Modelling BNS mergers and their aftermath: the role of neutrinos

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TU-Darmstadt \rightarrow INFN, Milano-Bicocca & Parma

14 August 2017 Talk at INT: Electromagnetic Signature of r-process Nucleosynthesis in NS Binary Mergers

in collaboration with A. Arcones, D. Martin, W. Kastaun, O. Korobkin, S. Rosswog, H. Yasin ...









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Outline

- Introduction and background
- The role of neutrinos in BNS merger aftermath
 - Neutrino-driven winds
 - Neutrinos and dynamic ejecta
 - Neutrino pair annihilations
- Conclusions and open questions

BNS merger remnant properties

merger remnant: final stages of BNS system evolution



- ► Massive NS (\rightarrow BH) $M \sim 2.2 - 2.8 M_{\odot},$ $\rho \gtrsim 10^{12} \text{g cm}^{-3}$ $T \sim a \text{ few } 10 \text{ MeV}$
- $\begin{array}{l} \blacktriangleright \mbox{ thick accretion disk } \\ M\sim 10^{-3}-10^{-1}M_{\odot} \\ Y_{e}\lesssim 0.20 \\ T\sim a\mbox{ few MeV} \\ \left(Y_{e}=\frac{n_{e}}{n_{B}}\approx \frac{n_{p}}{n_{p}+n_{n}}\right) \end{array}$
- intense ν emission $E_{\nu} \gtrsim 10 \text{ MeV}$ $L_{\nu,\text{tot}} \sim 10^{53} \text{erg s}^{-1}$ $L_{N,\nu,\text{tot}} \sim 10^{57} \text{ particles s}^{-1}$

Processes in BNS merger remnants

Relevant ingredients and their interplay:

- magnetic field (e.g., MRI)
 - relevant source of viscosity
- viscosity
 - accretion
 - angular momentum redistribution
- weak processes
- e.g. works by Just, Fernandez, Metzger, Roberts, Radice, Bovard ...
- neutrino luminosities
- matter composition (remnant & ejecta)
- remnant fate
 - BH formation
 - metastable MNS
 - stable MNS

Large variability expected

- dependence on the BNS properties
- influence of EOS

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e.g., Siegel& Metzger17, Ciolfi+2016

e.g. Radice 16, Shibata+16

e.g., works by Just, Fernandez, Metzger

e.g., Piro+17

Neutrino-matter interaction in BNS merger remnants

Roberts talk from last week

- v's are weakly interacting particles:
- both CC and NC channels
- production (and possibly absorption):
 - $\begin{array}{ll} p+e^- \to n+\nu_e \ ({\sf EC}) & e^-+e^+ \to \nu+\bar{\nu} \\ n+e^+ \to p+\bar{\nu}_e \ ({\sf PC}) & N+N \to N+N+\nu+\bar{\nu} \end{array}$
- scattering:

$$N + \nu \rightarrow N + \nu$$
 $e^{\pm} + \nu \rightarrow e^{\pm} + \nu$

Neutrino production rates:

- production boosted by high temperatures & densities:
 - $R_{EC} \propto \rho X_p T^5 F_4(\mu_e/T)$
 - $R_{PC} \propto \rho X_n T^5 F_4(-\mu_e/T)$

e.g. Rosswog & Liebendörfer 03

Neutrino opacity in BNS merger remnants

Neutrino absorption/scattering rates: neutrino opacity \leftrightarrow neutrino mean free path, λ_{ν}

$$\begin{aligned} \sigma_{\nu} &\sim \sigma_0 \left(\frac{E_{\nu}}{m_{\rm e}c^2}\right)^2 \quad \sigma_0 = \frac{4G_F^2(m_ec^2)^2}{\pi(\hbar c)^4} \approx 1.76 \times 10^{-44} \,{\rm cm}^2 \approx 2.6 \times 10^{-20} \sigma_t \\ \lambda_{\nu} &\approx \frac{1}{n_{\rm target}\sigma_{\nu}} \sim 2.36 \times 10^{19} {\rm cm} \left(\frac{\rho}{1\,{\rm g/cm}^3}\right)^{-1} \left(\frac{E_{\nu}}{1\,{\rm MeV}}\right)^{-2} \end{aligned}$$

