

TeV Scale LNV: $0\nu\beta\beta$ -Decay & the LHC

M.J. Ramsey-Musolf

U Mass Amherst

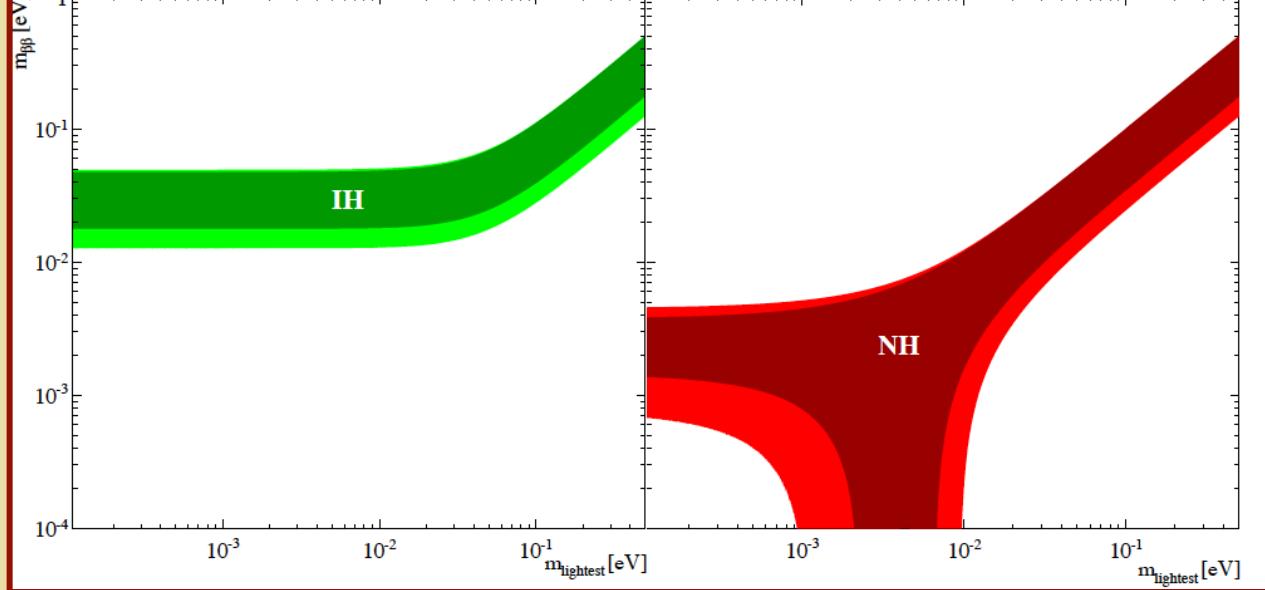


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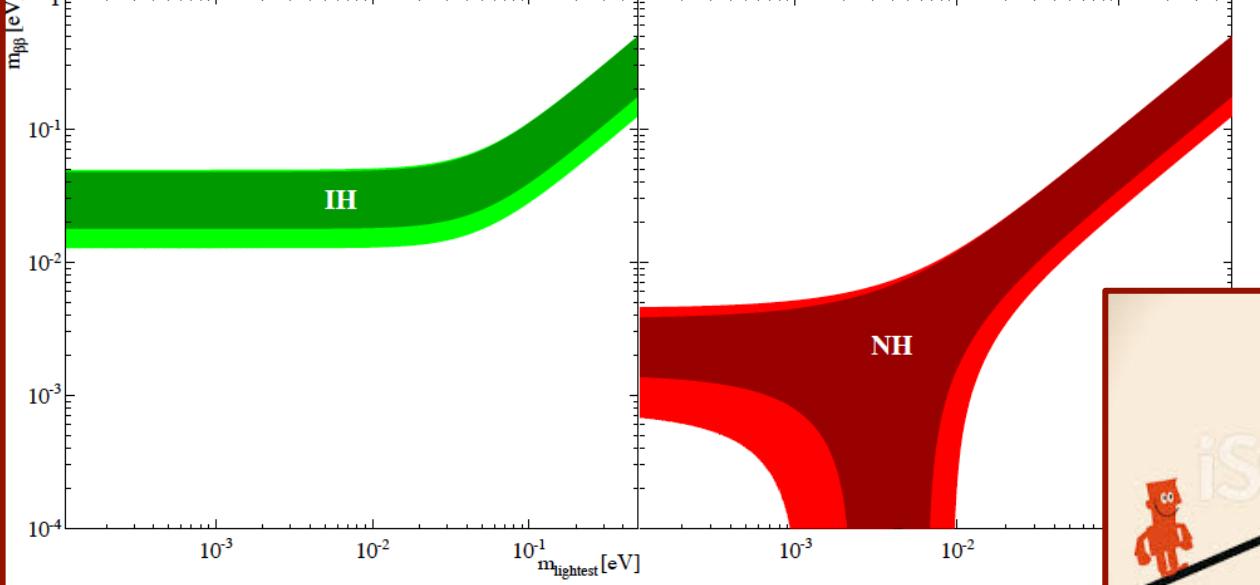
Collaborators: Tao Peng, Peter Winslow; V. Cirigliano, M. Graesser, M. Horoi, P. Vogel

INT $0\nu\beta\beta$ -Decay Workshop
June 2017

This talk: beyond the “poster child”



This talk: beyond the “poster child”



Themes for This Talk

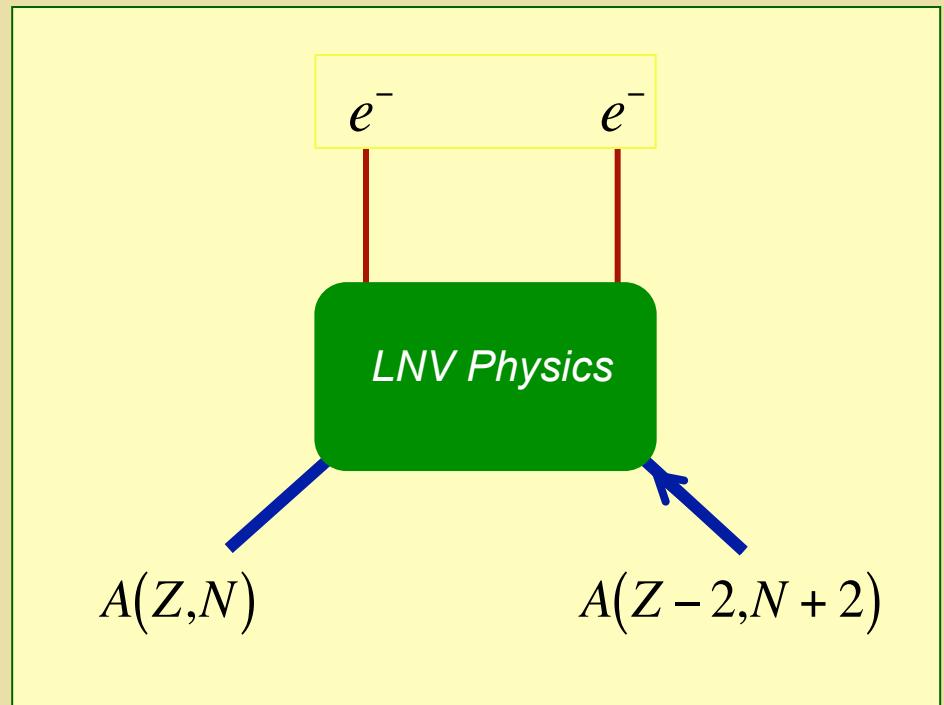
$0\nu\beta\beta$ -Decay: LNV? Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana



$0\nu\beta\beta$ -Decay: LNV? Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

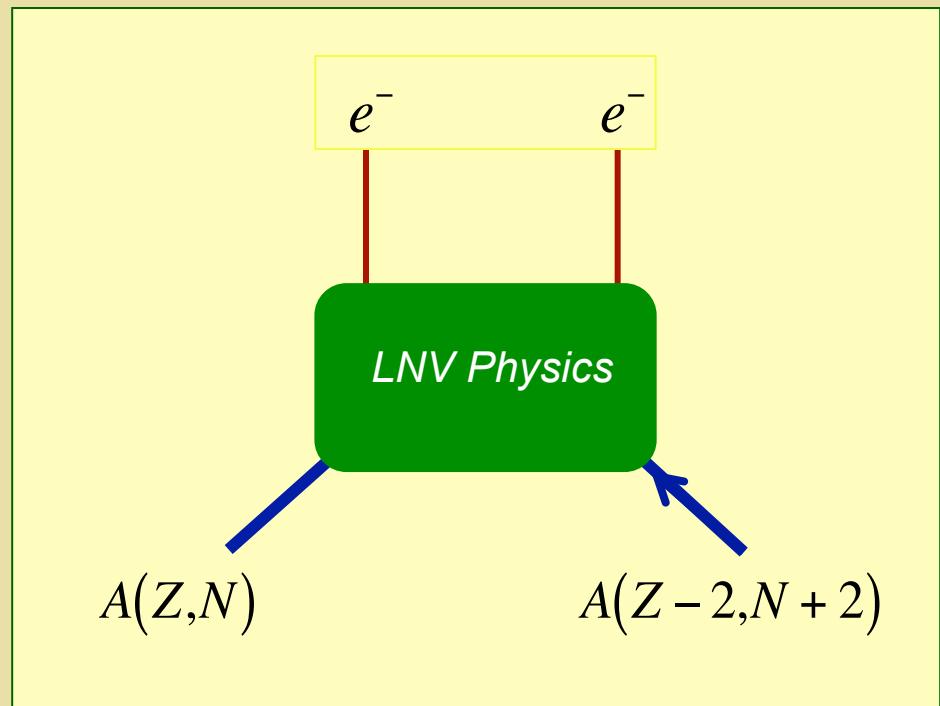
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$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Impact of observation

- Total lepton number not conserved at classical level
- New mass scale in nature, Λ
- Key ingredient for standard baryogenesis via leptogenesis

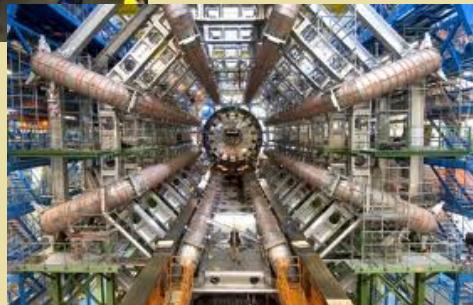


LNV: Discoverable at the Energy Frontier

LHC

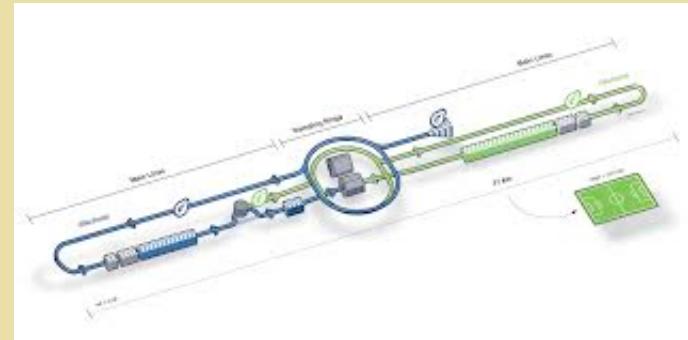


CMS

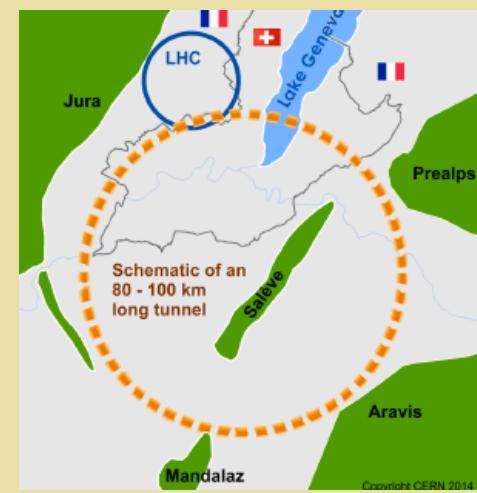
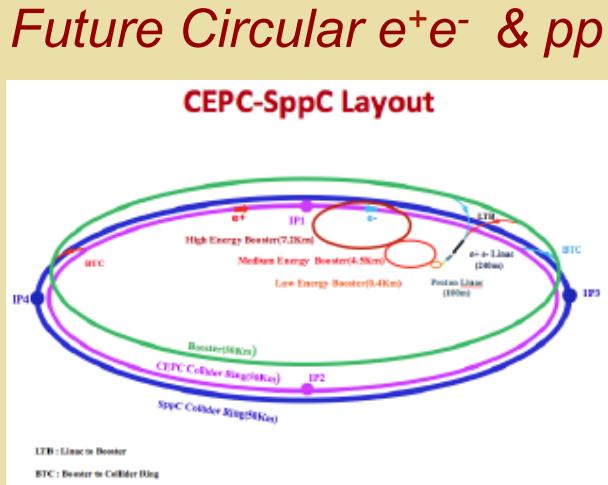


