

# Challenges in Quarkonium Production in p+p and p+Pb Collisions at the LHC

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U.S. DEPARTMENT OF  
**ENERGY**

Office of  
Science

# Outline:

- Quarkonium in  $p+p$
- Quarkonium in  $p+Pb$



# Quarkonium in p+p

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NRQCD factorization theorem for *e.g.*  $J/\psi$ :

$$\sigma_{J/\psi} = \sum_n \sigma_{c\bar{c}[n]} \langle \mathcal{O}^{J/\psi}[n] \rangle$$

# Quarkonium Production Schemes

- Color singlet model (CSM)
  - Assume physical color singlet state, normalization is quarkonium wavefunction at origin
  - Disagreement with  $p_T$  dependence apparent, including higher order terms does not make significant improvement (see Lansberg's talk)
- Nonrelativistic QCD (NRQCD)
  - Rigorous effective field theory based on factorization of soft and hard scales
  - Expansion of cross section in velocity and strong coupling
  - Not clear that NRQCD factorization agrees with data
- Color evaporation model (CEM)
  - Does not separate states into color or spin on average
  - Fewer parameters than NRQCD (one per state)
  - New results becoming available
- $k_T$  factorization
  - Off shell matrix elements and unintegrated gluon distributions



# Nonrelativistic QCD Approach

NRQCD factorization theorem for e.g.  $J/\psi$ :

$$\sigma_{J/\psi} = \sum_n \sigma_{ce[n]} \langle O^{J/\psi} [n] \rangle$$

$n$  sums over all Fock states, singlet and octet;  $\sigma_{ce[n]}$  is pair production rate for specific color and spin state, calculated in pQCD;  $\langle O^{J/\psi} [n] \rangle$  is long distance matrix element (LDME) describing conversion to final state  $J/\psi$  assuming that hadronization does not change spin or momentum

LDMEs are assumed to be universal

Cross section is a double expansion in relative velocity of the pair,  $v$ , and strong coupling constant  $\alpha_s$ , LDMEs scale with powers of  $v$

Color singlet,  $n = {}^3S_1^{[1]}$ , is leading term in  $v$ , color octet states ( ${}^1S_0^{[8]}$ ,  ${}^3S_1^{[8]}$ ,  ${}^3P_J^{[8]}$ ) are subleading, octet LDMEs determined by fitting LDMEs to data, these are then used to predict observables such as polarization, LDMEs of other states through heavy quark spin symmetry

NRQCD predicts strong transverse polarization at high  $p_T$



# Are the LDMEs Universal?

- The fit results depend on the energy scales of the process described, e.g. whether analysis is global or not and whether or not  $e^+e^-$  and  $ep$  data also included
- The fit results depend on the  $p_T$  scale, whether the minimum  $p_T$  is 3, 5, or 7 GeV
- The fit results depend on whether or not polarization is fitted or predicted
- Fits to  $p_T$  distributions do not describe the total cross sections
- Using LDMEs fitted to  $J/\psi$  results with heavy quark spin symmetry does not translate well to other states, e.g.  $\eta_c$



