Neutron Star Constraints From Mergers and Cold Atoms

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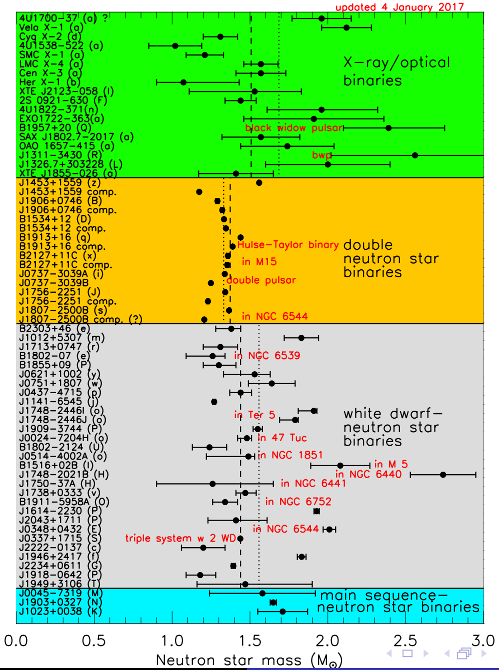
Outline

- The Unitary Gas Constraint on the Nuclear Symmetry Energy
- Neutron Star Universal Structure Relations
- Using Tidal Deformabilities to Infer the Equation of State
- Universal Structure Relations for Hybrid Stars



vanKerkwijk 2010

Romani et al. 2012

Although simple average mass of w.d. companions is 0.23 $M_\odot$ larger, weighted average is 0.07 $M_\odot$ larger

Demorest et al. 2010
Fonseca et al. 2016
Antoniadis et al. 2013
Barr et al. 2016

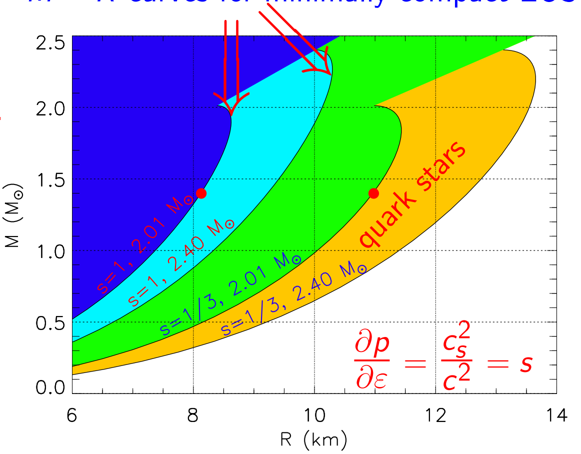
Causality + GR Limits and the Maximum Mass

A lower limit to the maximum mass sets a lower limit to the radius for a given mass.

Similarly, a precision upper limit to R, with a well-measured mass, sets an upper limit to the maximum mass.

$R_{1.4} > 8.15$ km if $M_{max} \geq 2.01 M_\odot$.

$M - R$ curves for minimally compact EOS



If quark matter exists in the interior, the minimum radii are substantially larger.

The Unitary Gas

The **unitary gas** is an idealized system consisting of fermions interacting via a pairwise zero-range s-wave interaction with an infinite scattering length:

As long as the scattering length $a >> k_F^{-1}$ (interparticle spacing), and the range of the interaction $R << k_F^{-1}$, the properties of the gas are universal in the sense they don't depend on the details of the interaction.

The sole remaining length scale is $k_F = (3\pi^2 n)^{1/3}$, so the unitary gas energy is a constant times the Fermi energy $\hbar^2 k_F^2/(2m)$:

$$E_{\rm UG} = \xi_0 \frac{3\hbar^2 k_F^2}{10m}.$$

$\xi_0 \simeq 0.37$ is known as the **Bertsch** parameter, measured in cold-atom experiments.

The Unitary Gas as Analogue of the Neutron Gas

A pure neutron matter (PNM) gas differs from the unitary gas:

- $|a| \simeq 18.5$ fm; $|ak_F|^{-1} \simeq 0.03$ for $n = n_s$.
- $R \simeq 2.7$ fm; $Rk_F \approx 4.5$ for $n = n_s$.
- Repulsive 3-body interactions are additionally necessary for neutron matter to fit the energies of light nuclei.
- Neutron matter has potentially atractive p-wave and higher-order interactions.

The first three imply $E_{\rm PNM} > E_{\rm UG}$:

- $\xi \simeq \xi_0 + 0.6|ak_F|^{-1} + \ldots \qquad |ak_F|^{-1} << 1$
- $\xi \simeq \xi_0 + 0.12Rk_F + \ldots \qquad Rk_F << 1$

A reasonable conjecture would appear to be ($u = n/n_s$)

$$E_{\rm PNM}(u) = E(u, Y_p = 0) \geq E_{{\rm UG},0}u^{2/3} \simeq 12.6u^{2/3}\ {\rm MeV}$$

Consequences for the Nuclear Symmetry Energy

$$S(u) = E_{\rm PNM} - E(u, Y_p = 1/2).$$

A good approximation for the Y_p−dependence of E is

$$S(u) \simeq \frac{1}{8}\frac{\partial^2 E(u, Y_p)}{\partial Y_p^2}.$$

Near n_s,

$$S(u) \simeq S_0 + \frac{L}{3}(u-1) + \frac{K_{sym}}{18}(u-1)^2 + \cdots$$

$$E(u, Y_p = 1/2) \simeq -B + \frac{K_s}{18}(u-1)^2 + \cdots$$

In this case, the unitary gas conjecture is

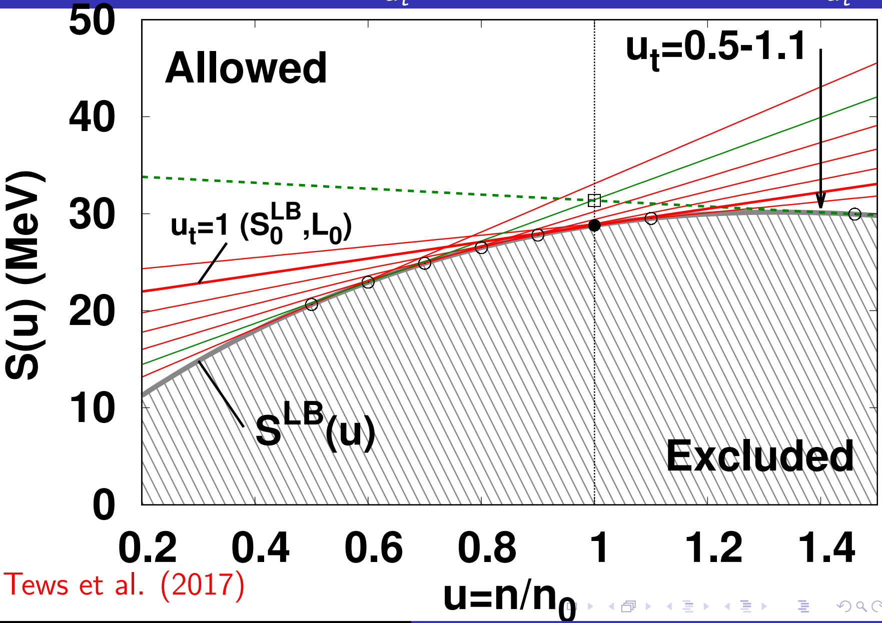
$$S(u) > E_{\rm UG,0}u^{2/3} - \left[-B + \frac{K_s}{18}(u-1)^2 + \cdots\right] \equiv S^{\rm LB}(u)$$

