

Ab initio computations of the nuclear spectral function



Vittorio Somà
CEA Saclay, France

Neutrino detection and interactions: challenges and opportunities for ab initio nuclear theory

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Collaborators:

Carlo Barbieri, Francesco Raimondi, Arnau Rios (University of Surrey),
Thomas Duguet (CEA Saclay), Petr Navrátil (TRIUMF)

Outline

- ◎ **Introduction**
- ◎ **Self-consistent Green's function approach**
- ◎ **Applications**
 - Ground-state properties
 - Spectral function
 - Response function
- ◎ **Conclusions**

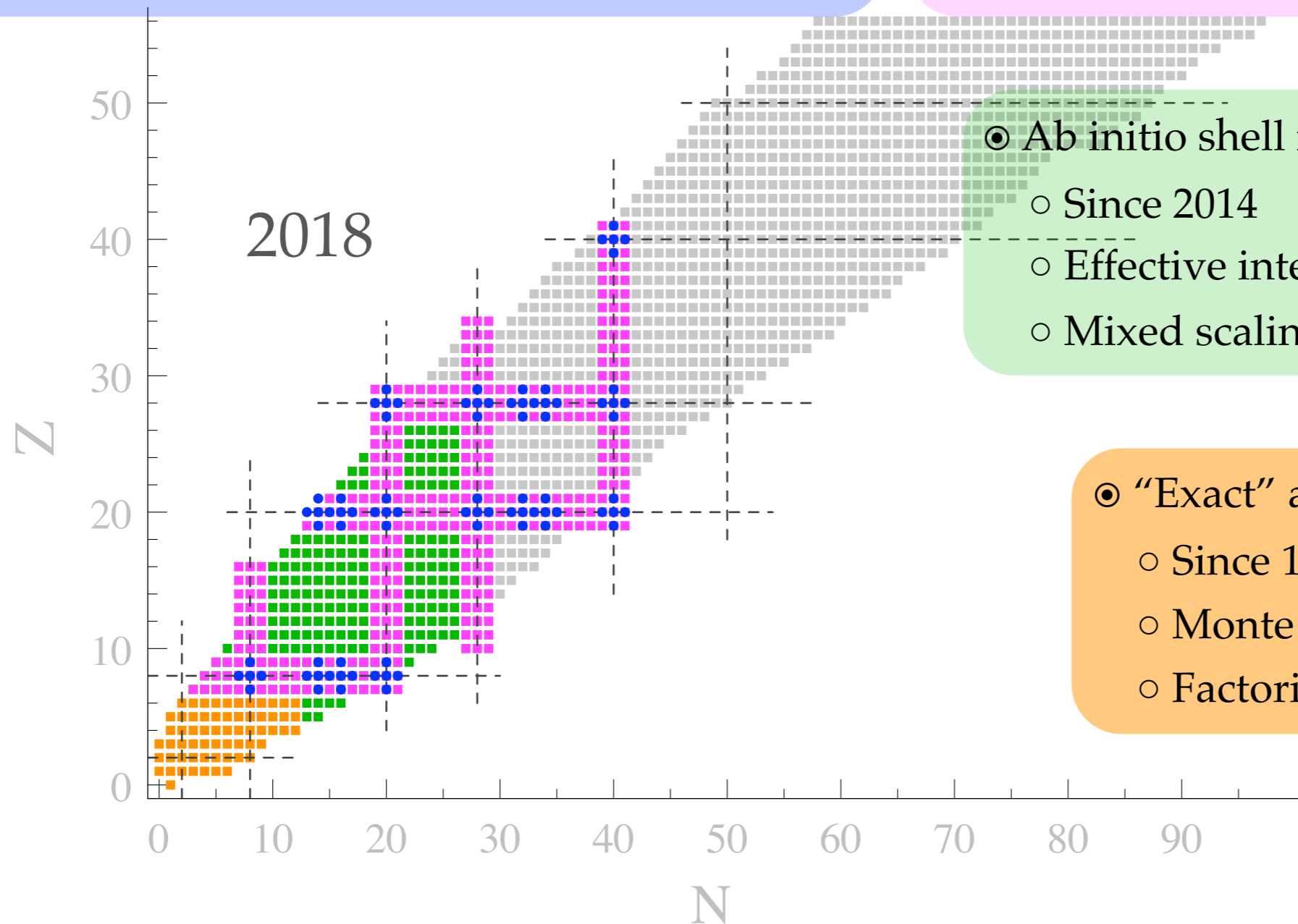
Evolution of ab initio nuclear chart

Approximate approaches for closed-shell nuclei

- Since 2000's
- SCGF, CC, IMSRG
- Polynomial scaling

Approximate approaches for open-shells

- Since 2010's
- GGF, BCC, MR-IMSRG
- Polynomial scaling



Ab initio shell model

- Since 2014
- Effective interaction via CC/IMSRG
- Mixed scaling

"Exact" approaches

- Since 1980's
- Monte Carlo, CI, ...
- Factorial scaling

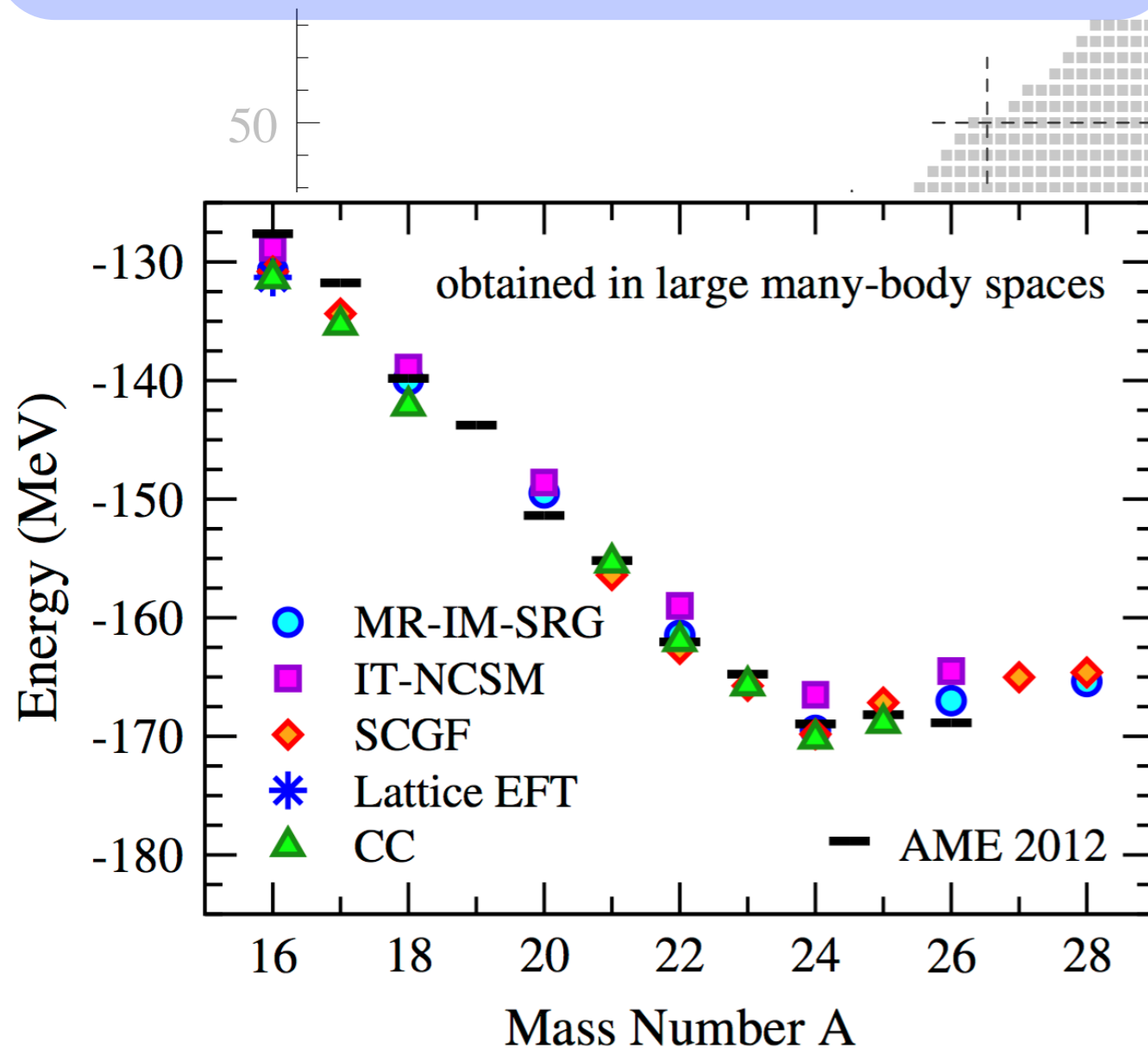
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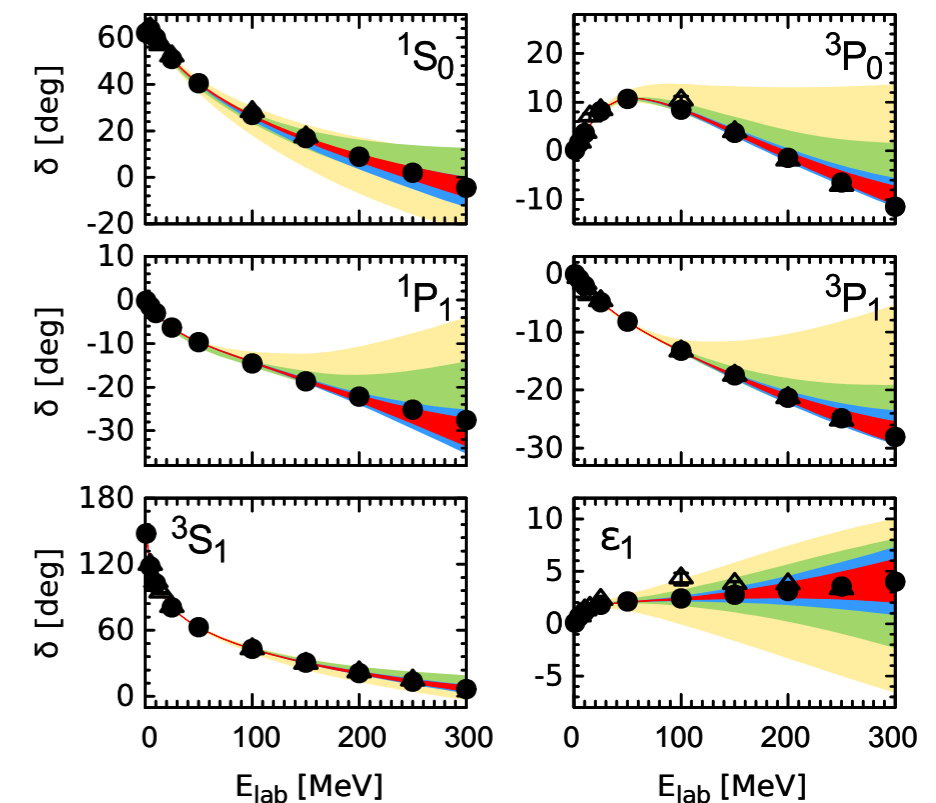
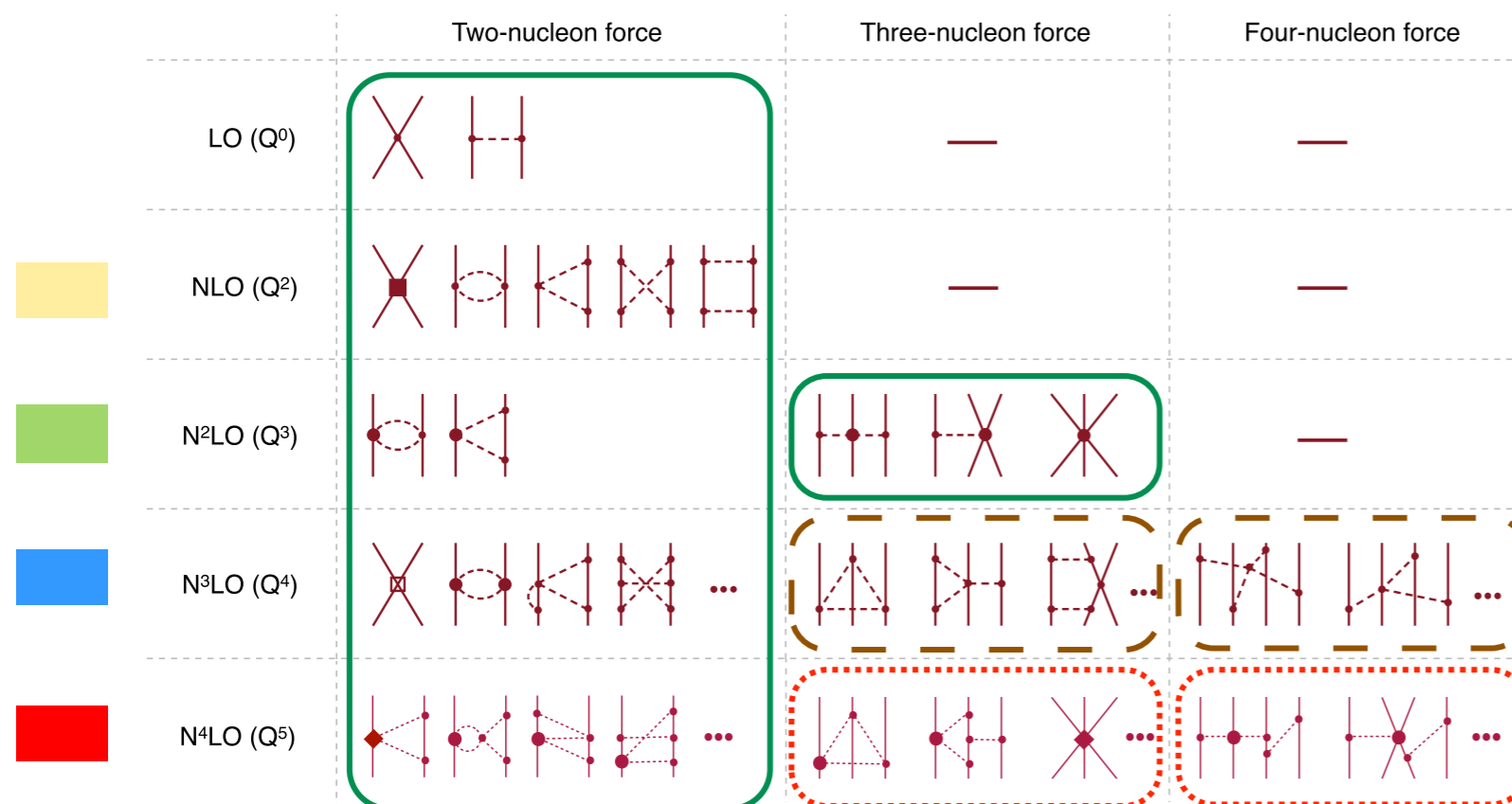
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Chiral effective field theory & nuclear interactions

Chiral EFT provides a **systematic** framework to construct AN interactions ($A=2, 3, \dots$)

Main features:

- High-energy physics unresolved \rightarrow **soft potentials** \rightarrow improved many-body convergence
- Many-body forces and currents consistently derived
- A **theoretical error** can be, in principle, assigned to each order in the expansion



[Meißner 2016]

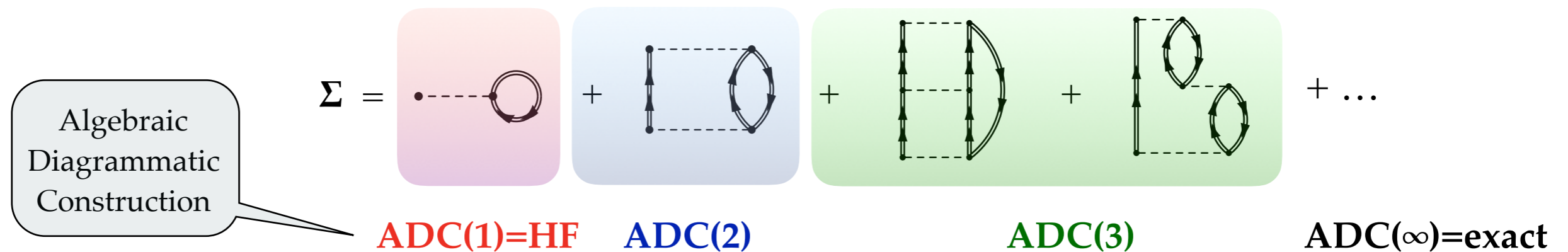
\Rightarrow Ideally: apply to the many-nucleon system (and propagate the theoretical error)

Self-consistent Green's function approach

◎ **Solution of the A -body Schrödinger equation** $H|\Psi_k^A\rangle = E_k^A|\Psi_k^A\rangle$ achieved by