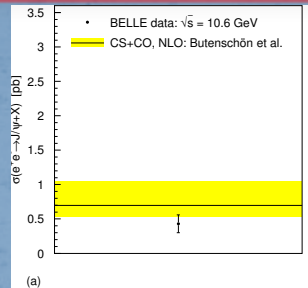
# LDMEs depend on process and scale

State of the art as of 1404.3723

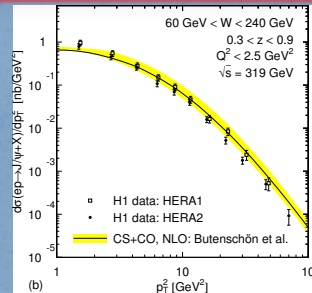
Included in fits

Butenschon  
& Kniehl  
 $p_T > 3 \text{ GeV}$

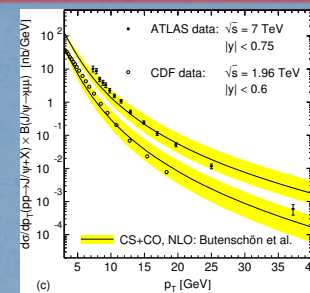
$e^+e^-$



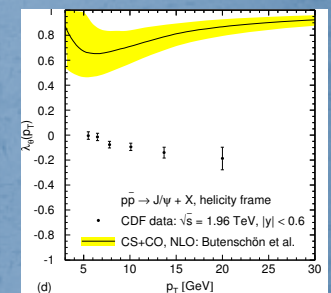
$ep$



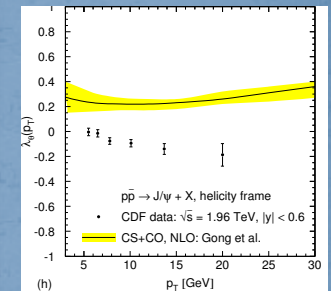
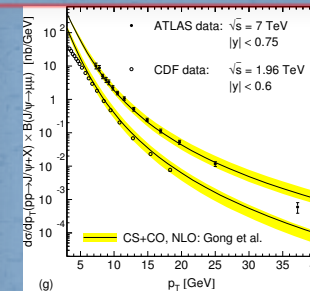
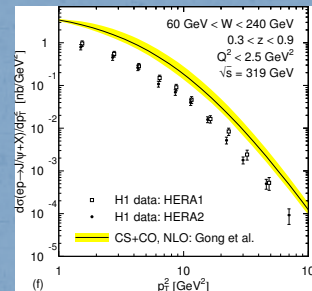
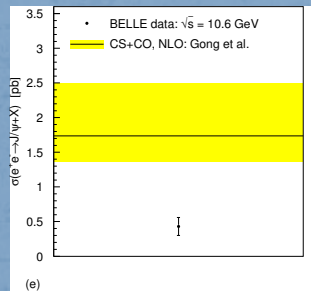
$pp$  distribution



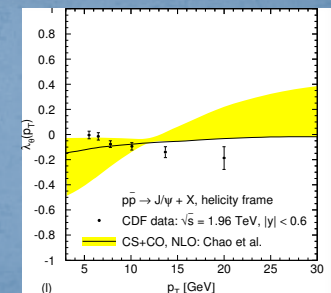
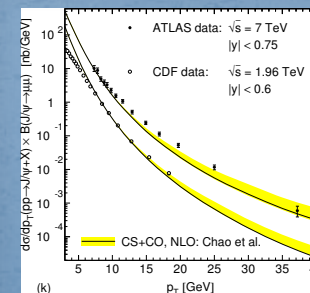
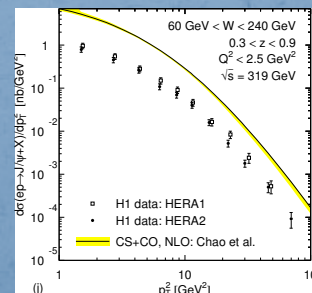
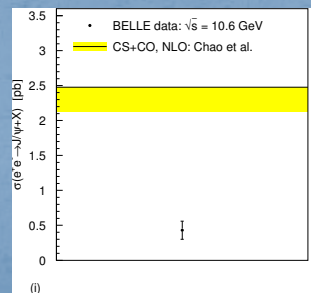
$pp$  polarization