Thus, the symmetry energy parameters S_0 and L must satisfy

$$S(u=1) = S_0 \geq S_0^{\rm LB} = E_{\rm UG,0} + B \simeq 28.5 \text{ MeV}$$

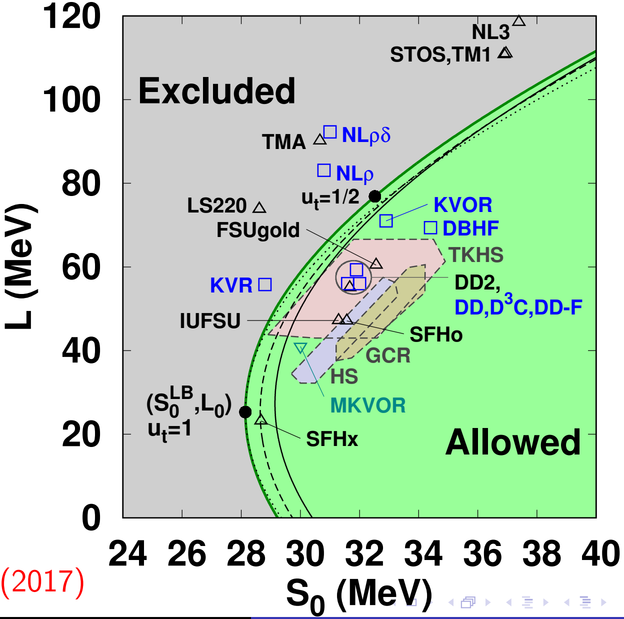
$$L(u=1) = L_0 = 3\left(udS/du\right)_{u=1} = 2E_{\rm UG,0} \simeq 25.2\text{MeV}$$

$S_0 > S_0^{\rm LB}: \quad \left(S = S^{\rm LB}\right)_{u_t}, \left(dS/du = dS^{\rm LB}/du\right)_{u_t}$



Tews et al. (2017)

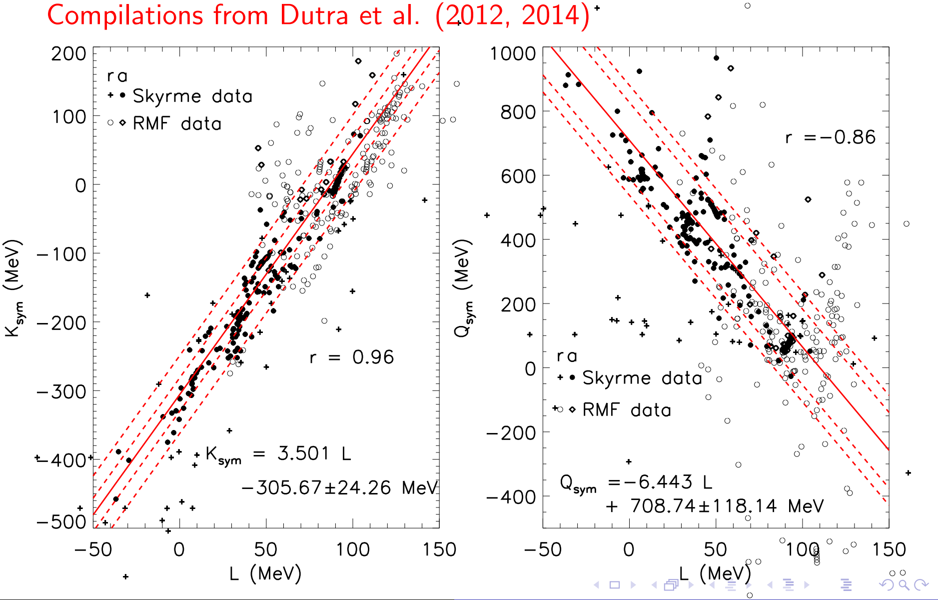
Symmetry Parameter Exclusions



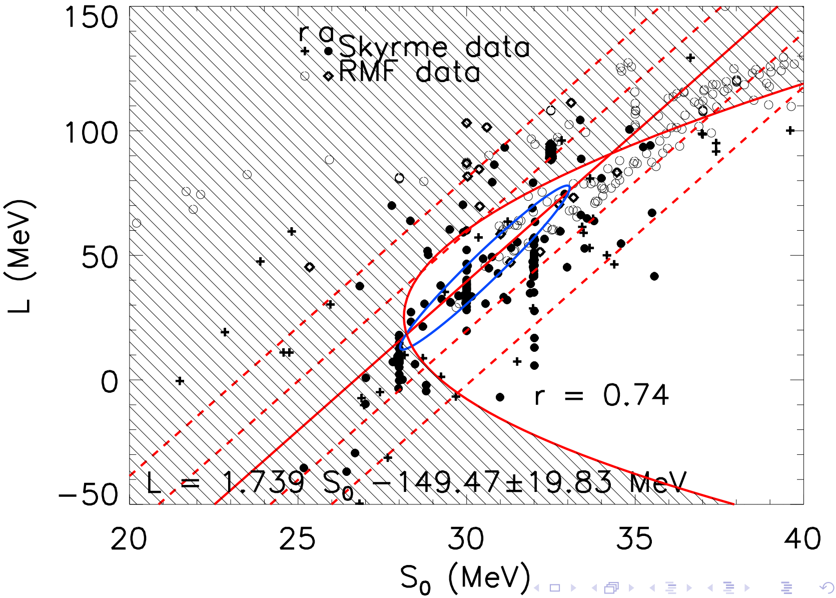
Tews et al. (2017)

Symmetry Parameter Correlations

Compilations from Dutra et al. (2012, 2014)



More Realistic Exclusion Region



Analytic Approximation for the Boundary

$$S(u_t) = S^{\rm LB}(u_t), \qquad \left(\frac{dS}{du}\right)_{u_t} = \left(\frac{dS^{\rm LB}}{du}\right)_{u_t}$$

gives

$$S_0 + \frac{L}{3}(u_t - 1) + \frac{K_{sym}}{18}(u_t - 1)^2 = E_{\rm UG,0}u_t^{2/3} + B - \frac{K_s}{18}(u_t - 1)^2$$

$$L + \frac{K_{sym}}{3}(u_t - 1) = 2E_{\rm UG,0}u_t^{-1/3} - \frac{K_s}{3}(u_t - 1)$$

Assume $K_n = 3L$ (i.e., $K_{\rm sym} \approx 3L - K_s$). Then

$$S_0 = \frac{E_{\rm UG,0}}{3u_t^{4/3}}(1 + 2u_t^2) - E_0, \qquad L = \frac{2E_{\rm UG,0}}{u_t^{4/3}}$$

or after eliminating u_t,

$$S_0 = \frac{L}{6}\left[1 + 2\left(\frac{2E_{\rm UG,0}}{L}\right)^{3/2}\right] - E_0$$

