

Heavy quark transport coefficient from a Bayesian analysis

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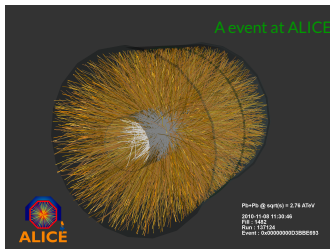
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Steffen A. Bass

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Heavy quarks in heavy-ion collision – a very complex system



outputs (y)

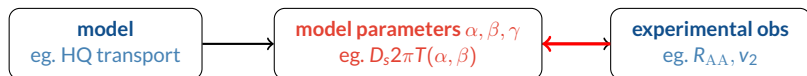
R_{AA}, v_2

Models

transport model:
Langevin, Boltzmann

calibration
parameters (x)

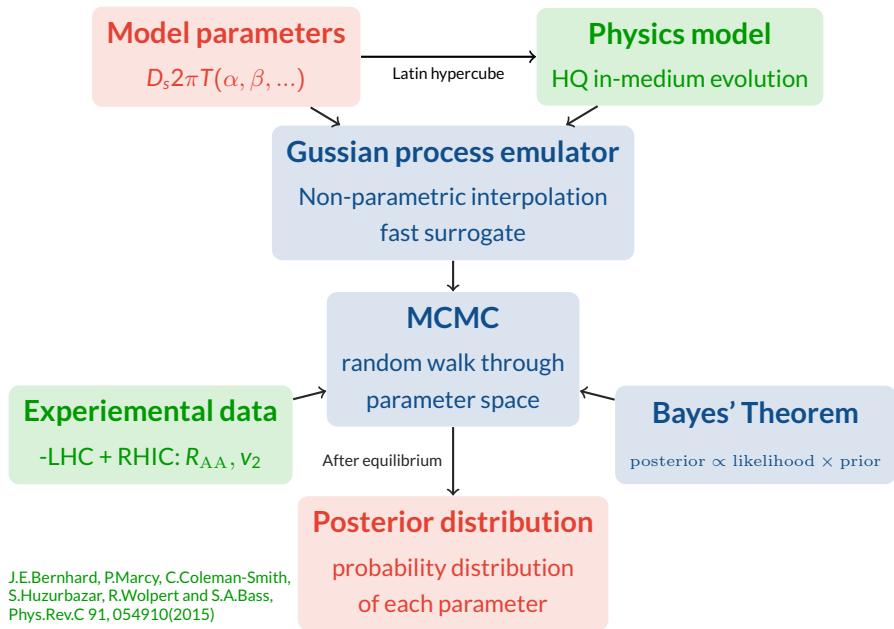
transport
coefficients



$$\vec{y} = \text{PM}(\vec{x}) + \delta_{\text{error}} \quad (1)$$

- We have a goal: to learn the unknowns (this case: heavy quark transport coefficients $\vec{x} = \hat{q}, D_s 2\pi T$) from the observables \vec{y}_{exp}
- We have some idea about \vec{x} (or not?), captured in our prior distribution $P(\vec{x})$
- We also have some idea how the outputs depends on the inputs \vec{x} , given in the form of likelihood $P(\vec{y}|\vec{x})$ (can be calculated by our model)
- Bayesian inference \Rightarrow posterior distribution:

$$P(\vec{x}|\vec{y}) = \frac{P(\vec{y}|\vec{x}) \cdot P(\vec{x})}{\int P(\vec{y}|\vec{x}) \cdot P(\vec{x}) d\vec{x}} \propto P(\vec{y}|\vec{x}) \cdot P(\vec{x}) \quad (2)$$



J.E.Bernhard, P.Marcy, C.Coleman-Smith,
S.Huzurbazar, R.Wolpert and S.A.Bass,
Phys.Rev.C 91, 054910(2015)

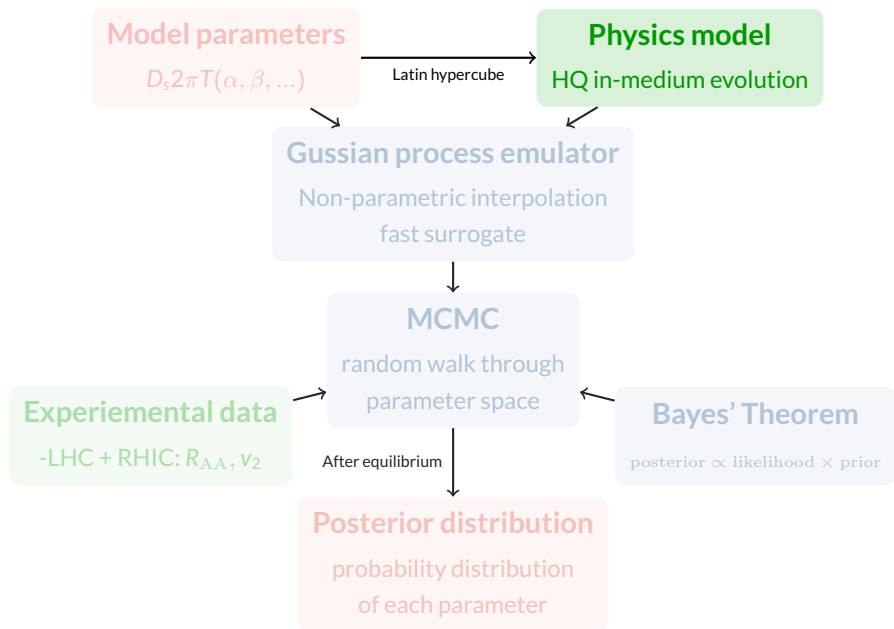
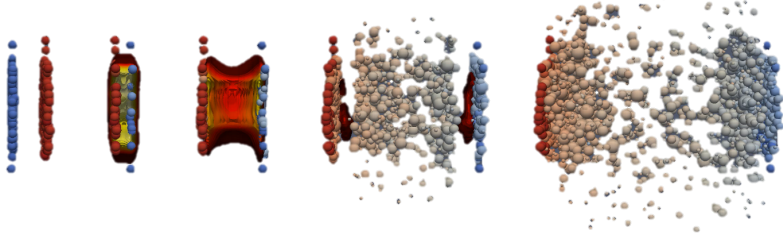