- 1) Rewriting it in terms of 1-, 2-, ... A -body objects $G_1=G, G_2, \dots G_A$ (**Green's functions**)
- 2) Expanding these objects in perturbation (in practise $\mathbf{G} \mapsto$ **one-body observables**, etc..)
 - **Self-consistent** schemes resum (infinite) subsets of perturbation-theory contributions

◎ **Self-energy expansion**



◎ **Access a variety of quantities**

- One-body GF \rightarrow Ground-state properties of even-even A + spectra of odd-even neighbours
- Two-body GF \rightarrow Excited spectrum of even-even A
- Self-energy \rightarrow Optical potential for nucleon-nucleus scattering

Gorkov-Green's functions for open-shell systems

⊙ Standard expansion schemes fail to account for superfluidity

⊙ **Gorkov scheme generalises GF theory to superfluid systems**

[Gorkov 1958]

○ Use **symmetry breaking** (particle number) to effectively include pairing correlations

○ Start expansion from symmetry-breaking reference $|\Psi_0\rangle \equiv \sum_A^{\text{even}} c_A |\psi_0^A\rangle$

○ 4 one-body Gorkov propagators

$$\mathbf{G}_{ab} = \begin{pmatrix} G_{ab}^{11} & G_{ab}^{12} \\ G_{ab}^{21} & G_{ab}^{22} \end{pmatrix} = \left(\begin{array}{c|c} \uparrow & \downarrow \\ \hline \downarrow & \uparrow \end{array} \right)$$

○ Symmetry must be eventually restored

⊙ **Current implementation: ADC(2)**

[Somà, Duguet & Barbieri 2011]

$$\Sigma_{ab}^{11(1)} = \begin{array}{c} a \\ \bullet \\ b \end{array} \text{---} \begin{array}{c} c \\ \bullet \\ d \end{array} \begin{array}{c} \circlearrowleft \\ \circlearrowright \end{array} \downarrow \omega'$$

$$\Sigma_{ab}^{12(1)} = \begin{array}{c} a \\ \bullet \\ c \end{array} \text{---} \begin{array}{c} \bar{b} \\ \bullet \\ \bar{d} \end{array} \begin{array}{c} \curvearrowright \\ \curvearrowleft \end{array} \leftarrow \omega'$$

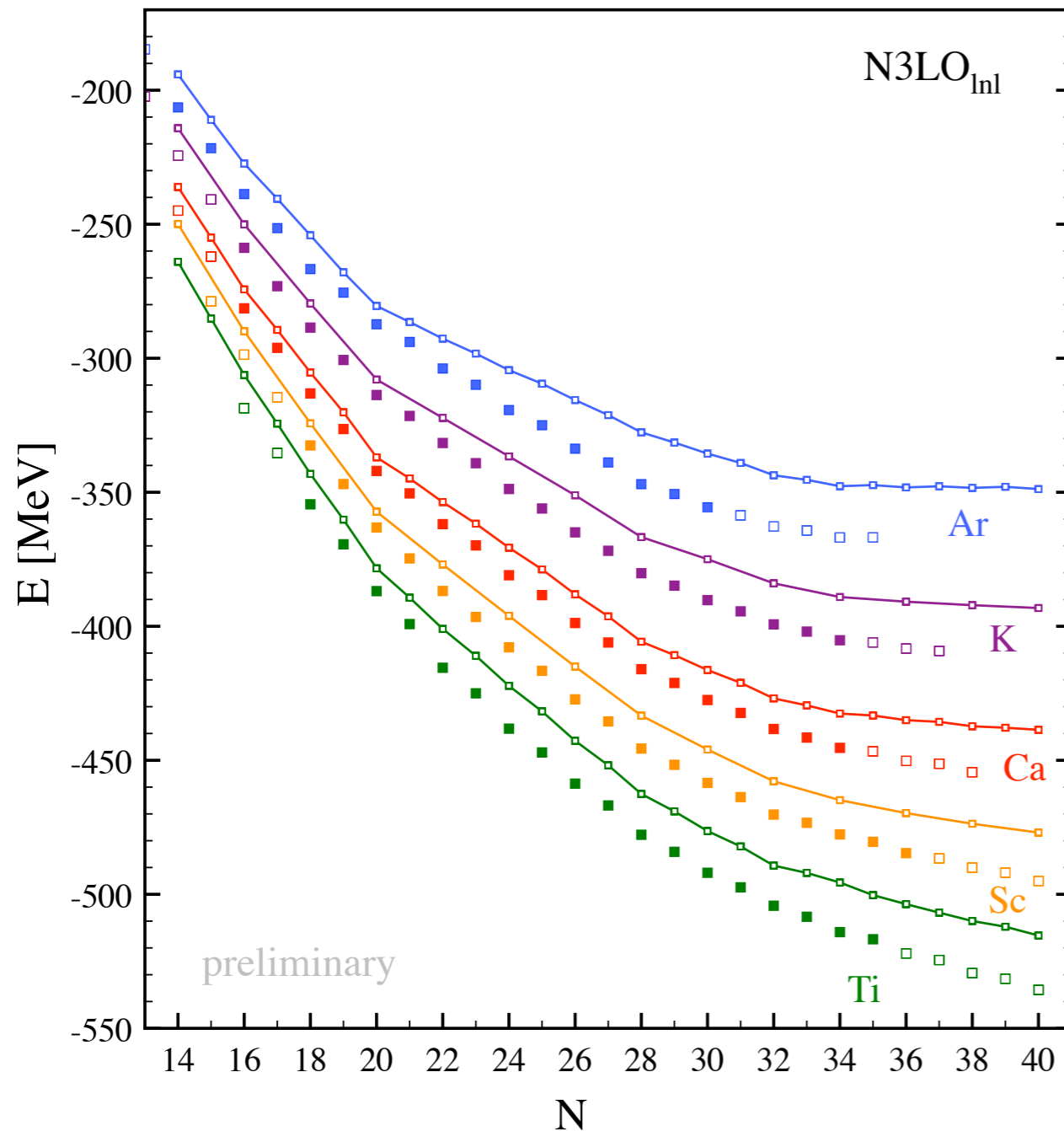
$$\Sigma_{ab}^{11(2)}(\omega) = \begin{array}{c} a \\ \bullet \\ c \\ \bullet \\ d \\ \bullet \\ b \end{array} \begin{array}{c} e \\ \bullet \\ f \\ \bullet \\ g \\ \bullet \\ h \end{array} \begin{array}{c} \uparrow \omega' \\ \uparrow \omega'' \\ \downarrow \omega''' \end{array} + \begin{array}{c} a \\ \bullet \\ c \\ \bullet \\ d \\ \bullet \\ b \end{array} \begin{array}{c} e \\ \bullet \\ f \\ \bullet \\ g \\ \bullet \\ \bar{h} \end{array} \begin{array}{c} \uparrow \omega' \\ \uparrow \omega'' \\ \uparrow \omega''' \end{array}$$

$$\Sigma_{ab}^{12(2)}(\omega) = \begin{array}{c} a \\ \bullet \\ c \end{array} \begin{array}{c} e \\ \bullet \\ f \\ \bullet \\ g \\ \bullet \\ h \end{array} \begin{array}{c} \uparrow \omega'' \\ \downarrow \omega''' \end{array} \begin{array}{c} \bar{b} \\ \bullet \\ \bar{d} \end{array} \begin{array}{c} \leftarrow \omega' \\ \leftarrow \omega' \end{array} + \begin{array}{c} a \\ \bullet \\ c \end{array} \begin{array}{c} e \\ \bullet \\ f \\ \bullet \\ g \\ \bullet \\ \bar{h} \end{array} \begin{array}{c} \downarrow \omega''' \\ \leftarrow \omega'' \end{array} \begin{array}{c} \bar{b} \\ \bullet \\ \bar{d} \end{array} \begin{array}{c} \leftarrow \omega' \\ \leftarrow \omega' \end{array}$$

Systematics of medium-mass nuclei

◎ Calculation of complete mid-mass isotopic chains possible

○ Chiral N3LO 2N (500) + N2LO 3N (400 local/non-local) interaction, further SRG-ed to 2.0 fm⁻¹



◎ Total energies

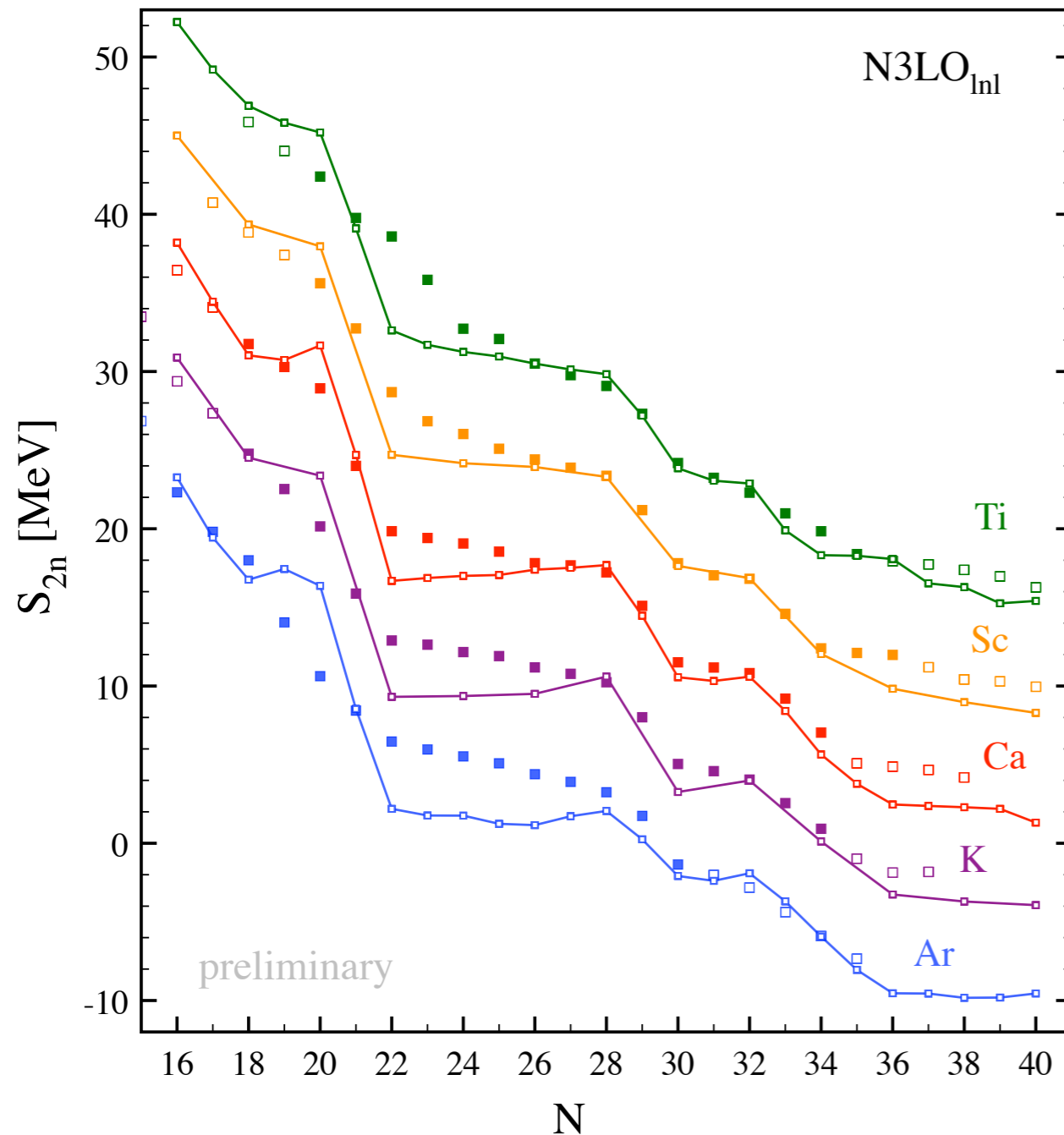
○ Missing correlation energy from higher-order diagrams

○ Overall trends reasonable

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◎ **Total energies**

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◎ **Energy differences (2N separation energies)**

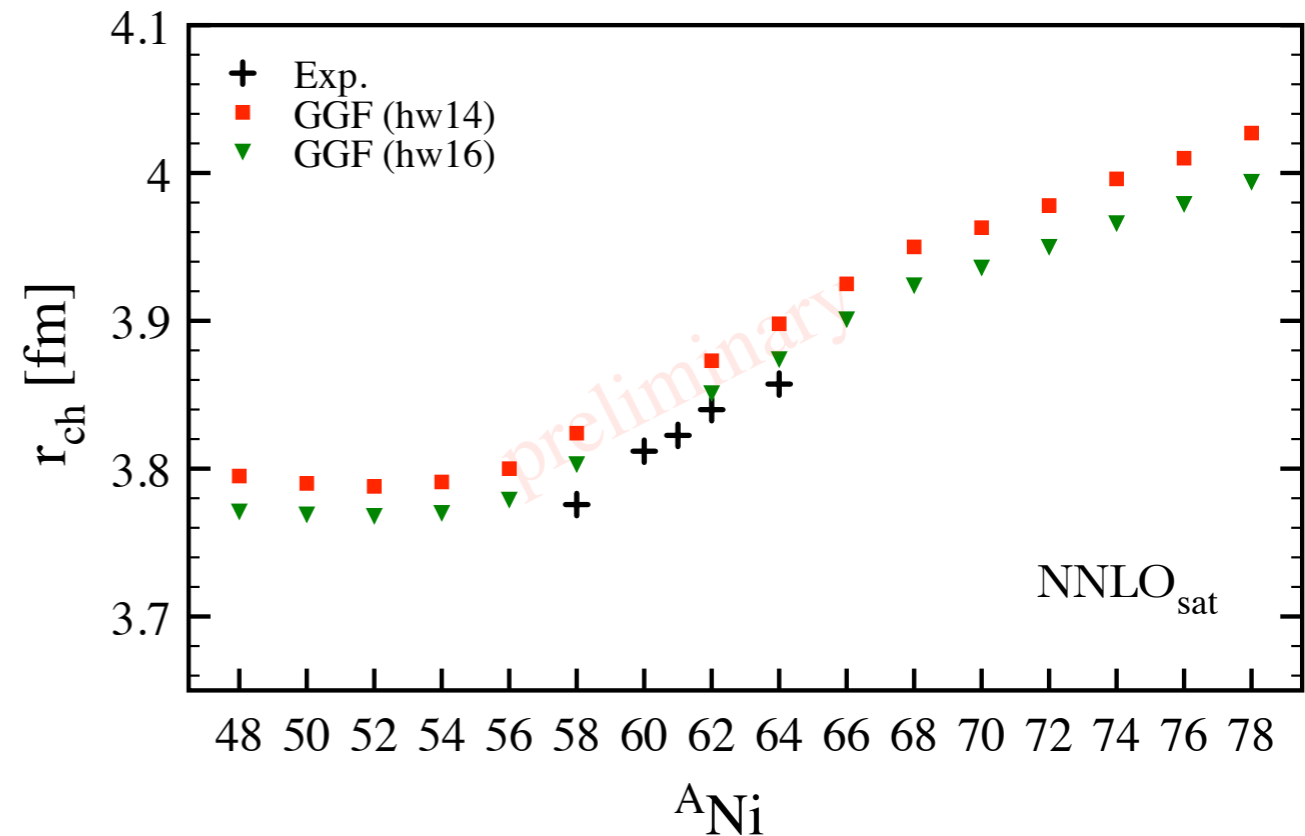
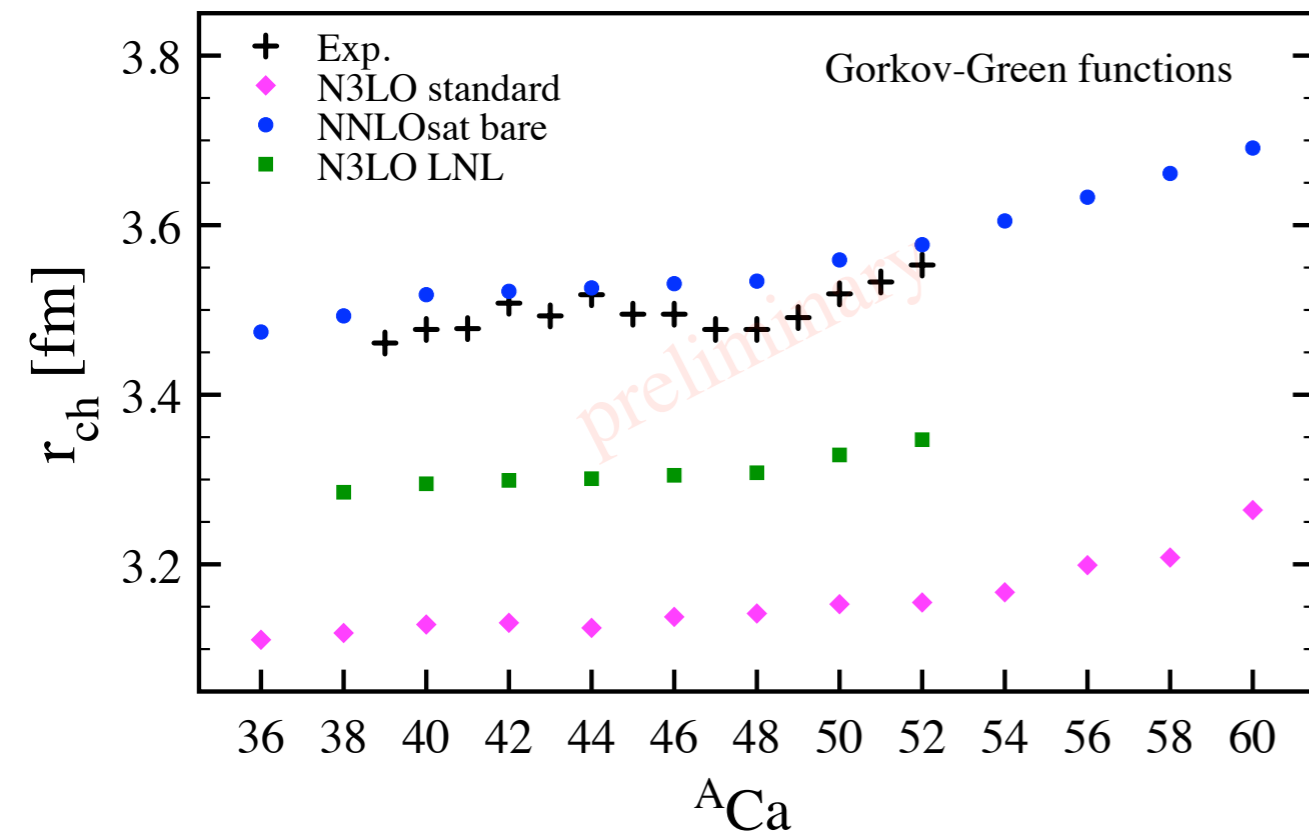
○ $N=20$ gap overestimated

○ $N=28$ gap OK, pf shell well reproduced

○ Drip lines?

