

Neutron-antineutron in nuclei

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Table of contents

- 1 History
- 2 Other examples of mixing
 - S - D mixing in the deuteron
 - S - D mixing in charmonium
 - S - D mixing in muonium or hydrogen
- 3 Deuteron lifetime
- 4 Lifetime of ^{16}O
- 5 Conclusions

History

- $n - \bar{n}$ oscillation, see Mohapatra and others
- Meeting Maurice Goldhaber: 1982, before and after, e.g. 2006

13-45-14-05	J. M. Richard	$\chi(3872)$ as a four-quark state	PDF
14-05-14-25	P. Wang	The $\Upsilon(4260)$ as an omega χ_{b1} molecular state	PDF
14-25-14-45	E. Kou	Suppressed decay into open charm for the $\Upsilon(4260)$ being an hybrid	PDF
14-45-15-05	F. Llanes Estrada	$\Upsilon(4260)$, to what extent a conventional charmonium state?	PPT
15-05-15-25	C.-F. Qiao	The theoretical understanding of $\Upsilon(4260)$	PPT
15-25-15-45	A. Szczepaniak	Examining the evidence for constituent gluons and implications for the spectrum of hybrids.	20'

- Controversy
 - Hot seminars on the subject
 - Nazarek and others

New limit on the neutron–antineutron transitions
 V.I. Nazarek¹
Institute for Nuclear Research of the Academy of Sciences of Russia, 600 October Anniversary Prospect 7a, 117312 Moscow, Russia
 Received 20 July 1991
 Editor P.V. Landshoff

Abstract

We reexamine the problem of extracting a lower limit on the free-space $n\bar{n}$ oscillation time $\tau_{n\bar{n}}$ from the nuclear annihilation lifetime. It is shown that the $n\bar{n}$ transitions in the molecules are suppressed only by a factor 0.5. As a result we get $\tau_{n\bar{n}} = 3 \times 10^{17}$ yr, which increases the previous limit by 2.1 orders of magnitude.

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PHYSICAL REVIEW

Limits on $n - \bar{n}$ Oscillations

A recent paper¹ reported an upper limit of $0.7 \times 10^{20} \text{ yr}^{-1}$ for the rate of $n - \bar{n}$ transitions² in oxygen nuclei and deduced a corresponding lower limit of 2×10^9 s for the free-neutron oscillation time, by use of a relation taken from Dover, Gal, and Richard.³ This is the latest in a series of papers (Refs. 11–17 of Ref. 1) which reflect the prevalent view that there is a direct relation between the $n - \bar{n}$ transition rates for free neutrons and for those inside a nucleus. This has led to the unjustified interpretation that improved tests

- Our results supported by Alberico et al., Kopeliovich et al.,

NEUTRON–ANTINEUTRON OSCILLATIONS IN NUCLEI

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Received 11 June 1990





S-D mixing in the deuteron

Rarita-Schwinger equations

$$\psi = \frac{u(r)}{r} |^3S_1\rangle + \frac{w(r)}{r} |^3D_1\rangle$$

$$-u''(r) + m V_{00} u(r) + m V_{02} w(r) = m E u(r) ,$$

$$-w''(r) + \frac{6}{r^2} w(r) + m V_{22} w(r) + m V_{02} u(r) = m E w(r) ,$$

$$V_{00} = V_c , \quad V_{22} = V_c - 2 V_T - 3 V_{LS} - 3 V_{LL} , \quad V_{02} = \sqrt{8} V_T .$$

Only about 5%, but crucial for the deuteron and many other nuclear states. See Ericson & Rosa-Clot and Blatt & Weisskopf



S - D mixing in charmonium

- At first, $J/\psi = 1S$, $\psi' = 2S$, $\psi'' = 1D$, etc.
- Leptonic coupling of ψ'' requires some S -wave admixture
- Usually

$$|S(^3D_1)\rangle \simeq \frac{\langle 2^3S_1 | \sqrt{8} V_T | 1^3D_1 \rangle}{E_0(2S) - E_0(1D)} |2^3S_1\rangle .$$

