



Canada's national laboratory  
for particle and nuclear physics  
and accelerator-based science

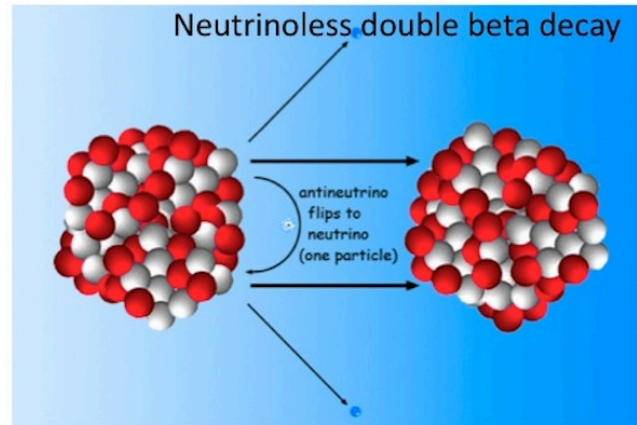
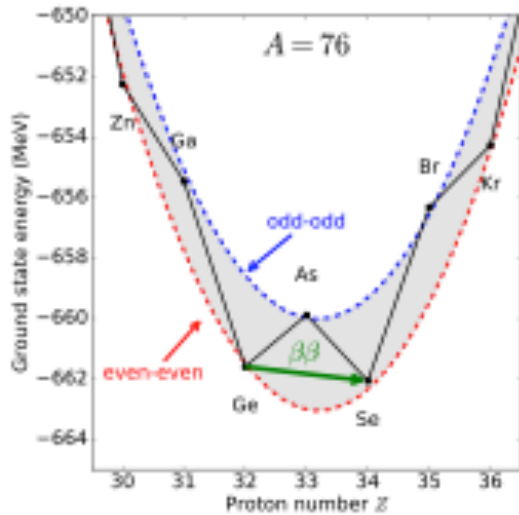
# Ab initio calculations of double-beta decay and WIMP-nucleus scattering

Jason D. Holt

INT, 6 March 2018



Neutrino own antiparticle  $\iff 0\nu\beta\beta$  decay

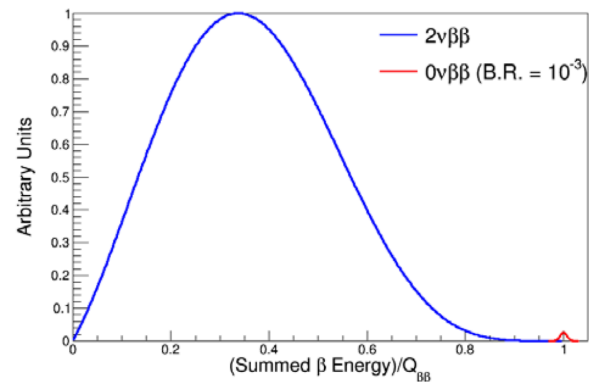


**Tremendous impact on BSM physics:**

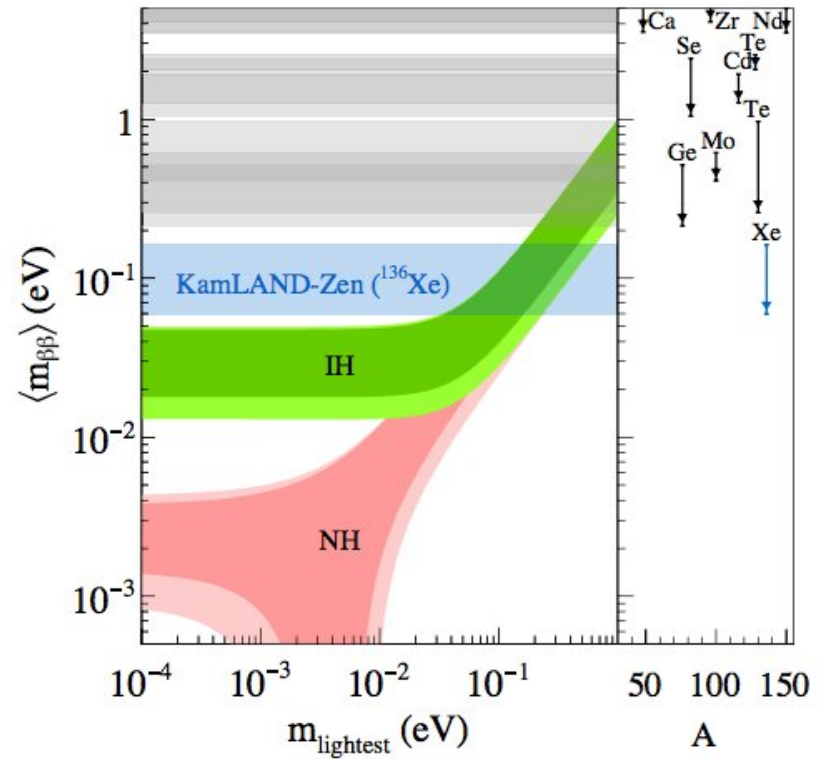
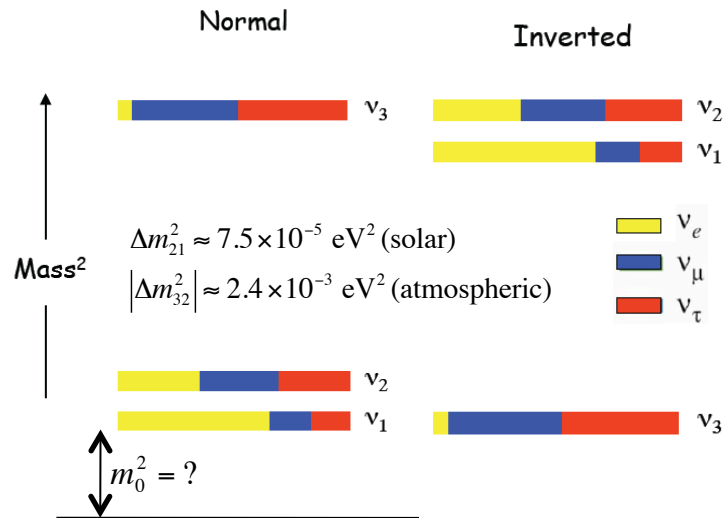
**Lepton-number violating process**

Majorana character of neutrino

Absolute neutrino mass scale

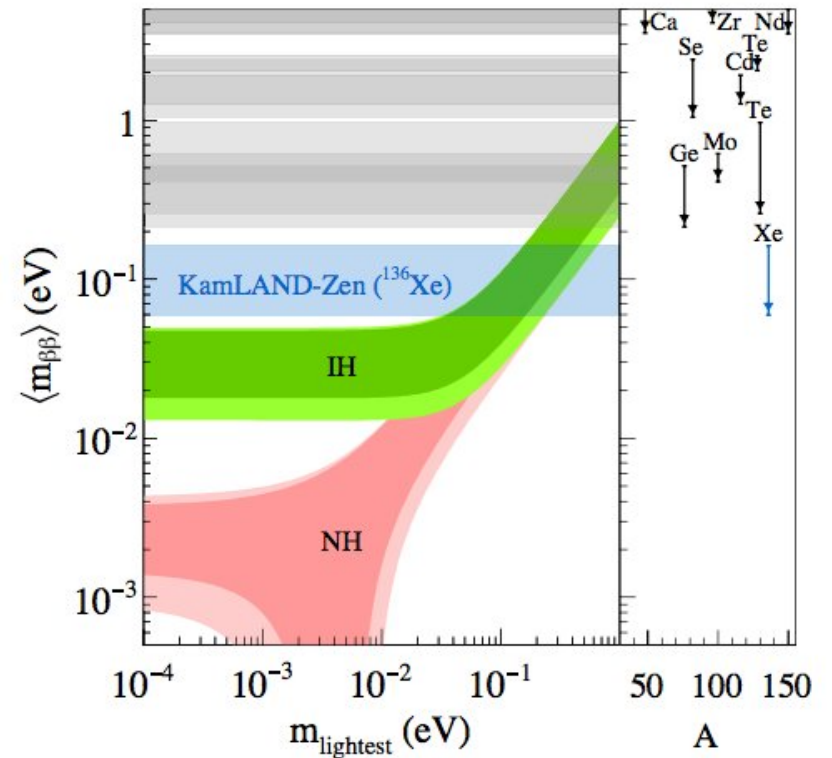
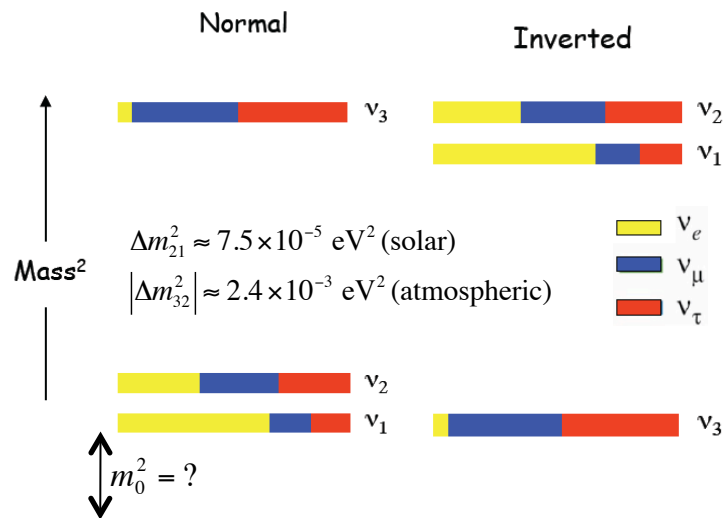


## Progress in large-scale searches pushing towards IH



$$\left(T_{1/2}^{0\nu\beta\beta}\right)^{-1} = G^{0\nu} |M^{0\nu}|^2 \langle m_{\beta\beta} \rangle^2 \langle m_{\beta\beta} \rangle = \left| \sum_{i=1}^3 U_{ei} m_i \right|$$

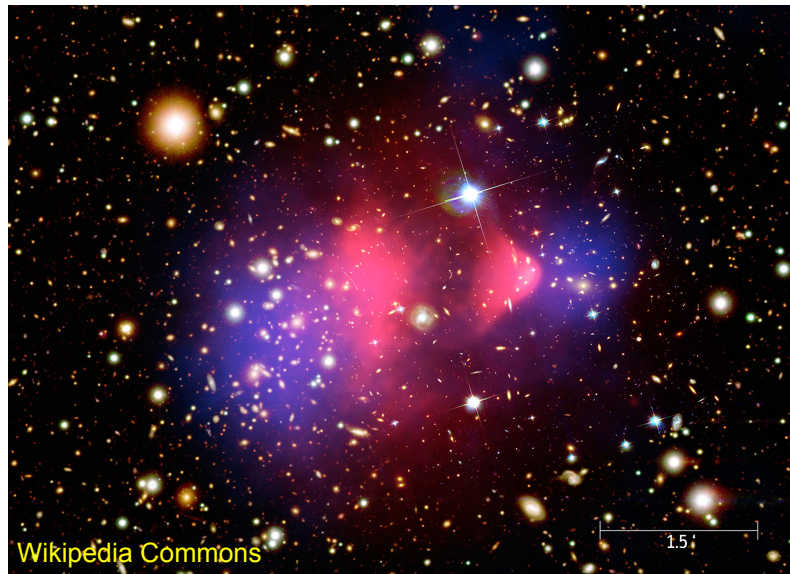
Progress in large-scale searches pushing towards **IH**



$$\left(T_{1/2}^{0\nu\beta\beta}\right)^{-1} = G^{0\nu} \left[M^{0\nu}\right]^2 \langle m_{\beta\beta} \rangle^2 \langle m_{\beta\beta} \rangle = \left| \sum_{i=1}^3 U_{ei} m_i \right|$$

Uncertainty from **Nuclear Matrix Element**; bands do not represent rigorous uncertainties

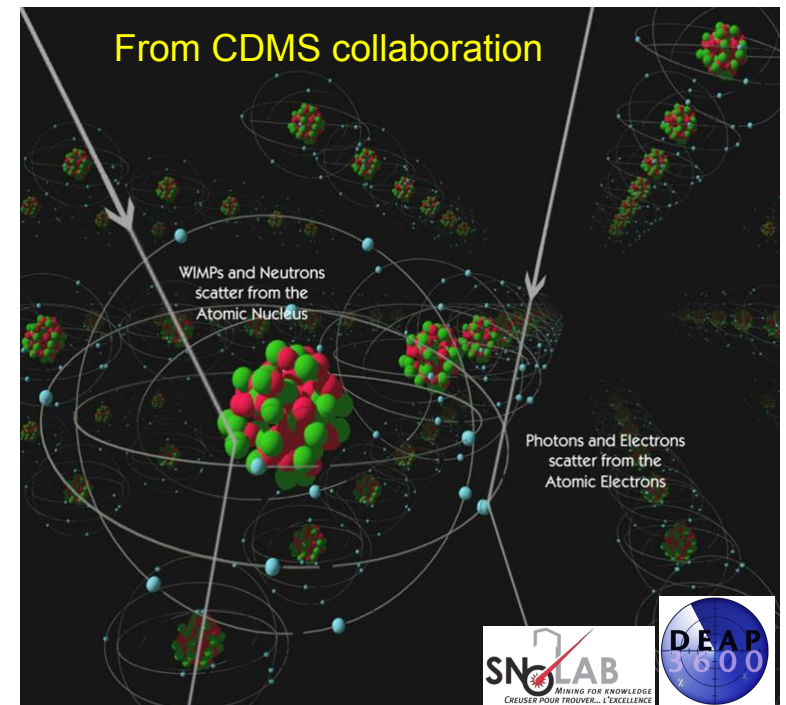
Many direct-detection searches underway worldwide



Direct detection:  $X_{SM} \rightarrow X_{SM}$

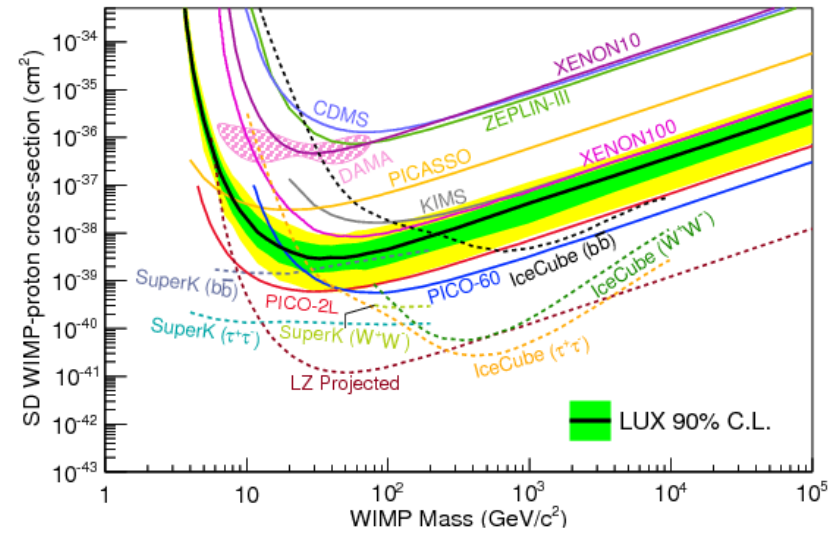
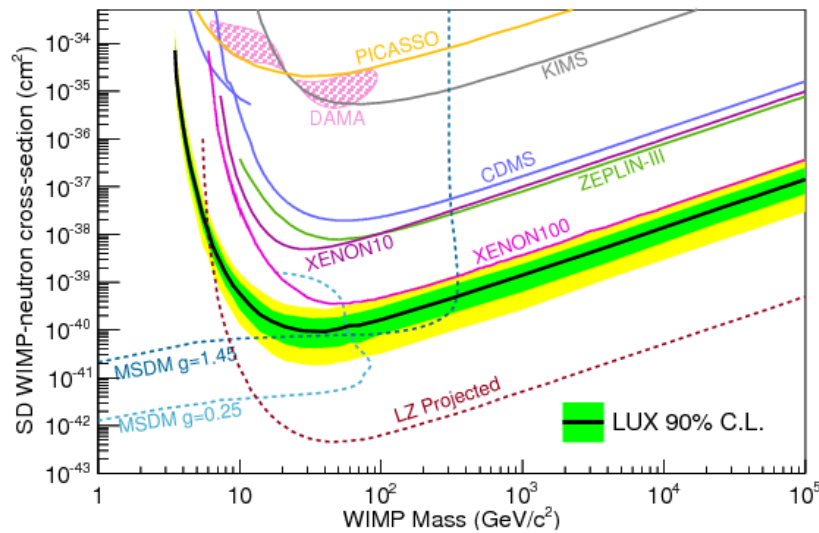
Leading candidates: neutralinos

Couples primarily to scalar and axial-vector currents in atomic nuclei



Observation of nuclear recoil

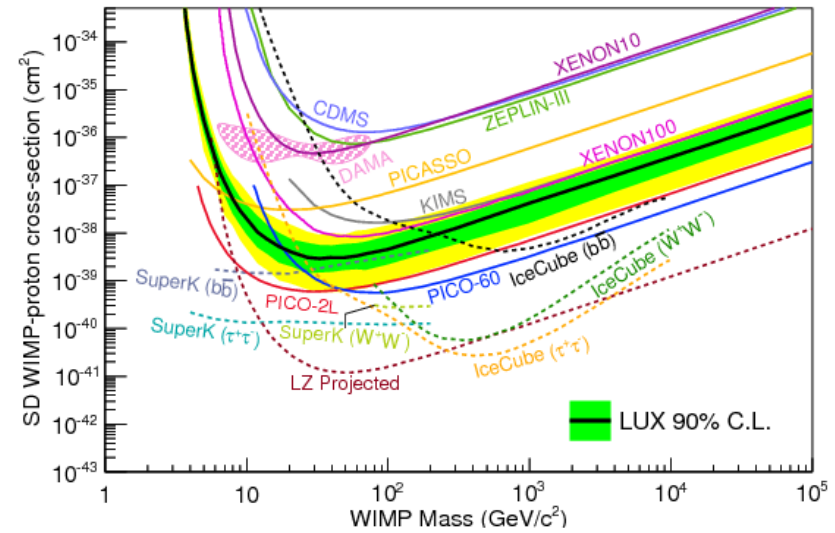
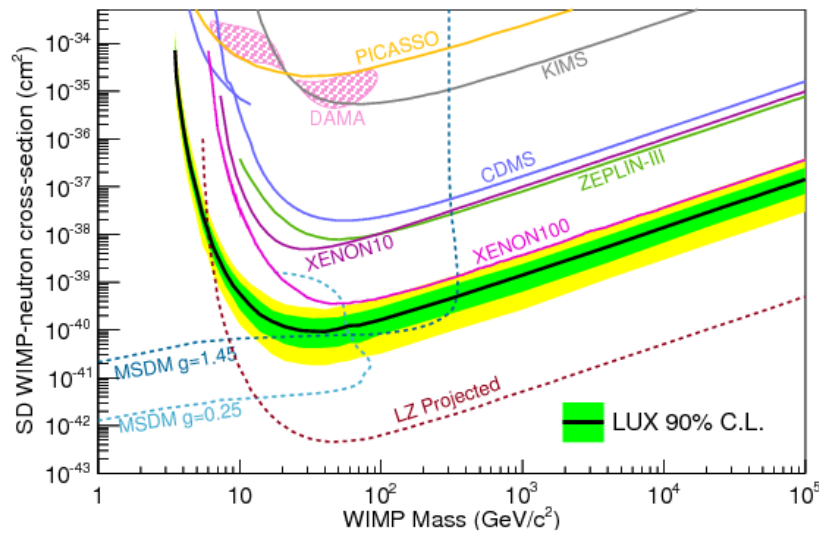
## Exclusion plots for WIMP-nucleon total cross section (spin-dependent)



## Differential cross section: compare results from different targets

$$\frac{d\sigma}{dp^2} = \frac{8G_F^2}{(2J_i + 1)v^2} S_A(p)$$

## Exclusion plots for WIMP-nucleon total cross section (spin-dependent)

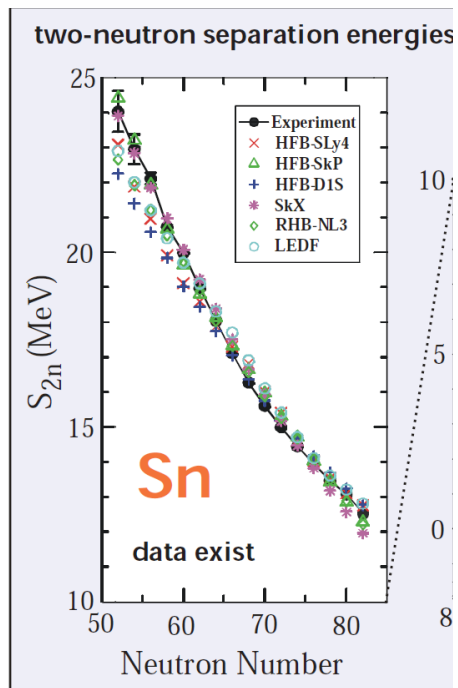


Differential cross section: compare results from different targets

$$\frac{d\sigma}{dp^2} = \frac{8G_F^2}{(2J_i + 1)v^2} S_A(p)$$

**Structure functions** required from nuclear theory

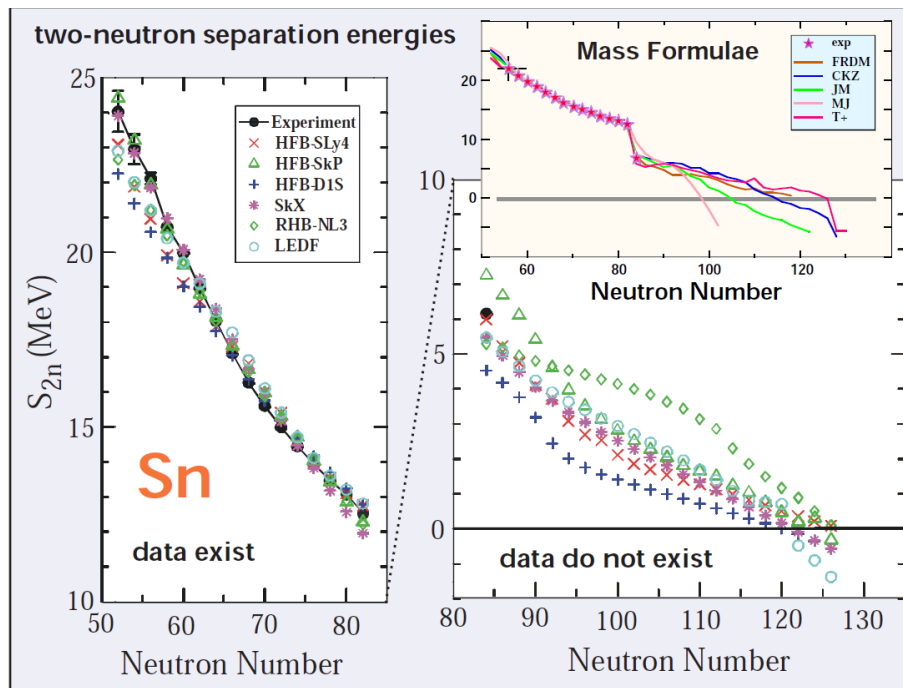
How well can nuclear models motivate experiments, predict beyond data?



Work well in regions where  
informed by data



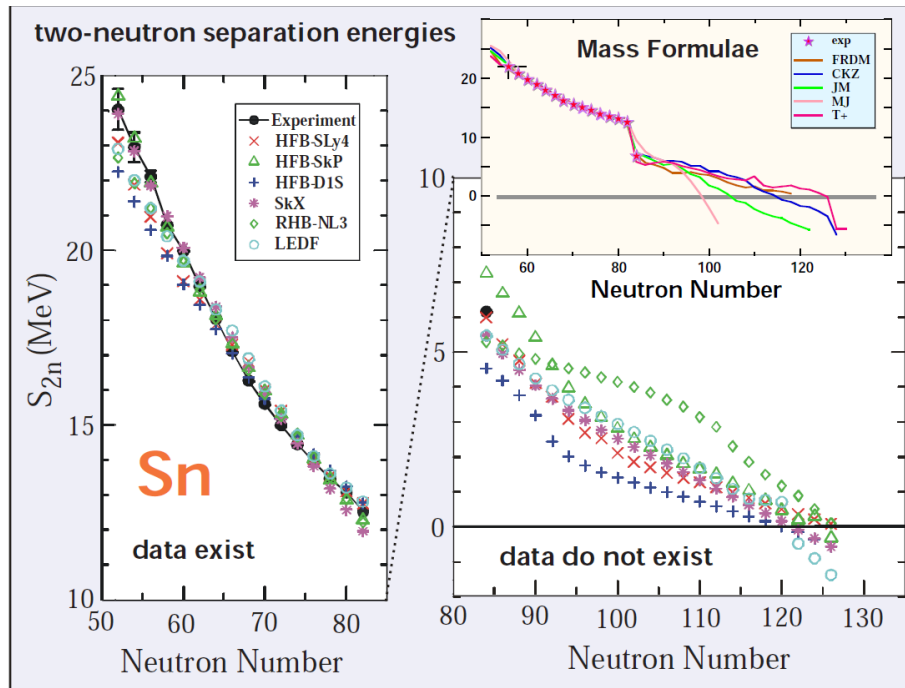
How well can nuclear models motivate experiments, predict beyond data?



Often extrapolate unreliably

Spread in results = meaningful uncertainty?

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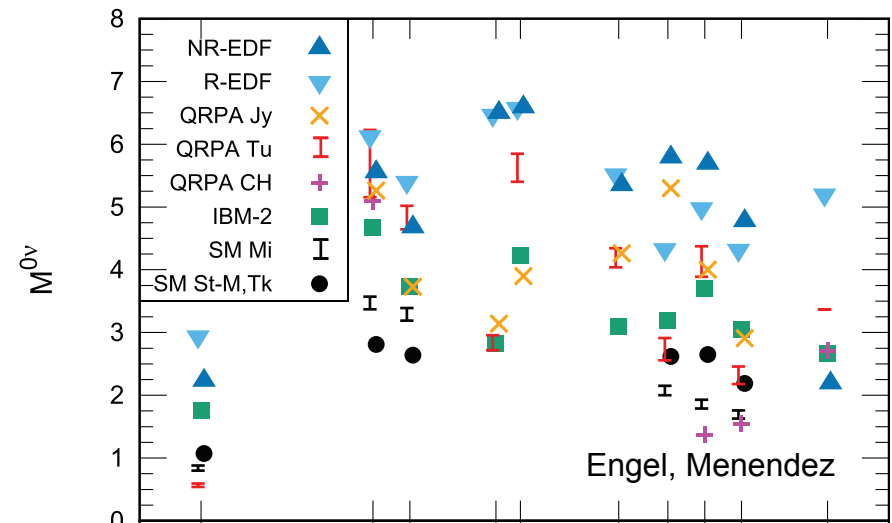


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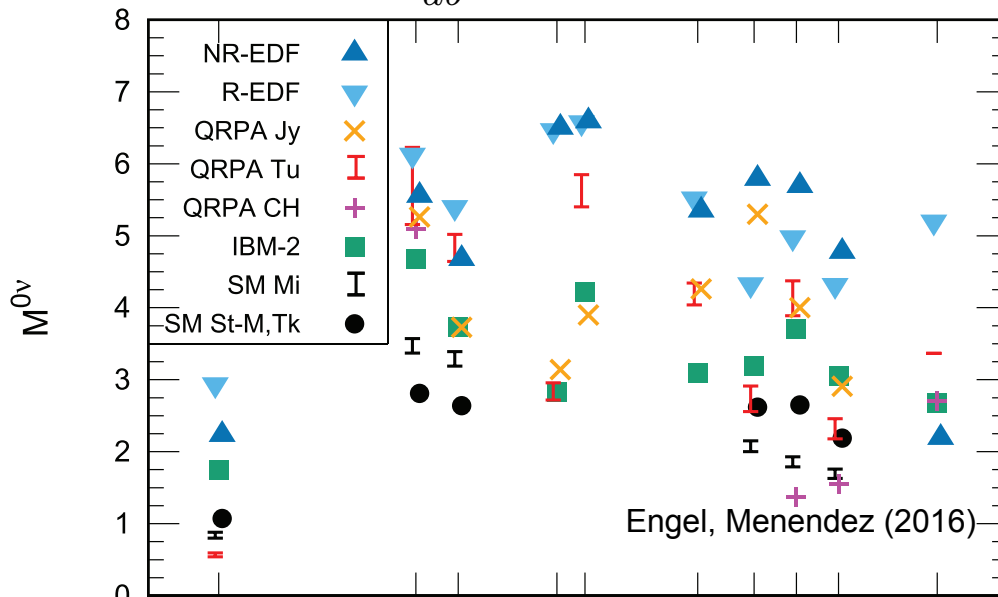
Analogous picture in  $0\nu\beta\beta$  decay

$$M_{GT}^{0\nu} = \langle f | \sum_{ab} H(r_{ab}) \sigma_a \cdot \sigma_b \tau_a^+ \tau_b^+ | i \rangle$$



All calculations to date from **extrapolated** phenomenological models; large spread in results

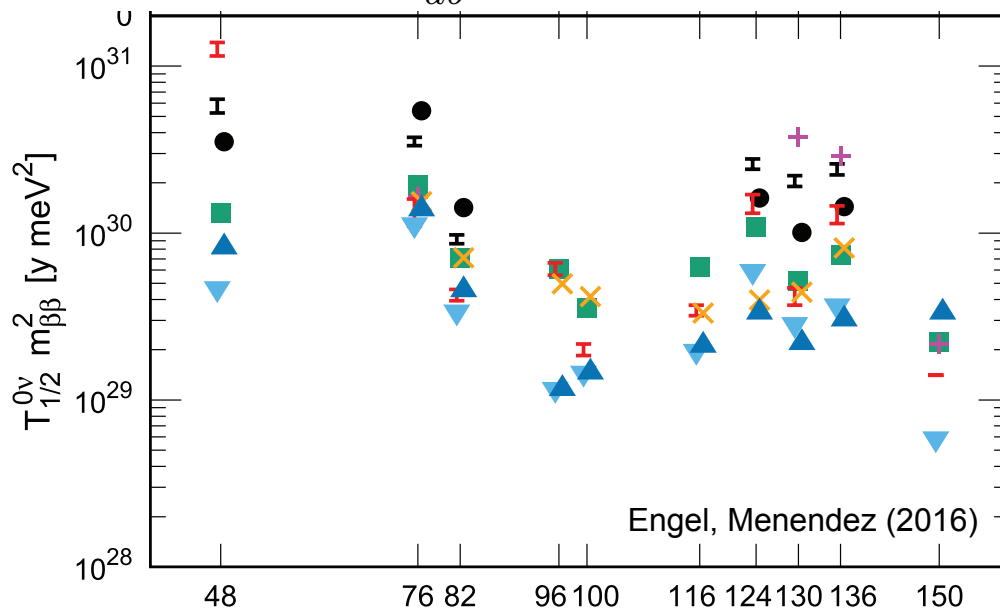
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**All models missing essential physics**

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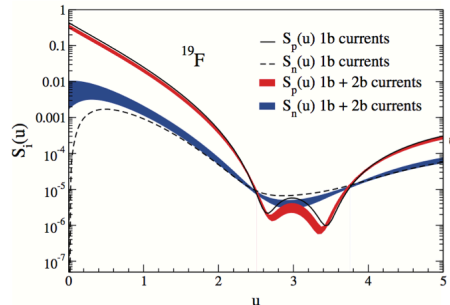
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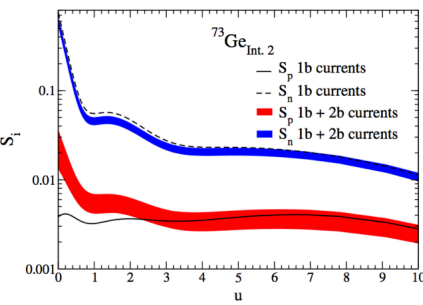
**All models missing essential physics**

Phenomenological wfs + inconsistent bare operator (**with two-body currents**)

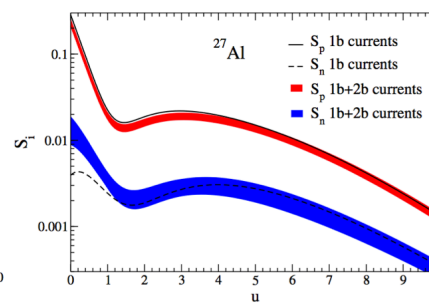
$$S_A(p) = \sum_{L \geq 0} |\langle J_f || \mathcal{L}_L || J_i \rangle|^2 + \sum_{L \geq 0} \left( |\langle J_f || \mathcal{T}_L^{\text{el}} || J_i \rangle|^2 + |\langle J_f || \mathcal{T}_L^{\text{mag}} || J_i \rangle|^2 \right)$$