Systematics of medium-mass nuclei

Charge radii along calcium and nickel chains

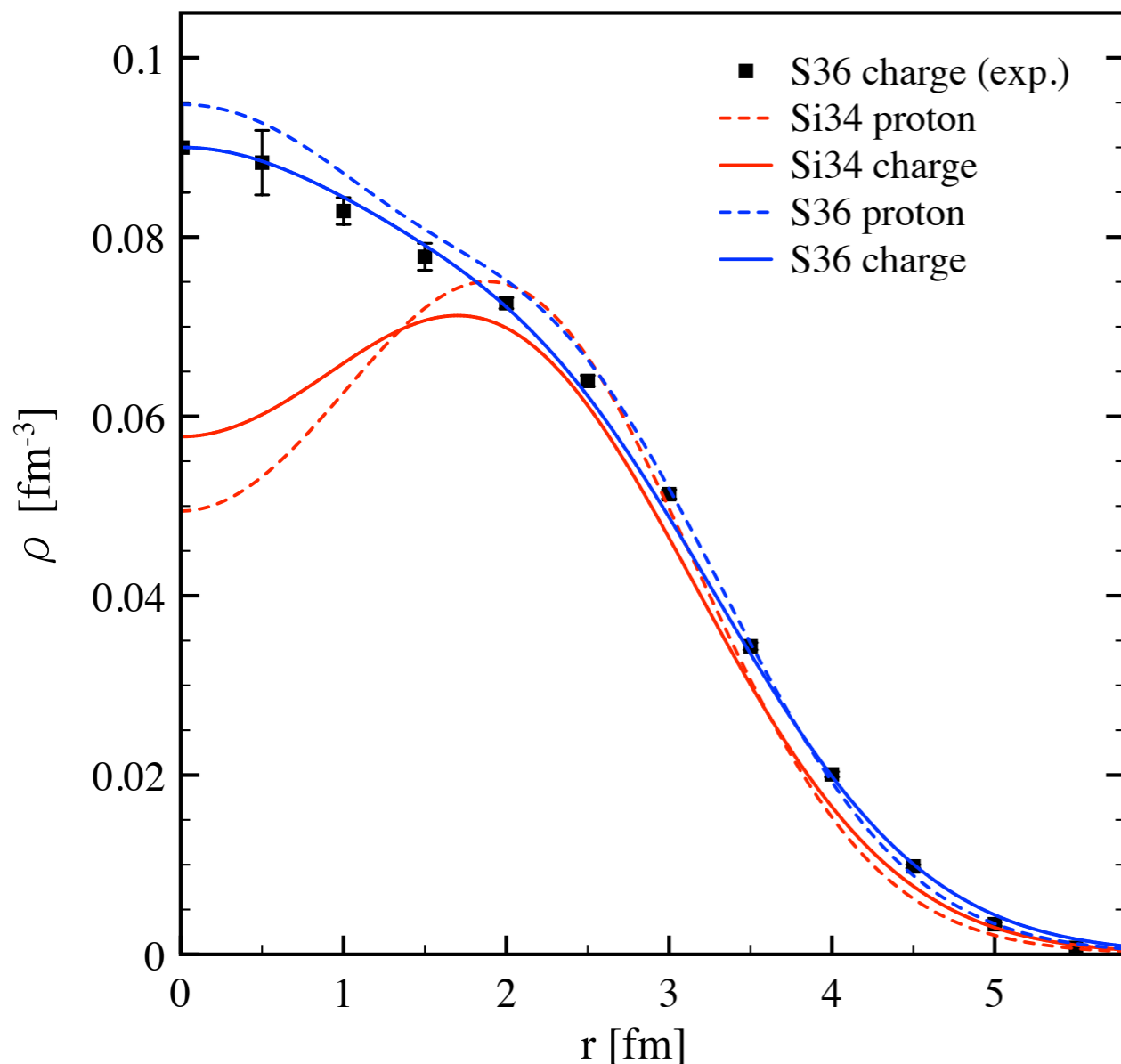


- Large sensitivity on the employed nuclear Hamiltonian
- Discrepancies between different Hamiltonians depend on the observable
- Good reproduction of nuclear radii with NNLOsat \leftrightarrow saturation properties

Charge density distribution of ^{34}Si

- ⊙ **Unconventional depletion** (“bubble”) in the centre of ρ_{ch} conjectured for ^{34}Si
 - Purely quantum mechanical effect (vacancy of $\ell = 0$ states embedded in larger- ℓ orbitals)
 - Conjectured associated effect on spin-orbit splitting (reduction for low- ℓ spin-orbit partners)

⊙ **Charge density** computed through folding with the finite charge of the proton [Duguet *et al.* 2017]



$$\rho_{\text{ch}}(r) = \sum_{i=1}^3 \frac{\theta_i}{r_i \sqrt{\pi}} \int_0^{+\infty} dr' \frac{r'}{r} \rho_{\text{p}}(r') \left[e^{-\left(\frac{r-r'}{r_i}\right)^2} - e^{-\left(\frac{r+r'}{r_i}\right)^2} \right]$$

- ⇒ Folding smears out central depletion
- ⇒ Good agreement with experiment for ^{36}S

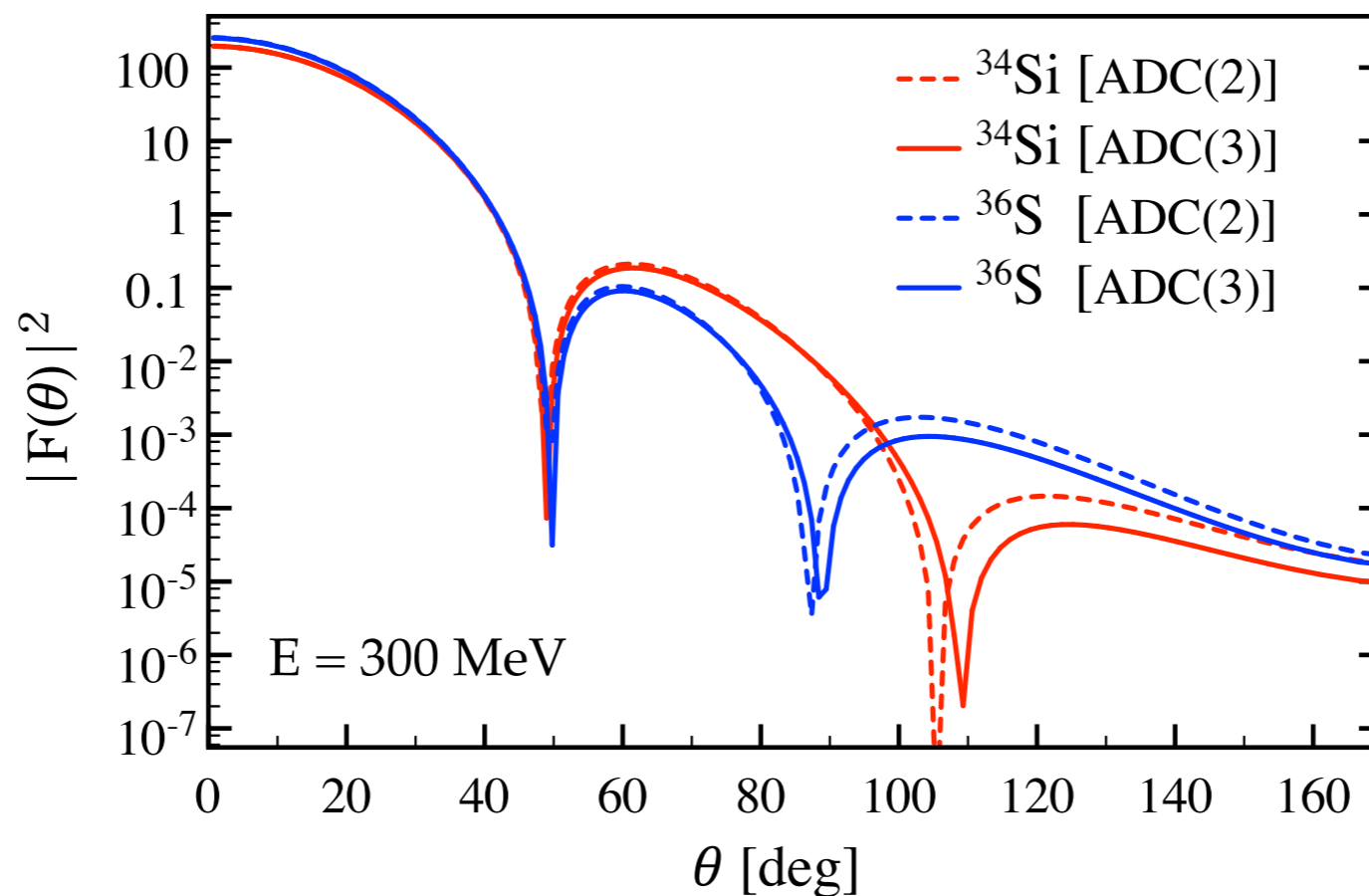
[Rychel *et al.* 1983]

- ⇒ Central depletion predicted for ^{34}Si

Charge form factor of ^{34}Si

- Charge form factor measured in (e,e) experiments sensitive to bubble structure?

$$F(q) = \int d\vec{r} \rho_{\text{ch}}(r) e^{-i\vec{q}\cdot\vec{r}} \quad \text{and} \quad q = 2p \sin \theta/2$$



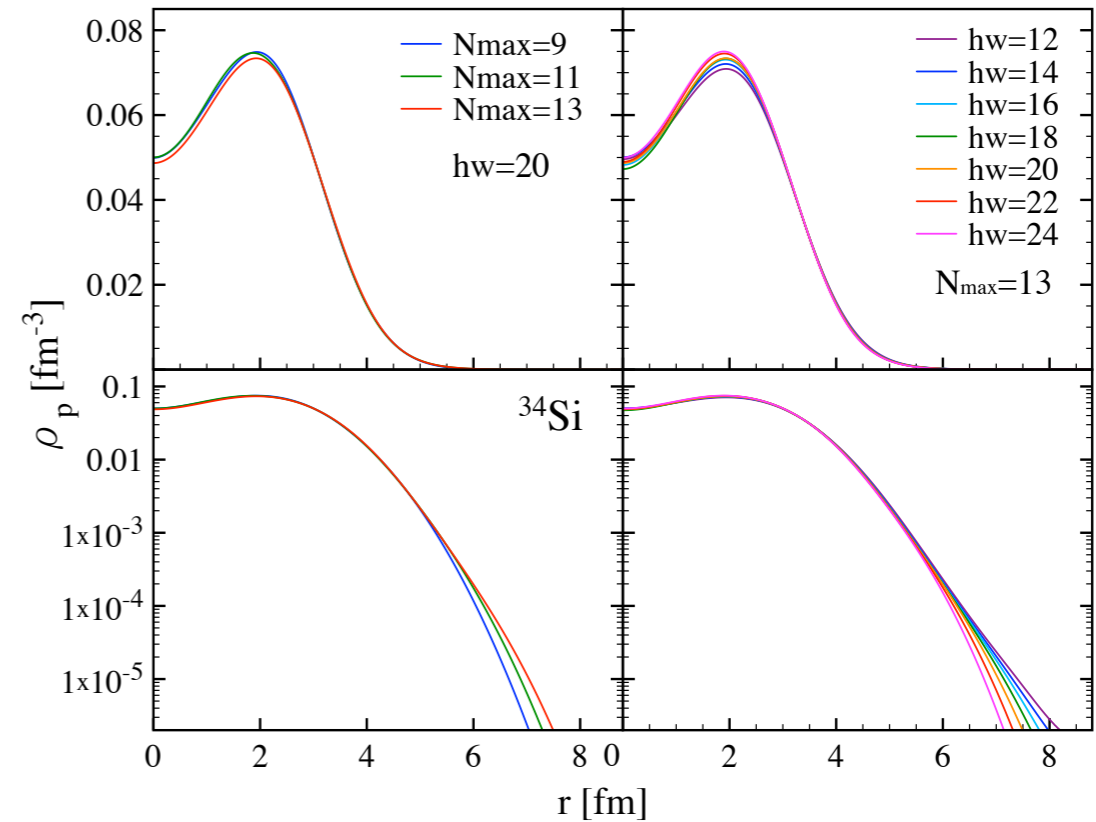
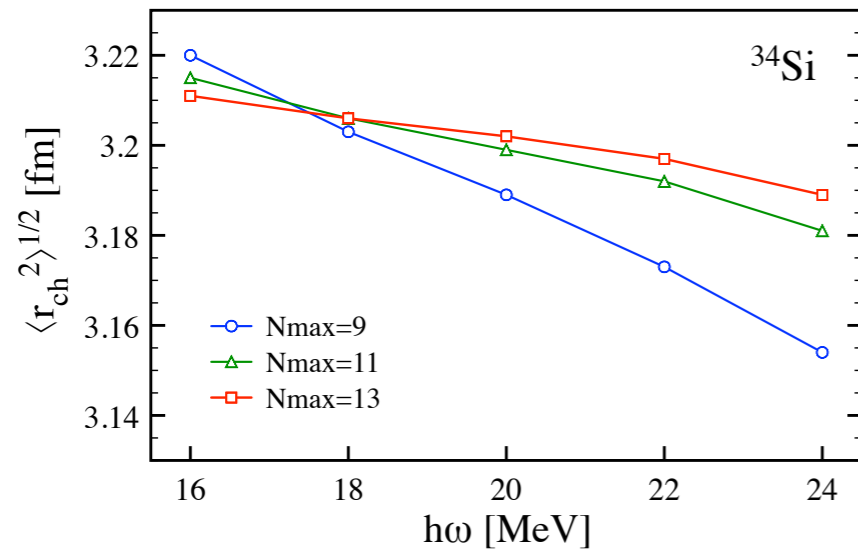
- Central depletion reflects in larger $F(\theta)$ for angles $\theta > 70^\circ$ and shifted 2nd minimum
- Future electron scattering experiments might be able to see its fingerprints**

On convergence

⊙ Calculations performed within different many-body truncations

○ ADC(1) = HF, ADC(2) & ADC(3)

⊙ Model space convergence



⊙ Many-body convergence

Binding energies

E [MeV]	ADC(1)	ADC(2)	ADC(3)	Experiment
^{34}Si	-84.481	-274.626	-282.938	-283.427
^{36}S	-90.007	-296.060	-305.767	-308.714



ADC(3) brings only $\sim 5\%$ additional binding

Charge radii

$\langle r_{\text{ch}}^2 \rangle^{1/2}$	ADC(1)	ADC(2)	ADC(3)	Experiment
^{34}Si	3.270	3.189	3.187	-
^{36}S	3.395	3.291	3.285	3.2985 ± 0.0024



