phi_S}$ in SIDIS integrated over P_T (Mulders, Tangerman (1996); Bacchetta, et al. (2007))

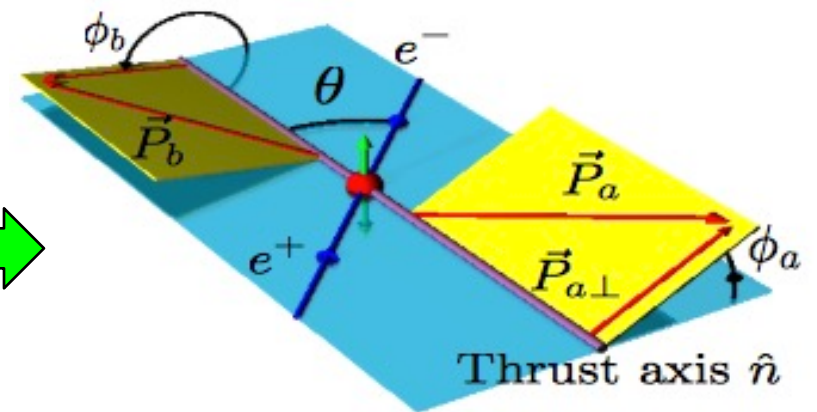
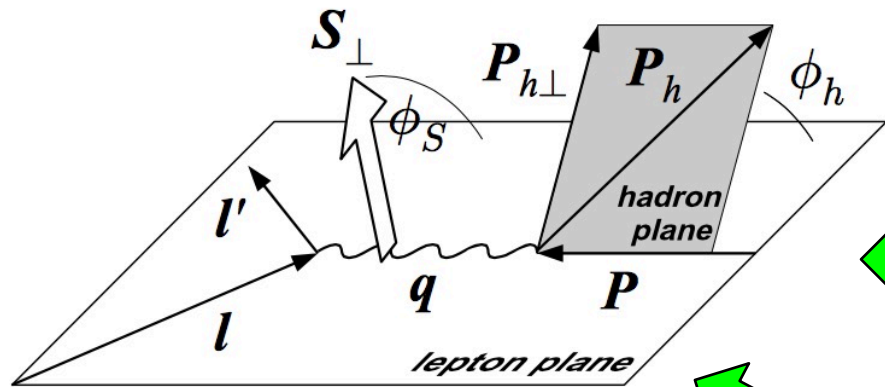
$$F_{UT}^{\sin\phi_S} \propto \sum_a e_a^2 \frac{2M_h}{Q} h_1^a(x) \frac{\tilde{H}^a(z)}{z}$$



$A_{UT}^{\sin\phi_S}$ in $e^+e^- \rightarrow h_1 h_2 X$ integrated over q_T (Boer, Jakob, Mulders (1997))

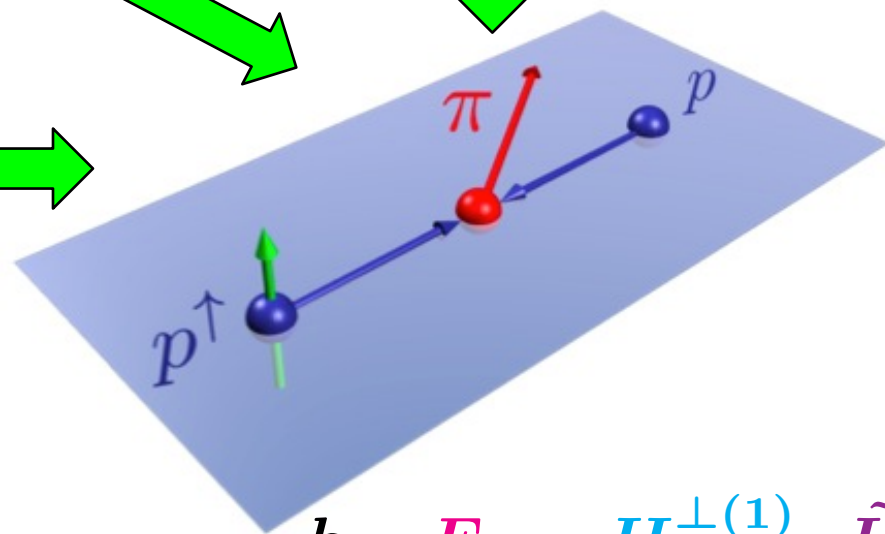
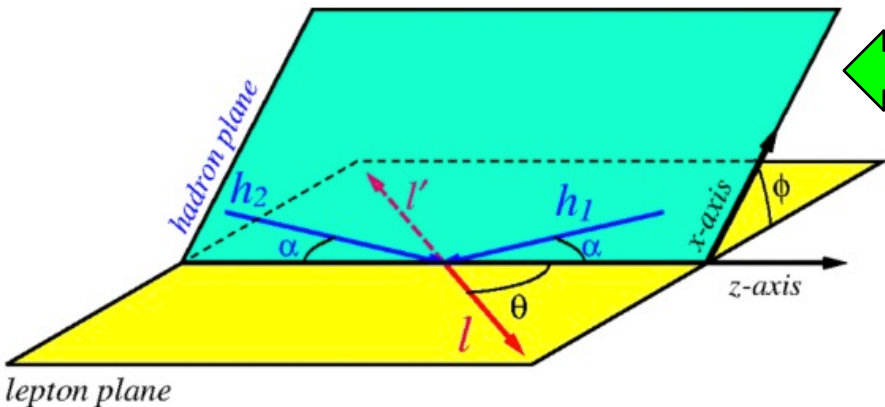
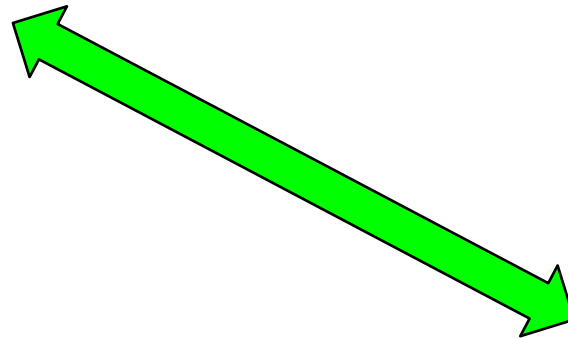
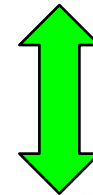
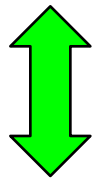
$$F_{UT}^{\sin\phi_S} \propto \sum_{a,\bar{a}} e_a^2 \left(\frac{2M_2}{Q} D_1^a(z_1) \frac{D_T^{\bar{a}}(z_2)}{z_2} + \frac{2M_1}{Q} \frac{\tilde{H}(z_1)}{z_1} H_1^{\bar{a}}(z_2) \right)$$