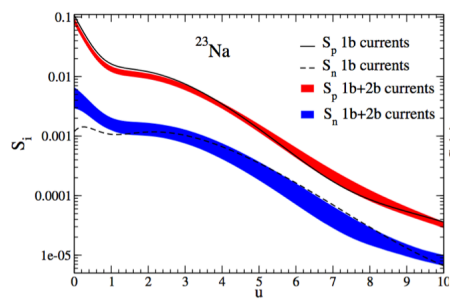
PICASSO, COUPP, SIMPLE



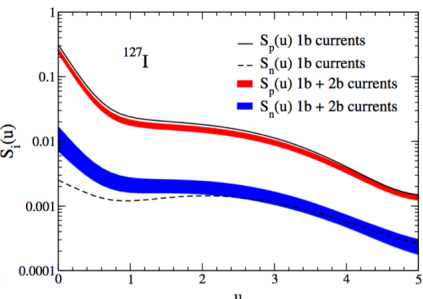
CDMS, EDELWEISS, EURECA



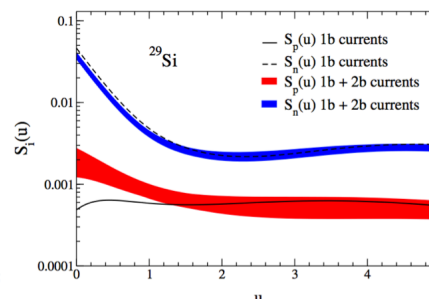
CRESST



DAMA, ANAIS, DM-Ice



DAMA, ANAIS, DM-Ice, KIMS

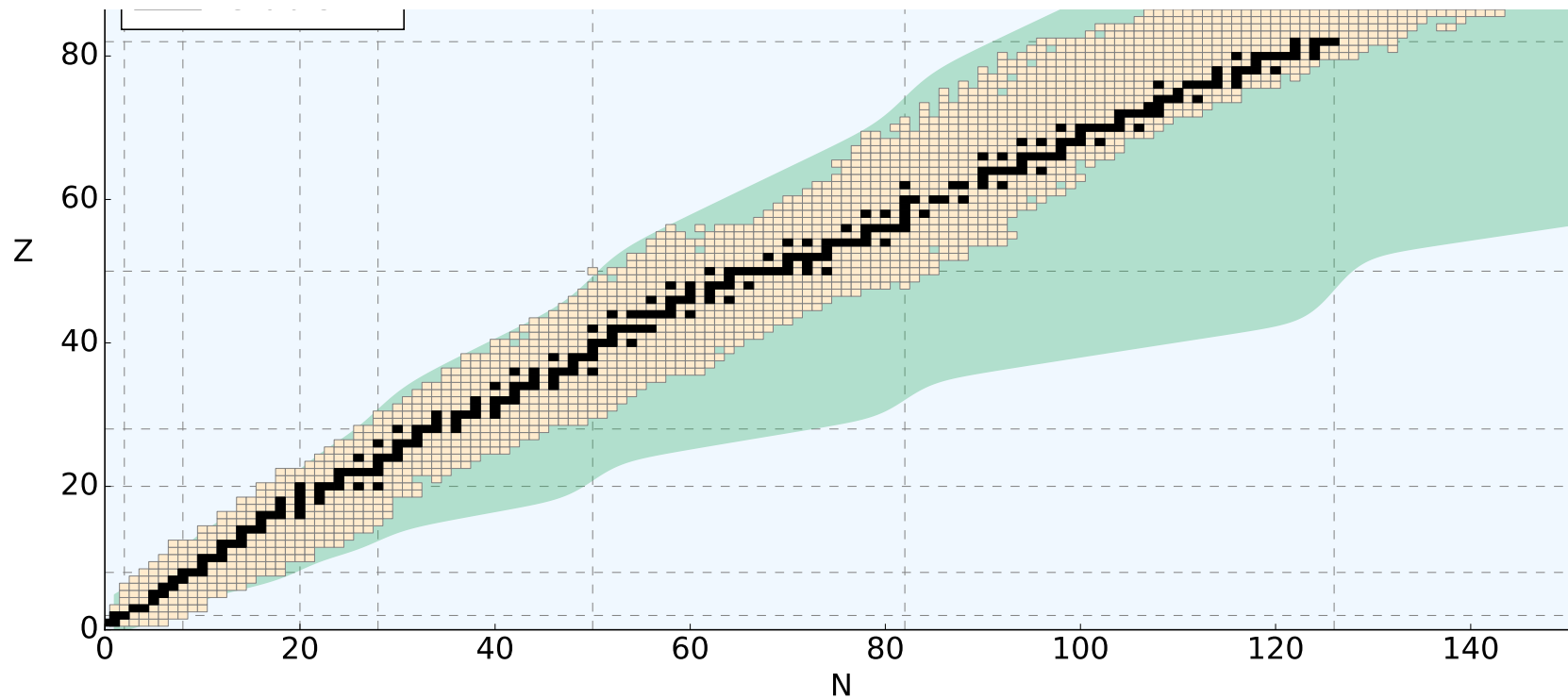


CDMS-II

Klos et al, PRD (2013)

Aim of modern nuclear theory: Develop unified *first-principles* picture of structure and reactions

$$H\psi_n = E_n\psi_n$$

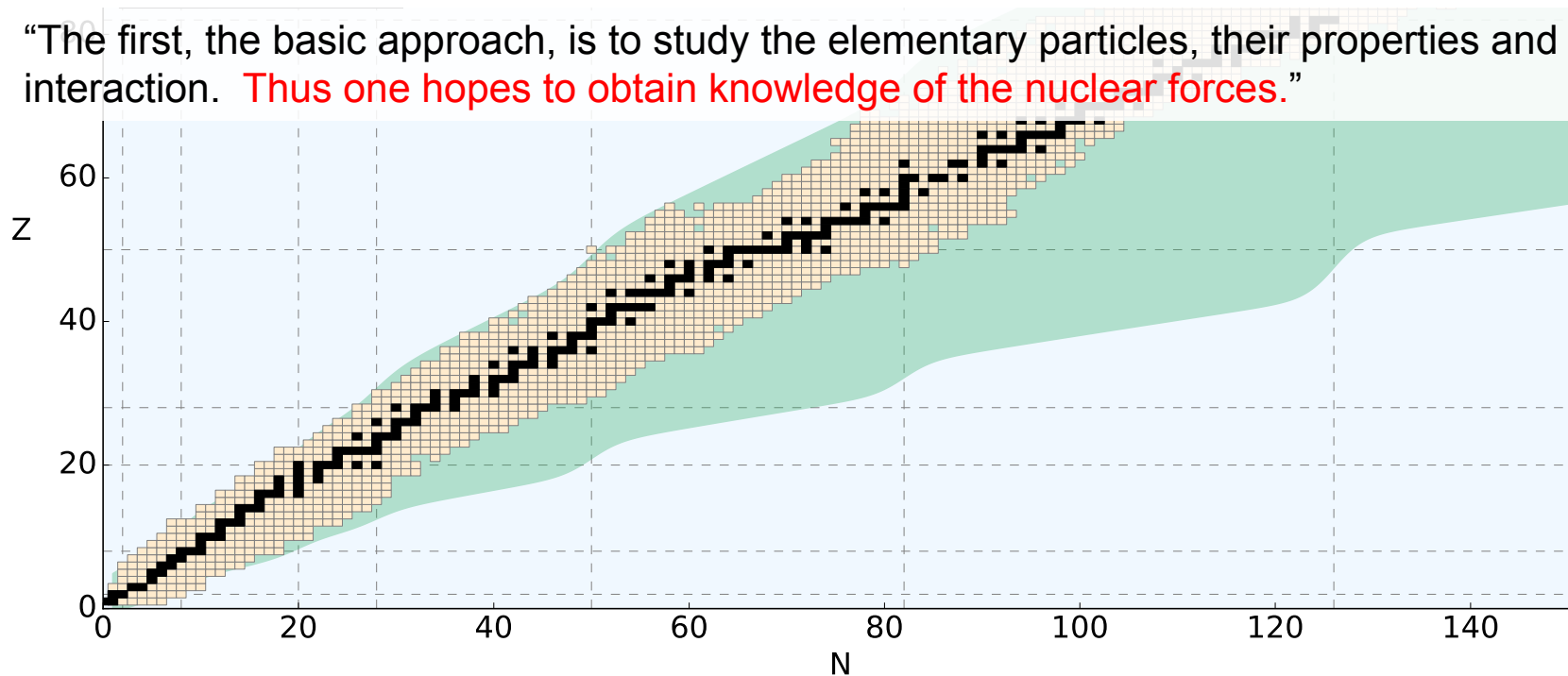


**Aim of modern nuclear theory:** Develop unified *first-principles* picture of structure and reactions

- Nuclear forces (low-energy QCD)
- Electroweak physics

$$\boxed{H}\psi_n = E_n\psi_n$$

“The first, the basic approach, is to study the elementary particles, their properties and mutual interaction. **Thus one hopes to obtain knowledge of the nuclear forces.**”



**Chiral effective field theory:** systematic expansion of nuclear interactions

$$\boxed{H}\psi_n = E_n\psi_n$$

	NN	3N	4N
LO $\mathcal{O}\left(\frac{Q^0}{\Lambda^0}\right)$		—	—
NLO $\mathcal{O}\left(\frac{Q^2}{\Lambda^2}\right)$		—	—
N <sup>2</sup> LO $\mathcal{O}\left(\frac{Q^3}{\Lambda^3}\right)$			—
N <sup>3</sup> LO $\mathcal{O}\left(\frac{Q^4}{\Lambda^4}\right)$			

derived in (1994/2002)

(2011) (2006)

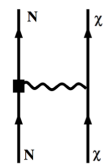
Nucleons interact via contact and pion exchanges

Undetermined **low-energy constants** fit to 2,3,4-body data

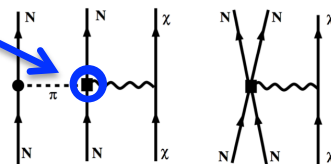
Consistent treatment of NN, 3N, 4N... forces

Quantitative estimation of neglected orders possible

one-body currents at  $Q^0$  and  $Q^2$



+ two-body currents at  $Q^3$



Consistent EW interactions

Quantifiable uncertainties possible

Best fitting strategy for ~30

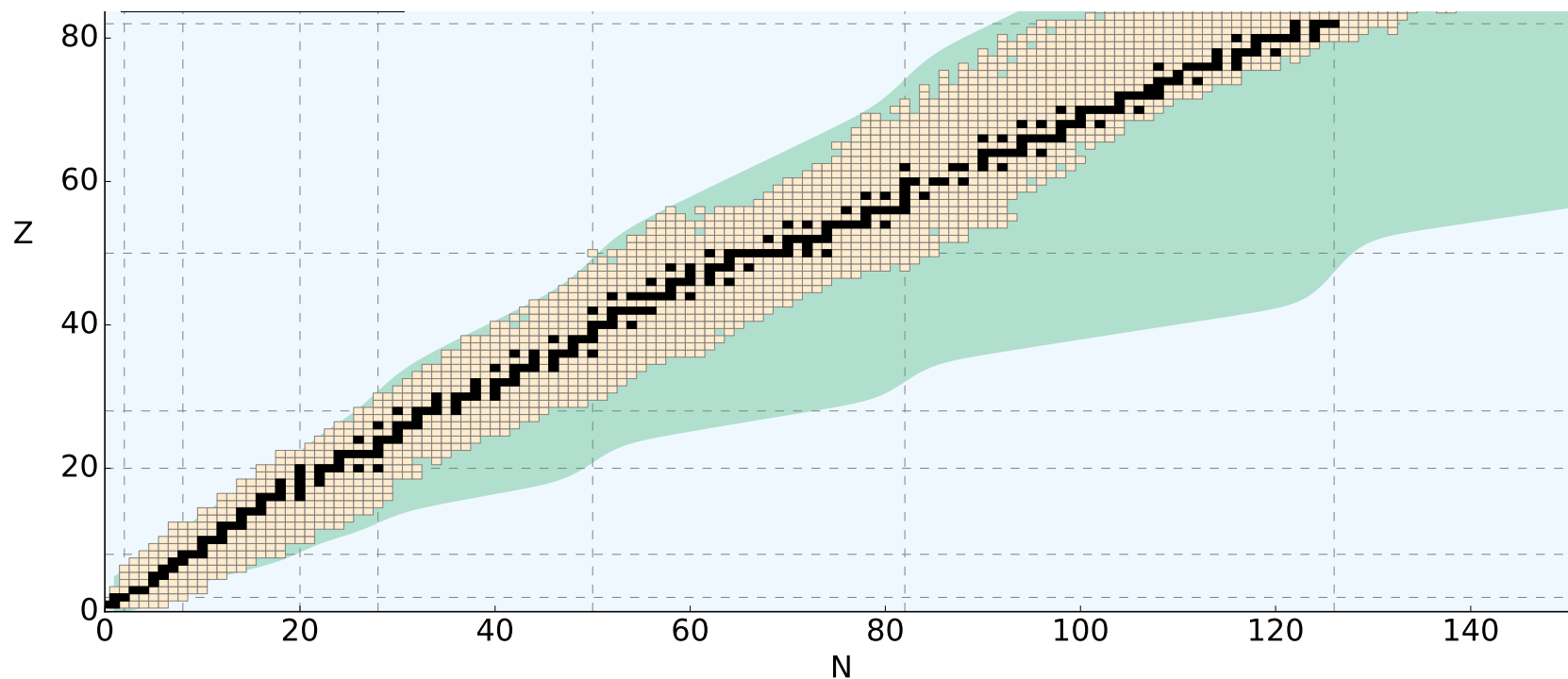
undetermined couplings debated



**Aim of modern nuclear theory:** Develop unified *first-principles* picture of structure and reactions

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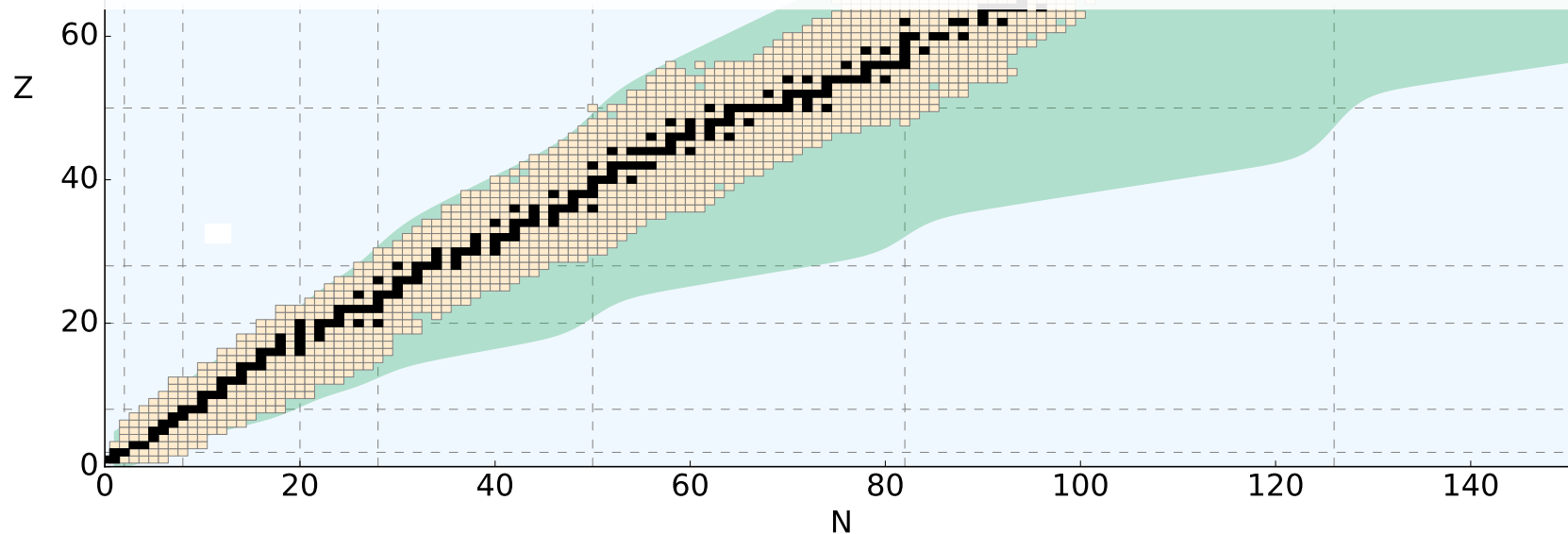
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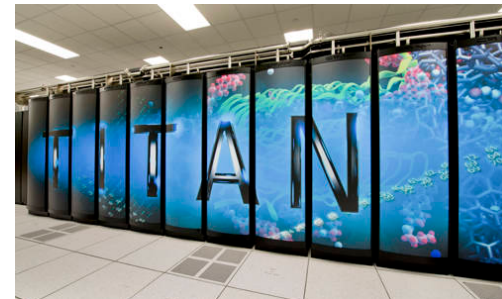
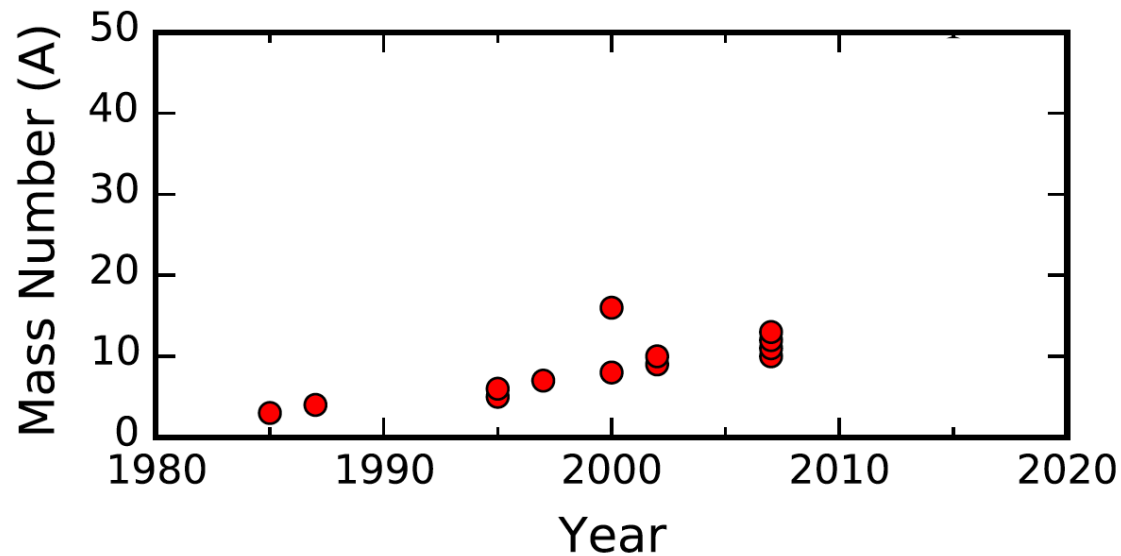
- **Nuclear many-body problem**

“If the forces are known, one should, in principle, be able to calculate deductively the properties of individual nuclei.”



Moore's law: exponential growth in computing power

Methods for light nuclei (QMC, NCSM) scale exponentially with mass

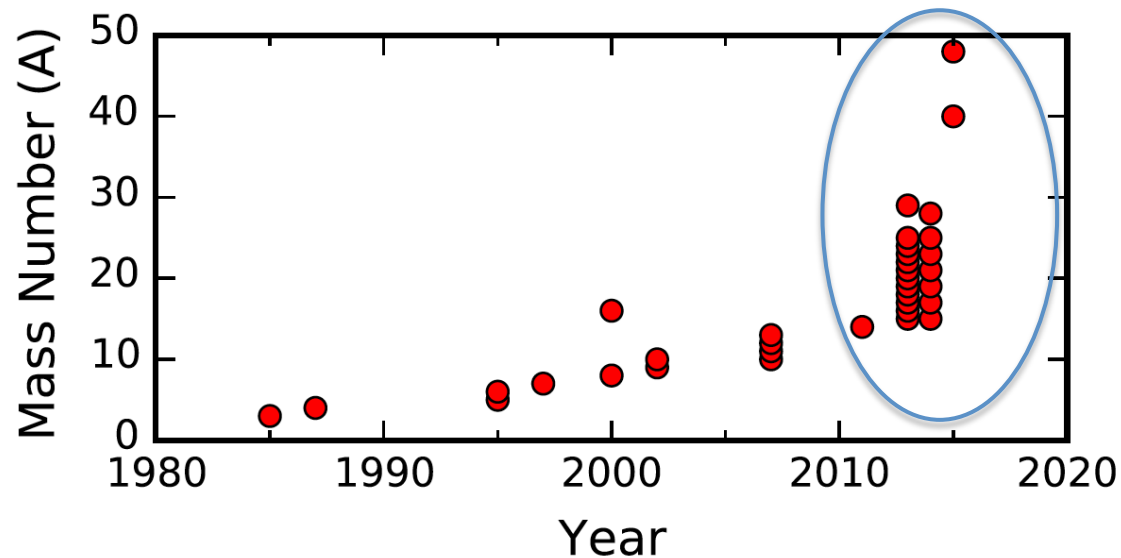


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Mid 2000's **polynomial scaling methods developed** (coupled cluster, in-medium SRG,...)

Explosion in limits of ab initio theory

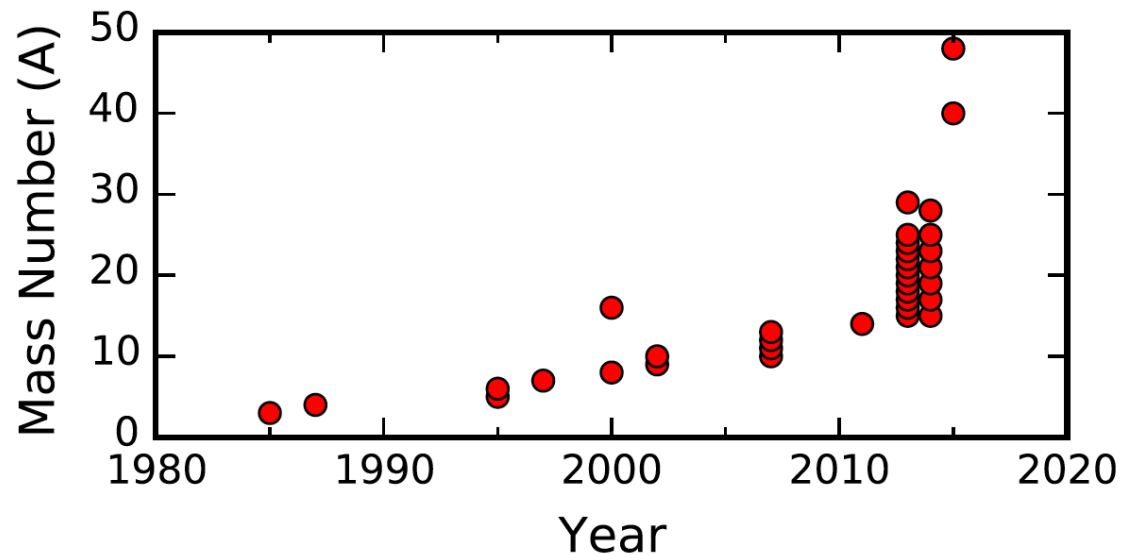


Moore's law: exponential growth in computing power  ← **2017:  $A > 100$**

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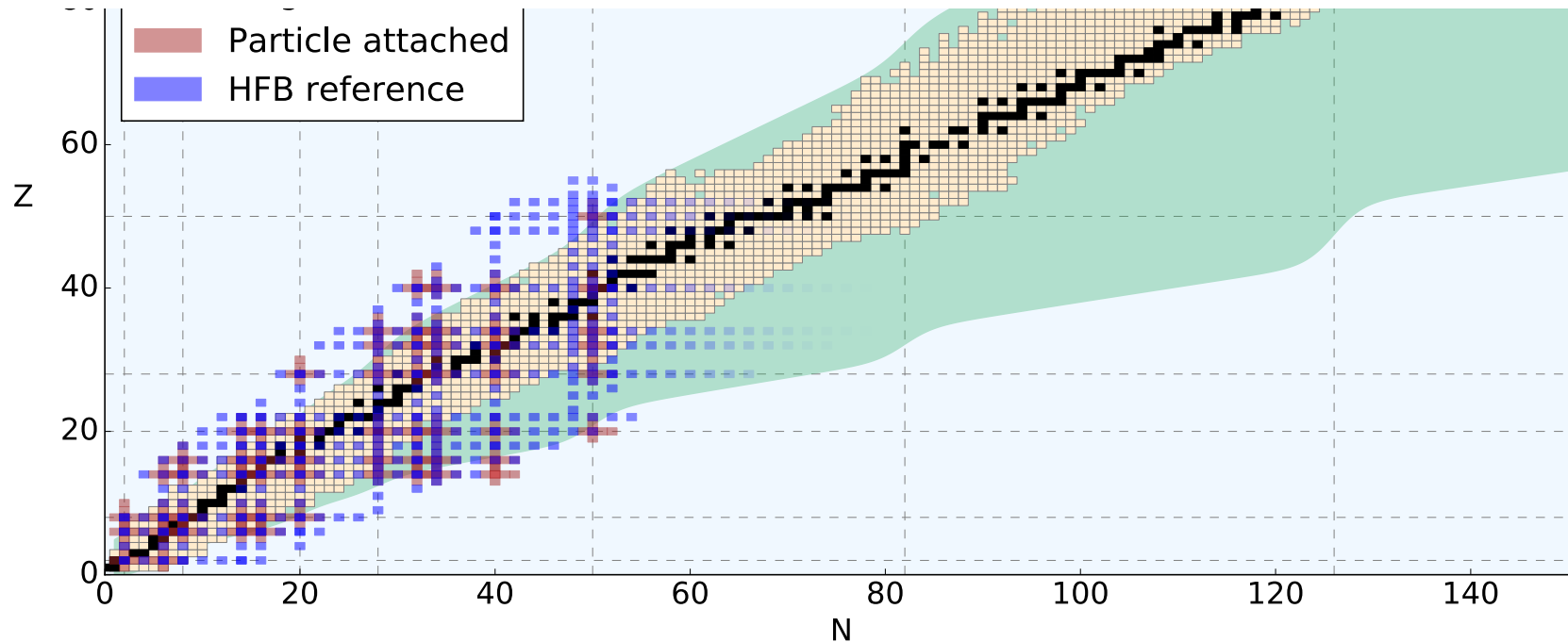


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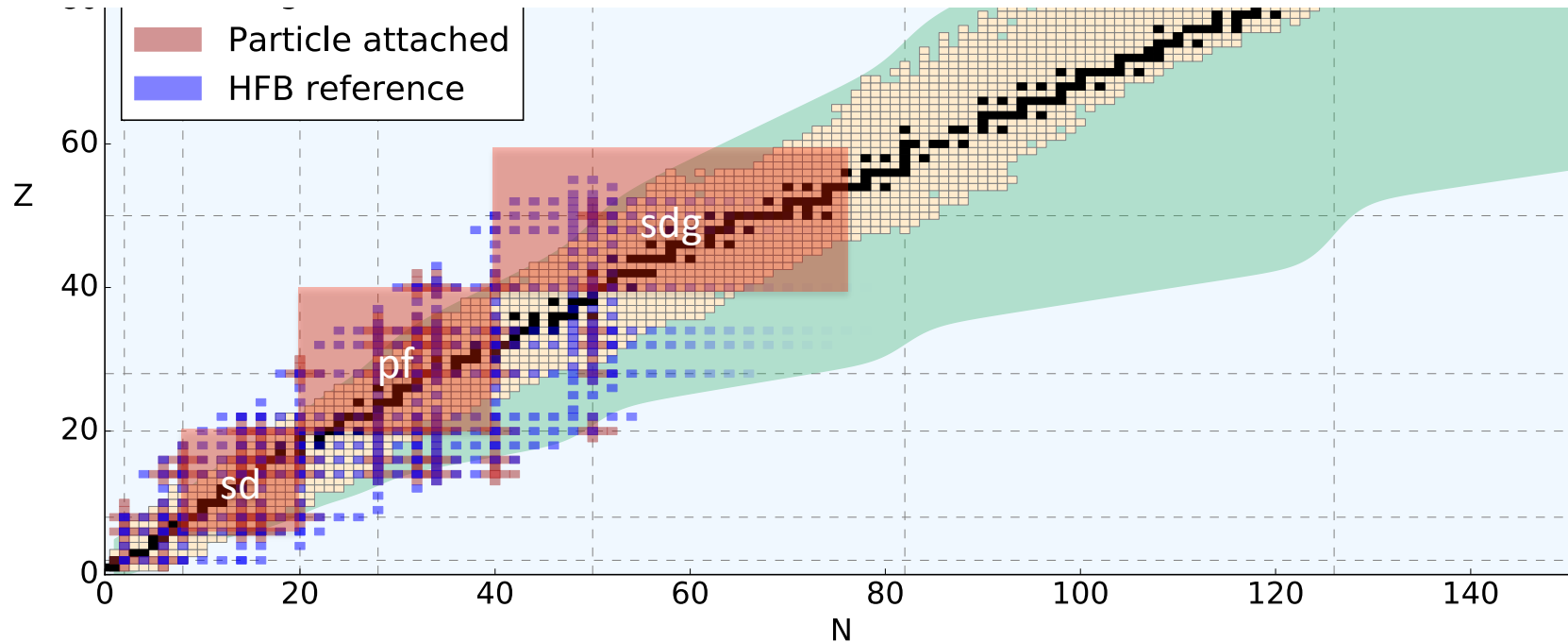


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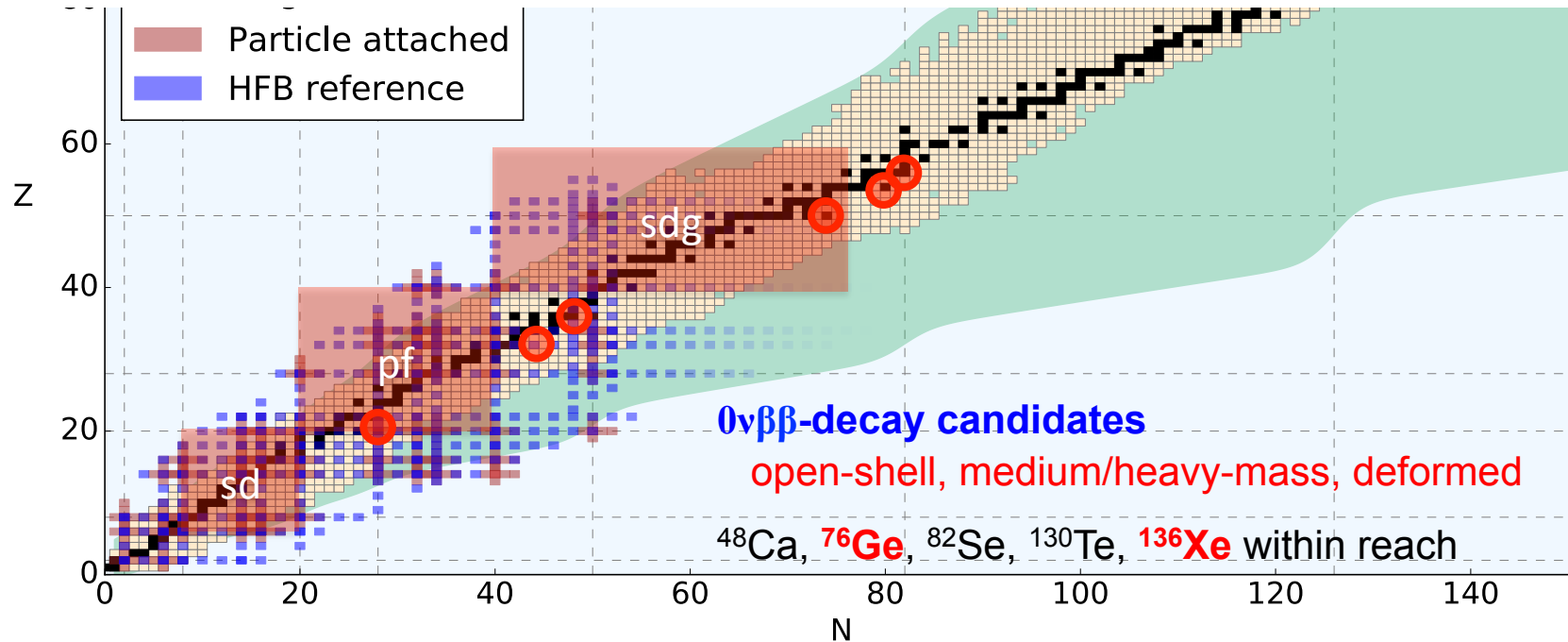


