

Experimentally determined nuclear-structure properties of 0ν2β decay candidates

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Nuclear ab initio theories and neutrino physics INT-18-1a Week 5, 29 March 2018

Overview

- o *Basic premises (ground-state nucleon occupancies, pairing)*
- o *Experiments (now a 10-year project, 4 candidates 'done')*
- o *Analysis techniques*
- o *Normalizations*
- o *Quenching*
- o *An overview of results, compared with theory*
- o *Comments on pairing*

Connecting half-life and mass

Engel and Menéndez, Rep. Prog. Phys. 80, 046301 (2017)

Nuclear matrix elements

$$
[T_{1/2}^{0\nu}]^{-1} = (\text{Phase Space Factor}) \times |\text{Nuclear Matrix Element}|^2 \times |\langle m_{\beta\beta} \rangle|^2
$$

$$
T_{1/2}^{0\nu} \propto a \cdot \epsilon \cdot \sqrt{\frac{M \cdot t}{B \cdot \Delta E}} \propto 1/|\text{NME}|^2
$$

Experimental searches are often discussed in terms of their sensitivity to a given half life, accounting for enrichment, efficiency, backgrounds, resolution, and mass, though NP has a significant role …

Figure taken from one of Jason Detwiler's talks found online

Mechanism, rationale

Ground states

- Single- and two-particle properties should be important:
	- § How do the *protons and neutrons rearrange themselves going from the initial to final state*? (we can probe that)
	- Are the ground states 'simple' BCS like states? (we can probe that too) *Bernadette Rebeiro talked about this earlier in the week*
- Can knowledge of the above inform or constrain theoretical calculations?
- How well are the uncertainties (in the analysis of the experimental data) understood?
- *(Are all these things not already known (after all, these are [essentially] stable isotopes?)*

Series of experiments

Single-nucleon and two-nucleon transfer on nuclei involved in the $76\text{Ge} \rightarrow 76\text{Se}$, 100 Mo→ 100 Ru, 130 Te→ 130 Xe, and 136 Xe→ 136 Ba decays

Original works, including cross sections and analyzed data:

S. J. Freeman et al., Phys. Rev. C 75, 051301(R) (2007): A = 76 neutron pairing J. P. Schiffer et al., Phys. Rev. Lett. 100, 112501 (2008): A = 76neutron occupancies B. P. Kay et al., Phys. Rev. C 79, 021301(R) (2009): A = 76 proton occupancies T. Bloxham et al., Phys. Rev. C 82, 027308 (2010): A = 130 neutron (and proton) pairing J. S. Thomas et al., Phys. Rev. C 86, 047304 (2012): A = 100 neutron pairing B. P. Kay et al., Phys. Rev. C 87, 011302(R) (2013): A = 130 neutron occupancies A. Roberts et al., Phys. Rev. C 87, 051305(R) (2013): A = 76 proton pairing J. P. Entwisle et al., Phys. Rev. C 93, 064312 (2016): A = 130 and A = 136 proton occupancies S. V. Szwec et al., Phys. Rev. C 94, 054314 (2016): A = 136 neutron occupancies S. J. Freeman et al., Phys. Rev. C 96, 054325 (2017): A = 100 proton and neutron occupancies

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D. K. Sharp et al., upcoming works on A = 116, 124, and 150 neutron occupancies

Collaborators

Initiative led by Argonne and Manchester groups *Experiments at MLL (Munich), WNSL (Yale), RCNP (Osaka), and IPN (Orsay) [Tandem facilities (except of RCNP), with beams around 10 MeV/u, high-resolution magnetic spectrographs. They are among the few facilities left in the world where these measurements are possible.]*