Neutrino opacity in BNS merger remnants

Neutrino absorption/scattering rates: neutrino opacity \leftrightarrow neutrino mean free path, λ_{ν}



Neutrino Surfaces

Optical depth:

$$\blacktriangleright \ \tau = \int_{\gamma} \frac{1}{\lambda} \, \mathrm{d}s$$

- measure of how optically thin/thick matter is moving along an escaping path γ
- $\lambda_{sc} \approx \sum \lambda_i \rightarrow \tau_{\nu,sc} \sim \# \text{ of } \nu\text{-matter interaction along } \gamma$
- $\lambda_{en} \approx \sqrt{\lambda_{sc}\lambda_{ab}} \rightarrow \tau_{\nu,en} \sim \text{# of inelastic } \nu$ -matter along γ

 $\tau_{\nu} = 2/3 \quad \Rightarrow \quad \nu \text{ surfaces, for } E_{\nu} = 4.6, 10.6, 16.2, 24.6, 57.0 \, \mathrm{MeV}$



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- $\lambda_{sc} \approx \sum \lambda_i \rightarrow \tau_{\nu,sc} \sim \#$ of ν -matter interaction along γ
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Properties of ν 's emission in BNS mergers

evolution of ν luminosities:

- inspiral phase: negligible
- merger: sudden increase due to T increase
- post merger: powered by accretion and cooling
- possible dependence on nuclear EOS e.g. Sekiguchi+15, Radice+16

properties of ν luminosities:

- ν gas formation and diffusion
- n richness $\rightarrow L_{\bar{\nu}_e} \gtrsim L_{\nu_e}$
- ► different decoupling depths: E_{νe} < E_{νe} < E_{νu,τ}

Rosswog+13 (up), Neilsen+15 (down)

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Role of ν 's in BNS mergers

Role of ν 's

- release energy (cooling effect)
- ► set *n*-to-*p* ratio $p + e^- \leftrightarrow n + \nu_e \& n + e^+ \leftrightarrow p + \bar{\nu}_e$
- exchange energy and momentum with matter

 ⇒ relevant to set EM couterpart properties

Rosswog+13 (up), Neilsen+15 (down)

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ν -driven winds from BNS merger remnats: goals

Goals

- characterize ν emission
- study wind development
- perform nucleosynthesis on ejecta
- compute EM counterparts

see also Dessart+09,Metzger&Fernandez14,Just+14,Sekiguchi+15

disk lifetime:

$$t_{\rm disk} \sim \alpha^{-1} \left(\frac{H}{R}\right)^{-2} \Omega_{\rm K}^{-1} \sim 0.31 \, {\rm s} \left(\frac{\alpha}{0.05}\right)^{-1} \left(\frac{H/R}{1/3}\right)^{-2} \left(\frac{R_{\rm disk}}{100 \, {\rm km}}\right)^{3/2} \left(\frac{M_{\rm ns}}{2.5 \, M_{\odot}}\right)^{-1/2} \, {\rm s} \left(\frac{M_{\rm ns}}{1.5 \, M_{\odot}}\right)^{-1} \, {\rm s} \left(\frac{M_{\rm ns}}{1.5 \, M_{\odot}}\right)^{-1$$

 α : viscosity coefficient $R_{\rm disk}$: disk typical radius H/R: disk aspect ratio Ω_K : Keplerian angular velocity

M_{ns}: MNS mass

disk lifetime:

$$t_{\rm disk} \sim 0.31 \,\mathrm{s} \,\left(\frac{\alpha}{0.05}\right)^{-1} \left(\frac{H/R}{1/3}\right)^{-2} \left(\frac{R_{\rm disk}}{100 \,\mathrm{km}}\right)^{3/2} \left(\frac{M_{\rm ns}}{2.5 \,M_{\odot}}\right)^{-1/2}$$

disk L:

$$\begin{split} L_{\nu,\text{disk}} &\sim \frac{\Delta E_{\text{grav}}}{2 \, t_{\text{disk}}} \quad \approx \quad 8.35 \times 10^{52} \, \text{erg} \, \text{s}^{-1} \left(\frac{M_{\text{ns}}}{2.5 \, M_{\odot}}\right)^{3/2} \left(\frac{M_{\text{disk}}}{0.2 \, M_{\odot}}\right) \left(\frac{R_{\text{disk}}}{100 \, \text{km}}\right)^{-3/2} \\ &\times \left(\frac{\alpha}{0.05}\right) \left(\frac{R_{\text{ns}}}{25 \, \text{km}}\right)^{-1} \left(\frac{H/R}{1/3}\right)^2 \end{split}$$