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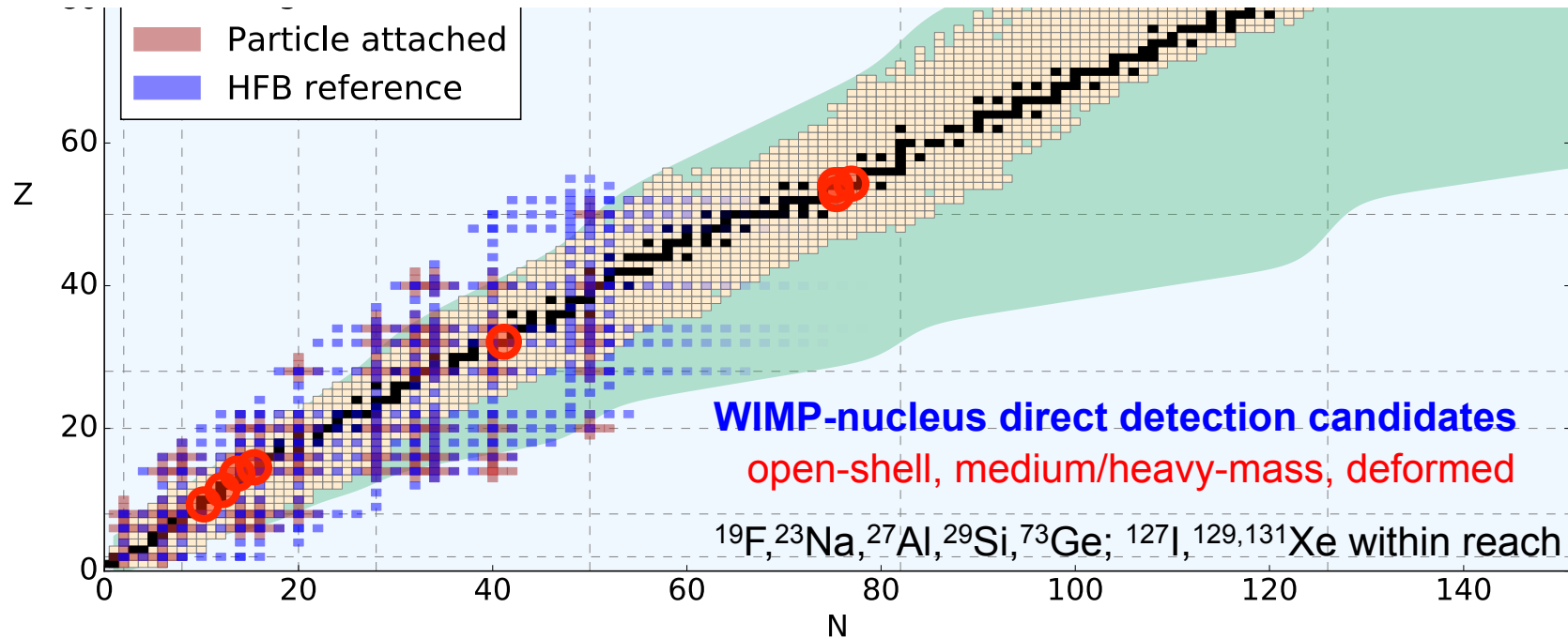
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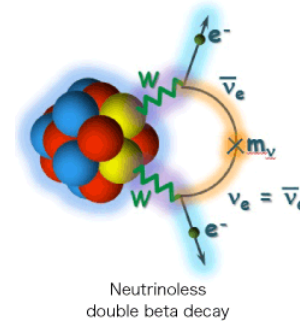


Conventional Shell Model: phenomenological wavefunctions

Ab initio valence-space: **wavefunctions based on NN+3N forces from chiral EFT**

$$M^{0\nu} = M_{GT}^{0\nu} - \frac{M_F^{0\nu}}{g_A^2} + M_T^{0\nu}$$

$$M_{GT}^{0\nu} = \langle f | \sum_{ab} H(r_{ab}) \sigma_a \cdot \sigma_b \tau_a^+ \tau_b^+ | i \rangle$$



- 1) Ab initio energies in medium/heavy-mass region

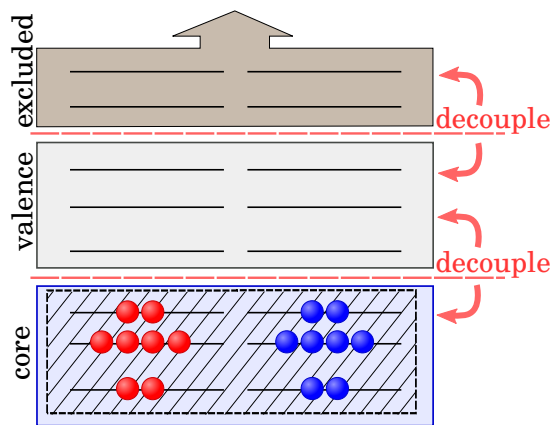
Explicitly construct unitary transformation from sequence of rotations

$$U = e^{\Omega} = e^{\eta_n} \dots e^{\eta_1} \quad \eta = \frac{1}{2} \arctan \left( \frac{2H_{\text{od}}}{\Delta} \right) - \text{h.c.}$$

$$\tilde{H} = e^{\Omega} H e^{-\Omega} = H + [\Omega, H] + \frac{1}{2} [\Omega, [\Omega, H]] + \dots$$

**All operators truncated at two-body level**

Tsukiyama, Bogner, Schwenk, PRC 2012  
Morris, Parzuchowski, Bogner, PRC 2015



$$|\Phi_0\rangle = |^{16}\text{O}\rangle$$

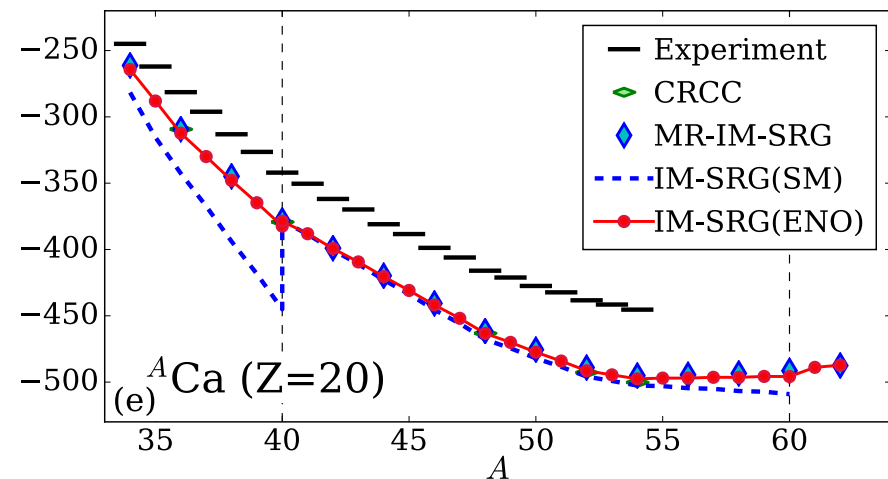
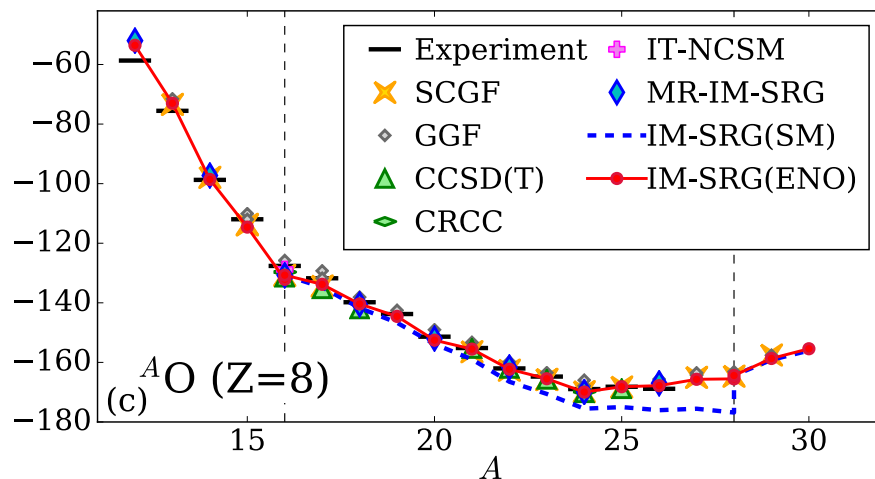
**Step 1: Decouple core**  
**Step 2: Decouple valence space**

Can we achieve accuracy  
of large-space methods?

$$\langle \tilde{\Psi}_n | P \tilde{H} P | \tilde{\Psi}_n \rangle \approx \langle \Psi_i | H | \Psi_i \rangle$$

$\langle P   H   P \rangle$	$\langle P   H   Q \rangle \rightarrow 0$
$\langle Q   H   P \rangle \rightarrow 0$	$\langle Q   H   Q \rangle$

**ENO agrees to 1% with large-space methods** (where calculations exist)



Stroberg et al., PRL (2017)

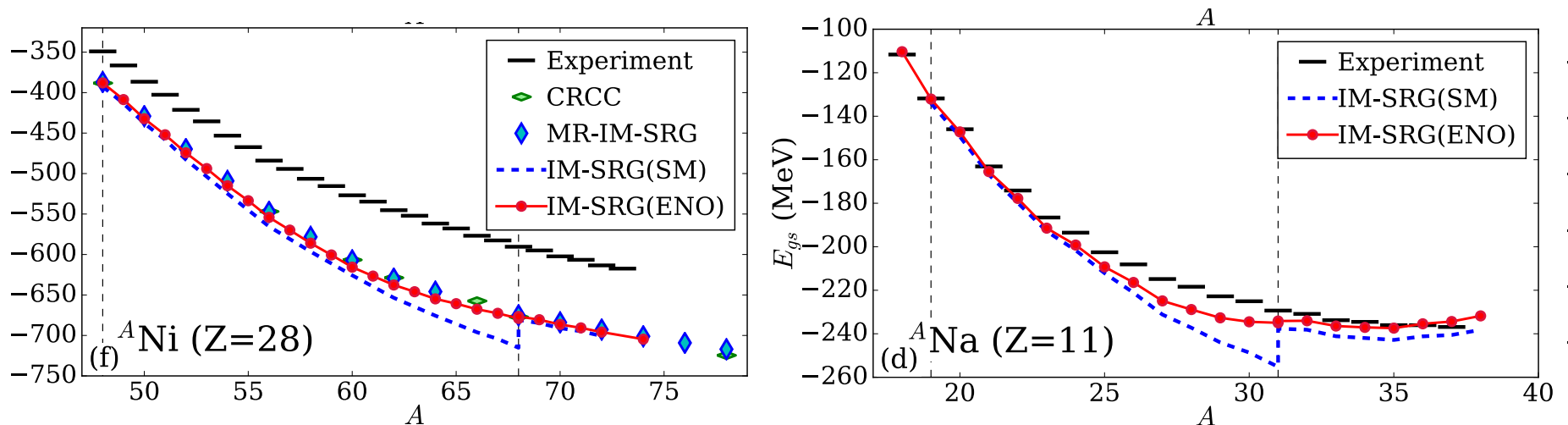
Extend beyond standard *sd*/*pf* shells

Agreement with experiment deteriorates for heavy chains (due to input Hamiltonian)

Significant gain in applicability with little/no sacrifice in accuracy

**Low computational cost:** ~1 node-day/nucleus

**Targeted valence space agrees to 1% with all large-space methods** (where calculations exist)



Stroberg et al., PRL (2017)

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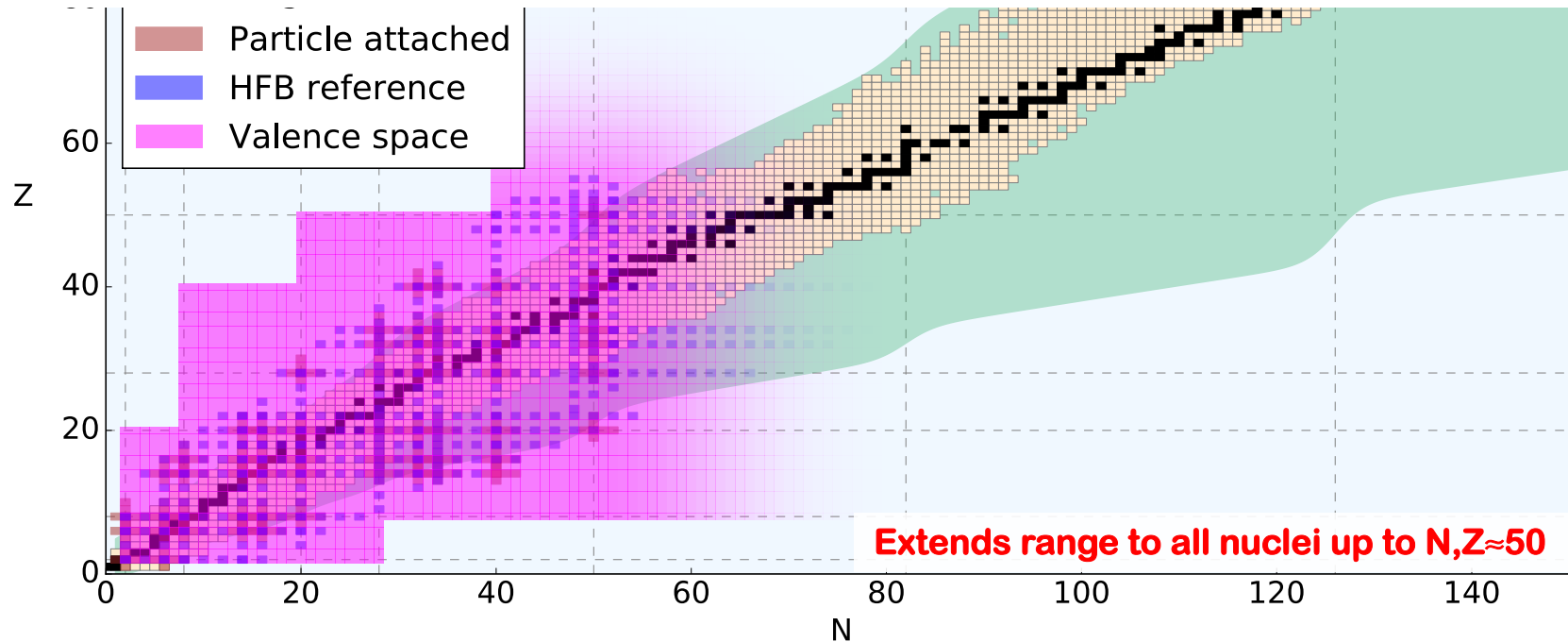
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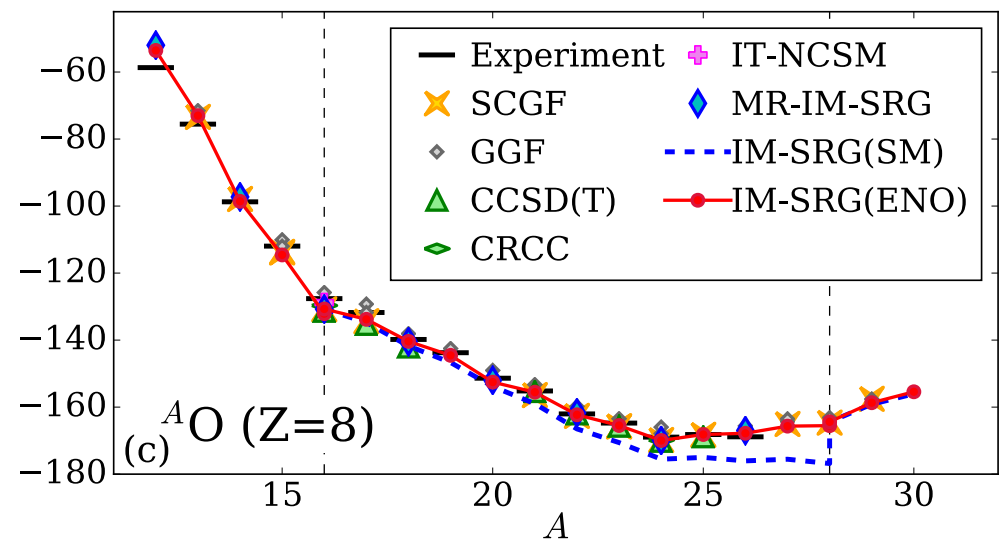
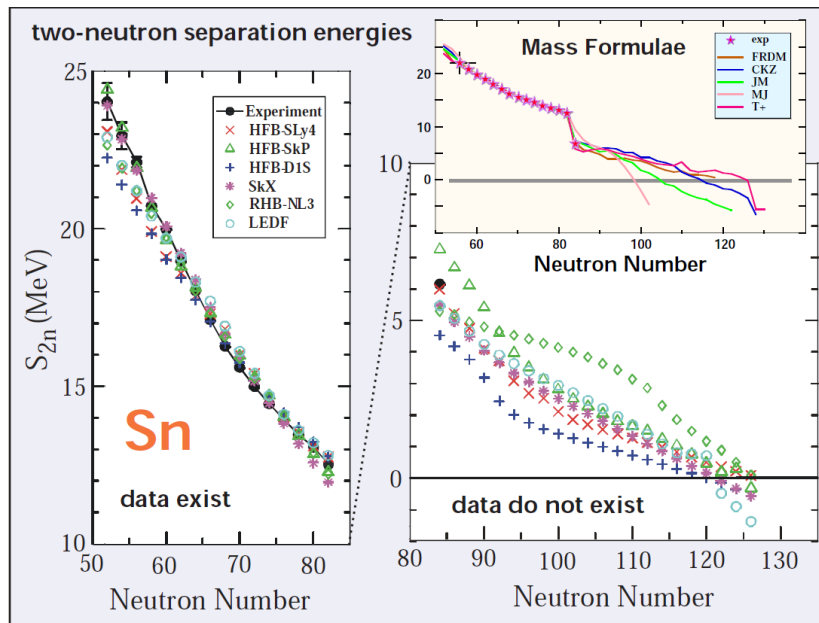
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- **Nuclear many-body problem**



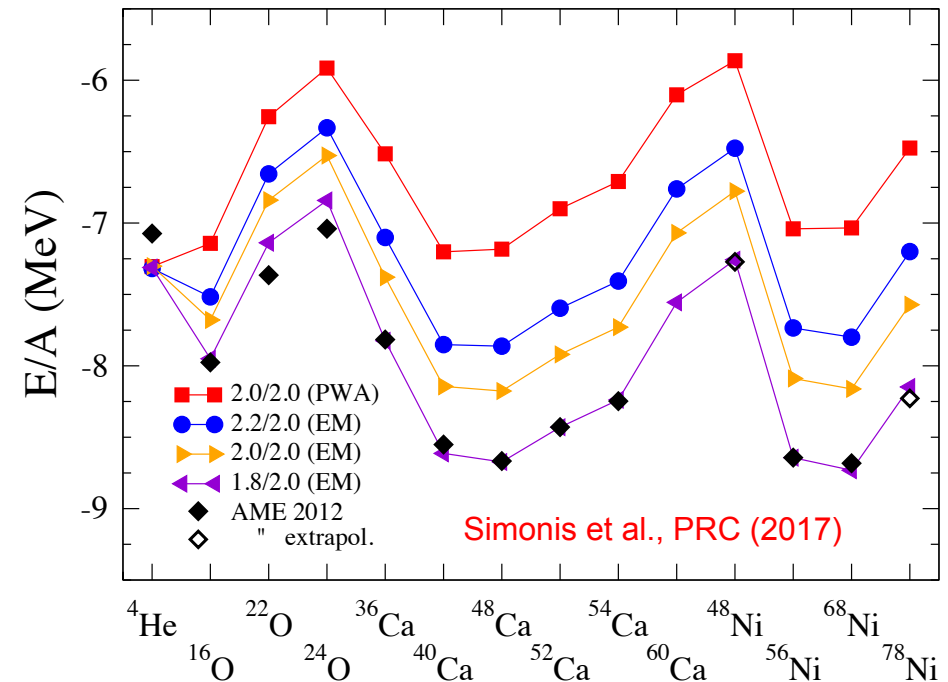
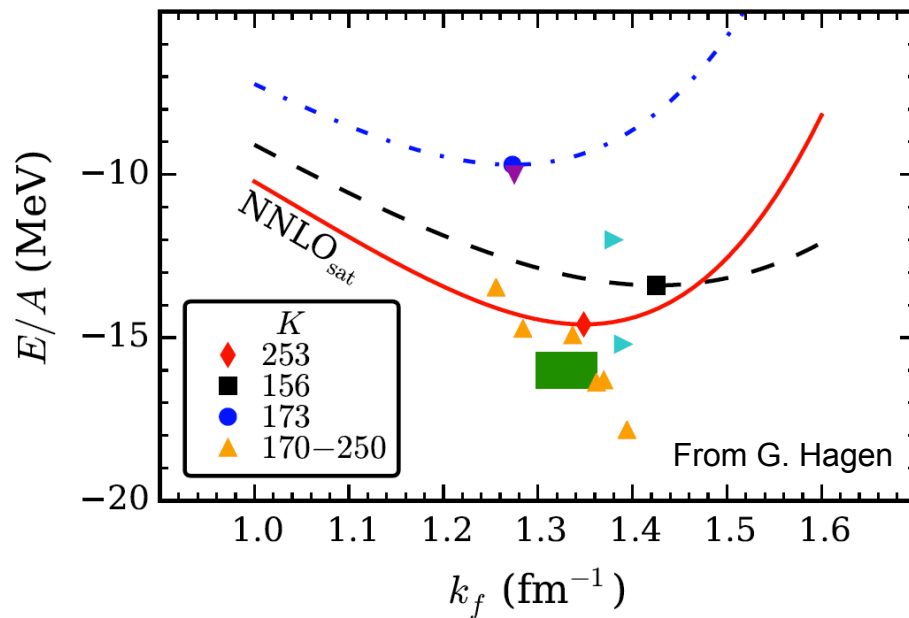
## Stark contrast in extrapolations between model extrapolations and ab initio



All ab initio methods in good agreement when starting from same input NN+3N forces

Only informed by 2,3-body data

## Hebeler/Simonis NN+3N forces with reasonable saturation properties

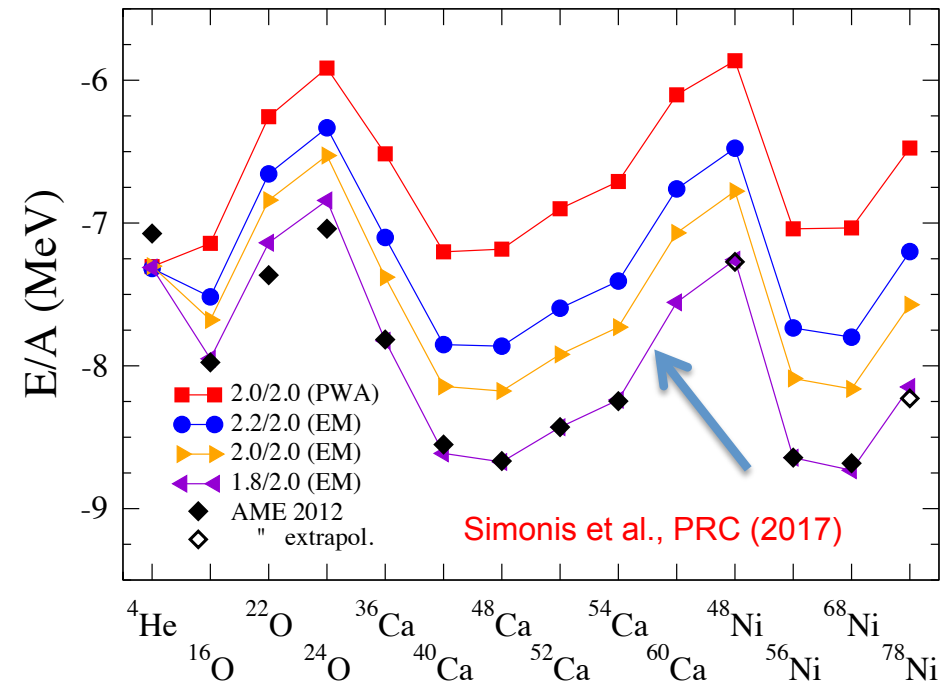
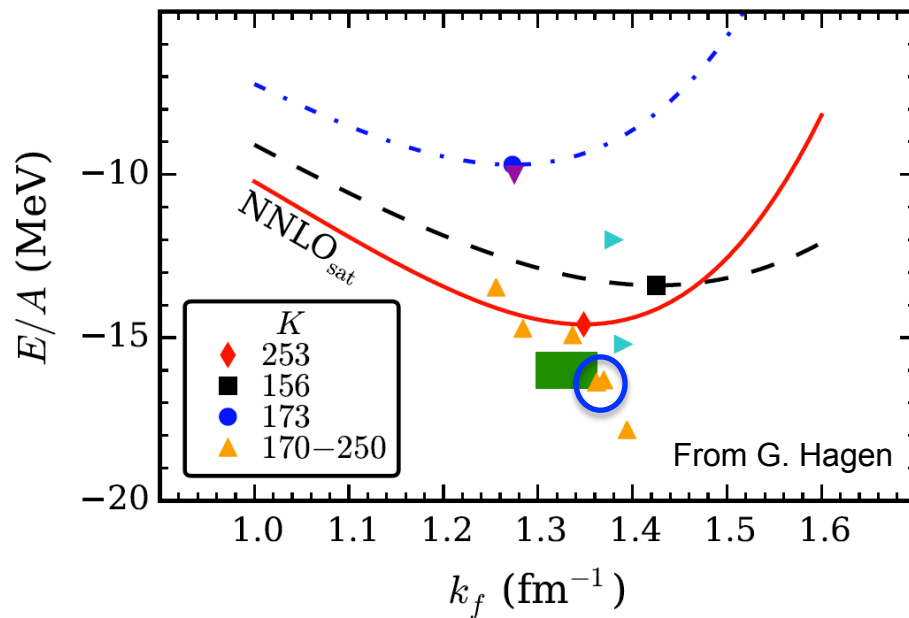


1.8/2.0 (EM) reproduces closed shells through  ${}^{78}\text{Ni}$

Only underbound for neutron-rich oxygen



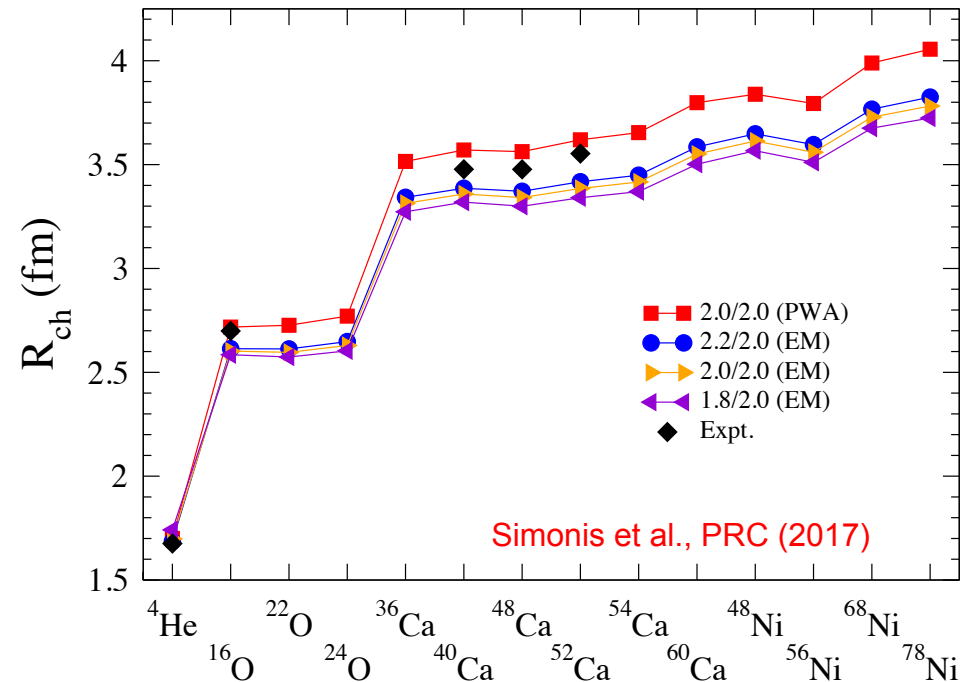
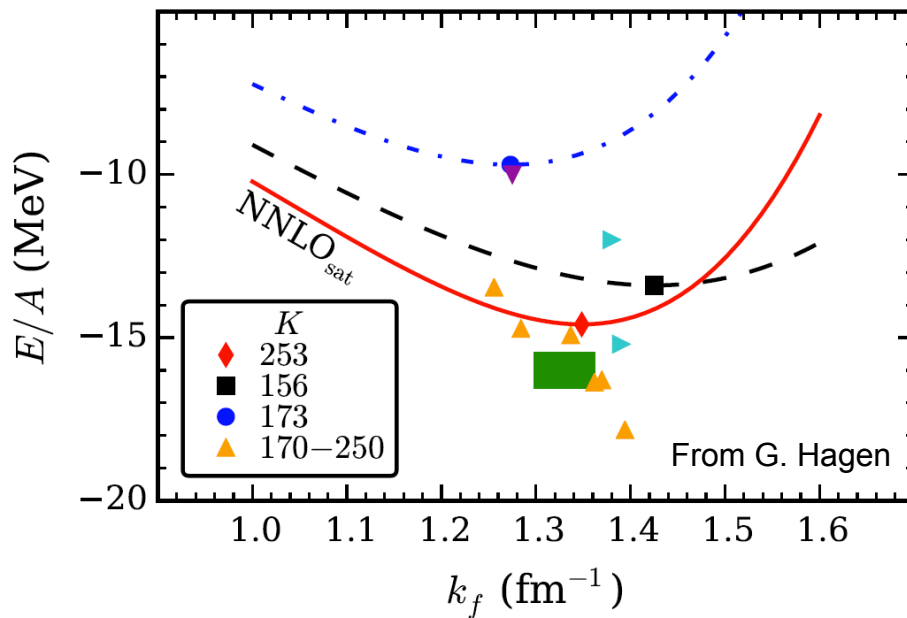
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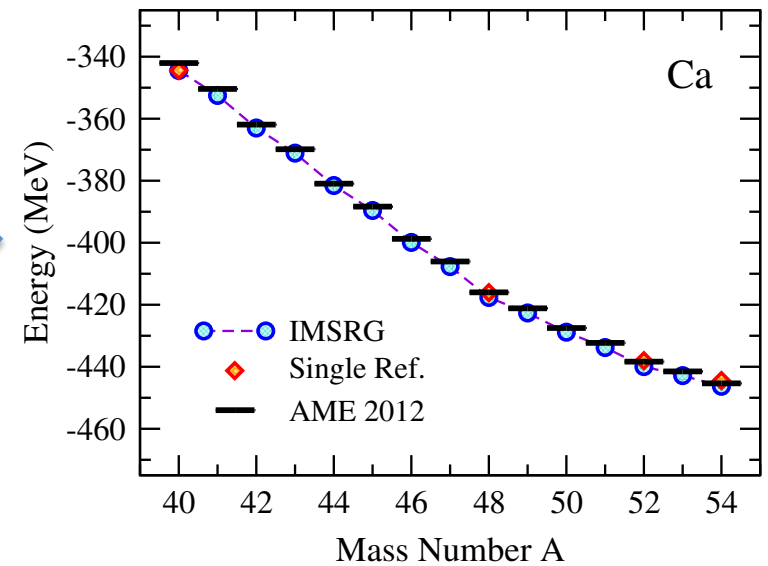
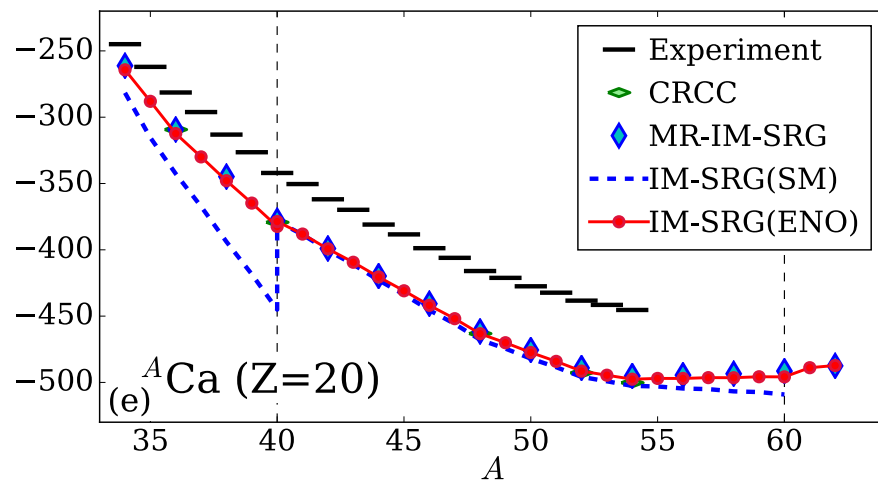
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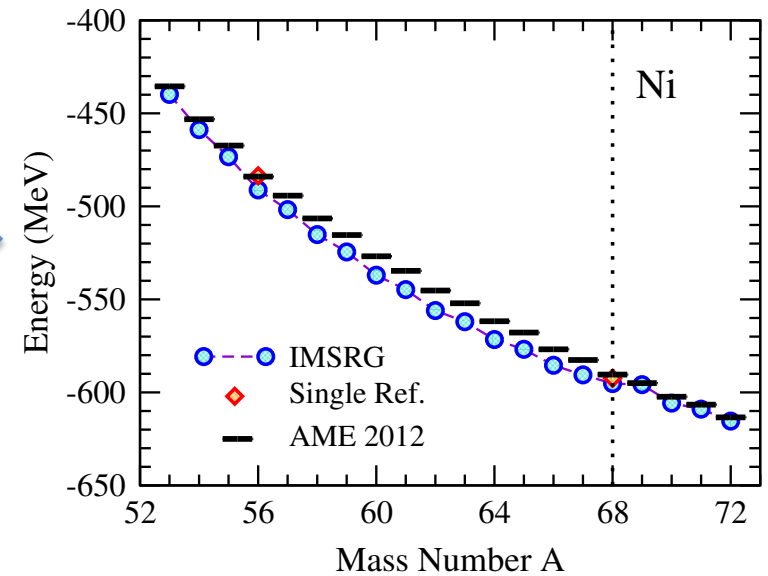
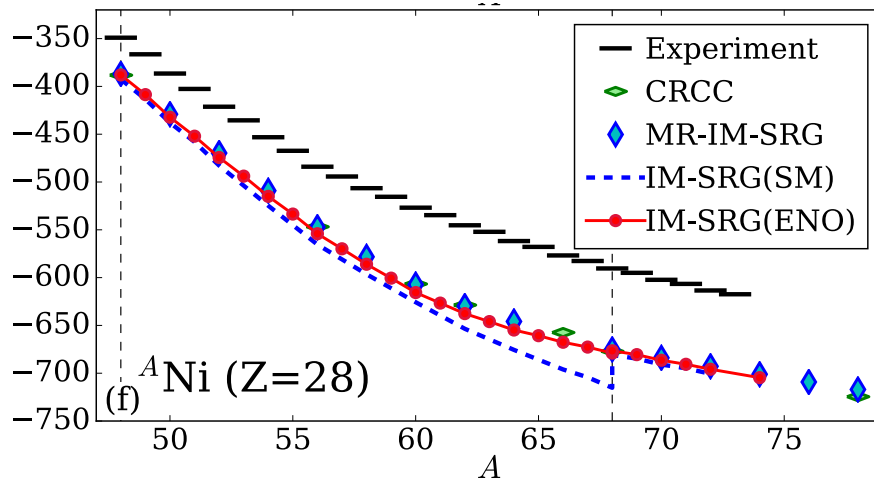
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Isotopic chains: dramatic improvement with respect to experimental data

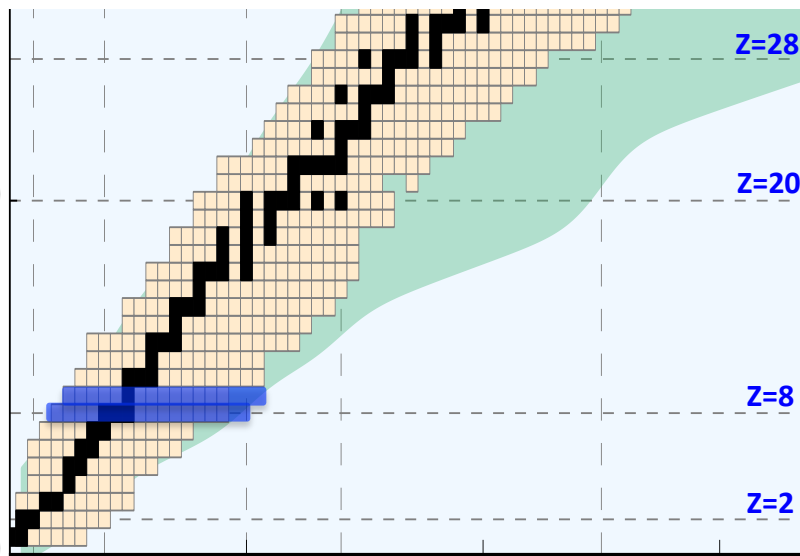


Isotopic chains: dramatic improvement with respect to experimental data



How well does it work across broad regions of nuclei?

