

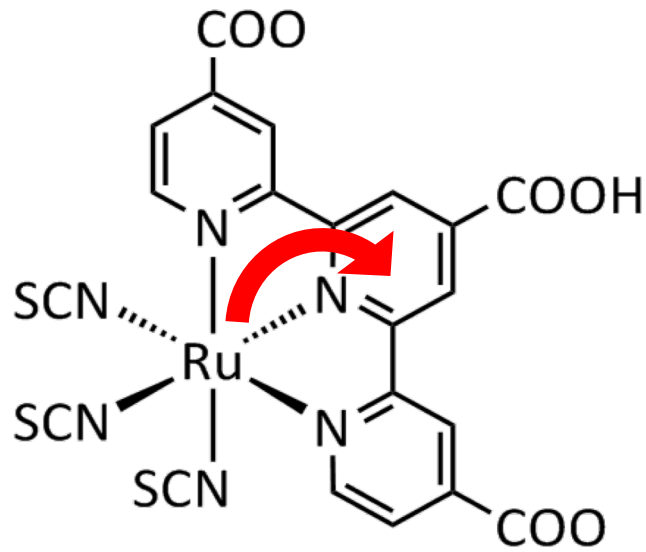
Variational Monte Carlo Optimization and Excited States

Eric Neuscamman

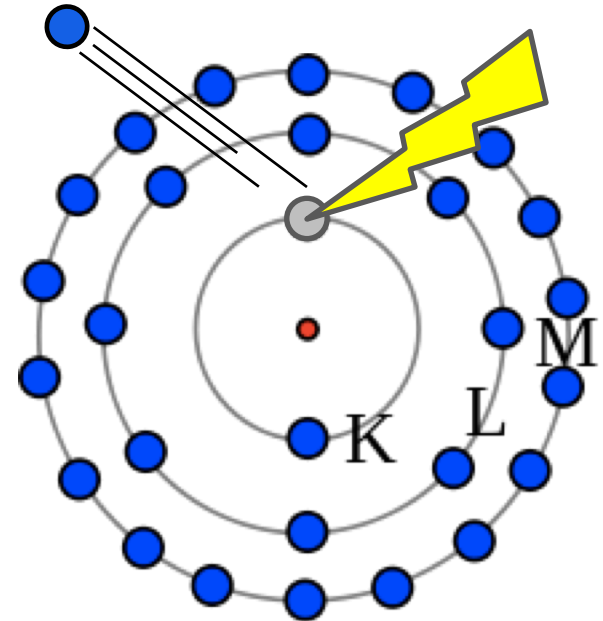
August 9, 2018



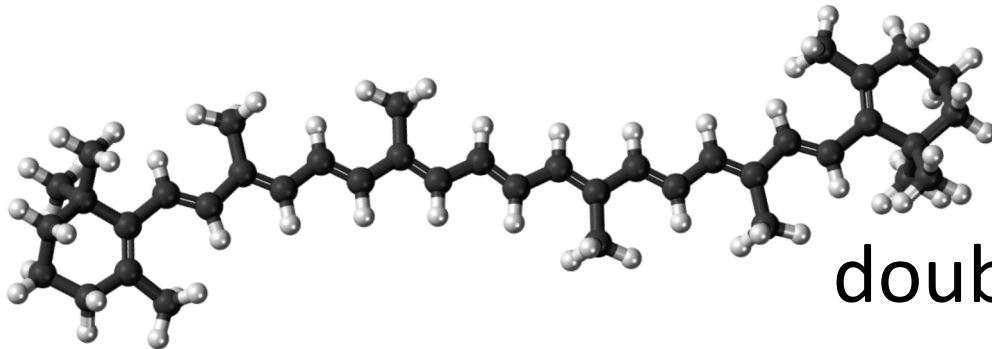
motivation



charge transfer



core spectroscopy



double excitations

the menu



aperitif: number counting Jastrows

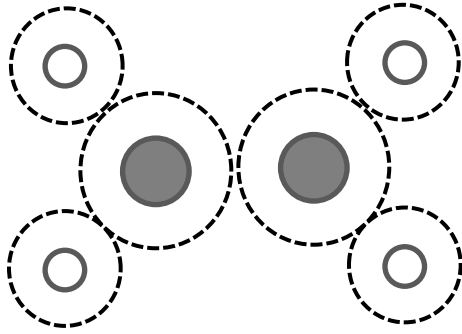


main course: excited states with QMC



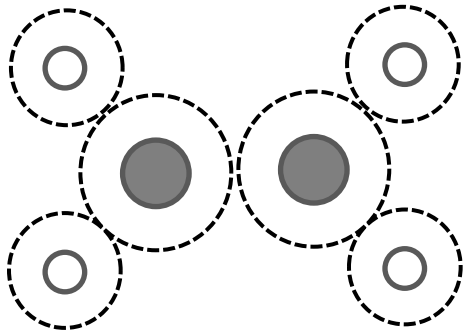
dessert: optimization methods

number counting Jastrow factors



$$C_i(\vec{r}) = 1 \text{ if } \vec{r} \text{ in region } i \\ = 0 \text{ otherwise}$$

number counting Jastrow factors

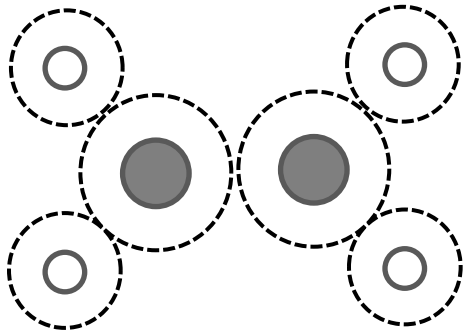


$$C_i(\vec{r}) = 1 \text{ if } \vec{r} \text{ in region } i \\ = 0 \text{ otherwise}$$

$$\exp\left(-\alpha\left(N - \sum_k C_i(\vec{r}_k)\right)^2\right)$$

modifies the balance
between these configs:
 CC C^+C^- C^-C^+

number counting Jastrow factors



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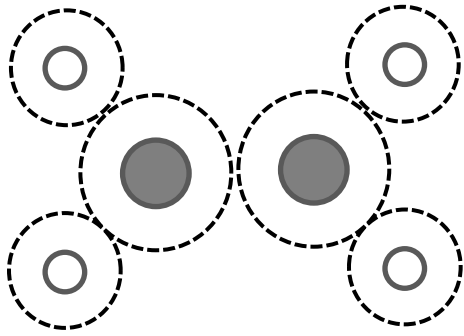
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general form:

$$\exp\left(\sum_{ijkl} F_{ij} C_i(\vec{r}_k) C_j(\vec{r}_l)\right)$$

number counting Jastrow factors



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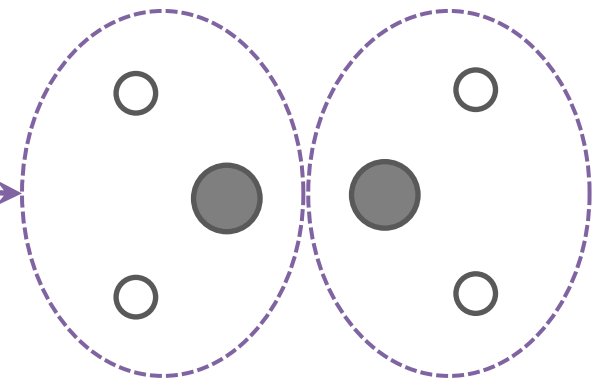
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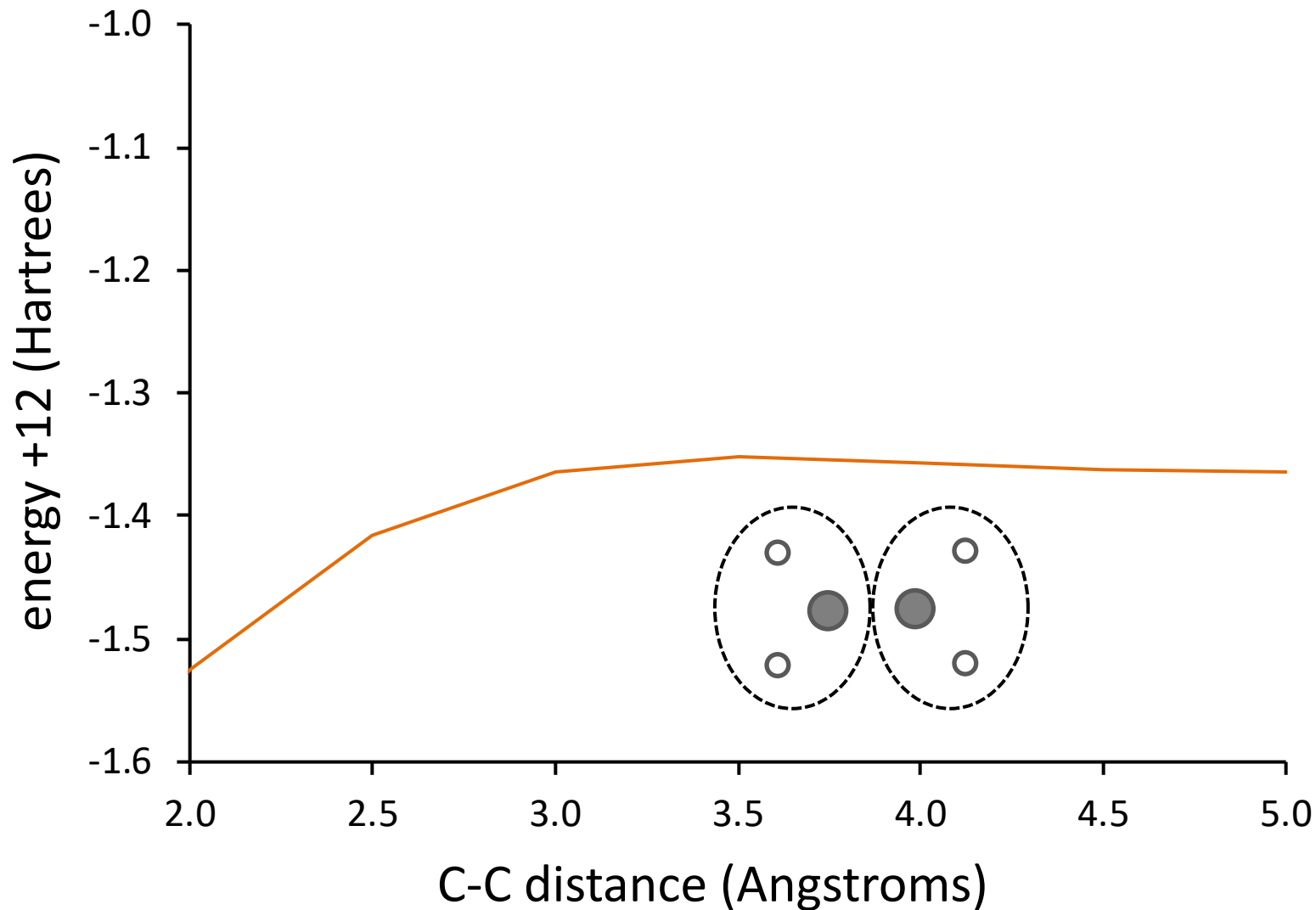
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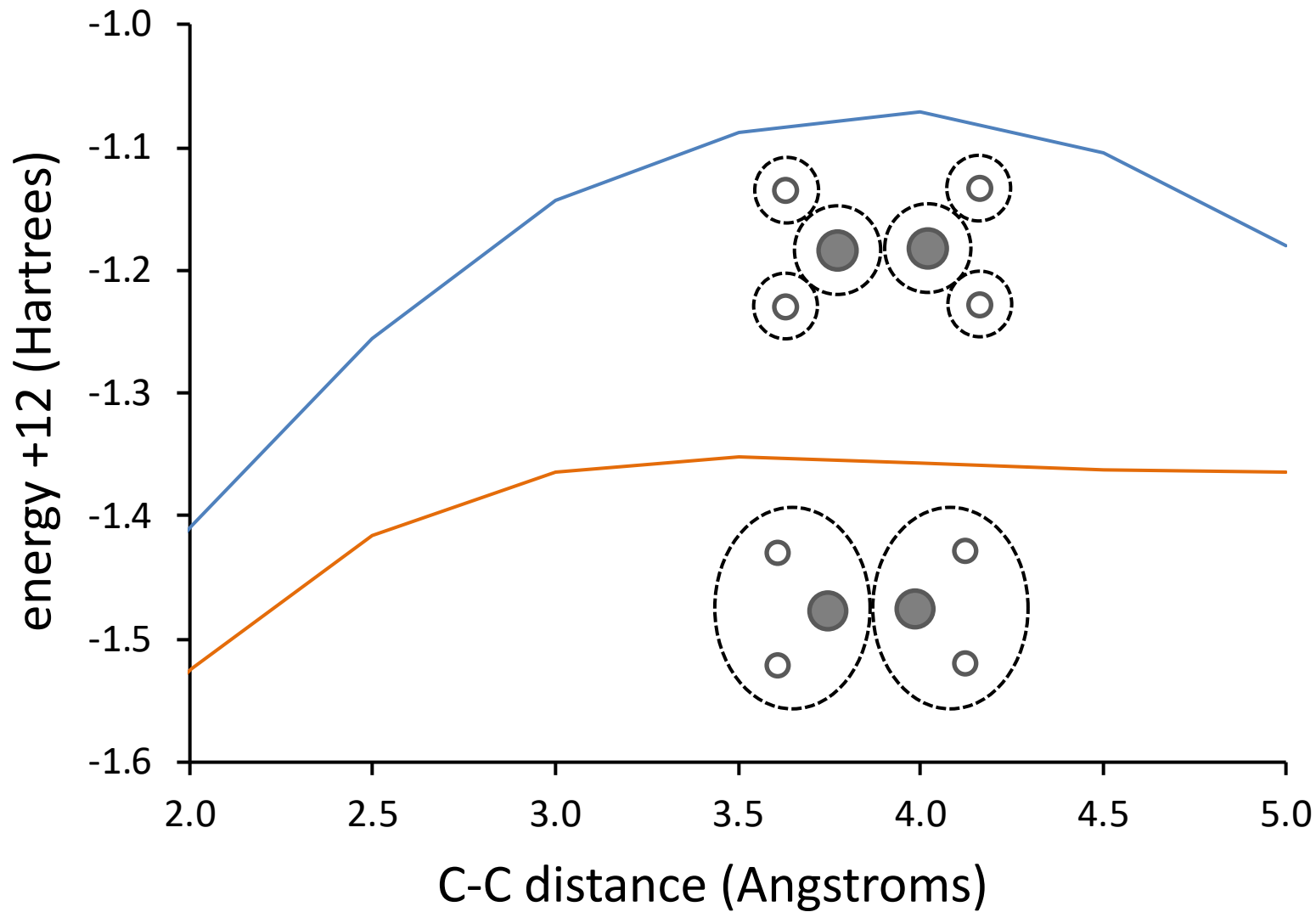
linear combinations?