 $\Delta E_{\rm grav}$: gravitational energy released during accretion

disk lifetime:

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MNS L:

$$\begin{split} L_{\nu,\mathrm{ns}} &\sim \frac{\Delta E_{\mathrm{ns}}}{t_{\mathrm{cool,ns}}} \quad \approx \quad 1.86 \times 10^{52} \, \mathrm{erg \, s^{-1}} \left(\frac{\Delta E_{\mathrm{ns}}}{3.5 \times 10^{52} \, \mathrm{erg}} \right) \left(\frac{R_{\mathrm{ns}}}{25 \, \mathrm{km}} \right)^{-2} \\ & \left(\frac{\rho_{\mathrm{ns}}}{10^{14} \, \mathrm{g cm^{-3}}} \right)^{-1} \left(\frac{k_{\mathrm{B}} T_{\mathrm{ns}}}{15 \, \mathrm{MeV}} \right)^{-2} \end{split}$$

 $\Delta E_{\rm ns}$: thermal energy $t_{\rm ns,cool} \sim 3\tau_{\nu,\rm ns}/(R_{\rm ns}c)$: diffusion time scale $\tau_{\nu,\rm ns}$: ν optical depth in MNS

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disk lifetime:

• disk L:

$$t_{\text{disk}} \sim 0.31 \,\text{s} \, \left(\frac{\alpha}{0.05}\right)^{-1} \left(\frac{H/R}{1/3}\right)^{-2} \left(\frac{R_{\text{disk}}}{100 \,\text{km}}\right)^{3/2} \left(\frac{M_{\text{ns}}}{2.5 \, M_{\odot}}\right)^{-1/2}$$

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MNS L:

$$L_{\nu,\rm ns} \sim 1.86 \times 10^{52} \,{\rm erg \, s^{-1}} \left(\frac{\Delta E_{\rm ns}}{3.5 \times 10^{52} \,{\rm erg}}\right) \left(\frac{R_{\rm ns}}{25 \,{\rm km}}\right)^{-2} \dots$$

wind time:

$$\begin{split} t_{\rm wind} \sim \frac{e_{\rm grav}}{\dot{e}_{\rm heat}} &\approx 0.072\,{\rm s}\,\left(\frac{M_{\rm ns}}{2.5\,M_\odot}\right) \left(\frac{R_{\rm disk}}{100\,{\rm km}}\right) \left(\frac{E_\nu}{15\,{\rm MeV}}\right)^{-2} \\ & \left(\frac{\xi L_{\nu_e}}{4.5\times10^{52}\,{\rm erg\,s^{-1}}}\right)^{-1} \end{split}$$

 e_{grav} : specific gravitational energy \dot{e}_{heat} : specific heating rate

 ξL_{ν_e} : isotropized ν_e luminosity at $\theta \approx \pi/4$, $\xi \sim 1.5$ and $L_{\nu_e} \sim (L_{ns} + L_{disk})/3$

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ν -driven winds from BNS merger: the model

initial conditions:

- final stage of high resolution BNS merger simulation performed with SPH code
 Rosswog&Price07
- 2 x 1.4 M_{\odot} , with Shen EOS
- > 3D hydro: FISH code
 κăppeli+11
 ν treatment: ASL
 Perego+16
 nuclear EoS: TM1
 Hempel+12
- tracer particles
- nucleosynthesis network: Winnet e.g. Winteler+12

The ASL scheme

Perego, Cabezon, Käppeli 2016

ASL scheme exploits potentiality of classical gray leakage scheme e.g. Ruffert et al. 1997, Rosswog & Liebendörfer 2003, O'Connor & Ott 2011