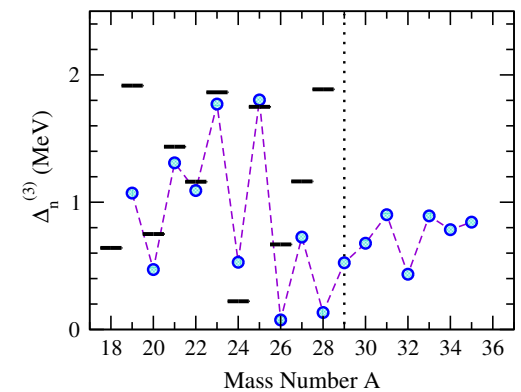
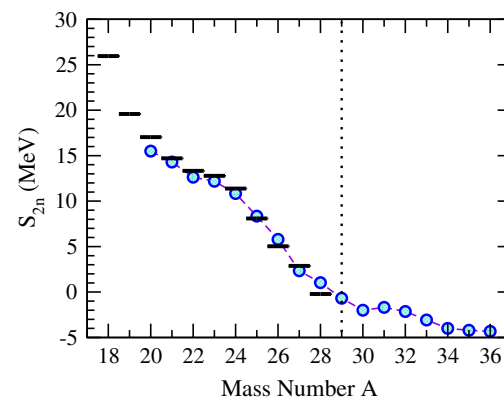
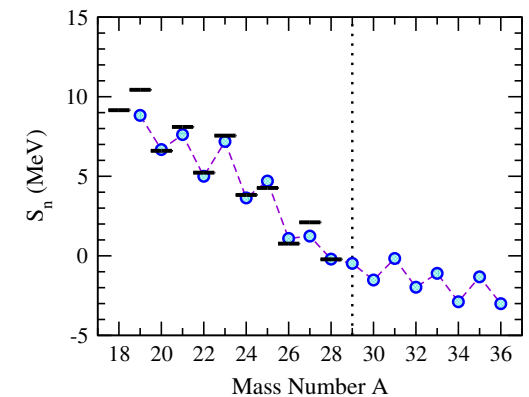
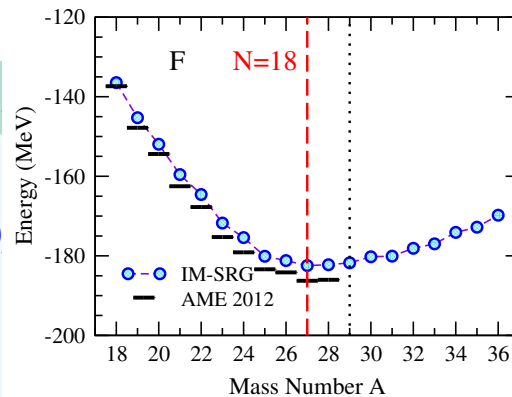
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

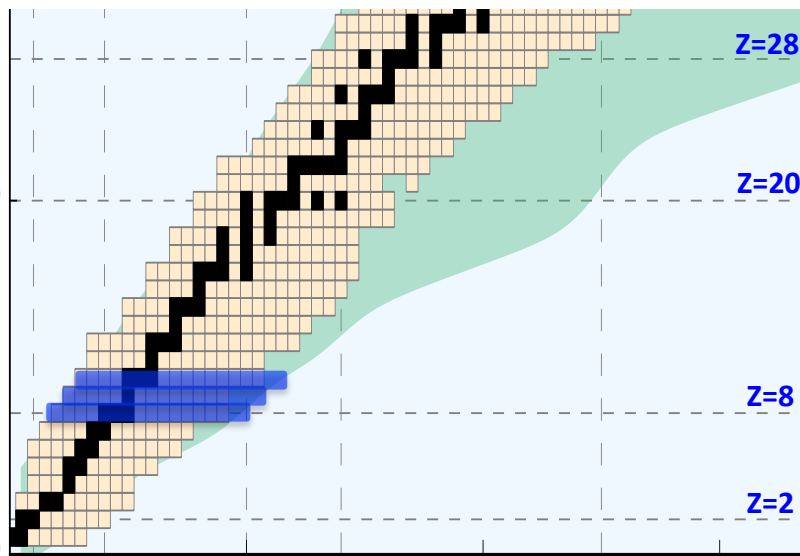
Probe dripline and beyond

Artifacts at neutron N=20 gap



JDH, Stroberg, et al., in preparation

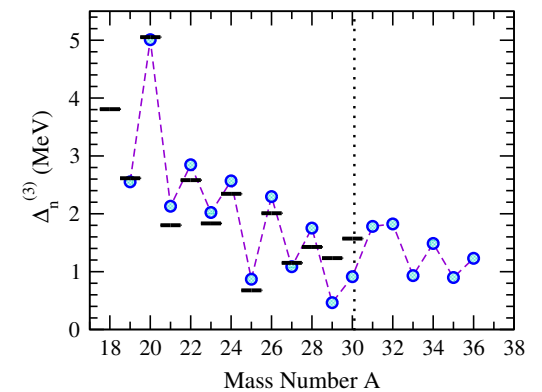
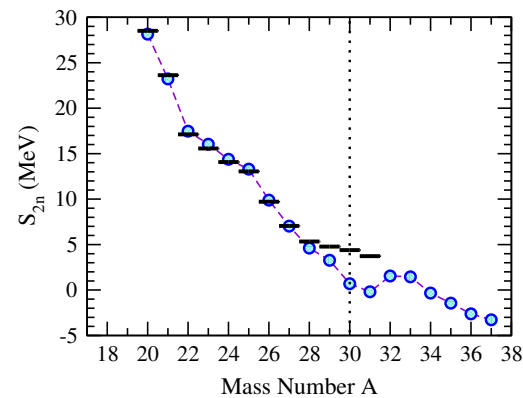
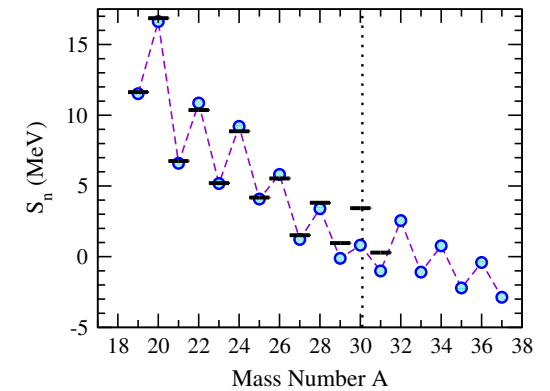
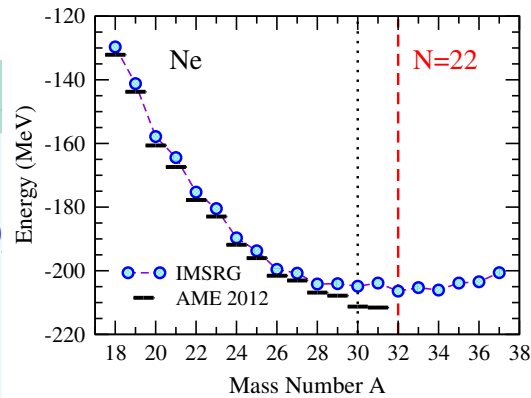
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

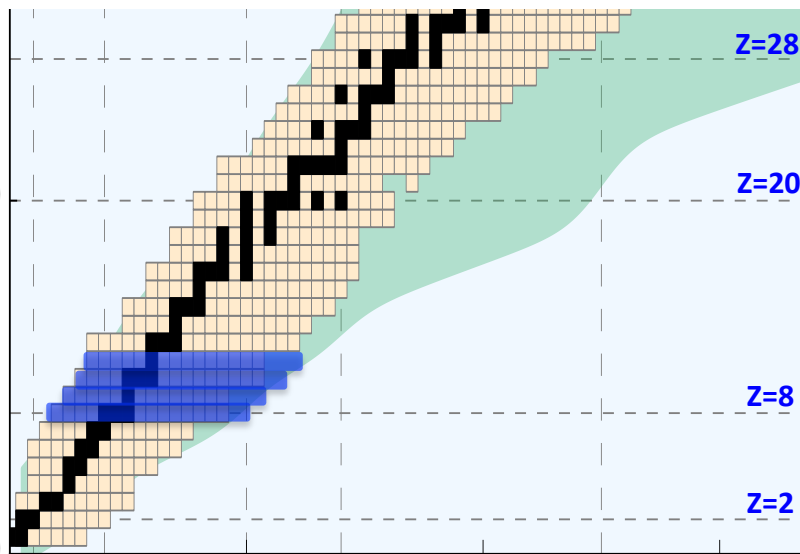
Probe dripline and beyond

Artifacts at neutron N=20 gap



JDH, Stroberg, et al., in preparation

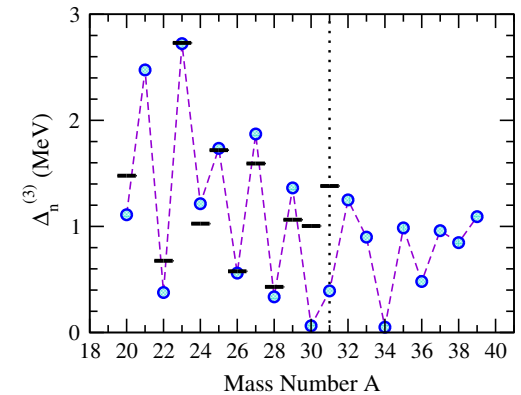
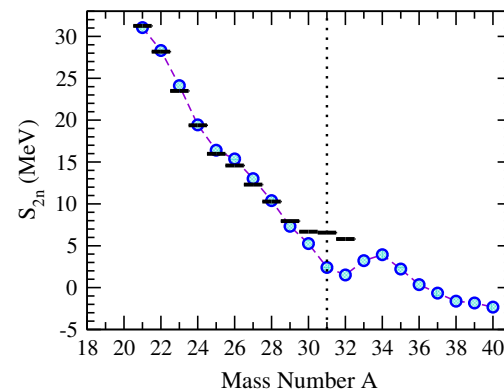
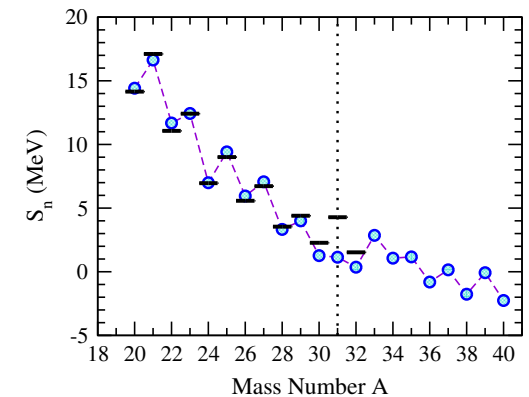
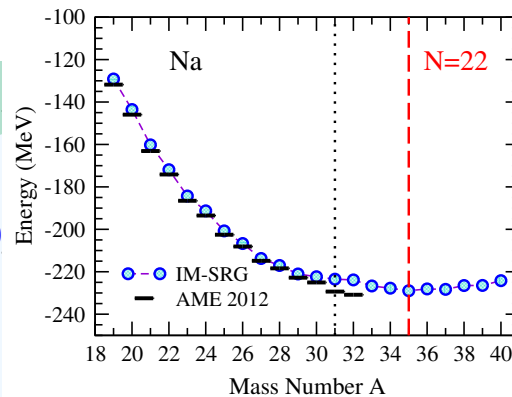
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

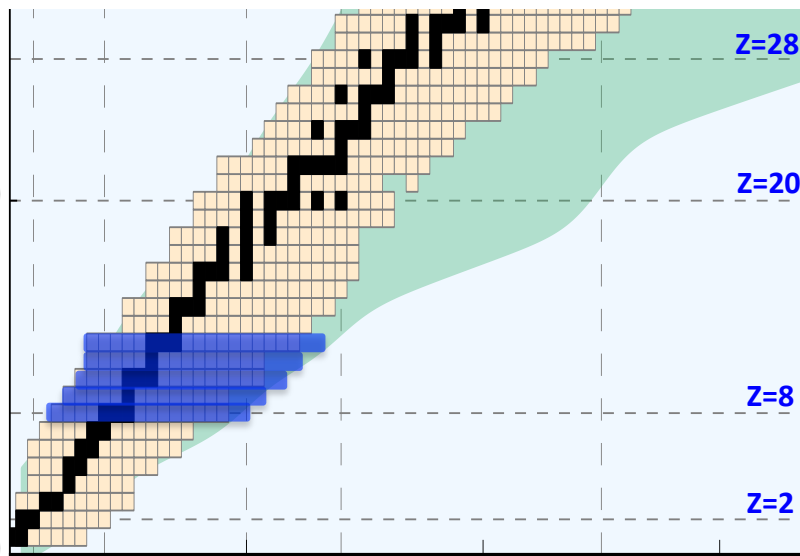
Probe dripline and beyond

Artifacts at neutron N=20 gap



JDH, Stroberg, et al., in preparation

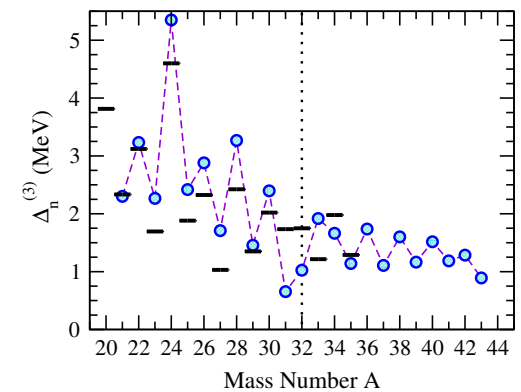
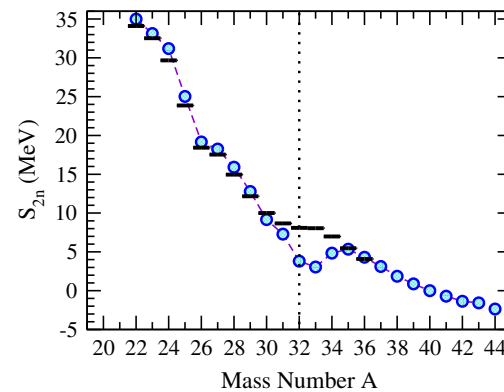
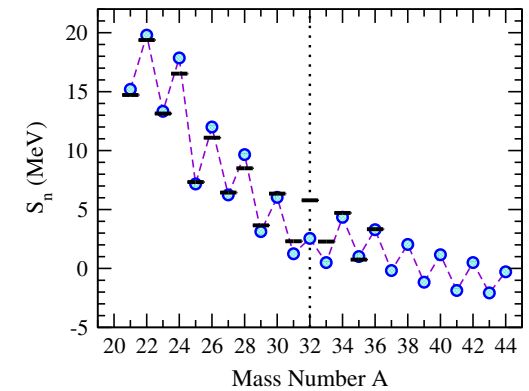
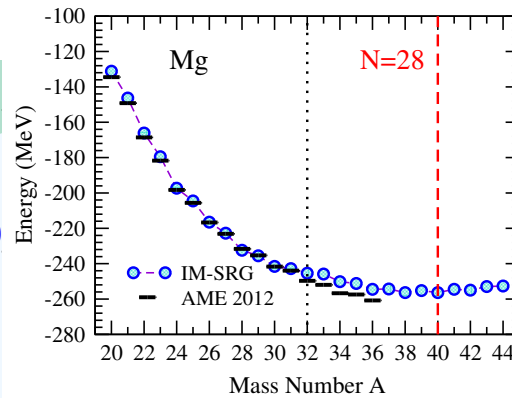
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

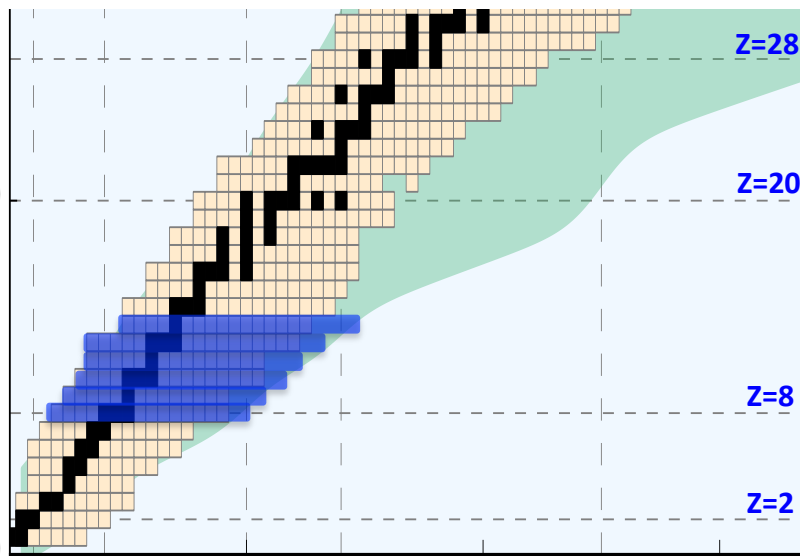
Probe dripline and beyond

Artifacts at neutron N=20 gap





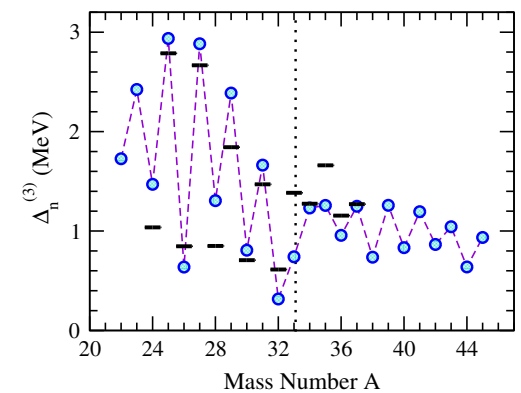
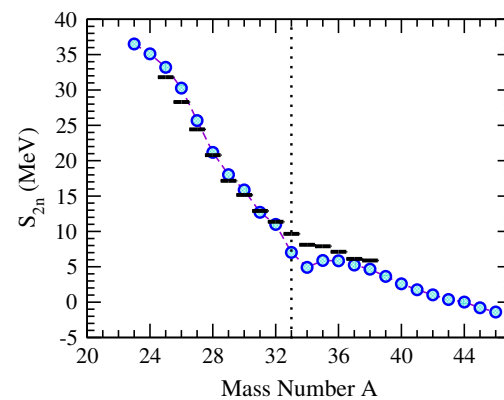
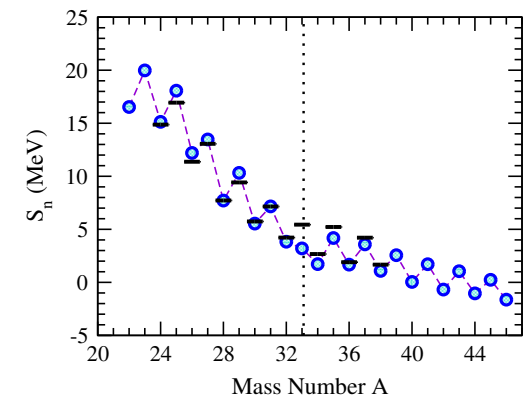
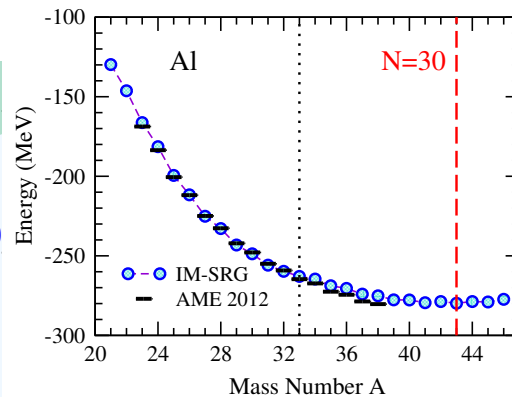
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

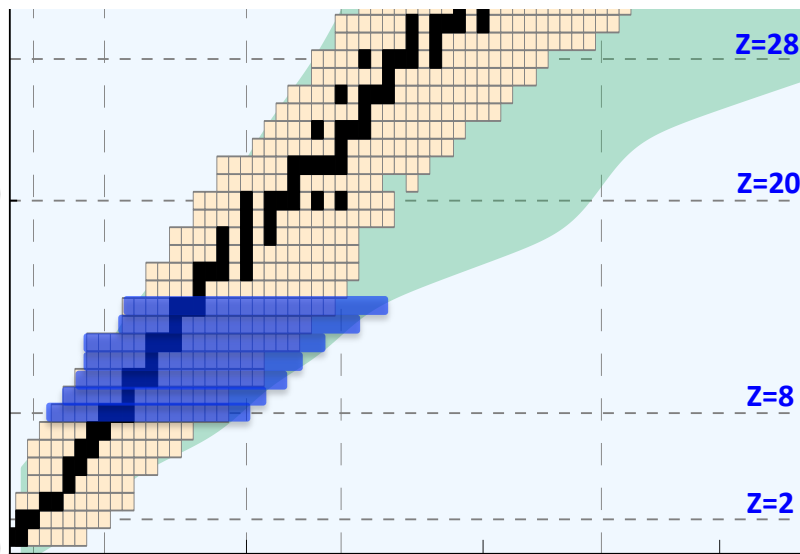
Probe dripline and beyond

Artifacts at neutron N=20 gap



JDH, Stroberg, et al., in preparation

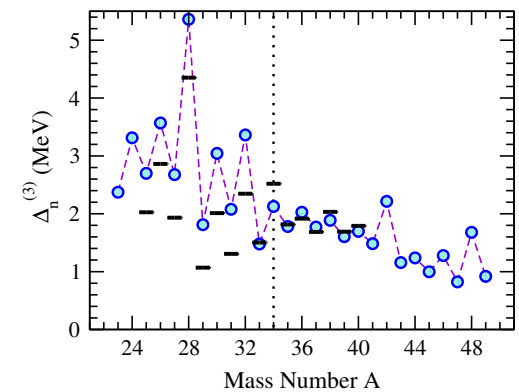
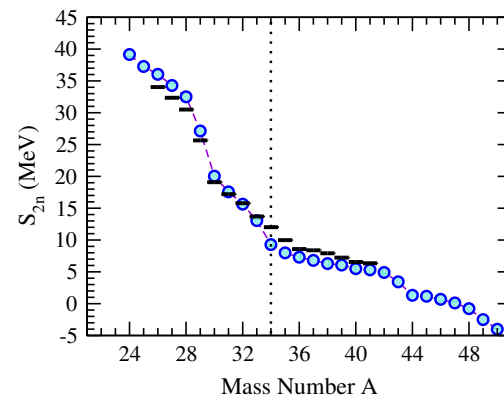
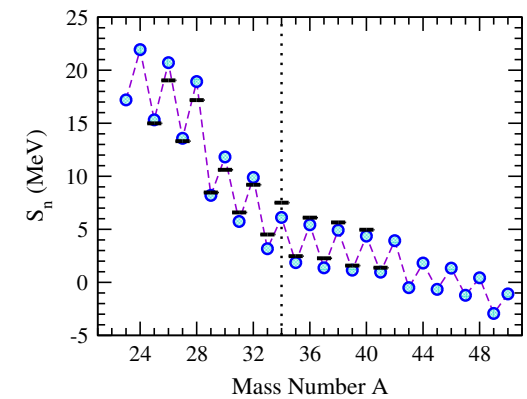
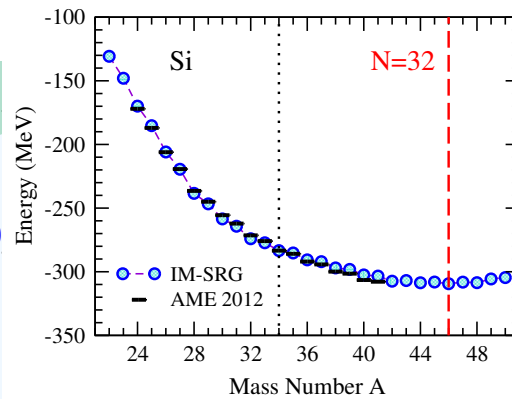
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

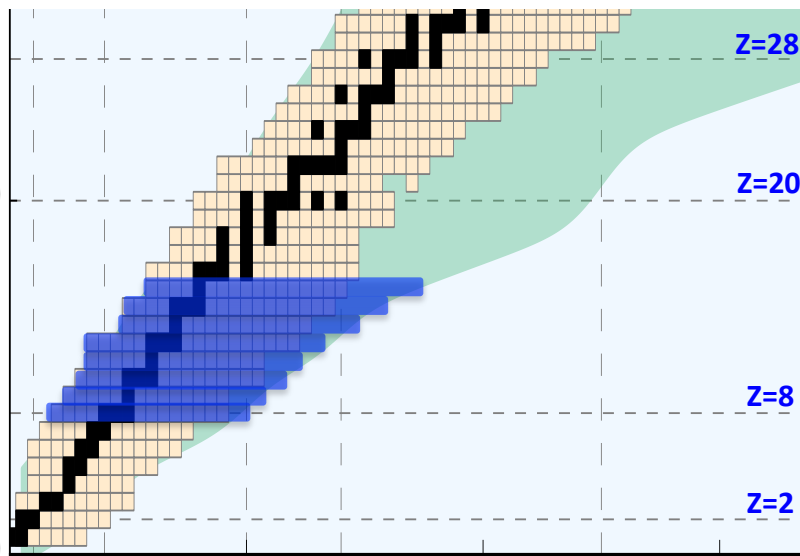
Probe dripline and beyond

Artifacts at neutron N=20 gap



JDH, Stroberg, et al., in preparation

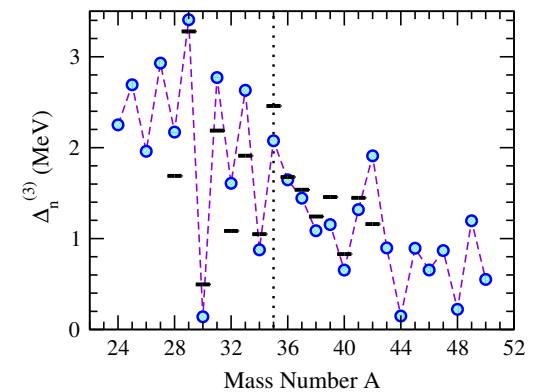
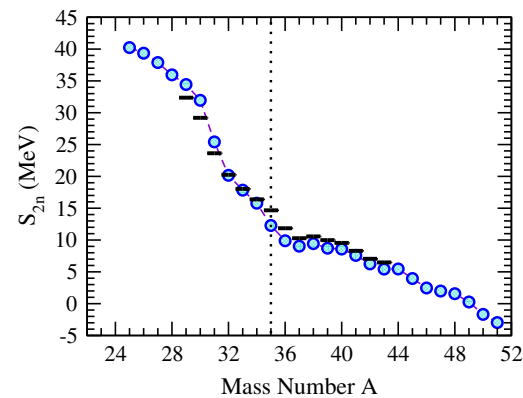
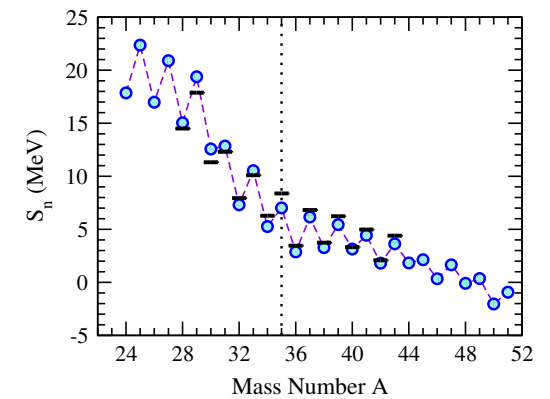
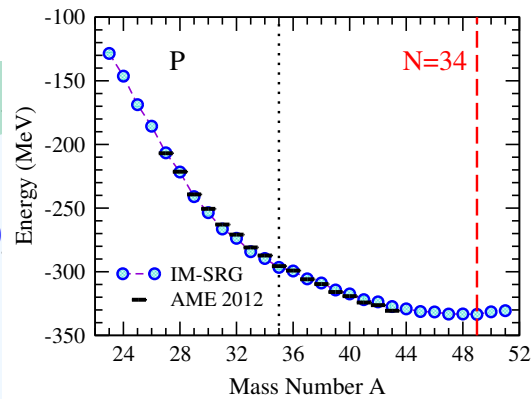
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

