Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion

Covariant extension of Generalized Parton Distributions

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Tomography of Hadrons and Nuclei at Jefferson Lab, Institute for Nuclear Theory, Seattle, August 31, 2017



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Introduction to Generalized Parton Distributions

- Definition and properties
- Experimental access

Representations of Generalized Parton Distributions

- Overlap of Light-cone wave functions
- Double Distributions

3 Covariant extention of Generalized Parton Distributions

- Motivation
- Inversion of Incomplete Radon Transform
- Results

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Definition of GF	PDs		

$$H^{q}(x,\xi,t) = \frac{1}{2} \int \frac{\mathrm{d}z^{-}}{2\pi} e^{i \times P^{+}z^{-}} \left\langle P + \frac{\Delta}{2} \left| \bar{q}(-z) \gamma^{+}q(z) \right| P - \frac{\Delta}{2} \right\rangle \Big|_{z^{+}=0, z_{\perp}=0}$$
(1)

with:

$$t=\Delta^2 \ , \ \xi=-rac{\Delta^+}{2\, P^+} \, .$$



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 $x+\xi$ $x-\xi$ **GPD** $P-\frac{\Delta}{2}$ $P+\frac{\Delta}{2}$

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• Similar matrix element for gluons.

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(2)

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Definition of			
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Theoretical o	onstraints on GPD)c	

Main properties:

• Physical region: $(x,\xi) \in [-1,1]^2$.

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Theoretical constraints on GPDs

Main properties:

- Physical region: $(x, \xi) \in [-1, 1]^2$.
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Image: A matrix

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Main properties:

- Physical region: $(x,\xi) \in [-1,1]^2$.
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- Link to PDFs and Form Factors:

$$\int \mathrm{d}x \, H^q(x,\xi,t) = F^q(t) \,, \tag{3}$$

$$H^{q}(x,0,0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x) .$$

$$(4)$$

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$$\int_{-1}^{1} \mathrm{d}x \, x^{m} \, H(x,\xi,t) = \text{Polynomial in } \xi \,. \tag{5}$$

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- From Lorentz invariance.
- Positivity (in DGLAP): (Pire et al., 1999; Radyushkin, 1999)

$$|H^{q}(x,\xi,t)| \leq \sqrt{q\left(\frac{x-\xi}{1-\xi}\right)q\left(\frac{x+\xi}{1+\xi}\right)}.$$
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Cauchy-Schwarz theorem in Hilbert space.

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Accessing GPDs			

• Exclusive processes:



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$$\mathcal{F}\left(\xi, t, Q^{2}\right) = \int_{-1}^{1} \mathrm{d}x \, C\left(x, \xi, \alpha_{S}\left(\mu_{F}\right), \frac{Q}{\mu_{F}}\right) F\left(x, \xi, t, \mu_{F}\right). \tag{7}$$

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- Observables are convolutions of:
 - a soft part, i.e. the GPD, with long distance interactions encoded (non-perturbative QCD).

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- Observables are convolutions of:
 - a soft part, i.e. the GPD, with long distance interactions encoded (non-perturbative QCD).
 - a hard-scattering kernel, calculated with perturbative QCD (short distance interactions).

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• A given hadronic state is decomposed in a Fock basis: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N,\beta} \int [\mathrm{d}x]_N \left[\mathrm{d}^2 \mathbf{k}_{\perp} \right]_N \Psi_{N,\beta}^{\lambda} \left(x_1, \mathbf{k}_{\perp 1}, \ldots \right) |N, \beta; k_1, \ldots, k_N\rangle , \qquad (10)$$

where the $\Psi_{N,\beta}^{\lambda}$ are the Light-cone wave-functions (LCWF).

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• For example, for the pion:

$$\left|\pi^{+}\right\rangle = \psi_{u\bar{d}}^{\pi} \left|u\bar{d}\right\rangle + \psi_{u\bar{d}g}^{\pi} \left|u\bar{d}g\right\rangle + \dots$$
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• GPD as an overlap of LCWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^{q}(x,\xi,t) = \sum_{N,\beta} \sqrt{1-\xi^{2-N}} \sqrt{1+\xi^{2-N}} \sum_{a} \delta_{a,q}$$
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$$\times \int [d\bar{x}]_{N} \left[d^{2}\bar{\mathbf{k}}_{\perp} \right]_{N} \delta(x-\bar{x}_{a}) \Psi^{*}_{N,\beta} \left(\hat{x}_{1}^{'}, \hat{\mathbf{k}}_{\perp 1}^{'}, ... \right) \Psi_{N,\beta} \left(\tilde{x}_{1}, \tilde{\mathbf{k}}_{\perp 1}, ... \right) ,$$

in the DGLAP region $\xi < x < 1$ (pion case).