What's new:

- inclusion of ν trapped components
- distinction between different ν energies ($\sigma_{\nu} \propto E_{\nu}^2$)
- *v*-absorption terms in optically thin conditions (ray-tracing)

Strengths:

- robust and flexible
- computationally inexpensive
- limited set of free parameters (calibrated with CCSNe)

Caveats:

- effective treatment
- diffusion part: order-of-magnitude

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Neutrino luminosities dependence on time

Neutrino luminosities dependence on time

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Neutrino luminosities

dependence on time

- mean energy hierarchy: $E_{\nu_{\mu,\tau}} > E_{\bar{\nu}_e} > E_{\nu_e}$
- $E_{\nu_e} \approx 11 \,\text{MeV}, E_{\bar{\nu}_e} \approx 15 \,\text{MeV}, E_{\nu_{\mu,\tau}} \approx 18 \,\text{MeV}$
- disk-shadow effect

Disk and wind dynamics

t = 0 ms

Disk and wind dynamics

t = 40 ms

Disk and wind dynamics

t = 90 ms

Ejection mechanism

- neutrino absorption $\rightarrow e_{tot} \approx 0 \rightarrow$ wind expansion
- ► nuclear recombination → matter ejection

Wind ejecta and nucleosynthesis

- $m_{\rm ej} \approx 0.05 \, M_{\rm disk} @\, 200 \, {\rm ms}$
- non-equatorial emission:
 θ < 60^o
- ► larger Y_e at small θ $Y_{e,eq} \sim \frac{\lambda_{\nu_e}}{\lambda_{\nu_e} + \lambda_{\bar{\nu}_e}} \approx 0.4$ $t_{wind}^{-1} \lesssim \lambda_{\nu_e} \lesssim \lambda_{\bar{\nu}_e} \approx 30 \text{ s}^{-1}$

e.g., Qian & Woosley 96

 ► thermodyn properties as input for nuclear network
 → nucleosynthesis abundances

wind ejecta (color lines) + dynamical ejecta (black lines)

from Korobkin+12

Discussion I

- which are the requirements for neutrino treatments in BNS merger?
- how relevant is the modelling of a v trapped component?
- which are the most relevant weak reactions and under which termodynamical conditions do they occur?

How robust are dynamic ejecta?

First dynamic ejecta simulations + nucleosynthesis calculations

e.g. Korobkin+2012, Bauswein+2013, Hotokezaka+2013

LK_QC

M0_OC

0.32 0.40

- extremely n-rich, $Y_e \lesssim 0.05$
- robust r-process nucleosynthesis, $A \gtrsim 120$
- at most, neutrino cooling included

GR simulations + neutrino treatment

e.g. Wanajo+14,Sekiguchi+15 (left figure),Foucart+15,Radice+16 (right figure)

- still robust r-process?
- effect of e^+ capture and/or ν_e absorption?

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Parametric study: the input

Perego+, in preparation; cfr. Goriely+2015

Kastaun+17

shock heated dynamic ejecta from GR simulation

- ▶ 1.4 M_{\odot} , equal mass merger, with stiff EOS
- rings of expanding matter: disk matter accelerated by shock produced by MNS oscillations
- Y_e : weak equilibrium at T = 0, simply advected
- $\blacktriangleright~\langle Y_e \rangle_{sim} \sim 0.044$

Parametric study: the model

Post-processing of tracer particles to include ν 's feedback

coupled Y_e & entropy evolution:

$$\frac{dY_{e}}{dt} = (\lambda_{\nu_{e}} + \lambda_{e^{+}}) Y_{n} - (\lambda_{\bar{\nu}_{e}} + \lambda_{e^{-}}) Y_{p}
\frac{ds}{dt} = \left(\frac{ds}{dt}\right)_{hydro} + \frac{1}{T} \left[\left(\frac{dQ}{dt}\right)_{\nu} - (\mu_{e} - \mu_{n} + \mu_{p}) \left(\frac{dY_{e}}{dt}\right)_{\nu} \right]$$