figure credit: Hannah Petersen (Au-Au collisions)



Initial condition:
Spatial IC: T_{RENTo}
Momentum IC: FONLL

In-medium evolution:
HQ transport: Langevin
(col + rad)
Medium: hydrodynamic

Hadronization:
fragmentation +
recombination

Position space: T_RENTo (A parametric IC model)

- Entropy deposition proportional to eikonal parameterization

$$\left. \frac{ds}{dy} \right|_{\tau=\tau_0} \propto \left(\frac{T_A^p + T_B^p}{2} \right)^{1/p}$$

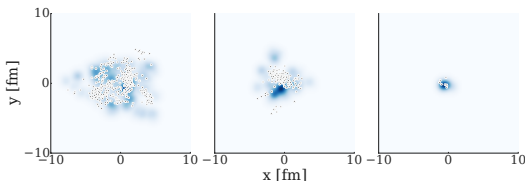
J.S.Moreland, J.Bernhard, and S.A.Bass,
Phys.Rev.C 92, 011901(2015)

- $p = 0 \Rightarrow ds/dy \propto \sqrt{T_A T_B}$ (mimic the behavior of IP-Glasma)

- Heavy quark initial production probability: $\left. \frac{dN}{dy} \right|_{\tau=\tau_0} \propto T_{AA}$

Momentum space: FONLL

- Parton distribution function: CTEQ6 M.Cacciari, S.Frixione, and P.Nason, arxiv:hep-ph/0102134
- Nuclear shadowing effect: EPS09 NLO



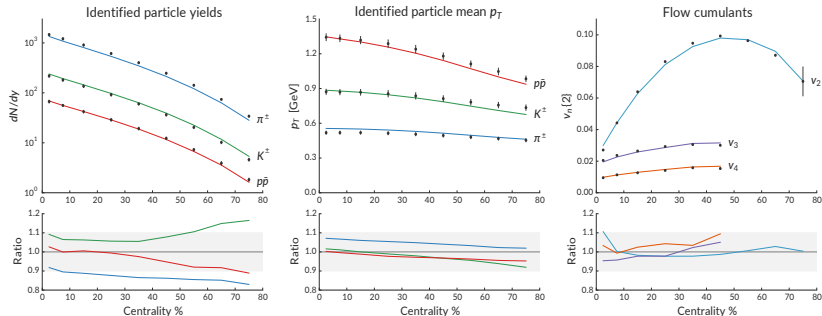
(2+1)D viscous hydro: iEbE-VISHNU

H.Song and U.W.Heinz,
Phys.Rev.C 77, 064901(2008)

- Equation of state from lattice QCD (HotQCD collaboration)
- Temperature-dependent shear + bulk vis correction

$$(\eta/s)(T) = (\eta/s)_{\min} + (\eta/s)_{\text{slope}}(T - T_c), T_c = 154\text{MeV}$$

$$(\zeta/s)(T) = (\zeta/s)_{\text{norm}} \times f(T)$$
- All the initial/medium related parameters (norm, p , η/s etc.) are calibrated by Bayesian analysis with experimental data



J.Bernhard, J.S.Moreland, S.A.Bass, J.Liu, and U.Heinz
Phys.Rev.C 94, 024907(2015)

S.Cao, G.Qin, and S.A.Bass,
Phys.Rev.C 92, 024907(2015)