J. P. Schiffer, S. J. Freeman, J. A. Clark, C. Deibel, C. R. Fitzpatrick, S. Gros, A. Heinz, D. Hirata, C. L. Jiang, B. P. Kay, A. Parikh, P. D. Parker, K. E. Rehm, A. C. C. Villari, V. Werner, C. Wrede, T. Adachi, H. Fujita, Y. Fujita, P. Grabmayr, K. Hatanaka, D. Ishikawa, H. Matsubara, Y. Meada, H. Okamura, Y. Sakemi, Y. Shimizu, H. Shimoda, K. Suda, Y. Tameshige, A. Tamii, T. Bloxham, S. A. McAllister, S. J. Freedman, K. Han, A. M. Howard, A. J. Mitchell, D. K. Sharp, J. S. Thomas, J. P. Entwisle, A. Tamii, S. Adachi, N. Aoi, T. Furuno, T. Hashimoto, C. R. Hoffman, E. Ideguchi, T. Ito, C. Iwamoto, T. Kawabata, B. Liu, M. Miura, H. J. Ong, G. Süsoy, T. Suzuki, S. V. Szwec, M. Takaki, M. Tsumura, T. Yamamoto T. E. Cocolios, L. P. Gaffney, V. Guimarães, F. Hammache, P. P. McKee, E. Parr, C. Portail, N. de Séréville, J. F. Smith, I. Stefan, ++

Now more 70 collaborators as participants in the various experiments

Focus of this talk

$$
\pi = \nu \qquad \qquad \pi = \nu
$$

$$
\pi = v
$$

 $136Ba$ 137Ba 138Ba

 $135C_S$ 136 C_S 137 C_S

134Xe 135Xe 136Xe

 $\pi \neq v$

Transfer reactions (what is measured) can be described by the stages, i.e., i.e

•• Around 10 MeV/u (direct reactions) • The Variety of reactions (momentum matching) Around 10 MeV/u (direct reactions)

dependence on the matrix element of the matrix element of the interaction. The interaction processes are action

neutron.

Spectra from BPK et al., Phys. Rev. C 87, 011302(R) (2013)

Transfer reactions (what is inferred)

Measure at several angles, shapes characteristic of ℓ

tions are tabulated in the Appendix. For cross sections

Nuclear structure (a parameterized model)

- >50 years experience / refinement
- Parameterized (Wood-Saxon potentials, derivatives)
- Lots of logical check points (e.g., *parameters are consistent with those derived from electron scattering* … *radii*, etc.), a wealth of nucleon-scattering data
- The spectroscopic factor is a 'reduced cross section' *modest corrections to account for kinematics* and spins

Distorted-wave Born Approximation, requires several ingredients and experimental consideration. ISPM cross sections.

Does it work? / model dependency? etc. \blacksquare ignored; both of the effect of slightly the effect of slightly slightl y? etc. $\sum_{i=1}^{n}$

- Need a normalization **broadening the plotted distribution**
- **EXECUTE:** Typical *uncertainty is between* **EXECUTE:** Stripping Register *+/-0.1-0.2 nucleons*
- Demonstrated in many systems (groups of isotopes/isotones) 3.5-MeV excitation is 2:8% of the total, providing an across the chart of nuclides ϵ across the chart of fluctudes ϵ

$$
S' \equiv \sigma_{\exp} / \sigma_{\text{DWBA}}
$$

$$
N_j \equiv S'/S
$$

$$
N_j \equiv \left(\Sigma G_+ S'_{\text{adding}} + \Sigma G_- S'_{\text{removing}} \right) / (2j+1) \Big| \overset{\delta}{\circ}{}_2 \Big|
$$

some of the spin assignments are ambiguous. They are

 T_{total} but is the normalization just arbitrary: • But is the normalization just arbitrary?

Ni occupancies from J. P. Schiffer et al., Phys. Rev. Lett. **108***, 022501 (2012)*

served fragments into more complicated states is likely. No

volume 32, NUMBER 3. $JULY$, 1960
 \blacksquare The diate is somewhat problematic because the problematic because the problematic because the problematic because the solution of the solution

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J. B. FRENCH
of Rochester, Rochester, New York

Analysis, e.g., 76Ge,Se

Analysis – sum rules and normalization

$$
N_j \equiv \left[\sum S'_{\rm removing} + \sum (2j+1)S'_{\rm adding}\right]/(2j+1)
$$

 $N_j \equiv \frac{[(0.45 + 0.12 + 1.29 + 0.10 + 0.11 + 0.37) + (0.44 + 0.15 + 0.12 + 0.04)]}{(2 + 4)} = 0.53$

Analysis – sum rules and normalization

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What's the significance of the normalization

PRL 111, 042502 (2013) PHYSICAL REVIEW LETTERS week ending The normalization appears meaningful, a *ubiquitous feature* of low-lying singleparticle strength, independent of *A*, *l*, nucleon type, reaction, *N*-*Z.*

Quenching of s. p. cross sections

*"Thus at any time only 2/3 of the nucleons in the nucleus act as independent particles moving in the nuclear mean field. The remaining third of the nucleons are correlated."**

Key points:

- *Academic in terms of change in occupancies*
- Arguably essential in terms of trusting the data
- How does theory handle it?

