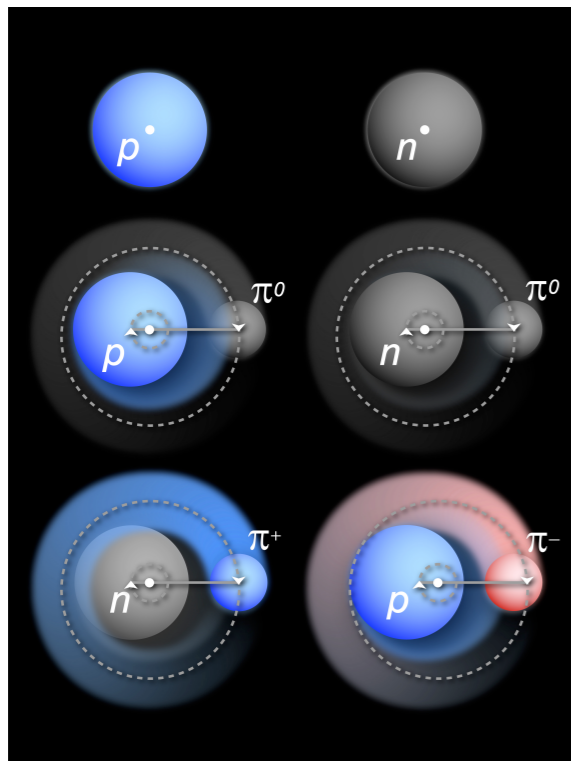


Taming the pion cloud

Gerald A. Miller, UW
Mary Alberg, Seattle U, UW

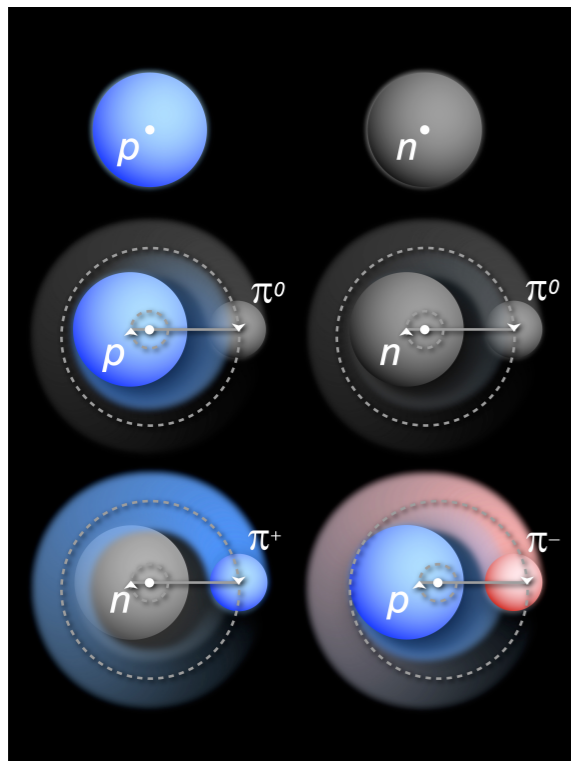


VS



Taming the pion cloud

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Mary Alberg, Seattle U, UW



VS

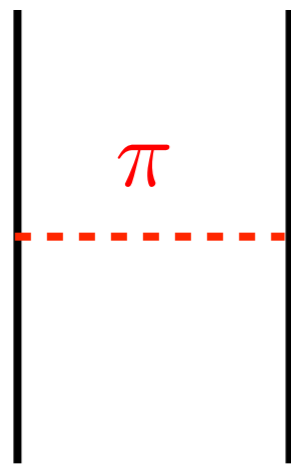


Ancient History

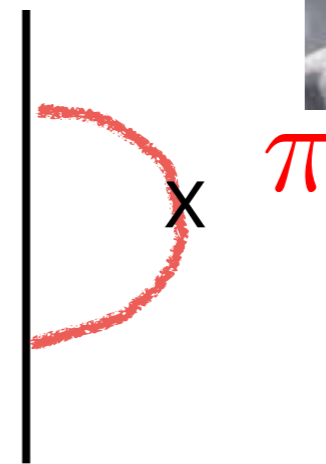
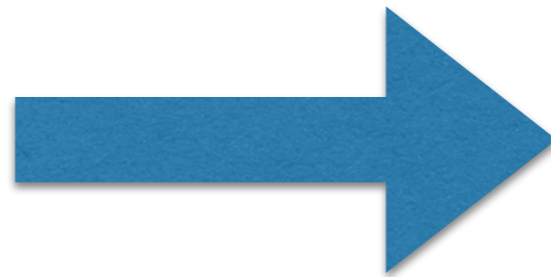