And also the TMD version of these (and other) observables (but with many more terms)



$h_1, F_{FT}, H_1^{\perp(1)}, \tilde{H}$

$H_1^{\perp(1)}, \tilde{H}$



F_{FT}

$h_1, F_{FT}, H_1^{\perp(1)}, \tilde{H}$

EOMR

$$H^q(z) = -2z H_1^{\perp(1),q}(z) + 2z \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\frac{1}{z} - \frac{1}{z_1}} \hat{H}_{FU}^{q,\mathfrak{S}}(z, z_1)$$

LIR

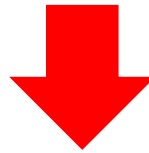
$$\frac{H^q(z)}{z} = - \left(1 - z \frac{d}{dz} \right) H_1^{\perp(1),q}(z) - \frac{2}{z} \int_z^\infty \frac{dz_1}{z_1^2} \frac{\hat{H}_{FU}^{q,\mathfrak{S}}(z, z_1)}{(1/z - 1/z_1)^2}$$

EOMR

$$H^q(z) = -2z H_1^{\perp(1),q}(z) + 2z \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\frac{1}{z} - \frac{1}{z_1}} \hat{H}_{FU}^{q,\mathfrak{S}}(z, z_1)$$

LIR

$$\frac{H^q(z)}{z} = - \left(1 - z \frac{d}{dz} \right) H_1^{\perp(1),q}(z) - \frac{2}{z} \int_z^\infty \frac{dz_1}{z_1^2} \frac{\hat{H}_{FU}^{q,\mathfrak{S}}(z, z_1)}{(1/z - 1/z_1)^2}$$



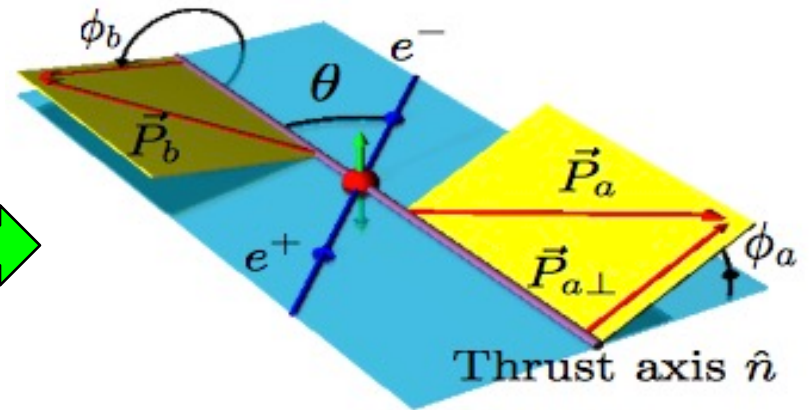
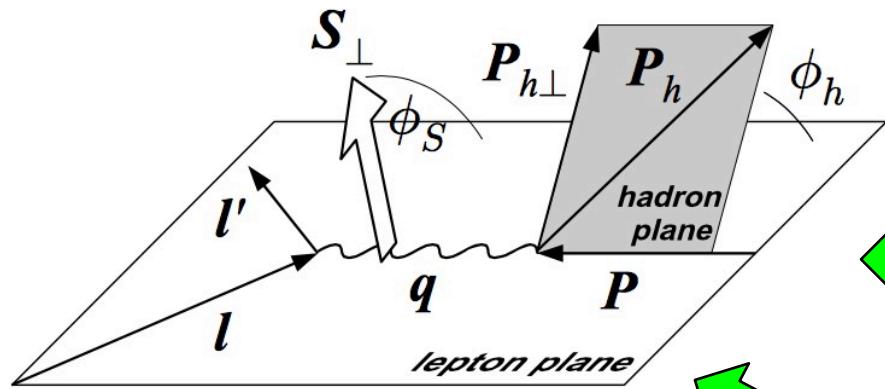
$$H(z) = \int_z^1 dz_1 \int_{z_1}^\infty \frac{dz_2}{z_2^2} 2 \left[\frac{\left(2\left(\frac{2}{z_1} - \frac{1}{z_2}\right) + \frac{1}{z_1} \left(\frac{1}{z_1} - \frac{1}{z_2}\right) \delta\left(\frac{1}{z_1} - \frac{1}{z}\right) \right)}{\left(\frac{1}{z_1} - \frac{1}{z_2}\right)^2} \hat{H}_{FU}^{\mathfrak{S}}(z_1, z_2) \right]$$

$$H_1^{\perp(1)}(z) = -\frac{2}{z} \int_z^1 dz_1 \int_{z_1}^\infty \frac{dz_2}{z_2^2} \frac{\left(\frac{2}{z_1} - \frac{1}{z_2}\right)}{\left(\frac{1}{z_1} - \frac{1}{z_2}\right)^2} \hat{H}_{FU}^{\mathfrak{S}}(z_1, z_2)$$

