The Flavor Structure of the Nucleon Sea, INT Workshop, Seattle October 2<sup>nd</sup> - 17<sup>th</sup>, 2017



#### Probing the sea quark content of the proton with SIDIS data

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in collaboration with I. Borsa and M. Stratmann 1708.01630 Motivation:

significant progress in final state hadron measurements



 $\frac{d\sigma}{dz}(e^+e^- \longrightarrow hX) = c_q(z) \otimes [D_q(z,Q^2) + D_{\overline{q}}(z,Q^2)] + c_g(z) \otimes D_g(z,Q^2)$ 

Motivation:

significant progress in final state hadron measurements





$$M_{e,p(d)}^{\pi^{\pm}} \equiv \frac{d\sigma^{\pi^{\pm}}/dx \, dQ^2 \, dz}{d\sigma/dx \, dQ^2}$$

$$\frac{d\sigma^{\pi}}{dxdzdQ^2} = c_{if}(x,z) \otimes f_i(x,Q^2) \otimes D_f^{\pi}(z,Q^2)$$
$$\frac{d\sigma}{dxdQ^2} = c_i(x) \otimes f_i(x,Q^2)$$

Motivation:

significant progress in final state hadron measurements



 $\frac{d\sigma}{d\eta dp_T} \sim f_i(x_1, p_T^2) \otimes f_j(x_2, p_T^2) \otimes C_{ijk} \otimes D_k^h(z, p_T^2)$ 







## a case for a combined PDFs and FFs analysis?

Quark elastic scattering as a source of high-transverse-momentum mesor flavor separation without nuclear targets? PHYSICAL REVIEW D

isospin symmetry?

nuclear effects in d?

nuclear effects in A?

QUARK ELASTIC SCATTERING AS A SOURCE OF ... the original momentum of the quark. This is given, in principle, by the hadrons produced by the rein principle, by the national produced by the re-colling quarks in deep-inelastic neutrino proton colling quarks in deep-menance neutrino pround scattering. Unfortunately, in both cases the data are incomplete and must be supplemented by the are incomplete and must be appremented by the oretical arguments that require much discussion. This first paper deals primarily with these func-(in fact, of the same flavor as the quarks that un sact, of the same there as the quarks that came in), which fragment or cascade down into several hadrons.<sup>1</sup> This is illustrated in Fig. 1. we disregard the theoretical argument that this 12 elastic cross section [which we write as d5/ exactle errors section (which we write as  $nT/d\hat{t}(\hat{s},\hat{t})$ , where  $\hat{s}$  and  $\hat{t}$  are the s,t invariants for the quark collision) must vary as  $3^{-2} \int (l/3)$  and, instead leaves the state terms of ter instead, leave it as an unknown function to be determined empirically by the data. It will vary Nore like s -f(t/s) with N about 4. We shall need the distributions  $G_{A \to t}(x)$  of quarks more like  $s^{-N}f(\hat{t}/\hat{s})$  with N about 4. we shall need the distributions  $O_{A \to A}(V)$  or quark q in the initial hadrons; for protons and neutrons this is sized to a large extent by descriptions d in the initial hadrons; for protons and neutron this is given to a large extent by deep-inelastic this is fiven to a targe extent by deep-inelastic ep (or  $\mu p$ ) scattering data. Also, we shall need to of (or  $\mu p$ ) scattering data. Also, we mail need to know what the chances,  $D_{4}^{h}(z)$ , are that a quark q going out at large momentum disintegrates into going out at large momentum unsumerance and various kinds of hadrons, k, with a fraction z of

(0)

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10+b-c+d)

ment from quarks. These effects have intue and fluence on the results of the present paper and we thence on me results or the present paper and thave omitted them in our calculations reported We are fully aware that all partons are not quarks, that half the momentum of the proton is something else (gluons?). And there is no good reason to exclude the possibility that some of the high-p, particles could result from gluon interhere. We are also aware that there is no good reason for the quark-quark cross section to var chosen to start here. Let us see what experin might exclude our specific choice, and indicat the presence of glaces, or some different me Before we begin, bowever, we must say of time in what region we expect our theor We must be careful, because we do not wi entirely. embarrassed later by appearing to think case as to why something does not fit. to allow, generally, any data outside the

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tions  $G_{A \rightarrow A}(x)$  and  $D_{A}^{A}(z)$  and with the behavior of entropy matrix matrix  $G_{A \rightarrow A}(x)$  and  $D_{A}^{A}(z)$ .