Figure 2. Schematic illustration of the distribution of the single-particle strength in stable closed-

Quenching factor (a bit more)

There are a handful of isotopes where reliable experimentally determined cross sections exist from numerous 'equivalent' probes, e.g., proton removal from ¹²C. A consistent picture emerges.

Evaluated Nuclear Structure Data File (www.nndc.bnl.gov)

Parage as Fig. 2. Same as Fig. 2. S excited state. The MFT wave function is pure 1*p*¹y², the total VMC wave function (circles) contains 1*p* and 1*f* (crosses) spectroscopic factor *S*, and the radius of the WS potential chin*a tactor the* ennig idetoij the

tion already starts at 250 MeVy*c* because in MFT the wave function is purely 1*p*1y2, whereas the VMC overlap contains four components (1*p*1y2, 1*p*3y2, 1*f*5y2, 1*f*7y2). In addition to the effect of correlations, these extra components cause an appreciable enhancement of the VMC wave function at high momentum relative to the MFT wave function. The experiment was performed with the 1% duty factor \mathcal{L}_max electron beam from the NIKHEF medium-energy accelerator and the high-resolution two-spectrometer setup in the EMIN end station [33]. The data were taken concurrently with those for the reaction ³²Ss*e*, *e*⁰ for which purpose a self-supporting disk of Li2S was used as a target (thickness roughly 25 mgycm2). The target could withstand maximum average currents of 6 mA when rotated continuously. The target thickness was monitored via frequent measurements of elastic scattering. The measurements were carried out in parallel kinematics for an outgoing proton energy of 90 MeV. As a result we needed two incident energies (329.7 and 454.7 MeV) to cover the missing momentum range of 270 to 260 MeVy*c*. Since the beam was tuned in dispersion matching mode [34] we could achieve an *Em* resolution of 180 keV (FWHM), sufficient to separate the

- Tempting to conclude it is well understood \mathbb{Z} is approximate in \mathbb{Z} ng to conclude it is well and
- **COLOGITE:** Not captured in, e.g., shell model (SM does not explicitly include SRC) absent in \mathbf{r} optiirad in a quebal modal. plared in, eigi, shen moder
- **EXECUTE: A Ab initio calculations do capture it (in light nuclei)** $\mathcal{L}_{\mathbf{w}}$ io calculations do canture it with the reduction of α dependent

discrete transitions from the two reactions. The data analysis was performed in a standard was performed in a standard way of the standard way of t α and Wiringa Phys Rev Lett **82** 4404 (1999) tions with MFT (solid) and VMC (dashed) wave functions. *Lapikas, Wessling, and Wiringa, Phys. Rev. Lett. 82, 4404 (1999)*summed strength (0.60) for both transitions and η error bars perfectly with the value of t

Quenching factor (…)

- No obvious change with neutron excess (*np* dominates) or binding energy (at least near stability) [though new results about high mtm fractions … and N/Z ratio]
- Note, there are very good (*e*,*e*'*p*) and (*e*,*e*'*n*) data on 48Ca
- Arguably not necessary to explore (*e*,*e*'*p*) [no obvious facilities … results agree with nucleon transfer]

• Does it relate to quenching of g_A ? In the sense that there is missing physics / *model space, and correlations, in the calculations*

O. Hen et al., Science 346, 614 (2014)

Aside: the g_A problem **PHYSICAL REVIEW C 97, 022501(R) (2018**)

(Shown earlier in the week)

FIG. 2. Ratios of GFMC to experimental values of the GT RMEs in the ${}^{3}H$, ${}^{6}He$, ${}^{7}Be$, and ${}^{10}C$ weak transitions. Theory predictions correspond to the χ EFT axial current in LO (blue circles) and up to N4LO (magenta stars). Green squares indicate "unquenched" shellmodel calculations from Ref. [17] based on the LO axial current.