	PDF (x)		PDF (x, x_1)	FF (z)		FF (z, z_1)
Hadron Pol.						
	<u>intrinsic</u>	<u>kinematical</u>	<u>dynamical</u>	<u>intrinsic</u>	<u>kinematical</u>	<u>dynamical</u>
U	h_1^U	$h_{1U}^{(1)}$	H_{FU}	$h_1^U, h_{1U}^{(1)}$	$H_{1U}^{(1)}$	$\hat{H}_{FU}^{\mathcal{R}, \mathcal{S}}$
L	h_1^L	$h_{1L}^{(1)}$	H_{FL}	$h_1^L, h_{1L}^{(1)}$	$H_{1L}^{(1)}$	$\hat{H}_{FL}^{\mathcal{R}, \mathcal{S}}$
T	g_T	$f_{1T}^{(1)}, g_{1T}^{(1)}$	F_{FT}, G_{FT}	I_T, G_T	$D_{1T}^{(1)}, G_{1T}^{(1)}$	$\hat{D}_{FT}^{\mathcal{R}, \mathcal{S}}, \hat{G}_{FT}^{\mathcal{R}, \mathcal{S}}$

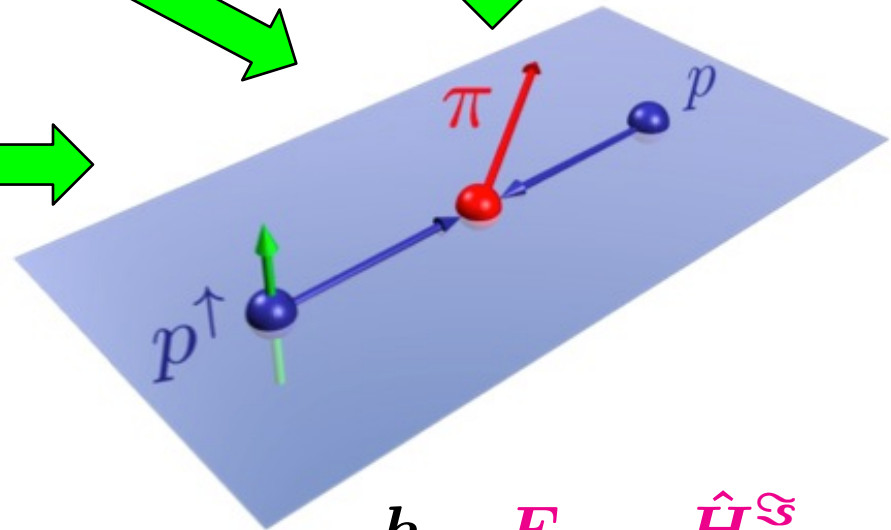
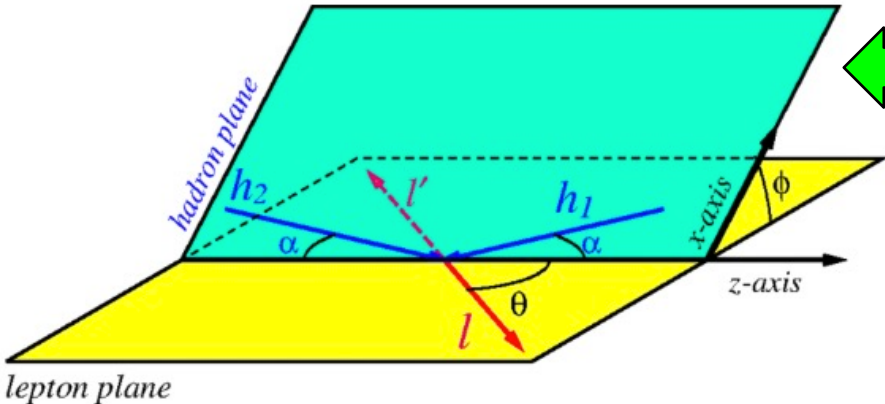
	PDF (x, x_1)	FF (z, z_1)
Hadron Pol.		
U	<u>dynamical</u> H_{FU}	<u>dynamical</u> $\hat{H}_{FU}^{\mathcal{R}, \mathcal{S}}$
L	H_{FL}	$\hat{H}_{FL}^{\mathcal{R}, \mathcal{S}}$
T	F_{FT}, G_{FT}	$\hat{D}_{FT}^{\mathcal{R}, \mathcal{S}}, \hat{G}_{FT}^{\mathcal{R}, \mathcal{S}}$

ALL transverse spin observables are driven by *multi-parton correlations*



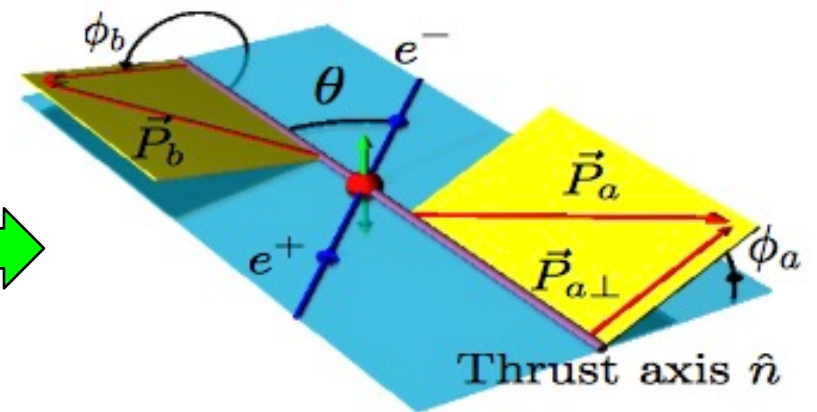
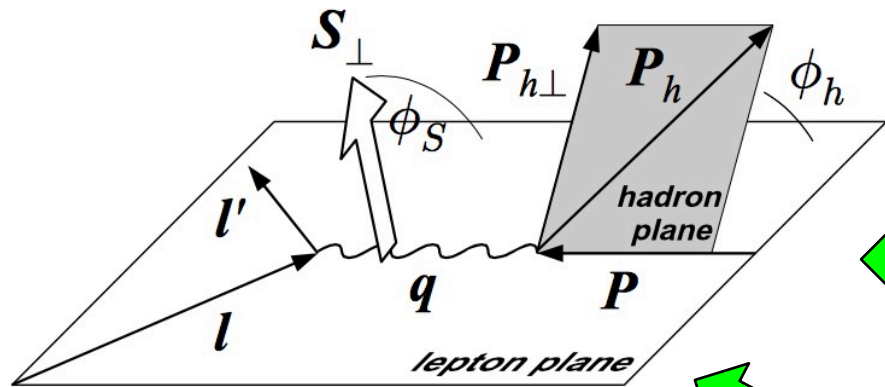
$h_1, F_{FT}, \hat{H}_{FU}^{\mathfrak{S}}$

