

Vector meson production in ultraperipheral collisions: accessing the small- x gluon

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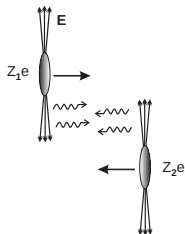
Probing QCD in Photon-Nucleus Interactions at RHIC and LHC:
the Path to EIC, February 14, 2017

- 1 Background
- 2 Pb+Pb collisions
- 3 p+Pb collisions

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Ultraperipheral collision as a deep inelastic scattering

$b \gtrsim 2R_A$: strong interactions are suppressed



J. Nystrand et al, nucl-ex/0502005

- 1 Nucleus creates a (real) photon flux $n(\omega)$
- 2 Photon-nucleus scattering

$$\sigma^{AA \rightarrow AA+V} \sim n(\omega) \sigma^{\gamma A \rightarrow VA}(\omega)$$

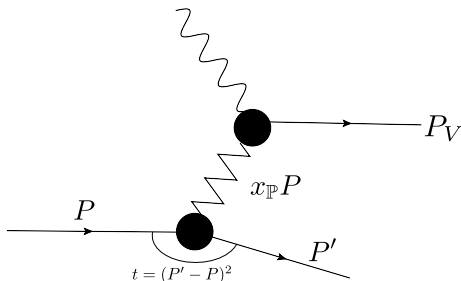
Interesting QCD part: high-energy γ -nucleus or γ -proton scattering.

Diffraction production of vector meson as a probe of small x

Exclusive production of vector meson

$$\gamma p \rightarrow Vp$$

$$\gamma A \rightarrow VA$$

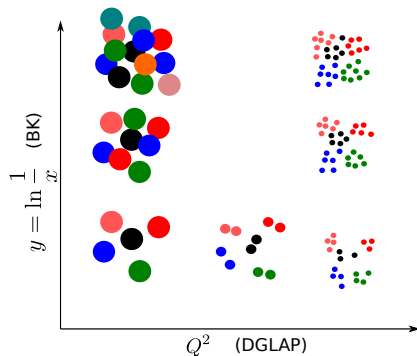


Pocket formula for diffraction (2-gluon exchange)

$$\frac{d\sigma^{\gamma^* H \rightarrow VH}}{dt} = \frac{16\pi^3 \alpha_s^2 \Gamma_{ee}}{3\alpha_{em} M_V^5} \left[xg(x, Q^2) \right]^2$$

- Diffraction is very sensitive to (small- x) gluons!
- UPC is $\gamma - p$ or $\gamma - A$ collision
 - Nuclear DIS before the EIC era - and at very high energies!

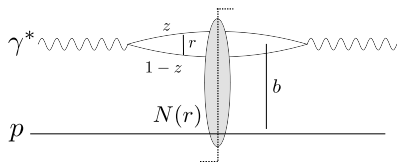
QCD at high energy: Color Glass Condensate



- CGC = QCD at high energies
- x (energy) dependence: BK/JIMWLK (perturbative)
- Saturation of gluon density at small x
- Saturation scale Q_s^2

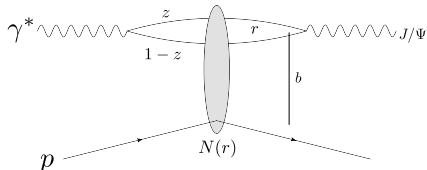
Natural framework to describe high energy scattering.

Deep inelastic scattering at high energy: dipole picture



Optical theorem:

$$\sigma^{\gamma^* p} \sim \text{dipole amplitude}$$



$$\sigma^{\gamma^* p \rightarrow V p} \sim |\text{dipole amplitude}|^2$$

Universal dipole amplitude

Same universal QCD evolved **dipole amplitude N** appears in calculations of

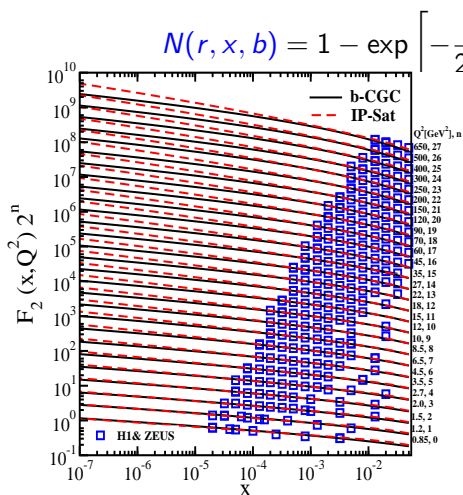
- DIS
- Diffraction
- Particle spectra in pp/pA
- ...

Non-perturbative input from a fit to HERA F_2 data.