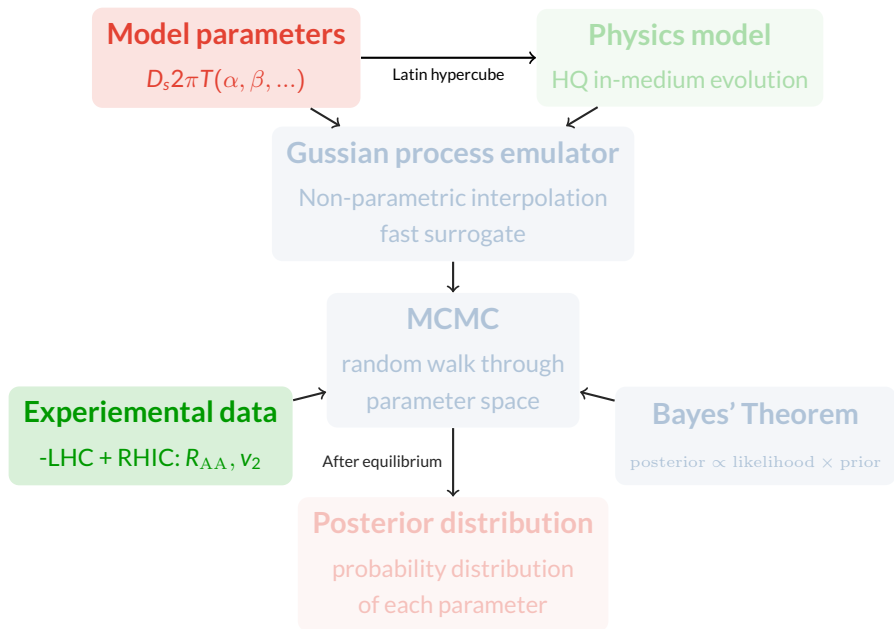
Improved Langevin transport model

$$\frac{d\vec{p}}{dt} = -\eta_D(p)\vec{p} + \vec{\xi} + \vec{f}_g \quad (3)$$

- Drag force: $\eta_D(p) = \kappa/(2TE)$
- Thermal random force: $\langle \xi^i(t)\xi^j(t') \rangle = \kappa\delta^{ij}\delta(t-t')$
- Recoil force from gluon radiation: $\vec{f}_g = -d\vec{p}_g/dt$
- Gluon emission probability:

$$\frac{dN_g}{dxdk_{\perp}^2 dt} = \frac{2\alpha_s P(x)\hat{q}_g}{\pi k_{\perp}^4} \sin^2\left(\frac{t-t_i}{2\tau_f}\right) \left(\frac{k_{\perp}^2}{k_{\perp}^2 + x^2 M^2}\right)^4 \quad (4)$$

- $\hat{q}_g = \hat{q}C_A/C_F = 2\kappa C_A/C_F, D_s = 2T^2/\kappa$
- **Diffusion coefficient $D_s=?$**



Model parameters and experimental data

HQ diffusion coefficient

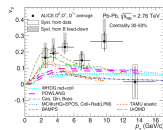
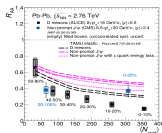
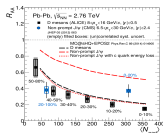
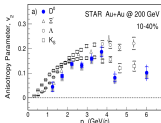
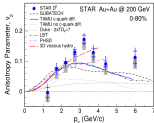
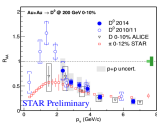
- Linear temperature dependence

$$D_S 2\pi T = A + B \cdot (T - T_C) \propto \eta/s$$

- + momentum dependence: $D_S 2\pi T = \frac{A(1+B \cdot T/T_C)}{1+C \cdot \log(E)}$

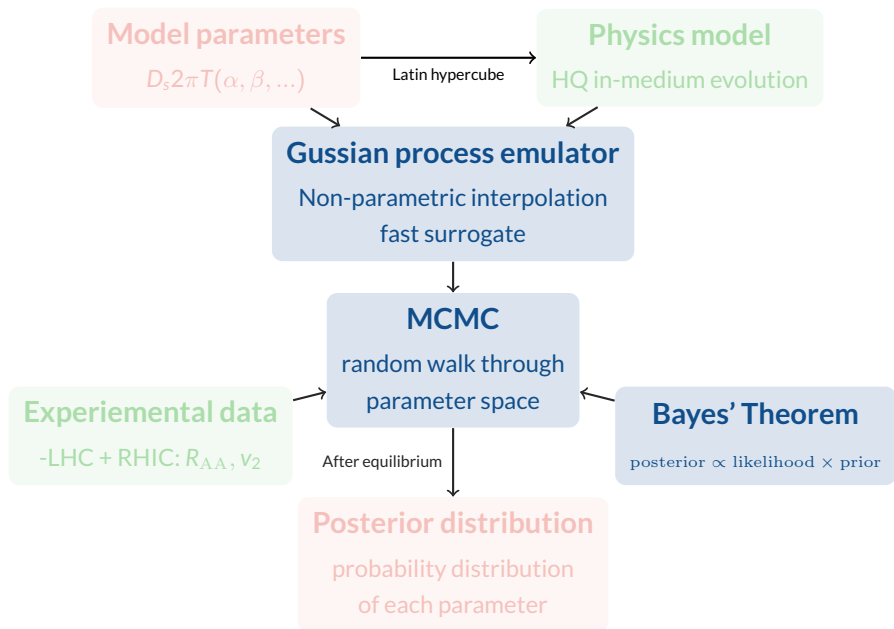
Experimental observables: $\vec{y}_{40} = (R_{AA}, v_2)$

- AuAu@ 200 GeV (STAR)
- PbPb@ 2.76 TeV (ALICE)