H. Yukawa

NN force

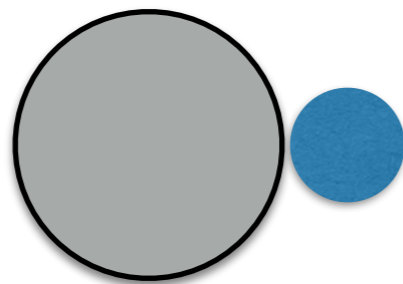


Longest range component



Pion cloud

Fundamental Question: meson cloud or $4q\bar{q}$



Hidden color

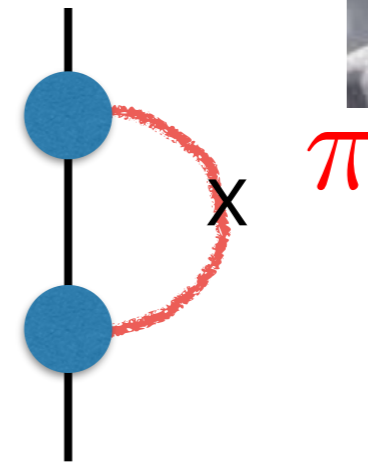
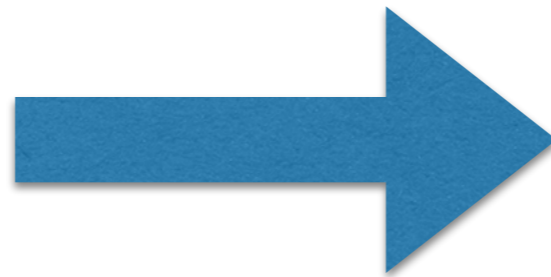
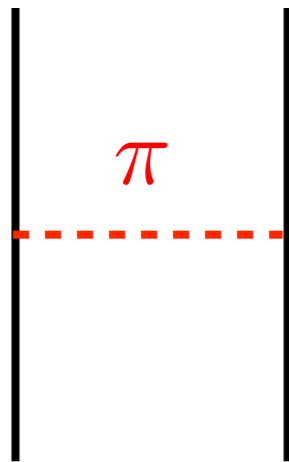
Important implications for nuclear force and nuclear structure if meson cloud picture is shown to fail

Ancient History



H. Yukawa

NN force



Pion cloud

Longest range component

Implications for sea

$$p \rightarrow n\pi^+, p \rightarrow \Delta^{++}\pi^-, \dots$$

$$\pi^+ \sim u\bar{d}, \pi^- \sim d\bar{u}$$

$$\text{Thus } \bar{d}(x) \neq \bar{u}(x)$$

A W Thomas PLB 126, 97 (1983)

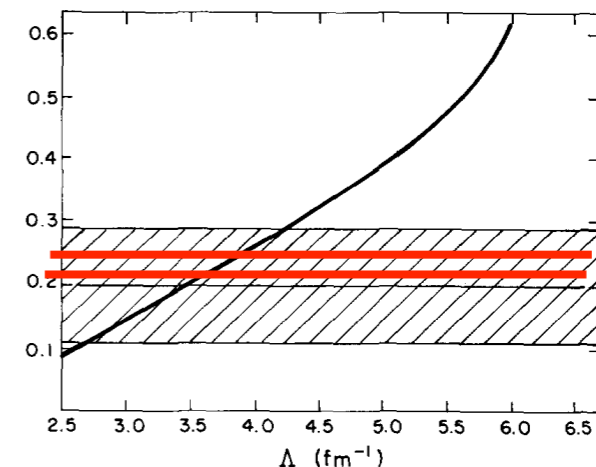
Wandmolders et al PRL 66,2712 (1991)

$$D \equiv \int_0^1 [\bar{d}(x) - \bar{u}(x)] dx = 0.136 \pm 0.060 \quad \mathbf{0.16 \pm 0.01}$$

Henley & Miller PLB 251, 453 (1990)

 Form factor $\exp(-k^2/\Lambda^2)$

$\frac{3}{2}D$

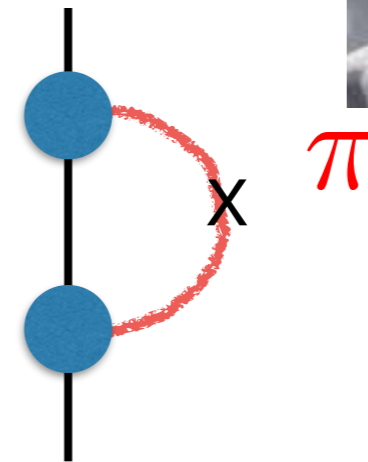
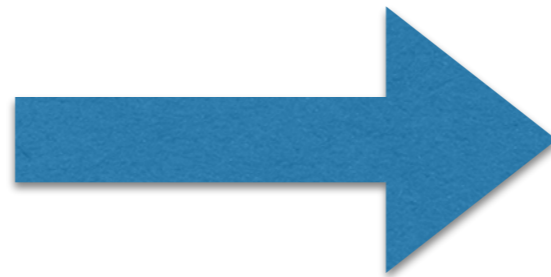
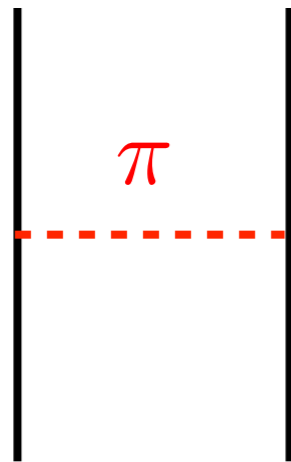