ATLAS

International Linear Collider



Future Circular e^+e^- & pp



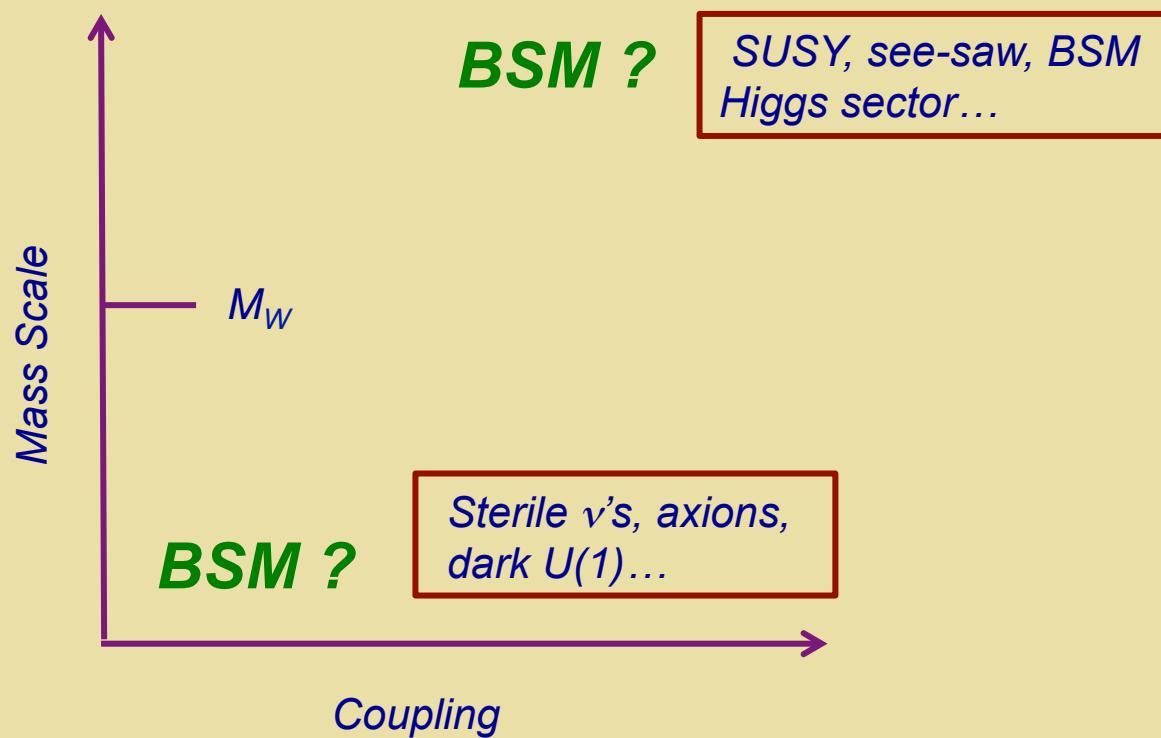
Thanks:
S.
Gascon-
Shotkin

Outline

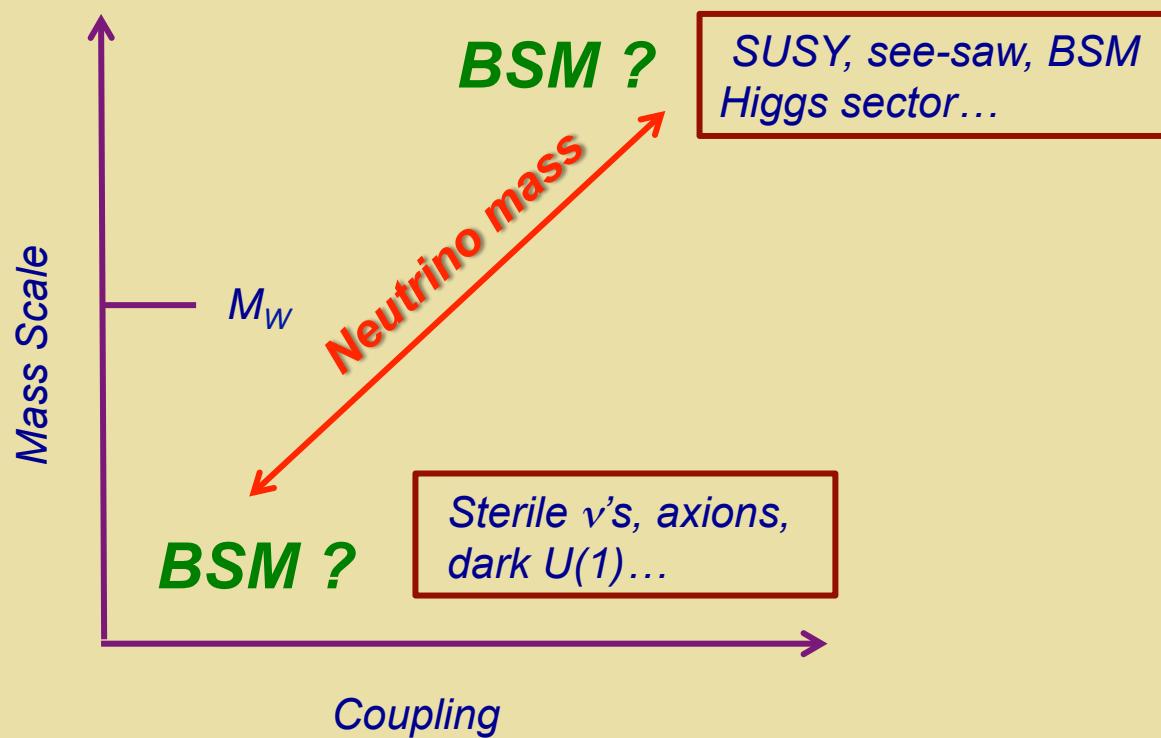
- I. The BSM Context*
- II. High Scale LNV*
- III. TeV Scale LNV*
- IV. Sub Weak Scale LNV (time permitting)*
- V. Summary*

I. The BSM Context

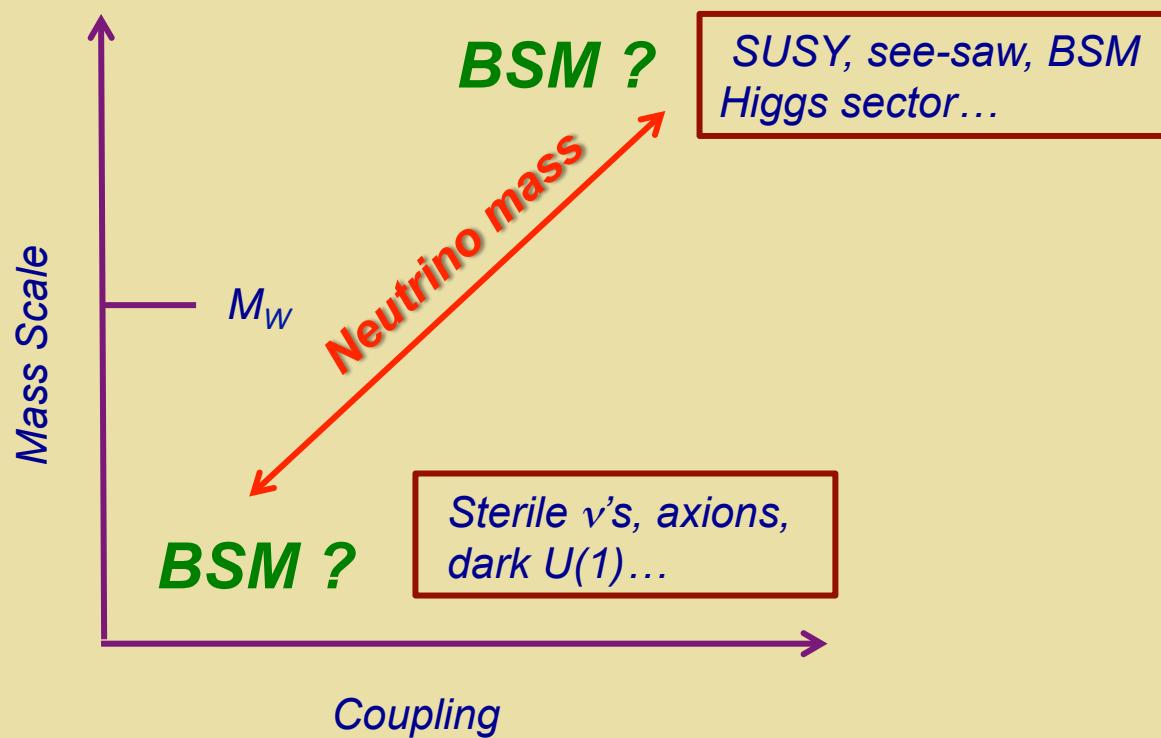
BSM Physics: Where Does it Live ?



BSM Physics: Where Does it Live ?



BSM Physics: Where Does it Live ?



Is the mass scale associated with m_ν far above M_W ? Near M_W ? Well below M_W ?

LNV Mass Scale & $0\nu\beta\beta$ -Decay



- *3 light neutrinos only: source of neutrino mass at the very high see-saw scale*
- *3 light neutrinos with TeV scale source of neutrino mass*
- *> 3 light neutrinos*

II. High Scale LNV

LNV Mass Scale & $0\nu\beta\beta$ -Decay

$$A(Z, N) \rightarrow \text{Underlying Physics} \rightarrow A(Z+2, N-2) + e^- e^-$$

- *3 light neutrinos only: source of neutrino mass at the very high see-saw scale*
- *3 light neutrinos with TeV scale source of neutrino mass*
- *> 3 light neutrinos*

$0\nu\beta\beta$ -Decay: LNV? Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

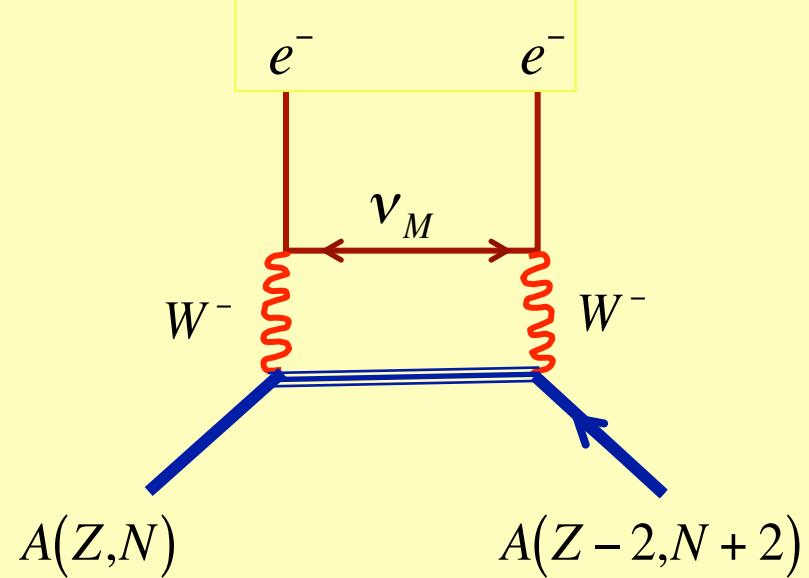
Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

“Standard” Mechanism

- Light Majorana mass generated at the conventional see-saw scale: $\Lambda \sim 10^{12} - 10^{15}$ GeV
- 3 light Majorana neutrinos mediate decay process



Details

Essentially the rest of this workshop...

III. TeV Scale LNV

LNV Mass Scale & $0\nu\beta\beta$ -Decay



- *3 light neutrinos only: source of neutrino mass at the very high see-saw scale*
- *3 light neutrinos with TeV scale source of neutrino mass*
- *> 3 light neutrinos*

Two parameters: Effective coupling & effective heavy particle mass

$0\nu\beta\beta$ -Decay: LNV? Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