Gong et al  
 $p_T > 5 \text{ GeV}$



Chao et al  
 $p_T > 7 \text{ GeV}$

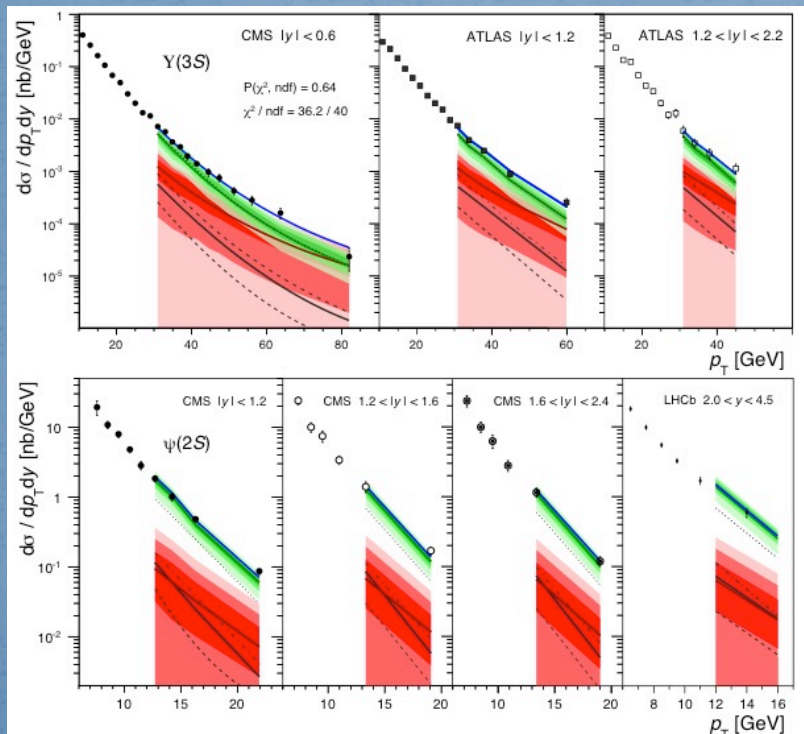


# Polarization: fitted or predicted?

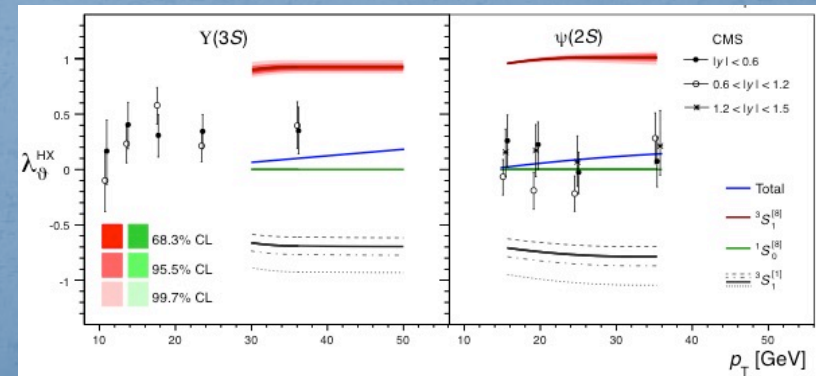
Fitting LDMEs to yields alone does not describe polarization

A combined fit to the two requires a higher  $p_T$  cut

This can be taken to the extreme, as in Faccioli et al where favorable  $p_T$  cut chosen



By looking only at excited states and  $p_T/m > 3$ , one can achieve a longitudinally polarized result