IPsat model and DIS

Use impact parameter dependent dipole amplitude (IPsat) fitted to HERA

(Kowalski, Motyka, Watt, 2006; Rezaeian et al, 2012)



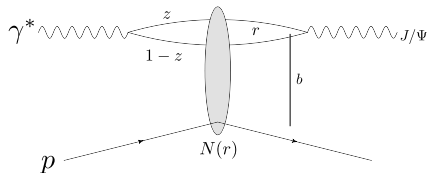
Rezaeian et al, 1307.0825

- DGLAP evolved gluon distribution $xg(x, \mu^2)$
- Proton profile T_p Gaussian
- Very good agreement with structure function data
- Generalization for nuclei:

$$S_A(r, b, x) = \prod_{i=1}^A S_p(r, b - b_i, x)$$
- Extremely good description of the precise HERA data

Diffractive vector meson production

- 1 $\gamma^* \rightarrow q\bar{q}$ splitting, wave function $\Psi^\gamma(r, Q^2, z)$
- 2 $q\bar{q}$ dipole scatters elastically
- 3 $q\bar{q} \rightarrow J/\Psi$, wave function $\Psi^V(r, Q^2, z)$



Diffractive scattering amplitude

$$\mathcal{A} \sim \int d^2b dz d^2r \Psi^{\gamma^*} \Psi^V(r, z, Q^2) e^{-ib \cdot \Delta} N(r, x, b)$$

- Fourier transfer from impact parameter to transverse momentum Δ
→ access to spatial structure

Still need to average over target configurations!

Coherent diffraction = target remains intact

Target is at the same quantum state before and after the scattering (Miettinen,

Pumplin, PRD 18, 1978, ...)

$$\frac{d\sigma^{\gamma^* p \rightarrow Vp}}{dt} \sim |\langle \mathcal{A}(x, Q^2, t) \rangle|^2$$

with

$$\mathcal{A} \sim \int d^2b dz d^2r \Psi^* \Psi^V(r, z, Q^2) e^{-ib \cdot \Delta} N(r, x, b)$$

- Coherent $t = -\Delta^2$ spectra is Fourier transfer of the **average density**

Incoherent diffraction = target breaks up

Total diffractive cross section – coherent cross section \Rightarrow target breaks up

$$\frac{d\sigma^{\gamma^* p \rightarrow V p^*}}{dt} \sim \langle |\mathcal{A}(x, Q^2, t)|^2 \rangle - |\langle \mathcal{A}(x, Q^2, t) \rangle|^2$$

with

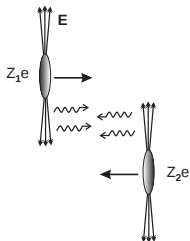
$$\mathcal{A} \sim \int d^2 b dz d^2 r \Psi^* \Psi^V(r, z, Q^2) e^{-ib \cdot \Delta} N(r, x, b)$$

- Incoherent cross section is proportional to the amount of fluctuations in the impact parameter space

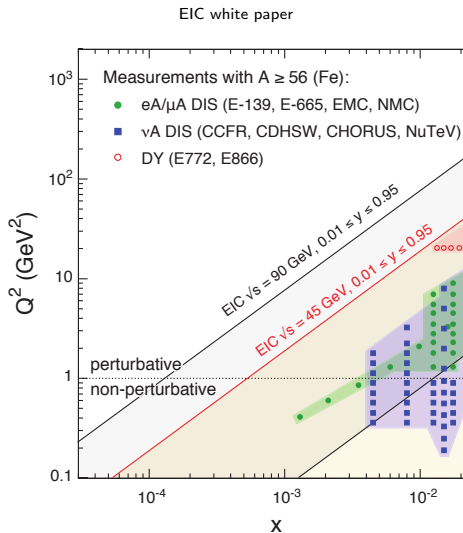
Cross section for incoherent γA diffraction: T. Lappi, H.M, 1011.1988

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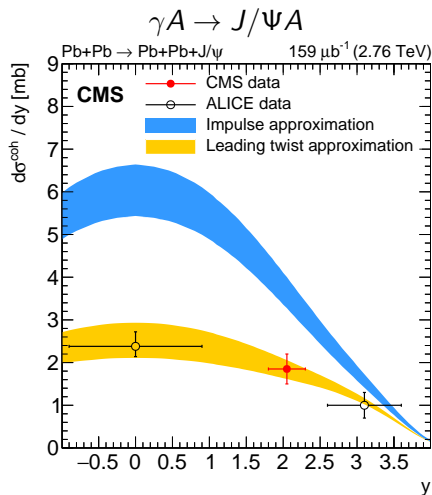
Going to small x with nuclei