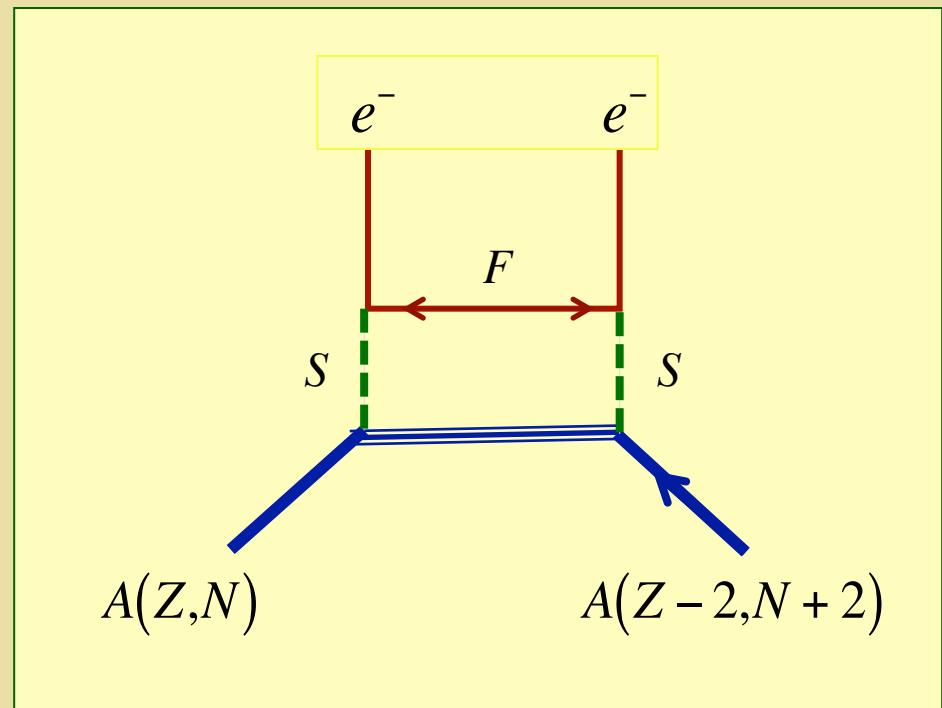
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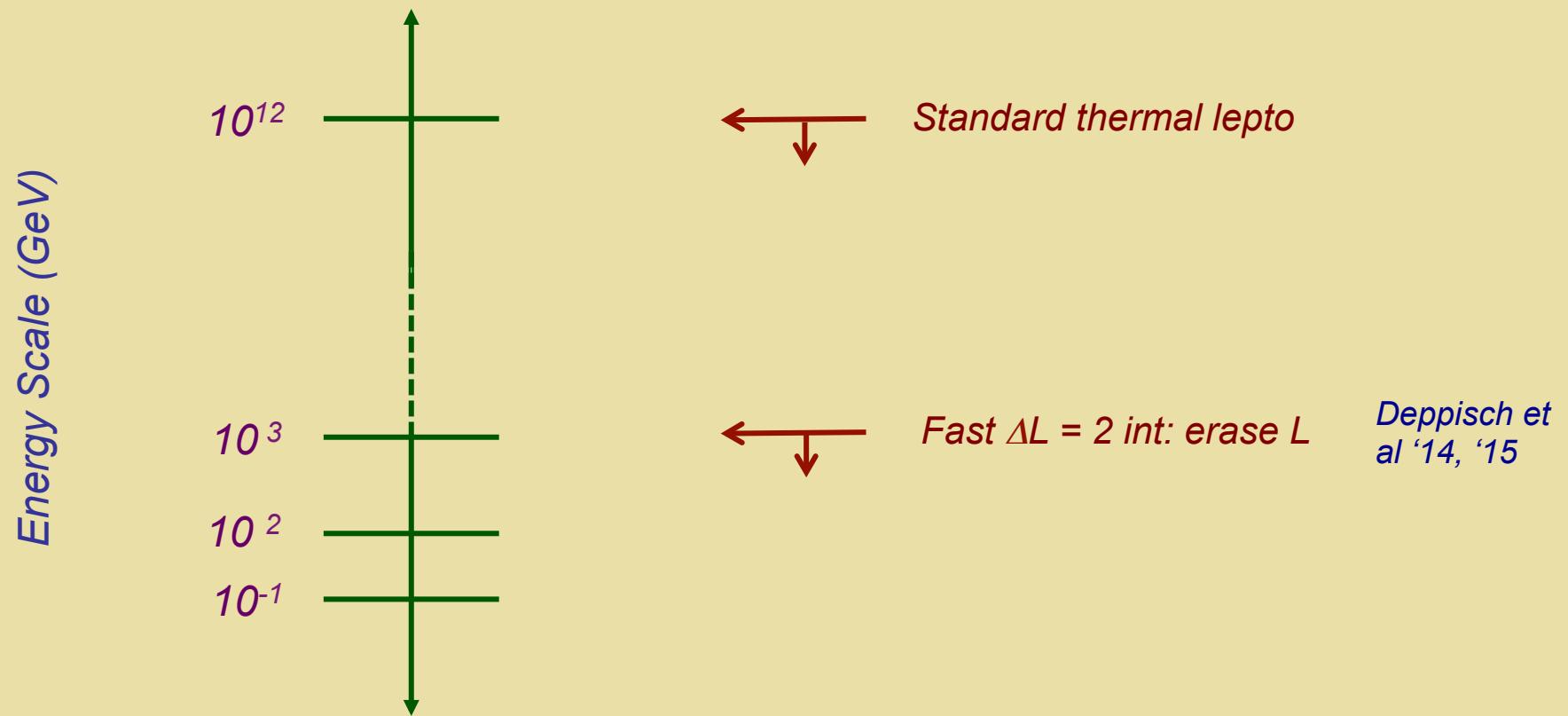
Majorana

TeV LNV Mechanism

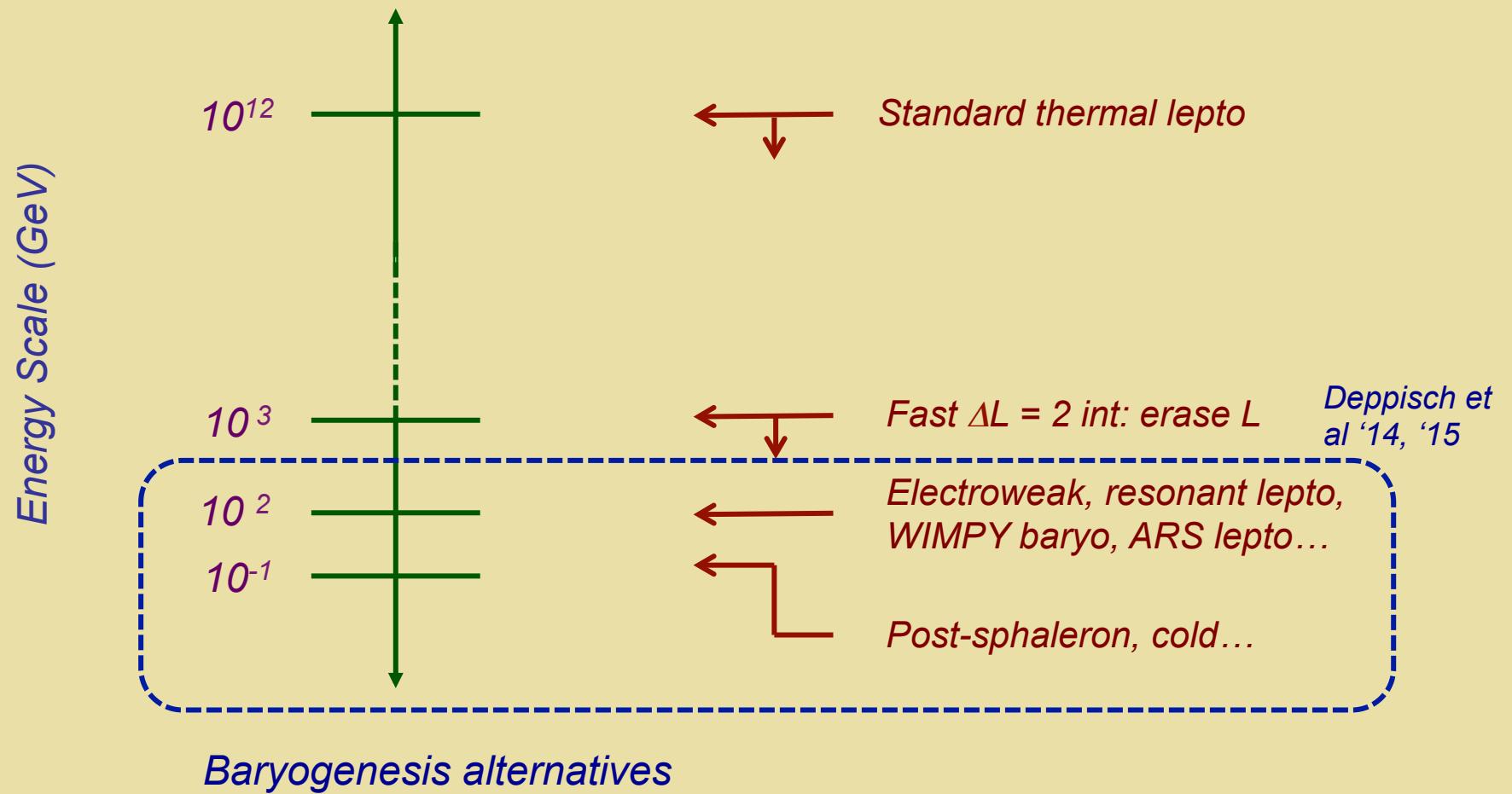
- Majorana mass generated at the TeV scale
 - Low-scale see-saw
 - Radiative m_ν
- $m_{\text{MIN}} \ll 0.01 \text{ eV}$ but $0\nu\beta\beta$ -signal accessible with tonne-scale exp'ts due to heavy Majorana particle exchange



TeV LNV & Leptogenesis



TeV LNV & Leptogenesis



$0\nu\beta\beta$ -Decay: TeV Scale LNV

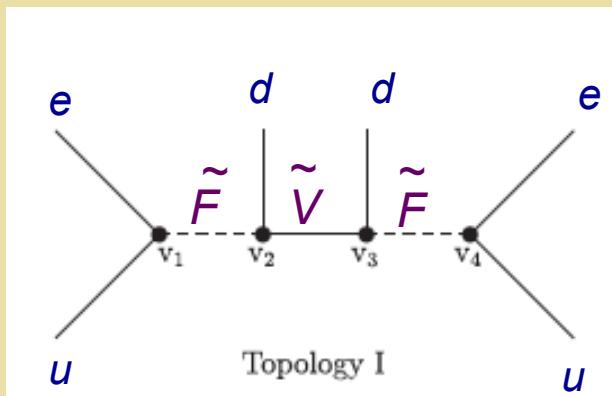
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$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

General Classification: Helo et al, PRD 88.011901, 88.073011



SUSY: R Parity-Violation

Sfermion \tilde{q}, \tilde{T}

Gaugino \tilde{g}, χ Majorana

$$W_{\Delta L=1} = \frac{1}{2} \lambda_{ijk} L_i L_j \bar{e}_k + \lambda'_{ijk} L_i Q_j \bar{d}_k + \mu'_i L_i H_u,$$

$0\nu\beta\beta$ -Decay: TeV Scale LNV

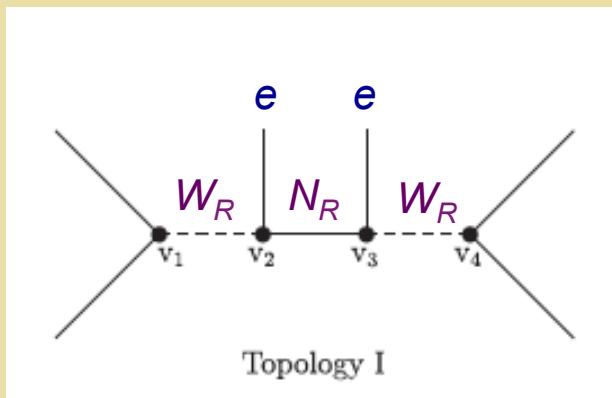
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LRSM: Type I See-Saw

Mass: standard see-saw but TeV scale

$0\nu\beta\beta$ -Decay: TeV Scale LNV

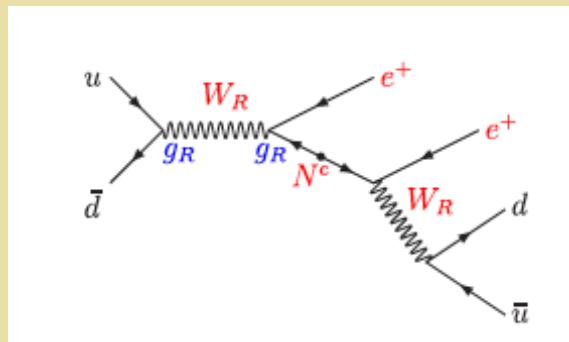
$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

$$\mathcal{L}_{\text{mass}} + \frac{y_{\text{Dirac}}}{\Lambda} \bar{L} H H^T L + \text{h.c.}$$

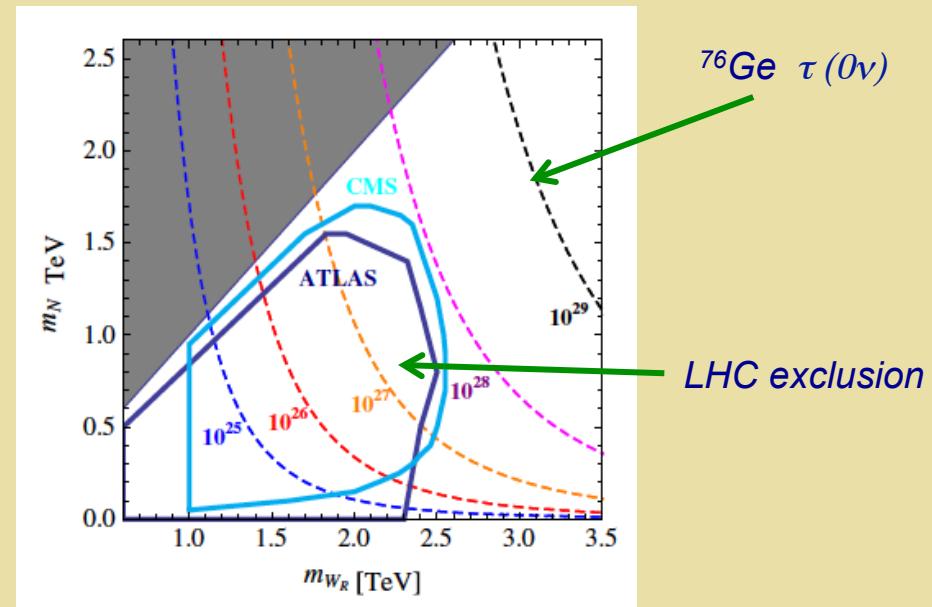
Dirac

Majorana

LHC Production & $0\nu\beta\beta$ -Decay



Helo et al, PRD 88.011901,
88.073011



$0\nu\beta\beta$ -Decay: TeV Scale LNV

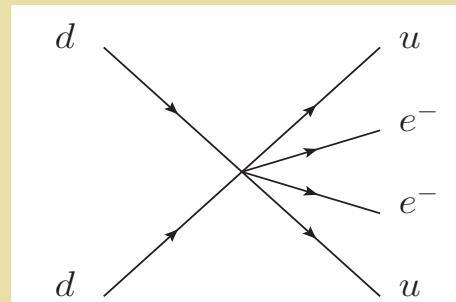
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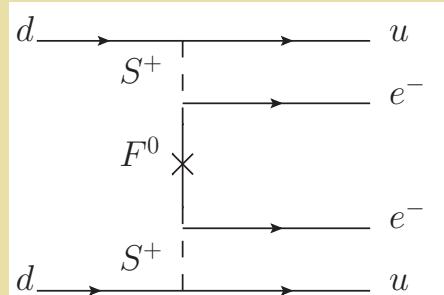
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Majorana