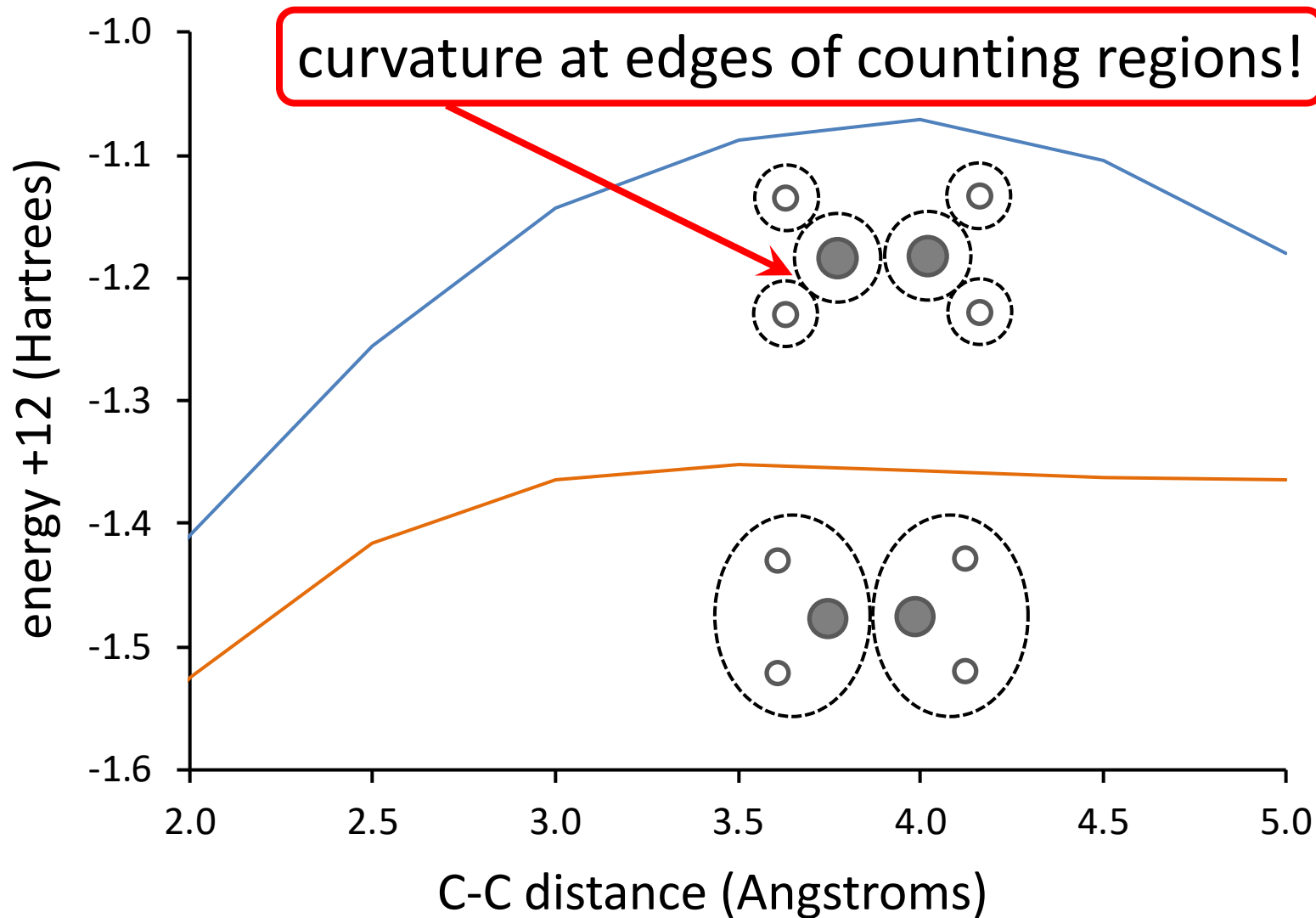
clean additivity



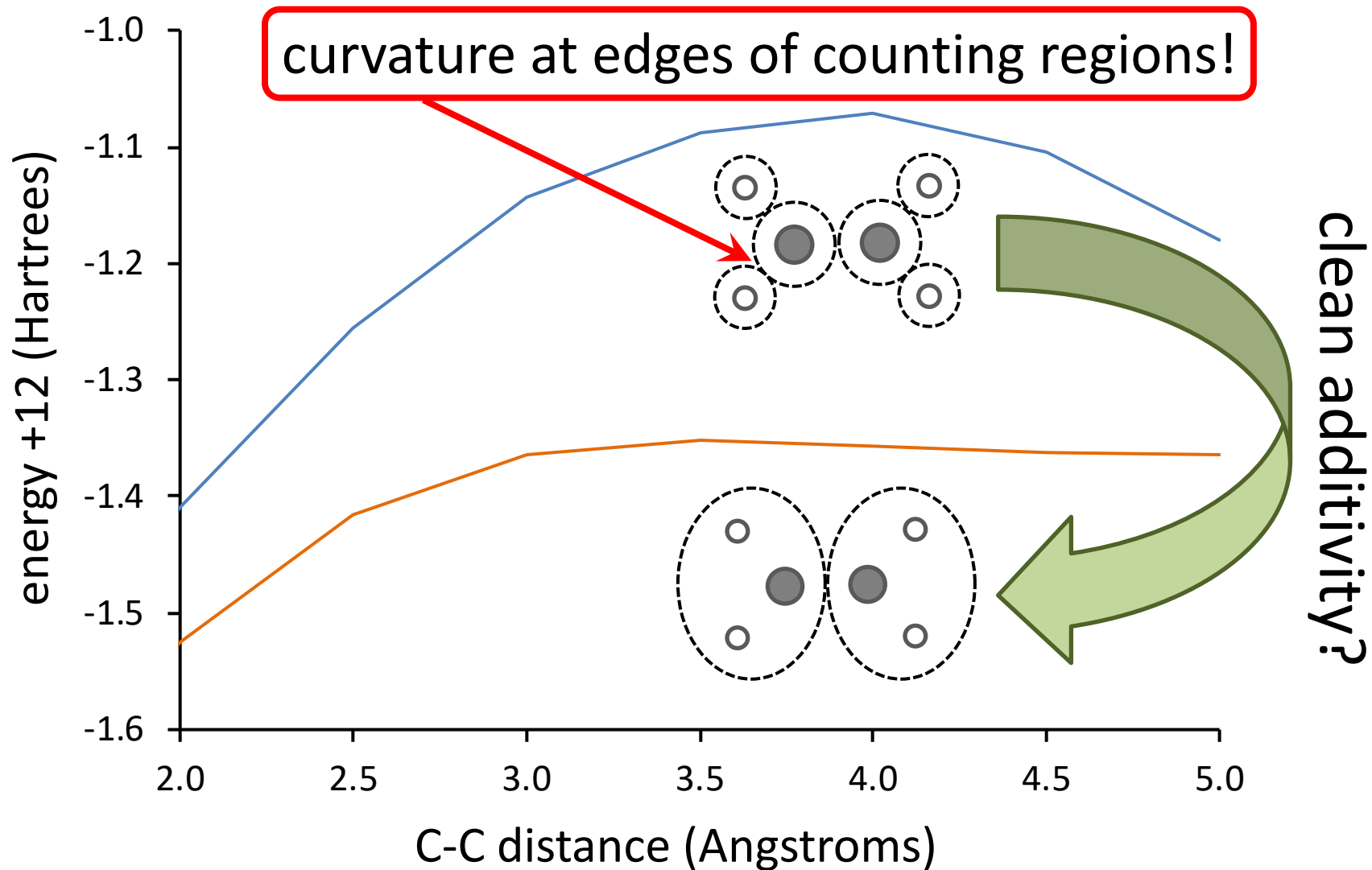
clean additivity



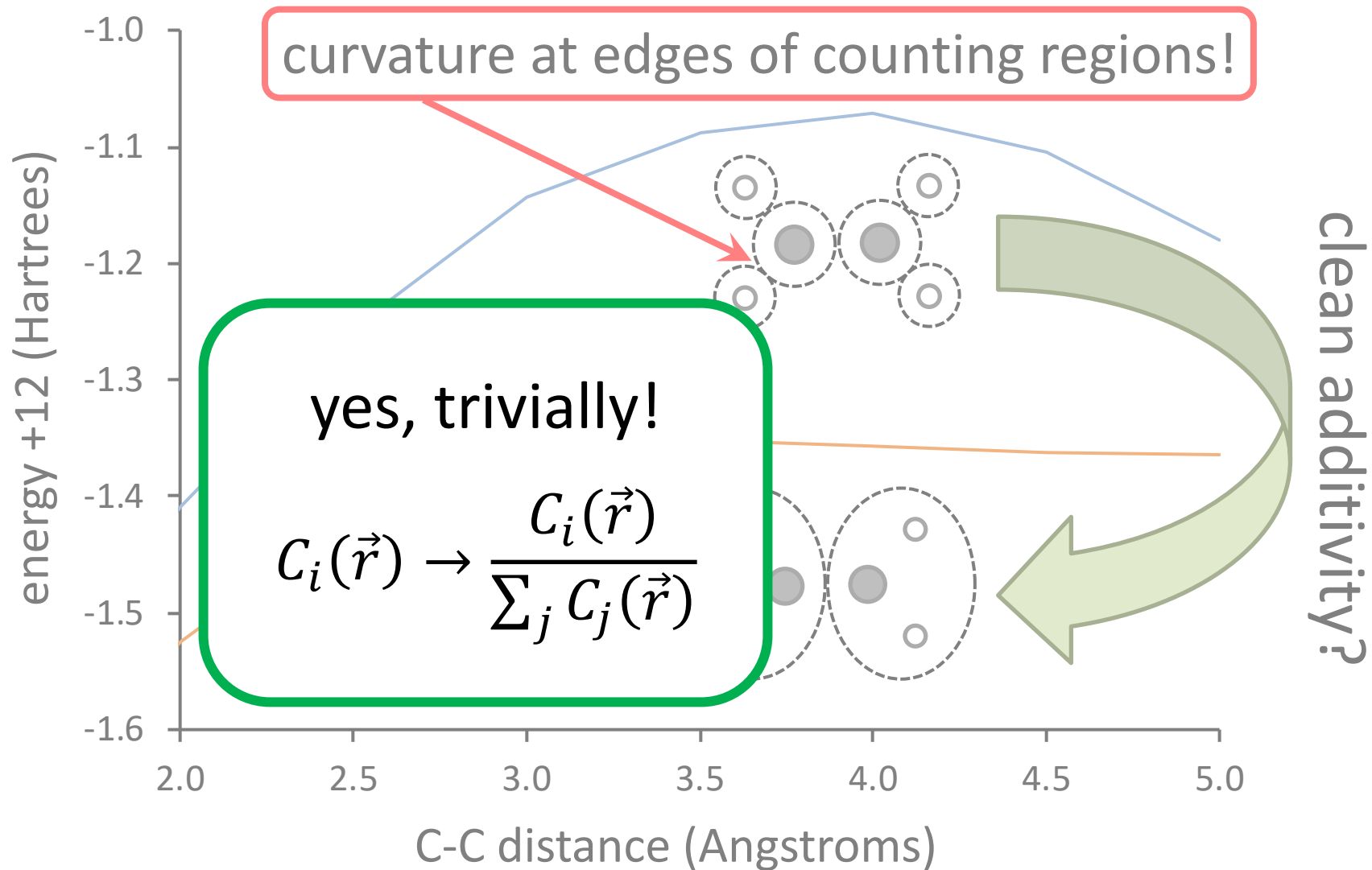
clean additivity



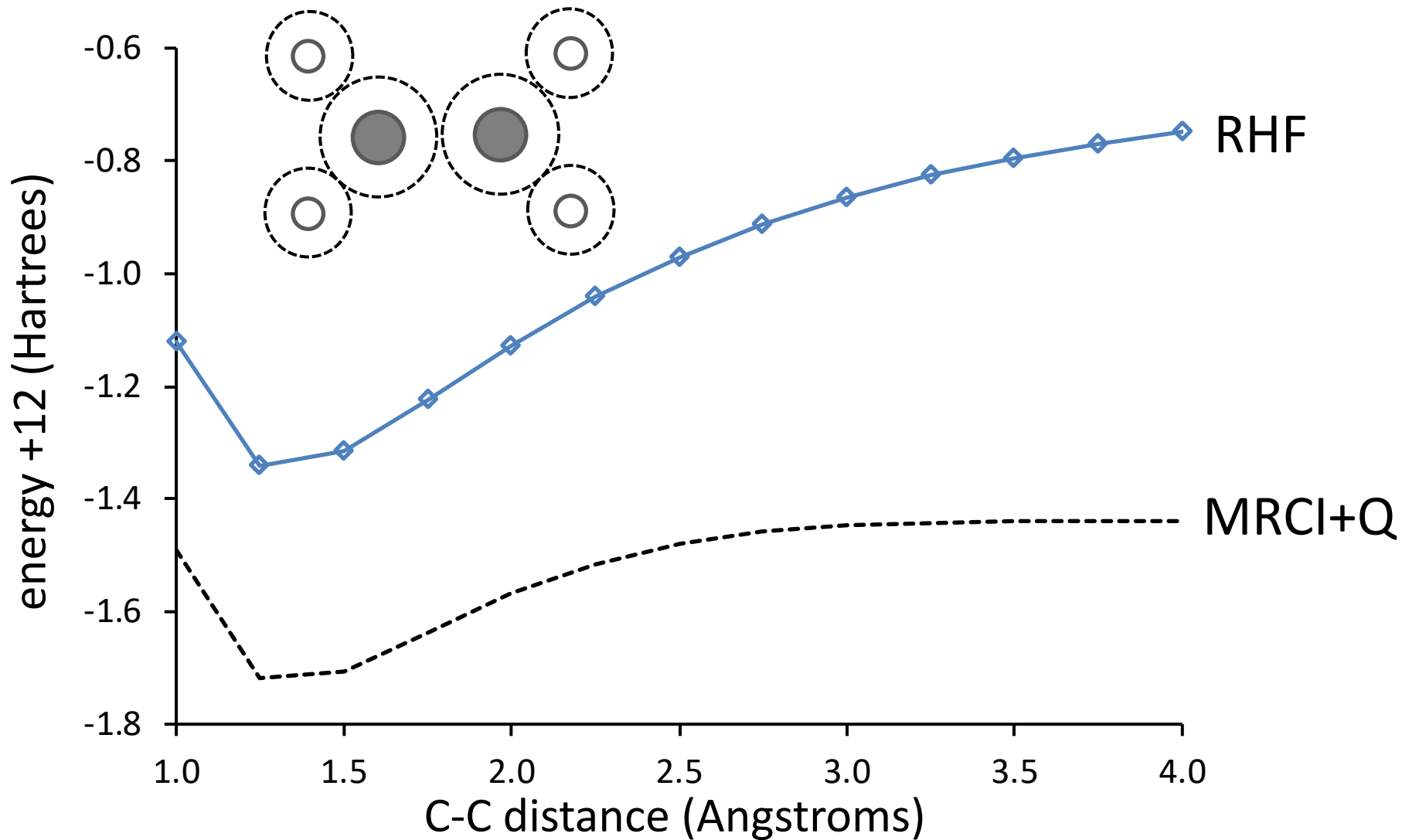
clean additivity



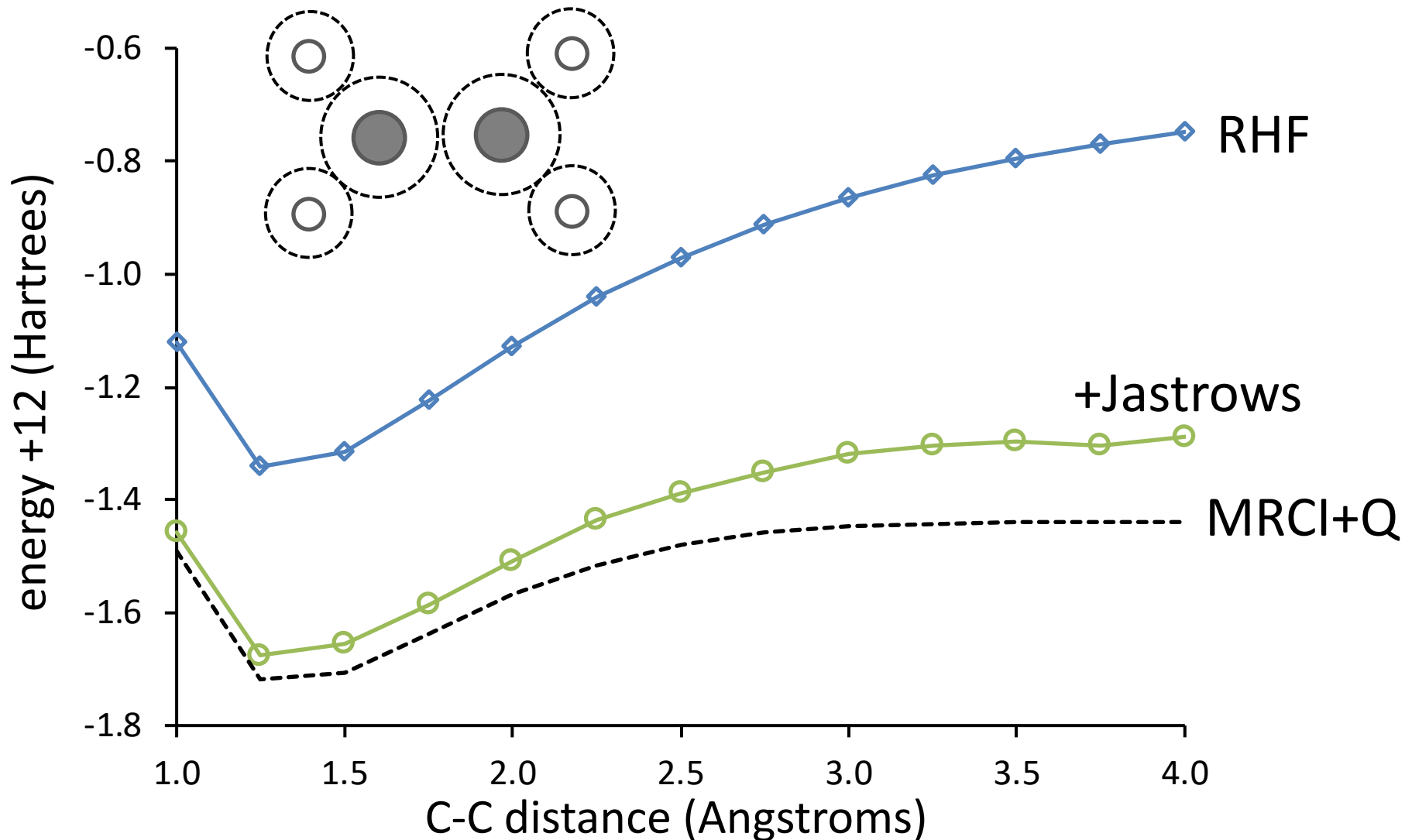
clean additivity



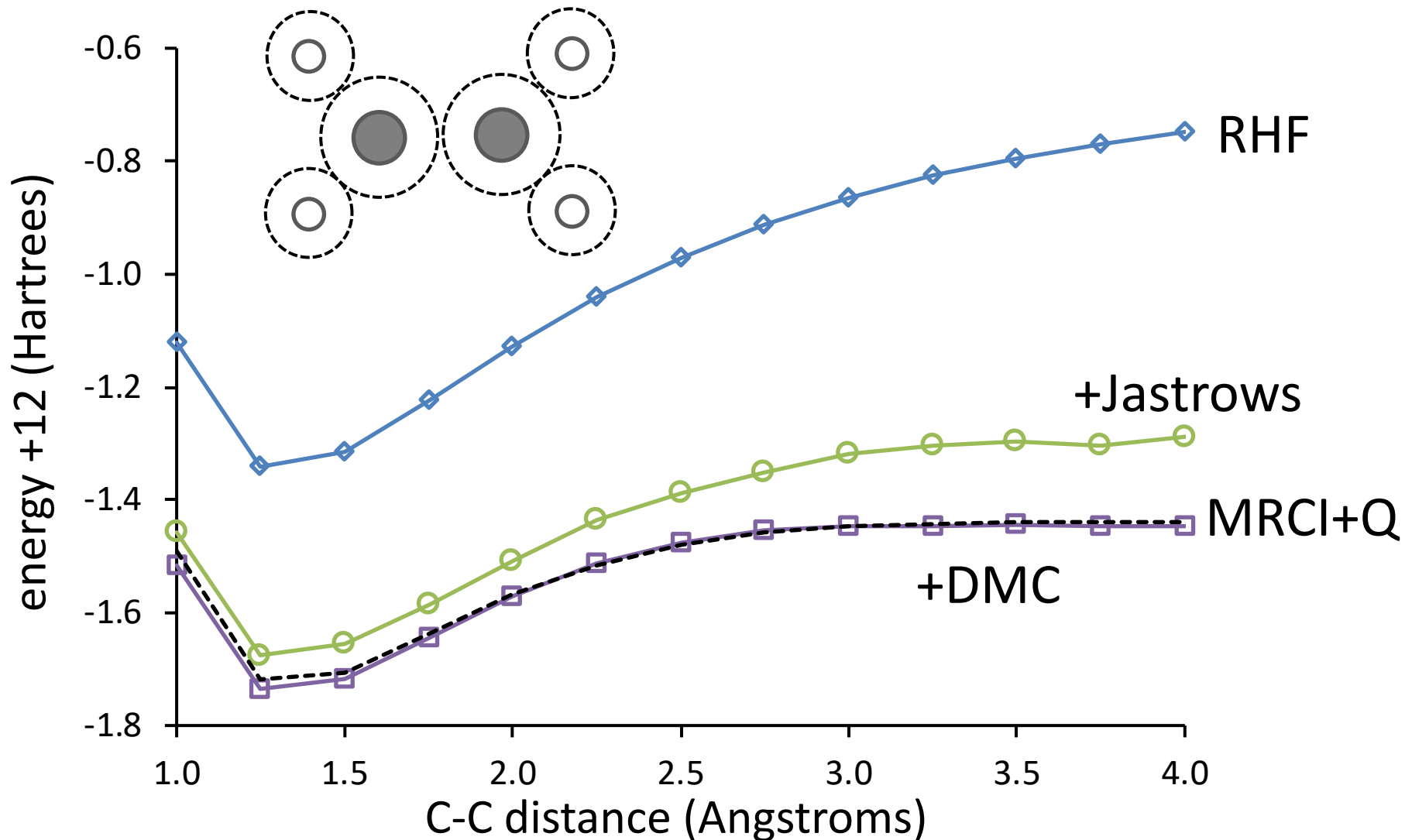
clean additivity works



clean additivity works



clean additivity works



the menu



aperitif: number counting Jastrows



main course: excited states with QMC



dessert: optimization methods

variational principles

ground state

$$\frac{\langle \Psi | H | \Psi \rangle}{\langle \Psi | \Psi \rangle}$$

any eigenstate

$$\frac{\langle \Psi | (H - E)^2 | \Psi \rangle}{\langle \Psi | \Psi \rangle}$$

Umrigar, Filippi, others

state near ω

$$\frac{\langle \Psi | (\omega - H)^2 | \Psi \rangle}{\langle \Psi | \Psi \rangle}$$

state above ω

$$\frac{\langle \Psi | (\omega - H) | \Psi \rangle}{\langle \Psi | (\omega - H)^2 | \Psi \rangle}$$

Choi, Lebeda, & Messmer, CPL 1970

Van Voorhis et al, JCP 2017