Rapid Commun

DOI: 10.1103/PhysRevC.97.022501

terms of many-body (primarily two- and three-body) effective interactions, and with external electroweak probes via effective currents describing the coupling of these probes to individual

accurate \mathcal{M} methods, the deceptively simple picture put \mathcal{M}

range there are few microscopic calculations of Gamow-Teller

Quantum Monte Carlo calculations of weak transitions in $A = 6-10$ **nuclei**

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(Received 11 September 2017; revised manuscript received 30 November 2017; published 26 February 2018)

Ab initio calculations of the Gamow-Teller (GT) matrix elements in the β decays of ⁶He and ¹⁰C and electron captures in ⁷Be are carried out using both variational and Green's function Monte Carlo wave functions obtained from the Argonne *v*¹⁸ two-nucleon and Illinois-7 three-nucleon interactions, and axial many-body currents derived from either meson-exchange phenomenology or chiral effective field theory. The agreement with experimental data is excellent for the electron captures in ⁷Be, while theory overestimates the ⁶He and ¹⁰C data by ∼2% and ∼10%, respectively. We show that for these systems correlations in the nuclear wave functions are crucial to explaining the data, while many-body currents increase by ∼2–3% the one-body GT contributions.

$PHYSICAL REVIEW C 97.014606 (2018)$ In such an approach the nucleons interact with each other interact with $\overline{}$ and approximation of the basic and basic and basic of the basic of

Neutrinoless double- β decay matrix elements in light nuclei *A* $\frac{1}{2}$ in the set of $\frac{1}{2}$ in the set of $\frac{1}{2}$ in the set of $\frac{1}{2}$

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We present the first abjuitie calculations of neutrinoless double- β decay matrix elements in $A = 6-12$ nuclei description of the structure and dynamics of the structure and dynamics of the required σ in shell-model and Hingis 7 three-nucleon interaction. We study both light Majorana neutrino exchange and potentials arising from a large
three-nucleon interaction. We study both light Majorana neutrino exchange and potentials arising from a large lace is class of multi-TeV mechanisms of lepton-number violation. Our results provide benchmarks to be used in testing many-body methods that can be extended to the heavy nuclei of experimental interest. In light nuclei we also study the impact of two-body short-range correlations and the use such as those corresponding to different orders in chiral effective theory. transitions in the corresponding to directed orders in this $\sum_{i=1}^{n}$

Carlo (VMC) method in Ref. [13]. It used nuclear axial currents

… old results

J. P. Schiffer et al., Phys. Rev. Lett. 100, 112501 (2008) [neutrons] BPK et al., Phys. Rev. C 79, 021301(R) (2009) [protons]

Argonne

POLYPROBY

Change in occupancy

J. P. Schiffer et al., Phys. Rev. Lett. 100, 112501 (2008) [neutrons] BPK et al., Phys. Rev. C 79, 021301(R) (2009) [protons] Rodin et al., Nucl. Phys. A 766, 107 (2006) [A] Suhonen et al., Phys. Lett. B 668, 277 (2006) [B] Caurier et al., Phys. Rev. Lett.100, 052503 (2008) [C]

Error bars are dominated by the systematic uncertainties relating to the analysis (Does not include more recent IBM results)

Impact?

Yes, some. Much discussed. A *40-70% reduction* in the well-known gap between QRPA and the ISM, resulted. This predated recent IBM work and newer calculations.

Šimkovic et al., Phys. Rev. C 79, 055501 (2009) [quote]

A = 100 occupancies Peaks corresponding to reactions on carbon and oxygen $\Lambda = 100$ secured m — TVV VLLUNAI their larger kinematic shift. These contaminant peaks obscured

J. Suhonen and O. Civitarese, Nucl. Phys. A 924, 1 (2014) [WS, WS ADJ] J. Kotila and J. Barea, Phys. Rev. C 94, 034320 (2016) [IBM]

groups of interest at some angles, but the difference in their kinematic shifts meant that angles were always available where EXPERIMENTAL STUDY OF THE REARRANGEMENTS OF . . . PHYSICAL REVIEW C **96**, 054325 (2017) *High level density, Munich Q3D (as good as 8-keV FWHM resolution*