STAR collaboration: arxiv 1701.06060, 1601.00695

ALICE collaboration:
arxiv 1506.06604,
PRC 90, 034904 (2014)



Bayes' Theorem

$$P(\vec{x}_* | X, Y, \vec{y}_{\text{exp}}) \propto P(X, Y, \vec{y}_{\text{exp}} | \vec{x}_*) P(\vec{x}_*) \quad (5)$$

- $P(A, B) = P(A)P(B|A) = P(B)P(A|B)$, $P(B|A) \propto P(A|B)P(B)$
- **Posterior distribution $P(\vec{x}_* | X, Y, \vec{y}_{\text{exp}})$: probability of \vec{x}_* given observing $(X, Y, \vec{y}_{\text{exp}})$**
- **Likelihood $P(X, Y, \vec{y}_{\text{exp}} | \vec{x}_*)$: probability of observing $(X, Y, \vec{y}_{\text{exp}})$ given \vec{x}_***
 $\propto \exp[-\frac{1}{2}(\vec{y} - \vec{y}_{\text{exp}})^T \Sigma^{-1}(\vec{y} - \vec{y}_{\text{exp}})]$ D.Foreman-Mackey,D.Hogg,D.Lang,J.Goodman
arXiv:1202.3665
 - \vec{y} is the model output for input parameter \vec{x}
 - covariance matrix $\Sigma = \text{diag}(\sigma_{\text{stat}}^2) + \text{diag}(\sigma_{\text{sys}}^2)$
- Prior distribution $P(\vec{x}_*)$: (simplest case) uniformly distributed in parameter space

However...

1. Too many CPU hours

- $O(1000)$ CPU hours for one evaluation of \vec{y} at given \vec{x}
- $O(1000)$ walkers separate Metropolis-Hastings chains,
- 500 burn-in step, 1000 production step per walker



Gaussian process emulator

2. Too many observables

- $\vec{y} = R_{AA}, v_2$
- Independent GP emulators for each output?
- Highly correlated



Principal component analysis

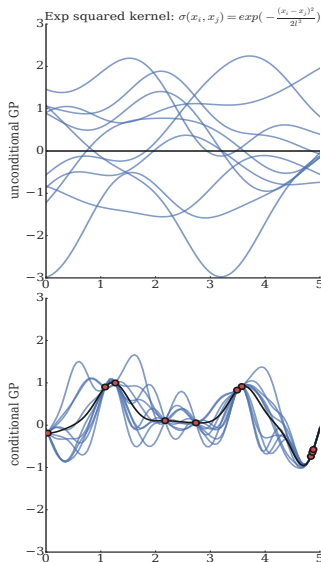
Gaussian process

- A collection of random variables, which have a joint Gaussian distribution
- Map inputs to normally-distributed outputs

$$\vec{y} \sim \mathcal{N}(\vec{\mu}, \sigma(X)) \quad (6)$$

- Only need to be specified by mean and covariance functions, for example:
$$\begin{pmatrix} y_1 \\ y_2 \end{pmatrix} \sim \mathcal{N} \left[\begin{pmatrix} \mu(\vec{x}_1) \\ \mu(\vec{x}_2) \end{pmatrix}, \begin{pmatrix} \sigma(\vec{x}_1, \vec{x}_1) & \sigma(\vec{x}_1, \vec{x}_2) \\ \sigma(\vec{x}_2, \vec{x}_1) & \sigma(\vec{x}_2, \vec{x}_2) \end{pmatrix} \right]$$

mean $\vec{\mu} = \vec{0}$, $\sigma(\vec{x}_1, \vec{x}_2) = \delta^2 \exp\left(-\frac{(\vec{x}_1 - \vec{x}_2)^2}{2\beta^2}\right)$
- Given (\vec{y}, X) GP parameters (hyper-parameters) δ^2, β can be estimated



GP as emulator

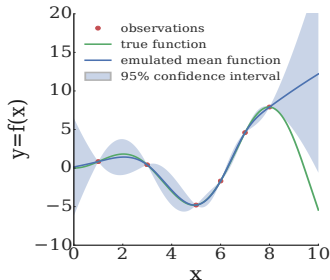
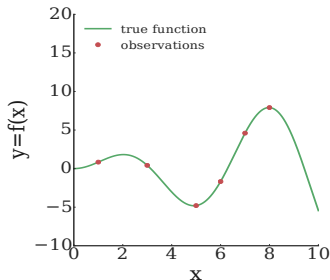
- Physics model: simulator $y = f(x) + \epsilon$

$$\begin{pmatrix} x_{11} & \dots & x_{1m} \\ \dots & & \dots \\ x_{n1} & \dots & x_{nm} \end{pmatrix} \Rightarrow \begin{pmatrix} y_1 \\ \dots \\ y_n \end{pmatrix} \quad (7)$$

- GP emulator: given a dataset (X, \vec{y}) , approximation of the simulator + probabilistic prediction
- This work: covariance(include a noise term)

$$\sigma(\vec{x}, \vec{x}') = \sigma_{\text{GP}}^2 \exp \left[- \sum_{k=1}^m \frac{(x_k - x'_k)^2}{2l_k^2} \right] + \sigma_n^2 \delta_{\vec{x}, \vec{x}'}$$

- Maximize the evidence $\log P(y_* | X, Y, \vec{x}_*) = -\frac{1}{2} Y^T \Sigma^{-1}(X, \vec{x}_*) Y - \frac{1}{2} \log |\Sigma(X, \vec{x}_*)| - N/2 \log(2\pi)$



Multiple observables, correlated?