Radii converged already at ADC(2) level

Spectral representation

Spectroscopic probabilities matrices

$$S_{\mu}^{+ab} \equiv \langle \Psi_0^A | a_a | \Psi_{\mu}^{A+1} \rangle \langle \Psi_{\mu}^{A+1} | a_b^{\dagger} | \Psi_0^A \rangle$$

$$S_{\nu}^{-ab} \equiv \langle \Psi_0^A | a_a^{\dagger} | \Psi_{\nu}^{A-1} \rangle \langle \Psi_{\nu}^{A-1} | a_b | \Psi_0^A \rangle$$

Spectroscopic factors

$$SF_{\mu}^{+} \equiv \text{Tr}_{\mathcal{H}_1} [\mathbf{S}_{\mu}^{+}] = \sum_{a \in \mathcal{H}_1} |U_{\mu}^a|^2$$

$$SF_{\nu}^{-} \equiv \text{Tr}_{\mathcal{H}_1} [\mathbf{S}_{\nu}^{-}] = \sum_{a \in \mathcal{H}_1} |V_{\nu}^a|^2$$

$$G_{ab}(z) = \sum_{\mu} \frac{\langle \Psi_0^A | a_a | \Psi_{\mu}^{A+1} \rangle \langle \Psi_{\mu}^{A+1} | a_b^{\dagger} | \Psi_0^A \rangle}{z - E_{\mu}^{+} + i\eta} + \sum_{\nu} \frac{\langle \Psi_0^A | a_b^{\dagger} | \Psi_{\nu}^{A-1} \rangle \langle \Psi_{\nu}^{A-1} | a_a | \Psi_0^A \rangle}{z - E_{\nu}^{-} - i\eta}$$

Eigenstates of A+1

One-nucleon addition
separation energies

$$E_{\mu}^{+} \equiv E_{\mu}^{A+1} - E_0^A$$

Eigenstates of A-1

One-nucleon removal
separation energies

$$E_{\nu}^{-} \equiv E_0^A - E_{\nu}^{A-1}$$

Combining numerator and denominator result in the spectral function

Spectral function

$$\mathbf{S}(z) \equiv \sum_{\mu \in \mathcal{H}_{A+1}} \mathbf{S}_{\mu}^{+} \delta(z - E_{\mu}^{+}) + \sum_{\nu \in \mathcal{H}_{A-1}} \mathbf{S}_{\nu}^{-} \delta(z - E_{\nu}^{-})$$

Spectral strength distribution

$$\begin{aligned} \mathcal{S}(z) &\equiv \text{Tr}_{\mathcal{H}_1} [\mathbf{S}(z)] \\ &= \sum_{\mu \in \mathcal{H}_{A+1}} SF_{\mu}^{+} \delta(z - E_{\mu}^{+}) + \sum_{\nu \in \mathcal{H}_{A-1}} SF_{\nu}^{-} \delta(z - E_{\nu}^{-}) \end{aligned}$$

Spectral representation

$$G_{ab}(z) = \sum_{\mu} \frac{U_a^{\mu} (U_b^{\mu})^*}{z - E_{\mu}^+ + i\eta} + \sum_{\nu} \frac{(V_a^{\nu})^* V_b^{\nu}}{z - E_{\nu}^- - i\eta}$$

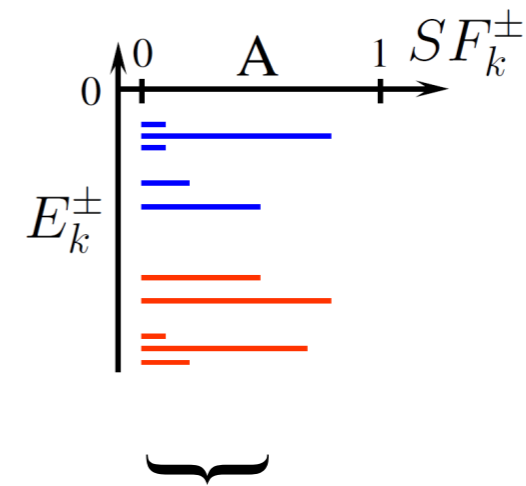
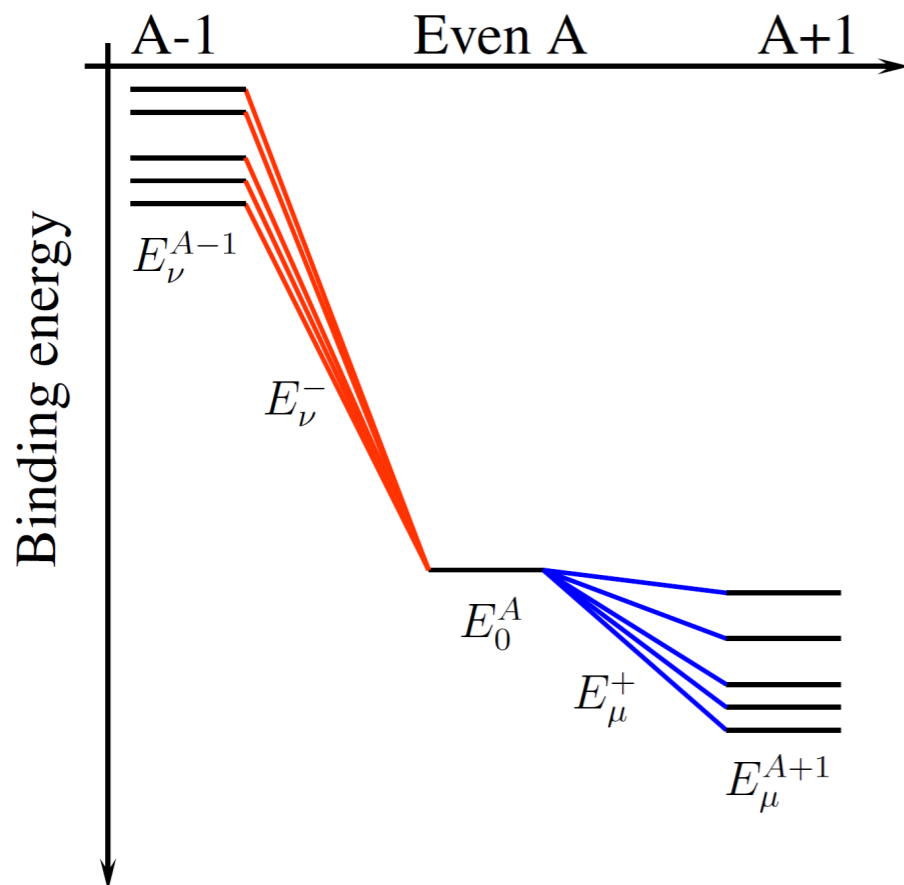
Separation energies

$$E_{\mu}^+ \equiv E_{\mu}^{A+1} - E_0^A$$

$$E_{\nu}^- \equiv E_0^A - E_{\nu}^{A-1}$$

Spectral strength distribution

$$\mathcal{S}(z) = \sum_{\mu \in \mathcal{H}_{A+1}} SF_{\mu}^+ \delta(z - E_{\mu}^+) + \sum_{\nu \in \mathcal{H}_{A-1}} SF_{\nu}^- \delta(z - E_{\nu}^-)$$



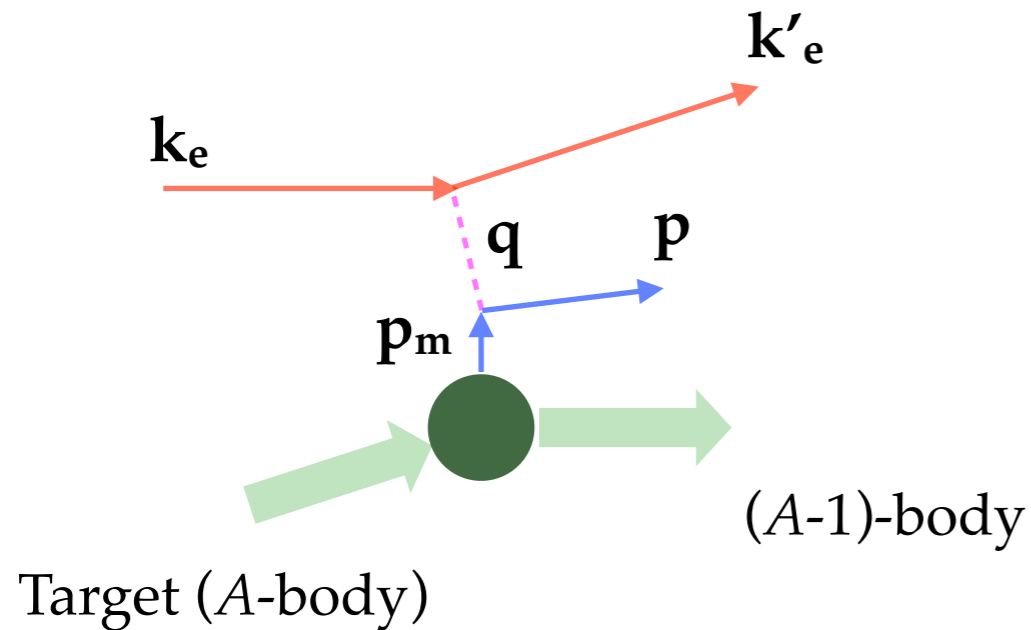
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Spectral strength in experiments

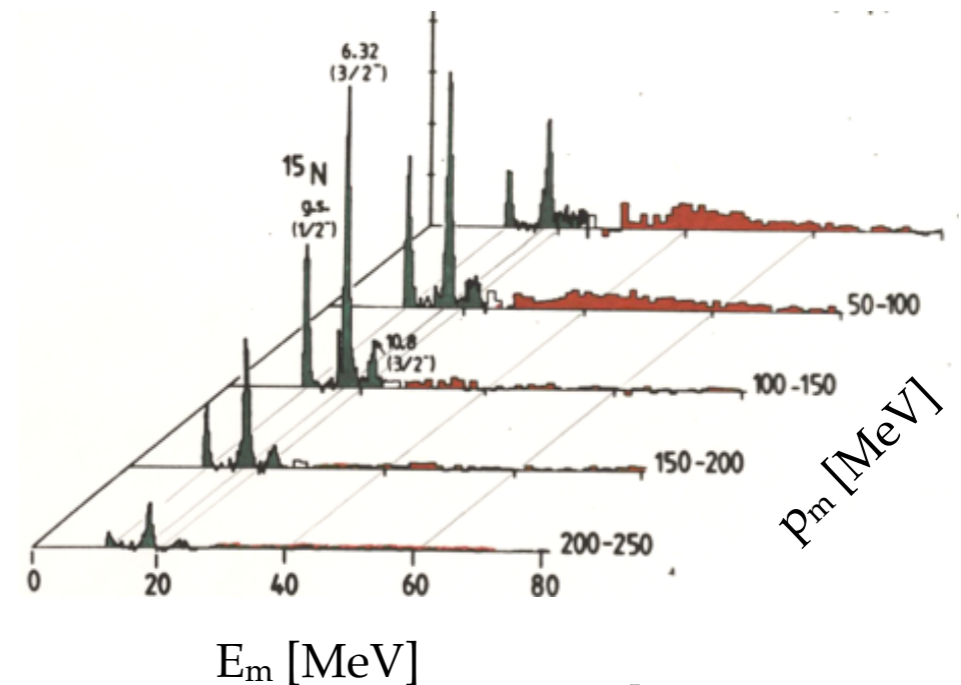
⊙ Clean connection to $(e,e'p)$ experiments



- Measuring \mathbf{q} and \mathbf{p} gives information on \mathbf{p}_m
- Similarly for missing energy E_m
- Spectral strength distribution $\leftrightarrow P(\mathbf{p}_m, E_m)$

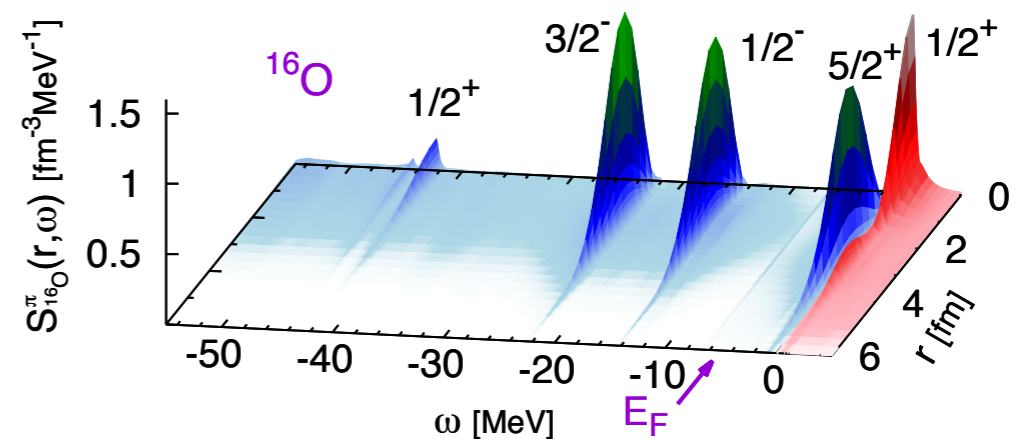
⊙ Spectroscopy via knockout/transfer exp.

Results from $(e,e'p)$ on ^{16}O (ALS in Saclay)



[Mougey *et al.* 1980]

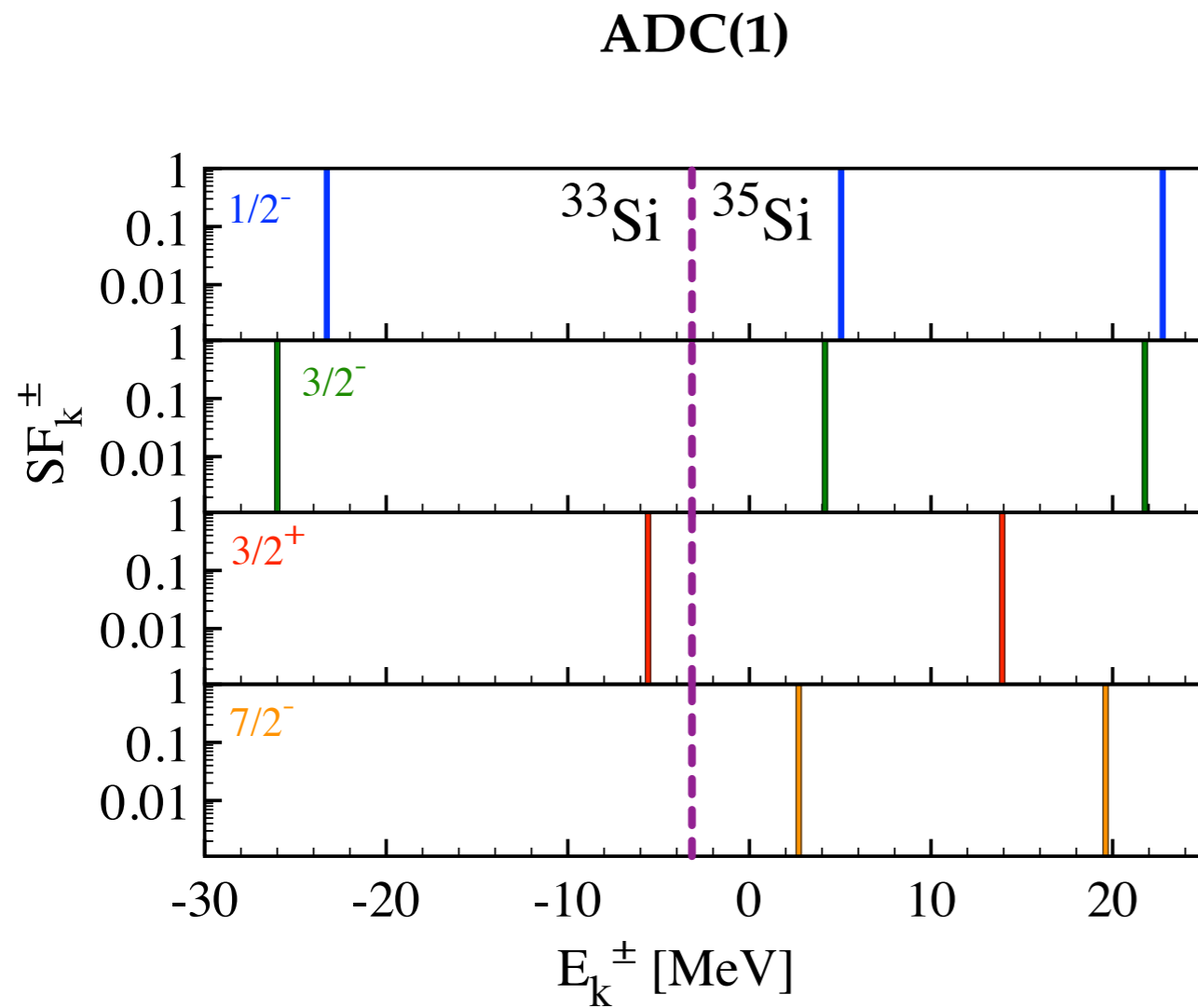
SCGF calculations



[Cipollone *et al.* 2015]