FIG. 2. Spectra of deuterons from the (*p,d*) reaction on targets of 98Mo, 100Mo, 100Ru, and 102Ru at a laboratory angle of 6◦ as a *Freeman et al., Phys. Rev. C 96, 054325 (2017)* function of the excitation energy in the residual nucleus. The position energy in the positions \mathcal{L}

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A. Neacsu and M. Horoi, Phys. Rev. C 91, 024309 (2015) [SM1] J. Menéndez et al., Nucl. Phys. A 818, 139 (2009) [SM2] J. Kotila and J. Barea, Phys. Rev. C 94, 034320 (2016) [IBM] J. P. ENTWISLE *et al.* PHYSICAL REVIEW C *J. Suhonen* **93**, 064312 (2016) *and O. Civitarese, Nucl. Phys. A 847, 207 (2010) [QRPA]*

$A = 136$ occupancies sections, $A = \int_{A K \text{ of } B} A \text{ and } B$ U dicies were targets the second secon There were estimated to be around 5%. Statistical uncertainties were estimated to be a

integrator. These are not trivial to estimate, and so we place a conservative estimate of ∼20%. The relative uncertainties **IV. DISCUSSION** *A. Neacsu and M. Horoi, Phys. Rev. C 91, 024309 (2015) [SM1]* t al., Nucl. Phys. A **818**, 139 (2009) [SM2] |
ea. Phys. Rev. C **94**, 034320 (2016) [IRM] | *J. Kotila and J. Barea, Phys. Rev. C 94, 034320 (2016) [IBM]* $\overline{}$ *J. Menéndez et al., Nucl. Phys. A 818, 139 (2009) [SM2]*

Taking advantage of ¹³⁶Xe being a good closed shell for neutrons removing and ∼5% for the neutron adding, becoming *>* 10% also show that the ν0*h*¹¹*/*² accounts for most of the vacancy. \boldsymbol{u} under \boldsymbol{v}

Comment on occupancies

• The agreement is perhaps qualitatively okay, in some instances within the uncertainties *(but not for both protons and neutrons)*, *but quite poor on the whole*

> Ø *We can ask whether it matters? … it does, regardless of how (in)sensitive the NME is to the change in occupancies*

While likely challenging theoretically, it would be interesting to know what the consequence of 'shutting off' part of the model-space would be (*e.g. the g_{7/2} neutrons above tin, or comparing e.g. JUN45 and GXPF1A – or even more simplistic approaches*)

Pairing properties

Can the ground states of the candidates be described as 'seas' of correlated 0+ paired protons and neutrons?

 \blacksquare

Pairing around A ~ 76

Pair-transfer reactions are a simple and effective probe of pairing correlations *No evidence of 'pairing vibrations' in the A = 76 region*

a. 8. Procinal et al., Phys. Rev. 899, 891861(R) (2007) [healtor
A. Roberts et al., Phys. Rev. C **87**, 051305(R) (2013) [protons] to *L* and *L*

Pairing around A ~ 100 the spectrometer by a multiwire gas proportional counter \mathbf{b} \mathbf{b} \mathbf{c} \mathbf{b} \mathbf{c} \mathbf{c} plane prosition position position of the projection of the position of the Particle of the Particle residual t identification was accomplished with the combination of the 1_a 100 peaked angular distribution, at θcm = 0◦, with all other *L*

 $s_{\rm 1D}$. The objective was to identify α and to accurately measure the intervalse \mathbb{R}^n high energy resolution. Any significant *differences* between the reactions on 100Mo and 100Ru would indicate different pairing properties of the nuclei connected through the 0ν2β decay $m_{\rm eff}$ element, which must be accounted for in calculations. The (*p,t*) reaction was also measured on a target $\sigma_{\rm 102R}$