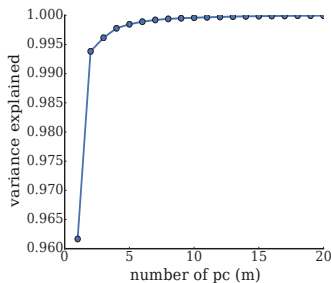
$$\begin{pmatrix} X_{11} & \dots & X_{1m} \\ \dots & & \dots \\ X_{n1} & \dots & X_{nm} \end{pmatrix} \Rightarrow \begin{pmatrix} Y_{11} & \dots & Y_{1k} \\ \dots & & \dots \\ Y_{n1} & \dots & Y_{nk} \end{pmatrix} \quad (8)$$

- Decompose into orthogonal linear principal components
- Singular value decomposition: n sets of k -dimension observables \Rightarrow n sets of l -dimensional PCs Z

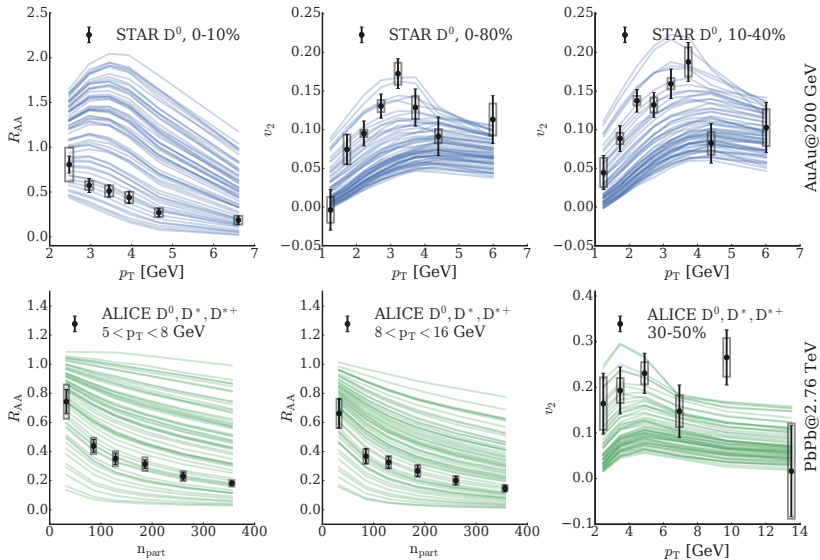
$$Y_{kn} = U_{kl} S_{ll} V_{ln}^T \quad (9)$$

$$Z = \sqrt{n} Y V \quad (10)$$

- Eigenvalue λ_i of Y represents the variance explained by PC \vec{z}_i



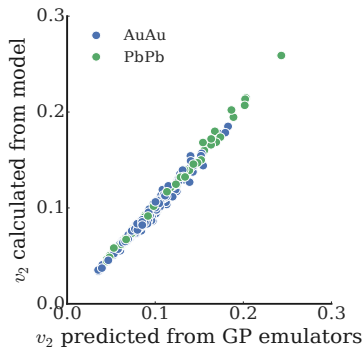
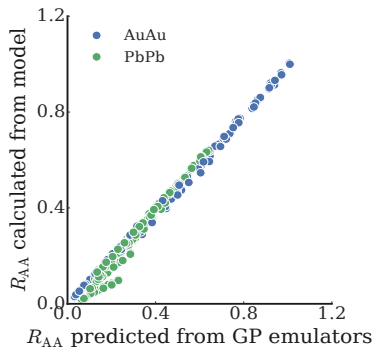
III: Prior (training data)