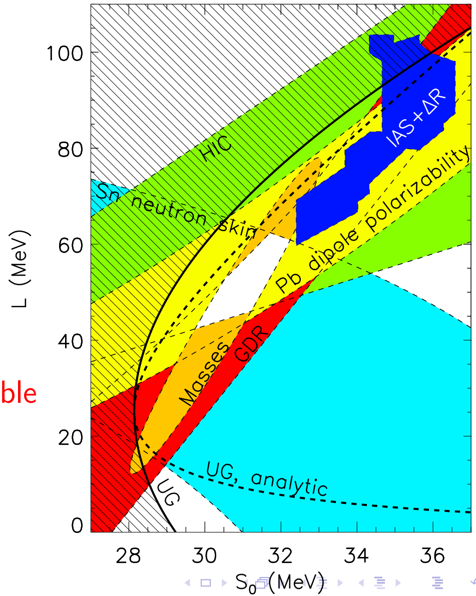
Experimental Constraints

Isovector Skins and
Isobaric Analog States
from Danielewicz et al. (2017)

Other experimental constraints
from Lattimer & Lim (2013)

Unitary gas constraints from
Tews et al. (2017)

Experimental and neutron
matter constraints are compatible
with unitary gas bounds.



Piecewise Polytropes

Crust EOS is known: $n < n_0 = 0.4 n_s$.

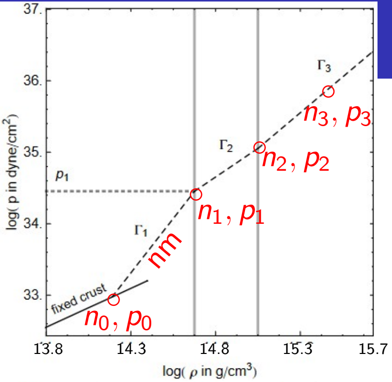
Read, Lackey, Owen & Friedman (2009) found high-density EOS can be modeled as piecewise polytropes with 3 segments.

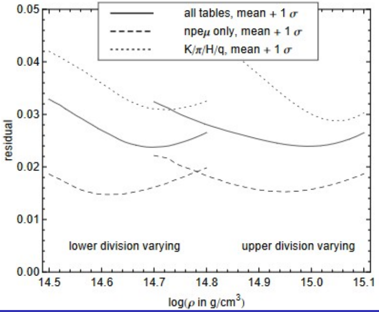
They found universal break points ($n_1 \simeq 1.85 n_s$, $n_2 \simeq 3.7 n_s$) optimized fits to a wide family of modeled EOSs.

For $n_0 < n < n_1$, assume neutron matter EOS. Arbitrarily choose $n_3 = 7.4 n_s$.

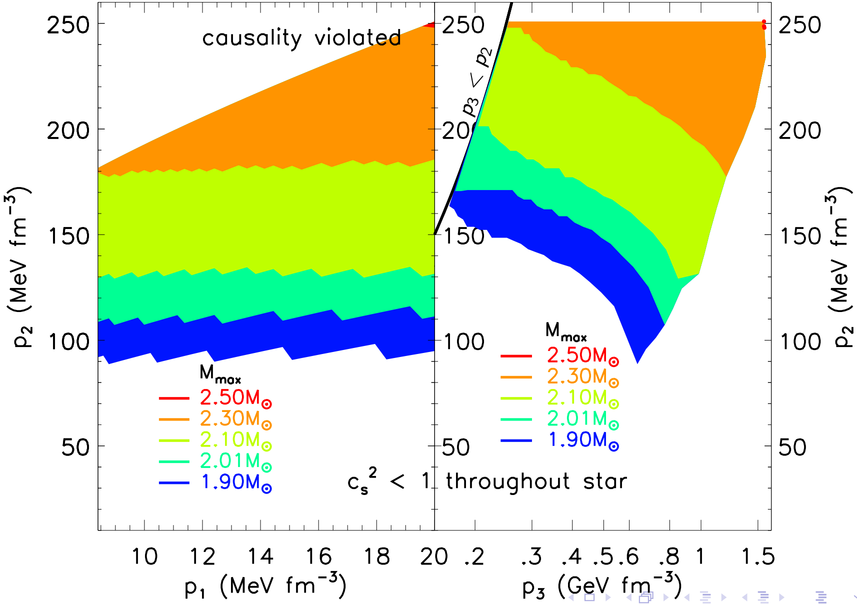
For a given p_1 (or Γ_1):
$0 < \Gamma_2 < \Gamma_{2c}$ or $p_1 < p_2 < p_{2c}$.
$0 < \Gamma_3 < \Gamma_{3c}$ or $p_2 < p_3 < p_{3c}$.

Minimum values of p_2, p_3 set by M_{max}; maximum values set by causality.