The (*p,t*) reaction was measured on four isotopically enriched targets of 100Mo (97*.*39%), 98Mo (97*.*18%), 100Ru (96*.*95%), and 102Ru (99*.*38%). The proton beam was delivered by the MP tandem accelerator of the Maier-Leibnitz- $L_{\rm{H}}$ of the Ludwig-Maximilians-Universitation and Technische Universität M α The typical beam current on target was ∼450 nA and was recorded by a Faraday cup. The tritons were momentum analyzed using a one quadrupole lens and three dipole \mathbb{R}^3 magnetic spectrograph. Separate elastic scattering measure $m = 25$ on each target with a 12 MeV 3He beam to determine the product of target thickness and the solid angle subtended by $t_{\rm eff}$ spectrograph aperture. Such a measurement is α energy regime of Rutherford elastic scattering and is necessary to convert triton yields from the (*p,t*) reaction to absolute cross

Charged particles were detected at the focal plane of the spectrometer by a multiwire gas proportional counter backed by a scintillator, providing measurements of focalplane position, energy loss, and residual energy. Particle identification was accomplished with the combination of the magnetic-field settings of the spectrograph \mathcal{A} charged particles from competing reactions have sufficiently different rigidities—and the focal-plane energy signals. The focal-plane position was determined from the readout of 255 cathode pads on the gas proportional counter. Each pad has

Λ transitional region with deformation $\overline{\Lambda}$ grated beams and target thicknesses for, respectively, (a) 102Ru(*p,t*), $\overline{}$ A transitional region with deformatic corresponding integrated beam currents and target thicknesses. FIG. 2. (Color online) Cross-section ratios of the states populated <u>A transitional region with deformation playing a role in the nuclear structure:</u>

- \overline{P} , **Popetions loading to and from** 100E **transitional region) 988 Proportions reading to and notified the 100 motified to t** excitation-energy region. <u>bration are 4 in the definition</u> $\frac{1}{2}$ and $\frac{1}{2}$ are $\frac{1}{2}$ in this interest are for $\frac{1}{2}$ in the letter $\frac{1}{2}$ are for $\frac{1}{2}$ in the letter $\frac{1}{2}$ in the letter $\frac{1}{2}$ in the letter $\frac{1}{2}$ • Reactions leading to and from 100 Ru show ~95% of the L=0(p,t) strength is in the g.s. (on the spherical side of the $\,$ $\tau_{\rm{0}}$ state becomes the band, which does the band, which does not a deformed band, which
	- For $^{100}\rm{Mo}$ about 20% of the L=0(p,t) strength is an excited 0⁺, a shape-transitional nucleus when protons occupy the 1*g*⁹/₂ orbital and neutrons occupy the 1*g*⁹/₂ orbital and neutrons occupy the 1*g*⁹/₂ or
	- *No evidence for pairing vibrations, but structure is complicated (proton work remains to be done)*

J. S. Thomas et al., Phys. Rev. C 86, 047304 (2012)

101

Pairing around A ~ 130, 136 uncertainty of ∼ 7%, while those quoted for proton-pair addition are $\mathbf{1} \bullet \bullet \mathbf{1}$

From the proton-pair adding Te(³He,n) reactions by Alford et al., significant strength is seen in ℓ = 0 transitions to excited states … $\ddot{}$ between $\ddot{}$ and $\ddot{}$

A *classic case of pair vibration* and likely a consequence of a sub-shell gap at *Z* = 64 **Consequences for QRPA?** (Does the shell-model include this feature also?) -model include this feature also ℓ)

T. Bloxham et al., Phys. Rev. C 82, 027308 (2010) [neutrons] W. P. Alford et al., Nucl. Phys. A 323, 339 (1979) [protons] \overline{a} $\begin{array}{ccc} \n\hline\n\vdots \\
\hline\n\end{array}$ $\frac{p}{\sqrt{2}}$

Summary

Experimental nuclear-structure data is an essential part of the story of the NME challenge

The candidates are *not 'generically similar' systems* (pairing, e.g. *Z* = 64, closed shells, deformation, etc., all different in each case)

'Traditional' calculations do not reliably reproduce information extracted from experiments *(what level of agreement should we expect?)*

New ab initio calculations likely essential (model space, interactions, Hamiltonians, correlations, weak currents, *all still being worked on*)

*E*0 lifetimes, Ge and Se (approved exp. at TRIUMF) Revisiting two-proton transfer [Xe(3He,*n*) at ANL] Other programs (e.g. *RCNPs CE reactions*, Catania DCE) And … many more (most covered this week)