Zhao & Neuscamman, JCTC 2016, 2017

Shea & Neuscamman, JCTC 2017

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evaluate $\langle H^2 \rangle$ by Monte Carlo integration

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poor properties

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Umrigar, Filippi, others

state near ω

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not size-consistent

state above ω

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Choi, Lebeda, & Messmer, CPL 1970

Van Voorhis et al, JCP 2017

Zhao & Neuscamman, JCTC 2016, 2017

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interconversion

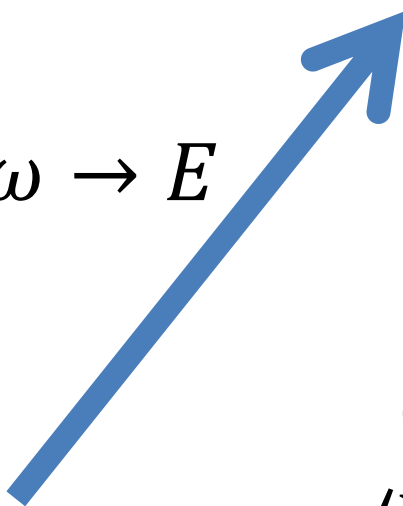
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$\omega \rightarrow E$



state near ω

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Umrigar, Filippi, others



$\omega \rightarrow E - \sigma$

state above ω

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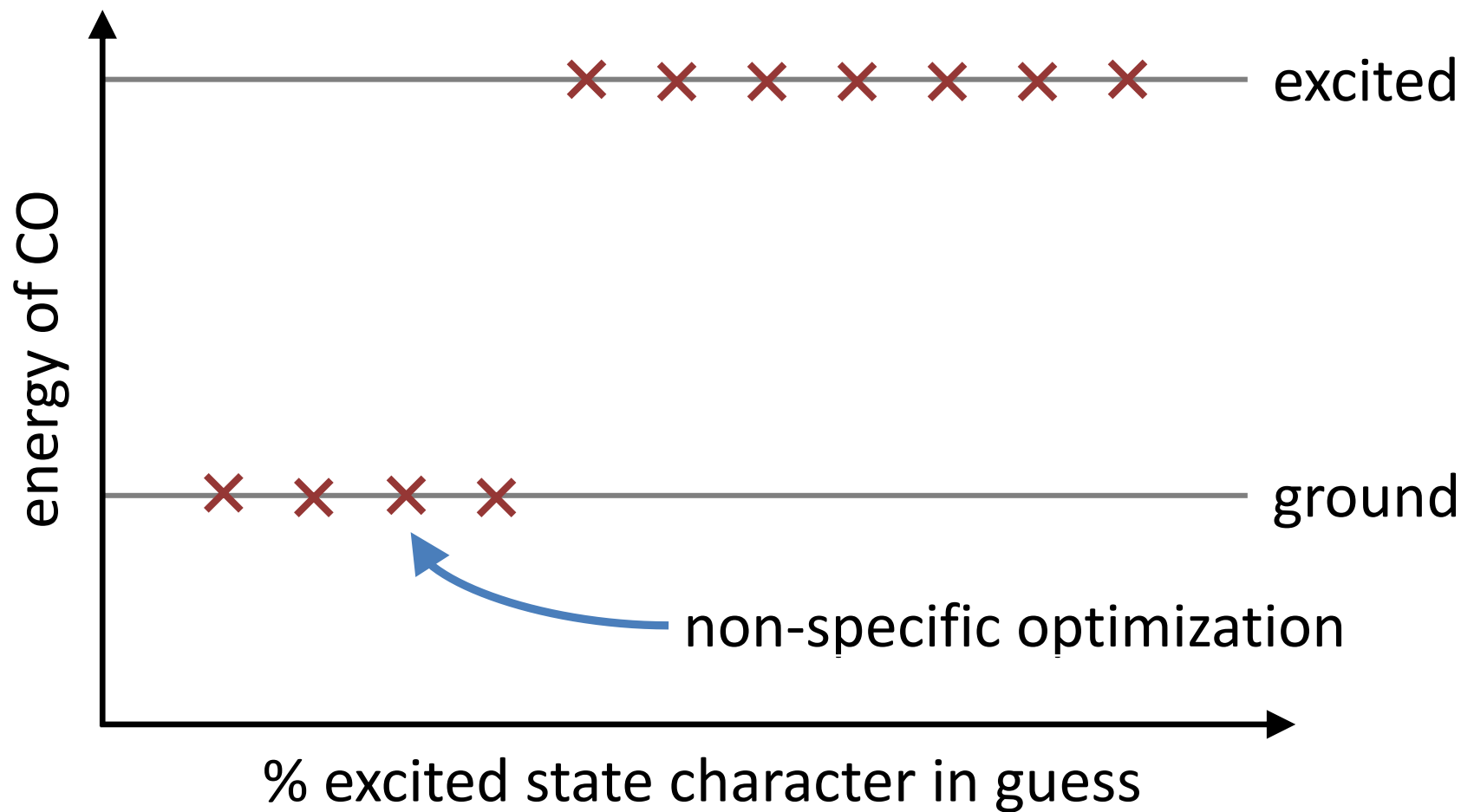
Choi, Lebeda, & Messmer, CPL 1970