- optically thin conditions ($\rho < 10^{12} \, {
 m g/cm^3}$)
- consistent ν emission, $\lambda_{e^{\pm}}$

Bruenn 1985 + Horowitz 2002

Perego+, in preparation: cfr. Goriely+2015

► parametrized L_{ν} for ν absorption, $\lambda_{\nu} \propto \frac{L_{\nu}}{\langle E_{\nu} \rangle R^2} (\cos \theta)^{\alpha}$

Name	$L_{\nu_{e}, \max}$	$L_{\overline{\nu}_{e}, \max}$	$E_{\nu_{e}, \max}$	$E_{\bar{\nu}_{e}, \max}$
	[10553 erg/s]	[1053 erg/s]	[MeV]	[MeV]
capture	0.0	0.0	0.0	0.0
low	0.86	1.0	11.5	16.2
medium	1.0	1.5	12.0	16.3
high	1.2	2.4	13.0	16.7

Dynamic ejecta: Representative tracer

Profiles from a representative tracer

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Dynamic ejecta: Distributions

Property distributions at 8 GK (2 different initial conditions):

- ▶ weak equilibrim down to $10^{12} \text{ g/cm}^3 \rightarrow \text{solid lines}$
- $\blacktriangleright\,$ simulation Y_e (i.e., $Y_e=0.044) \rightarrow$ dashed lines

$\alpha = 0$, i.e. isotropic ν emission

Dynamic ejecta: Distributions

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$\alpha = 2$, i.e. unisotropic ν emission

Dynamic ejecta: Nucleosynthesis

Nucleosynthesis from post-processed tracers

- mostly independent from initial conditions
- significant dependence on luminosity spatial distribution

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Discussion II

- which are the requirements for neutrino treatments in BNS merger?
- how relevant is the modelling of a v trapped component?
- which are the most relevant weak reactions and under which termodynamical conditions do they occur?
- do we understand the properties of the neutrino emission and its variability?
- which ρ , T and Y_e conditions occur in and after BNS mergers?

ν - $\bar{\nu}$ annihilation and short GRB engine

 ν - $\bar{\nu}$ annihilation rate as sGRB central engine?

Eichler et al. 1989

$$u_{e,\mu,\tau} + \bar{\nu}_{e,\mu,\tau} \rightarrow e^- + e^+$$

net effect: $\boldsymbol{\nu}$ energy and momentum deposition in plasma

 $q_{\rm ann}(t, \mathbf{x}) \approx \frac{1}{6} \frac{\sigma_0(C_A^2 + C_V^2)}{c(m_e c^2)} \int I_{\nu} I_{\bar{\nu}} \left(E_{\nu} + E_{\bar{\nu}} \right) \left(1 - \cos \Phi \right)^2 \, \mathrm{d}\Omega_{\nu} \mathrm{d}\Omega_{\bar{\nu}} \mathrm{d}E_{\nu} \mathrm{d}E_{\bar{\nu}}$

why relevant in the funnel?

- high intensity I_{ν} & good collision angle Φ
- Iower baryonic pollution
- open questions: interplay with mag field, MNS VS BH ...

Accretion rate and luminosities

Presence of the MSN VS BH \rightarrow relevant differences

- $L_{\nu} \not\propto \dot{M}$
- MNS cooling as additional v source
- structure of the innermost disk
- thick disk cooling

Energy deposition rates

post-processing of previous 3D, long term remnant simulations

Perego, Yasin & Arcones 2017

cfr. Dessart+09,Richers+15,Just+16,Fujibayashi+17

Energy deposition by $\nu_e - \bar{\nu}_e$

•
$$Q_{\nu,\bar{\nu}}(t) = \int_V q_{\nu,\bar{\nu}}(\mathbf{x},t) \, \mathrm{d}V \sim 10^{49-50} \, \mathrm{erg/s}$$

- $E_{\nu,\bar{\nu}}(t) = \int_t Q_{\nu,\bar{\nu}}(t') dt', E_{\nu,\bar{\nu}}(400 \text{ms}) \sim 2 \times 10^{49} \text{ erg}$
- $q_{\nu_x} \sim q_{\nu_e}/60$, due to lower coupling and luminosities
- GR effects (e.g., Doppler boost, red/blue shift) cause only a marginal decrease