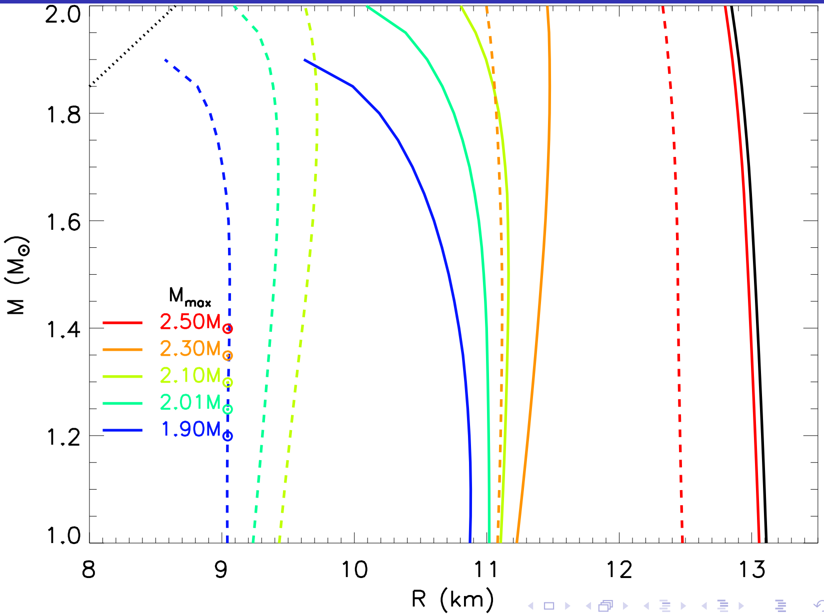




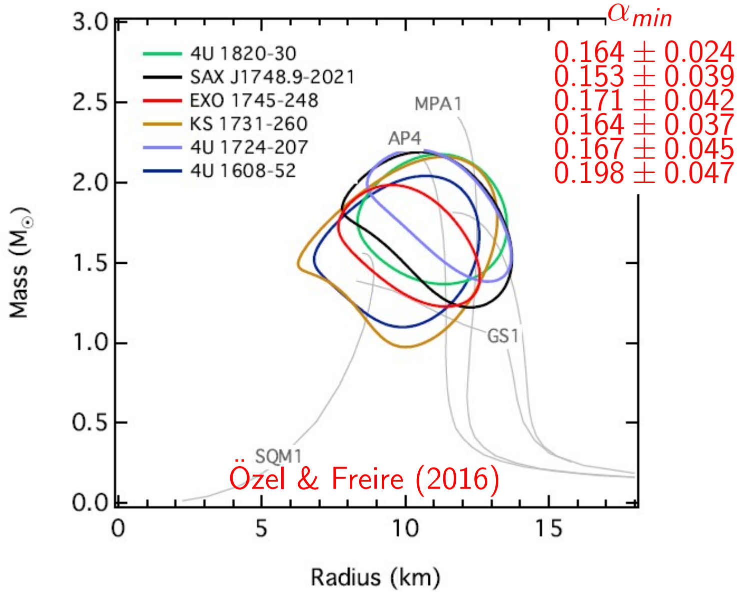
Maximum Mass and Causality Constraints



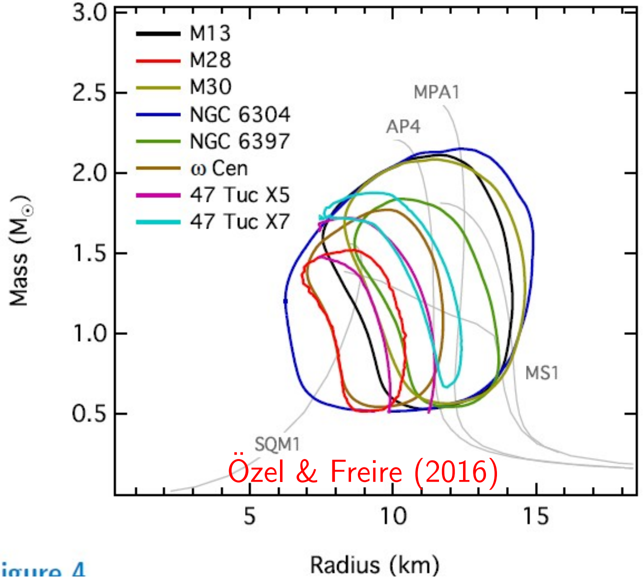
Mass-Radius Constraints from Causality



PRE $M - R$ Estimates



QLMXB $M - R$ Estimates



Özel & Freire (2016)

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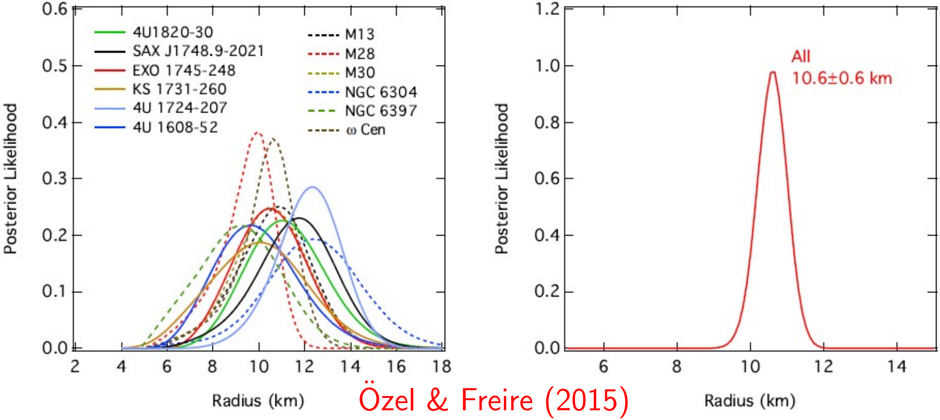
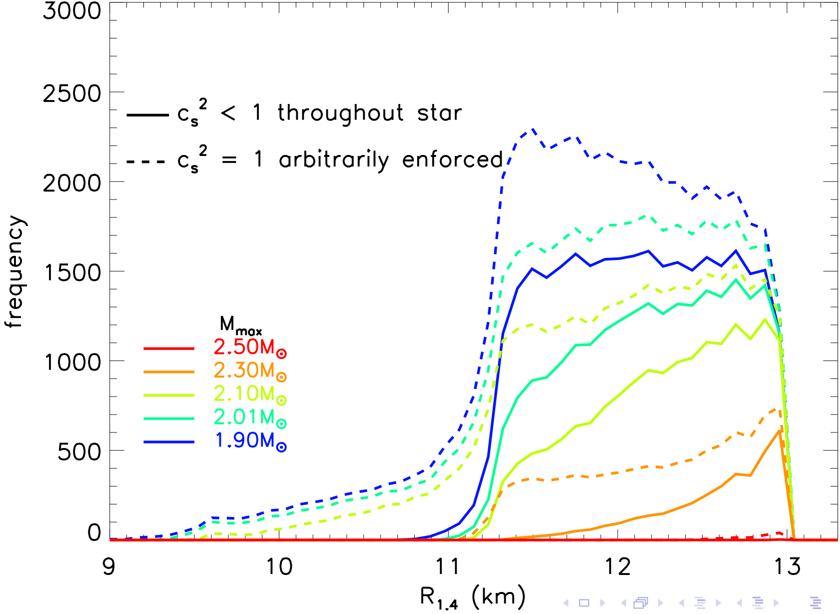
Combined R fits

Assume $P(M)$ is that measured from pulsar timing ($\bar{M} = 1.4 M_\odot$).

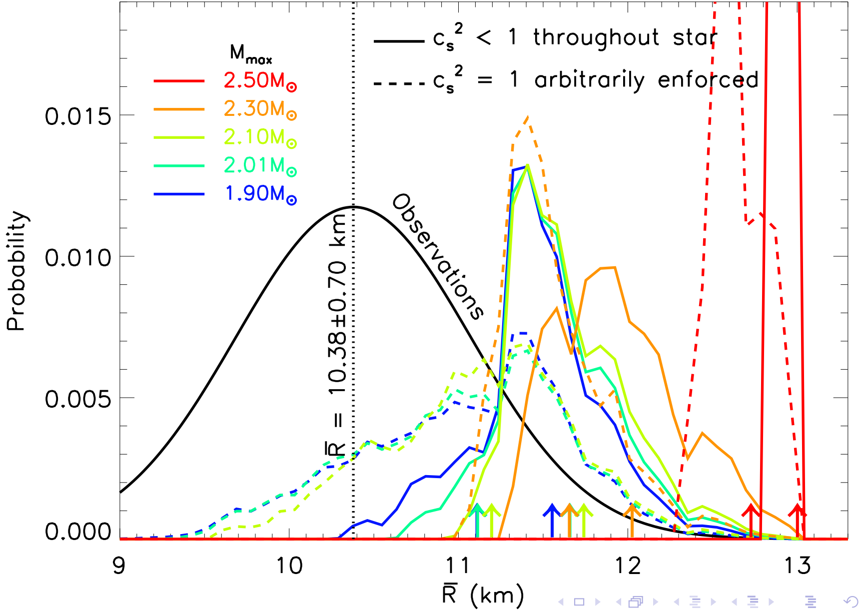


Özel & Freire (2015)