$\hat{H}_{FU}^{\mathfrak{S}}$



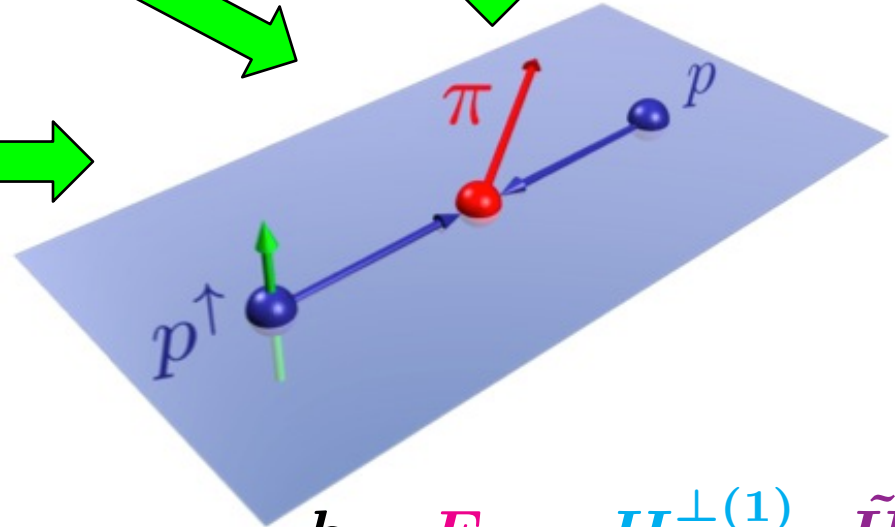
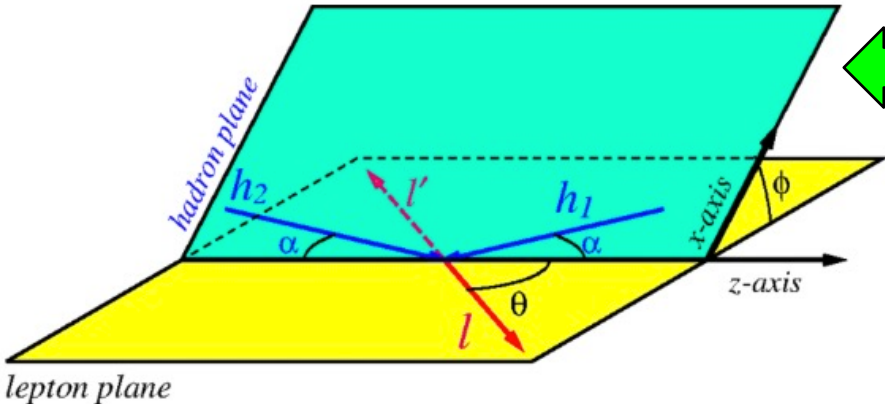
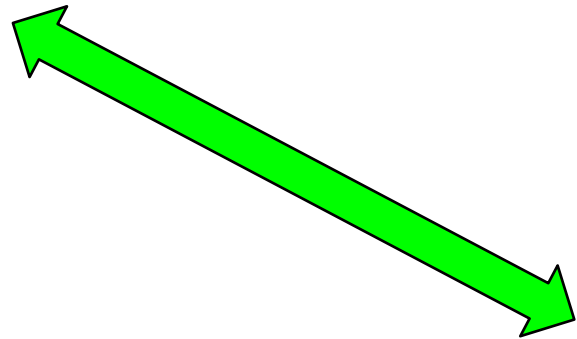
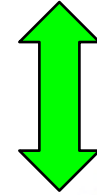
F_{FT}