Energy deposition rates

post-processing of previous 3D, long term remnant simulations

Perego, Yasin & Arcones 2017

cfr. Dessart+09,Richers+15,Just+16,Fujibayashi+17

$$I_{\nu} = (I_{\nu})_{\rm NS} + (I_{\nu})_{\rm DS} \Rightarrow$$
$$q_{\nu,\bar{\nu}} = (q_{\nu,\bar{\nu}})_{\rm NS,NS} + (q_{\nu,\bar{\nu}})_{\rm DS,DS} + (q_{\nu,\bar{\nu}})_{\rm NS\&DS}$$

- ▶ NS & Disk contribution \approx Disk-Disk contribution
- NS & NS contribution locally intense, but globally small

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Parametrization and efficiency

$$Q_{\nu,\bar{\nu}} \approx \mathcal{C} L_{\nu} L_{\bar{\nu}} \left[\frac{\langle \epsilon_{\nu}^2 \rangle}{\langle \epsilon_{\nu} \rangle} + \frac{\langle \epsilon_{\bar{\nu}}^2 \rangle}{\langle \epsilon_{\bar{\nu}} \rangle} \right] \, G_{\nu,\bar{\nu}}$$

•
$$G_{\nu,\bar{\nu}}$$
 geometrical factor

 constant during 400ms (remarkably)

$$\eta_{\nu,\bar{\nu}} = \frac{Q_{\nu,\bar{\nu}}}{L_{\nu} + L_{\bar{\nu}}}$$

• low efficiency, $\eta_{\nu,\bar{\nu}} \sim 10^{-3}$

• if
$$L_{\nu} \approx L_{o} \exp -t/\tau_{c} \& L_{\bar{\nu}}/L_{\nu} \approx \beta$$

 $\eta_{\nu,\bar{\nu}} \sim G_{\nu,\bar{\nu}} CL_{o} \beta/(\beta+1) \times \left[\frac{\langle \epsilon_{\nu}^{2} \rangle}{\langle \epsilon_{\nu} \rangle} + \frac{\langle \epsilon_{\bar{\nu}}^{2} \rangle}{\langle \epsilon_{\bar{\nu}} \rangle}\right] \exp(-t/\tau)$

Discussion III

- which are the requirements for neutrino treatments in BNS merger?
- how relevant is the modelling of a v trapped component?
- which are the most relevant weak reactions and under which termodynamical conditions do they occur?
- how can we understand the change in Y_e for the dynamic ejecta?
- do we understand the properties of the neutrino emission and its variability?
- which ρ , T and Y_e conditions occur in and after BNS mergers?
- \Box can we use parametrized ν annihilation rates?
- \Box which uncertainties affect ν annihilation process?

Conclusions

neutrinos play an important role in BNS mergers

- they set the properties of the ejecta and potentially of the EM couterparts
- they can drive an outflow (ν-driven wind)
- they can extend nucleosynthesis from I to III r-process peak
- they can inject significant energy in the funnel (MNS can double energy)
- many improvements to be done (e.g., transport, v oscillations)
- large parameter space exploration still required

- calibrated against 1D models with Boltzmann neutrino transport (AGILE-BOLTZTRAN)
 e.g. Liebendörfer+2004
- 15 M_{\odot} progenitor as reference case

Woosley+2002

► calibrated parameters tested for two different progenitor (12 M_☉ and 40 M_☉) with satisfactory results

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Flexibility = successful implementation in different codes

- ASL + 2D Eulerian code
- ASL + 3D SPH code

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Electromagnetic transient

 γ emission powered by radioactive material in the ejecta

Lanthanides and Actinides mass fraction,

$$\kappa_{A>130}\approx 10\,{\rm cm}^2/{\rm g}\sim 10\kappa_{Ni}$$

 1D model for photon propagation and emission

e.g. Kulkarni 05, Grossman+14

- potentially different from emission coming from dynamical/viscous ejecta
 - earlier and bluer
 - less contaminated by lanthanides and actinides

cf Metzger&Fernandez14

 possible dependence from viewing angle