- Solving RS eqs. in specific models indicate some important $1S$ admixture: states with same node structure mix better
- Also

$$\psi^{(n)} \leftrightarrow D^{(*)} \bar{D}^{(*)} \leftrightarrow \psi^{(m)}$$

e.g., Cornell model



S-D mixing in muonium or hydrogen

- Quadrupole deformation of an atom such as (μ^+, e^-)
- Small effect in principle measurable in a gradient of electric field
- More delicate than in a $(Q\bar{Q})$ potential model, as the sum on intermediate states (if performed!), extends over the continuum

$$|D(\text{g.s.})\rangle = \sum_n \frac{\langle n^3 D_1 | \sqrt{8} V_T | 1^3 S_1 \rangle}{E_0(1S) - E_0(nD)} |n^3 D_1\rangle .$$

- Sternheimer (Dalgarno & Lewis) equation

$$-w''(r) + \frac{6}{r^2} w(r) + m V_{22} w(r) + m V_{02} u_0(r) = m E_0 w(r) ,$$

- Good surprise: can be solved analytically, leading to a **compact expression** for the quadrupole moment of the ground state.



Deuteron lifetime

- Simplest nucleus. We restrict to S -wave, but including D -wave is straightforward
- Hulthen wave function

$$u(r) = N [\exp(-ar) - \exp(-br)] ,$$

with $a = 0.04570$ and $b = 0.2732 \text{ GeV}^{-1}$, and the proper behavior at $r \rightarrow 0$ and $r \rightarrow \infty$.

- **antineutron** component given by the Sternheimer equation

$$-w''(r) + m W w(r) - m E_0 w(r) = -m \gamma u(r) ,$$

with $E_0 = -0.0022 \text{ GeV}$ deuteron energy, $\gamma = 1/\tau(n\bar{n})$ strength of transition, and W complex potential of the $N\bar{N}$ interaction.

- **width** given by

$$-\frac{\Gamma}{2} = \int_0^\infty \text{Im } W |w(r)|^2 dr = -\gamma \int_0^\infty u(r) \text{Im } w(r) dr ,$$

Deuteron lifetime-2

- One gets (valid for other nuclei)

$$\Gamma \propto \gamma^2, \quad T = T_r \tau (n\bar{n})^2,$$

- where T_r is the **reduced lifetime** (in s^{-1} !).
- $N\bar{N}$ potential by Dover-Richard-Sainio (Khono-Weise, for instance, give similar results)

$$W(r) = -\frac{V_0 + i W_0}{1 + \exp[(r - R)/a]},$$

$$V_0 = W_0 = 0.5 \text{ GeV}, \quad a = 0.2 \text{ fm}, \quad R = 0.8 \text{ fm},$$

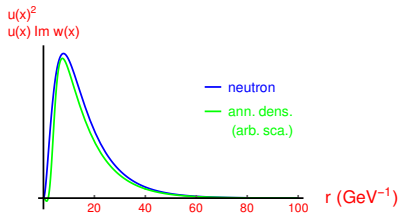
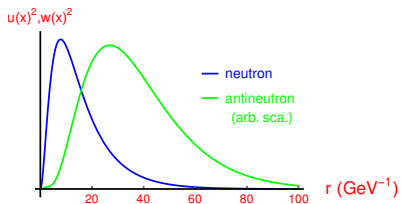
-

$$T_r \simeq 3 \cdot 10^{22} \text{ s}^{-1}.$$

Thus $T \gtrsim 10^{33} \text{ yr}$ for the deuteron $\Rightarrow \tau(n\bar{n}) \sim 10^9 \text{ s}$.

Lifetime of the deuteron-3

- Spatial extension of n , \bar{n} and annihilation density
 $\propto \gamma u(r) \text{Im } w(r)$.