Ancient History



H. Yukawa

NN force



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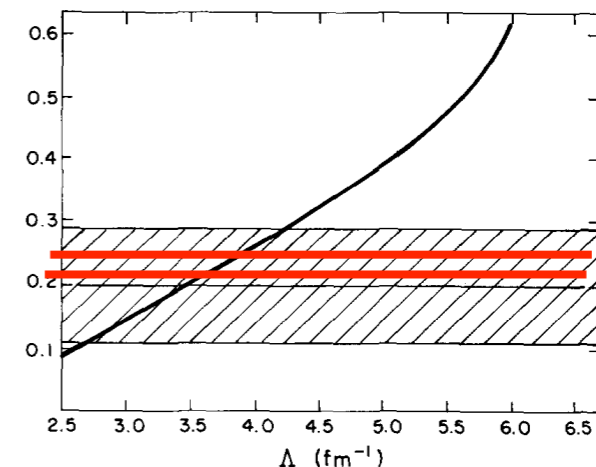
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$$\Lambda \rightarrow R_{\pi N} = R_{\pi N}^{\text{experiment}} = 0.66 \text{ fm}$$

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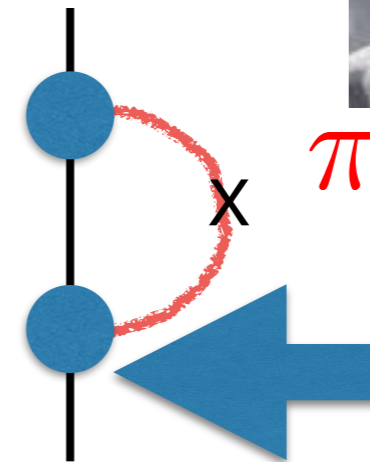
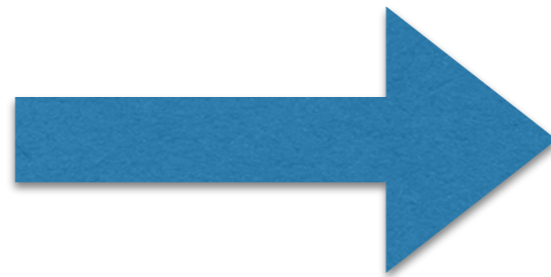
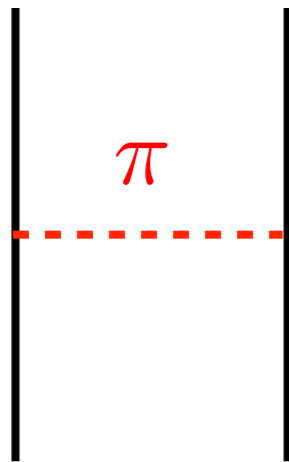


Ancient History



H. Yukawa

NN force



Must be tamed

Pion cloud

Longest range component

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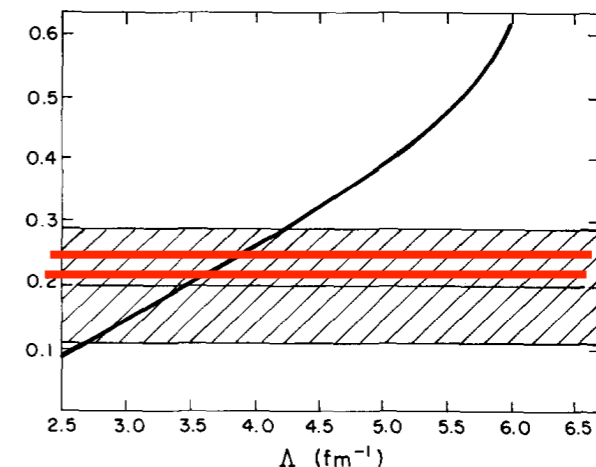
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$\frac{3}{2}D$



Experimental Progress Drell-Yan

- NMC measured integral quantity for Gottfried sum
- E866 FermiLab measured x-dependence