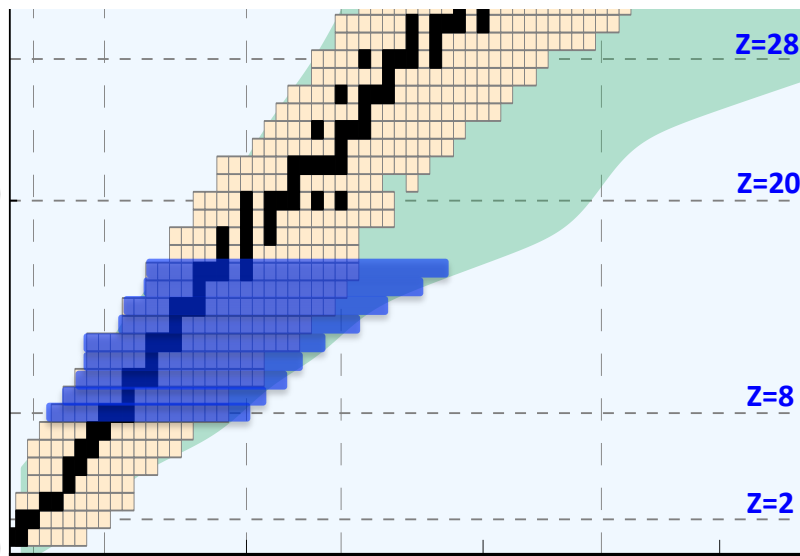
Probe dripline and beyond

Artifacts at neutron N=20 gap



JDH, Stroberg, et al., in preparation

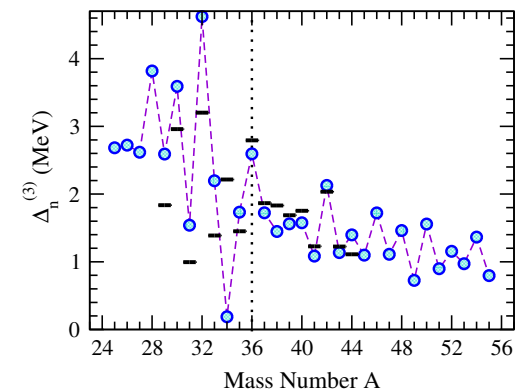
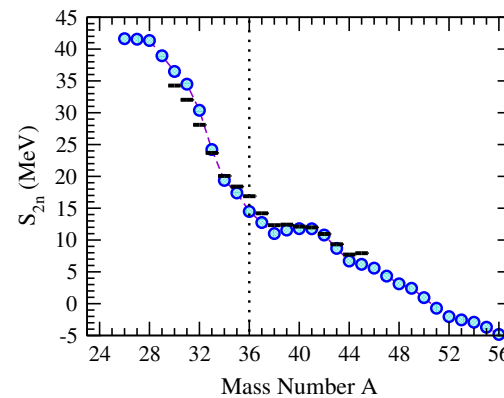
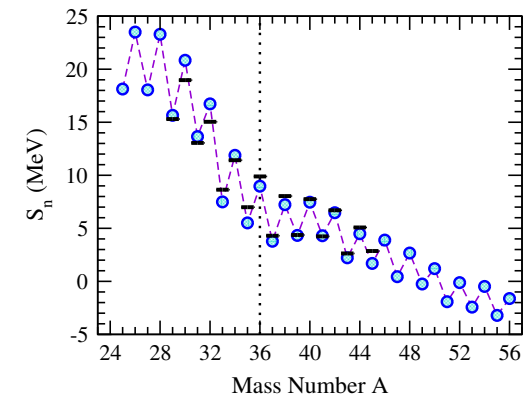
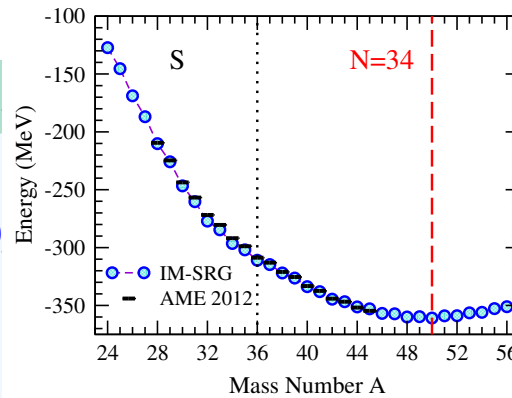
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

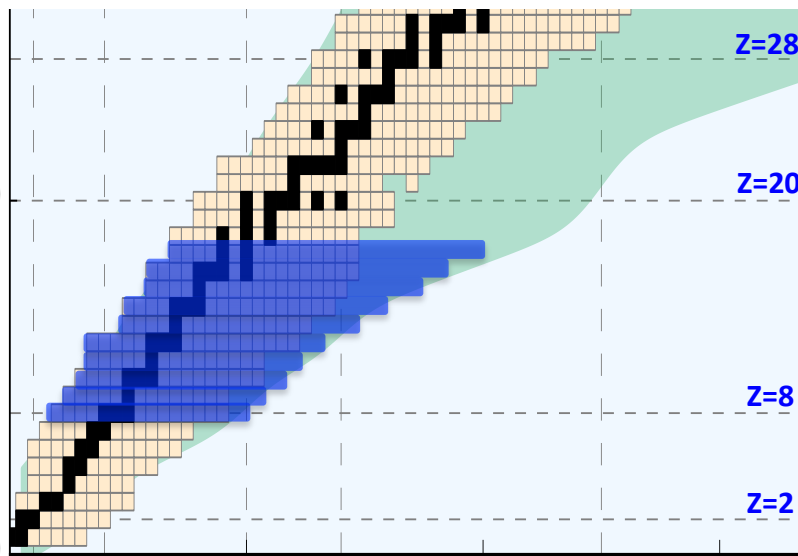
Probe dripline and beyond

Artifacts at neutron N=20 gap



JDH, Stroberg, et al., in preparation

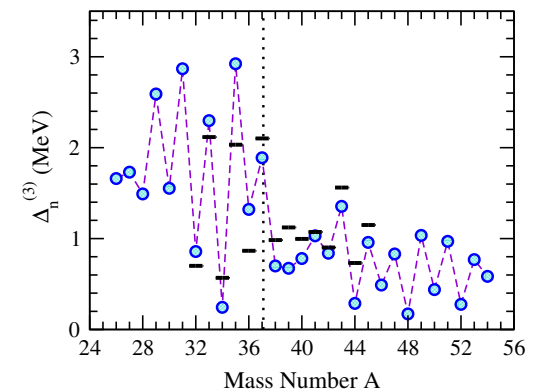
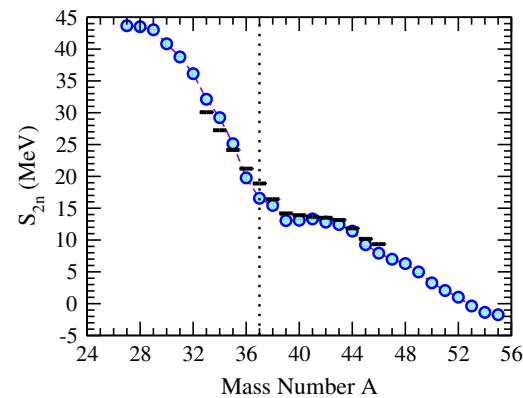
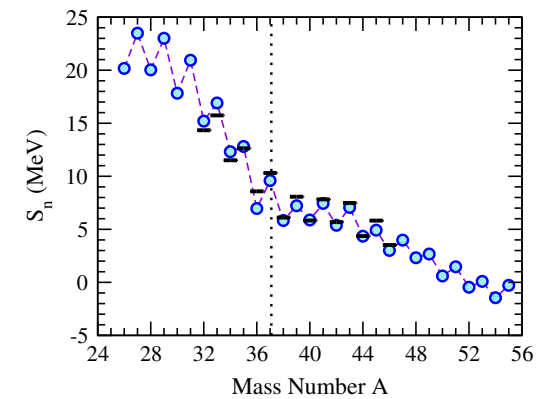
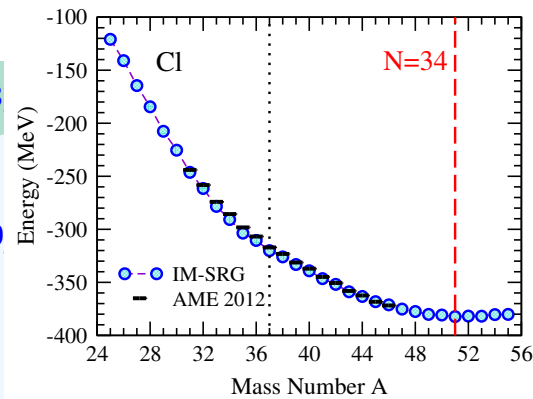
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

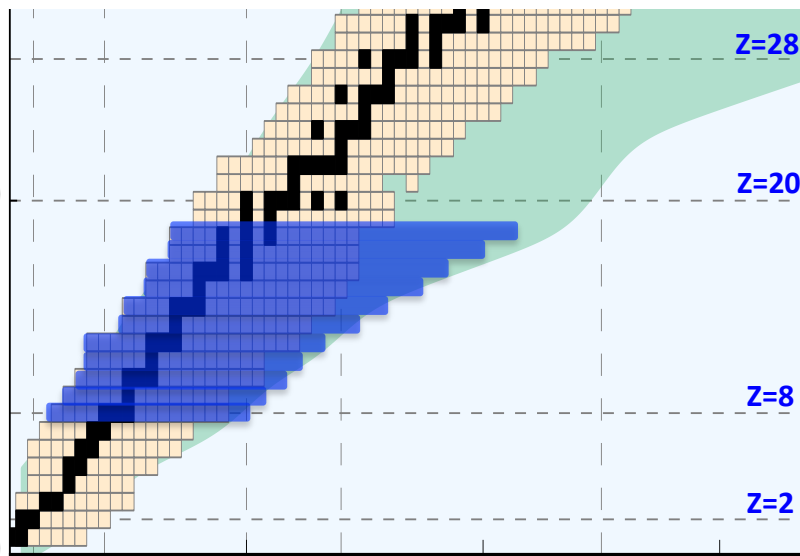
Probe dripline and beyond

Artifacts at neutron N=20 gap



JDH, Stroberg, et al., in preparation

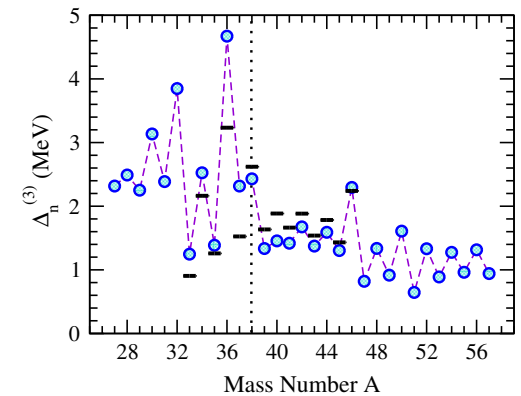
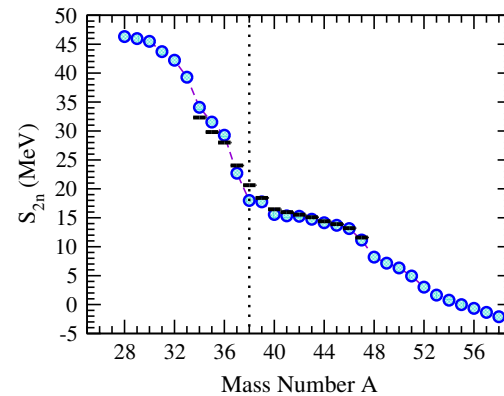
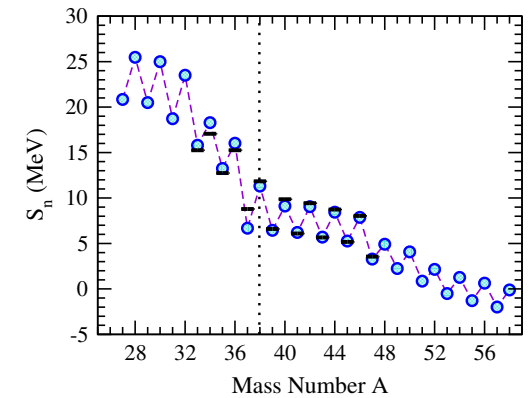
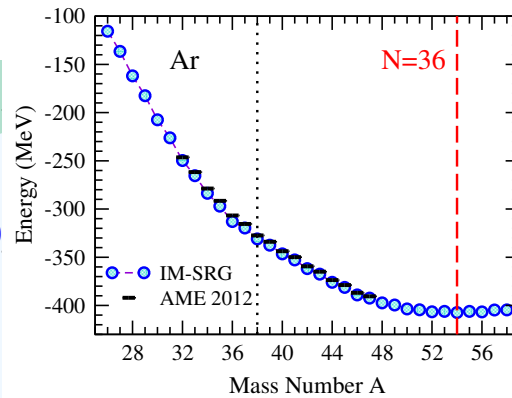
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

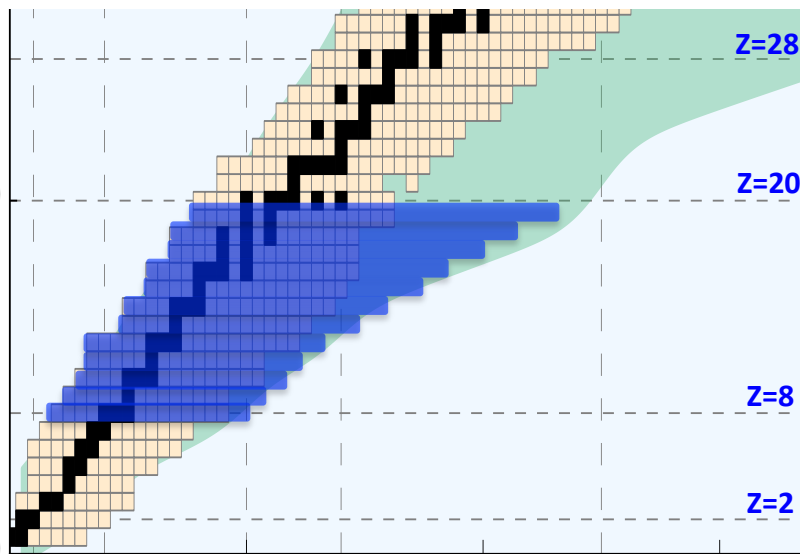
Probe dripline and beyond

Artifacts at neutron N=20 gap



JDH, Stroberg, et al., in preparation

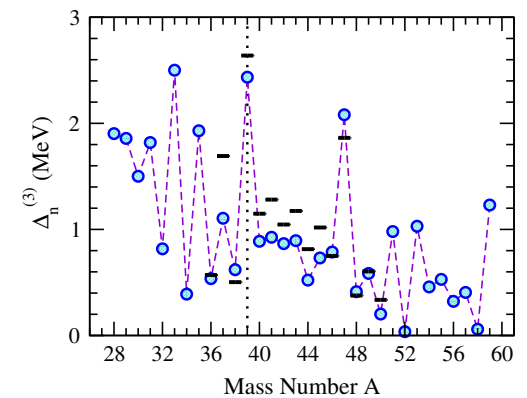
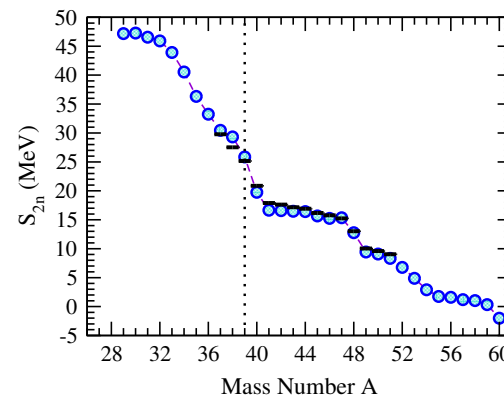
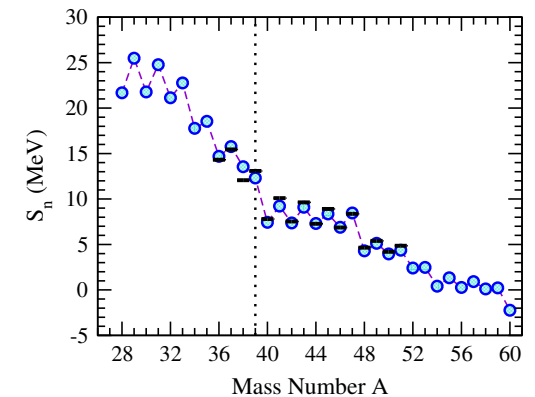
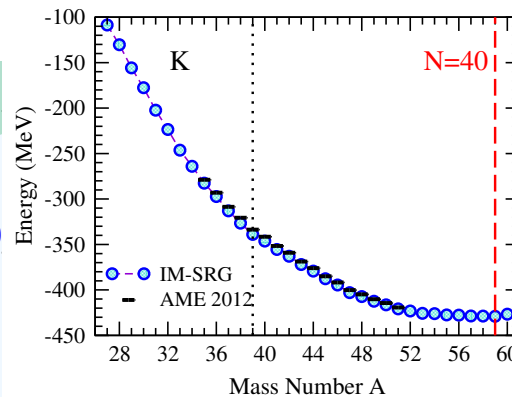
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

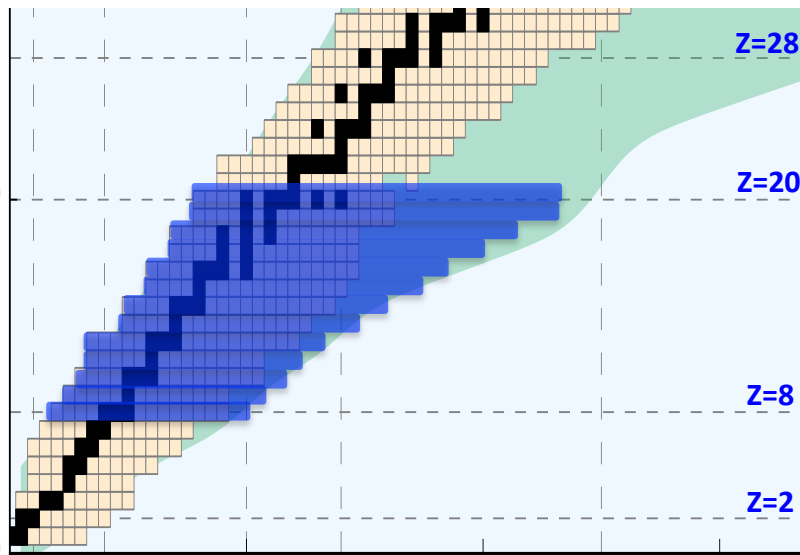
Probe dripline and beyond

Artifacts at neutron N=20 gap



JDH, Stroberg, et al., in preparation

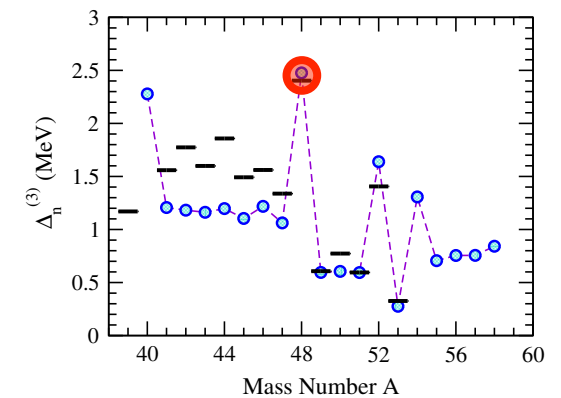
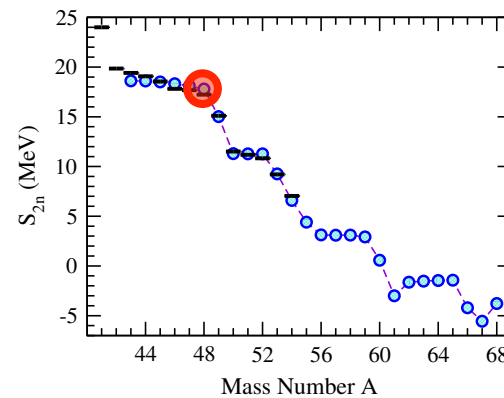
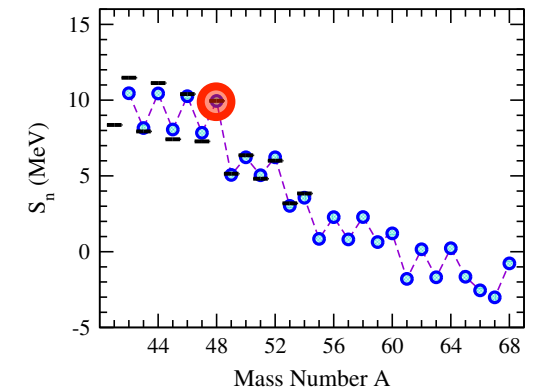
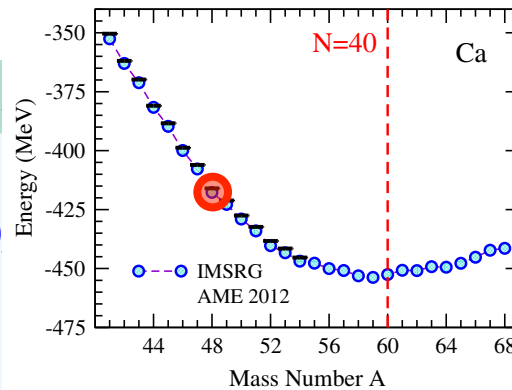
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

Probe dripline and beyond

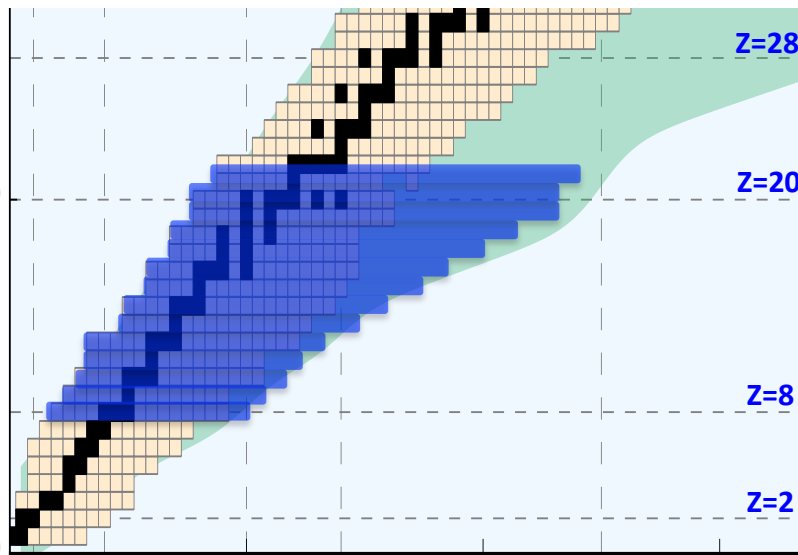
Artifacts at neutron N=40 gap



JDH, Stroberg, et al., in preparation



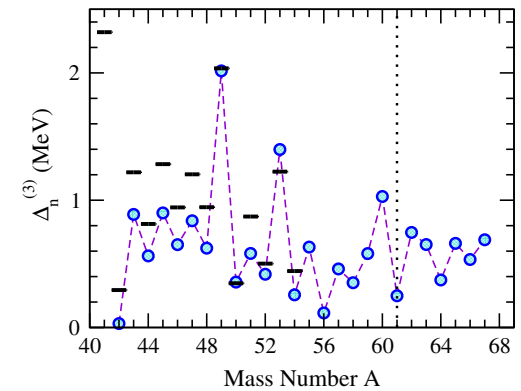
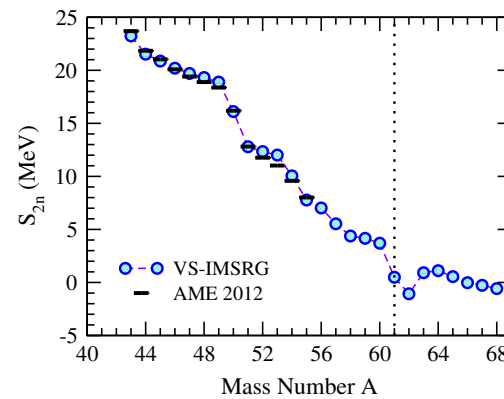
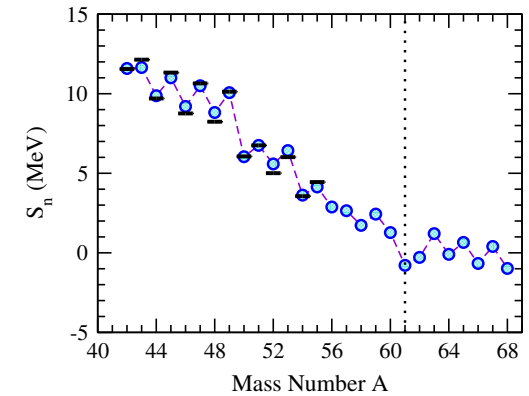
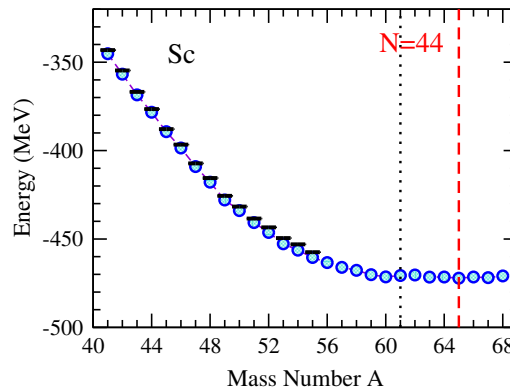
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

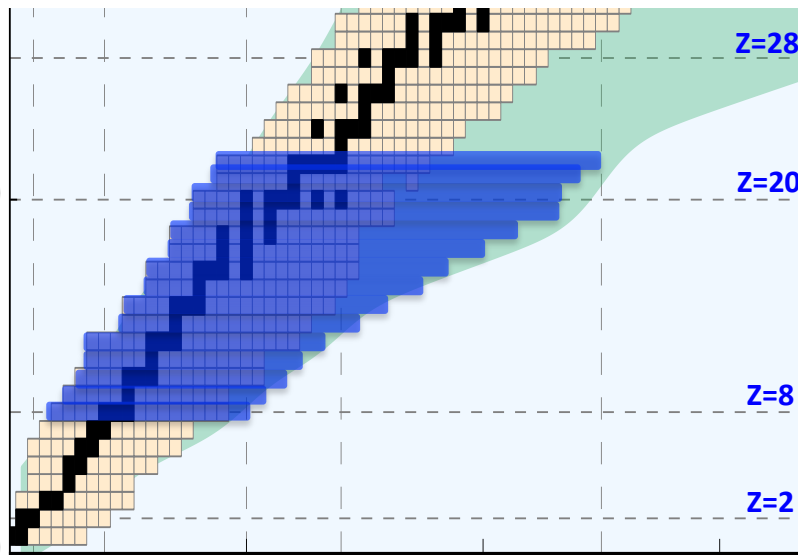
Probe dripline and beyond

Artifacts at neutron N=40 gap



JDH, Stroberg, et al., in preparation

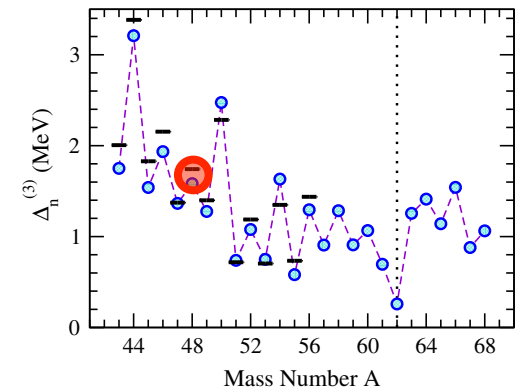
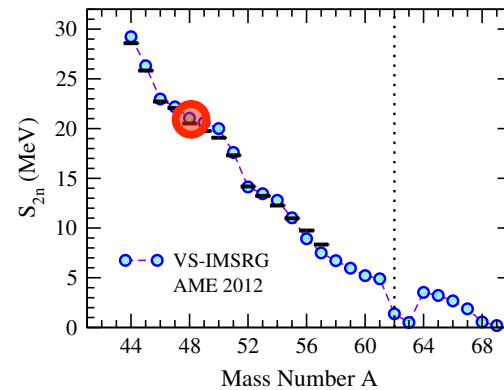
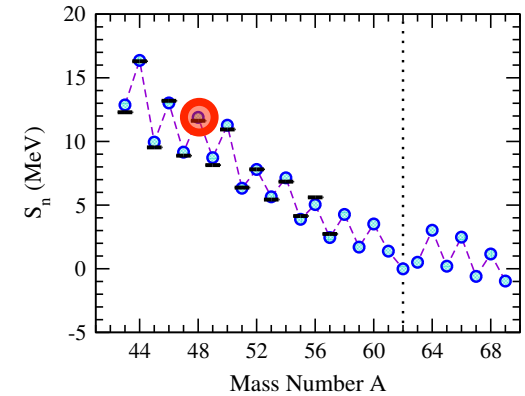
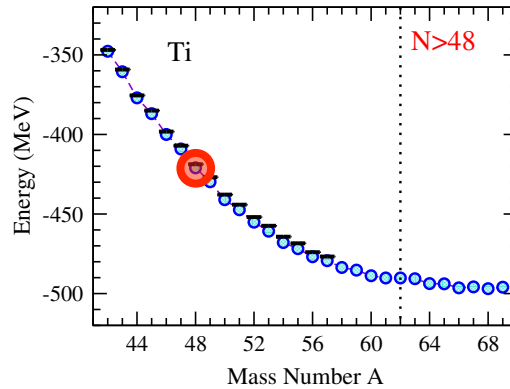
## Explore ground-state properties throughout medium-mass region