Spectral strength distribution

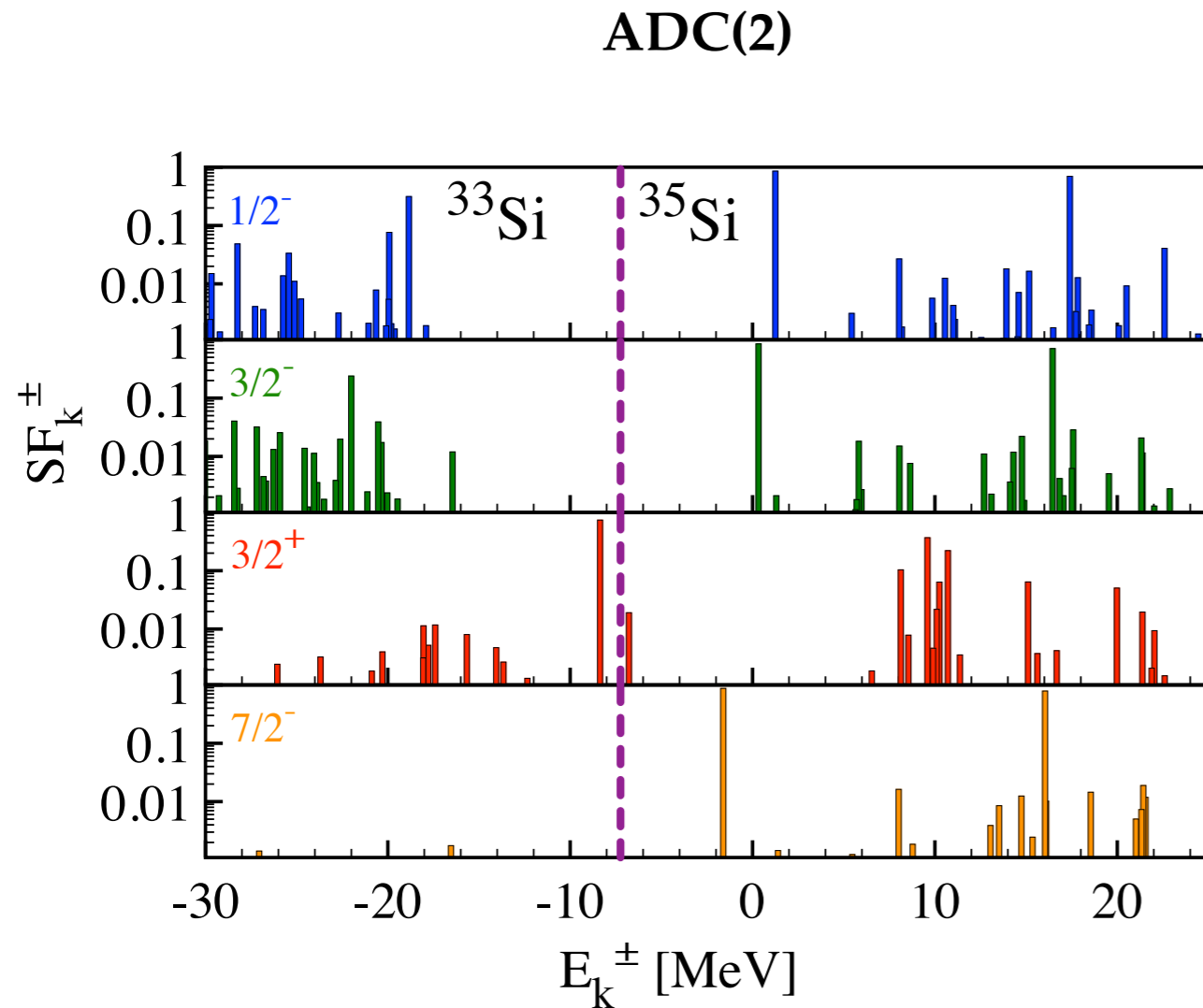
⊙ ^{34}Si neutron addition & removal strength



○ Independent-particle picture

Spectral strength distribution

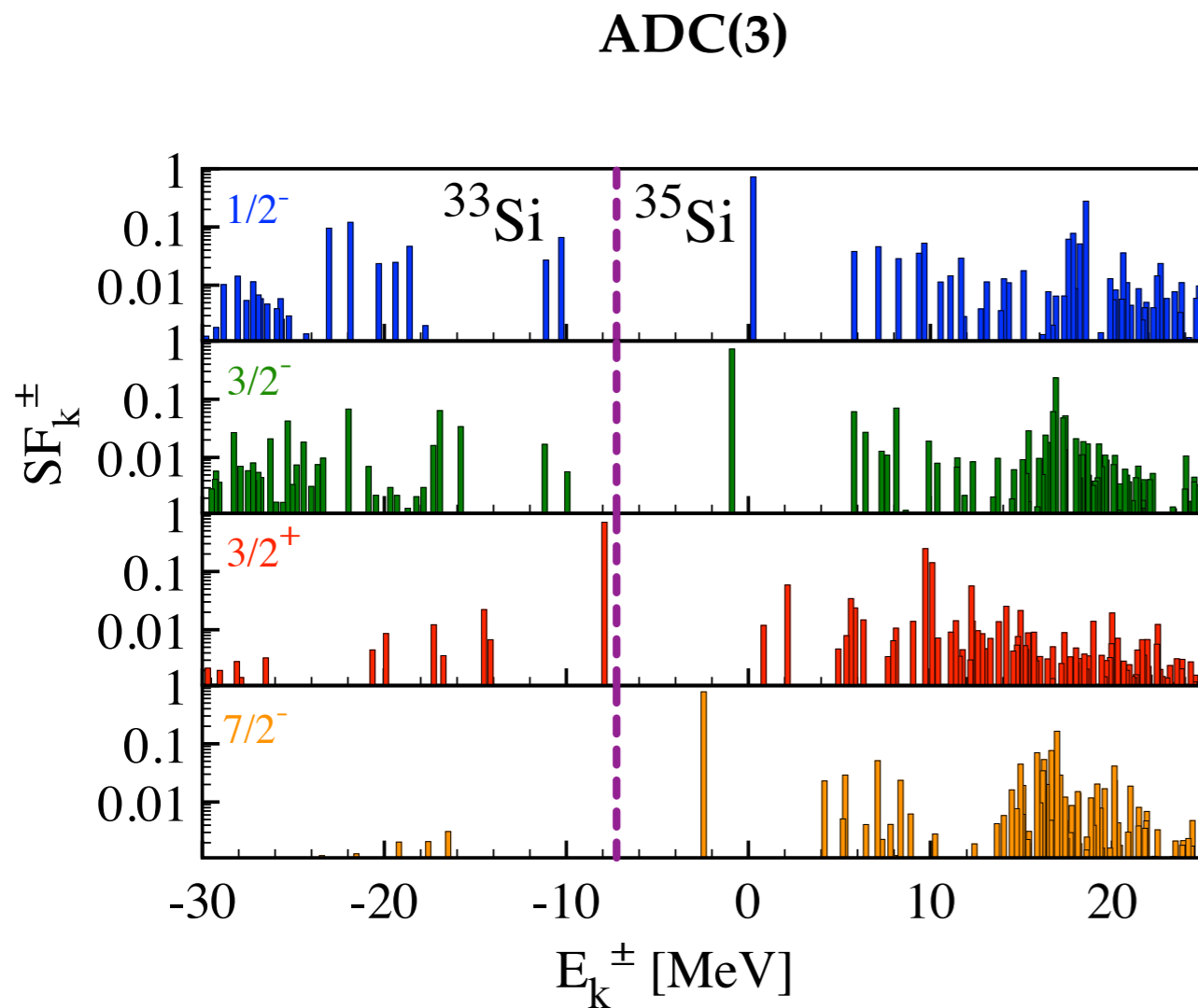
⊙ ^{34}Si neutron addition & removal strength



○ Second-order dynamical correlations fragment IP peaks

Spectral strength distribution

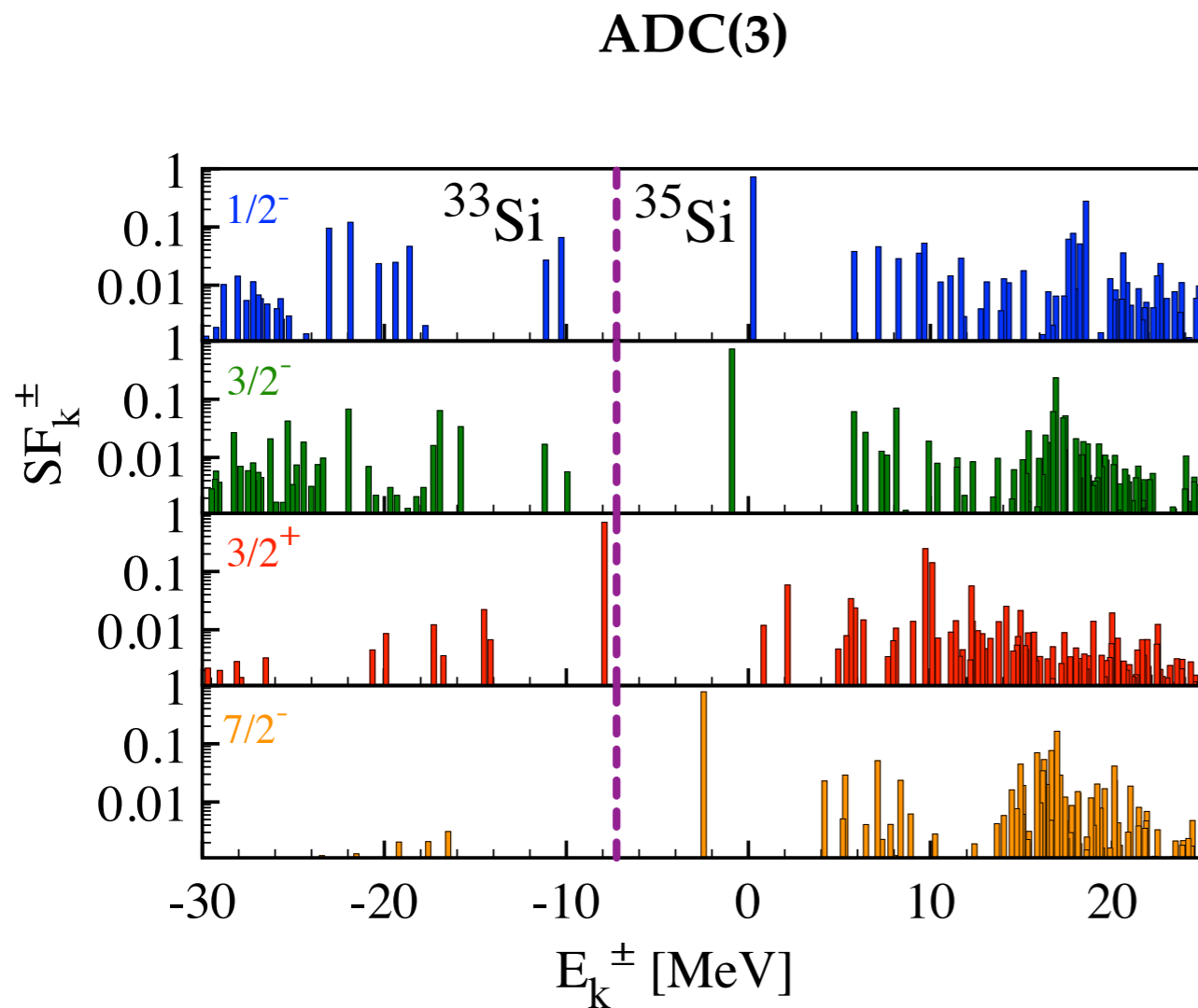
○ ^{34}Si neutron addition & removal strength



- Third-order compresses the spectrum (main peaks)
- Further fragmentation is generated

Spectral strength distribution

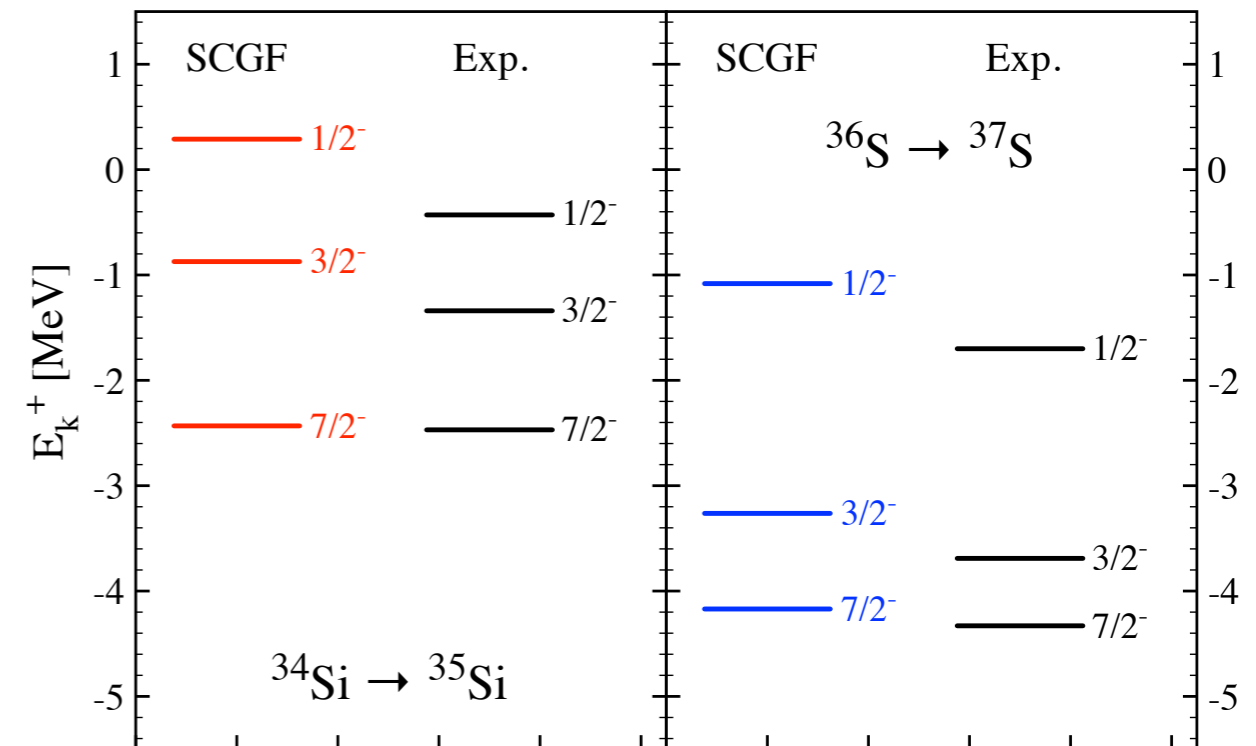
○ ^{34}Si neutron addition & removal strength



- Third-order correlations compress the spectrum
- Further fragmentation is generated

One-neutron addition

Exp. data: [Thorn *et al.* 1984]
 [Eckle *et al.* 1989]
 [Burgunder *et al.* 2014]