J. C. PENG *et al.* PHYSICAL REVIEW D **58** 092004

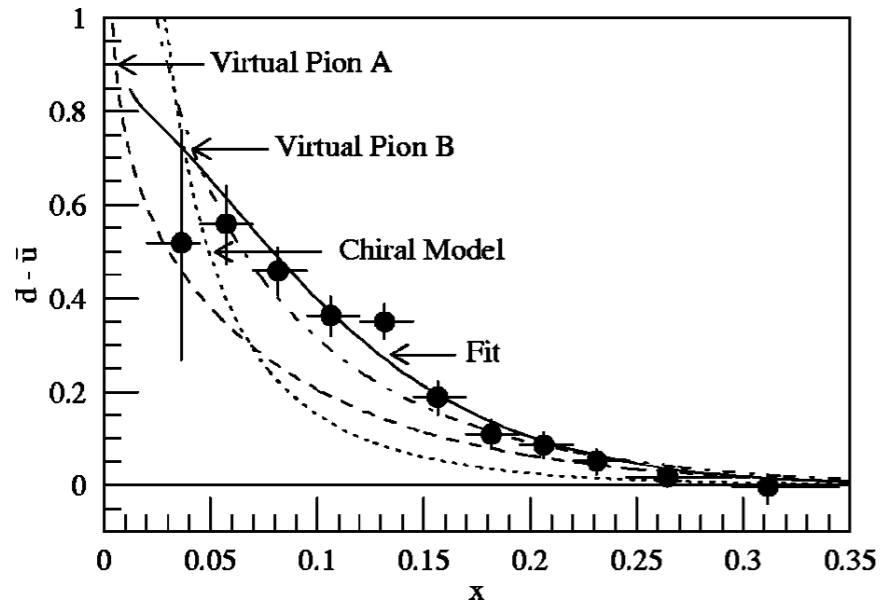
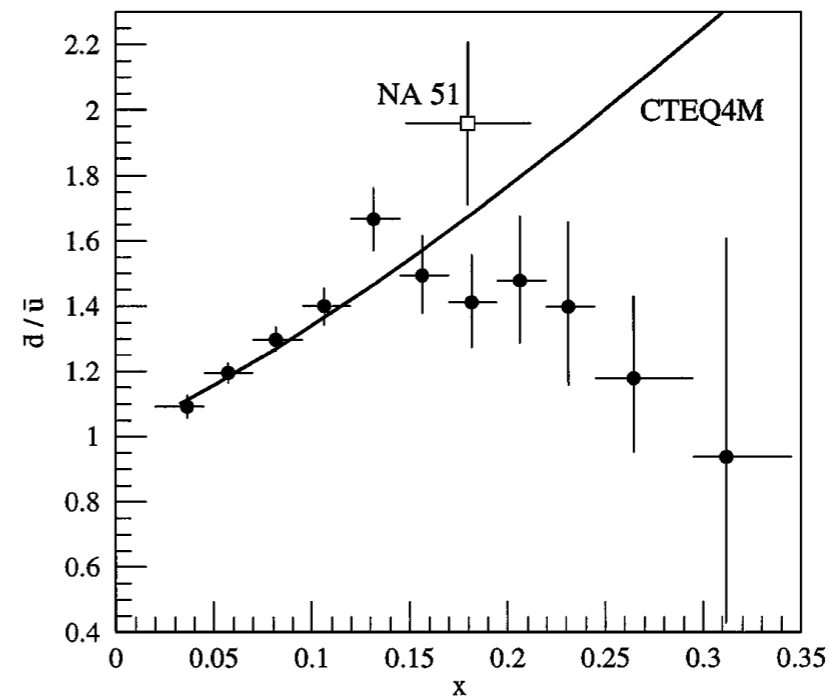


FIG. 1. Comparison of the E866 $\bar{d}-\bar{u}$ results at $Q=7.35$ GeV with the predictions of various models as described in the text.

$\bar{d}(x) > \bar{u}(x)$ what about \bar{d}/\bar{u} ?

Expect large ratio at large x

Hawker et al PRL 80, 3715



More data- E866 (1999)

Drell-Yan Measured $\bar{d} - \bar{u}$ **and** $\frac{\bar{d}}{\bar{u}}$

Alberg, Henley and Miller PLB 471, 396 (2000)