$0\nu\beta\beta$ - decay



LHC: $pp \rightarrow jj e^- e^-$



TeV Scale LNV

*Can it be discovered
with combination of
 $0\nu\beta\beta$ & LHC searches ?*

Simplified models

Simplified Models: Illustrative Case

$$\mathcal{L}_{\text{INT}} = g_1 \bar{Q}_i^a d^a S_i + g_2 \epsilon^{ij} \bar{L}_i F S_j^* + \text{H.c.}$$

$S:$ $(1, 2, \frac{1}{2})$
 $F:$ $(1, 0, 0)$ Majorana

$0\nu\beta\beta$ -Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

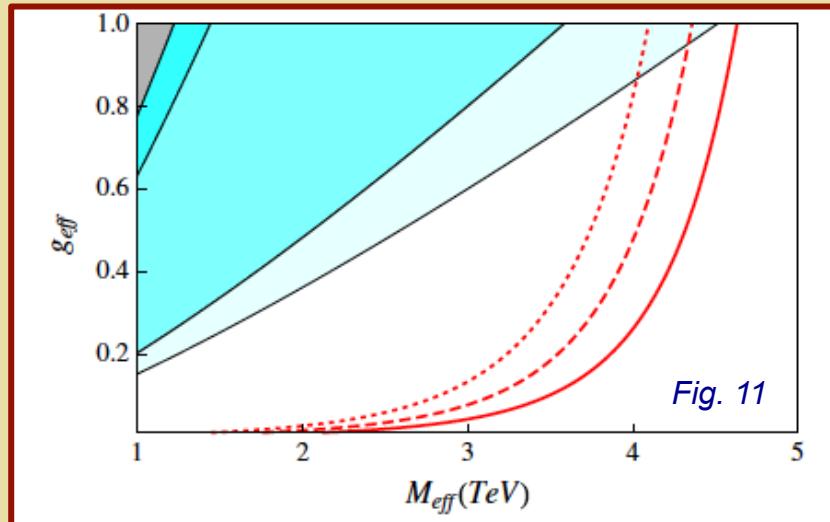
$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Helo et al claim:

$$\mathcal{L}_{\text{INT}} = g_1 \bar{Q}_i^\alpha d^\alpha S_i + g_2 \epsilon^{ij} \bar{L}_i F S_j^* + \text{H.c.}$$

$$g_{\text{eff}(S)} = (g_1 g_2)^{1/2}$$



$$M_{\text{eff}(S)} = (m_S^4 m_\psi)^{1/5}$$

$0\nu\beta\beta$ -Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

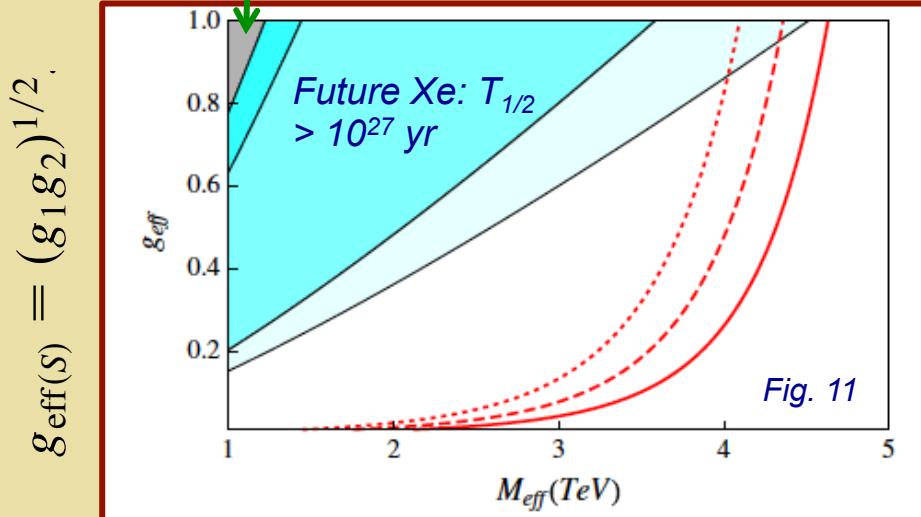
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Majorana

Helo et al claim:

EXO exclusion

$$\mathcal{L}_{\text{INT}} = g_1 \bar{Q}_i^\alpha d^\alpha S_i + g_2 \epsilon^{ij} \bar{L}_i F S_j^* + \text{H.c.}$$



$$M_{\text{eff}(S)} = (m_S^4 m_\psi)^{1/5}$$

$0\nu\beta\beta$ -Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

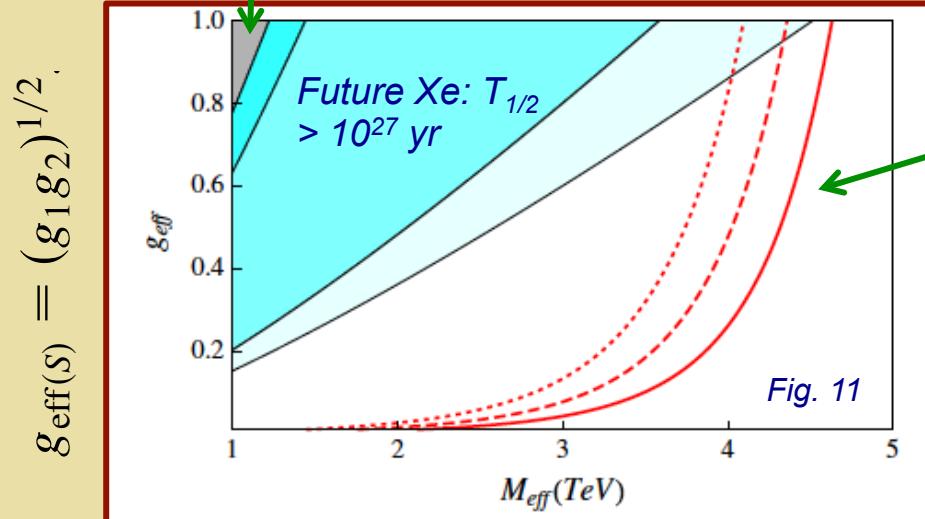
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LHC: $pp \rightarrow jj e^- e^-$

300 fb^{-1} :

— < 3 events

$$M_{\text{eff}(S)} = (m_S^4 m_\psi)^{1/5}$$

$0\nu\beta\beta$ -Decay: TeV Scale LNV

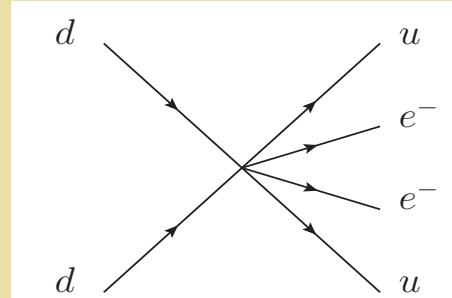
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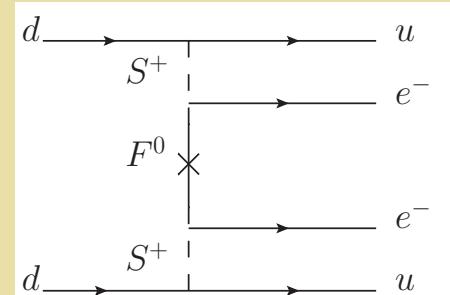
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Majorana

$0\nu\beta\beta$ - decay



LHC: $pp \rightarrow jj e^- e^-$



TeV Scale LNV

Comparing $0\nu\beta\beta$ & LHC sensitivities (our work):

- LHC backgrounds
- Running effective op's to low energy
- Matching onto hadronic d.o.f.
- Long range NME contributions

$0\nu\beta\beta$ -Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Backgrounds:

- *Charge flip*
- *Jet faking electron*

$0\nu\beta\beta$ -Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

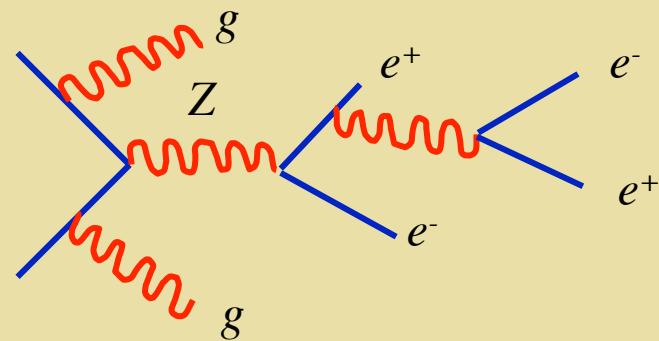
Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Backgrounds:

- Charge flip
- Jet faking electron



e^+ transfers most of p_T to conversion e^- ;
 $Z/\gamma^* + \text{jets} \rightarrow \text{apparent } e^- e^- jj \text{ event}$

$0\nu\beta\beta$ -Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

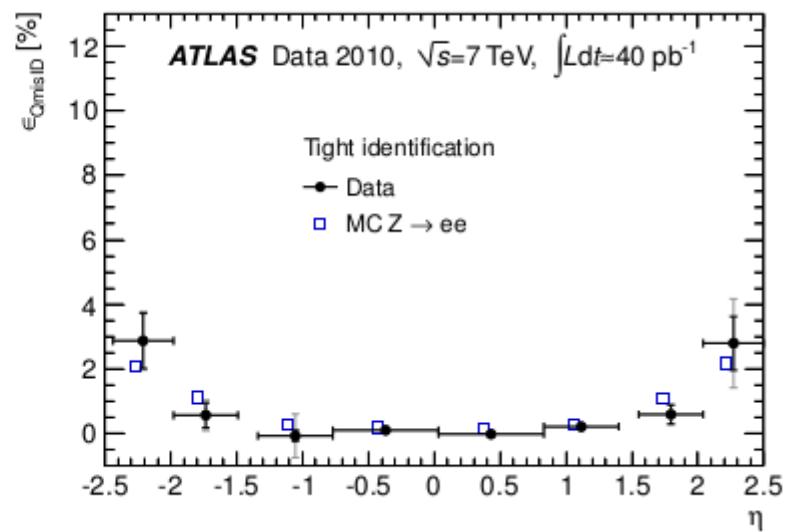
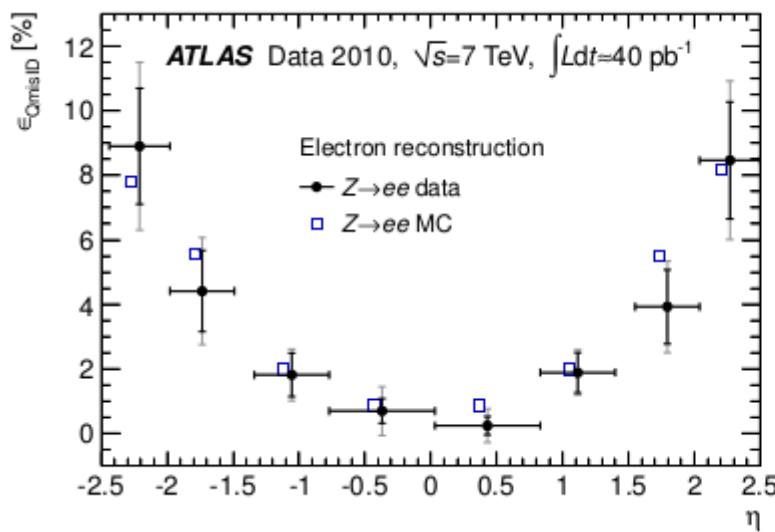
Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Backgrounds:

Bin in η and apply charge flip prob



$0\nu\beta\beta$ -Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Backgrounds: *Jet fakes (e.g., π^+ looks like e^+)*

$$\sigma_{JF} \text{ before cuts} = \sigma_{JF, MG+Pythia+PGS} \times (1/5000 \times 1/2)^{\# \text{ of jet-fakes}} \times \binom{\# \text{ of jets}}{\# \text{ of jet-fakes}}$$

0νββ-Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Backgrounds: Cuts

- H_T
- MET
- M_{\parallel}

$0\nu\beta\beta$ -Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Backgrounds: Cuts