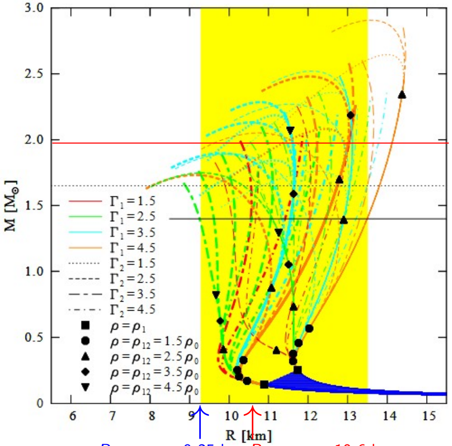
Piecewise-Polytrope Average Radius Distributions



Folding Observations with Piecewise Polytropes



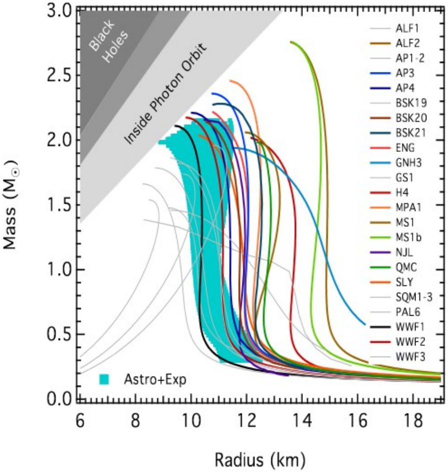
Other Studies



$R_{1.4,min} = 9.25$ km
$M_{max} = 1.66 M_\odot$

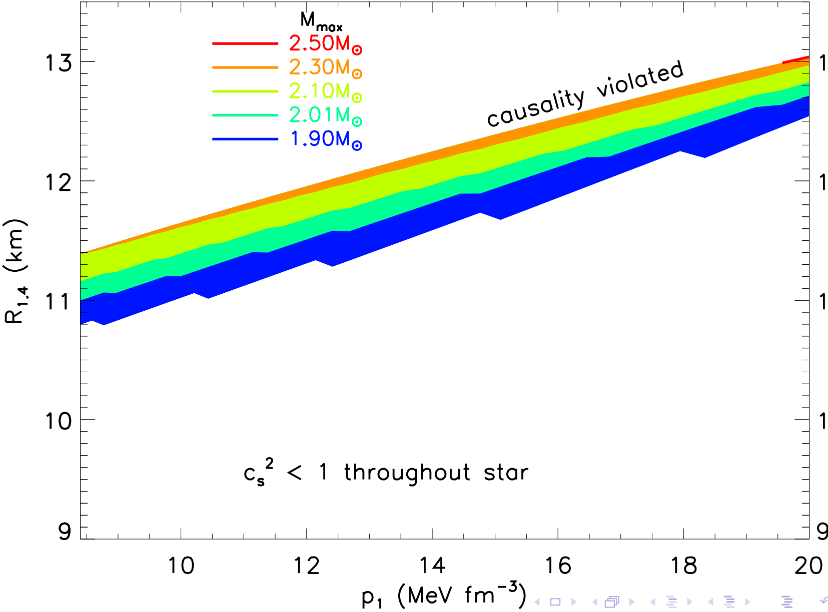
$R_{1.4,min} = 10.6$ km
$M_{max} = 1.97 M_\odot$