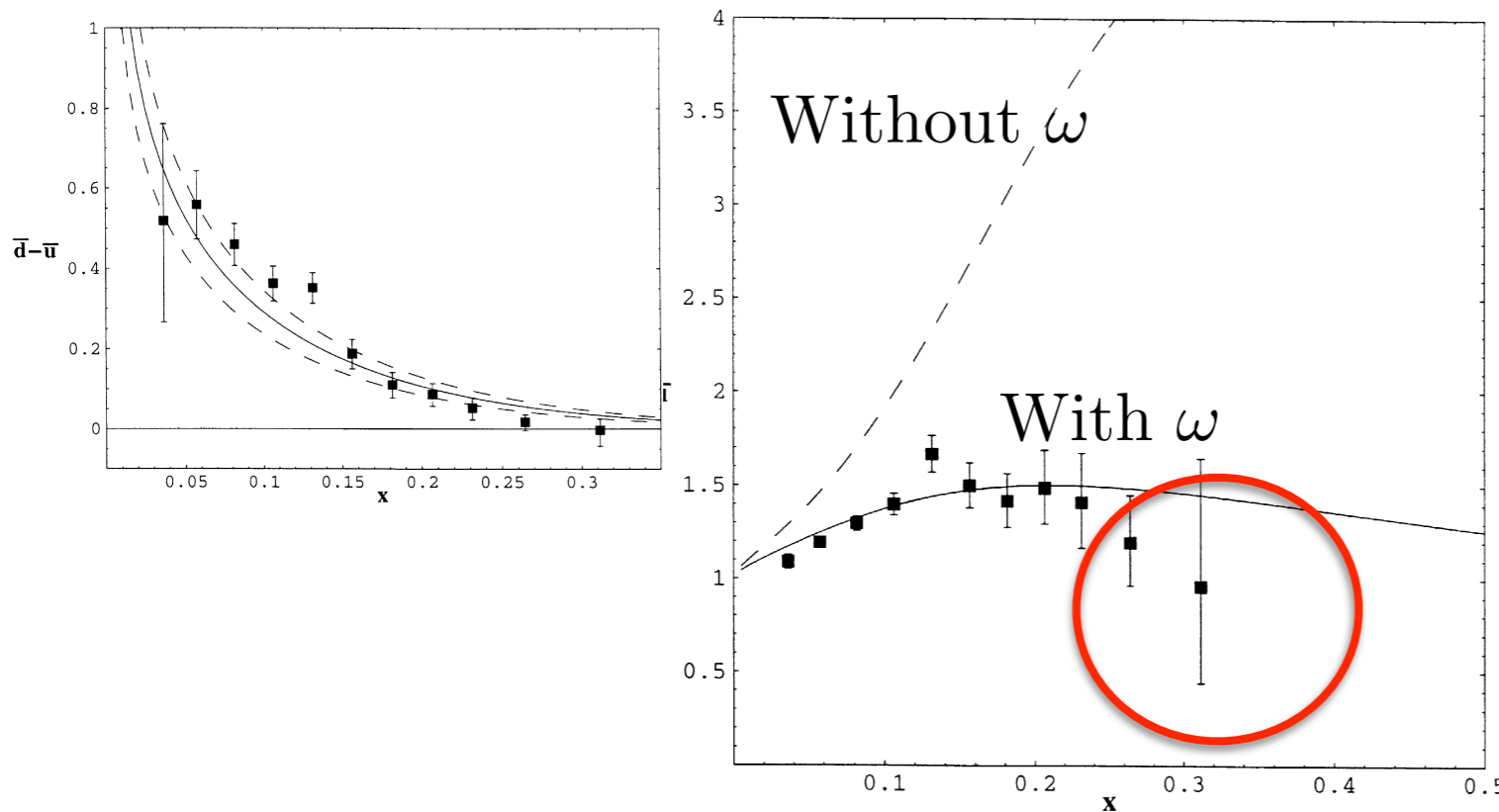
With pions get too large a ratio $\frac{\bar{d}}{\bar{u}}$

Form factor

$$G_M(t, u) = \exp\left(\frac{t - m_M^2}{2 \Lambda_M^2}\right) \exp\left(\frac{u - m_B^2}{2 \Lambda_M^2}\right),$$

J. Speth, A.W. Thomas, Advances in Nuclear Physics, vol. 24, J.W. Negele, E.W. Vogt (Eds.), Plenum Press, New

Is there an isoscalar non-perturbative sea (omega meson)?



The omega represents any non-perturbative isoscalar sea

What's going on at high x?

SeaQuest aims at better measurement, so we try to improve

Theory problems

- Results depend on form factor parameter Λ
- form factors enter as three dimensional functions even though expressed in terms of t and u
- how to derive ????
- Why do we need form factors? **Form factors oppose chiral perturbation theory**

Why do we need form factors?

Using form factors opposes chiral perturbation theory

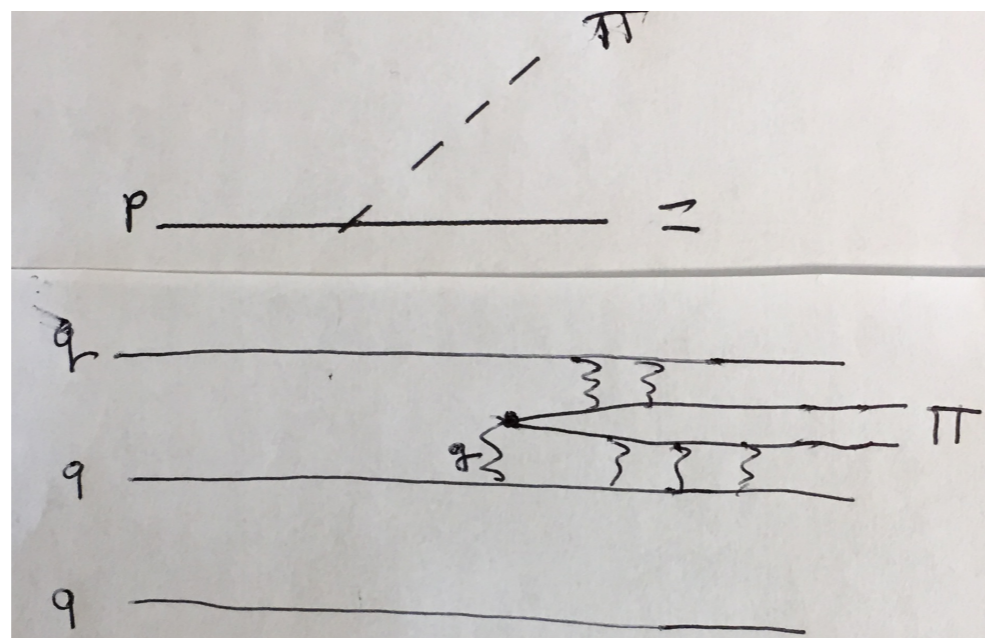
$$\mathcal{L}_N^{(1)} = \bar{\psi}(i\gamma \cdot \partial - M)\psi - \frac{g_A}{2f_\pi} \bar{\psi} \gamma_\mu \gamma_5 \tau^a \psi \partial_\mu \pi^a - \frac{1}{f_\pi^2} \bar{\psi} \gamma_\mu \tau^a \psi \epsilon^{abc} \pi^b \partial_\mu \pi^c$$