$\sigma(\text{fb})$	Signal	Backgrounds									$\frac{S}{\sqrt{S+B}} (\sqrt{\text{fb}})$	
		Diboson			Charge Flip		Jet Fake					
		$W^- W^- + 2j$	$W^- Z + 2j$	$ZZ + 2j$	$Z/\gamma^* + 2j$	$t\bar{t}$	$t\bar{t}$	$\bar{t}+3j$	$W^- + 3j$	$4j$		
Before Cuts	0.142	0.541	6.682	0.628	903.16	68.2	6.7	0.45	15.09	362.352	0.0038	
Signal Selection	0.091	0.358	4.66	0.435	721.7	28.9	2.37	0.22	11.73	72.03	0.0031	
$H_T(\text{jets}) > 650 \text{ GeV}$	0.054	0.04	0.187	0.015	5.6	0.266	0.025	0.0003	0.102	0.027	0.0213	
$m_{\ell_1 \ell_2} > 130 \text{ GeV}$	0.039	0.029	0.105	0.008	0.163	0.127	0.024	3×10^{-4}	0.101	0.027	0.0493	
$E_T < 40 \text{ GeV}$	0.036	0.005	0.036	0.007	0.126	0.014	0.005	3×10^{-5}	0.03	0.017	0.0684	
$(\eta_{j_{1,2}} - \eta_{\ell_{1,2}})_{\text{max}} < 2.2$	0.033	0.003	0.022	0.005	0.093	0.009	0.004	2×10^{-5}	0.019	0.011	0.0738	

$0\nu\beta\beta$ -Decay: TeV Scale LNV

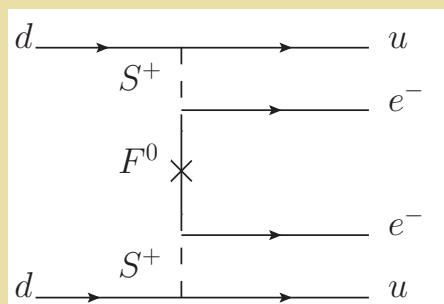
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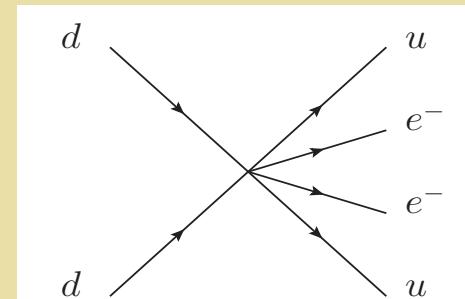
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Majorana

Low energy: *Matching*



Match onto \mathcal{O}_{eff} at Λ_{BSM}



$$\mathcal{L}_{\text{LNV}}^{\text{eff}} = \frac{C_1}{\Lambda^5} \mathcal{O}_1 + \text{h.c.}$$

$$\mathcal{O}_1 = \bar{Q} \tau^+ d \bar{Q} \tau^+ d \bar{L} L^C$$

$$g_{\text{eff}} = C_1(\Lambda)^{1/4}$$

$0\nu\beta\beta$ -Decay: TeV Scale LNV

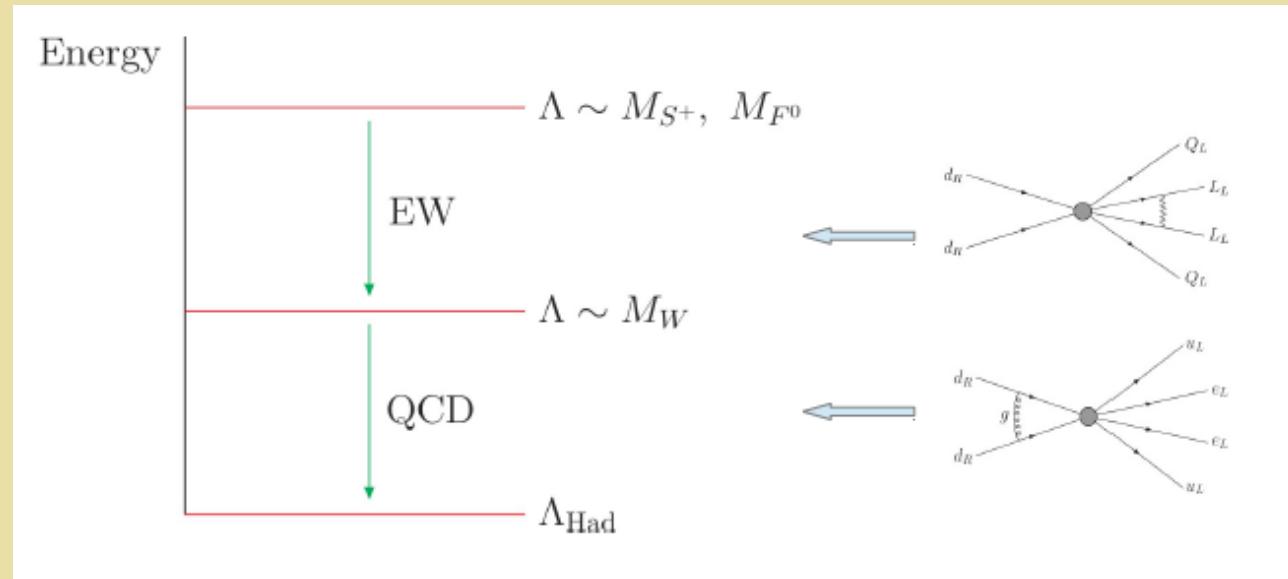
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Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Low energy: Running



$0\nu\beta\beta$ -Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Low energy: QCD Running

$$\begin{aligned}\mathcal{O}_1 &= (\bar{u}_L d_R)(\bar{u}_L d_R)(\bar{e}_L e_R^c), \\ \mathcal{O}_2 &= (\bar{u}_L \sigma^{\mu\nu} d_R)(\bar{u}_L \sigma_{\mu\nu} d_R)(\bar{e}_L e_R^c), \\ \mathcal{O}_3 &= (\bar{u}_L t^a d_R)(\bar{u}_L t^a d_R)(\bar{e}_L e_R^c), \\ \mathcal{O}_4 &= (\bar{u}_L t^a \sigma^{\mu\nu} d_R)(\bar{u}_L t^a \sigma_{\mu\nu} d_R)(\bar{e}_L e_R^c).\end{aligned}$$

$$\gamma^{ij} = -\frac{\alpha_s}{2\pi} \begin{pmatrix} 8 & 0 & 0 & 1 \\ 0 & -8/3 & 48 & 0 \\ 0 & 2/9 & -1 & 5/12 \\ 32/3 & 0 & 20 & 19/3 \end{pmatrix}$$

$$\mathcal{L}_{\text{eff}} = \sum_j \frac{C_j(\mu)}{\Lambda^5} \mathcal{O}_j(\mu) + \text{h.c.},$$

$$\mu \frac{d}{d\mu} C = \gamma^T C$$

$0\nu\beta\beta$ -Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

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Majorana

Low energy: QCD Running

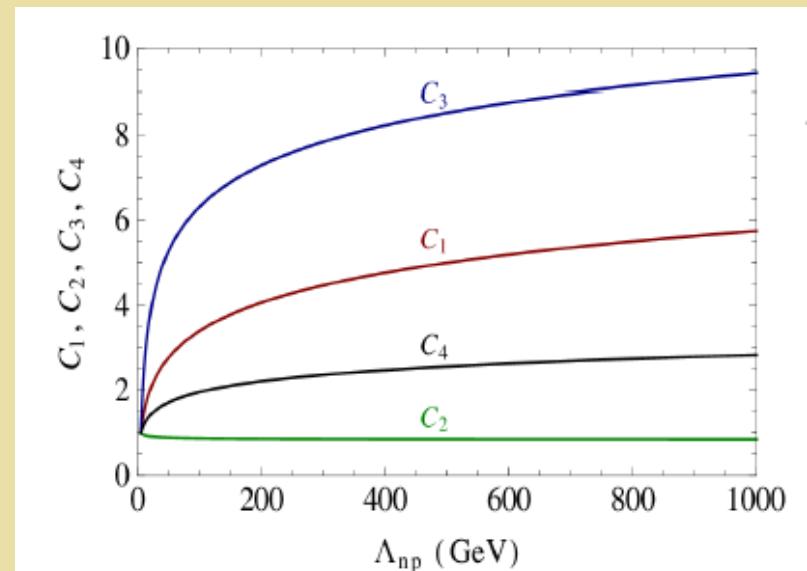
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$$\mathcal{O}_2 = (\bar{u}_L \sigma^{\mu\nu} d_R)(\bar{u}_L \sigma_{\mu\nu} d_R)(\bar{e}_L e_R^c),$$

$$\mathcal{O}_3 = (\bar{u}_L t^a d_R)(\bar{u}_L t^a d_R)(\bar{e}_L e_R^c),$$

$$\mathcal{O}_4 = (\bar{u}_L t^a \sigma^{\mu\nu} d_R)(\bar{u}_L t^a \sigma_{\mu\nu} d_R)(\bar{e}_L e_R^c).$$

Assuming $C_k = 1$ at $\mu = 5 \text{ GeV} \rightarrow$
 Effective DBD amplitude for \mathcal{O}_1
 substantially weaker for given
 LHC constraints



$0\nu\beta\beta$ -Decay: TeV Scale LNV

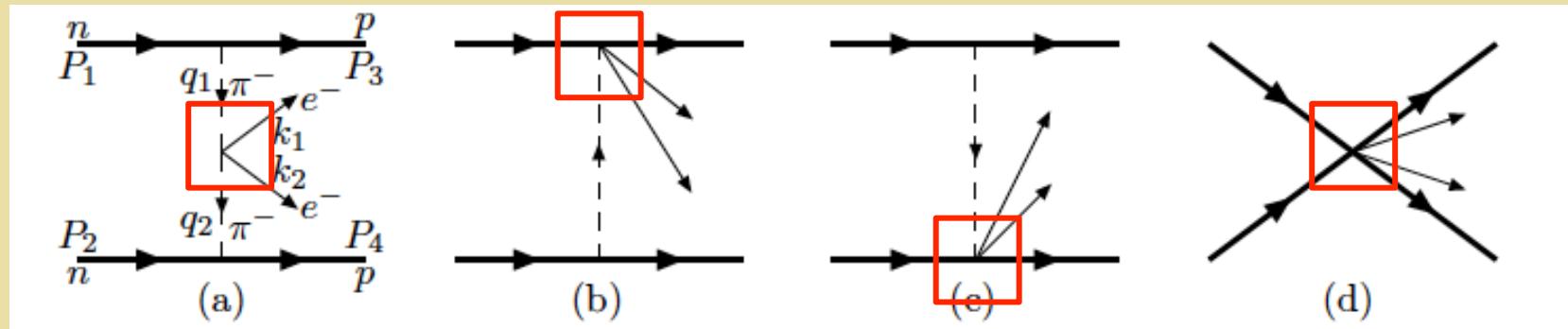
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Majorana

Low energy: Nuclear Matrix Elements: Long Range Effects



Exploit Chiral Symmetry & EFT ideas

$0\nu\beta\beta$ -Decay: TeV Scale LNV

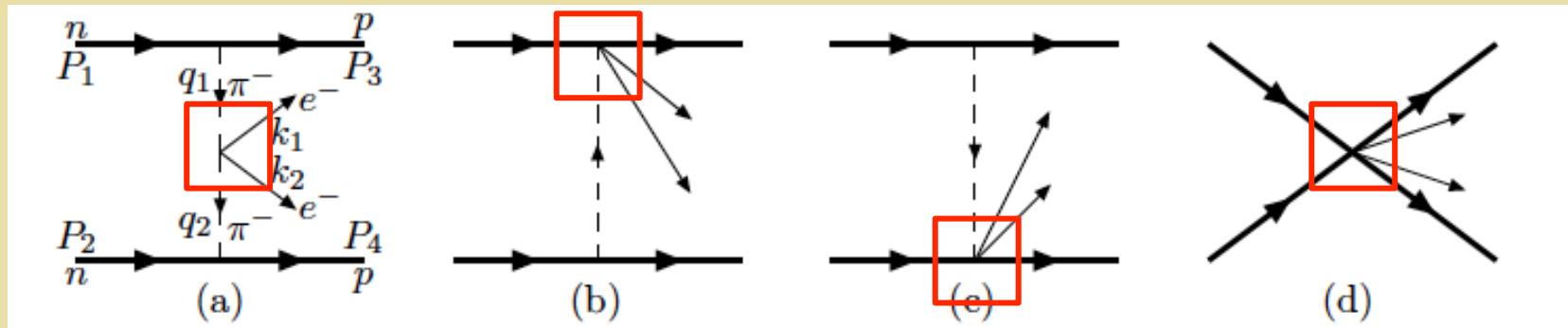
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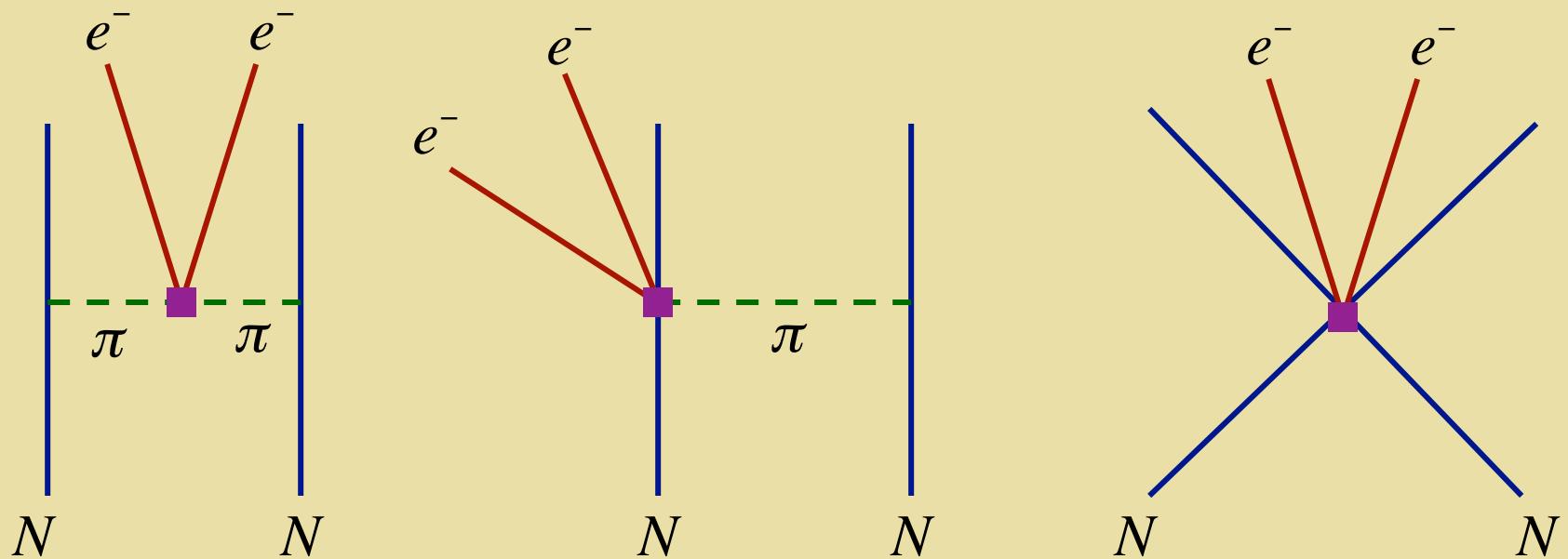