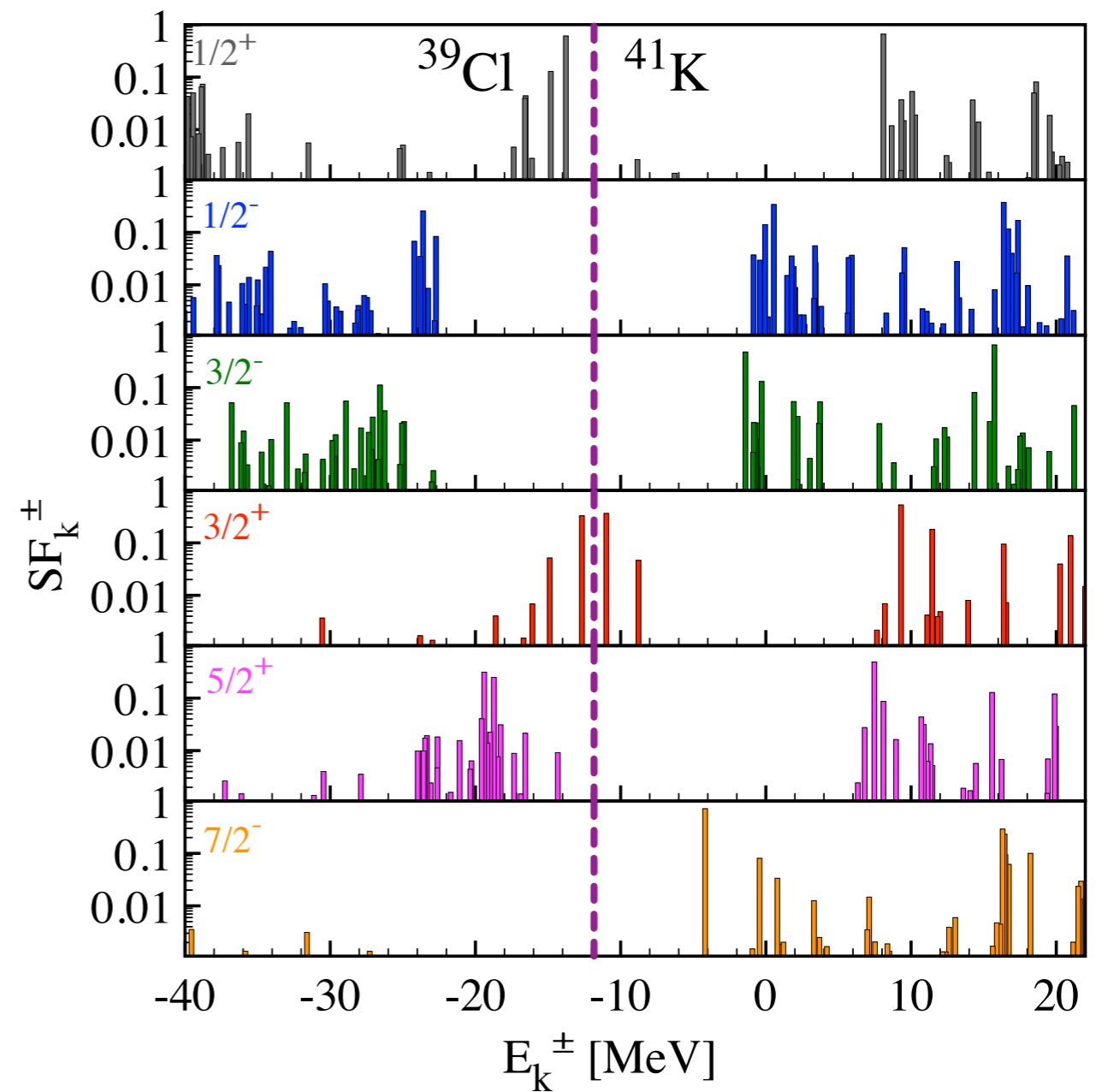
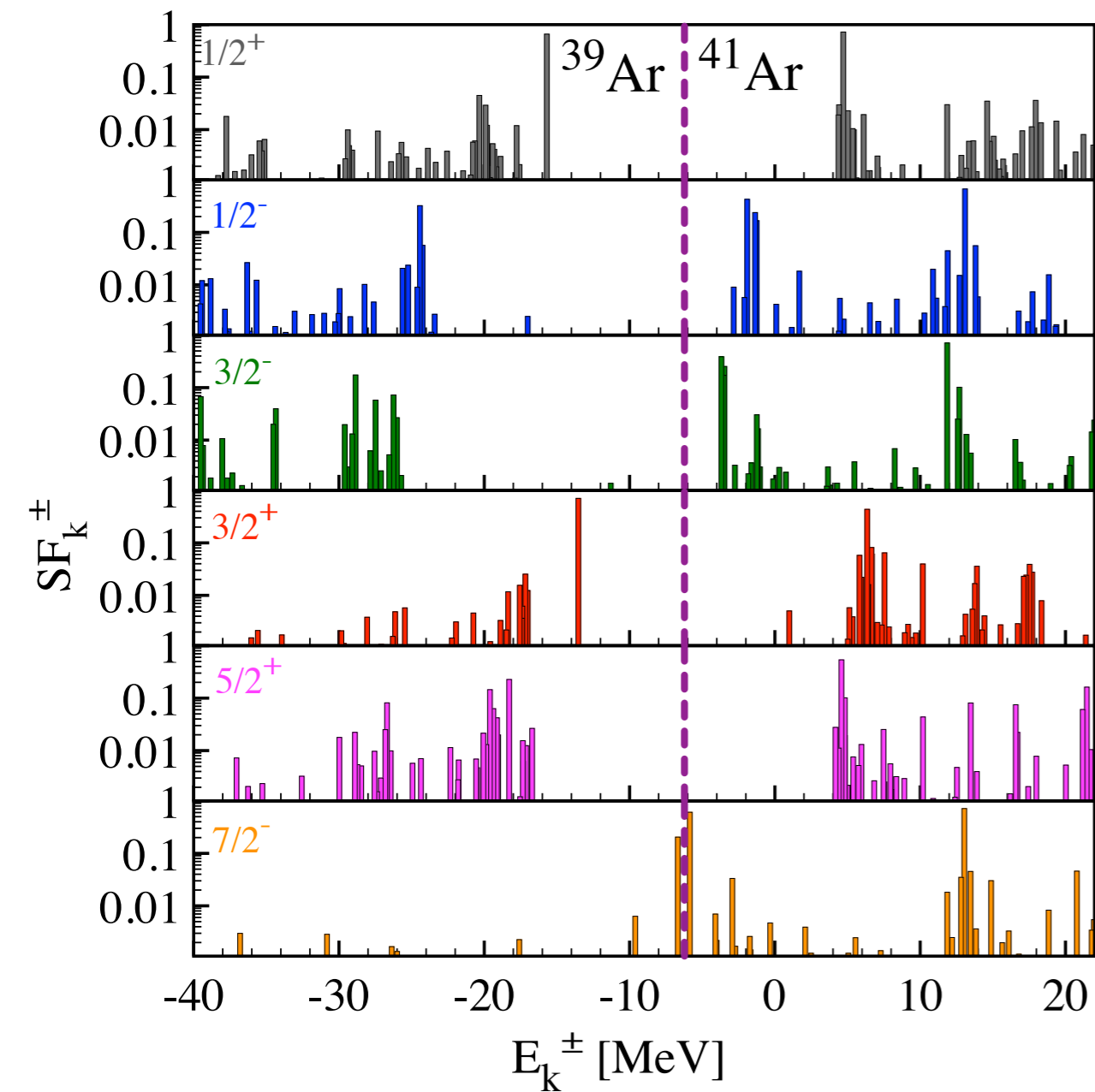
Reduction of $E_{1/2^-} - E_{3/2^-}$ spin-orbit splitting
 (unique in the nuclear chart) well reproduced

Spectral function of ^{40}Ar

© ADC(2) truncation, N3LO local-nonlocal interaction

Neutrons

Protons

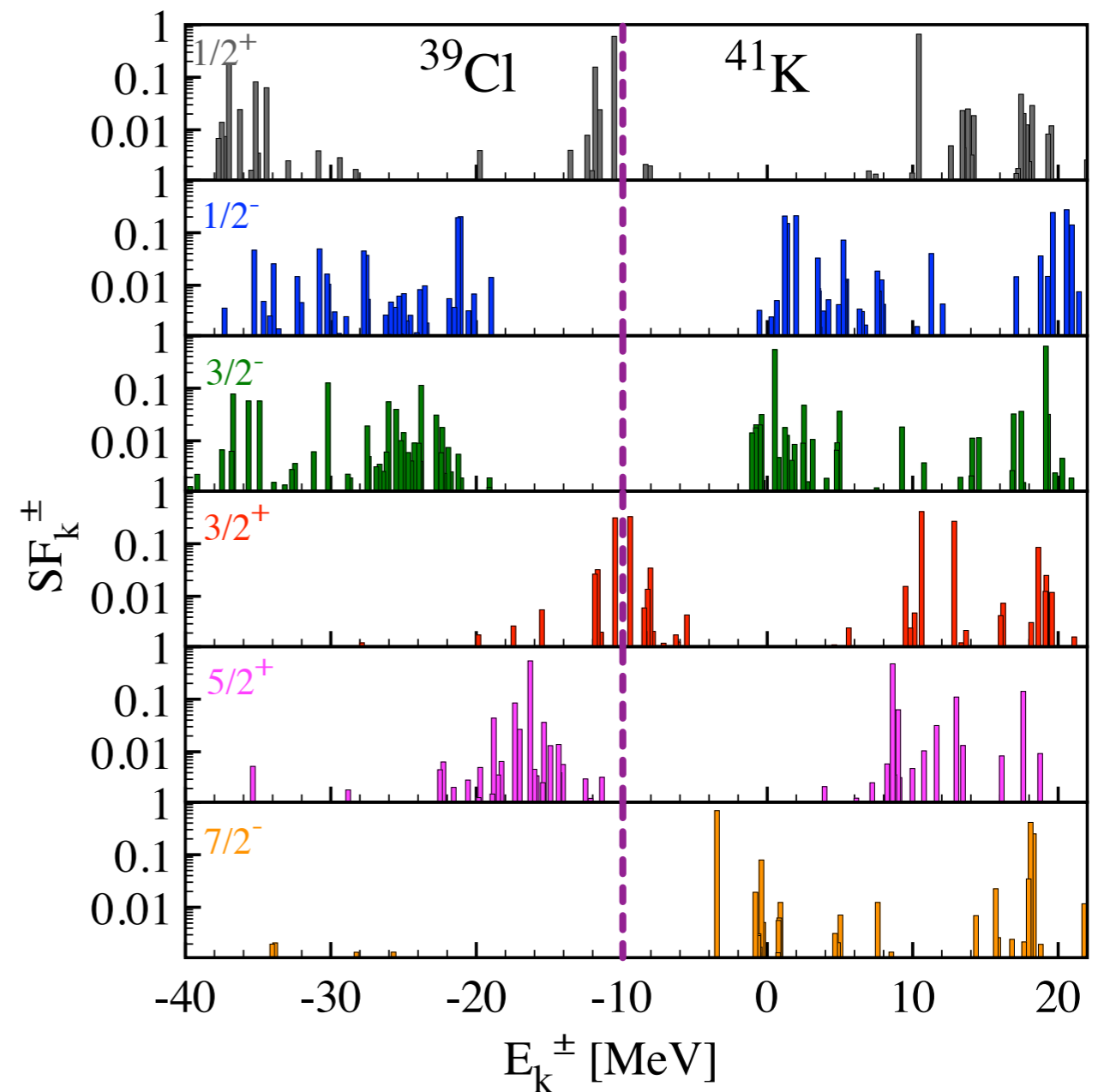
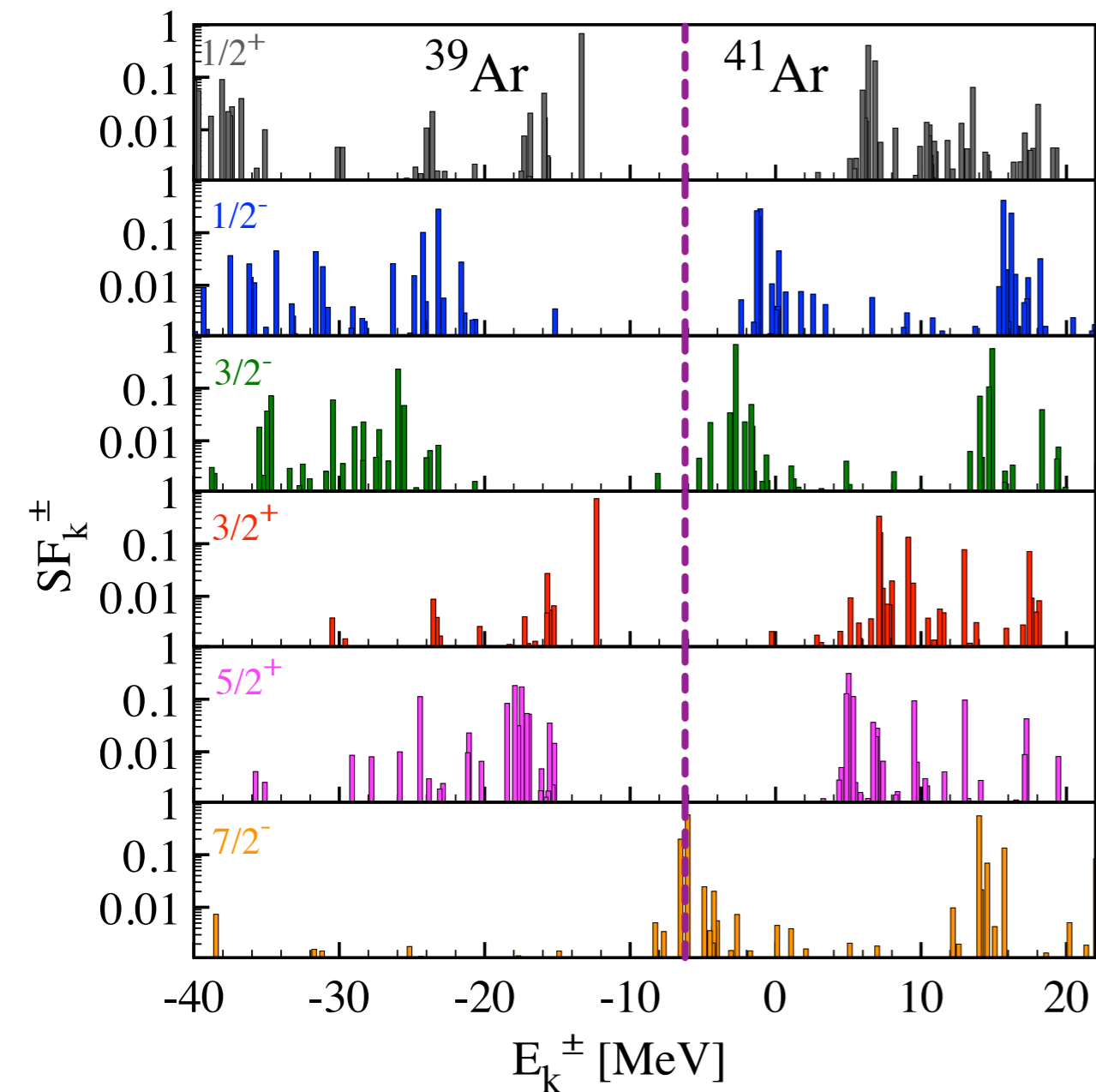