$h_1, F_{FT}, \hat{H}_{FU}^{\mathfrak{S}}$



$h_1, F_{FT}, H_1^{\perp(1)}, \tilde{H}$

$H_1^{\perp(1)}, \tilde{H}$



F_{FT}

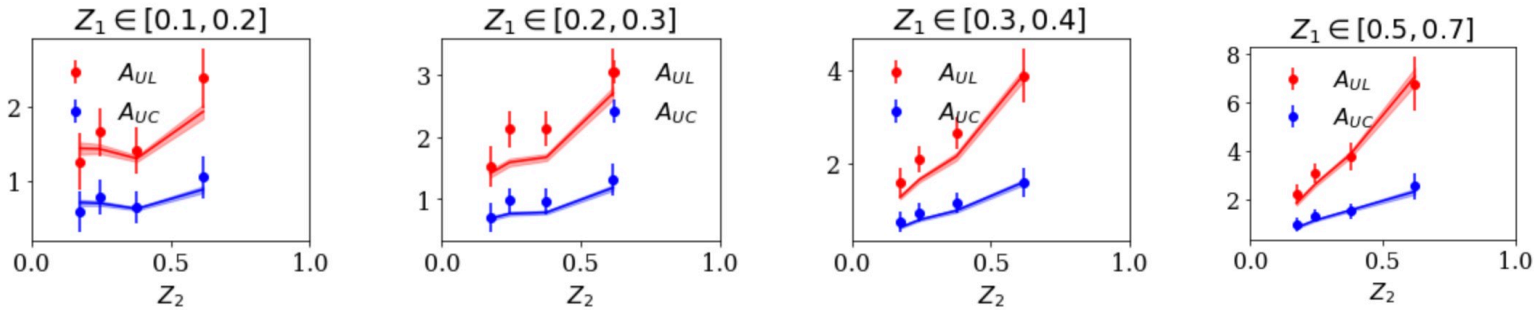
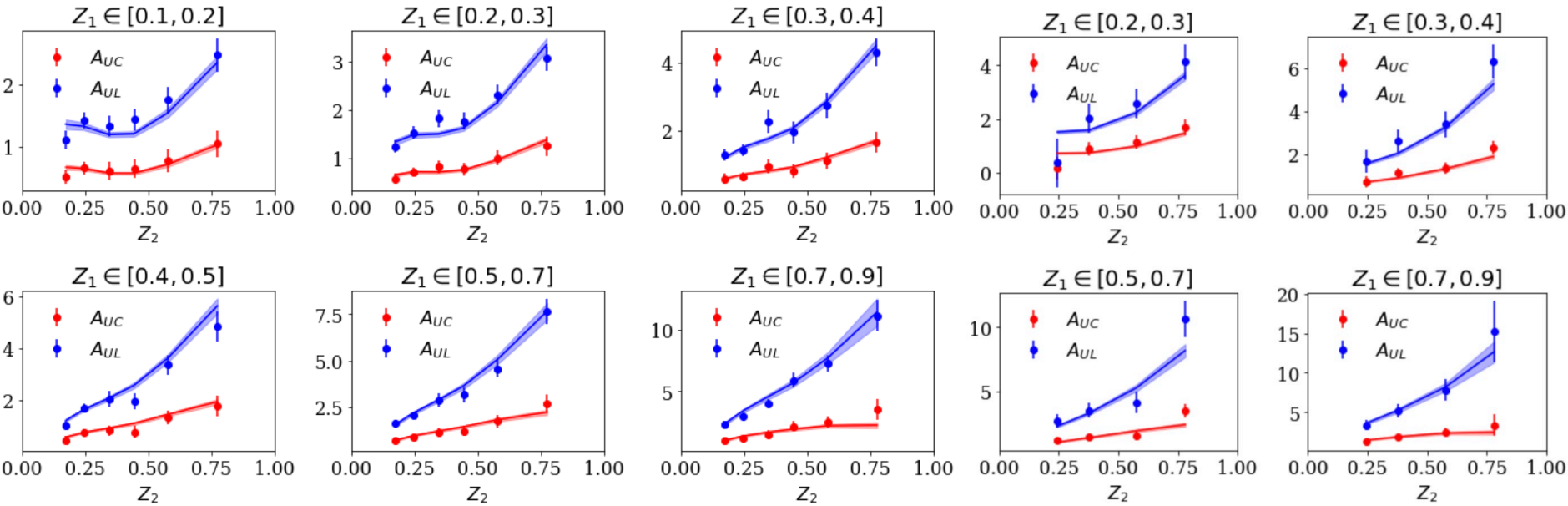
$h_1, F_{FT}, H_1^{\perp(1)}, \tilde{H}$

- What follows are *very preliminary* results of a global fit of
 - 1) Collins effect in e^+e^-
 - 2) Collins effect in SIDIS
 - 3) (Integrated) $A_{UT}^{\sin \phi_s}$ in SIDIS
 - 4) A_N in proton-proton collisions (fragmentation term)

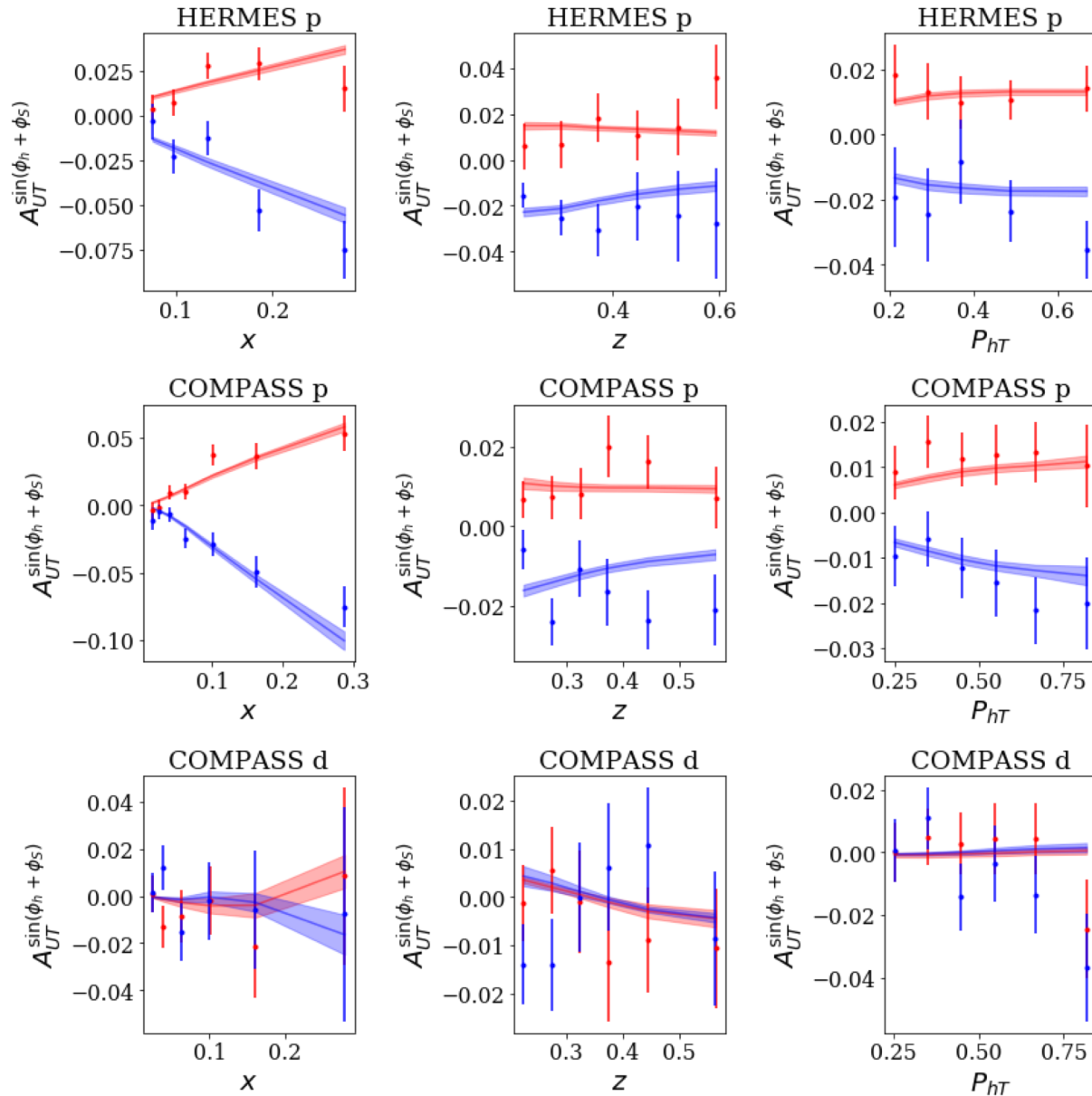
*Also will add Sivers and QS term of A_N to the analysis
- Monte Carlo (MC) sampling was used to determine error bands. For now, we use a simple Gaussian ansatz for TMDs.
- We have found solutions for the relevant non-perturbative functions (including \tilde{H} !) that describe simultaneously a non-trivial amount of observables.
- Large errors in the (transversely polarized) deuteron SIDIS data make flavor separation subject to significant correlations which can only be estimated by MC – an EIC can hopefully deliver more accurate data.