Our work

Exploit Chiral Symmetry & EFT ideas

Helo et al

$0\nu\beta\beta$ - decay in effective field theory



Tractable nuclear operators

Systematic operator classification

Prezeau, MJRM, Vogel
PRD 68 (2003) 034016

0ν ββ - decay in effective field theory

Operator classification

$$\mu = M_{WEAK}$$

$$\mathcal{L}(q,e) = \frac{G_F^2}{\Lambda_{\beta\beta}} \sum_{j=1}^{14} C_j(\mu) \hat{O}_j^{++} \bar{e} \Gamma_j e^c + h.c.$$

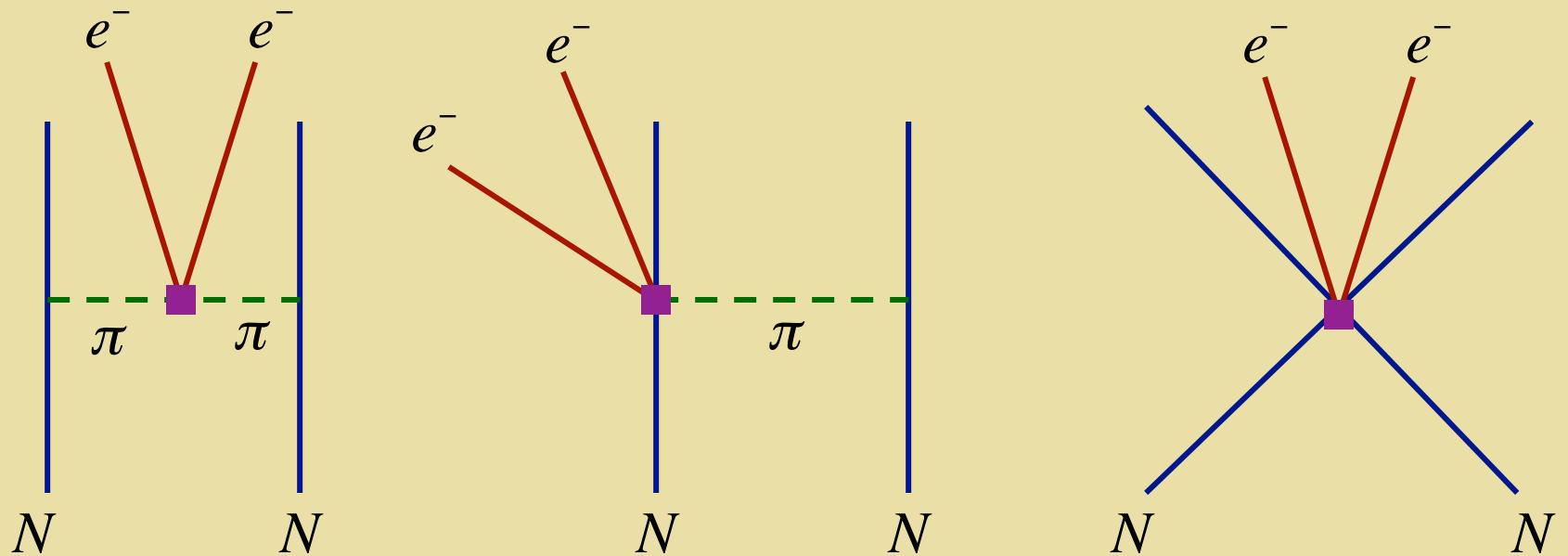
e.g.

$$\hat{O}_{1+}^{ab} = \bar{q}_L \gamma^\mu \tau^a q_L \bar{q}_R \gamma_\mu \tau^b q_R$$

0ν ββ - decay: **a = b = +**

Prezeau, MJRM, Vogel
PRD 68 (2003) 034016

$0\nu\beta\beta$ - decay in effective field theory



$K_{\pi\pi} \ p^{-2}$

$K_{\pi NN} \ p^{-1}$

$K_{NNNN} \ p^0$

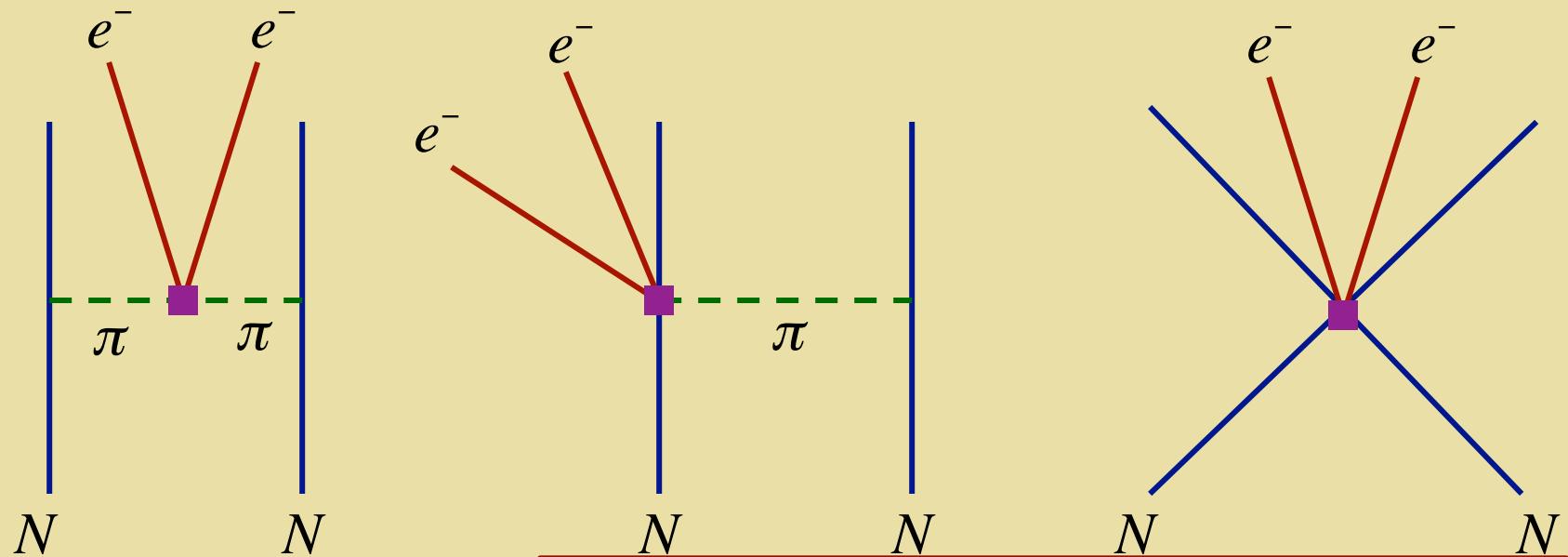
$\mathcal{O}(p^{-2})$ for

\hat{O}_{1+}^{++}

$\mathcal{O}(p^0)$ for

\hat{O}_{3+}^{++}

$0\nu\beta\beta$ - decay in effective field theory



$$K_{\pi\pi} p^{-2}$$

Hadronic matrix elements: A. Nicholson, M. Graesser talks

$O(p^{-2})$ for

$$\hat{O}_{1+}^{++}$$

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Putting Pieces Together

$0\nu\beta\beta$ -Decay: TeV Scale LNV

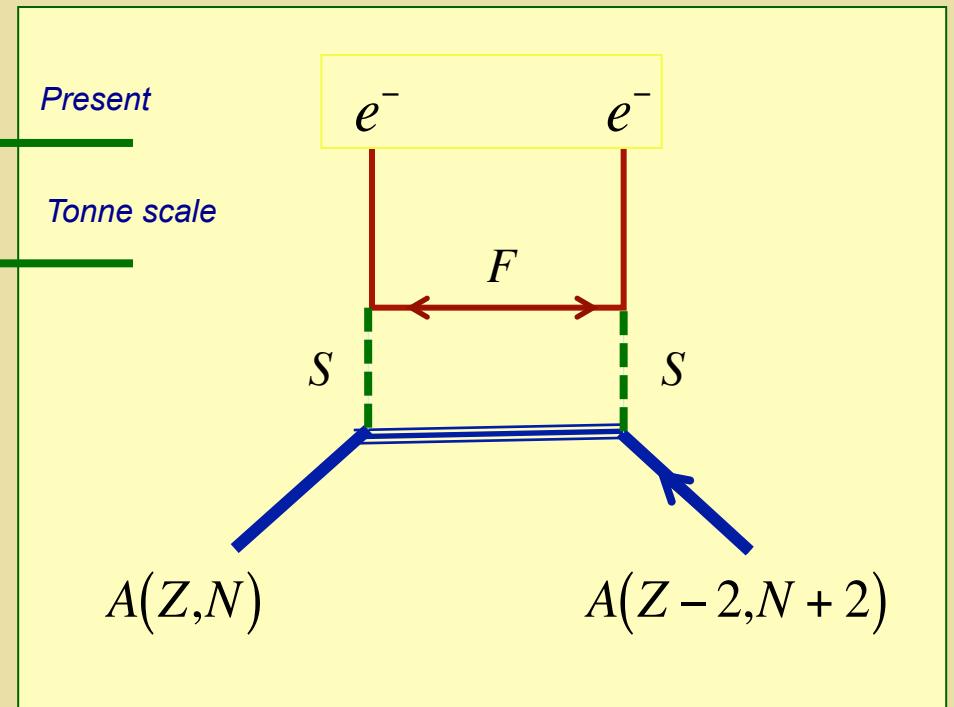
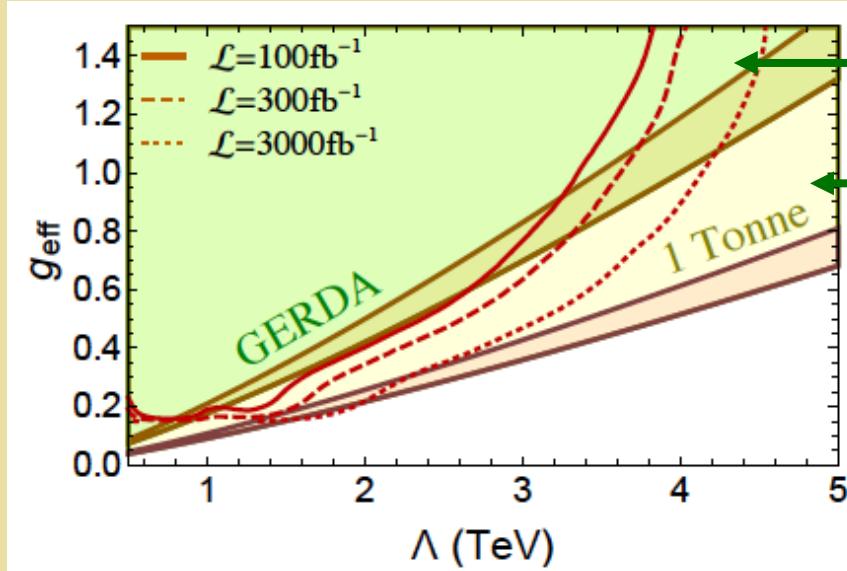
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Majorana