Hebeler, Lattimer, Pethick & Schwenk 2010

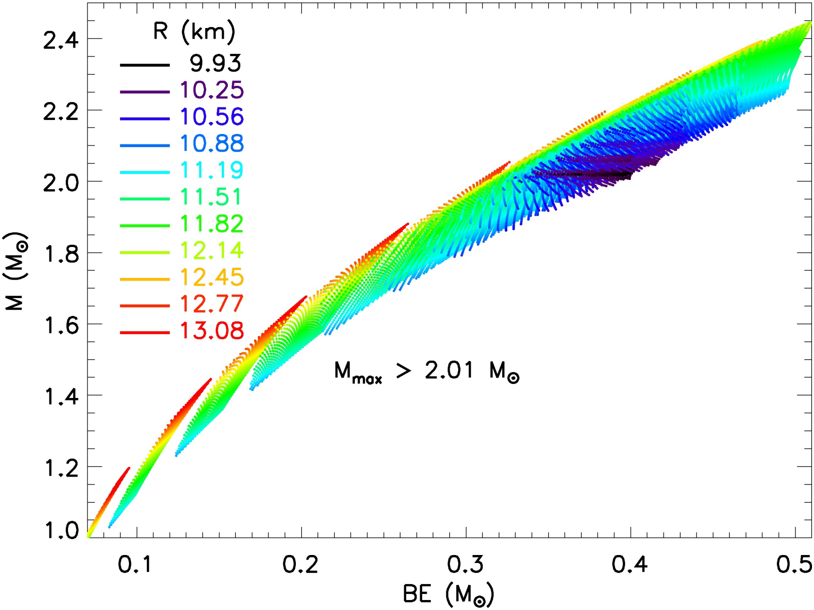


Özel & Freire 2016

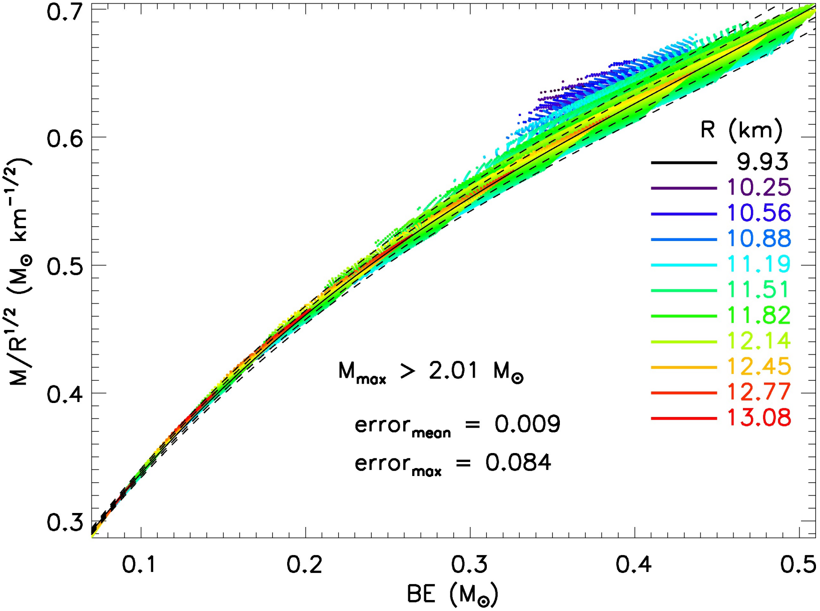
Radius - p_1 Correlation



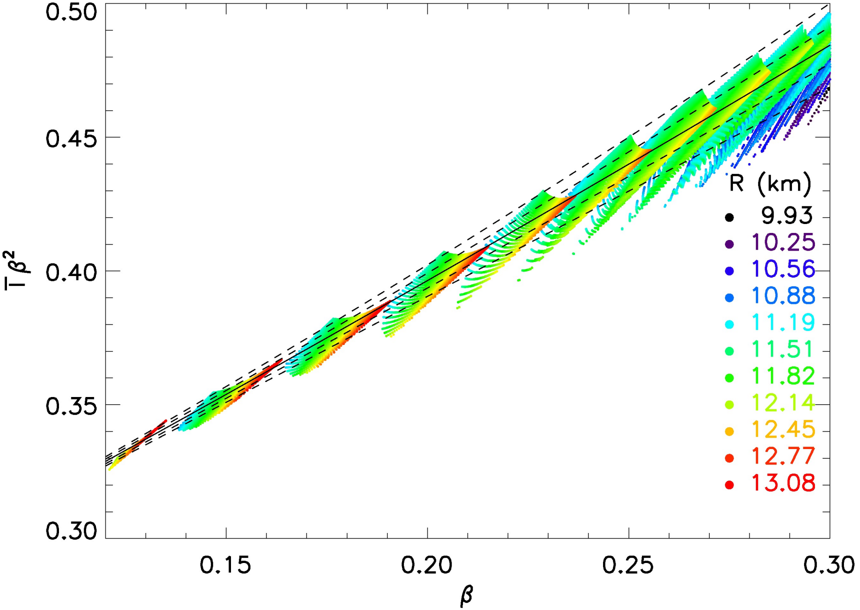
Binding Energy - Mass Correlations



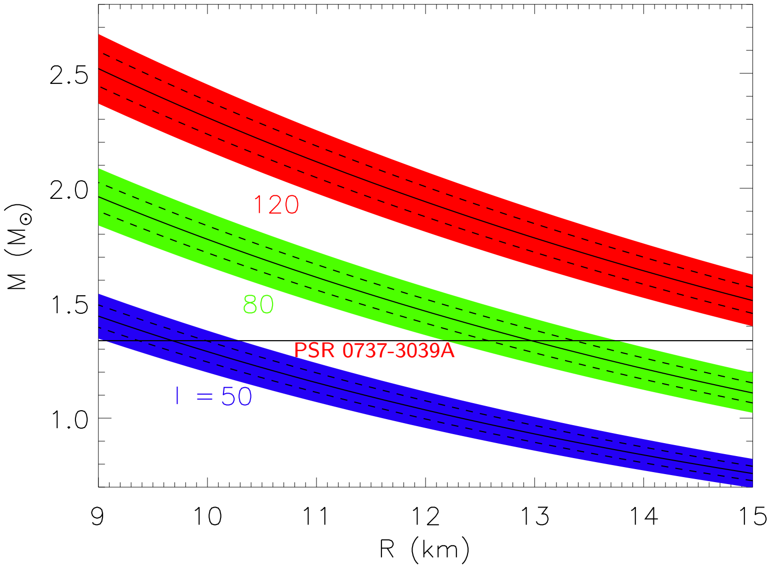
Binding Energy - Mass - Radius Correlations



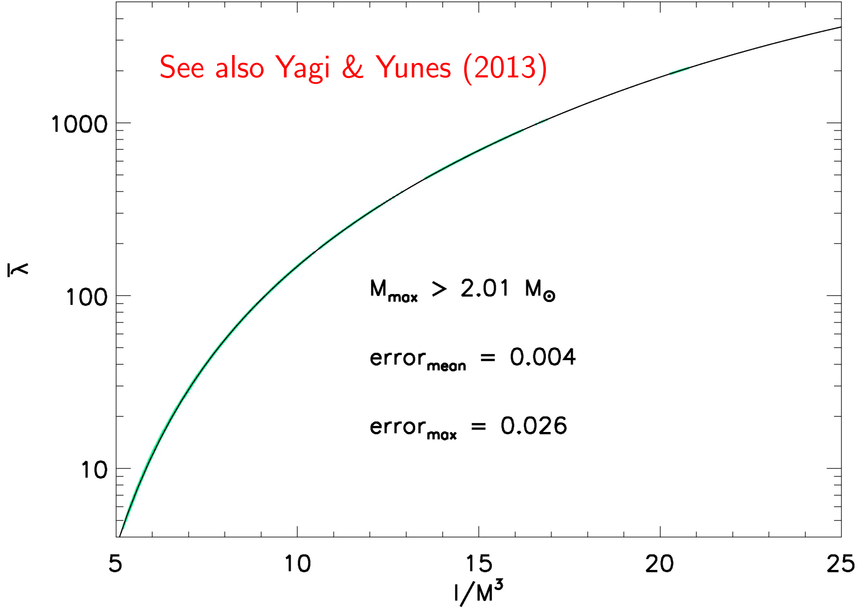
Moment of Inertia - Mass - Radius Correlations



