## **Hot QCD and Quantum Simulations**

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The system created at RHIC behaves like perfect liquid (2005) How does the system thermalize ?

1

Is there is a critical point on the QCD phase diagram ? (2019-2021)

INT, Novemver 14-15, 2017

### Viscous hydrodynamics and flow

Assume that a thermal system is created shortly after the collisions that expands hydrodynamically.

To describe the experimental data very small shear viscosity to entropy ratio is needed

RHIC Scientists Serve Up "Perfect" Liquid, New state of matter more remarkable than predicted -raising many new questions April 18, 2005







How small is the shear viscosity ?

Validity of the hydrodynamics is governed by *η/s* Hadron gas and QGP at very high temperature have large value *η/s*

Super-symmetric gauge theories at strong coupling have small *η/s* with lower bound dictated by quantum mechanics  $\eta/s > 1/(4 \pi)$  (Kovtun, Son Starinets 2005)  $\Rightarrow$  QGP near the transition temperature *T<sub>c</sub>* has close to minimal  $\eta$ /s



## Initial time dynamics and thermalization in heavy ion collisions

Classical-statistical calculations of gluon distribution at early times (large gluon occupation numbers)

Berges Schenke, Schlichting, Venugopalan, Nucl. Phys. A 931 (2014) 348



The gluon occupation number decreases at later times reaching *O(1)*, the system becomes quantum and strongly coupled

quantum simulations are needed

Early time dynamics is important event-by-event fluctuations in *AA*, and high multiplicity *pA* and *AA* collisions

# Strongly coupled QGP and heavy quarks

Heavy quarks *(* $M_c \sim 1.5 \text{ GeV}$ *)* flow in the strongly coupled QGP





$$
\frac{p(T,\mu)}{T^4} = \sum_{n=1}^{\infty} \frac{1}{(2n)!} \chi_{2n}(T) \mu^{2n}
$$

Calculable in LQCD but the computational difficulty increases with *n* ! (noise problem vs. sign problem) Current calculations exist only to  $n=6$ . 6 LQCD Accomplishments: The Transition in QCD

The QCD transition at zero net baryon density is a smooth crossover (not a phase transition) the cross-over temperature has been determined in the continuum limit

Fluctuations of the order parameter:  $\chi_{disc} = VT^{-1} \left( \langle (\bar{\psi}\psi)^2 \rangle - \langle \bar{\psi}\psi \rangle^2 \right)$ 

Bhattacharya et al (HotQCD), PRL 113 (2014)082001



#### Bazavov et al, PRD85 (2012) 054503  $T_c = (154 \pm 8 \pm 1 (scale))$ MeV



no increase with the volume

 Crossover transition Aoki et al, Nature 443 (2006) 675 <sup>7</sup> Equation of state in the continuum limit

Equation of state has be calculated in the continuum limit up to *T=400* MeV

 $p \sim \rho T$ ,  $\rho \sim T^3$ in ultra-relativistic case



HRG agrees with the lattice for *T< 145* MeV

Calculations that use two different discretization schemes agree:

#### Bazavov et al, PRD 90 (2014) 094503



Thermodynamics at non-zero net baryon density

6 th order Taylor expansion, Bazavov et al, PRD 95 (2017) 054504



Truncation errors of the 6<sup>th</sup> order Taylor expansions are small for *μ<sub>B</sub>/T<2.5* 

Critical point is strongly disfavored for  $\mu_B$ /T<2.0

## Correlation functions and transport coeffocients

Transport coefficients are encoded in the spectral functions:

$$
\sigma(\omega, p, T) = \frac{1}{2\pi} \operatorname{Im} \int_{-\infty}^{\infty} dt e^{i\omega t} \int d^3x e^{ipx} \langle [J(x, t), J(x, 0)] \rangle_T
$$

In LQCD one can calculated the Euclidean time

transport coefficients =  $\lim_{\omega \to 0} \frac{\sigma(\omega)}{\omega}$ 

10

$$
G(\tau,p,T)=\int d^3x e^{ipx}\langle J(x,-i\tau),J(x,0)\rangle_T
$$

Due to analytic continuation

$$
G(\tau,T)=D^{>}(-i\tau)
$$



## Heavy quark diffusion constant from quenched LQCD

Direct method: determine the width of the transport peak, Ding et al, arXiv:1204:4954, quenched *128*×*N<sup>τ</sup>* lattices, *Nτ=24-48*

Integrate out the heavy quark fields:  $\langle J_i(\tau) J_i(0) \rangle \Rightarrow \langle E_i^a(\tau) E_i^a(0) \rangle$ Banarjee et al, arXiv:1109.5738, Kaczmarek et al, arXiv:1109:3941, *Nτ=16-24*





the width of the transport peak is potentially overestimated



- How does the system created at RHIC thermalizes at time scales < 1fm/c ? What is the effect of the early time out of equilibrium stage on experimental observables?
- What is the shear viscosity and heavy quark diffusion coefficient as function of temperature
- Where is the critical point of the QCD phase diagram ? Or is the QCD transition is always a crossover ?
- $\Rightarrow$  Quantum Simulations of QCD !

# Hamiltonian formulation of gauge theories

Consider continuous time and discretized space :

$$
H = H^G + H^M
$$

$$
H^G = \frac{g^2}{2} \sum_{n,k} L_{n,k}^2 - \frac{1}{2g^2} \sum_{plaq} \text{tr}\left(U_1 U_2 U_3^\dagger U_4^\dagger + h.c.\right) \qquad L_{\mathbf{n},k} = -i \dot{U}_{\mathbf{n},k} U_{\mathbf{n},k}^\dagger
$$

$$
H^M = \sum_{n,k} (\psi_n^{\dagger} U_{n,k} \psi_{n+k} + h.c.) + m \sum_n (-1)^{\sum_k n_k} \psi_n^{\dagger} \psi_n^{\dagger} \psi_n
$$

Simulate using interacting cold atom systems in optical lattices:

Zohar, Cirac, Reznik, Rep. Prog. Phys. 79 (2016) 014401



# Realizations of gauge theories on optical lattices and challenges

Concrete proposal for quantum simulations:

- $Z_N$  gauge theories in 2+1 dimensions: Zohar, Brunello, PRD 91 (2015) 054506, Zohar, Farace, Reznik, Cirac, PRL 118 (2017) 070501, PRA 95 (2017) 023604
- U(1) Higgs model in 2+1 dimensions: Gonzales-Cuadra, Zohar, Cirac, New. J. Phys. 19 (2017) 063038

Theoretical challenges: Finite Hilbert space, continuum limit ? Higher dimensions ?

Experimental:

- 1) Generating complex lattices and superlattice structures
- 2) Better control of scattering parameters
- 3) Longer lived lattices
- 4) Improved measurement techniques

# **Summary**

- The are compelling questions in hot QCD that require quantum computations:
- 1) What is the QCD phase diagram at high baryon density ? Is there critical point ?
- 2) How does thermalization in ultra-relativistic heavy ion collisions happen ?
- 3) What are the QCD transport coefficients ?
- Quantum simulations using optical lattices may provide an avenue addressing these questions but many open challenges remain