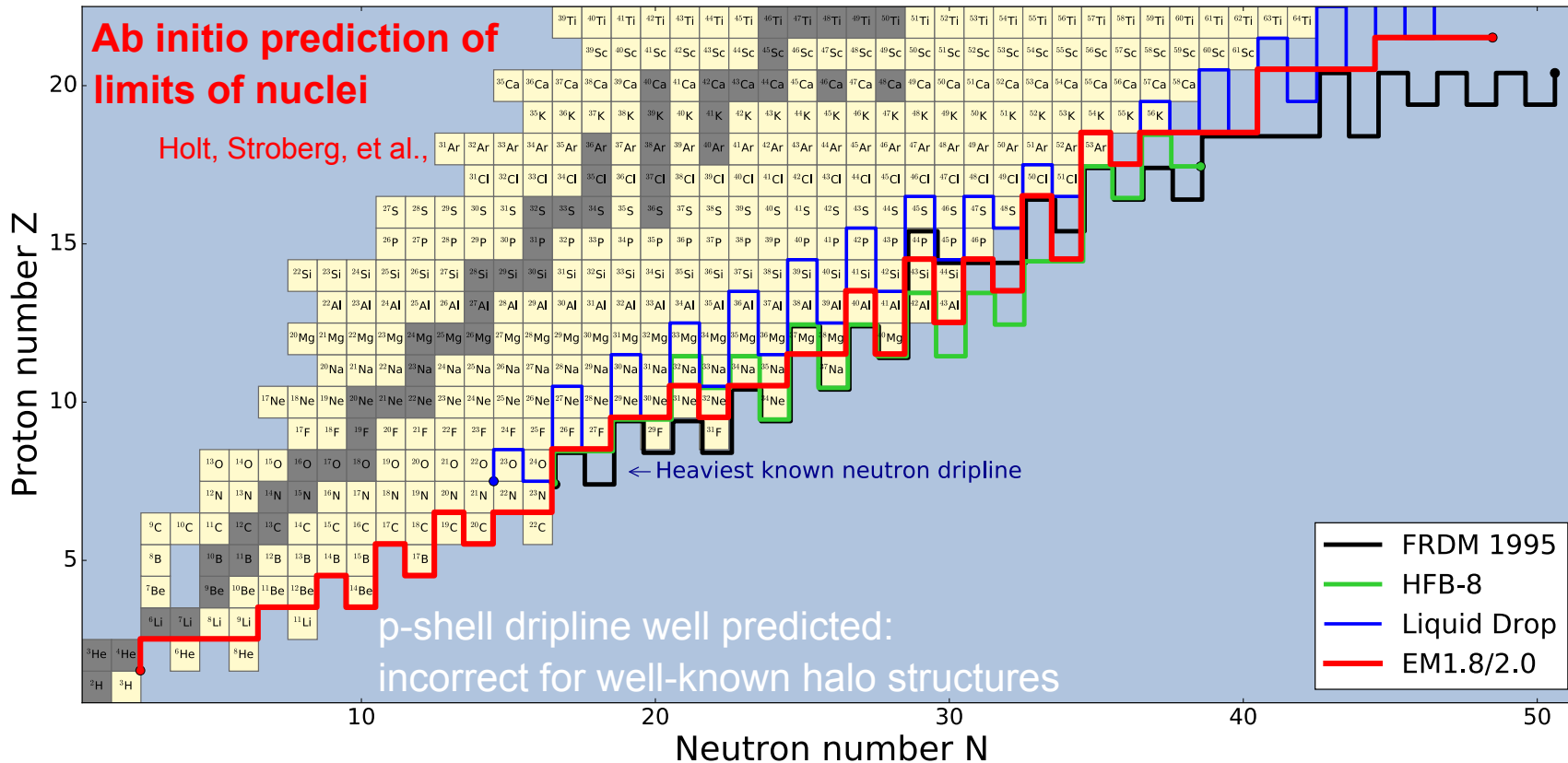


Remarkable agreement with experiment

Probe dripline and beyond

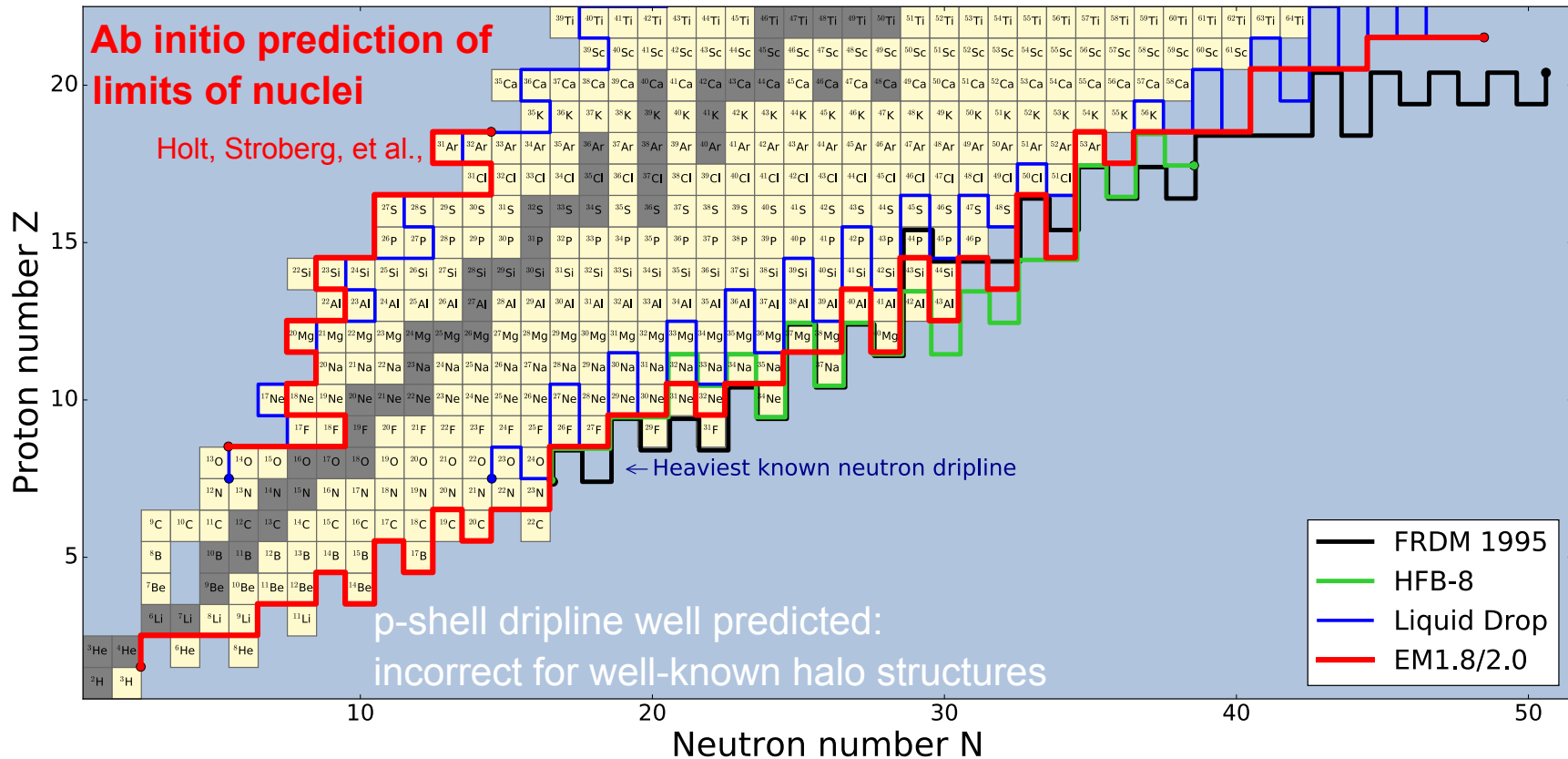
Artifacts at neutron N=40 gap





General agreement with model predictions

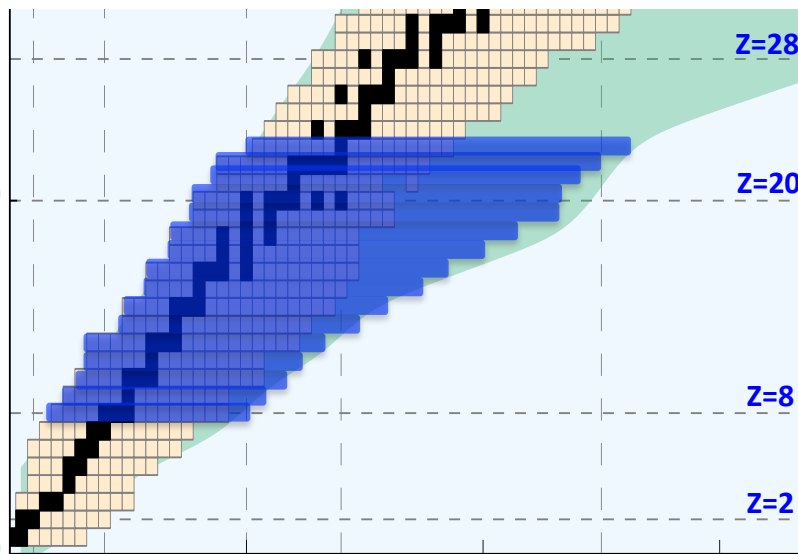
Significant differences arise for heavy nuclei



General agreement with model predictions

Proton dripline: very good agreement with experiment

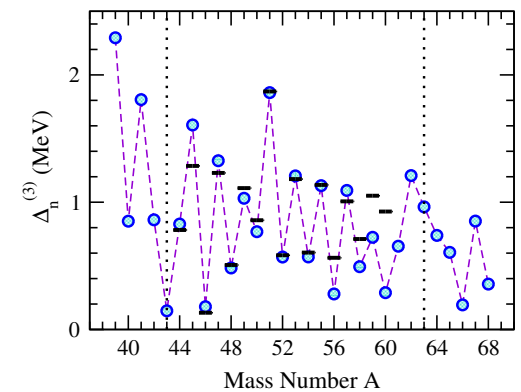
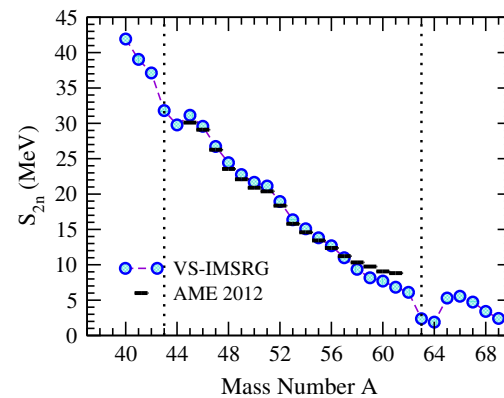
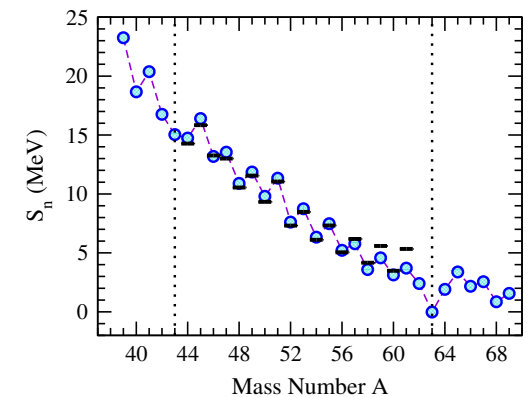
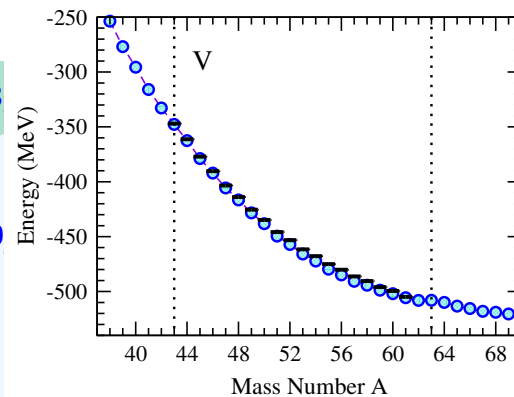
## Explore ground-state properties throughout medium-mass region



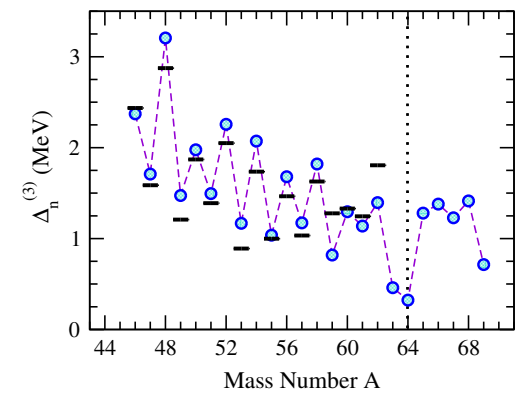
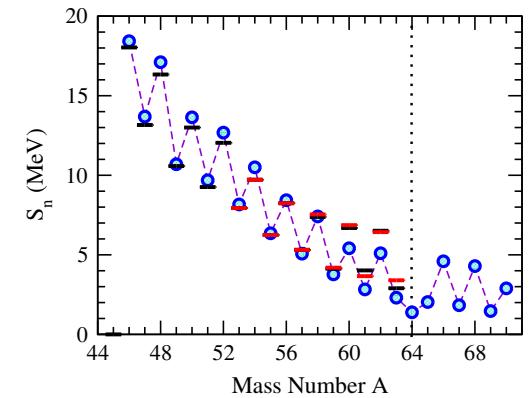
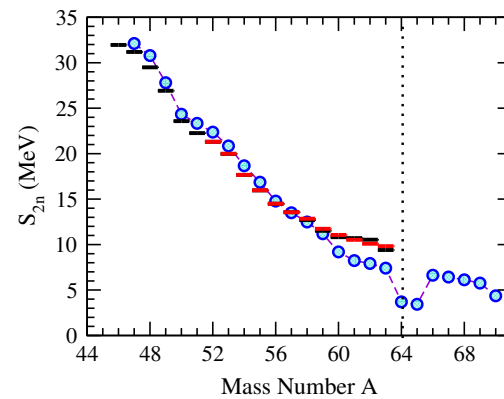
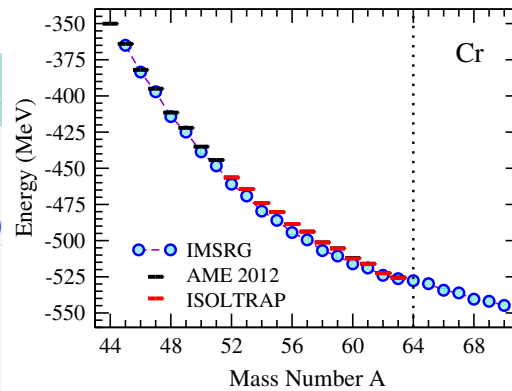
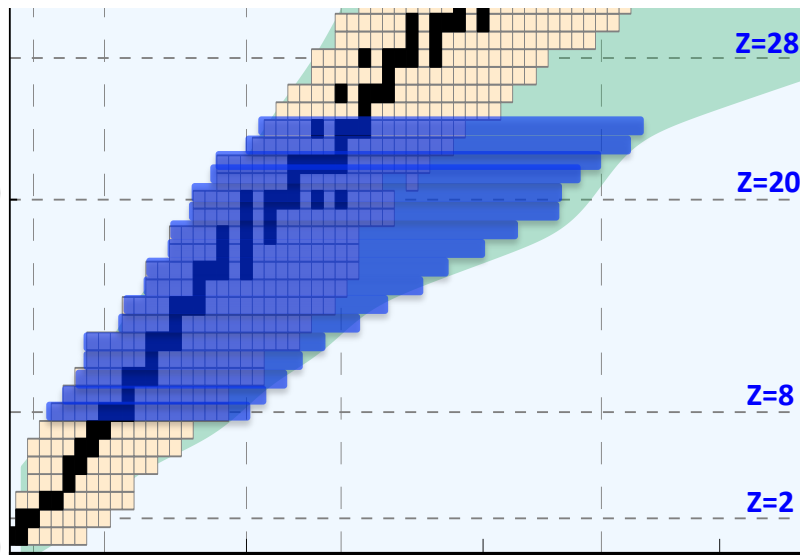
Remarkable agreement with experiment

Probe dripline and beyond

Artifacts at neutron N=40 gap

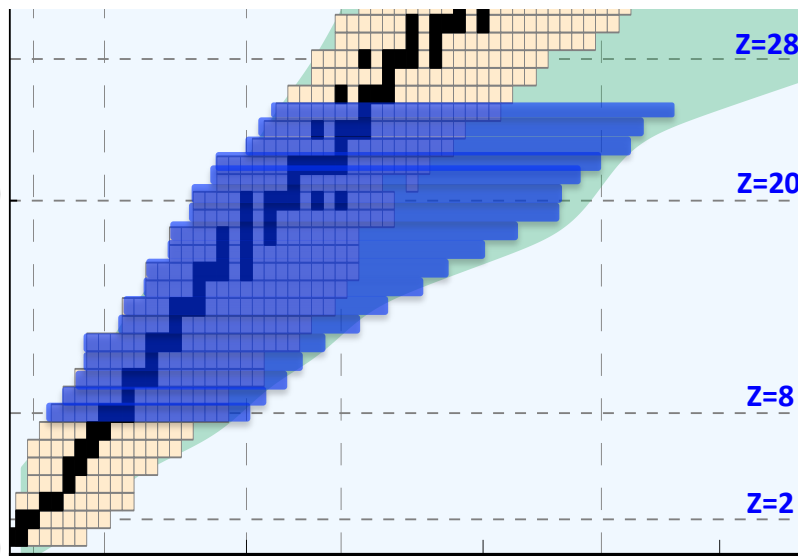


Generally deformed, new data from ISOLTRAP



JDH, Stroberg, et al., in preparation

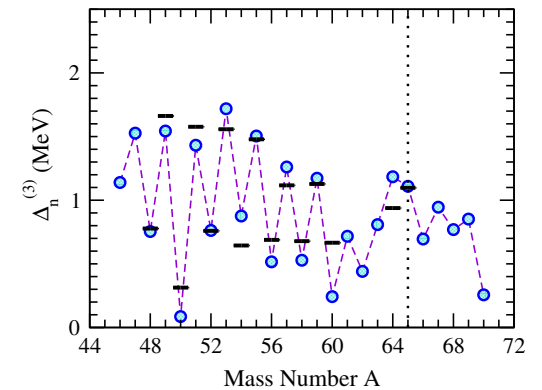
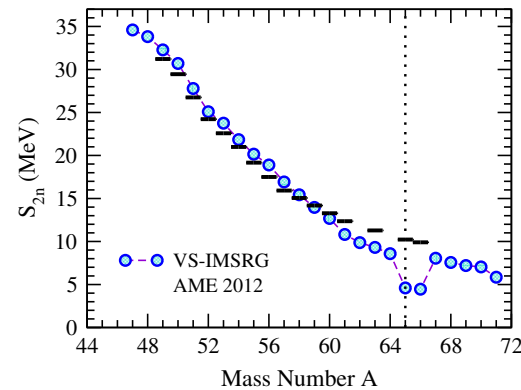
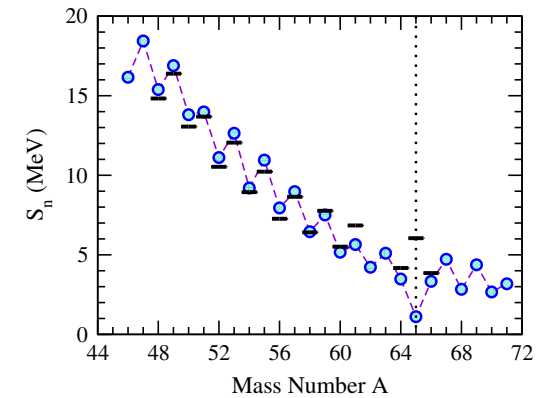
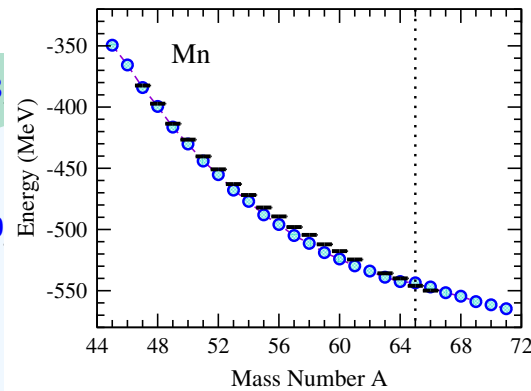
## Explore ground-state properties throughout medium-mass region



Remarkable agreement with experiment

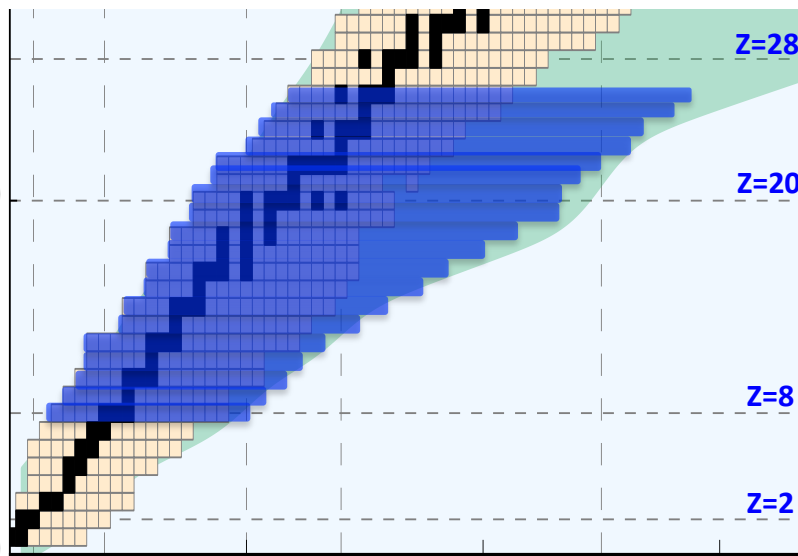
Probe dripline and beyond

Artifacts at neutron N=40 gap



JDH, Stroberg, et al., in preparation

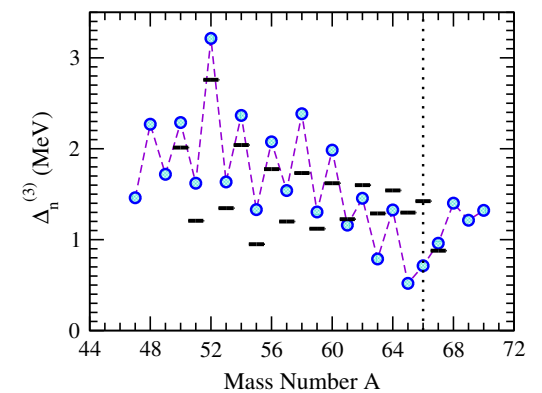
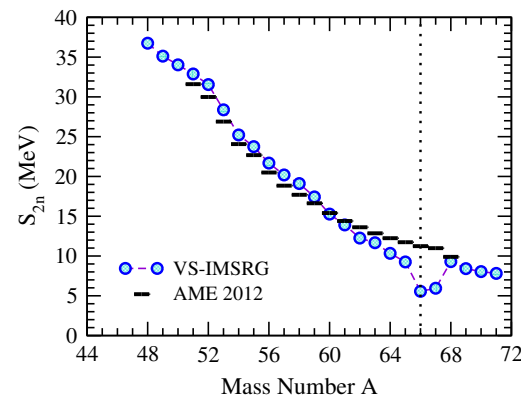
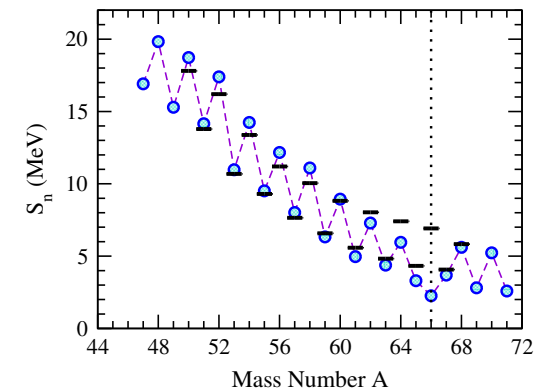
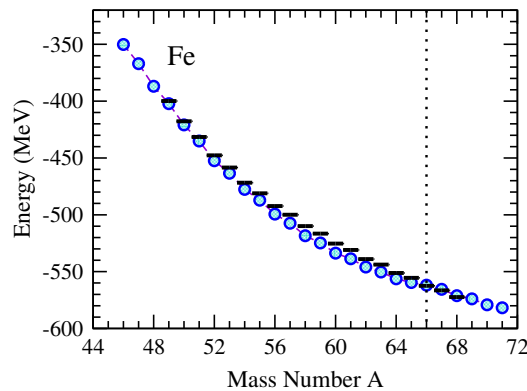
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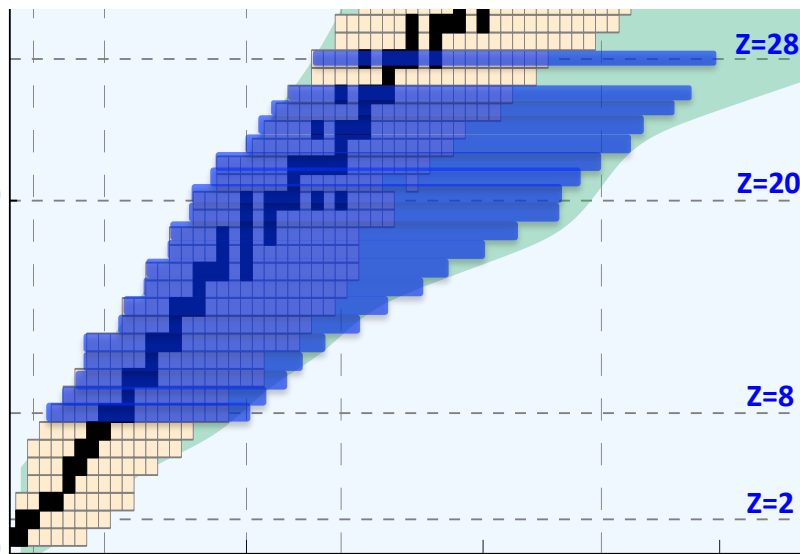
Artifacts at neutron N=40 gap



JDH, Stroberg, et al., in preparation



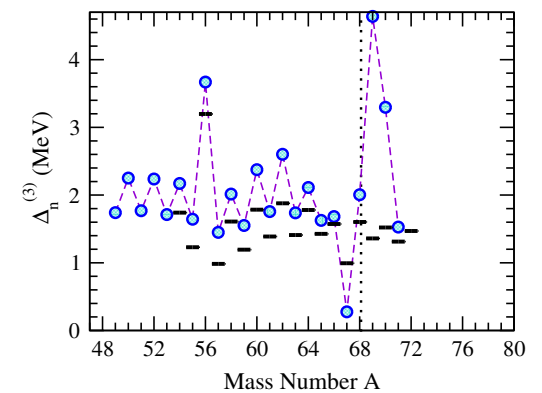
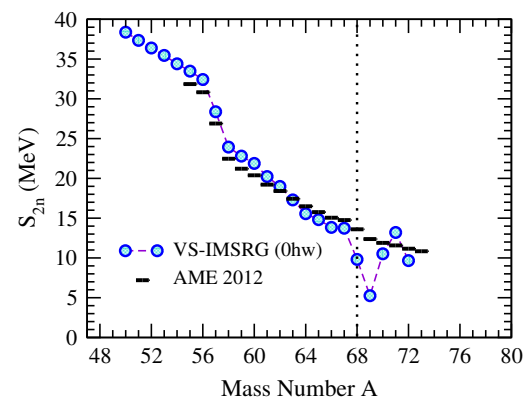
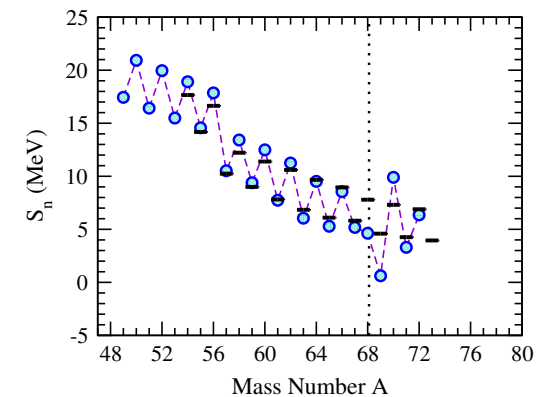
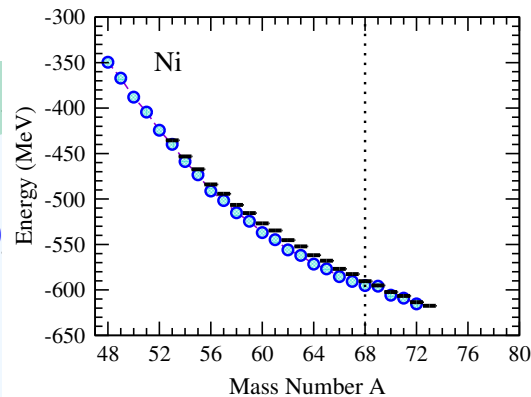
## Explore ground-state properties throughout medium-mass region



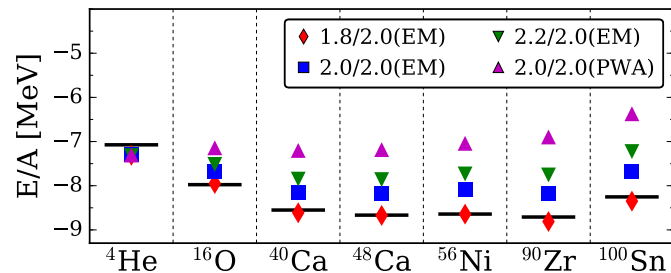
Remarkable agreement with experiment

Probe dripline and beyond

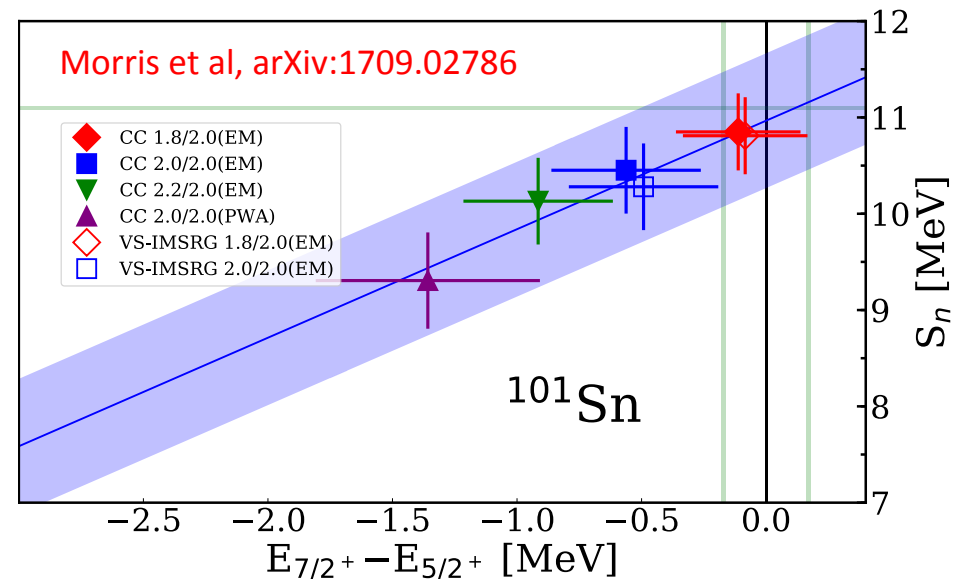
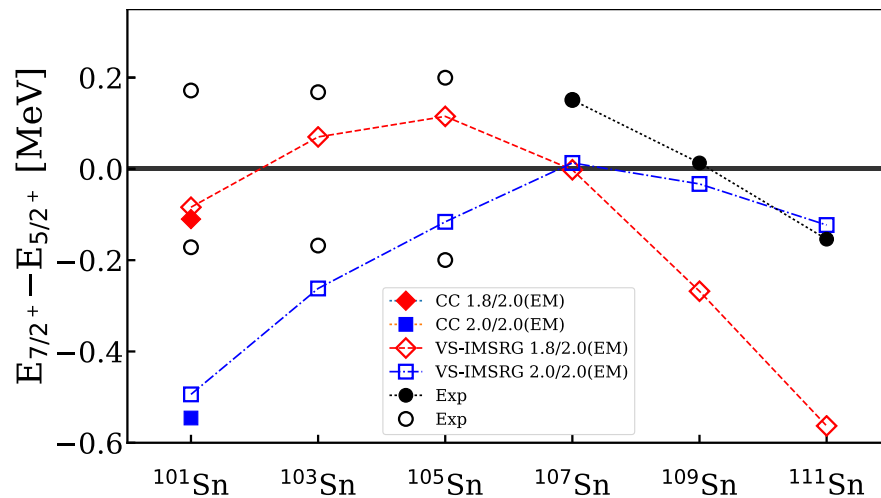
Artifacts at neutron N=40 gap



Level ordering near  $^{101}\text{Sn}$  controversial and unknown: insights from ab initio valence-space IMSRG



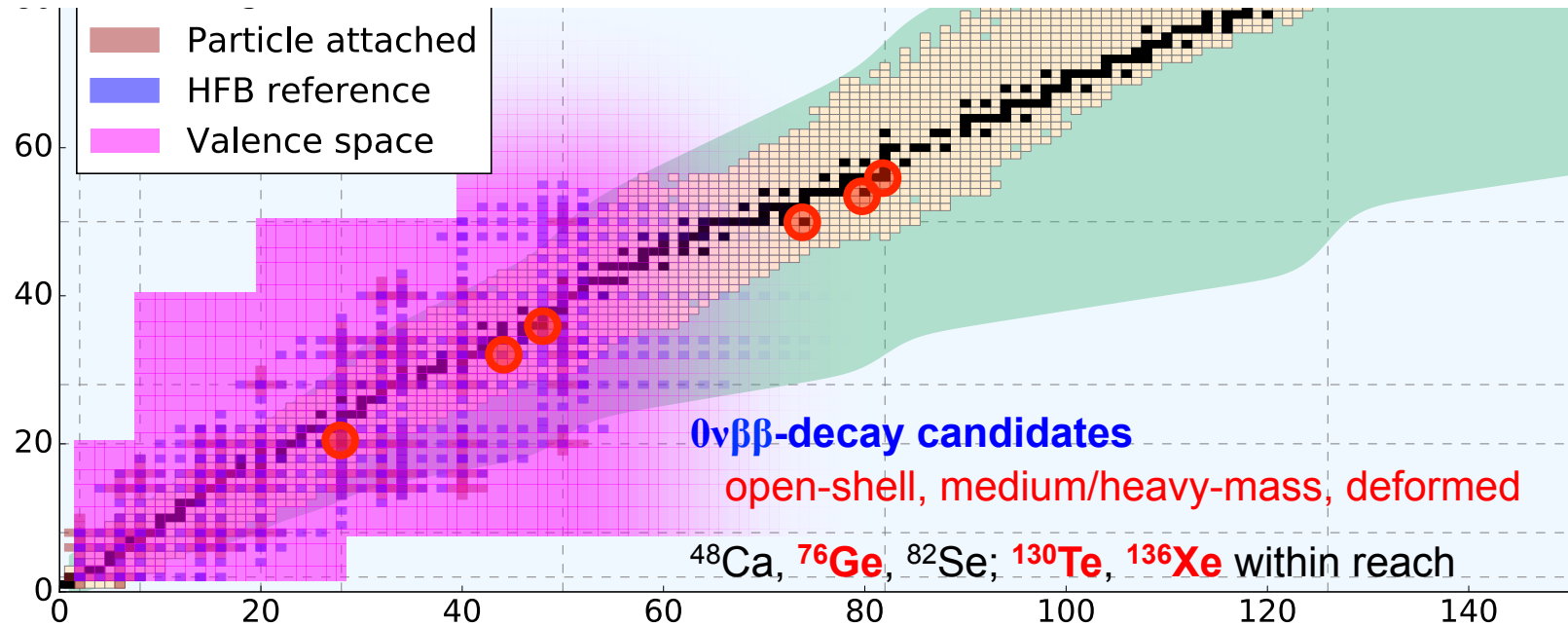
1.8/2.0(EM) reproduces gs energies across chart  
Use in study of Sn isotopes