- Currently there is basically no small- x nuclear DIS data
- Momentum fraction
 $x = M_V e^y / \sqrt{s}$
- Midrapidity J/ψ at the LHC:
 $x \sim 10^{-3}$
- Forward J/ψ at the LHC $x \sim 10^{-2}$ and $x \sim 10^{-5}$



Signal of nuclear effects seen in γA diffraction

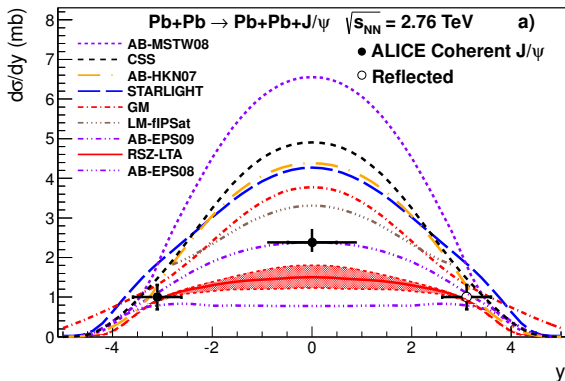


CMS, 1605.06966

Impulse approximation = scaled γp

- Clear nuclear effects seen (e.g. saturation, shadowing)

Coherent diffraction, model comparison

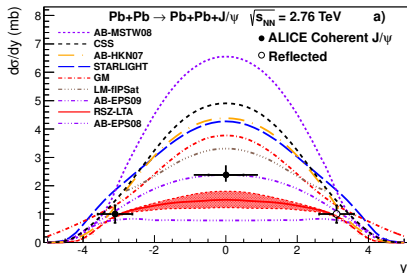


ALICE, 1305.1467

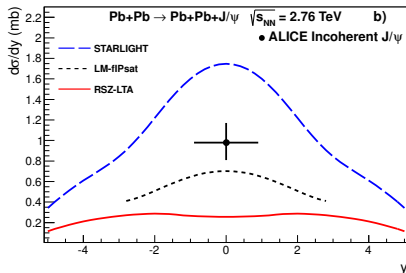
Shadowing/saturation needed, compare e.g.
AB-MSTW08 and AB-EPS09 (nuclear pdf) / LM-fIPsat (saturation)

Coherent and incoherent diffraction

Coherent



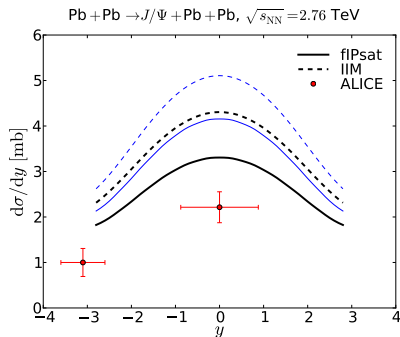
Incoherent



Dipole model calculation (LM-fIPsat): ok simultaneous description

- More LHC data coming
 - $\sqrt{s} = 5.02$ TeV: x dependence of the gluon density
 - Spectra differentially in $t \rightarrow$ Geometric structure and fluctuations
 - **QM2017**: Neutron tagging \Rightarrow large x /small x separation

This is becoming precision physics (linear scale!)

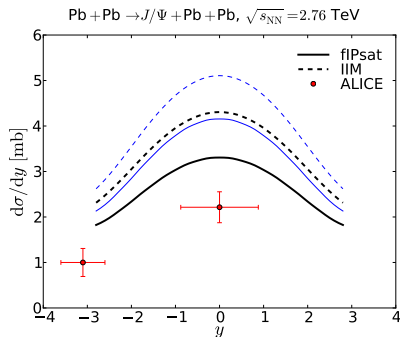


T. Lappi, H.M., 1301.4095

“Theory uncertainties” are still large

- Dipole-nucleus amplitude
 - No nuclear DIS data to fit
 - But other data, e.g. R_{pA}
- Vector meson wave function (thin-thick lines)
 - Constrained mainly by the leptonic decay width = wave function at origin!
- Large phenomenological corrections
 - Especially skewedness (2 gluons, $x \ll x'$) is large, $\sim 50\%$
- NLO???