Spectral function of ^{40}Ar

© ADC(2) truncation, NNLO_{sat} interaction

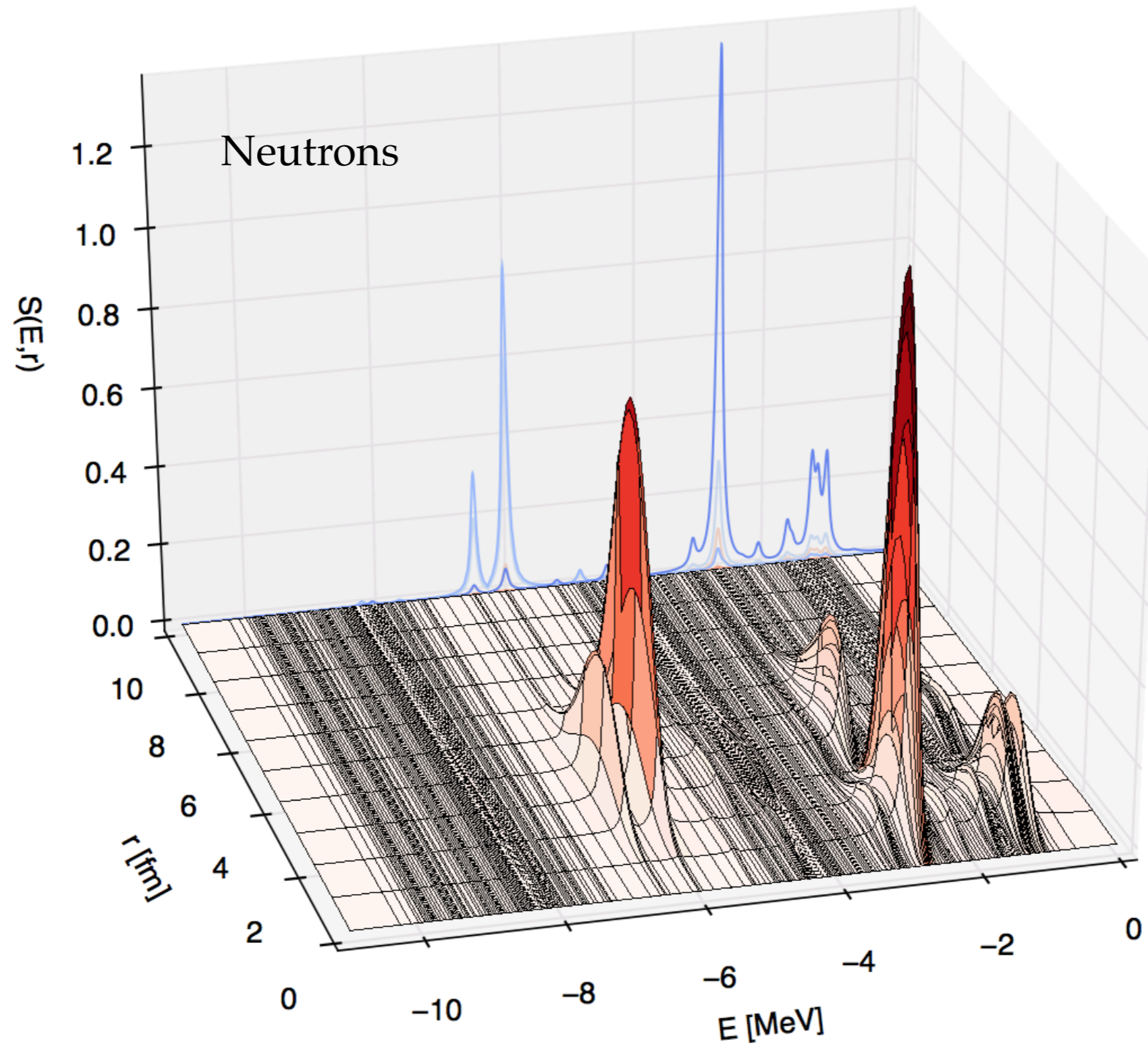
Neutrons

Protons



Spectral function of ^{40}Ar

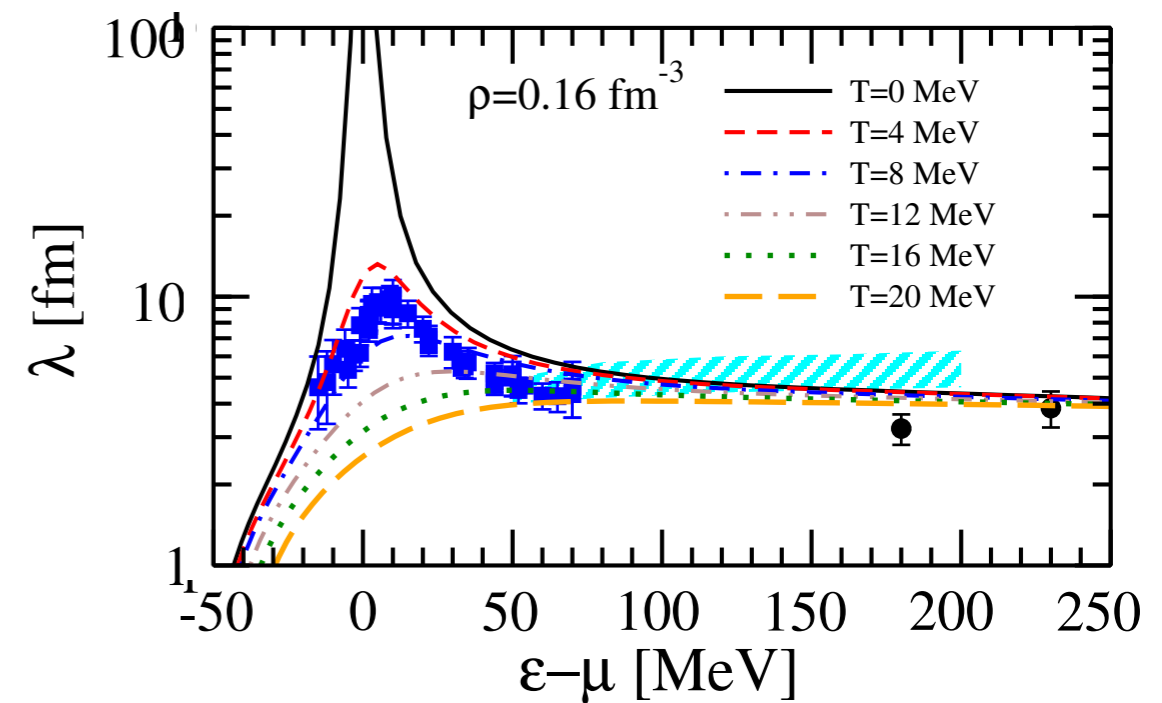
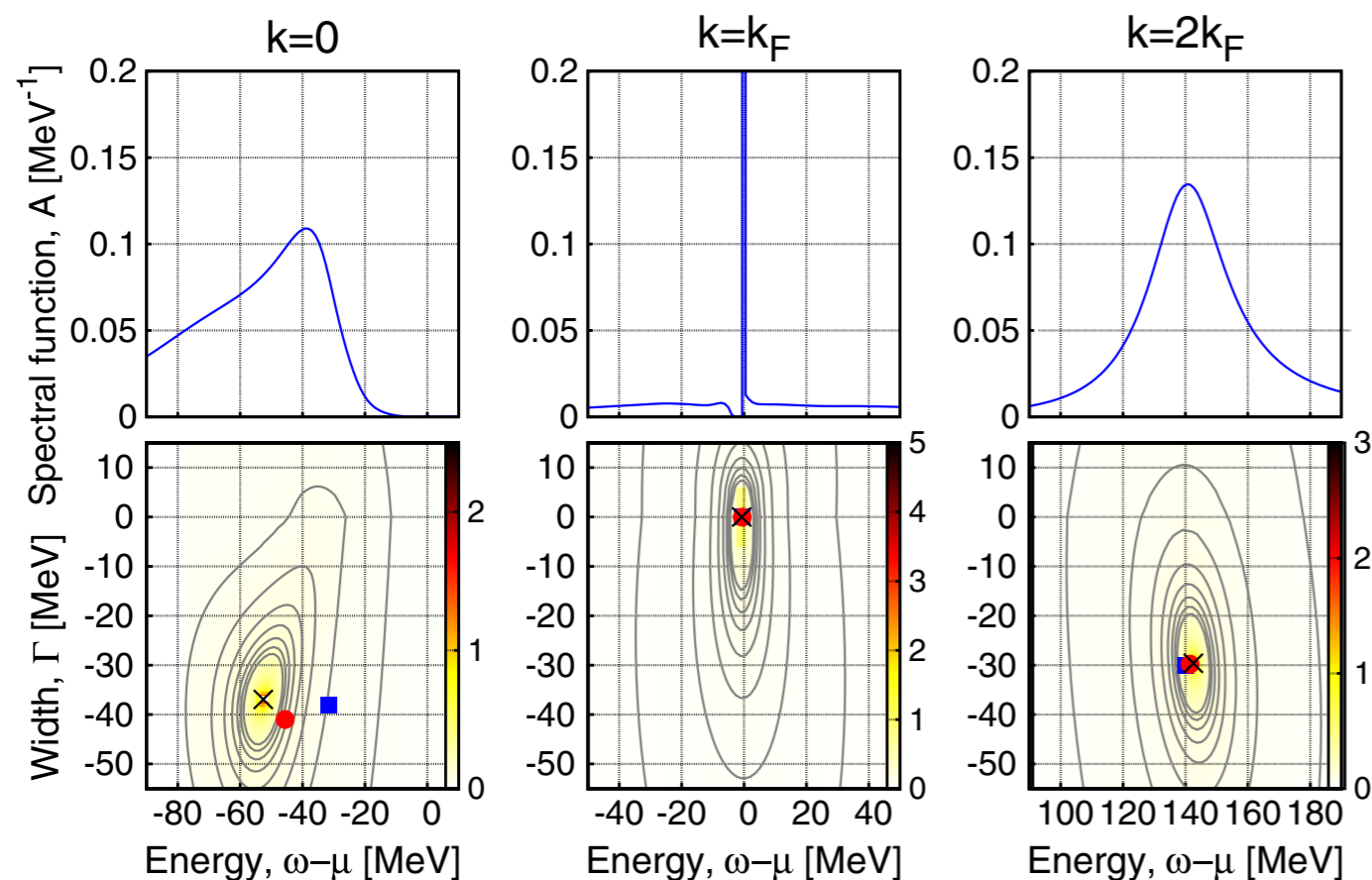
© ADC(2) truncation, NNLO_{sat} interaction



Fragmentation of single-particle strength in infinite matter

◉ Spectral function depicts correlations

- ◉ Broad peak signals depart from mean-field / independent particle picture
- ◉ Well-defined (long-lived) quasiparticles at the Fermi surface
- ◉ Long mean free path for $E < E_F$



Electromagnetic response

- ⊙ Nuclear response produced by an isovector dipole operator

$$R(E) = \sum_{\nu} |\langle \Psi_{\nu}^A | \hat{D} | \Psi_0^A \rangle|^2 \delta(E_{\nu} - E)$$



$$\sum_{ab} \langle a | \hat{D} | b \rangle \langle \Psi_{\nu}^A | c_a^{\dagger} c_b | \Psi_0^A \rangle$$

Electric dipole excitation operator

Polarisation propagator computed in RPA

- ⊙ Total photoabsorption cross section

$$\sigma_{\gamma}(E) = 4\pi^2 \alpha E R(E)$$

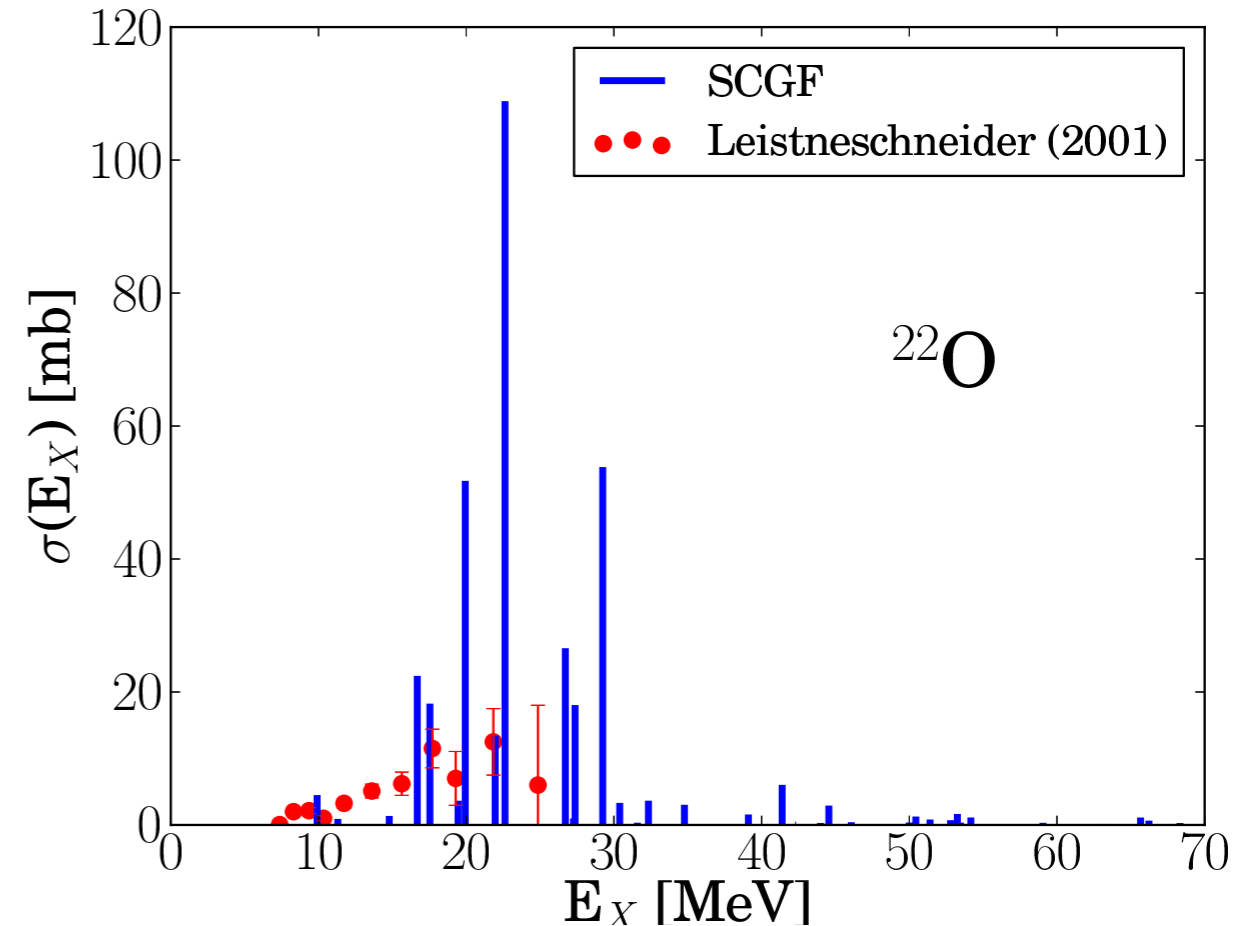
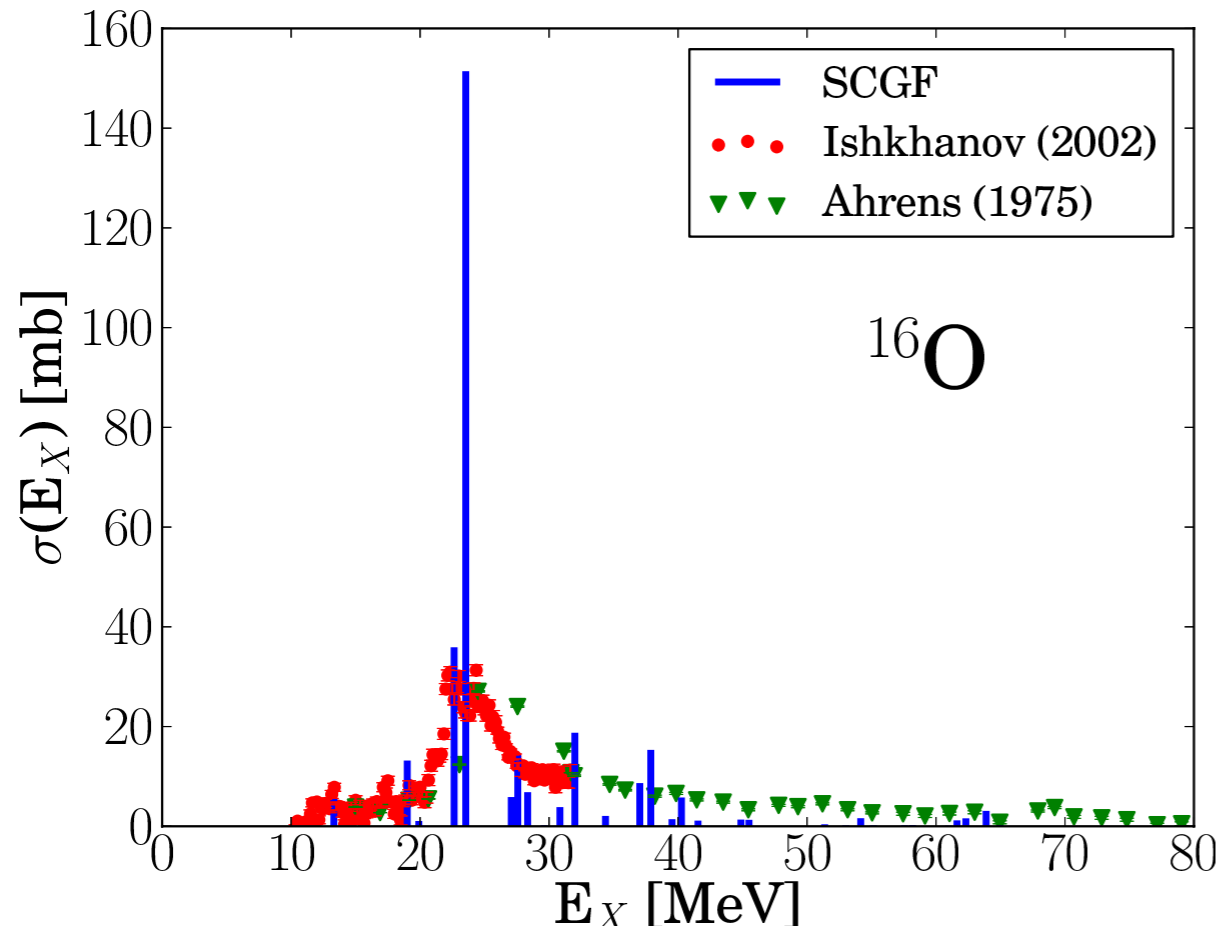
- ⊙ Dipole polarisability

$$\alpha_D = 2\alpha \int dE \frac{R(E)}{E}$$

Electromagnetic response

⊙ Computed σ from RPA response vs. σ from photoabsorption and Coulomb excitation

[Raimondi *et al.* in preparation]