Lifetime of the deuteron-4

- Alternative formula

$$T_R \approx \frac{\langle V_n - \text{Re } V_{\bar{n}} \rangle^2 + \langle \text{Im } V_{\bar{n}} \rangle^2}{-2 \langle \text{Im } V_{\bar{n}} \rangle},$$

- is not too bad, but not too good either
- does not distinguish inner from outer neutrons
- works in the limit of **deep** binding!
- underestimates the rate of decay, especially in case of weakly-bound external neutrons

Lifetime of ^{16}O

- As an example of medium-size nucleus proton-decay exp.
- See Dover, Gal, R., and Friedman & Gal, ...
- Shell-model with individual wave function for $S_{1/2}$, $P_{1/2}$, ... to reproduce the observed properties (mainly r.m.s.)
- Summarized as an **effective neutron potential** for each shell, $V_n = V(n - ^{15}\text{O})$
- While $V_{\bar{n}}$ taken from \bar{p} -nucleus phenomenology (exotic atoms, low-energy scattering)
- **Same inhomogeneous** eqn. as for deuteron, for each shell

Results for ^{16}O and ^{56}Fe

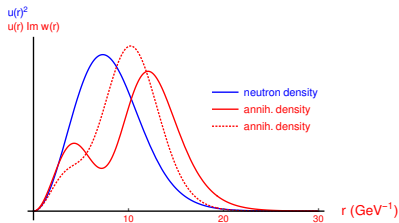
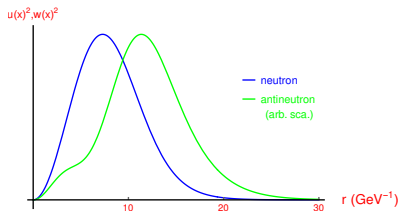
TABLE I. Reduced lifetime T_R (in units of 10^{23} sec^{-1}) for the neutrons in ^{16}O .

Orbit lj	$s_{1/2}$	$p_{3/2}$	$p_{1/2}$	Average
Model I (Ref. 18)	1.63	1.11	0.94	1.2
Model II (Ref. 19)	1.21	0.85	0.75	0.8

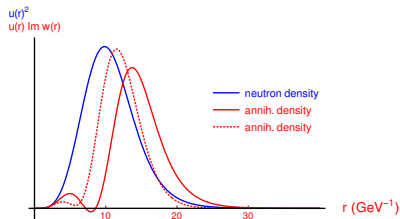
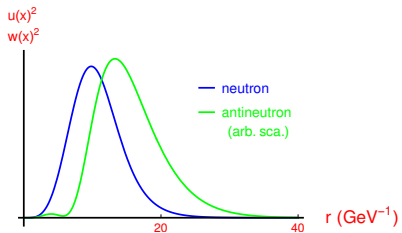
TABLE II. Reduced lifetime T_R (units 10^{23} sec^{-1}) for the neutrons in ^{56}Fe

Orbital lj	T_R (Model II)	T_R (Model I)
$s_{1/2}$	1.68	3.32
$p_{3/2}$	1.50	2.75
$p_{1/2}$	1.54	2.92
$d_{5/2}$	1.26	2.04
$2s_{1/2}$	1.09	1.60
$d_{3/2}$	1.33	2.29
$f_{7/2}$	0.98	1.34
$2p_{3/2}$	0.57	0.64
Average	1.13	1.69

Results for ^{16}O and ^{56}Fe



Results for ^{16}O and ^{56}Fe



Conclusions

- Oscillations **mainly outside**
- Subsequent annihilation **mainly at the surface**
- So minimal risk of dramatic medium renormalization of the basic process
- Good knowledge of the antinucleon-nucleus interaction in this region
- Nuclei with **neutron skin** or neutron halo favored
- $T_R \sim 10^{23} \text{ s}$ in $T(\text{nucleus}) = T_r \tau_{n\bar{n}}^2$