Van Voorhis et al, JCP 2017

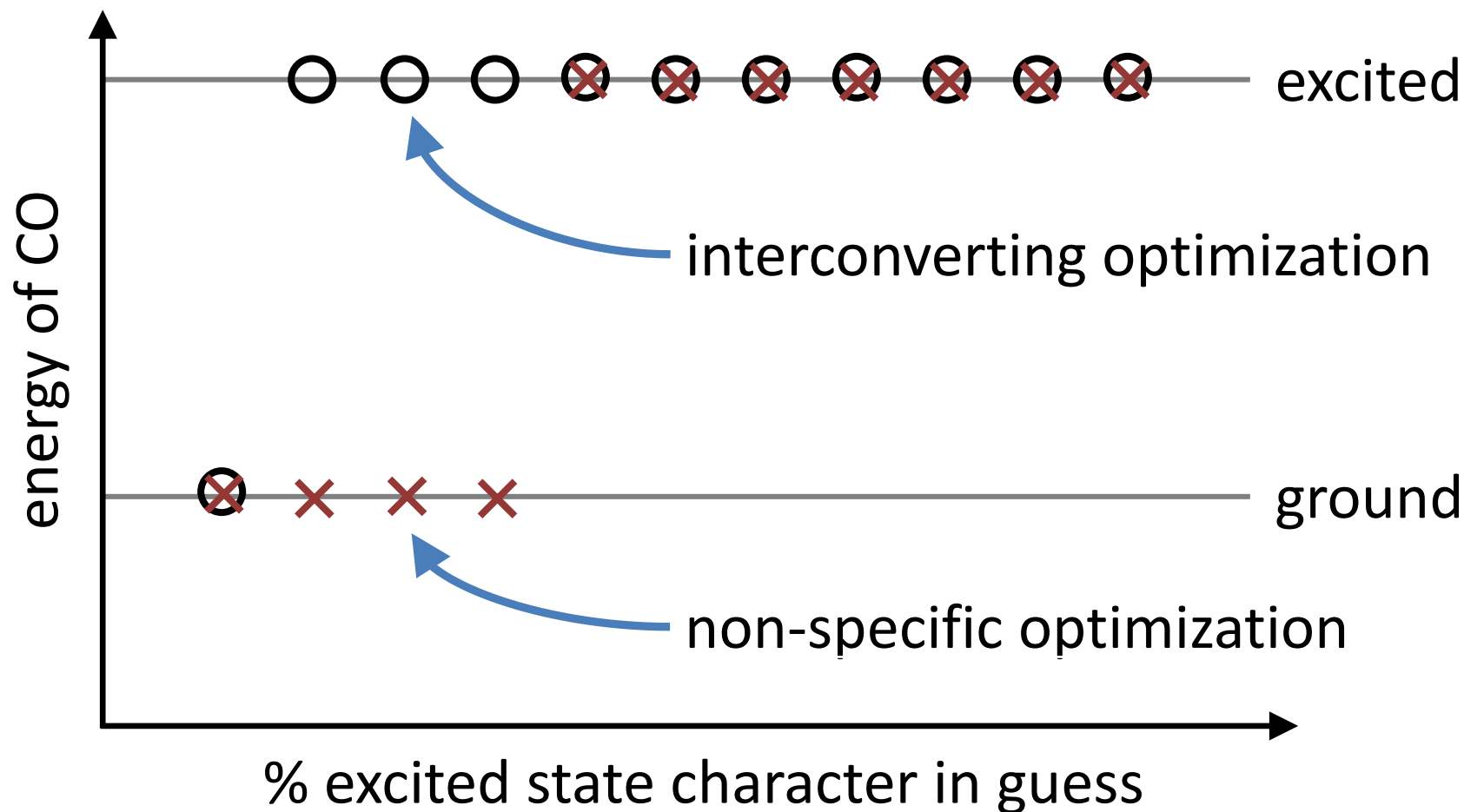
Zhao & Neuscamman, JCTC 2016, 2017

Shea & Neuscamman, JCTC 2017

having it both ways



having it both ways

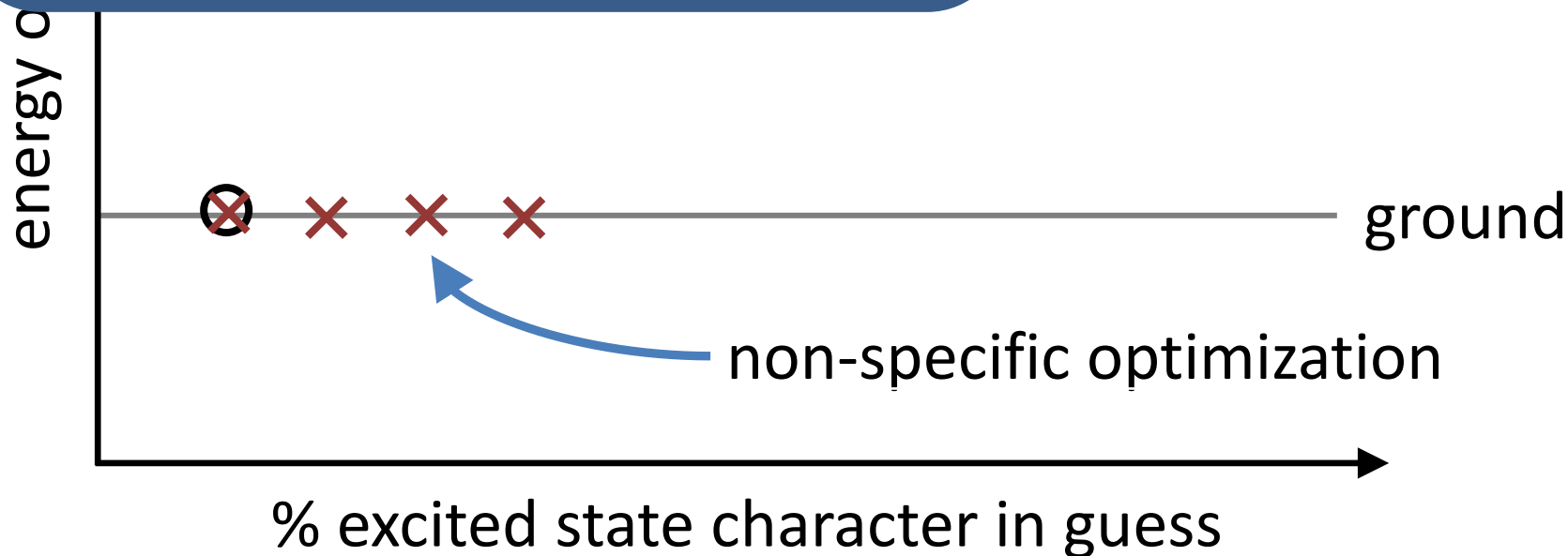


having it both ways

CO + N₂ size consistency error

fixed ω : $2.60 \pm 0.2 \text{ mE}_h$

adaptive ω : $0.04 \pm 0.2 \text{ mE}_h$



charge transfer orbital relaxation

state-averaged CASSCF:

2.7 eV to 5.9 eV (yikes!)

state-averaged CASPT2:

intruder states

state-specific CASSCF:

4.4 eV

state-specific CASPT2:

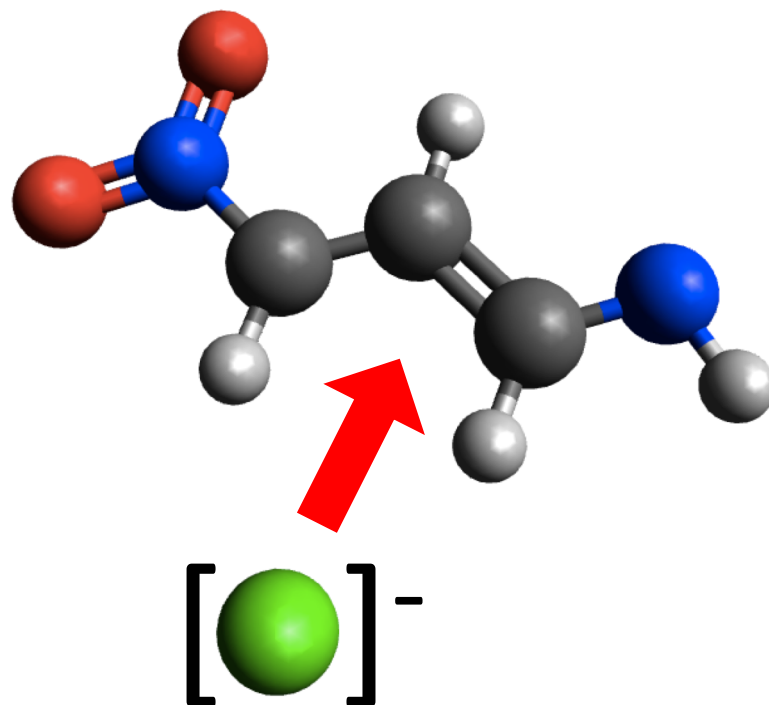
3.5 eV

EOM-CCSD:

3.8 eV

σ^2 -matched VMC:

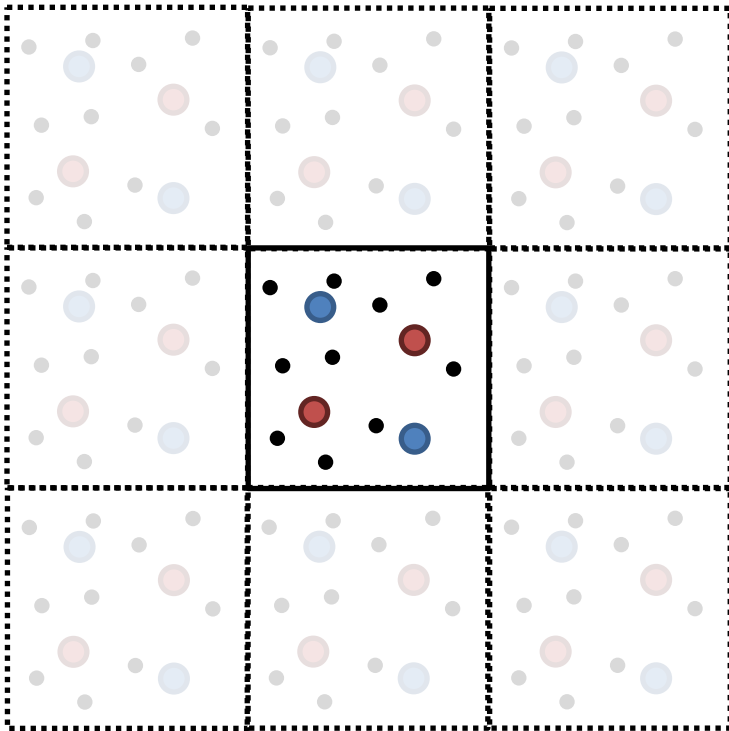
3.8 eV



Ψ_0 = multi-Slater from CISPI

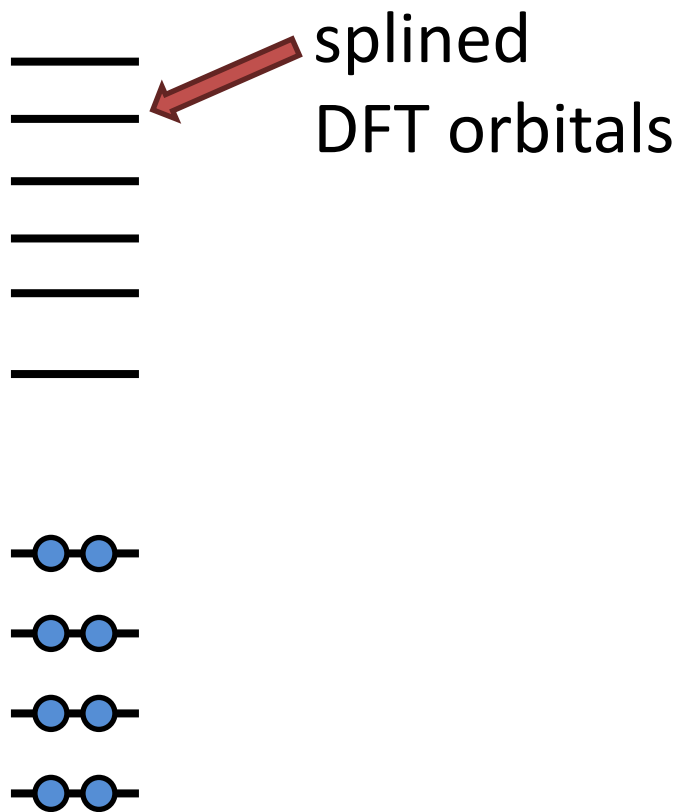
solids

in QMC, solids are molecules with funny electrons

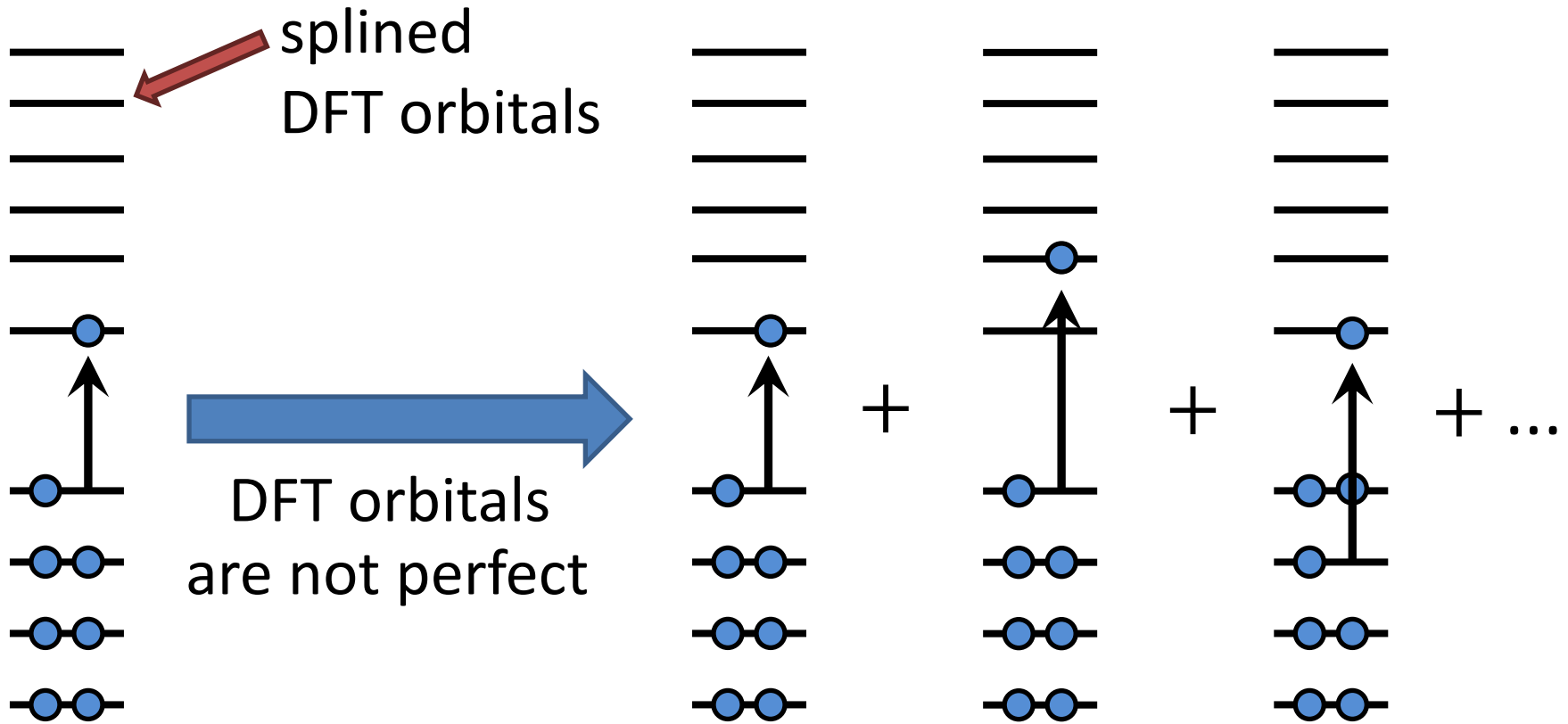


$$H = T + \underbrace{V_{nn} + V_{en} + V_{ee}}_{\text{periodic sums, but still a finite number of particles}}$$