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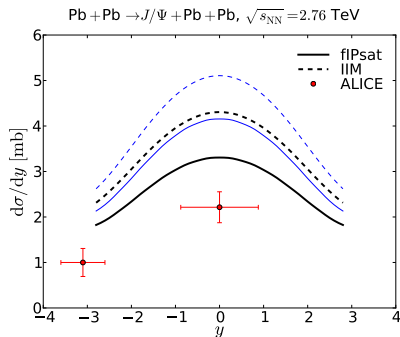


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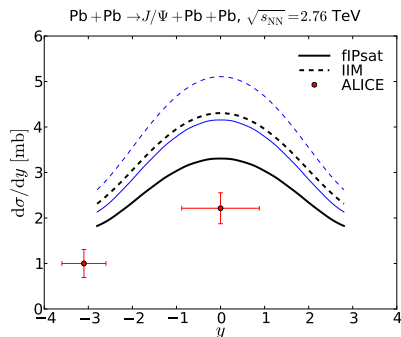


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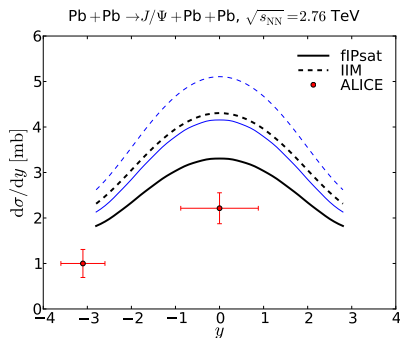


T. Lappi, H.M., 1301.4095

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T. Lappi, H.M., 1301.4095

- Model uncertainties mainly affect normalization, not rapidity (Bjorken- x) dependence.(?)

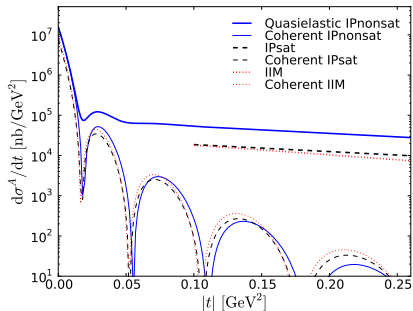
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More differential measurements in the near future

$$\gamma A \rightarrow J/\psi A$$

$$A=197, Q^2=0 \text{ GeV}^2, x_p=0.001$$



T. Lappi, H.M., 1011.1988

Solid line: no saturation

Dashed lines: with saturation

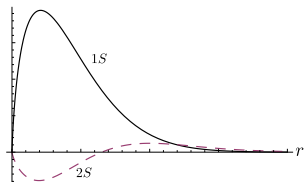
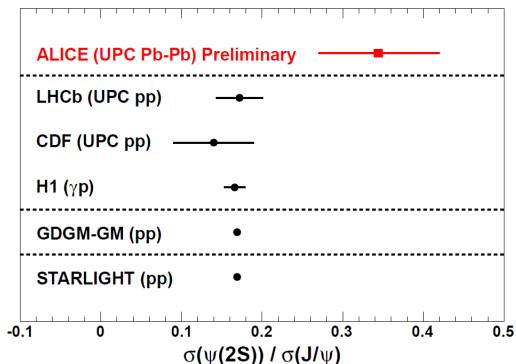
Measure differentially in vector meson momentum $t = -\Delta^2$

- Incoherent: more sensitive to saturation effects
 - **QM2017, CMS:** No/small incoherent cross section at small $x!$
 - Qualitatively predicted:
T. Lappi, H. M., 1011.1988
- Coherent: extract transverse density profile of the small- x gluons in the nucleus

Toll, Ullrich, 1211.3048

Excited states: $\Psi(2S)/J/\Psi$

Species dependence could test our understanding of the wave function

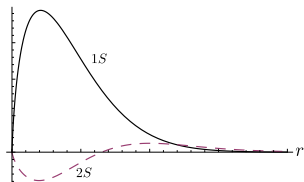
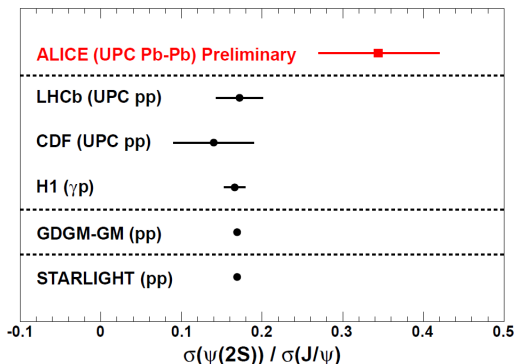


ALICE, 1508.05076

Qualitative agreement with dipole model: node effect is damped at large Q_s^2 (contribution from large dipoles is suppressed).

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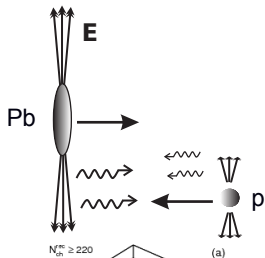
ALICE, 1508.05076

Qualitative agreement with dipole model: node effect is damped at large Q_s^2 (contribution from large dipoles is suppressed).