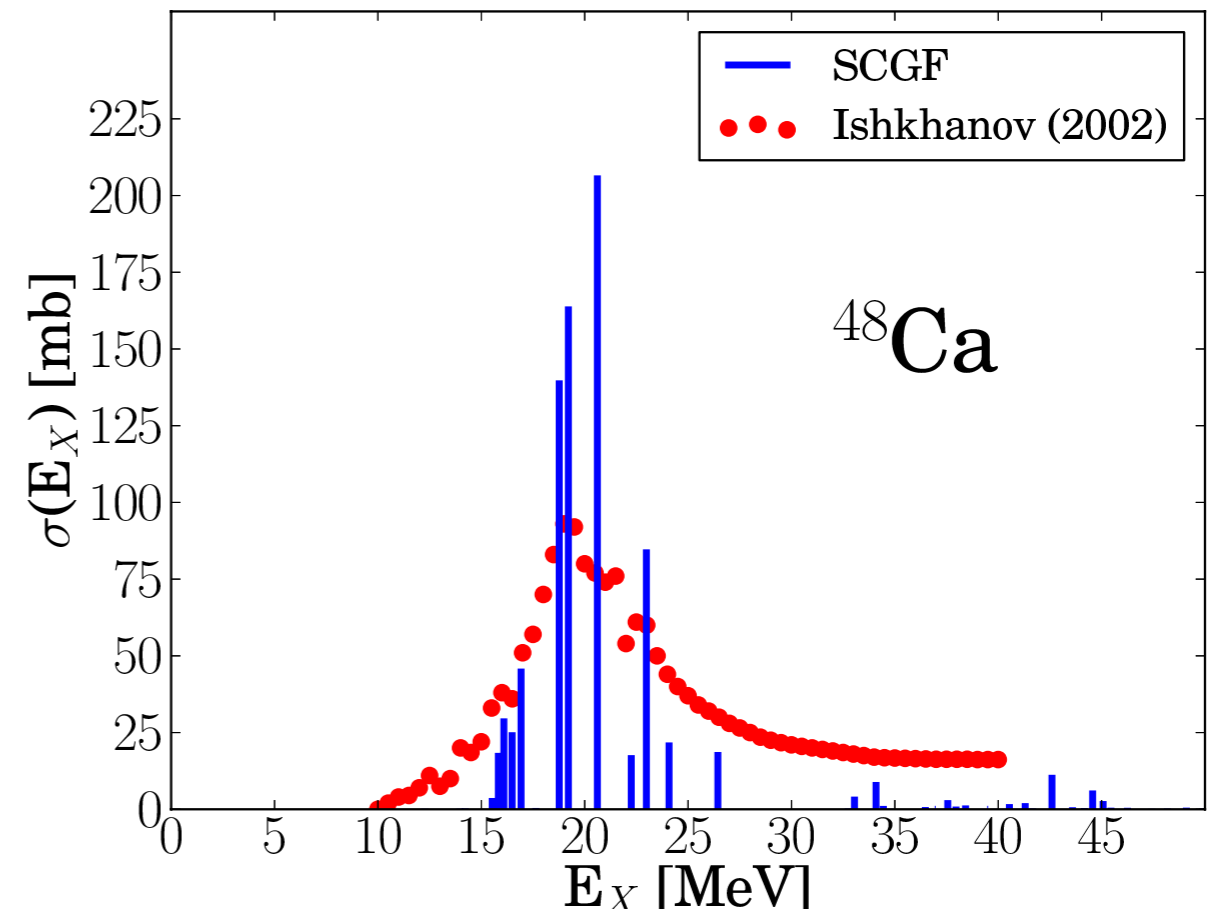
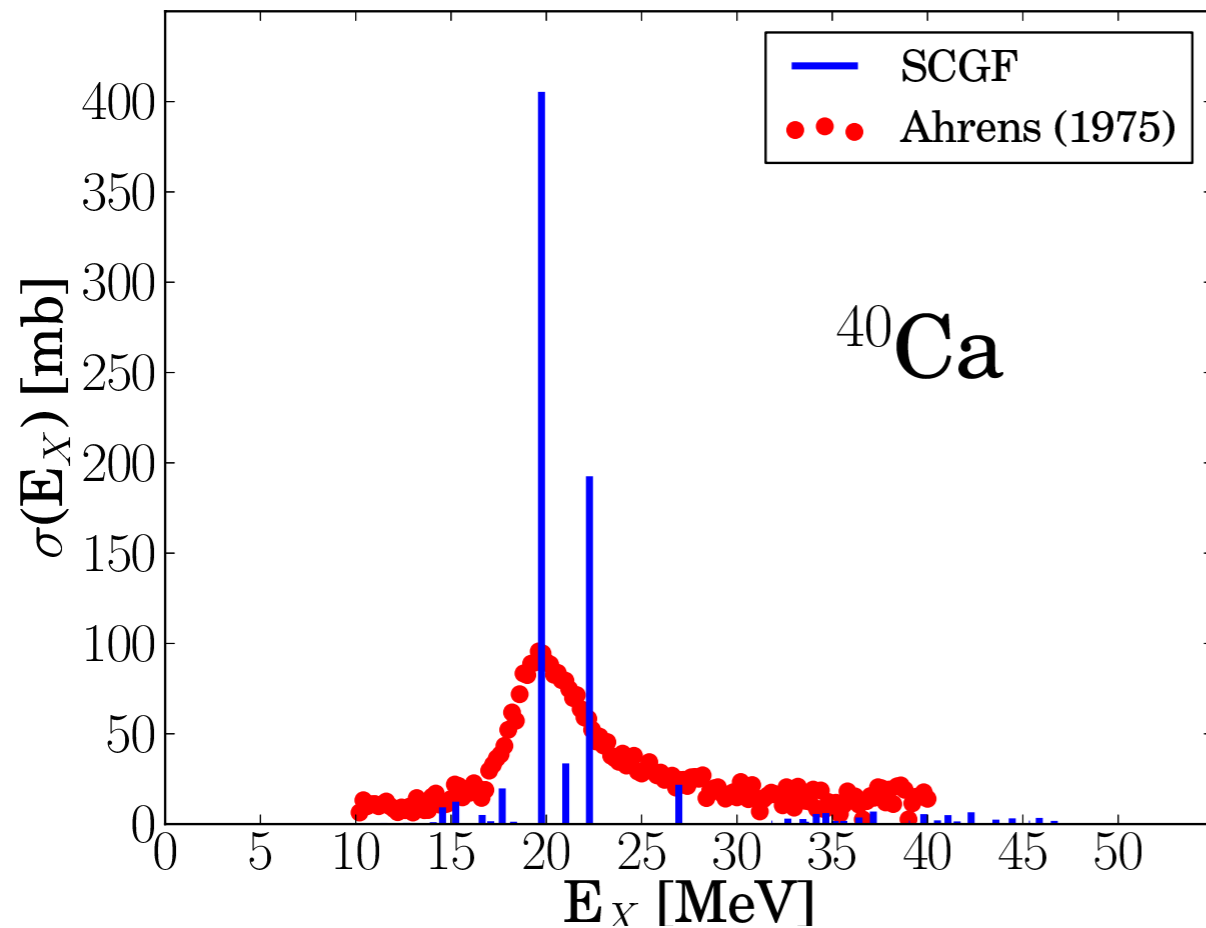
- GDR position of ^{16}O well reproduced
- Hint of a soft dipole mode in ^{22}O
- Comparison with CC LIT results for α_D

Nucleus	Dipole polarizability α_D (fm^3)		
	SCGF	CC/LIT	Exp
^{16}O	0.50	0.57(1)	0.585(9)
^{22}O	0.72	0.86(4)	0.43(4)

Electromagnetic response

⊙ Computed σ from RPA response vs. σ from photoabsorption and Coulomb excitation

[Raimondi *et al.* in preparation]



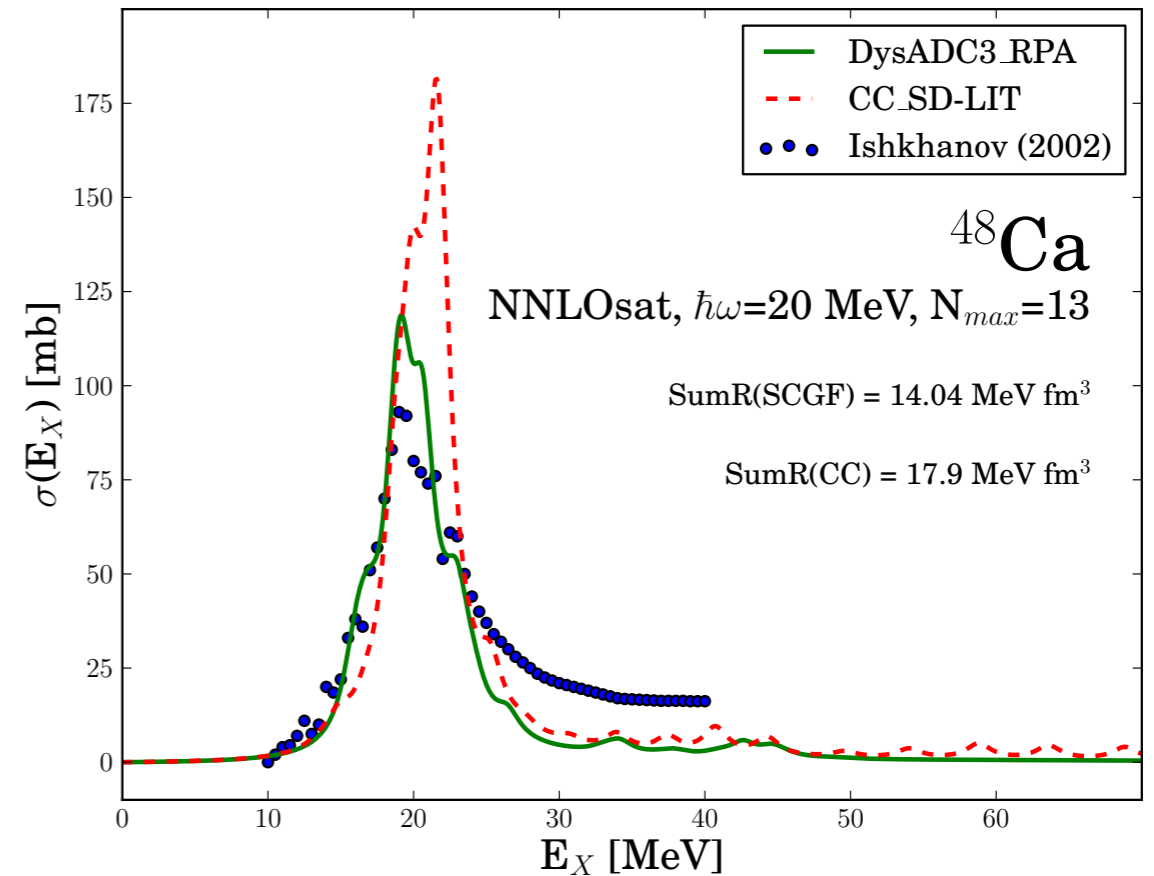
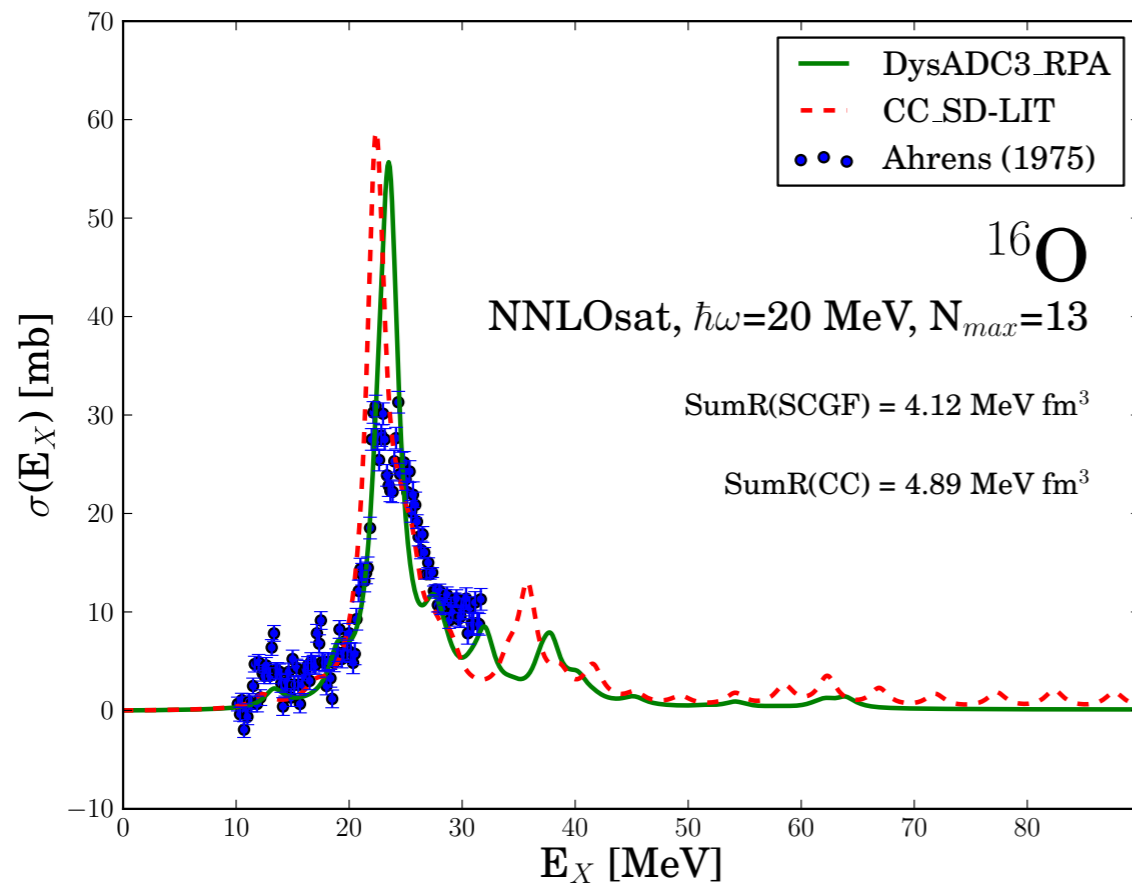
- GDR positions reproduced
- Total sum rule OK but poor strength distribution
- Comparison with CC LIT results for α_D

Nucleus	Dipole polarizability α_D (fm^3)		Exp
	SCGF	CC/LIT	
^{40}Ca	1.79	1.47 (1.87) _{thresh}	1.87(3)
^{48}Ca	2.08	2.45	2.07(22)

Electromagnetic response

- Comparison with coupled-cluster Lorentz integral transform (CC-LIT)

[Raimondi *et al.* in preparation]



- Different ways of including correlations

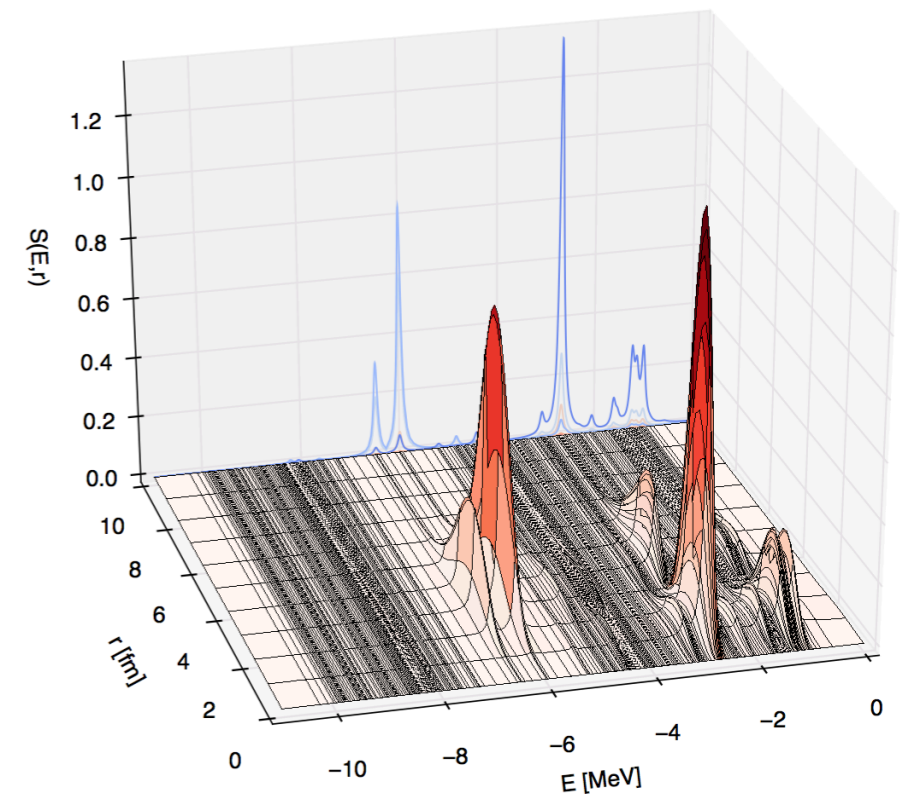
GF \rightarrow RPA (first-order 2-body correlator) on top of fully correlated reference state

CC \rightarrow SD (analogous to second RPA) on top of HF reference state

Conclusions

- ◎ **Ab initio calculations routinely access mid-mass nuclei**

- Many-body formalism well grounded
- Interactions currently represent largest source of uncertainty
- Pragmatic choices lead to successful applications
- Thorough assessment of theoretical error?



- ◎ **Self-consistent Green's function approach gives access to a variety of observables**

- Ground-state properties along isotopic chains, spectral functions, response functions
- Case of ^{40}Ar to be investigated in ADC(3) scheme (work in progress...)