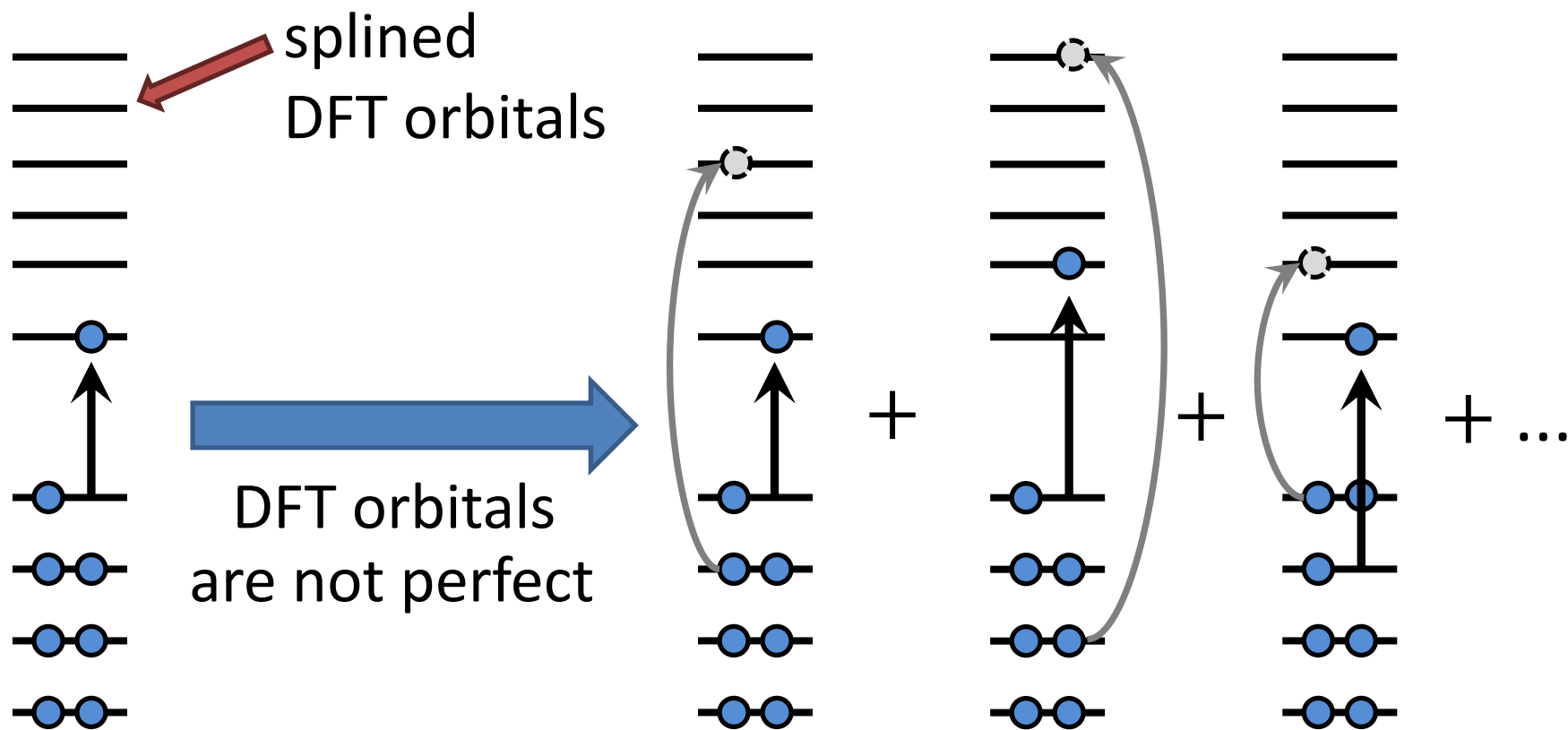
the wave function



the wave function

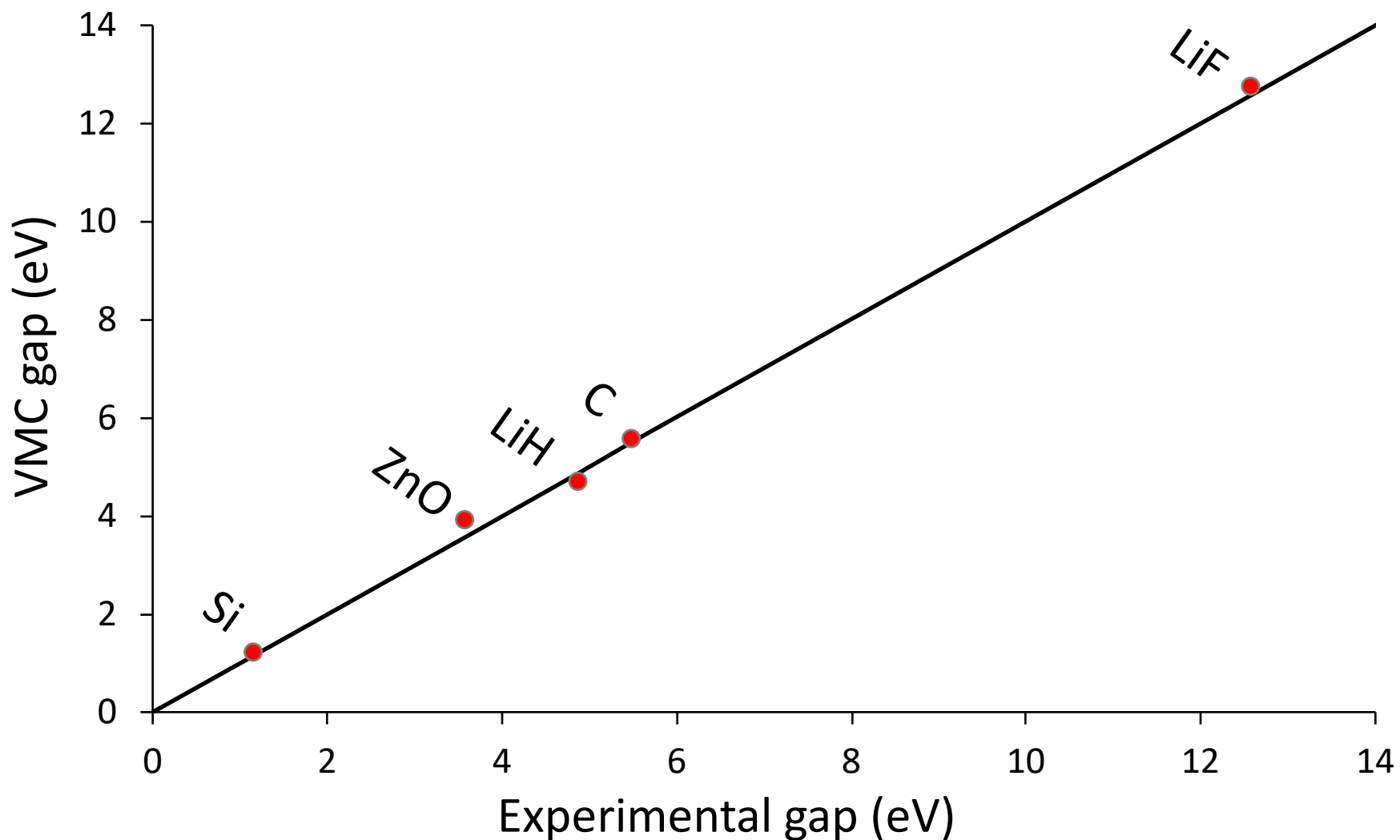


the wave function

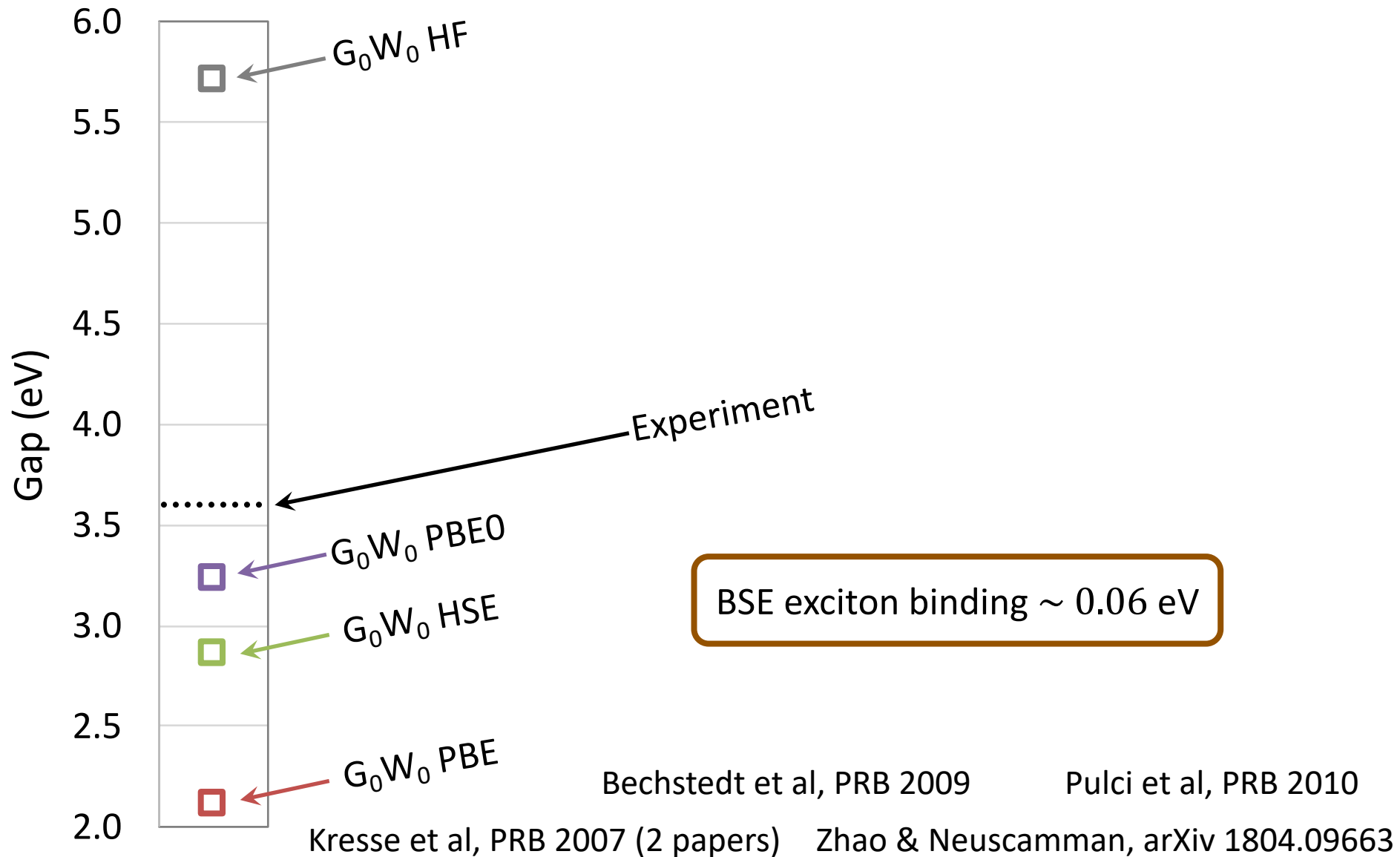


$\Psi = \text{all singles} + \text{doubles on important singles}$
(then cut off small singles to variance match)

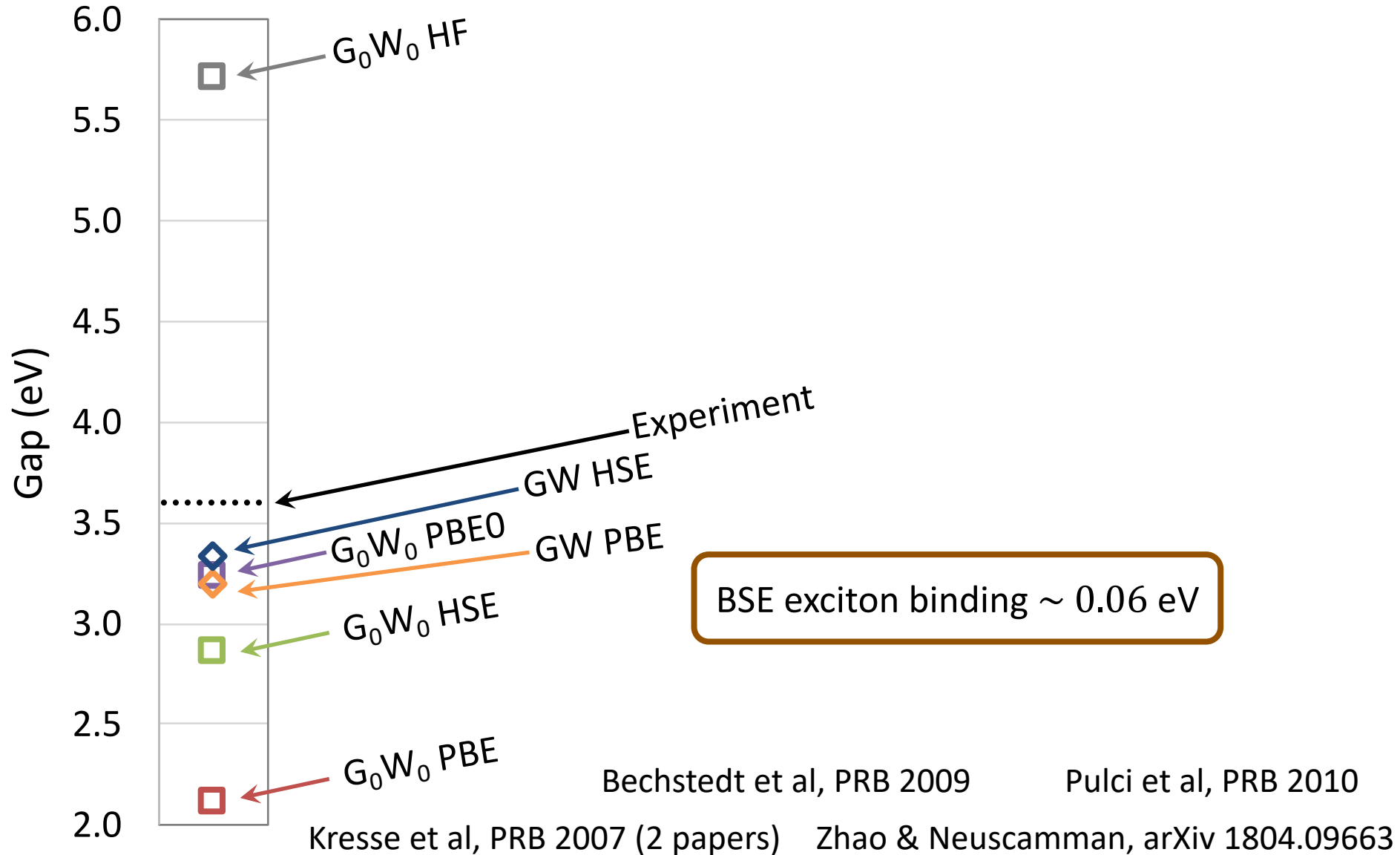