Ab initio predicts  $5/2^+$  ground state, but within theoretical uncertainties

**Aim of modern nuclear theory:** Develop unified *first-principles* picture of structure and reactions

- Nuclear forces (low-energy QCD)
- Electroweak physics
- Nuclear many-body problem

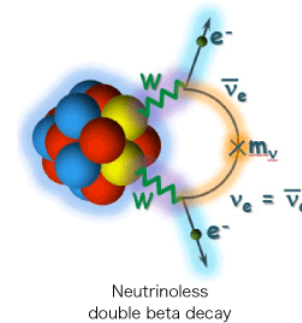


Conventional SM: phenomenological wavefunctions

Ab initio SM: wavefunctions from chiral NN+3N forces

$$M^{0\nu} = M_{GT}^{0\nu} - \frac{M_F^{0\nu}}{g_A^2} + M_T^{0\nu}$$

$$M_{GT}^{0\nu} = \langle f | \sum_{ab} H(r_{ab}) \sigma_a \cdot \sigma_b \tau_a^+ \tau_b^+ | i \rangle$$



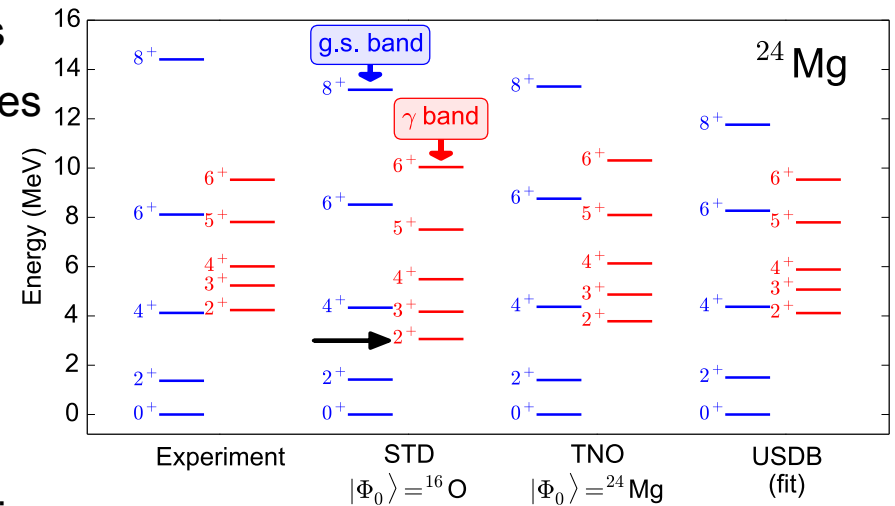
- 1) Ab initio energies in medium/heavy-mass region  
Valence-space IM-SRG for all medium-mass nuclei

Conventional SM: phenomenological wavefunctions

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Valence-space IM-SRG for all medium-mass nuclei


**Deformation challenging for large-space methods**

Conventional SM: phenomenological wavefunctions

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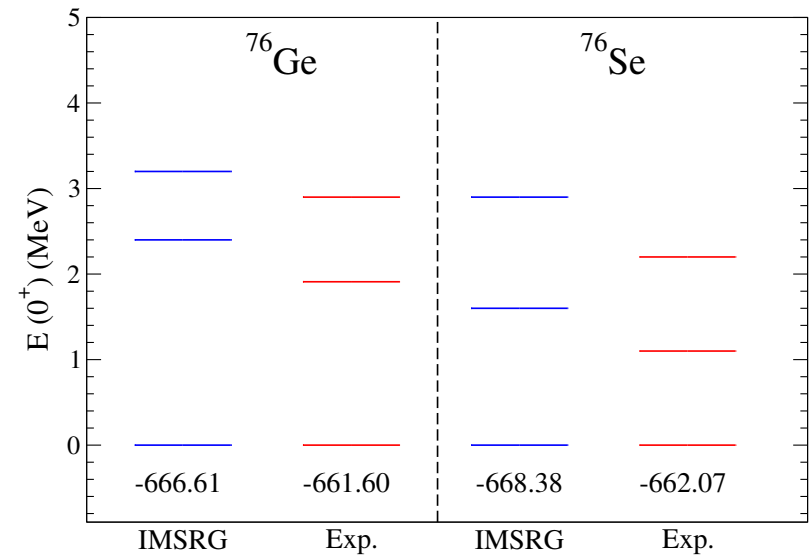
$$M_{GT}^{0\nu} = \langle f | \sum_{ab} H(r_{ab}) \sigma_a \cdot \sigma_b \tau_a^+ \tau_b^+ | i \rangle$$

1)  Ab initio energies in medium/heavy-mass region

Valence-space IM-SRG for all medium-mass nuclei

Deformation challenging for large-space methods

First ab initio calculation of  $^{76}\text{Ge}/^{76}\text{Se}$

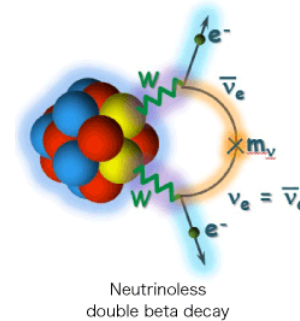


Conventional SM: phenomenological wavefunctions + **bare operator**

Ab initio SM: wavefunctions from chiral NN+3N forces + **consistent effective operator**

$$M^{0\nu} = M_{GT}^{0\nu} - \frac{M_F^{0\nu}}{g_A^2} + M_T^{0\nu}$$

$$M_{GT}^{0\nu} = \langle f | \sum_{ab} H(r_{ab}) \mathbf{M}_{\text{eff}} \cdot \sigma_b \tau_a^+ \tau_b^+ | i \rangle$$



- 1) ✓ Ab initio energies in medium/heavy-mass region
- 2) Effective decay operator: decouple valence-space operator (analogous to Hamiltonian)



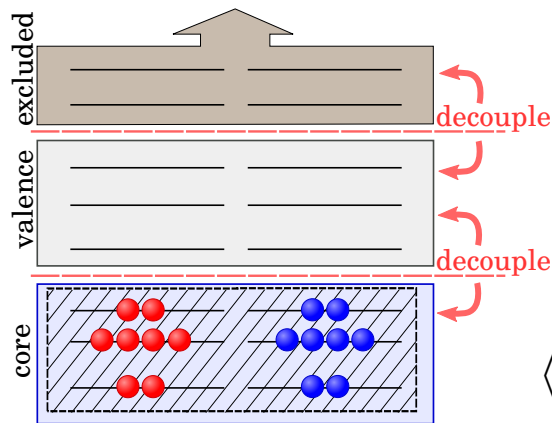
Payne, Stroberg, JDH, Menendez, in preparation

Explicitly construct unitary transformation from sequence of rotations

$$U = e^{\Omega} = e^{\eta_n} \dots e^{\eta_1}$$

$$\tilde{H} = e^{\Omega} H e^{-\Omega} = H + [\Omega, H] + \frac{1}{2} [\Omega, [\Omega, H]] + \dots$$

$$\tilde{\mathcal{O}} = e^{\Omega} \mathcal{O} e^{-\Omega} = \mathcal{O} + [\Omega, \mathcal{O}] + \frac{1}{2} [\Omega, [\Omega, \mathcal{O}]] + \dots$$



$$|\Phi_0\rangle = |^{16}\text{O}\rangle$$

**Step 1: Decouple core**

**Step 2: Decouple valence space**

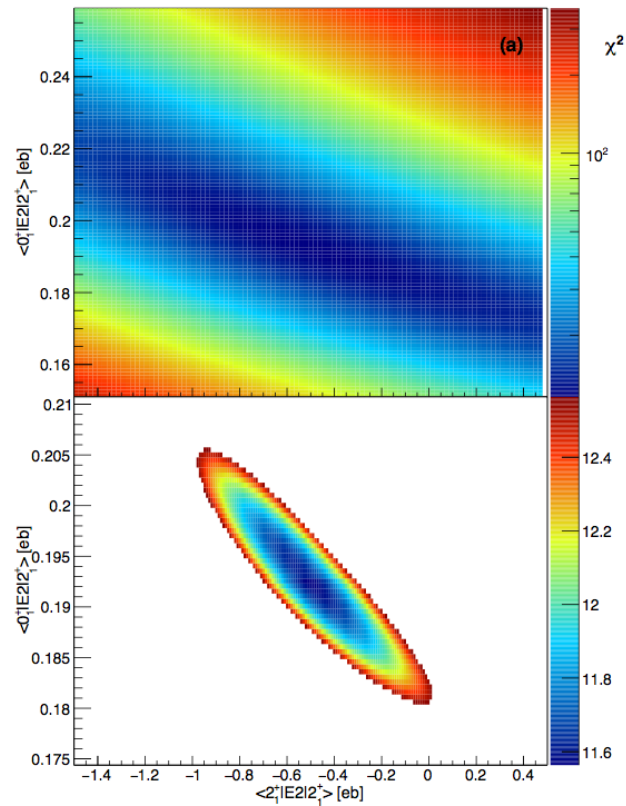
**Step 3: Decouple additional operators**

$$\langle \tilde{\Psi}_n | P \tilde{H} P | \tilde{\Psi}_n \rangle \approx \langle \Psi_i | H | \Psi_i \rangle$$

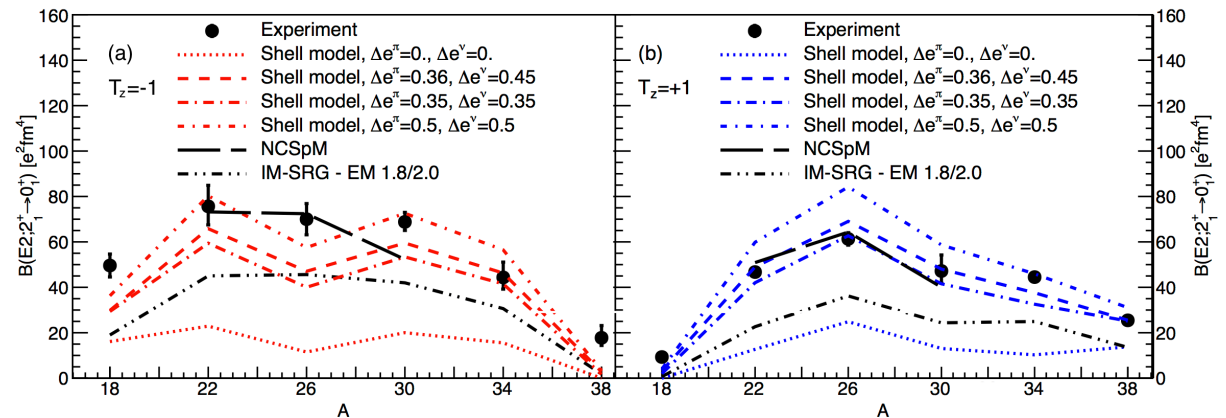
$$\langle \tilde{\Psi}_n | P \tilde{M}_{0\nu} P | \tilde{\Psi}_n \rangle \approx \langle \Psi_i | M_{0\nu} | \Psi_i \rangle$$

$\langle P   H   P \rangle$	$\langle P   H   Q \rangle \rightarrow 0$
$\langle Q   H   P \rangle \rightarrow 0$	$\langle Q   H   Q \rangle$





- Use GOSIA Coulomb-excitation code to extract matrix elements
- Compare with NCSpM (LSU) and VS-IM-SRG (TRIUMF)
  - NCSpM - does excellent job - expensive calculations
  - VS-IM-SRG underpredicts strength - relatively inexpensive - qualitative description excellent

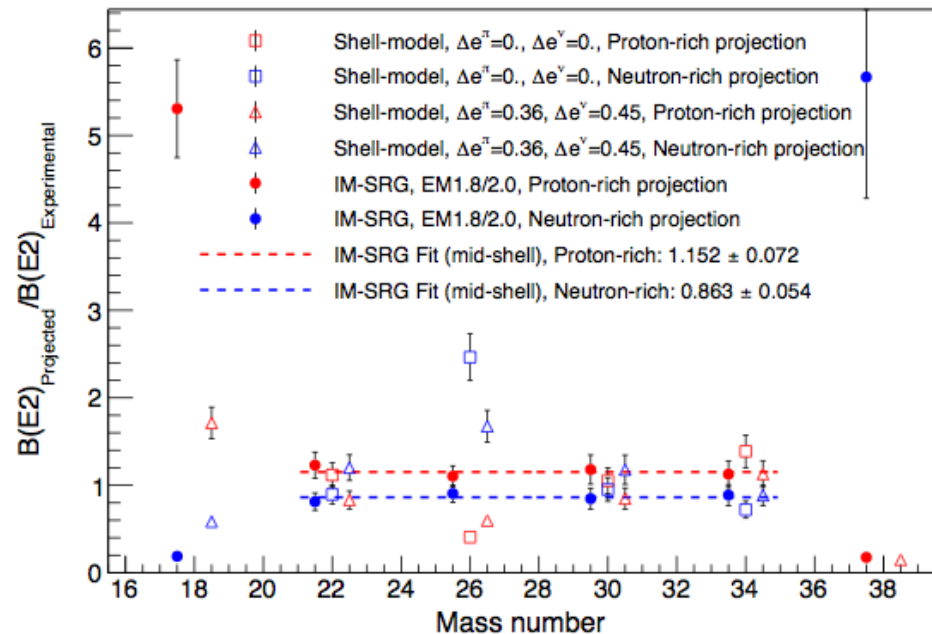


Assess nature of missing VS-IM-SRG E2 strength:

$$B(E2)_{T_z=-1}^{\text{Proj.}} = B(E2)_{T_z=-1}^{\text{Theory}} \times \frac{B(E2)_{T_z=+1}^{\text{Exp}}}{B(E2)_{T_z=+1}^{\text{Theory}}},$$

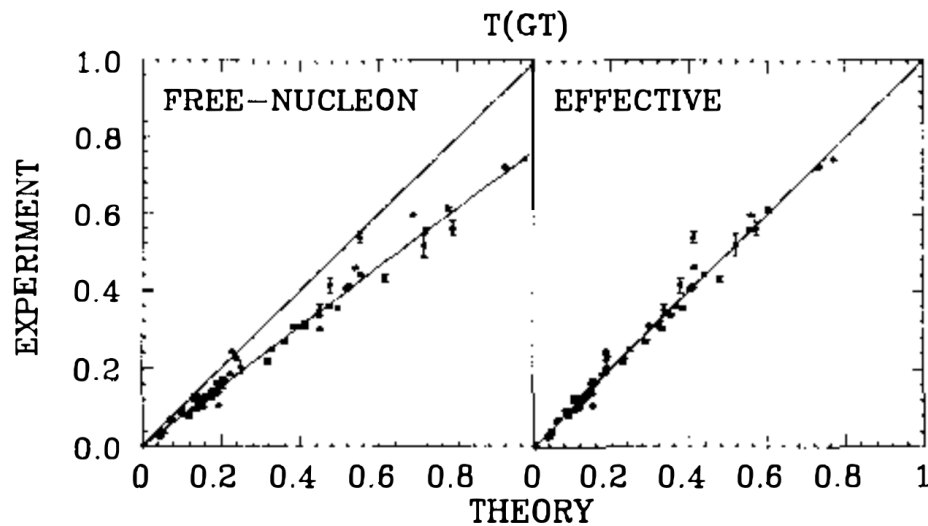
If missing E2 strength isoscalar, expected “projected” B(E2) to match experiment

- Projected B(E2) consistently 15% over/under predicted by VS-IM-SRG
  - Missing strength has consistent isovector component
  - Promising for future development
- Shell model (USDB) shows no consistent behaviour

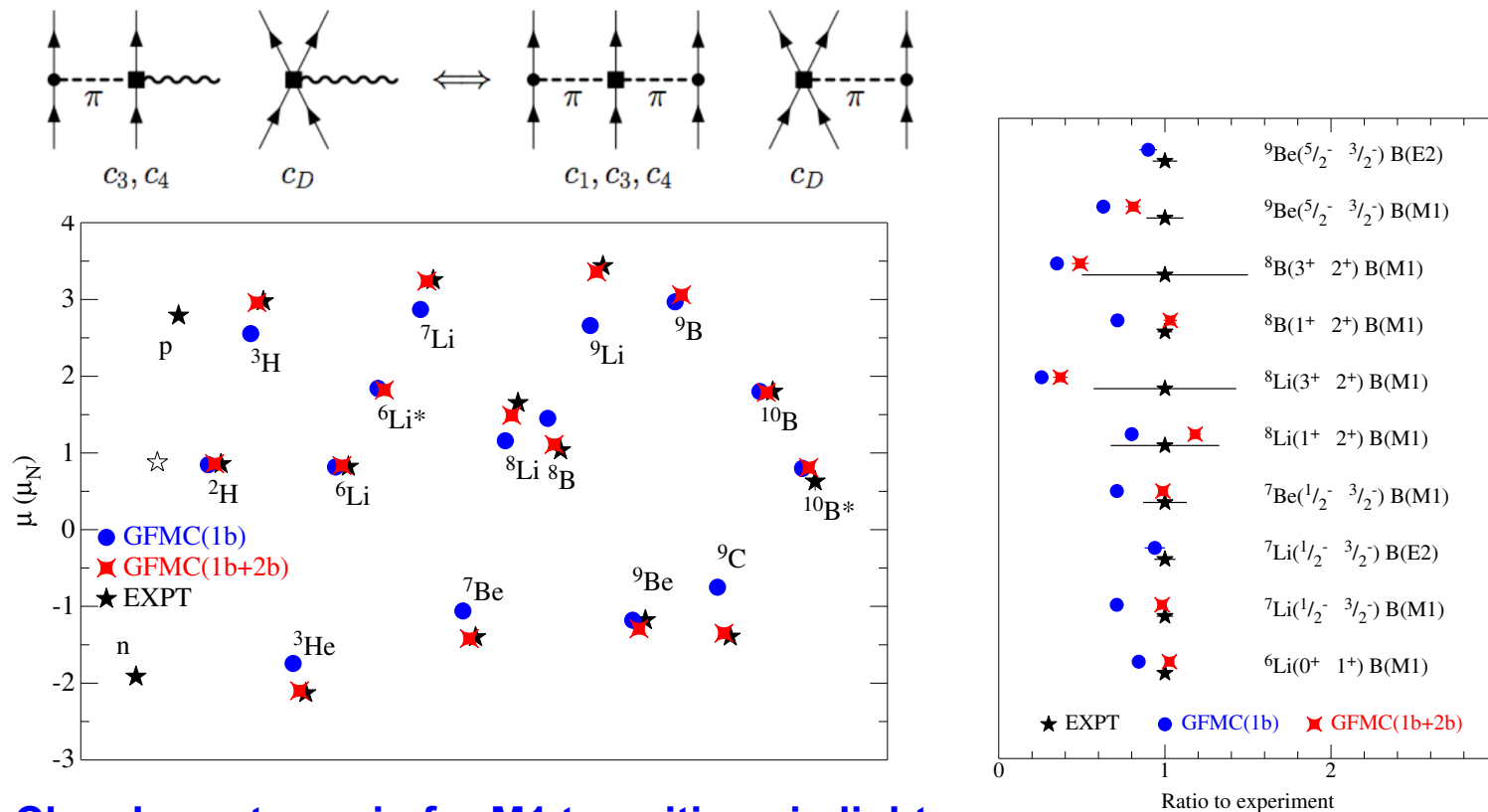


Long-standing problem in weak decays of nuclei: should  $g_A$  be “quenched”?

Using  $g_A^{\text{eff}} \approx 0.77 \times g_A^{\text{free}}$  agrees with data

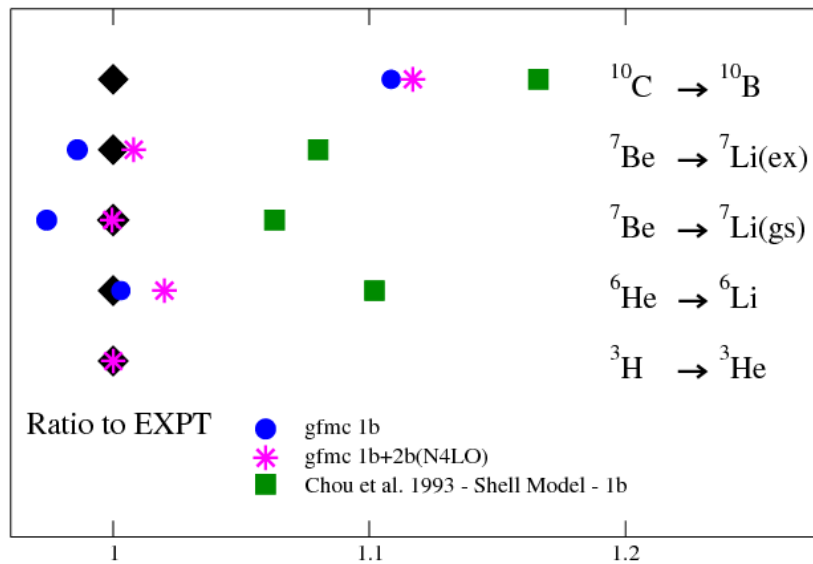
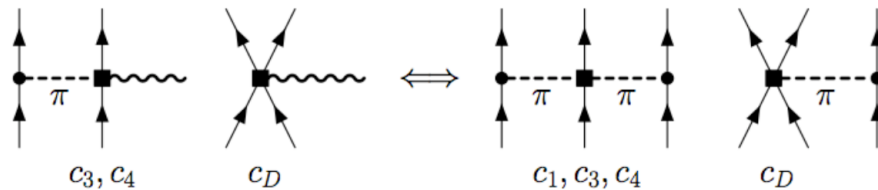


Chiral Effective Field Theory – electroweak currents consistent with nuclear forces



Clear Importance in for M1 transitions in light nuclei

Chiral Effective Field Theory – electroweak currents consistent with nuclear forces

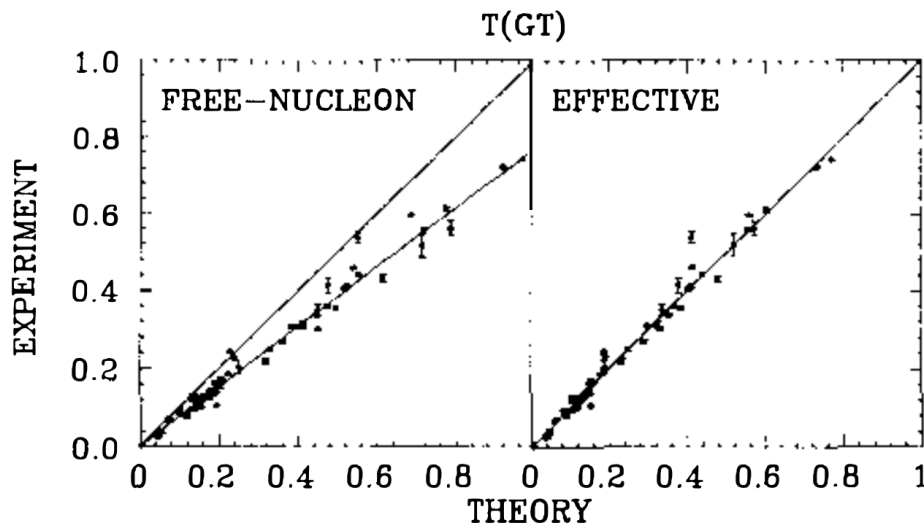


**Clear increase for GT transitions in GFMC – inconsistent forces/currents**

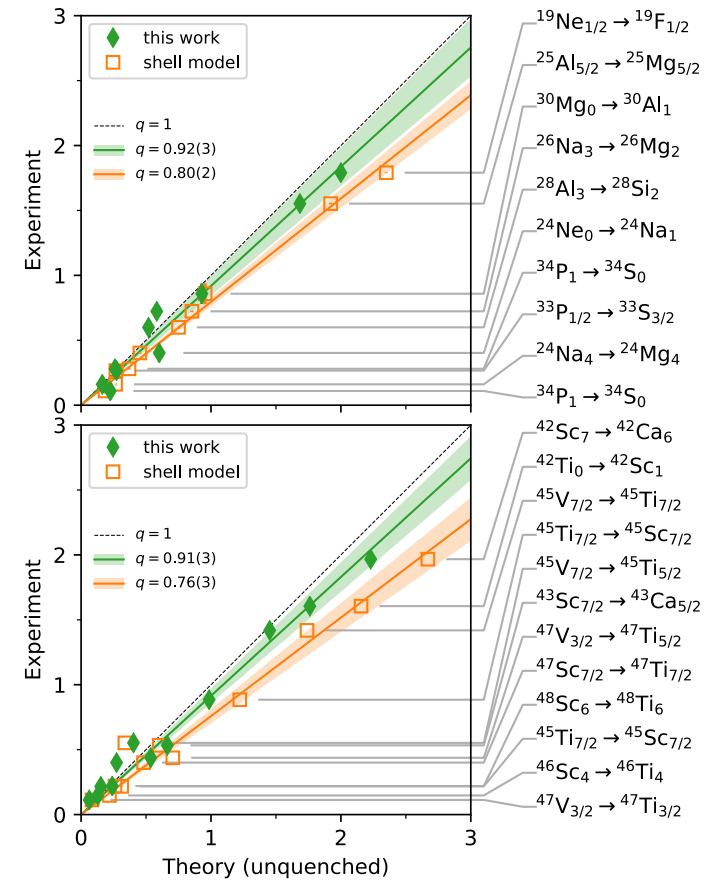
VS-IMSRG calculations of GT transitions in sd, pf shells

**Minor effect from consistent effective operator**

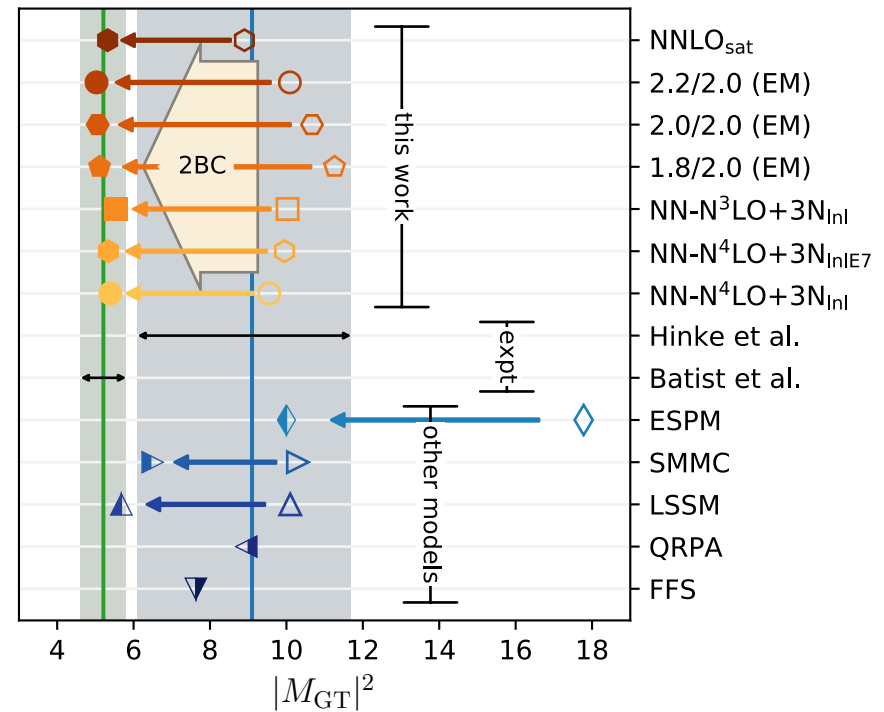
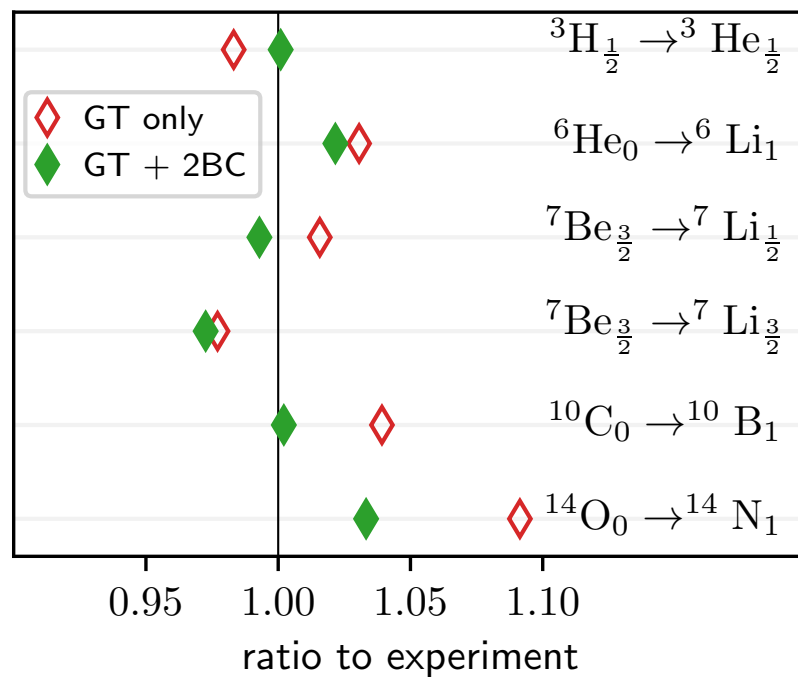
**Significant effect from neglected 2-body currents**



**Ab initio calculations explain data with unquenched  $g_A$**

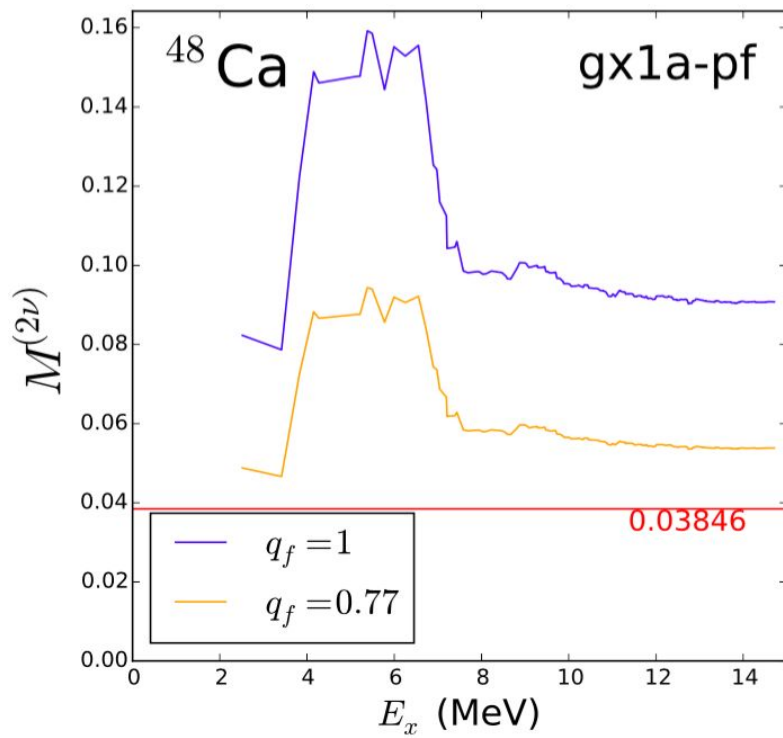


Prediction from light nuclei to super allowed GT transition in  $^{100}\text{Sn}$