Benchmark Sensitivity: TeV LNV



$0\nu\beta\beta$ -Decay: TeV Scale LNV

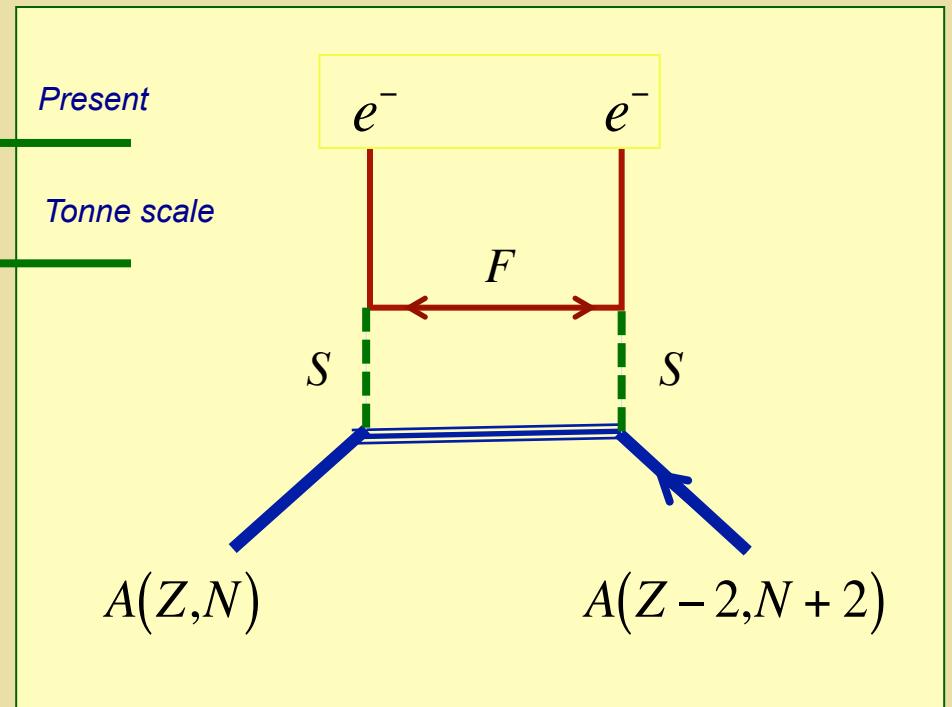
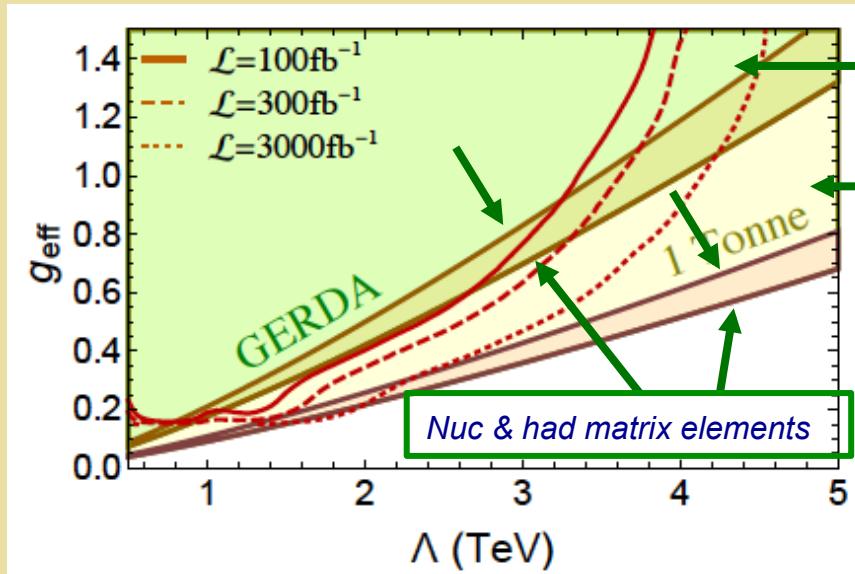
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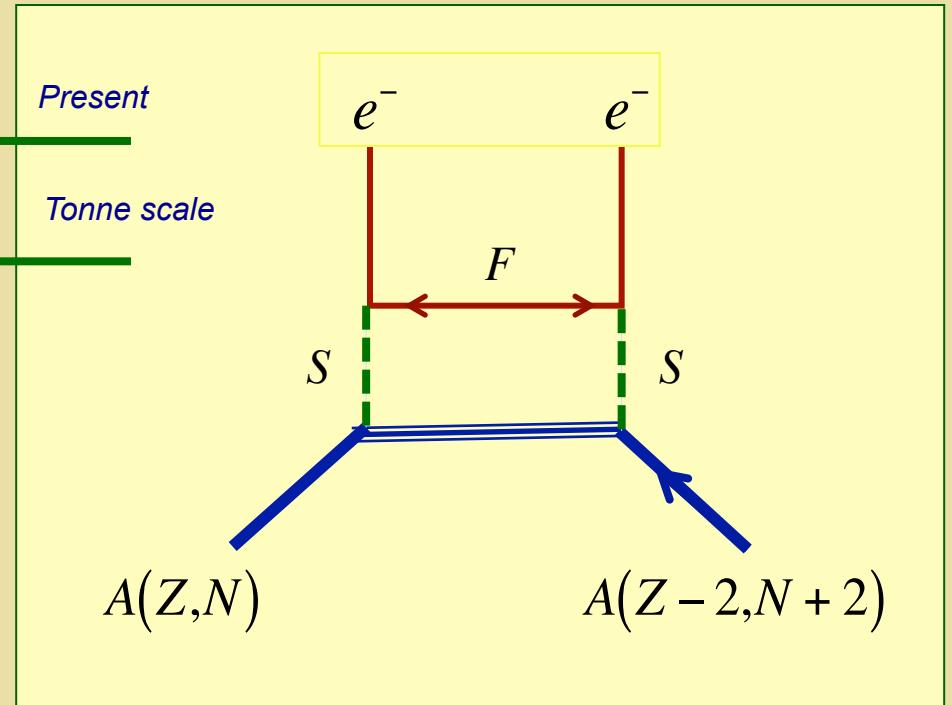
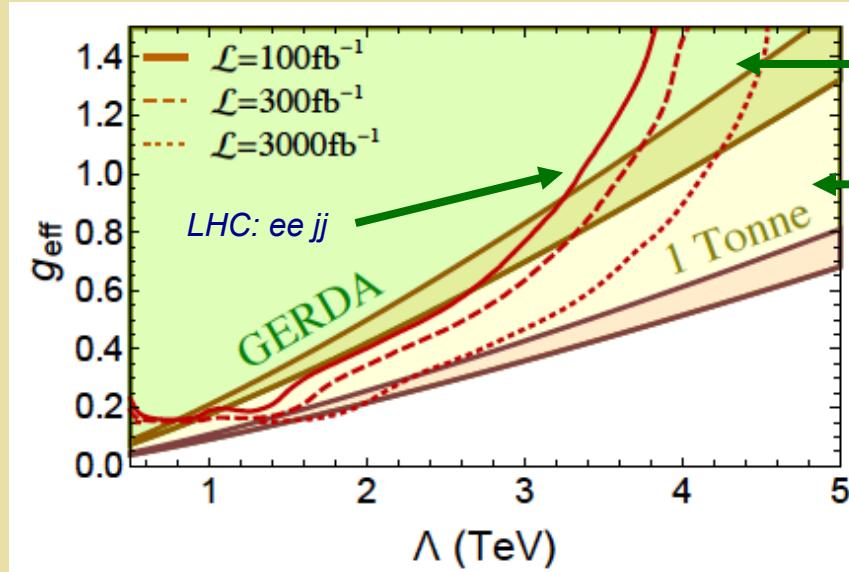
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$0\nu\beta\beta$ -Decay: TeV Scale LNV

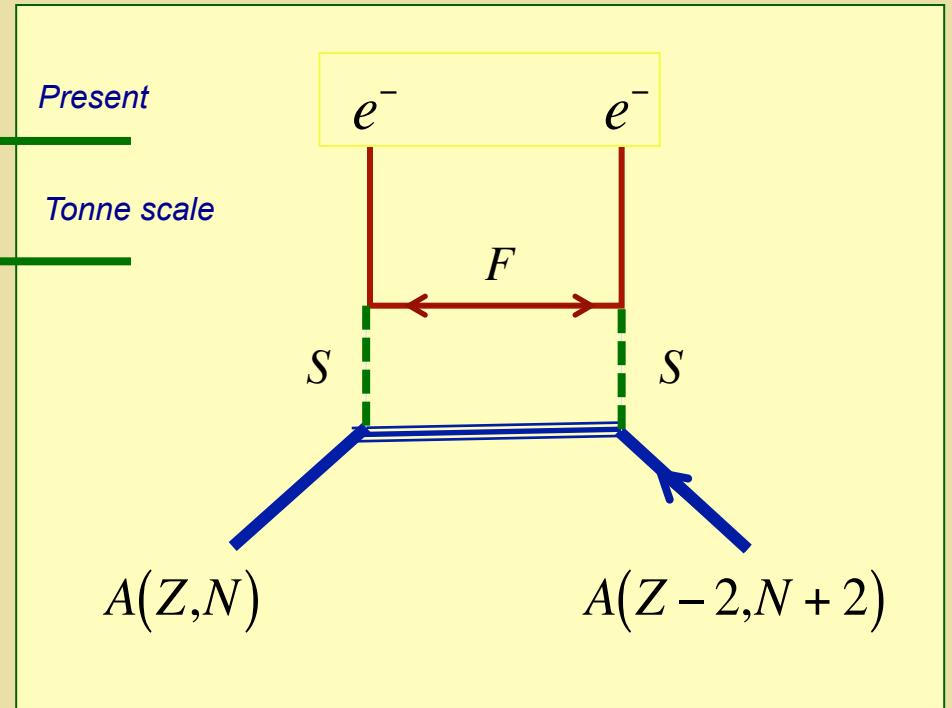
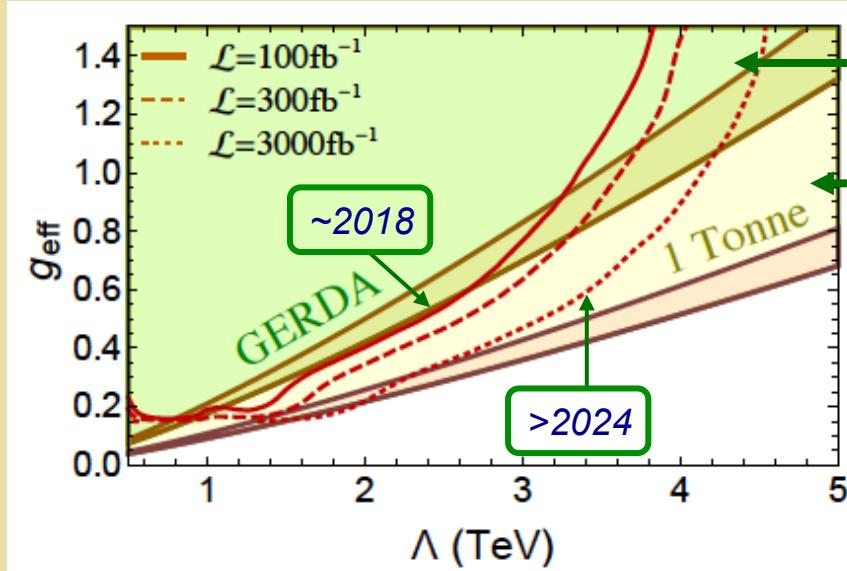
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Majorana

Benchmark Sensitivity: TeV LNV



$0\nu\beta\beta$ -Decay: TeV Scale LNV & m_ν

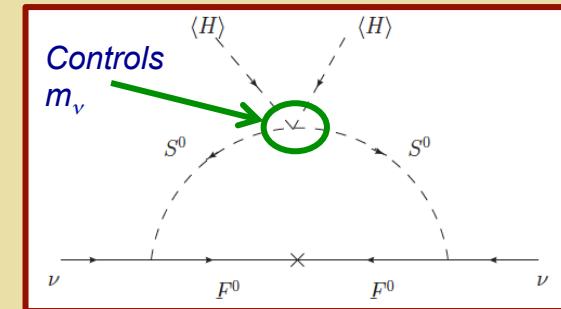
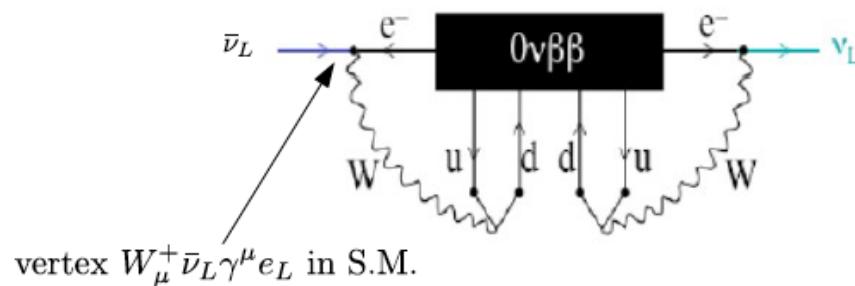
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Majorana

Implications for m_ν :