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 - Polynomiality not manifest...

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Double Distrib	utions (DDs)		

$$H(x,\xi,t) \propto \int_{\Omega} d\beta \, d\alpha \, h(\beta,\alpha,t) \, \delta(x-\beta-\alpha\xi) \, . \tag{13}$$

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- DD h is defined on the support $\Omega = \{|\beta| + |\alpha| \le 1\}.$
- **Polynomial** in ξ :

$$\int_{-1}^{1} dx \, x^{m} H(x,\xi,t) \quad \propto \quad \int dx \, x^{m} \int_{\Omega} d\beta \, d\alpha \, h(\beta,\alpha,t) \, \delta(x-\beta-\alpha\xi)$$
$$\propto \quad \int_{\Omega} d\beta \, d\alpha \, (\beta+\xi\alpha)^{m} \, h(\beta,\alpha,t) \, . \tag{14}$$

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Double Distri	ibutions (DDs)		

$$H(x,\xi,t) \propto \int_{\Omega} d\beta \, d\alpha \, h(\beta,\alpha,t) \, \delta(x-\beta-\alpha\xi) \, . \tag{13}$$

- DD h is defined on the support $\Omega = \{ |\beta| + |\alpha| \le 1 \}.$
- **Polynomial** in ξ :

$$\int_{-1}^{1} dx \, x^{m} H(x,\xi,t) \quad \propto \quad \int dx \, x^{m} \int_{\Omega} d\beta \, d\alpha \, h(\beta,\alpha,t) \, \delta(x-\beta-\alpha\xi)$$
$$\propto \quad \int_{\Omega} d\beta \, d\alpha \, (\beta+\xi\alpha)^{m} h(\beta,\alpha,t) \, . \tag{14}$$

• Positivity not manifest...

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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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Double Distri	butions (DDs)		

$$H(x,\xi,t) \propto \int_{\Omega} d\beta \, d\alpha \, h(\beta,\alpha,t) \, \delta(x-\beta-\alpha\xi) \, . \tag{13}$$

• Radon Transform: (Deans, 1983; Teryaev, 2001)



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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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Double Distri	butions (DDs)		

• DD representation of GPDs:

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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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Outline			

Introduction to Generalized Parton Distributions

- Definition and properties
- Experimental access

Representations of Generalized Parton Distributions

- Overlap of Light-cone wave functions
- Double Distributions

3) Covariant extention of Generalized Parton Distributions

- Motivation
- Inversion of Incomplete Radon Transform
- Results

Conclusion

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Covariant ext	ension		

• What do we want?



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Covariant ext	ension	

- What do we want?
 - GPD at $\xi = 0$, for nucleon tomography.



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Introduction to GPDs	00	●000000	O
Covariant ext	ension		

- What do we want?
 - GPD at $\xi = 0$, for nucleon tomography.
- What do we have?



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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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Covariant ext	ension		

- Covariant extension
 - What do we want?
 - GPD at $\xi = 0$, for nucleon tomography.
 - What do we have?
 - Experimental access: integrals over x of GPD at $\xi > 0$.



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Covariant ext	ansion		

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 - Covariant approach?



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Covariant ext	ension		

- What do we want?
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 - Experimental access: integrals over x of GPD at $\xi > 0$.
- Covariant approach?
 - Positivity?







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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion

- What do we want?
 - GPD at $\xi = 0$, for nucleon tomography.
- What do we have?
 - Experimental access: integrals over x of GPD at $\xi > 0$.
- Covariant approach?
 - Positivity?
 - Loss of symmetries...







Covariant extension of GPDs

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Introducti 000	on to G	SPDs	Representations of GPDs 00	Covariant extention of GPDs ●○○○○○○	Conclusion O
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Covariant extension

- What do we want?
 - GPD at $\xi = 0$, for nucleon tomography.
- What do we have?
 - Experimental access: integrals over x of GPD at $\xi > 0$.
- Link to first principles through Light cone wave functions:



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Introduction to GPDs	Representations of GPDs	Covariant extention of GPD:
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 - Derive all parton distributions (not just GPDs)!