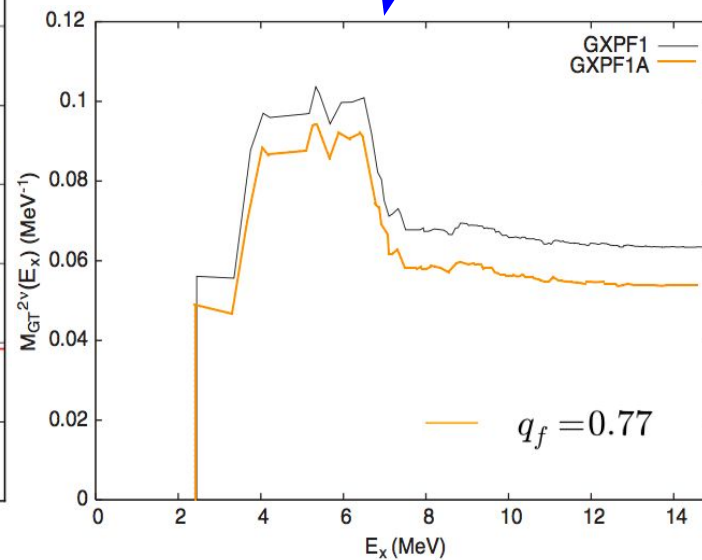


**Agreement with data with no need for quenching**

First benchmark to reproduce known shell-model results

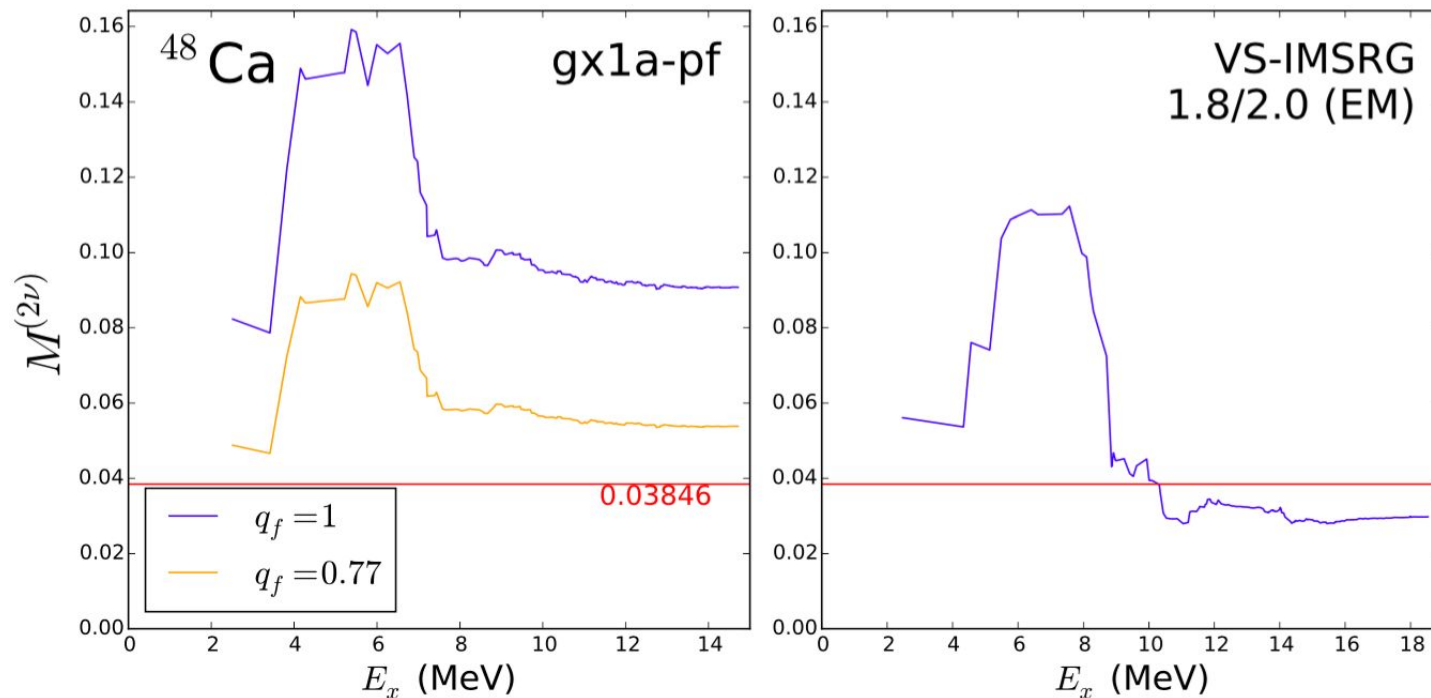


VS = pf-shell, with **no** IMSRG evolution reproduces ShM results [6]





Consistent many-body wfs/operators from chiral NN+3N forces (no 2b currents)

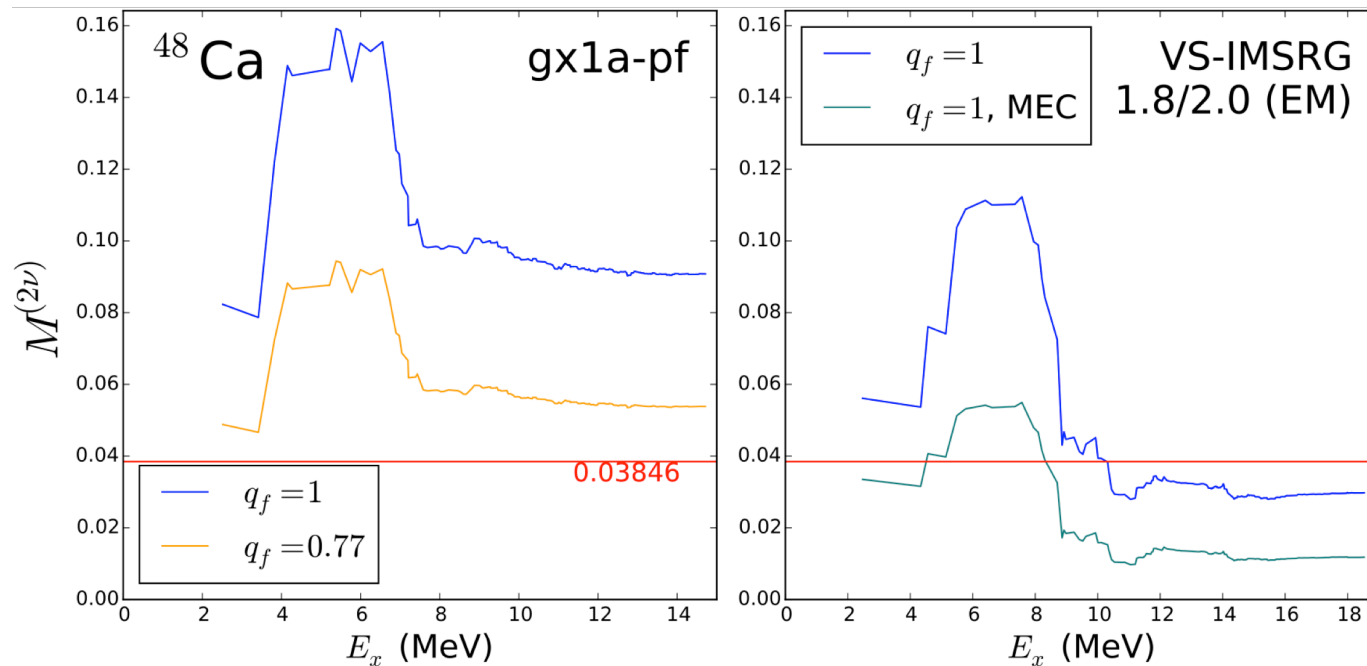


Payne, Stroberg, JDH, et al., in prep

**VS-IMSRG: decrease in final matrix element**

Likely missing contributions from intermediate states outside valence space

Consistent many-body wfs/operators from chiral NN+3N forces (**with 2b currents**)

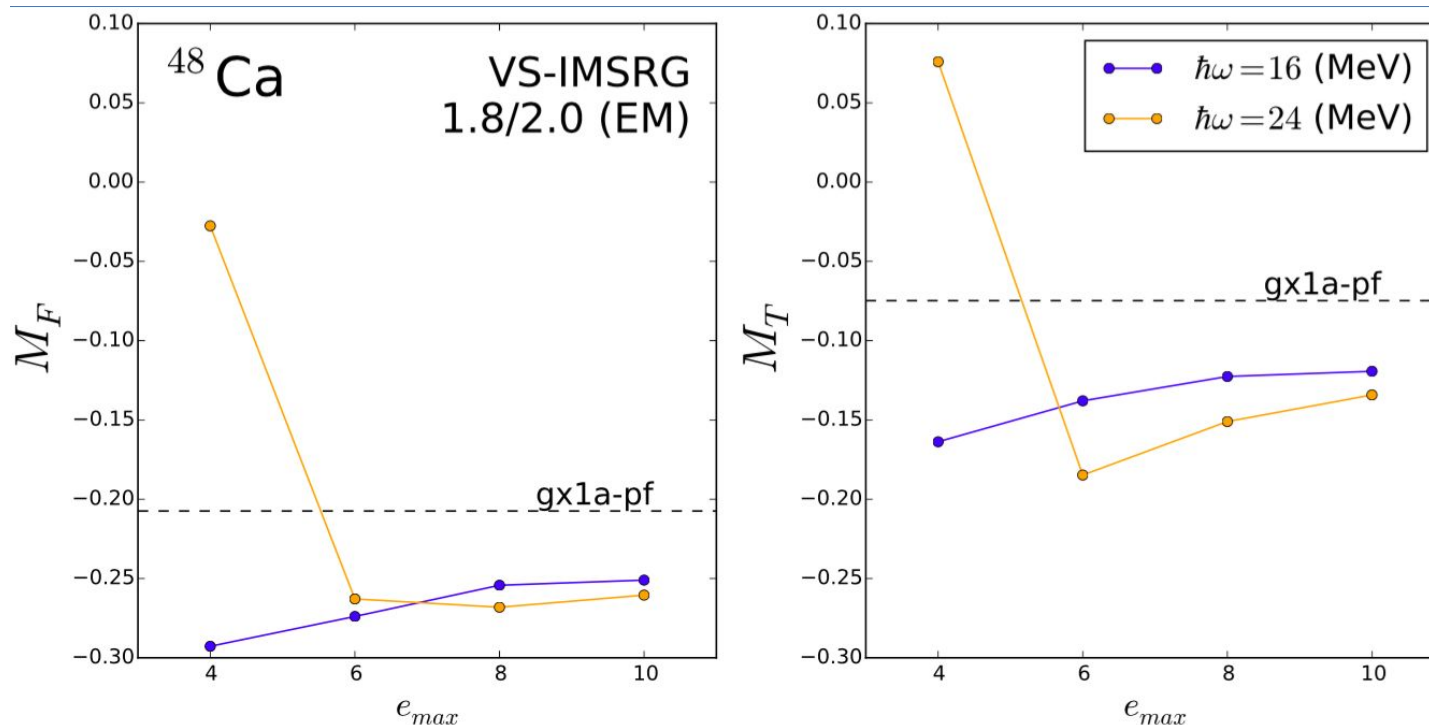


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**VS-IMSRG: decrease in final matrix element**

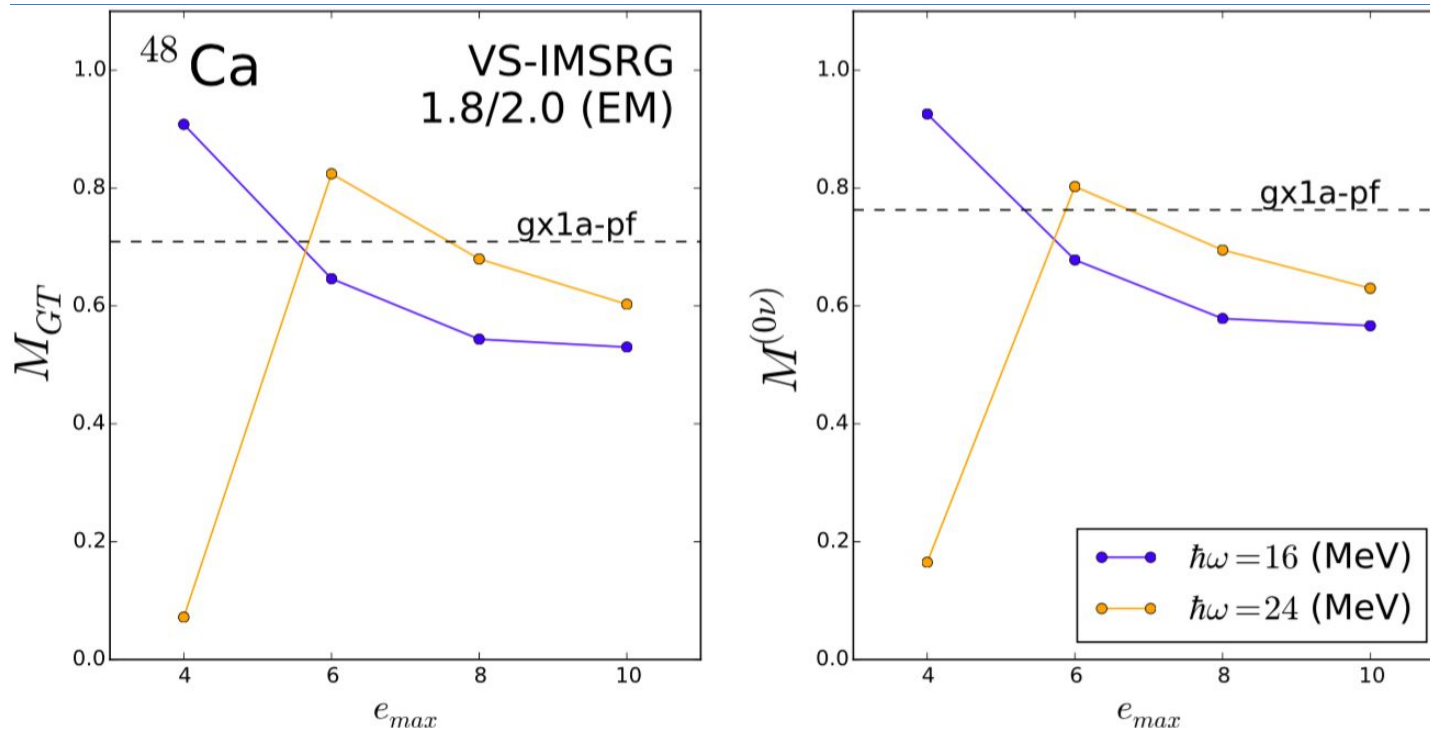
Likely missing contributions from intermediate states outside valence space

Consistent many-body wfs/operators from chiral NN+3N forces (no 2b currents)



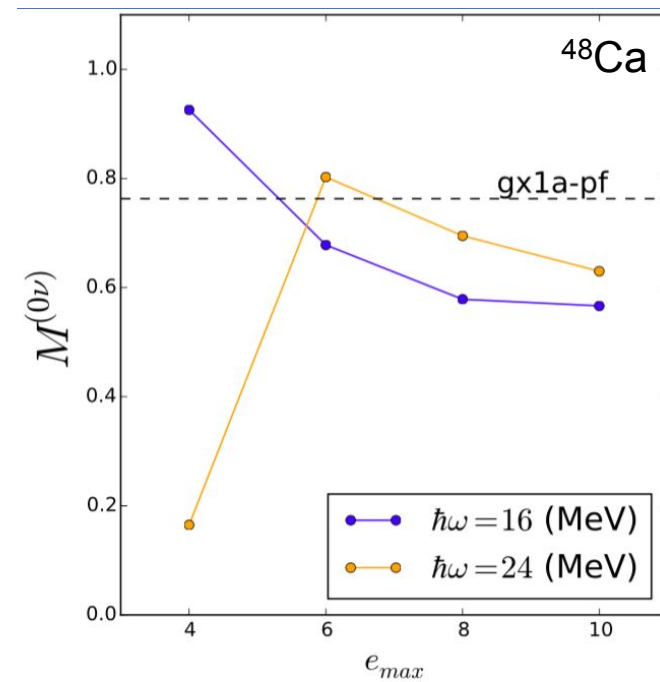
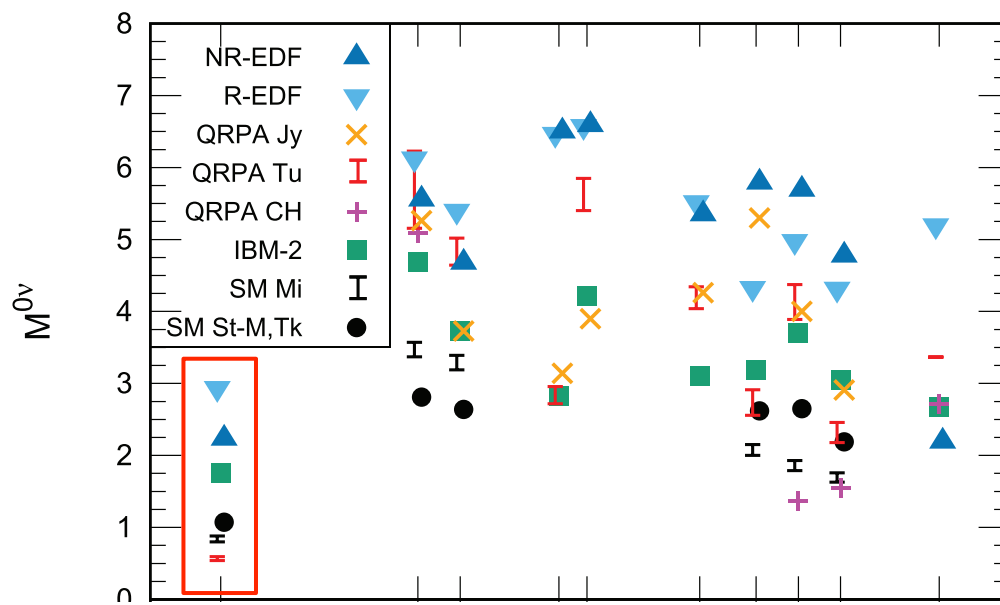
General cancellation between Fermi and Tensor contributions

Consistent many-body wfs/operators from chiral NN+3N forces (no 2b currents)

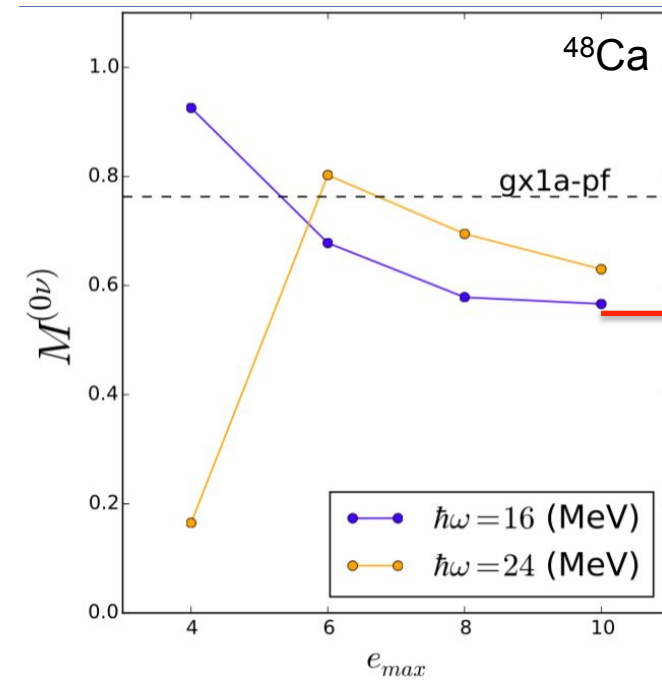
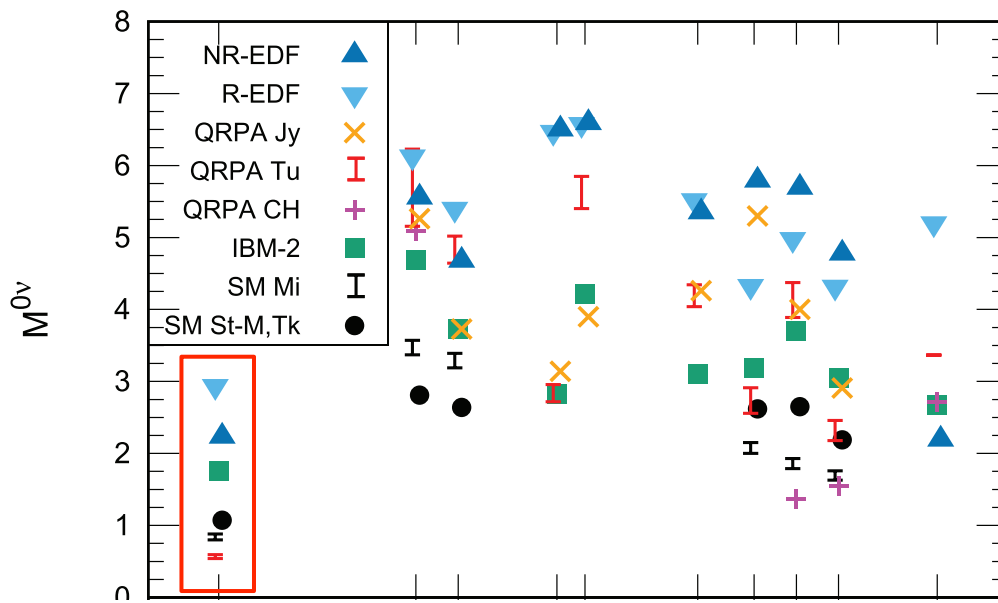


Final matrix element converged – significant decrease from phenomenology

Ab initio: Consistent many-body wfs/operators from chiral NN+3N forces



Ab initio: Consistent many-body wfs/operators from chiral NN+3N forces



**Consistent prediction from independent method**

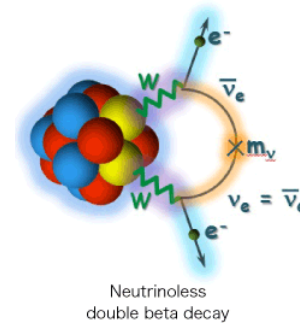
Two-body currents in progress – typically decrease NME

Standard SM: phenomenological wavefunctions + bare operator

Ab initio SM: wavefunctions from chiral NN+3N forces + consistent effective operator

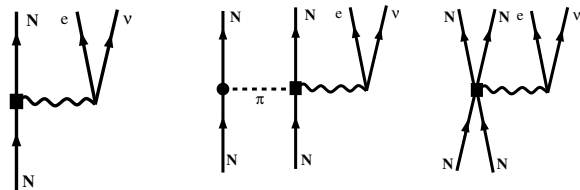
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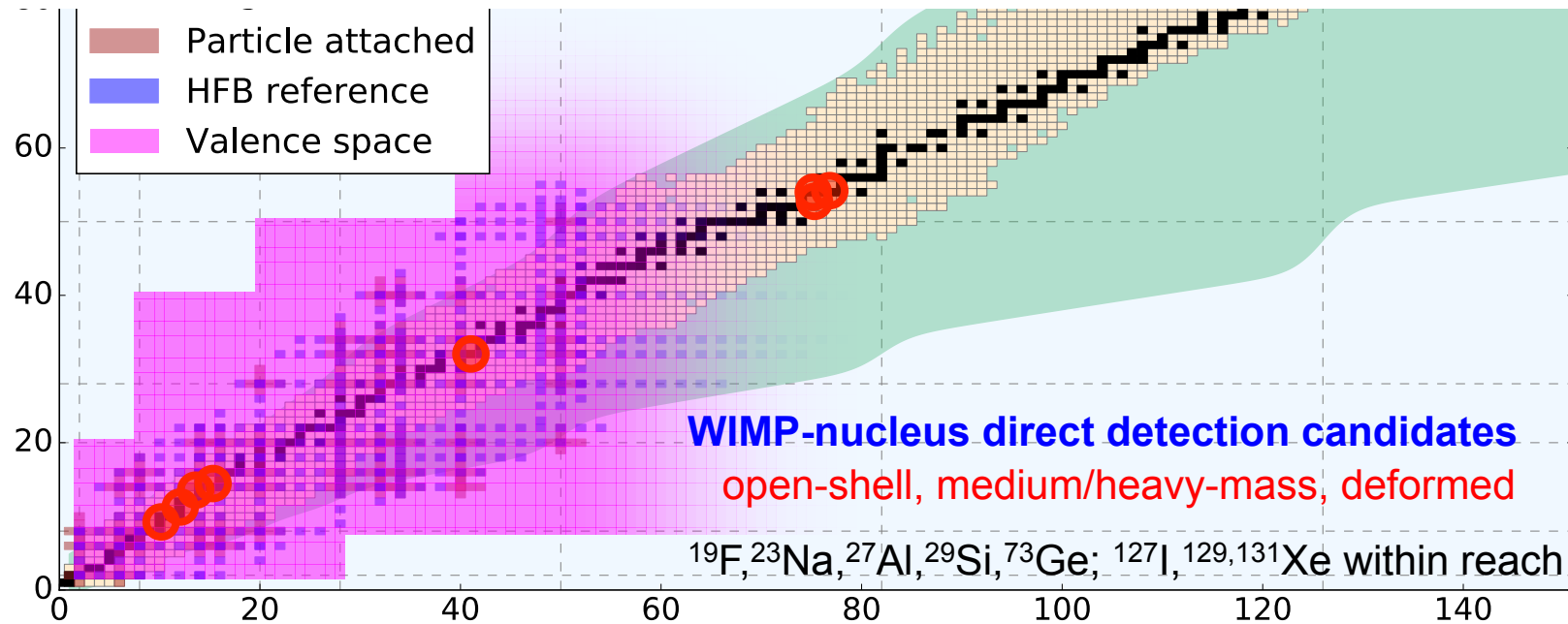
- 1) ✓ Ab initio energies in medium/heavy-mass region
- 2) ✓ Effective decay operator: decouple valence-space operator
- 3) ⚠ Operator corrections

**Two-body currents** S. Leutheusser (UBC/MIT)



**Aim of modern nuclear theory:** Develop unified *first-principles* picture of structure and reactions

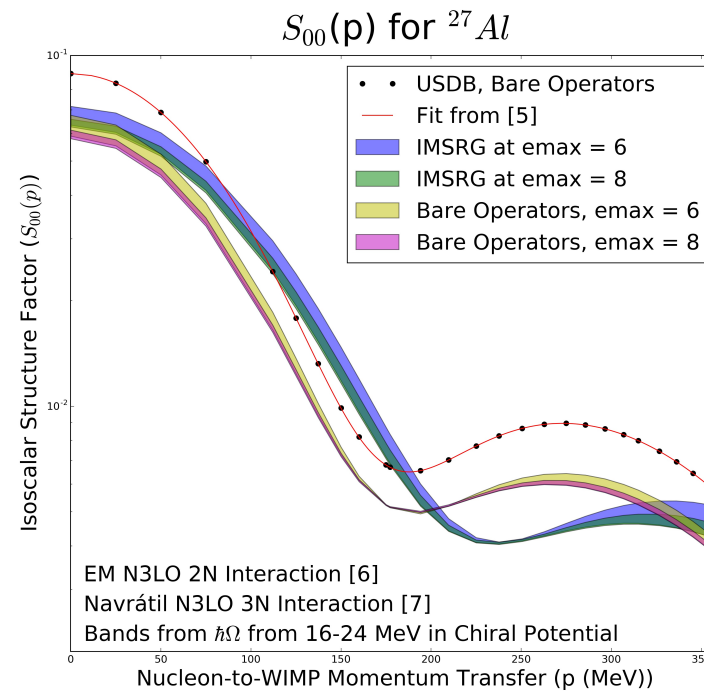
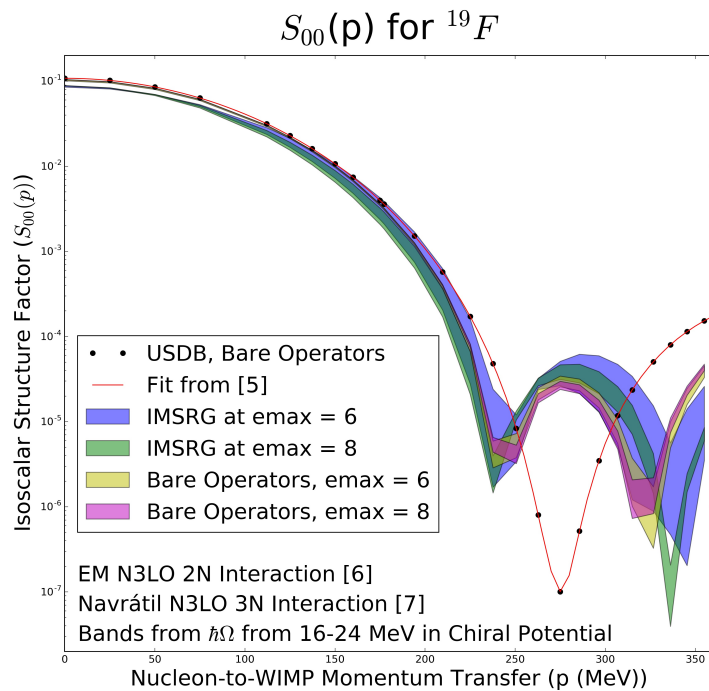
- Nuclear forces (low-energy QCD)
- Electroweak physics
- Nuclear many-body problem





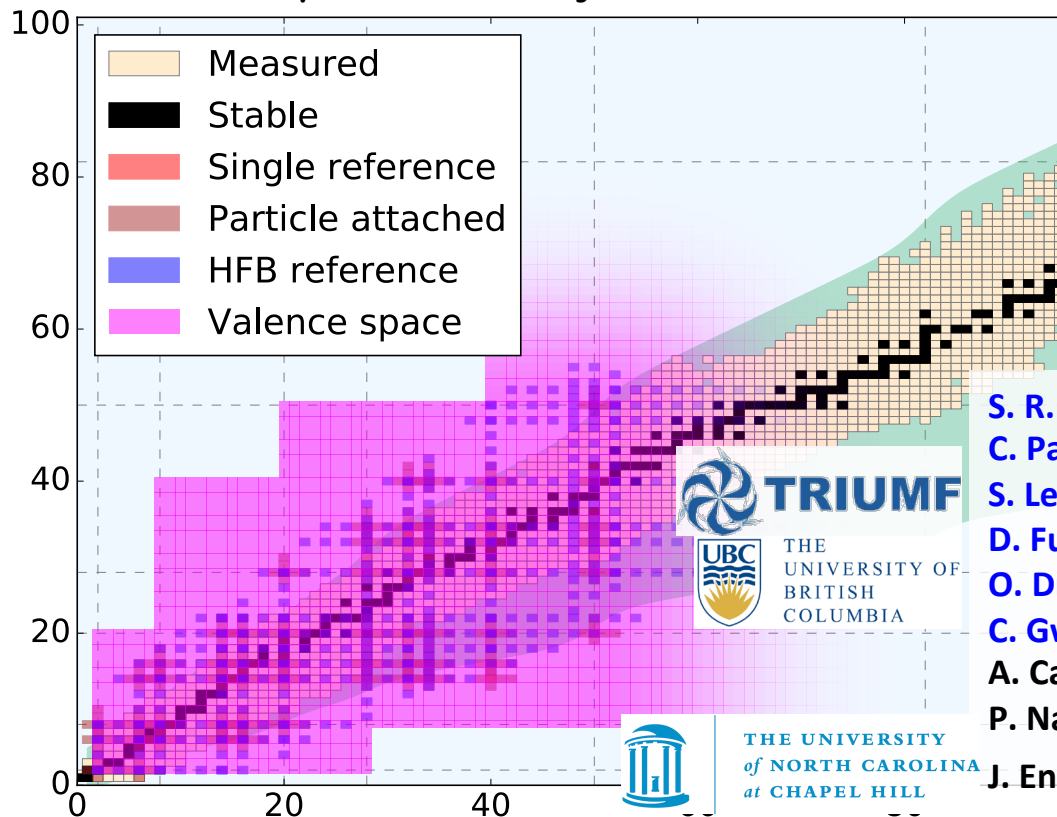
Ab initio: Consistent many-body wfs/operators from chiral NN+3N forces + 2b currents

$$S_A(p) = \sum_{L \geq 0} |\langle J_f \| \mathcal{L}_L \| J_i \rangle|^2 + \sum_{L \geq 0} \left( |\langle J_f \| \mathcal{T}_L^{\text{el}} \| J_i \rangle|^2 + |\langle J_f \| \mathcal{T}_L^{\text{mag}} \| J_i \rangle|^2 \right)$$



### Ab initio valence-shell Hamiltonians

First ab initio prediction of nuclear driplines  
 Cross-shell spaces underway: Island of inversion



### Fundamental physics

Effective electroweak operators: M1, GT, ...  
 Effective  $0\nu\beta\beta$  decay operator  
 WIMP-Nucleus scattering

### Outstanding issues

Controlled IMSRG(3) approximation  
 E2 operators problematic  
 Continuum essential beyond stability  
 Too-high 2+ energies at closed shells  
 Quantify uncertainties

S. R. Stroberg  
 C. Payne  
 S. Leutheusser  
 D. Fullerton  
 O. Drozdowski  
 C. Gwak  
 A. Calci  
 P. Navrátil  
 J. Engel



H. Hergert  
 S. Bogner  
 N. Parzuchowski



J. Simonis  
 A. Schwenk



T. Morris  
 G. Hagen, T. Papenbrock



J. Menéndez