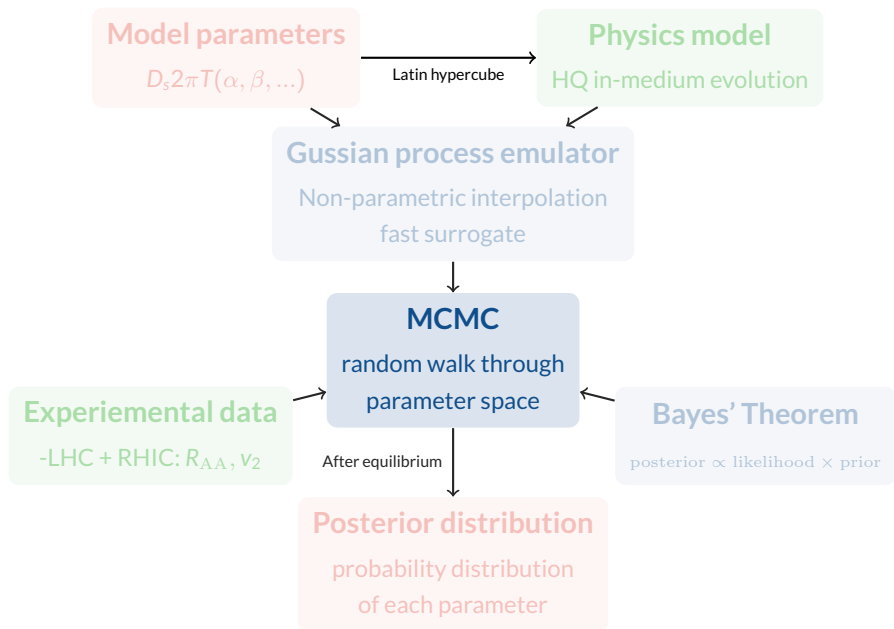


60 sets of inputs ($X = (\vec{x}_1, \dots, \vec{x}_{60})$) \Rightarrow 60 sets of outputs ($Y = (\vec{y}_1, \dots, \vec{y}_{60})$)

Another 10 sets of validation inputs \vec{x}

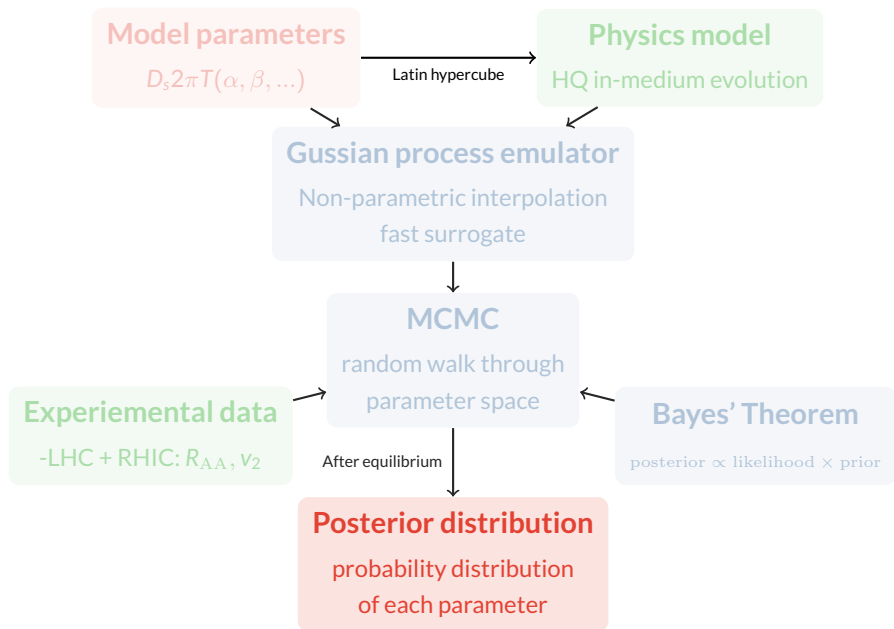
Compare between physics model calculation \vec{y} and GP emulator predicted \vec{y}_{pred}





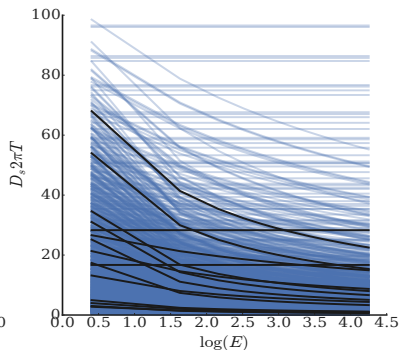
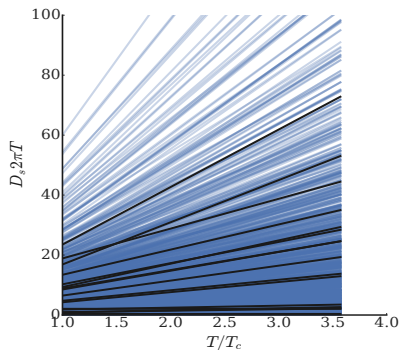
Affine invariance MCMC ensemble sampler: emcee

- random walk through parameter space, weighted by likelihood
- equilibrium \Rightarrow posterior distribution
- An ensemble \vec{X} (consists of many independent walkers)
- Each walker takes random walk by Metropolis-Hasting algorithm, where each step is accepted or rejected based on the likelihood $P(X, Y, \vec{y}_{\text{exp}} | \vec{X})$
- **Acceptance rate:** α
 - $\alpha = 1$ purely random walk
 - $\alpha = 0$ walker stuck
 - Normally optimal proposal has $\alpha = 44\%$ for 1d, 25% larger than 4d
- This work: $\sim 30 - 40\%$