Non-renormalizable \mathcal{L} , expand in powers of momentum
add counter terms order-by order -LECs.

results **INDEPENDENT OF CUTOFF** Not sufficient for DIS, IMHO because momenta $\mathcal{O} \sim m_N$

Form factor relates the LECs in a very specific way
different philosophy

Comment talk last week titled "... with chiral perturbation theory"
is **NOT** – no LEC's, but yes to cutoff dependence



pion-nucleon
Form factor takes composite
nature of pion and nucleon
into account

Taming I -Alberg & Miller

PRL 108 (2012) 172001

Constrain form factor using experimental input info from Thomas and Weise book

$$q_\mu \langle p(P') | A_+^\mu(0) | n(P) \rangle = 2\bar{u}_p(P') \left[M G_A(Q^2) - \frac{Q^2 f_\pi g_{\pi NN}(Q^2)}{Q^2 + m_\pi^2} \right] \gamma_5 u_n(P).$$

Thus the matrix element of the divergence of the axial current vanishes as $m_\pi^2 \rightarrow 0$ if $G_A(Q^2)$ and the pion-nucleon form factor $g_{\pi NN}(Q^2)$ are related by

$$M G_A(Q^2) = f_\pi g_{\pi NN}(Q^2). \quad (3.10)$$

At $Q^2 = 0$ this is known as the Goldberger-Treiman relation and it is satisfied at the level of 3 % (with $g_A = G_A(0) = 1.267 \pm 0.004$, $g_{\pi NN} = g_{\pi NN}(0) = 13.2 \pm 0.1$ and $M =$

$$G_A(Q^2) = \frac{G_A(0)}{\left(1 + \frac{Q^2}{M_A^2}\right)^2},$$

Exp.	M_A [GeV]
BNL	1.07 ± 0.06
ZGS	1.00 ± 0.05
Fermilab	$1.05 \pm_{-0.16}^{+0.12}$
Average:	1.03 ± 0.04

Pion-Nucleon form factor determined for on-mass-shell nucleons, off shell pion

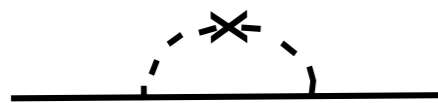
Taming I -Alberg & Miller

PRL 108 (2012) 172001

$$g_{\pi NN}(Q^2) \propto G_A(Q^2) = \frac{G_A(0)}{\left(1 + \frac{Q^2}{M_A^2}\right)^2}$$

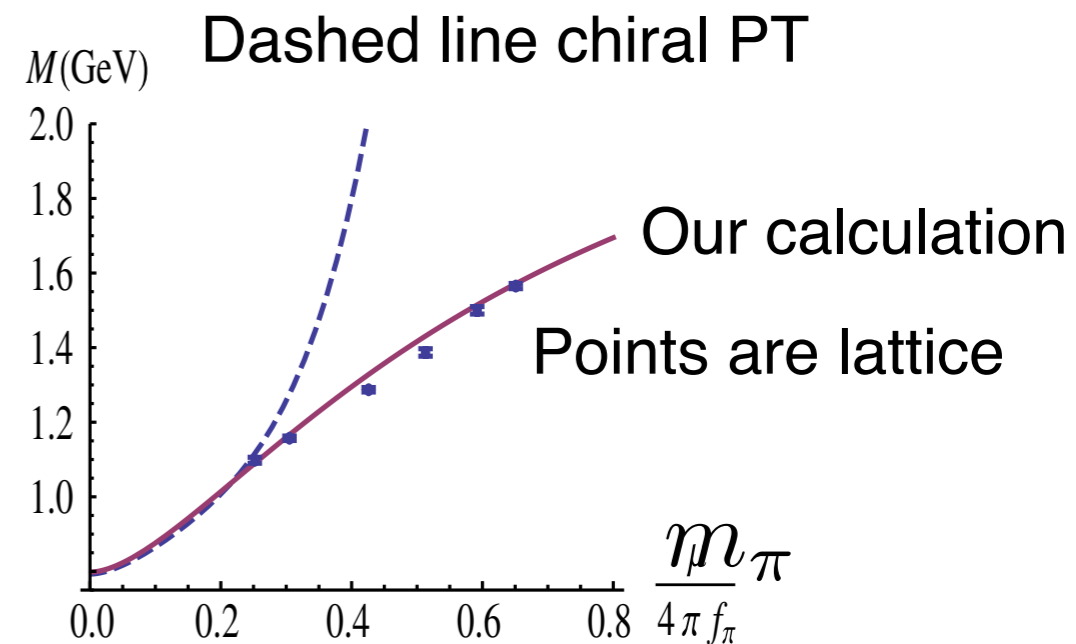
Pion-Nucleon form factor determined for on-mass-shell nucleons, off shell pion