QM2017: $\sigma(2S)/\sigma(J/\Psi) = 0.166 \pm 0.011$

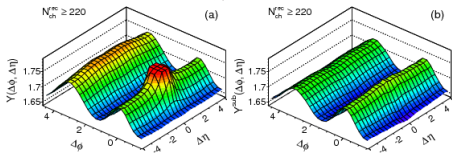
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Diffraction in pA collisions

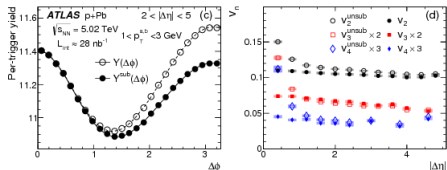


Photon flux $\sim Z^2$

- Ultrapерipheral pA collision $\approx \gamma p$ collision
- Access small- x structure of the proton

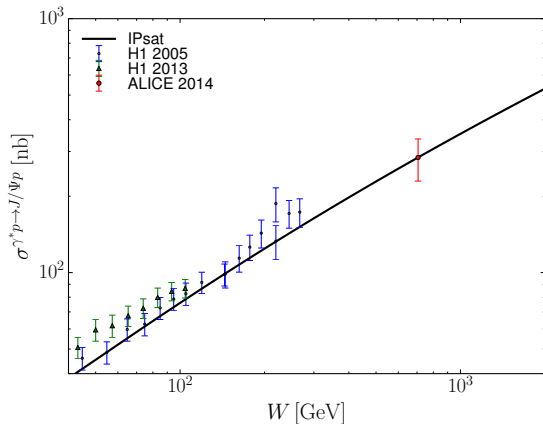


Collective phenomena seen in pp&pA



- Have to understand role of the initial state geometry
- Diffraction: average structure and fluctuations

Total coherent diffractive cross section



Total coherent cross section follows the same W^γ power law as HERA

- No significant saturation effect expected, described by IPsat

Recall:

- Coherent diffraction probes average structure
- Incoherent diffraction is sensitive to amount of fluctuations

Strategy

Simultaneous description of HERA coherent and incoherent data allow us to constrain event-by-event proton structure fluctuations.

Constraining proton fluctuations

Start with a simple constituent quark inspired picture:

- Sample quark positions from a Gaussian distribution, width B_{qc}
- Small- x gluons are located around the valence quarks (width B_q).
- Combination of B_{qc} and B_q sets the degree of geometric fluctuations
- Dipole-target scattering: IPsat model fitted to F_2 data

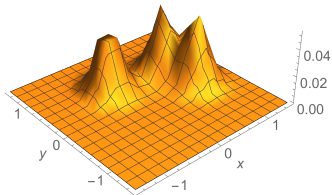
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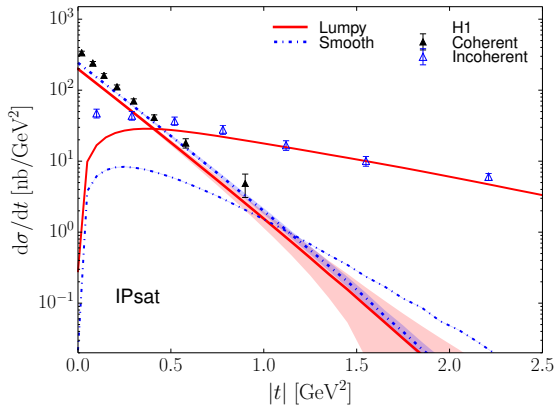
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Now proton = 3 overlapping hot spots.

$$T_{\text{proton}}(b) = \sum_{i=1}^3 T_q(b - b_i) \quad T_q(b) \sim e^{-b^2/(2B_q)}$$



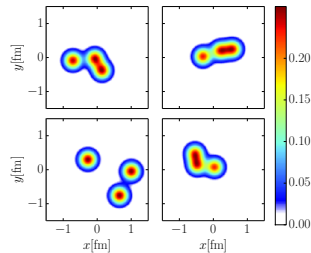
Lessons from the HERA data



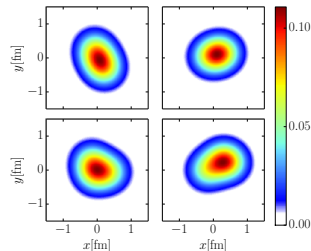
H.M, B. Schenke, PRL 117 (2016), 052301 and PRD94 (2016), 034042

- H1 incoherent data requires large fluctuations
- Proton-photon center-of-mass energy $W = 75 \text{ GeV}$, probing $x \approx 10^{-3}$