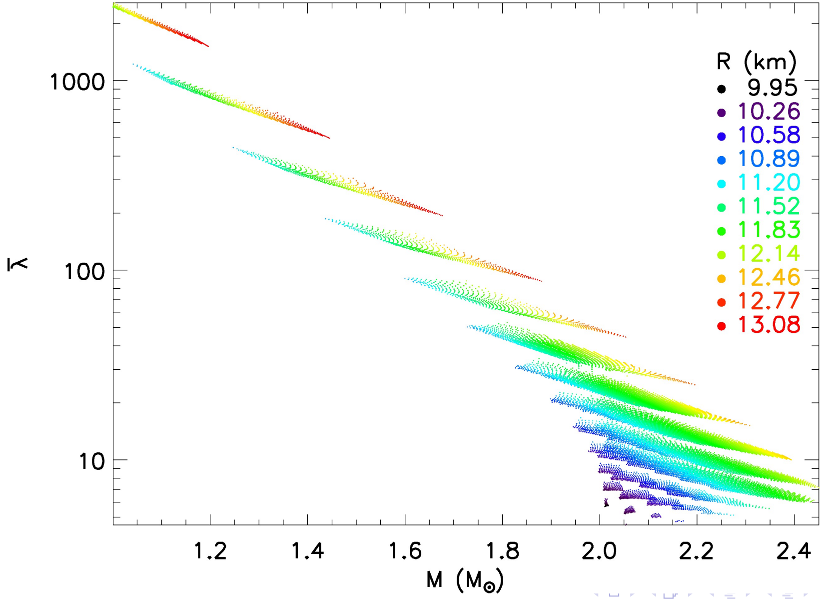
Moment of Inertia - Radius Constraints



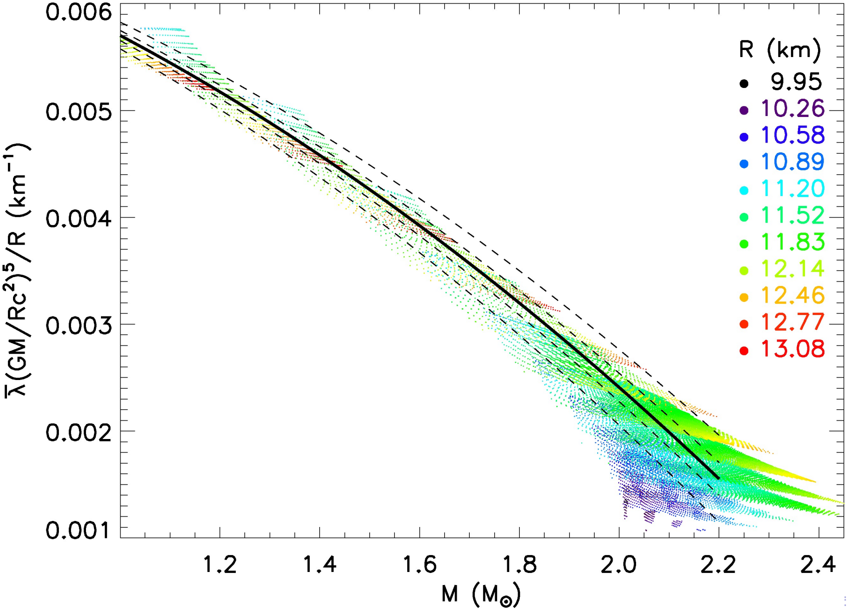
Tidal Deformatibility - Moment of Inertia



Tidal Deformatibility - Mass



Tidal Deformatibility - Mass - Radius



Binary Tidal Deformability

In a neutron star merger, both stars are tidally deformed. The most accurately measured deformability parameter is

$$\bar{\Lambda} = \frac{16}{13}\left[\bar{\lambda}_1 q^4(12q+1) + \bar{\lambda}_2(1+12q)\right]$$

where

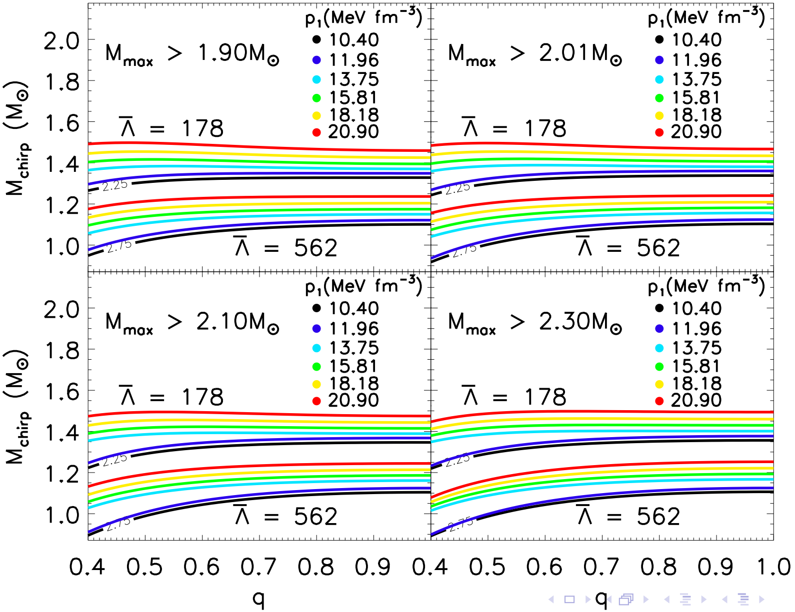
$$q = \frac{M_1}{M_2} < 1$$

For $S/N \approx 20 - 30$, typical measurement accuracies are expected to be (Rodriguez et al. 2014; Wade et al. 2014):

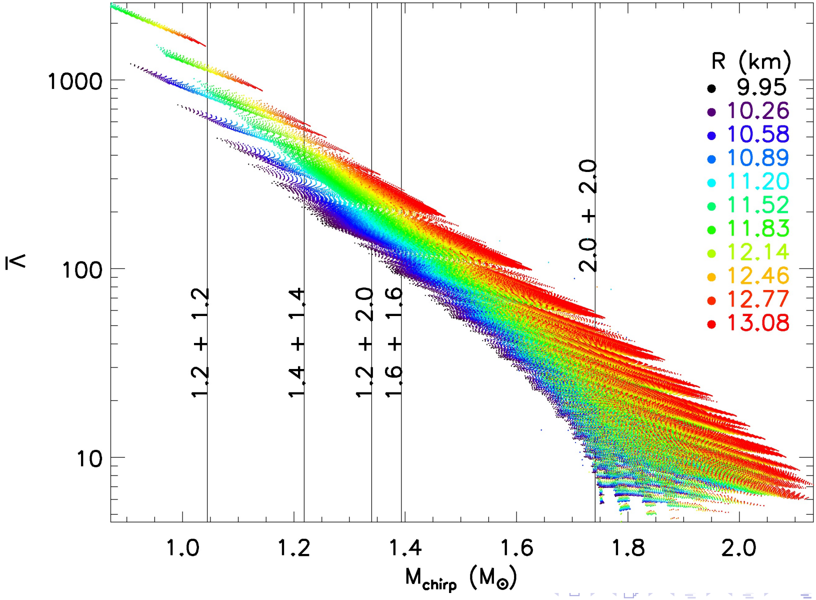
$$\Delta M_{chirp} \sim 0.01 - 0.02\%, \qquad \Delta\bar{\Lambda} \sim 20 - 25\%$$

$$\Delta(M_1 + M_2) \sim 1 - 2\%, \qquad \Delta q \sim 10 - 15\%$$

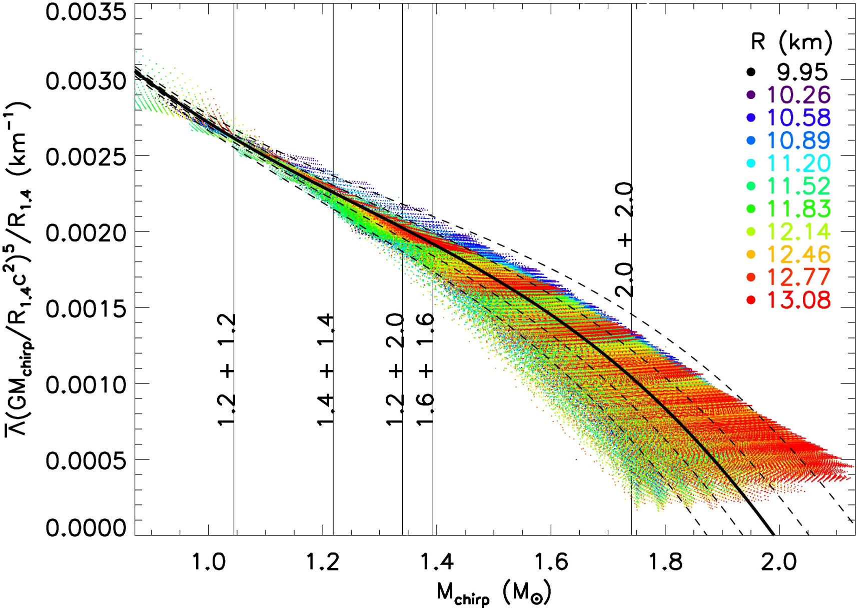
Binary Tidal Deformatibility - $\bar{\Lambda}$