optical gaps



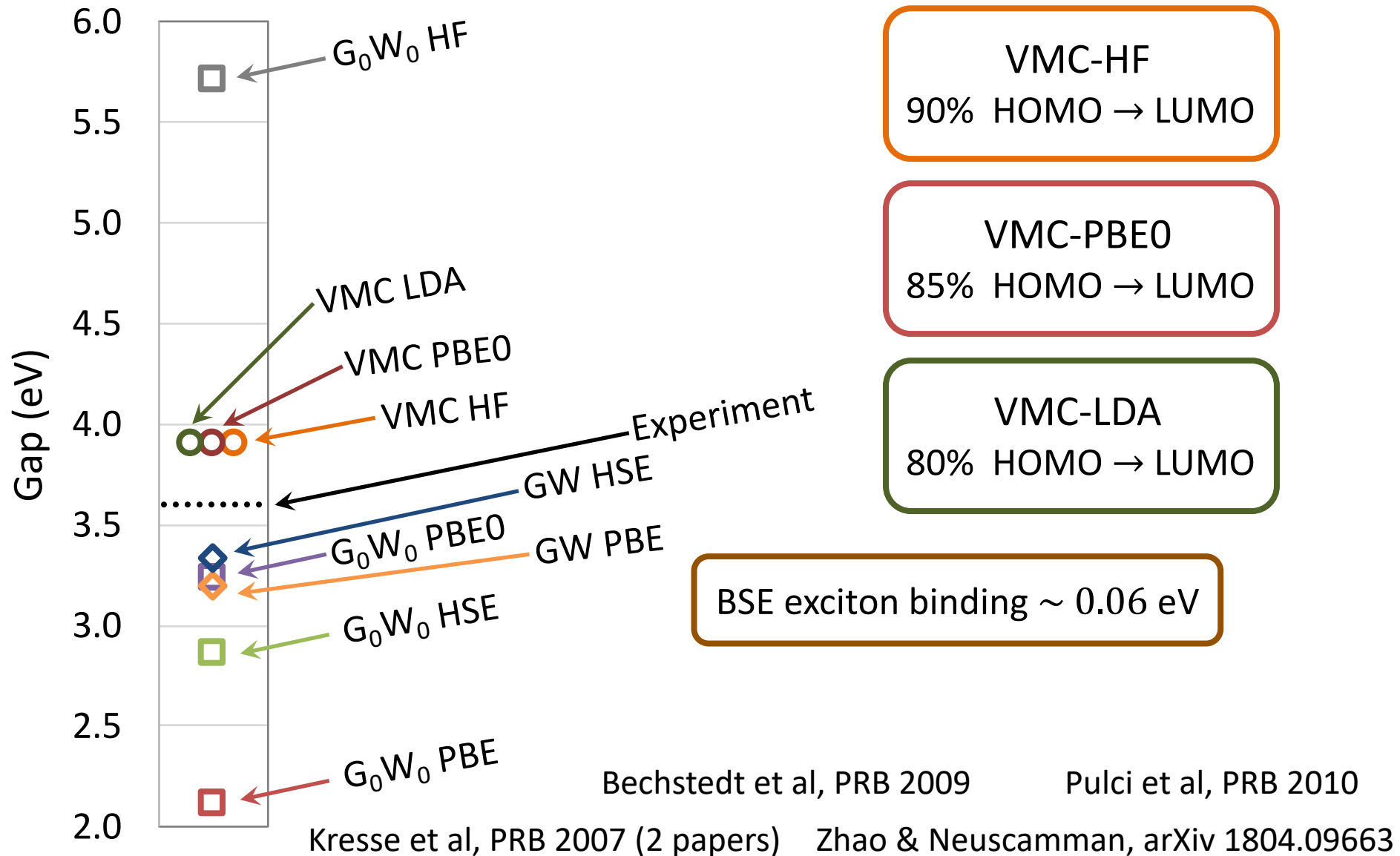
Zinc Oxide



Zinc Oxide



Zinc Oxide



the menu



aperitif: number counting Jastrows



main course: excited states with QMC



dessert: optimization methods

the linear method

$$\Psi(\vec{c}) \rightarrow a_0 \Psi(\vec{c}) + \sum_{i=1}^N a_i \frac{\partial \Psi}{\partial c_i}$$