Introduction to GPDs	Representations of GP

Covariant extention of GPDs ●○○○○○○

- What do we want?
 - GPD at $\xi = 0$, for nucleon tomography.
- What do we have?
 - Experimental access: integrals over x of GPD at ξ > 0.
- Link to first principles through Light cone wave functions:
 - Derive all parton distributions (not just GPDs)!
 - But how to truncate?





Introduction	GPDs	

Covariant extention of GPDs ●○○○○○○

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Introduction	GPDs	

Covariant extention of GPDs

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- Use Lorentz invariance to extend from DGLAP!





Introduction	GPDs	

Covariant extention of GPDs

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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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Inversion			

Problem

Find $h(\beta, \alpha)$ on square $\{|\alpha| + |\beta| \le 1\}$ such that

$$\left| H(x,\xi) \right|_{\mathrm{DGLAP}} \propto \int \mathrm{d}\beta \,\mathrm{d}\alpha \,h(\beta,\alpha) \,\delta\left(x-\beta-\alpha\xi
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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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Introduction to GPDs 000	Representations 00	of GPDs	Covariant extention o	of GPDs	Conclusion O
Inversion					
Quark CDD. I		1	$ \dot{c} \rightarrow h(\theta, z)$	$0 f_{au} \partial < 0$	





Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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Inversion			

- Quark GPD: $H(x,\xi) = 0$ for $-1 < x < -|\xi| \Longrightarrow h(\beta,\alpha) = 0$ for $\beta < 0$.
- Domains $\beta < 0$ and $\beta > 0$ are uncorrelated in the DGLAP region.



Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion

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 - Divide and conquer:
 - Better numerical stability.
 - Lesser complexity: $O(N^p + N^p) \ll O((N + N)^p)$.



Inversion			
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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion

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- Divide and conquer:
 - Better numerical stability.
 - Lesser complexity: $O(N^p + N^p) \ll O((N + N)^p)$.
- α -parity of the DD: $h(\beta, -\alpha) = h(\beta, \alpha)$.



Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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Discretization			

• Expansion of the DD into basis functions $\{v_j\}$:

$$h(\beta,\alpha) = \sum_{j} h_{j} v_{j}(\beta,\alpha) , \qquad (15)$$

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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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Piece-wise constant, piece-wise linear, etc.



Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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- n columns of the matrix.



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- Sampling:



Introduction to GPDs Representations of GPDs Covariant extention of GPDs Conclusion 000 00 00 00 0				
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Discretization

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 - Random couples $(x,\xi) \longrightarrow m \ge n$ lines of the matrix.



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Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion O
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- Linear problem: AX = B where $B_i = H(x_i, \xi_i)$ and $A_{ij} = \mathcal{R}v_j(x_i, \xi_i)$.



Introduction to GPDs	Representations of GPDs	Covariant extention of GPDs	Conclusion
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- Regularization necessary: discrete ill-posed problem.



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- Regularization necessary: discrete ill-posed problem.
 - Trade-off between noise and convergence.





Figure: Extension of GPDs for the pion DSE model of Refs. (Mezrag, 2015; Mezrag et al., 2016). Comparison to the analytical result. Left: Plot for fixed ξ values 0, 0.5 and 1, at t = 0 GeV². Right: Plot for fixed t values 0, -0.25 and -0.5 GeV², at $\xi = 0.5.$

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E(x, ξ)

Figure: Extension of GPD E for the nucleon model of Ref. (Hwang and Mueller, 2008). Comparison to the analytical result of the authors. Left: Plot for fixed ξ values 0, 0.5 and 1, at t = 0 GeV². Right: Plot for fixed t values 0, -0.25 and -0.5 GeV², at

-0.75 -0.5 -0.25

0

x

0.25 0.5 0.75

0.8 0.6 لم

0.4

0.2

0

0.25 0.5 0.75 1



-1 -0.75 -0.5 -0.25 0

 $\xi = 0.5.$

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/Gel:

-0.1 -0.2 -0.3

-0.5 . 1

Representations of GPDs Covariant extention of GPDs 0000000

Some examples (gaussian model)



Figure: Extension of GPD for a gaussian pion model (in the vein of AdS/QCD). Left: Plot for fixed ξ values 0, 0.5 and 1, at t = 0 GeV². Right: Plot for fixed t values 0. -0.25 and -0.5 GeV², at $\xi = 0.5$.