Binary Tidal Deformatibility - $M_{\rm chirp}$

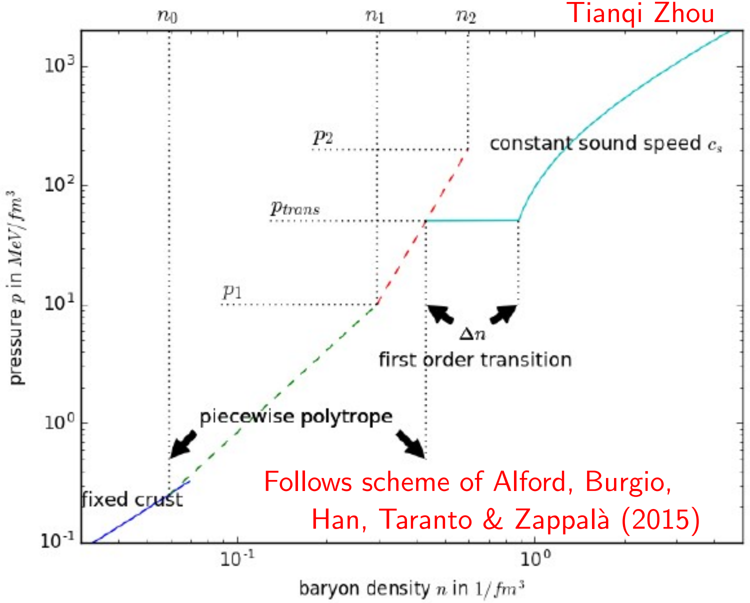


Binary Tidal Deformatibility - $M_{\rm chirp}$ - $R_{1.4}$



Hybrid Stars With First-Order Phase Transition

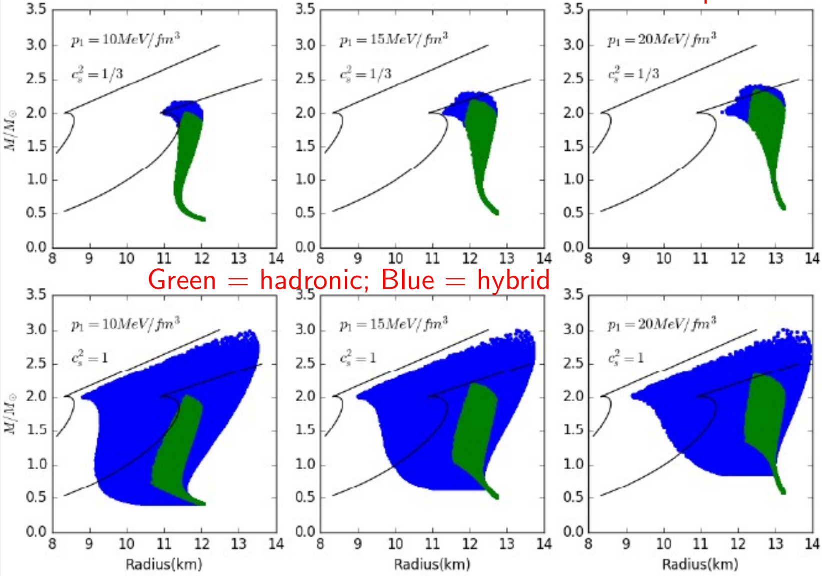
Tianqi Zhou



Follows scheme of Alford, Burgio, Han, Taranto & Zappalà (2015)

Mass-Radius Comparisons

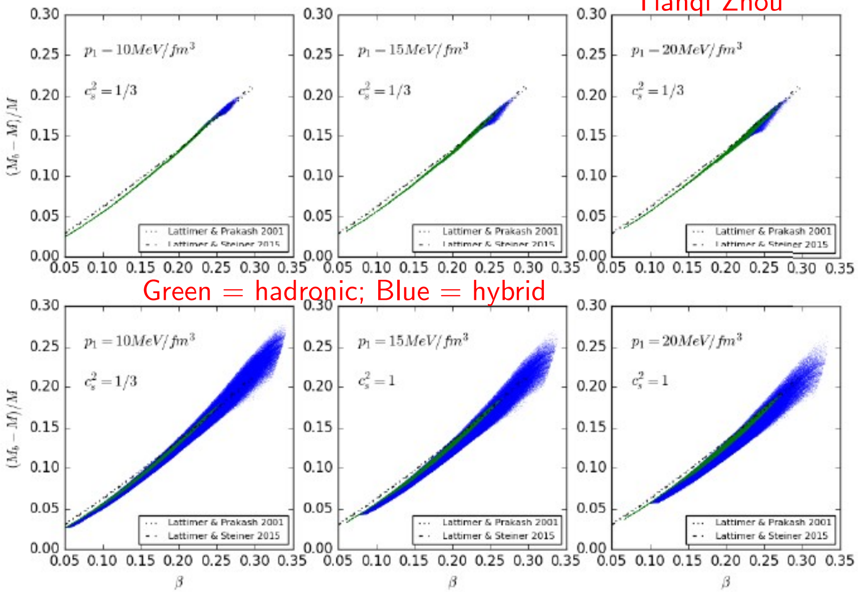
Tianqi Zhou



Green = hadronic; Blue = hybrid

Binding Energy - Compactness Comparisons

Tianqi Zhou



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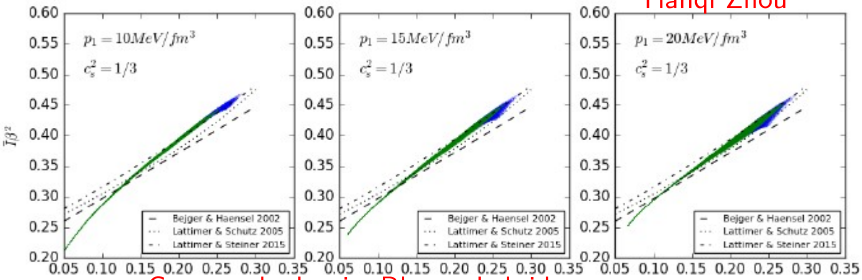
Moment of Inertia - Compactness Comparisons

Tianqi Zhou



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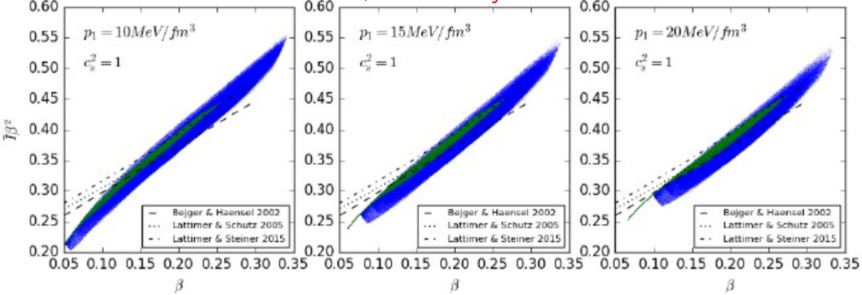

Tidal Deformability Comparisons