Lumpy: $B_{qc} = 3.3, B_q = 0.7$

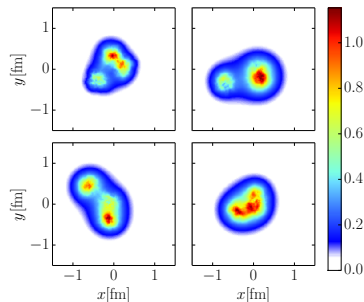
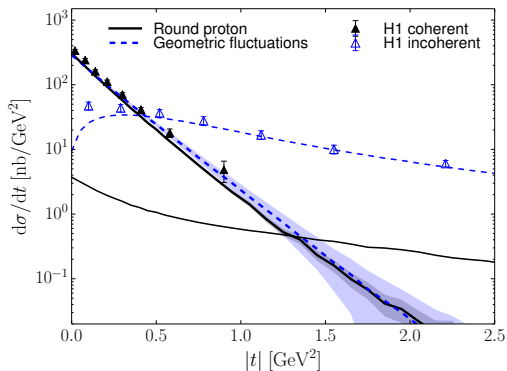


Smooth: $B_{qc} = 1.0, B_q = 3.0$



Units: GeV^{-2}

Include color charge fluctuation, parameters fitted to H1 data

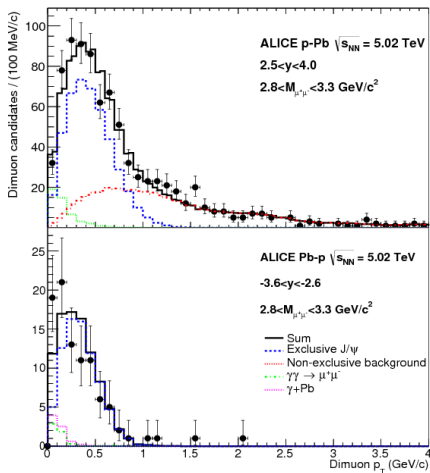


H.M., B. Schenke, PRD94 (2016), 034042

- Initial condition for pA hydro, good description of v_2 and v_3 data!

Towards smaller x / larger W in γp

ALICE measurement in $\gamma + p \rightarrow J/\psi + p(p^*)$ collisions

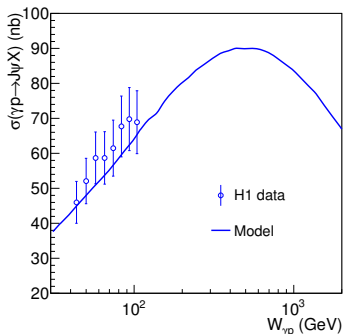


$$x \sim 10^{-2} \rightarrow 2 \cdot 10^{-5}$$

- Incoherent cross section not observed at small x
- Signature of smoothening at small x ?
- Proton grows, diffractive slope $B_p : 4 \text{ GeV}^{-2} \rightarrow \sim 6.7 \text{ GeV}^{-2}$

Proton smoothening at small x ?

Cepila, Contreras, Takaki (1608.07559): Number of constituent quarks $\sim x^a(1 + b\sqrt{x})$, parameters fitted to HERA data

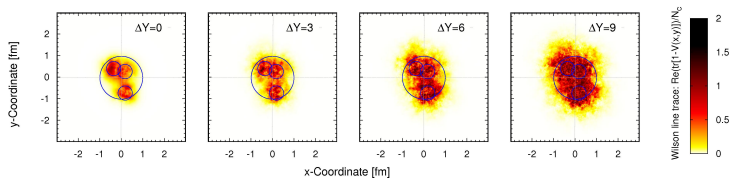


- Qualitatively expect incoherent cross section to decrease at high W
- Still expect to see significant incoherent contribution at ~ 1 TeV

Evolution to small x (work in progress)

HERA data constrains proton structure at $x \sim 10^{-3}$.

Evolve to smaller x by perturbative CGC evolution equation (JIMWLK)

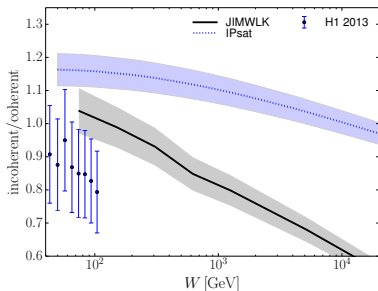


B. Schenke, S. Schlichting, Phys.Lett. B739 (2014) 313-319

- Proton grows
- Proton gets smoother

Energy evolution of diffractive J/ψ production

Work in progress / preliminary (qualitative results at this point)



- Incoherent cross section grows more slowly
 - Proton gets smoother
 - Dipole must not scatter off other constituent quarks (included in IPsat calculation T. Lappi, H.M., 1011.1988)
- We would expect to still see a large incoherent contribution at $W \sim 700 \text{ GeV}$