Nucleon self energy -intermediate nucleon and Delta



$$D \equiv \int_0^1 [\bar{d}(x) - \bar{u}(x)] dx = 0.118 \pm 0.012.$$

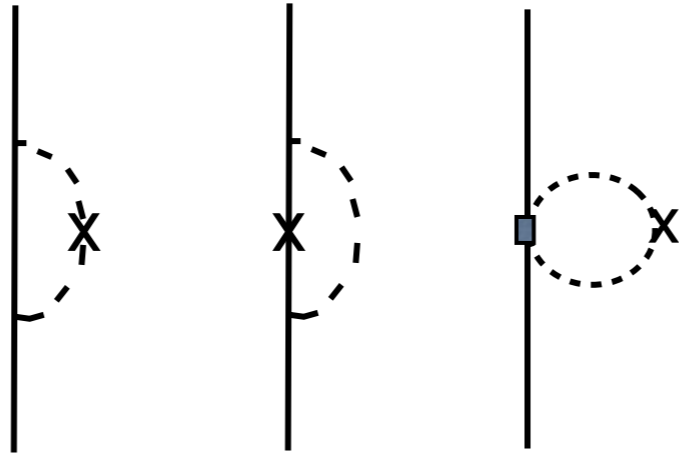
Ours 0.109



Parameter free calculation!
So what's the problem?

Taming II- recent

To get \bar{u} and \bar{d} need to calculate the graphs:



Both pion and nucleon are off-shell in the Feynman graphs
need to reconsider the formalism

$$q_N^f(x) = Z_2 q_{N0}^f(x) + \sum_{B,M} \int_x^1 \frac{dy}{y} f_{MB}(y) q_M^f\left(\frac{x}{y}\right) + \sum_{B,M} \int_x^1 \frac{dy}{y} f_{BM}(y) q_B^f\left(\frac{x}{y}\right)$$

$$Z_2^{-1} - 1 = \sum_{B,M} \int dy f_{BM}(y),$$

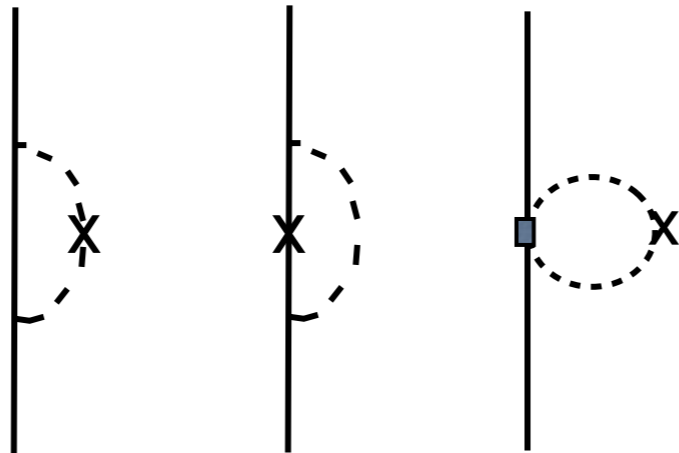
Brodsky-Lepage Fock space representation:

$$|\pi N\rangle \propto \int_0^1 \frac{dy}{\sqrt{y}} d^2 k_{\perp\pi} \int_0^1 \frac{dy_N}{\sqrt{y_N}} d^2 k_{\perp N} \delta(1-y-y_N) \delta(\vec{k}_{\perp\pi} + \vec{k}_{\perp N}) \psi_{\pi N}(y, \vec{k}_{\perp\pi}; y_N, \vec{k}_{\perp N}) |\cdots\rangle$$

$$f_{\pi N}(y) = \int d^2 k_{\perp\pi} \left| \psi_{\pi N}(y, \vec{k}_{\perp\pi}; 1-y, -\vec{k}_{\perp\pi}) \right|^2 = f_{N\pi}(1-y)$$