• Integrable singularity for the GPD at $x \sim 0$: $H(x,\xi) \propto \frac{1}{\sqrt{x}}$.

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- Integrable singularity for the GPD at $x \sim 0$: $H(x,\xi) \propto \frac{1}{\sqrt{x}}$.
- Equivalent to an integrable singularity for the DD at $\beta \sim 0$: $h(\beta, \alpha) \propto \frac{1}{\sqrt{\beta}}$.



- Integrable singularity for the GPD at $x \sim 0$: $H(x,\xi) \propto \frac{1}{\sqrt{x}}$.
- Equivalent to an integrable singularity for the DD at $\beta \sim 0$: $h(\beta, \alpha) \propto \frac{1}{\sqrt{\beta}}$.

• We solve for
$$\sqrt{\beta} h(\beta, \alpha)$$
 instead of $h(\beta, \alpha)$!

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Summary			

• Generalized Parton Distributions

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Summary			

- Generalized Parton Distributions
 - encode information about the 3D structure of a hadron.

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• Generalized Parton Distributions

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- encode information about the 3D structure of a hadron.
- are accessible with exclusive processes in experiments: JLab, COMPASS, etc.

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 - $\blacktriangleright \ \mathsf{LCWFs} \underset{\mathrm{Overlap}}{\longrightarrow} \mathsf{GPD} \text{ in } \mathsf{DGLAP} \underset{\mathrm{Inverse \ Radon \ Transform}}{\longrightarrow} \mathsf{DD} \underset{\mathrm{RT}}{\longrightarrow} \mathsf{GPD}.$

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 - Both polynomiality and positivity!

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- Unified phenomonelogy of GPDs and TMDs at the level of LCWFs?

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 - Both polynomiality and positivity!
 - Compromise with respect to noise and convergence.
- Unified phenomonelogy of GPDs and TMDs at the level of LCWFs?
- Thank you!
 - Any questions?



Ill-posed problems and Regularization

- Ill-posed problems?
 - For example the inversion of a Fredholm equation of the first kind:

$$\int \mathcal{K}(x,y)f(y)\,\mathrm{d}y = g(x)\,. \tag{16}$$

- The inverse is not continuous: an arbitrarily small variation Δg of the rhs can lead to an arbitrarily large variation Δf of the solution.
- The corresponding discrete problem needs to be regularized.

• E.g Tikhonov regularization: min $\{ \|AX - B\|^2 + \epsilon \|X\|^2 \}$.





Theoretical "L-curve": curve parameterized by the regularization factor.

(fig. taken from Ref. (Hansen, 2007))

Nabil Chouika

Covariant extension of GPDs

L-curve with the iteration number as

regularization factor.

D-term considerations

Polynomiality property:

$$\int_{-1}^{1} \mathrm{d}x \, x^{m} H\left(x, \xi, t\right) = \sum_{\substack{k=0\\k \, even}}^{m+1} c_{k}^{(m)}(t) \xi^{k} \,. \tag{17}$$

Recast polynomiality property for H - D:

$$\int_{-1}^{1} \mathrm{d}x \, x^{m} \left(H\left(x, \xi, t\right) - D\left(\frac{x}{\xi}, t\right) \right) = \sum_{\substack{k=0\\k \text{ even}}}^{m} c_{k}^{(m)}(t) \xi^{k} \,, \tag{18}$$

where $D\left(\frac{x}{\xi}, t\right)$ is the so-called D-term with support on $-\xi < x < \xi$. H - D is a Radon Transform:

$$H(x,\xi,t) - D\left(\frac{x}{\xi},t\right) = \int_{\Omega} d\beta \, d\alpha \, h_{\rm PW}\left(\beta,\alpha\right) \delta\left(x-\beta-\alpha\xi\right) \,. \tag{19}$$

The DGLAP region gives no information on the D-term.

With other DD representations, we can generate intrinsic D-terms, e.g. Pobylitsa representation:

$$H(x,\xi,t) = (1-x) \int_{\Omega} d\beta \, d\alpha \, h_{\rm P}(\beta,\alpha) \, \delta(x-\beta-\alpha\xi) \,.$$
⁽²⁰⁾

Still freedom of extra D-term.

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