Collins effect e^+e^-

A_{UC} A_{UL}

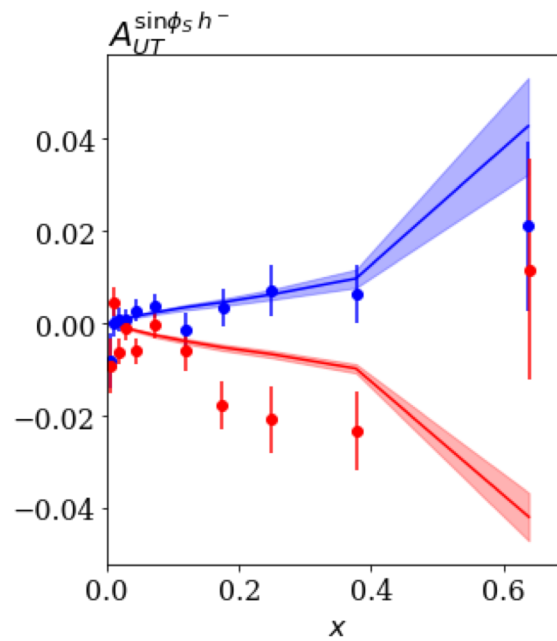
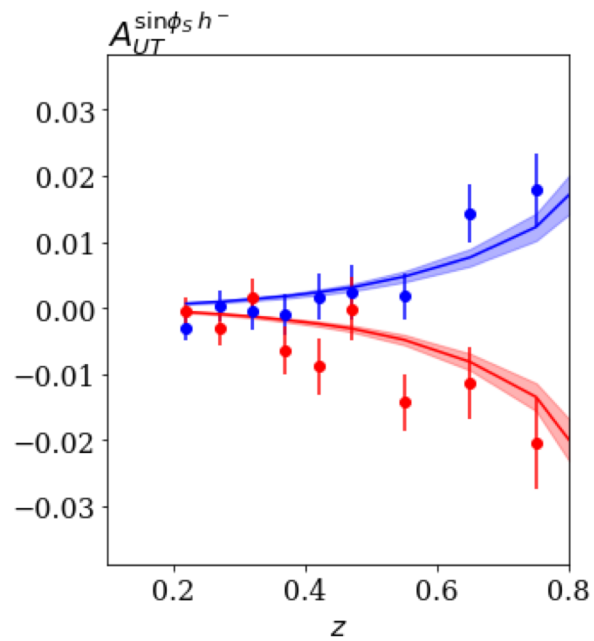