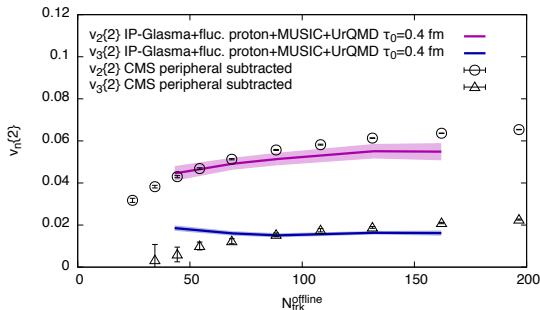
- UPC: diffractive vector meson production at very high energy
- CGC calculations compatible with both coherent and incoherent J/ψ measurements from the LHC
- Signatures of saturation/shadowing in heavy ion collisions
 - Still largish model dependence
- pA collisions: study proton structure at high energies
 - Description of HERA data requires large proton structure fluctuations
 - LHC pA data hints for smoothening at small x
 - Fluctuations applied to hydro calculations of pA collisions: good description of the v_n measurements (backup)
- Lots of new data coming!

BACKUPS

Hydro calculations with proton fluctuations from HERA

Hydro numbers

- $\tau_0 = 0.4$ fm
- $T_{fo} = 155$ MeV
- Shear and bulk viscosity
- Initial $\pi^{\mu\nu}$
- $\eta/s = 0.2$

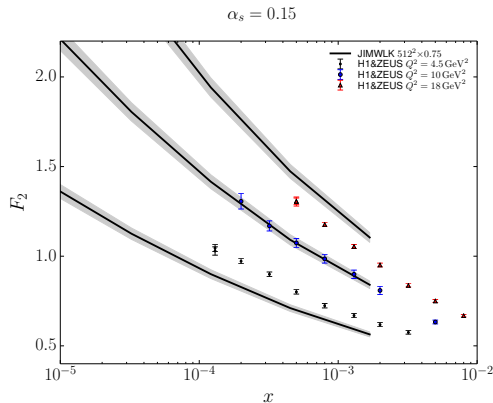


Large v_2 and v_3 at largest centrality bins reproduced well.

In preparation with B. Schenke, C. Shen, P. Tribedy

Constrain evolution speed

Work in progress / preliminary



- F_2 data at $Q^2 = 4.5 \dots 18 \text{ GeV}^2 \sim M_{J/\psi}^2$ constrain α_s
- MV model does not give exactly correct Q^2 dependence

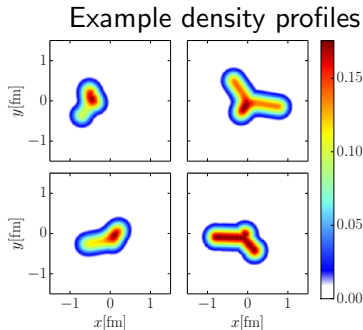
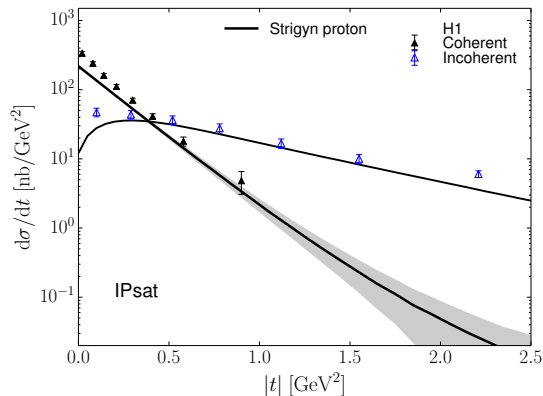
- Obtain saturation scale $Q_s(x_T)$ from IPsat (with fluctuations)
- MV-model: Sample color charges, density $\sim Q_s(x_T)$
- Solve Yang-Mills equations to obtain the Wilson lines

$$V(x_T) = P \exp \left(-ig \int dx^- \frac{\rho(x^-, x_T)}{\nabla^2 + m^2} \right)$$

- Dipole amplitude: $N(x_T, y_T) = 1 - \text{Tr} V(x_T) V^\dagger(y_T) / N_c$
- Fix parameters B_{qc} , B_q and m with HERA data