minimize E w.r.t. \vec{a} , or, in other words, solve

$$\mathbf{H} \vec{a} = E \mathbf{S} \vec{a}$$

good:

fast convergence

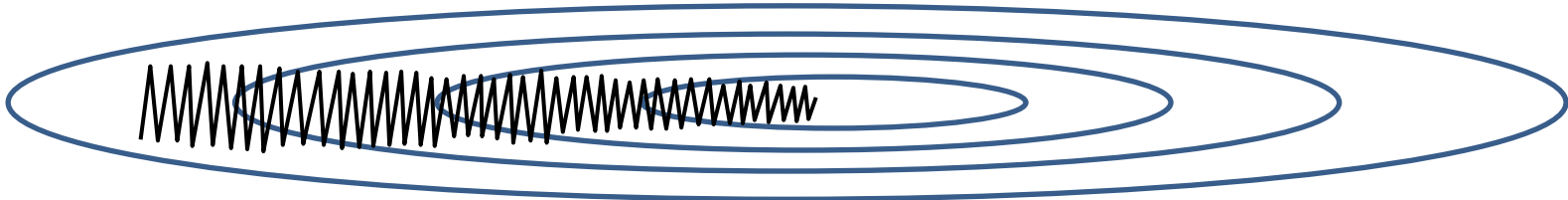
well established

bad:

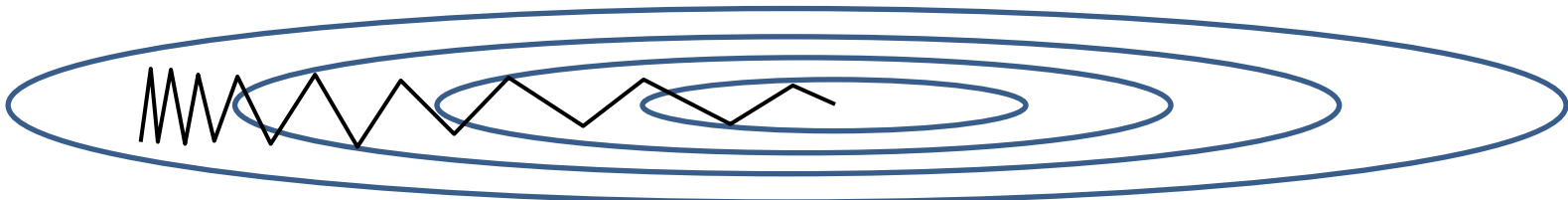
matrices get huge

highly nonlinear

accelerated descent



remember the average movement and add momentum...



good:

low memory

nearly linear

bad:

first order method

little VMC data so far

blocked linear method

$\Psi(\vec{c})$ with $\vec{c} = \{c_1, \dots, c_n, c_{n+1}, \dots, c_{2n}, c_{2n+1}, \dots, c_{3n}, \dots\}$

do linear method for first set...

blocked linear method

$\Psi(\vec{c})$ with $\vec{c} = \{c_1, \dots, c_n, c_{n+1}, \dots, c_{2n}, c_{2n+1}, \dots, c_{3n}, \dots\}$

do linear method for first set...

then second set...

blocked linear method

$\Psi(\vec{c})$ with $\vec{c} = \{c_1, \dots, c_n, c_{n+1}, \dots, c_{2n}, c_{2n+1}, \dots, c_{3n}, \dots\}$

do linear method for first set...

then second set...

then third set...

blocked linear method

$\Psi(\vec{c})$ with $\vec{c} = \{c_1, \dots, c_n, c_{n+1}, \dots, c_{2n}, c_{2n+1}, \dots, c_{3n}, \dots\}$

do linear method for first set...

then second set...

then third set...

now restrict variables to best directions from each set

do LM on $\Psi(\vec{c}) \rightarrow \Psi(\vec{z})$ with \vec{z} much shorter than \vec{c}

good:

fast convergence

low memory

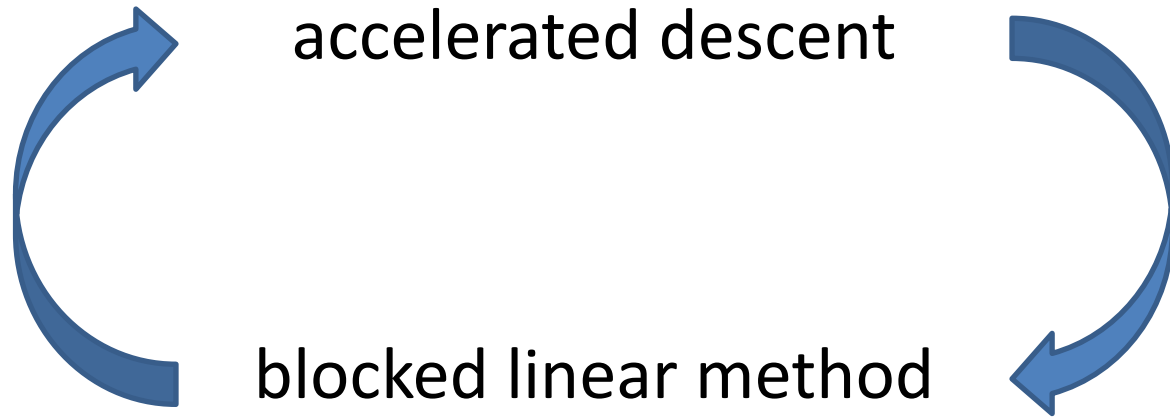
bad:

extra samples

little VMC data so far

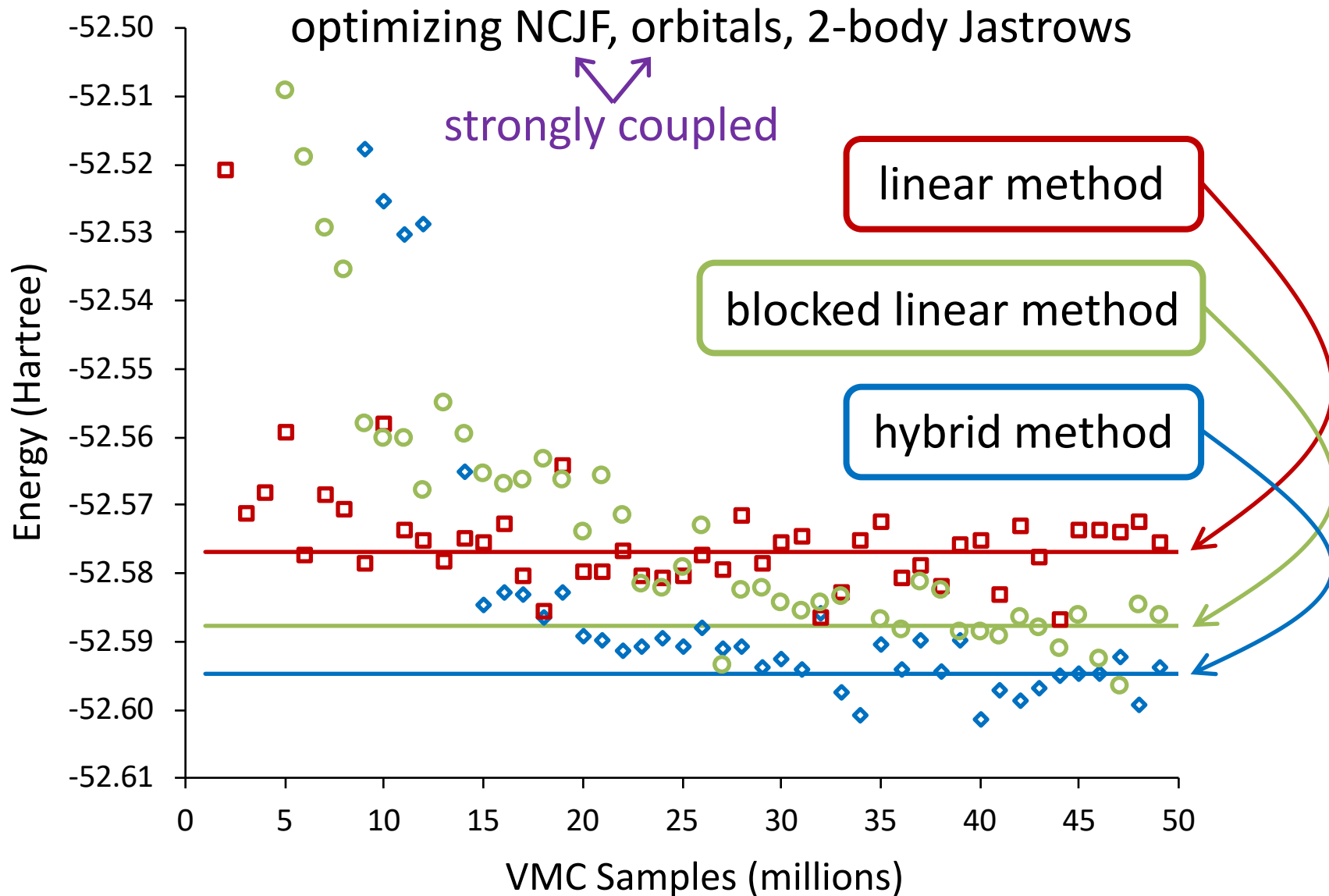
hybrid method?

hard to couple stiff parameters w/o 2nd order derivatives



needs info about other blocks to get coupling right

comparing methods in CaO



future directions

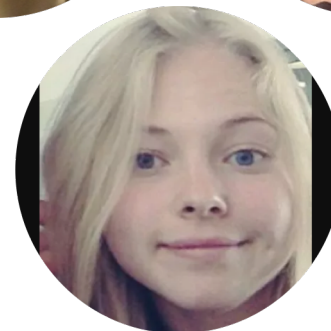
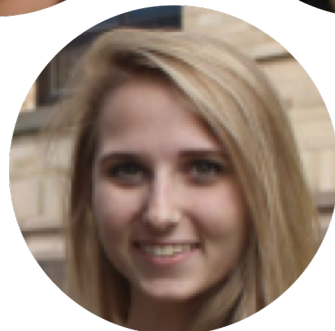
oo-MSJ in solids

core excitations

selective CI \rightarrow VMC \rightarrow DMC

black-box hybrid optimizer

acknowledgements



CPSFM

Center for Predictive Simulation
of Functional Materials

NERSC



- Gas Phase Chemical Physics
- Computational & Theoretical Chemistry