Collins effect SIDIS



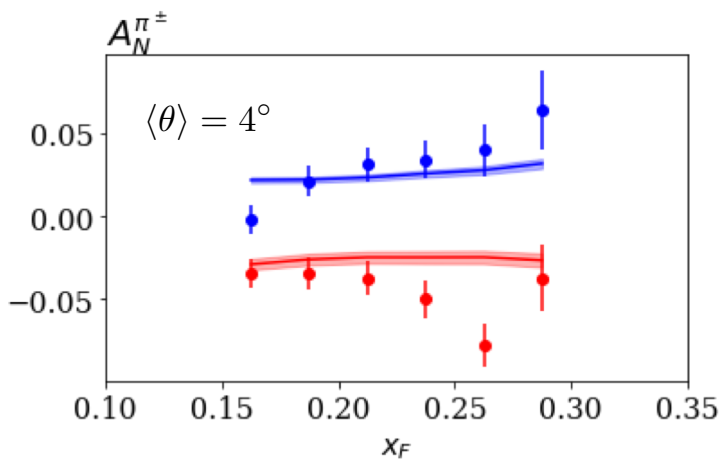
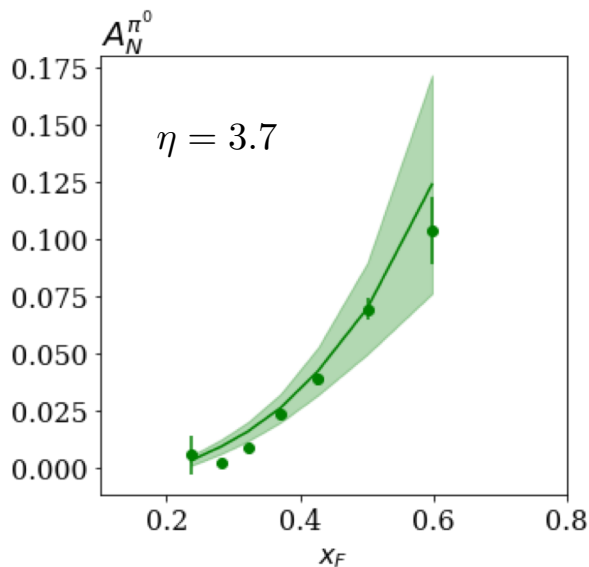
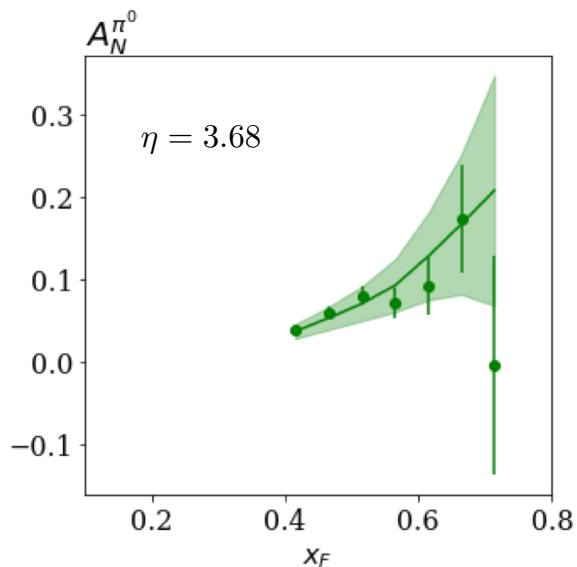
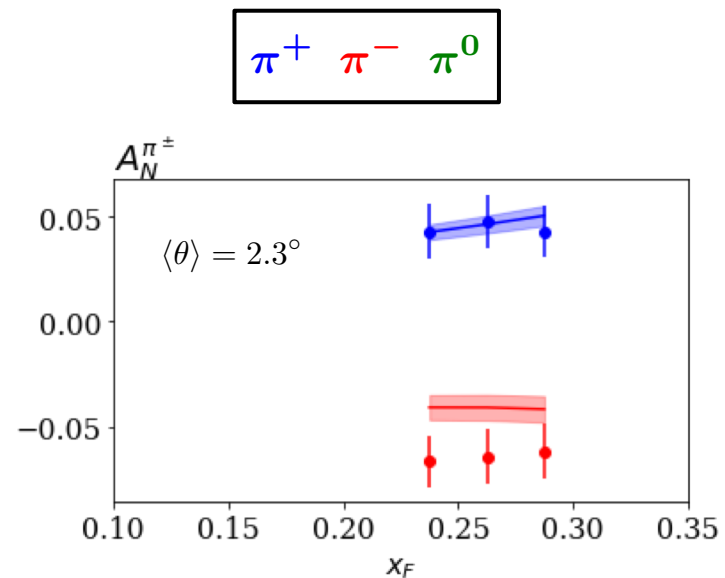
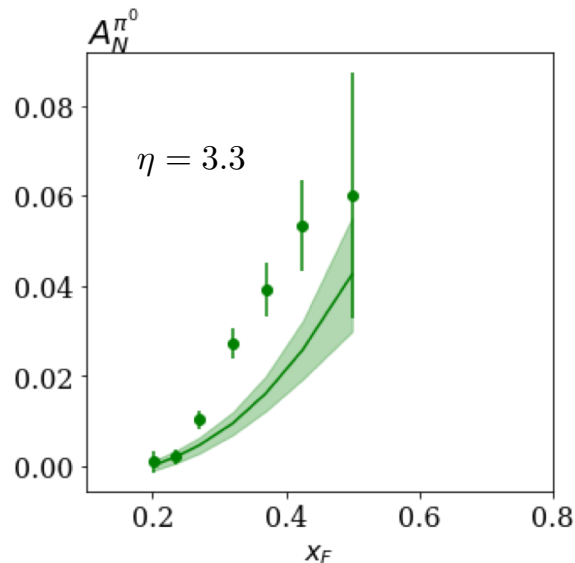
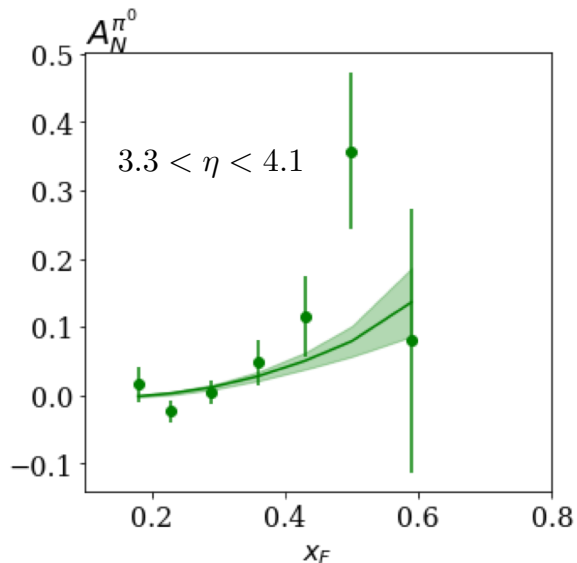
π^+ π^-