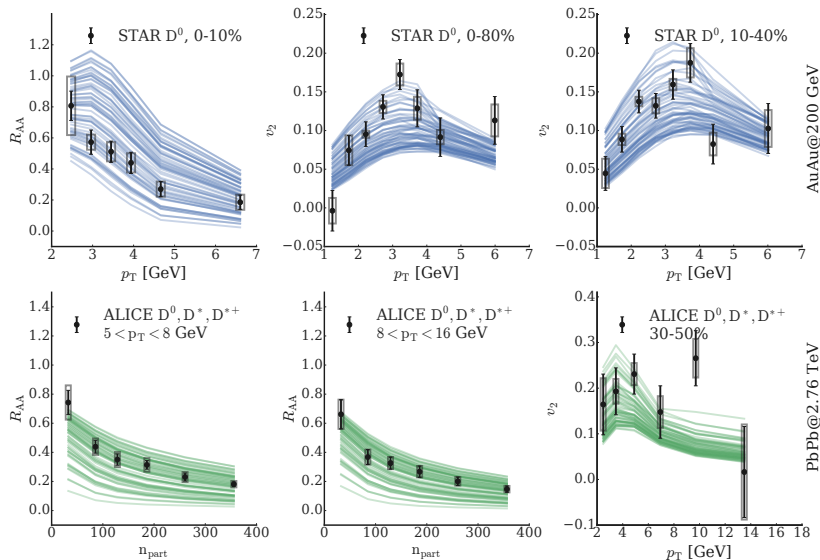
Parametrization

- $D_s 2\pi T = A + B \cdot (T - T_c)$
- $D_s 2\pi T = \frac{A(1+B \cdot T/T_c)}{1+C \cdot \log(E)}$

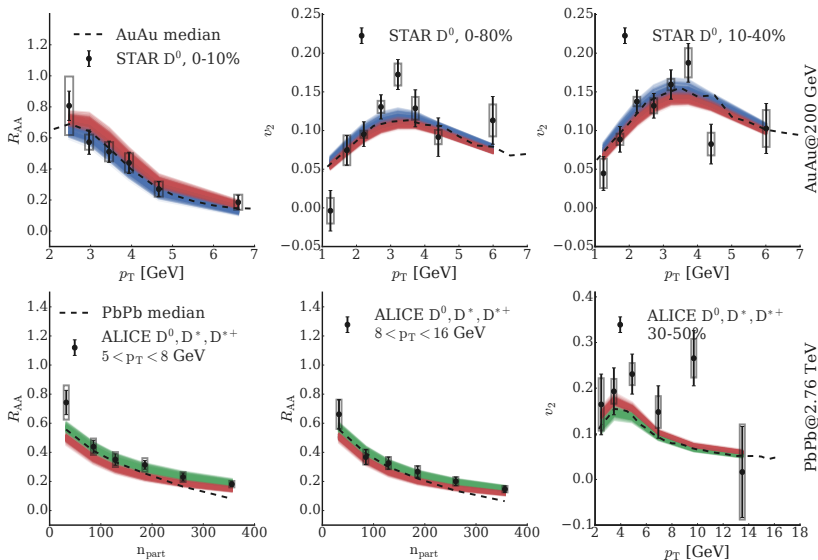


$$\text{IV: } D_S 2\pi T = A + B \cdot (T - T_C)$$

Prior: (60 sets of training data)

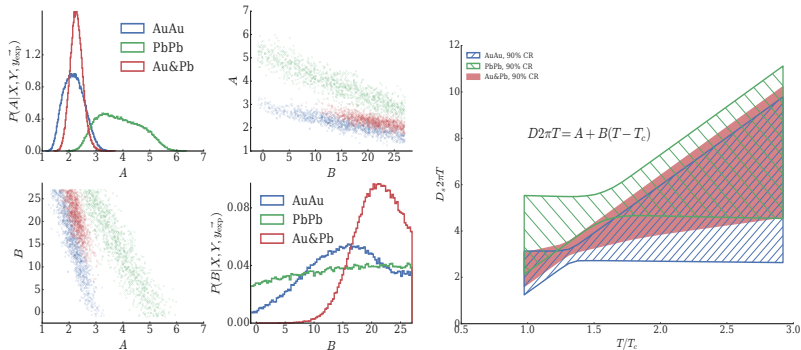


Posterior results: (200 sets of random posterior outputs)



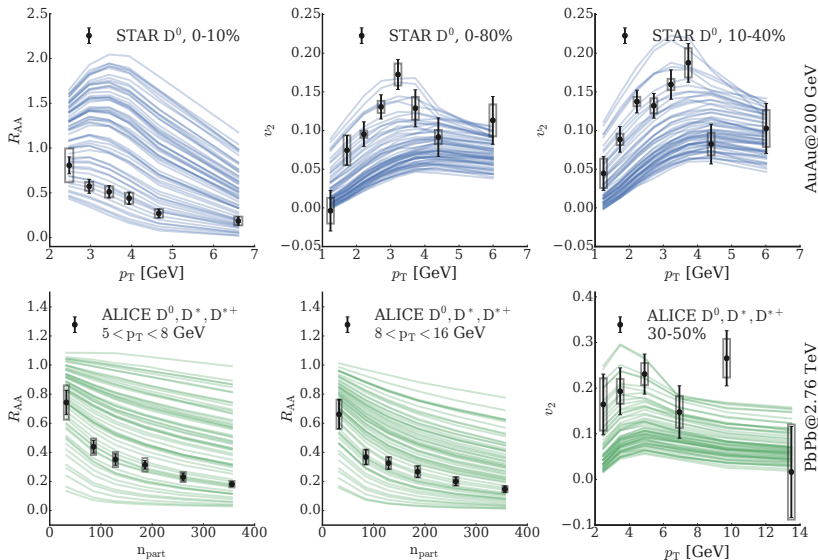
$$\text{III: } D_S 2\pi T = A + B \cdot (T - T_c)$$