outgoing particle and incoming beam ratios for

various forms for  $d\partial/d\hat{i}$  and make predictions that

various forms for an/ar and make Predictions mat are insensitive to its detailed form. Since the be-

havior of  $G_{3 \rightarrow 4}(x)$  and  $D_{4}^{h}(z)$  is inferred from lepton-

hadron and lepton-lepton processes, much of this

dict properties of badron-badron collisions from

information gained studying lepton-initiated re-

perimental quantities that depend more strongly

bermanial quantities that append more strongly on the precise form of d3/d2 (e.2., two-particle

on me precise form of *d9/dE* (e.E., two-particle correlation data in large-*p*<sub>1</sub> hadron collisions). Then it will be necessary to include the effects of

the transverse momentum spread of the quarks within the hadrons and of the hadrons that frag-

These effects have little in-

A subsequent paper will investigate ex-

tirst paper can be viewed as an attempt to pre-

single-particle production. We examine

## are PDFs that good?

![](_page_7_Figure_1.jpeg)

![](_page_7_Figure_2.jpeg)

#### are PDFs that good?

![](_page_8_Figure_1.jpeg)

![](_page_8_Figure_2.jpeg)

number/type data number parameter/unknowns topography contamination

PDFs reweighting: 1012.0836

 $M_{ep}^{\pi^+}(z,Q^2)$  $M_{ed}^{\pi^{+}}(z,Q^{2})$ not fitted  $0.2 \le z \le 0.3$ 0.2 ≤ z ≤ 0.3 0.3 ≤ z ≤ 0.4  $0.3 \le z \le 0.4$ 0.4 ≤ z ≤ 0.6  $0.4 \le z \le 0.6$  $0.6 \le z \le 0.8$  $0.6 \le z \le 0.8$  $M_{ep}^{\pi}(z,Q^2)$  $M_{ed}^{\pi^-}(z,Q^2)$ 0.2 ≤ z ≤ 0.3  $0.2 \le z \le 0.3$ 0.3 ≤ z ≤ 0.4 0.3 ≤ z ≤ 0.4  $0.4 \le z \le 0.6$ 0.4 ≤ z ≤ 0.6  $0.6 \le z \le 0.8$  $0.6 \le z \le 0.8$ 10 THIS FIT with 68 and 90% C.L. bands DSS <sup>o</sup> HERMES  $\pi^+$  multiplicities  $O^2$  [GeV] O<sup>2</sup> [GeV]  $f_i(x) \quad w(\chi^2)$   $f_i(x) \quad w(\chi^2)$   $f_i(x) \quad w(\chi^2)$  $f_i^{best}(x) = \frac{1}{N_{rep}} \sum f_i(x)$  $f_i^{reweight}(x) = \frac{1}{N_{rep}} \sum w f_i(x)$ 

 $\chi^2(\mathrm{FFs}(\mathrm{PDFs}))$ 

 $1000 \ times$ 

number/type data number parameter/unknowns topography contamination

![](_page_10_Picture_2.jpeg)

iterative FFs & PDFs determination:

![](_page_10_Figure_4.jpeg)

![](_page_11_Figure_1.jpeg)

#### similar results with CT14 replicas

![](_page_12_Figure_1.jpeg)

![](_page_13_Figure_1.jpeg)

# optimized FFs

![](_page_14_Figure_1.jpeg)

## reweighting esoterica:

Distribution of weights  $w_k$ Distribution of  $\chi^2/N_{data}$ 100 600 consistency checks: 68% CL FFs error 50% CL FFs error 75 no FFs error 400 **#** surviving replicas 50 distribution of w 200 25 0 0 10 -22 10 -12 10 -7 10<sup>-2</sup> 10 -17 10<sup>2</sup> 3 2 4 0 1 0.3 300 Reweighted Distribution of  $\chi^2/N_{data}$ Probability distribution  $P(\alpha)$ double counting 0.2 200 0.1 100 0 0 2 3 0 1 4 2 1  $\chi^2$  /  $N_{data}$ α

## SIDIS revisited:

![](_page_16_Figure_1.jpeg)

### SIDIS revisited:

![](_page_17_Figure_1.jpeg)

## SIDIS revisited:

![](_page_18_Figure_1.jpeg)

![](_page_19_Picture_0.jpeg)

FFs are coming on age: from rough pictures to precision tools

not yet as precise as current PDFs, but still can be useful

as **fundamental** as PDFs in the pQCD description

combined analysis PDFs & FFs works

key to **flavor separation** without nuclear targets?