$A_{UT}^{\sin \phi_S}$ in SIDIS

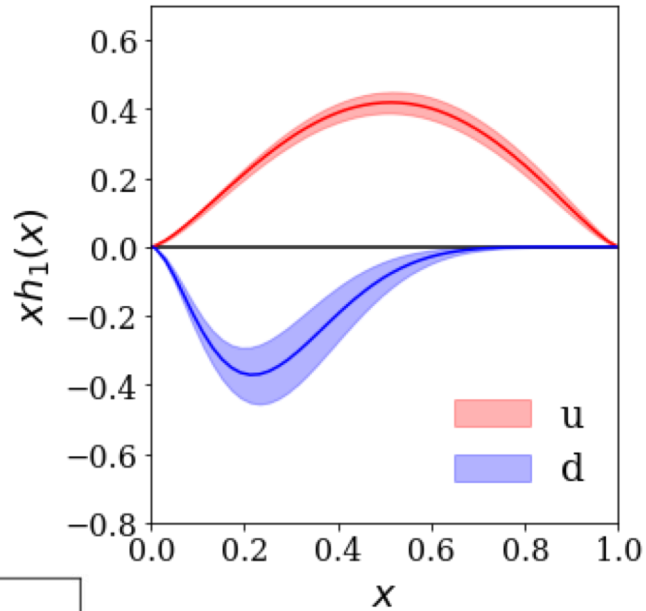


h^+ h^-

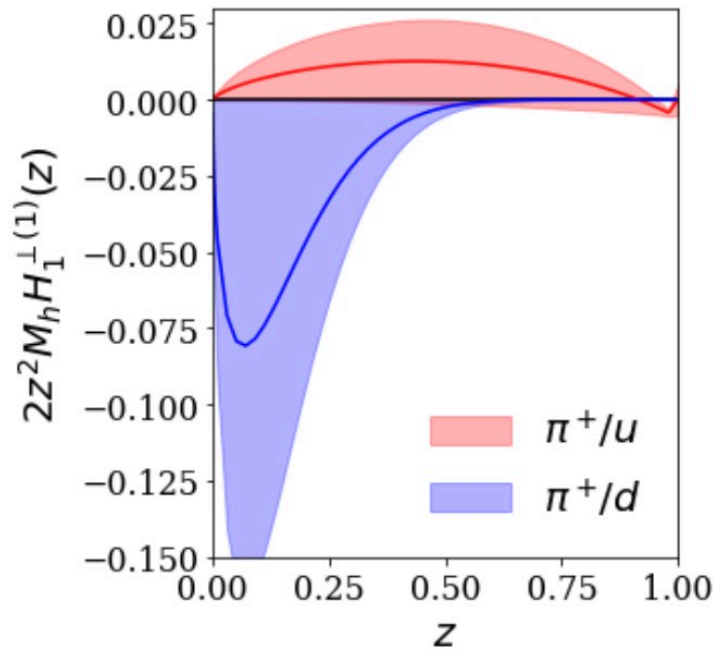
A_N in pp



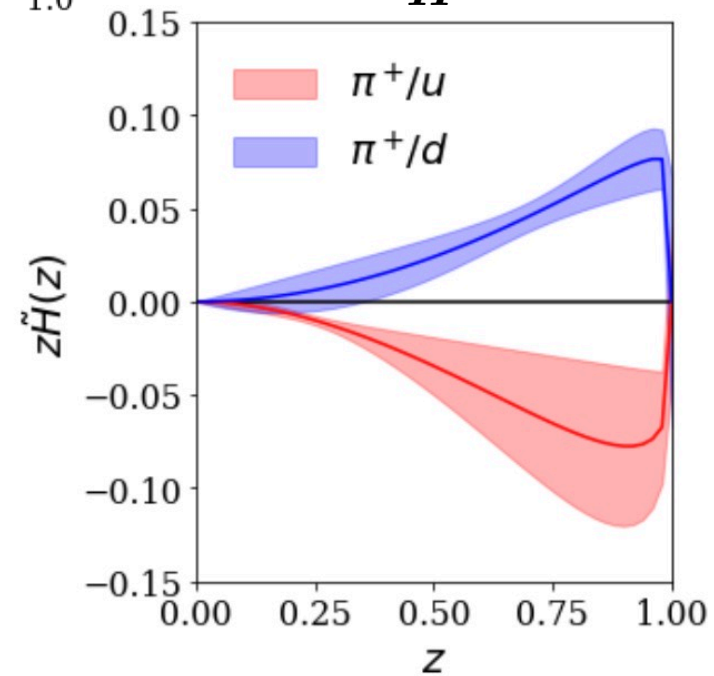
Transversity



Collins



\tilde{H}

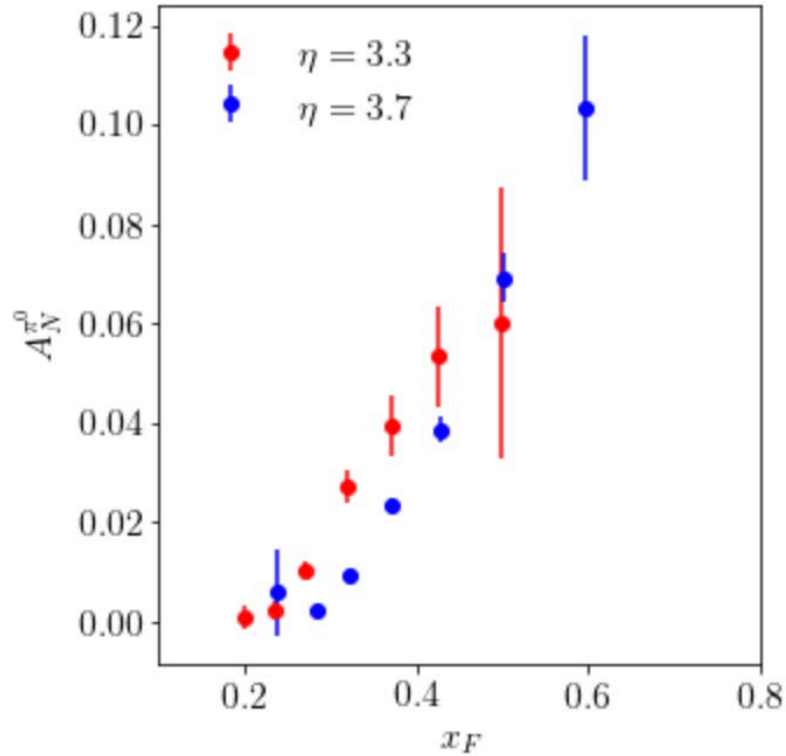


Summary and Outlook

- TMD and collinear functions are highly interconnected, especially for reactions involving transverse spin, and we should treat both types of observables on the same footing.
- A global analysis can be performed of TMD (Sivers and Collins effects) *AND* collinear twist-3 (A_N in pp , $A_{UT}^{\sin\phi_s}$ in SIDIS) transverse-spin observables.
- In addition to the Sivers and Collins effects that will be measured at a future EIC (with improved statistics needed for deuterium), we must also include measurements of A_N in electron-nucleon collisions.



Back-up Slides



There is an *increase* in A_N with P_T for $P_T < 2$ GeV. Need to see if evolution effects can account for this.