Posterior distributions of (A', B')



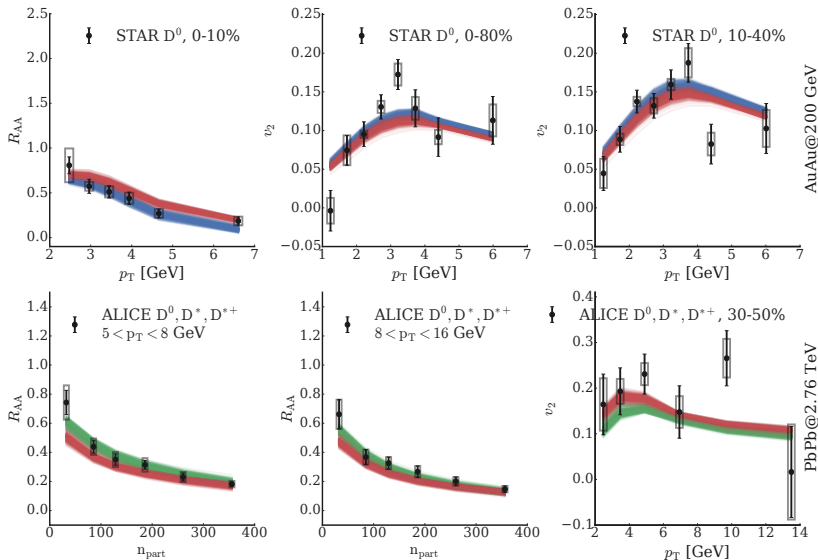
$$IV: D_S 2\pi T = \frac{A(1+B \cdot T/T_c)}{1+C \cdot \log(E)}$$

Prior: (60 sets of training data)



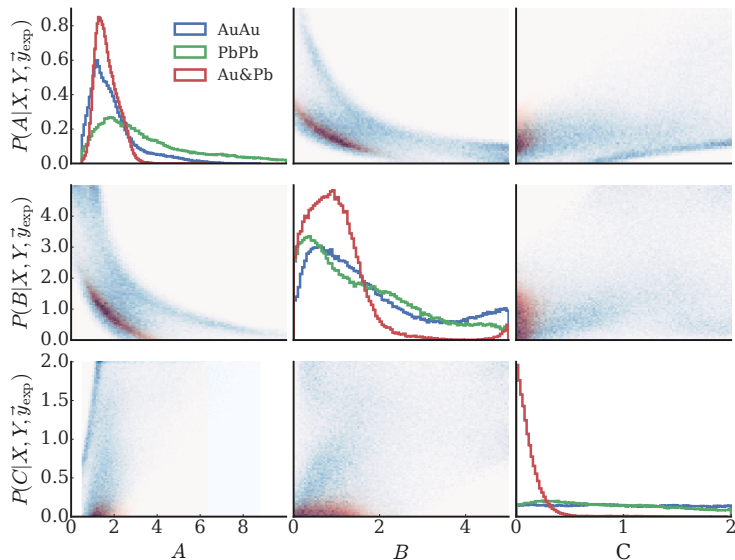
$$IV: D_S 2\pi T = \frac{A(1+B \cdot T/T_c)}{1+C \cdot \log(E)}$$

Posterior results: (200 sets of random posterior outputs)



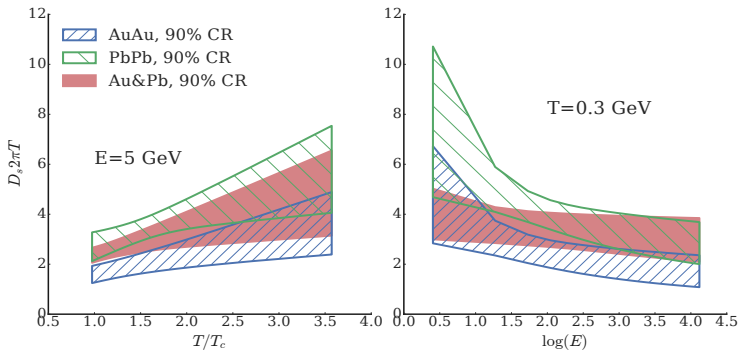
$$\text{IV: } D_S 2\pi T = \frac{A(1+B \cdot T/T_c)}{1+C \cdot \log(E)}$$

Posterior probability distribution



$$\text{IV: } D_s 2\pi T = \frac{A(1+B \cdot T/T_c)}{1+C \cdot \log(E)}$$

Posterior probability distribution



- Bayesian analysis provides a rigorous method to estimate the optimal parameters (with uncertainties) to simultaneously describing the experimental data
- Heavy quark diffusion coefficients can be extracted from the analysis and be constrained with some precision, ($D_s 2\pi T \sim 2 - 6$ near T_c)
- A higher precision of experimental data?
- A more sophisticated momentum dependent inspired by pQCD calculation $D_s = D_{\text{spQCD}} \cdot K(p)$; test on different transport models