Lumpiness matters, not details of the density profile

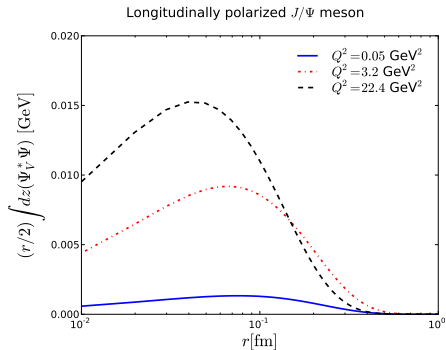
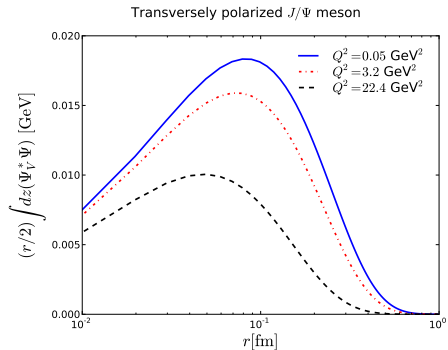
3 valence quarks that are connected by "color flux tubes" (Gaussian density profile, width B_q). Also good description of the data



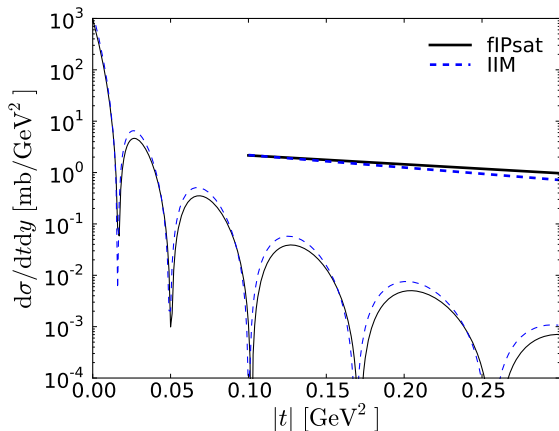
H.M., B. Schenke, PRD94 034042

Flux tubes implementation following results from hep-lat/0606016, used also e.g. in 1307.5911

Wave function overlap in J/Ψ production:

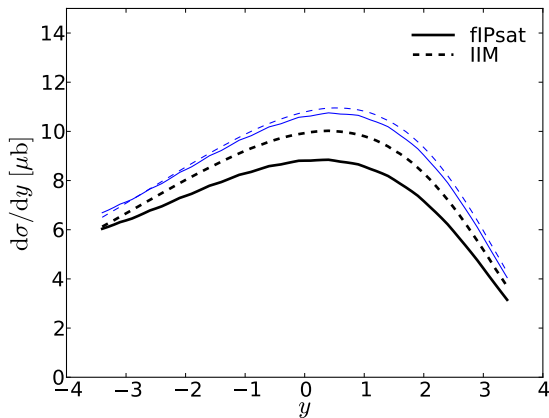


Differential cross section



T. Lappi, H. Mäntysaari, 1301.4095

Assuming proton profile function $T_p(b) \sim e^{b^2/(2B_p)} \Rightarrow$ incoherent cross section $\sim e^{-B_p t}$: probes spatial distribution of gluons in proton!



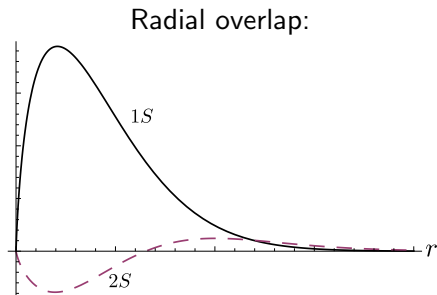
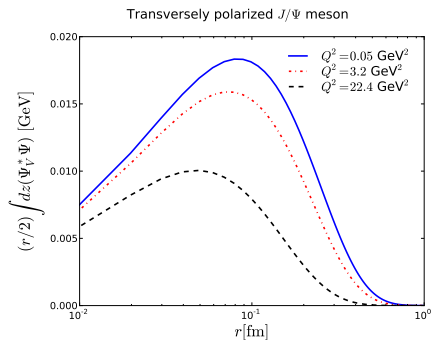
T. Lappi, H. Mäntysaari, 1301.4095 CMS frame

As the photon flux $\sim Z^2$, dominant process is the one where the nucleus emits the photon \Rightarrow probes mostly proton structure.

Vector meson wave functions

$\gamma^* \rightarrow q\bar{q}$ can be computed from QED, but $q\bar{q} \rightarrow$ vector meson requires some modelling, parameters fit to reproduce decay width.

Excited states: $\Psi(2S)$ wave function has a node (orthogonal to J/Ψ).
Cross section $\sim \int d^2r \Rightarrow$ large suppression compared to J/Ψ



Generalization for nuclei

S matrix \sim probability not to scatter [recall: $S = 1 - N$]:

$$S_A(r, b, x) = \prod_{i=1}^A S_p(r, b - b_i, x)$$

Average over nucleon configurations

$$\langle \mathcal{O}(\{b_i\}) \rangle_N = \int \prod_{i=1}^A [d^2 b_i T_A(b_i)] \mathcal{O}(\{b_i\})$$