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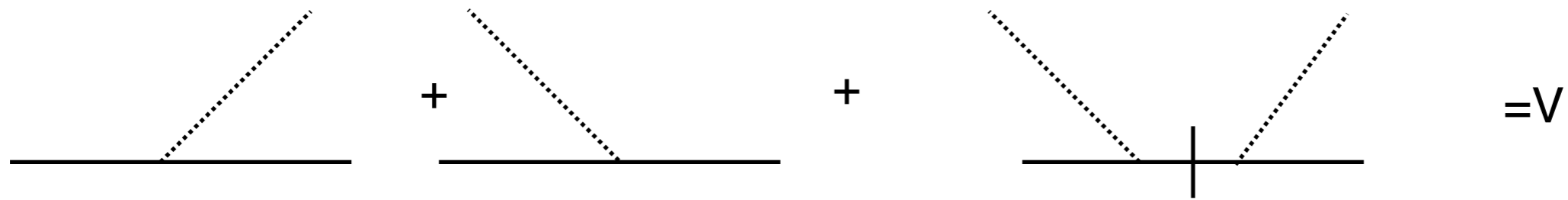
If f_{MB} has $\delta(y)$
 Z_2 would change,
but **NO** delta functions here!

Brodsky-Lepage Fock space representation:

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\hat{P}^- is Hamiltonian operator, construct from energy-momentum tensor $T^{+-} =$
 free particle kinetic energy M_0^2 plus interactions V



Schroedinger eq: $(\hat{P}^- - \hat{P}^- - \hat{P}_\perp^2)|p\rangle = M_p^2|p\rangle = (M_0^2 + V)|p\rangle$

$|p\rangle \approx Z_2 \left(|p\rangle_0 + \frac{1}{M^2 - M_0^2} V |p\rangle_0 \right)$

$|\pi N\rangle$ component

$$\mathcal{L}_N^{(1)} = \bar{\psi}(i\gamma \cdot \partial - M)\psi - \frac{g_A}{2f_\pi} \bar{\psi}\gamma_\mu\gamma_5\tau^a\psi \partial_\mu\pi^a - \frac{1}{f_\pi^2} \bar{\psi}\gamma_\mu\tau^a\psi \epsilon^{abc}\pi^b\partial_\mu\pi^c$$

Form factors absent

Form factors

- Including form factors goes beyond usual LF treatment
- Need form factors in frame independent manner (4-space)
- Maintain momentum conservation, **unique** LF wave function
- **Keep experimental input**
- For use in light front wave function-virtual N, π

- Product of pion and nucleon form factors:

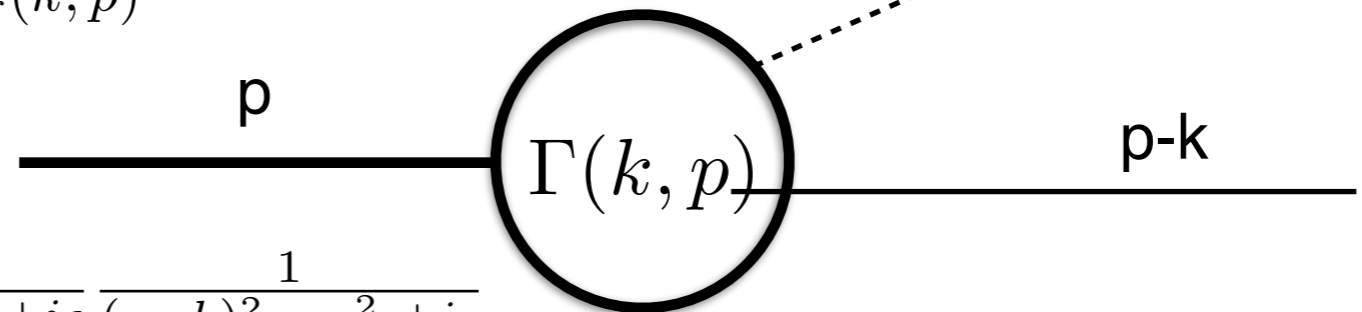
~Pauli-Villars reg.

$n = \text{integer}$

$$F(k, p, y) = \frac{1}{\left[1 - \frac{(k^2 - m_\pi^2)}{\Lambda^2}\right]^n} \frac{1}{\left[1 - \frac{y}{1-y} \frac{(p-k)^2 - m_N^2}{\Lambda^2}\right]^n}$$

$$\psi_{LF} \propto \int_{-\infty}^{\infty} dk^- \psi_{\text{Bethe-Salpeter}}(k, p)$$

$$\psi_{\text{Bethe-Salpeter}}(k, p) =$$



$$\psi_{LF} \propto \int dk^- \Gamma(p, k) \frac{1}{k^2 - m_\pi^2 + i\epsilon} \frac{1}{(p-k)^2 - m_N^2 + i\epsilon}$$

Γ contains form factor.

Integrate over UH k^- plane = integrate over LH k^- plane w. stated form factor

$n = 1$ gives form factor very close to dipole, maintain experimental input!

Summary

- Have formalism to get light front wave functions and meson distribution functions needed for light flavor nucleon sea
- Meson-nucleon coupling constants are known
- Form factors included in frame independent manner that incorporates experimental input
- Given the meson cloud model can make calculations with reasonably-well understood uncertainties
- True test of meson cloud model!
- See Alberg's talk