Schechter-Valle: non-vanishing Majorana mass at (multi) loop level

Simplified model: possible (larger) one loop Majorana mass

$0\nu\beta\beta$ -Decay: TeV Scale LNV & m_ν

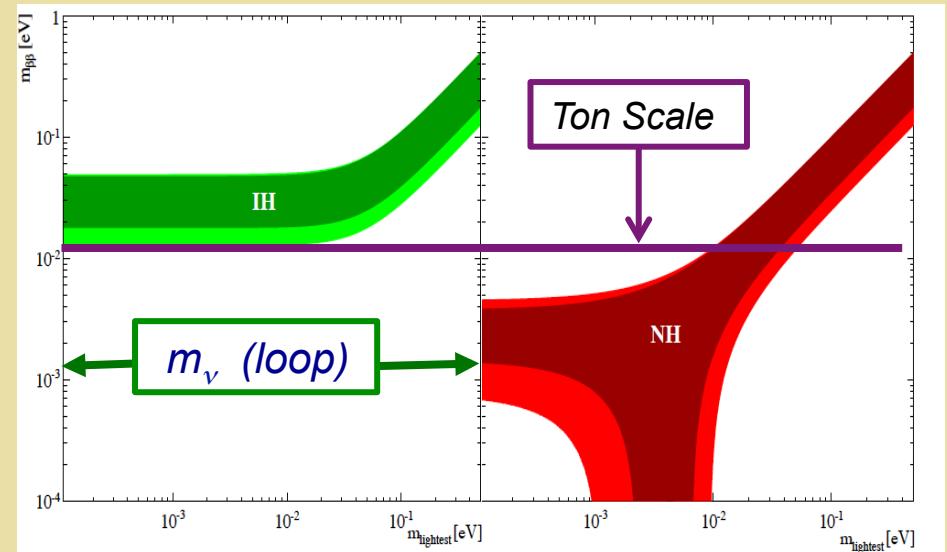
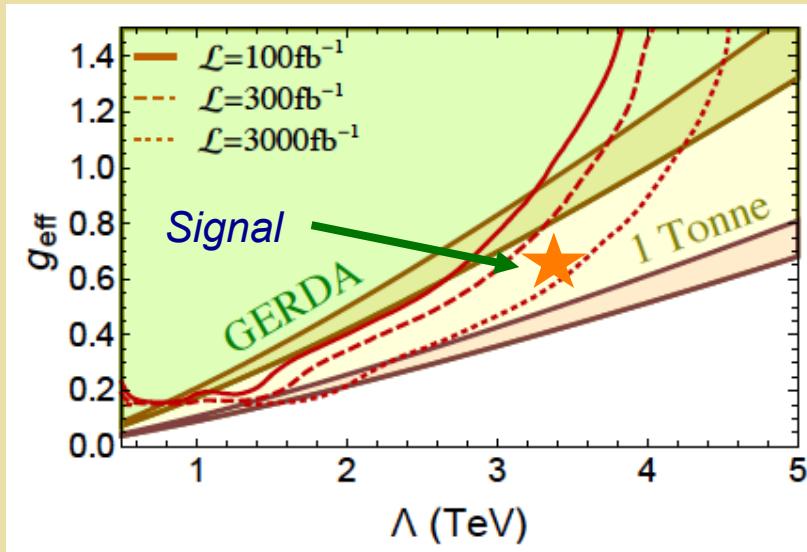
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Majorana

Implications for m_ν :



A hypothetical scenario

$0\nu\beta\beta / LHC$ Interplay: Matrix Elements

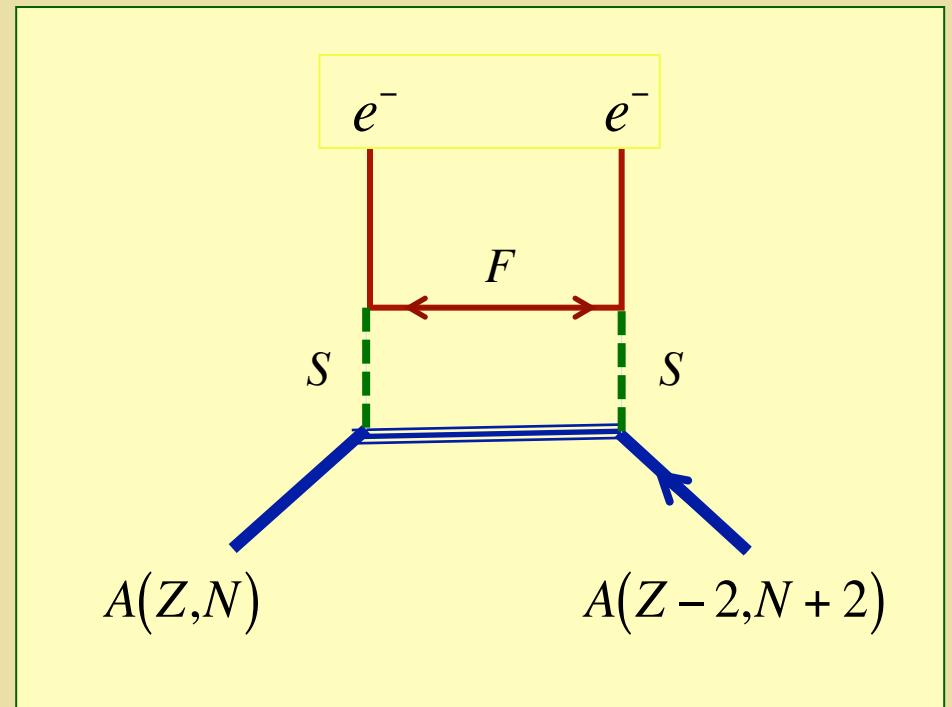
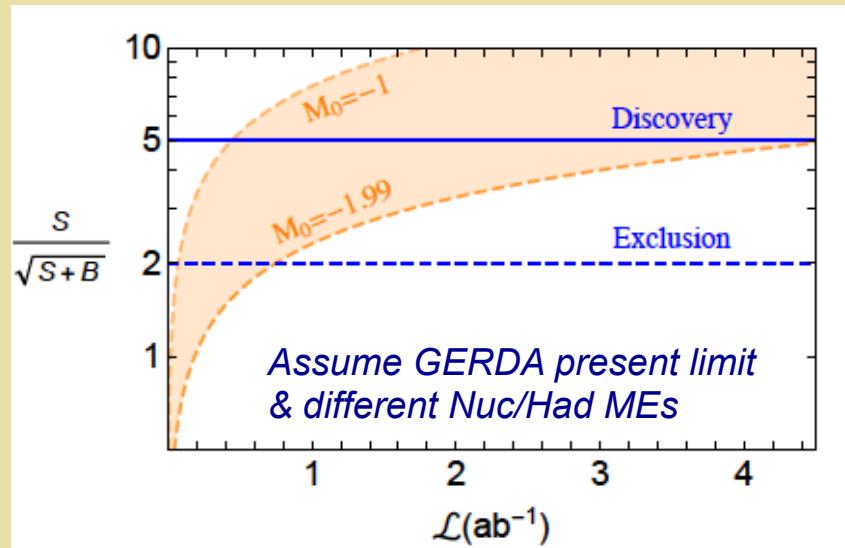
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Majorana

Benchmark Sensitivity: TeV LNV



$0\nu\beta\beta$ / LHC Interplay: Matrix Elements

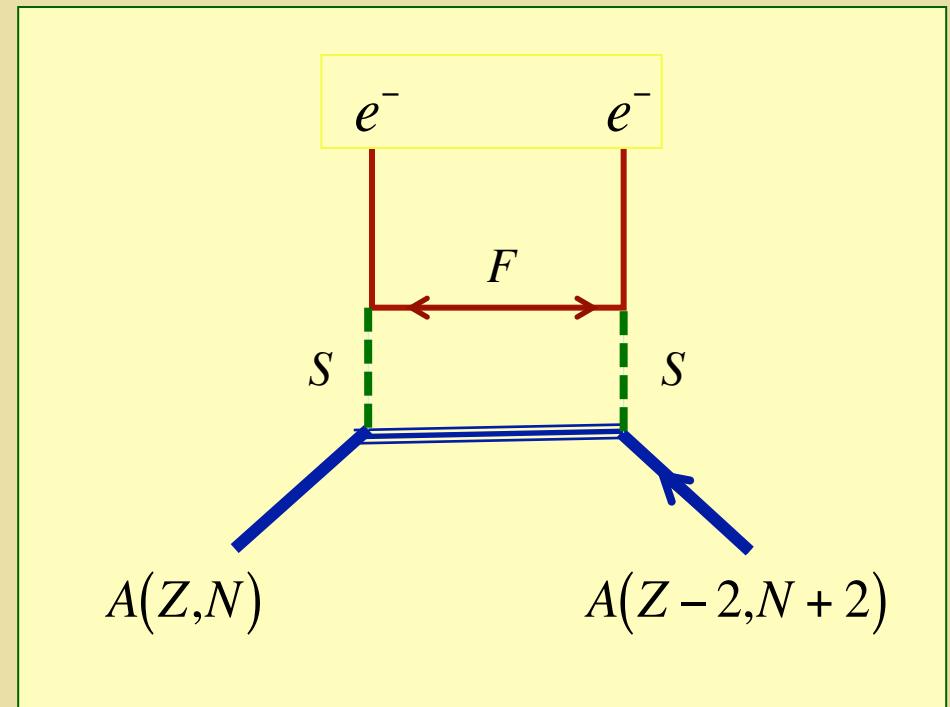
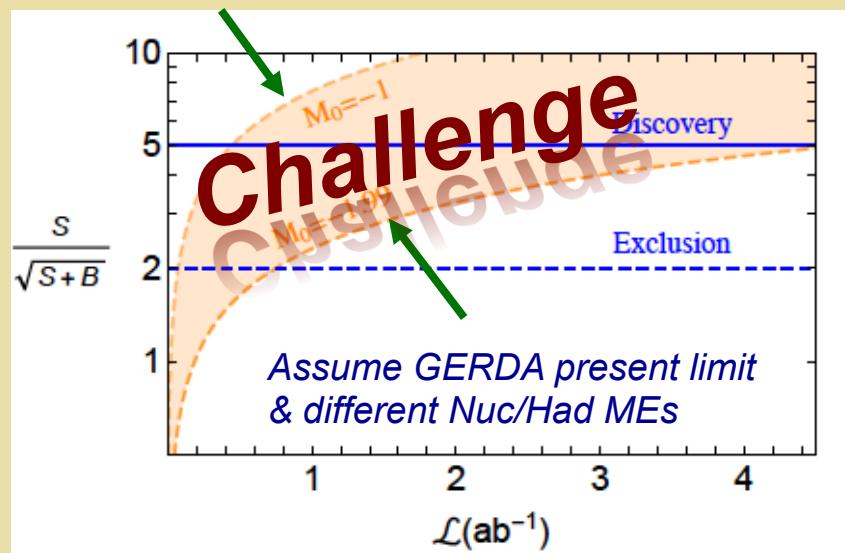
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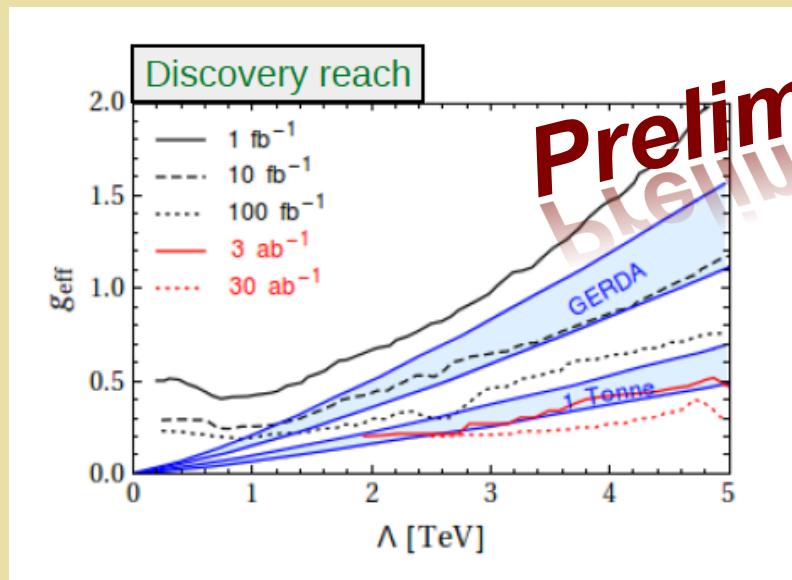
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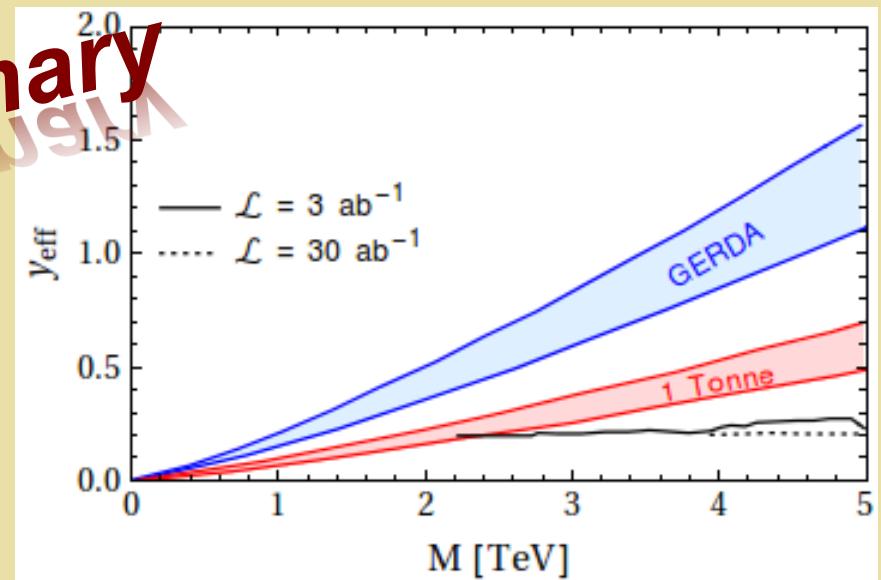
Benchmark Sensitivity: TeV LNV



LNV: pp at 100 TeV



Cut based analysis



Machine learning

M. Graesser, T. Peng,
MJRM, P. Winslow in prog...

V. Summary

- *LNV interactions responsible for m_ν may live at any scale from the conventional see-saw scale to the sub-GeV scale*
- *TeV scale LNV is theoretically well-motivated and would have important implications for baryogenesis if it exists*
- *$0\nu\beta\beta$ -decay and collider searches provide complementary probes of this scenario*
- *Fully exploiting this inter-frontier interface poses new challenges for hadronic & nuclear structure theory*