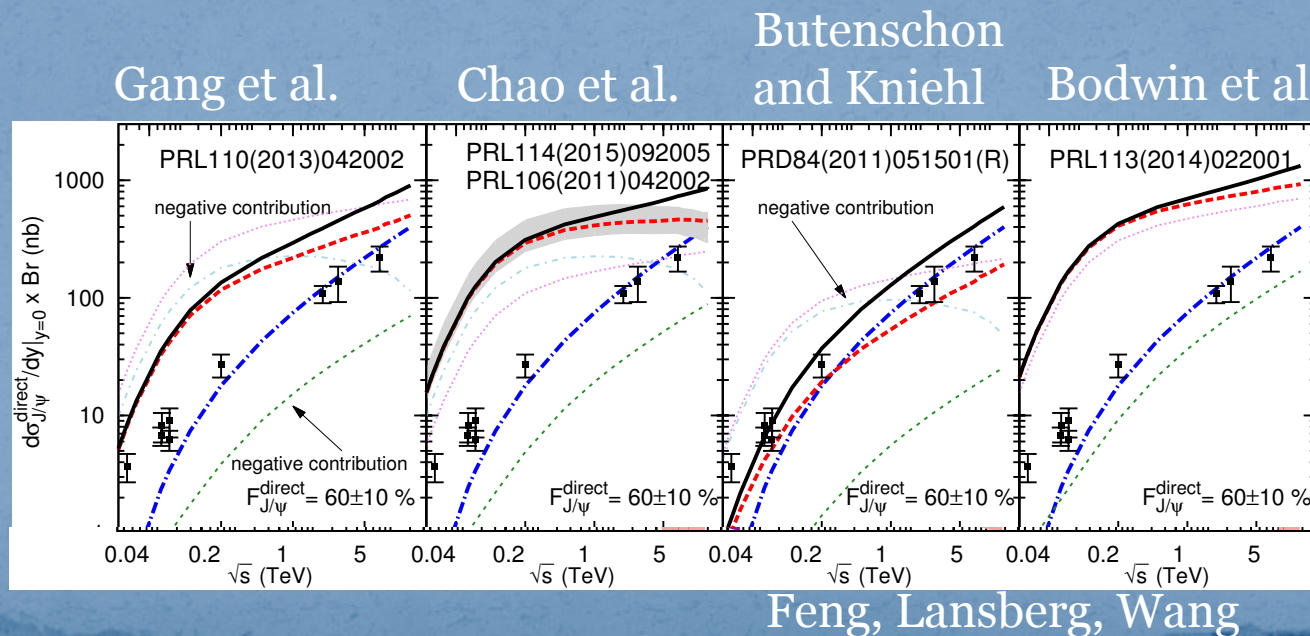


# LDMEs fit to yields fail to describe the $p_T$ -integrated rate (no big surprise)

LDMEs extracted from  $p_T$  distributions cannot describe center-of-mass energy dependence of  $y=0$  cross section

Lowest  $p_T$  cut ( $p_T > 3$  GeV) comes closest to data here yet is furthest off on polarization

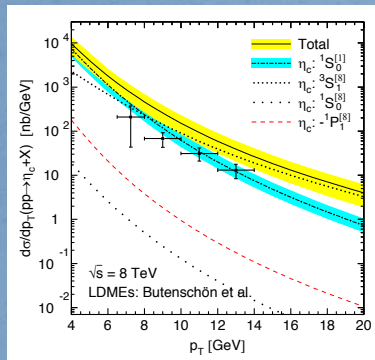
No low  $p_T$  resummation of logarithms included so far



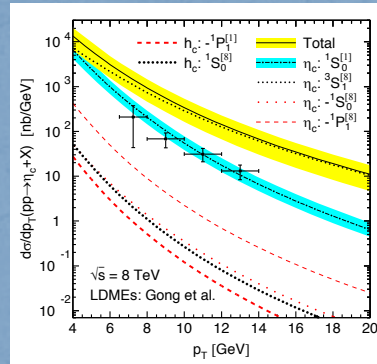
# J/ $\psi$ LDMEs do not describe $\eta_c$

If one takes heavy-quark spin symmetry LDMEs to apply to  $\eta_c$  production, all results so far overpredict LHCb  $\eta_c$  yields

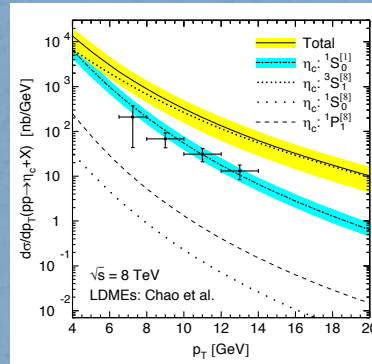
Butenschön  
& Kniehl



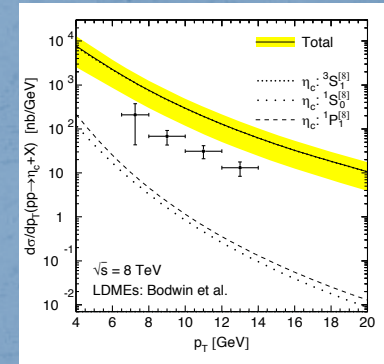
Chao et al



Gong et al



Bodwin et al



Calculations by Butenschön & Kniehl

If heavy-quark spin symmetry is given up, one can fit  $\eta_c$  LDMEs independently but then LDMEs are not universal, do not describe other processes



# Where do we go from here with NRQCD?

- Go beyond current NLO analyses
  - Adopt more parameters such as quark masses and scales
  - Resum logs at high and/or low  $p_T$
  - Look at associated production (although if single inclusive production not described why should associated production be better?)
  - Go to higher order
- Problem with NRQCD factorization?
  - Does not hold for polarized quantities (but this is a pillar of NRQCD)
  - Velocity expansion is too slow
- Some of the data are wrong (are theorists allowed to cherry pick data?)

# What about Y?

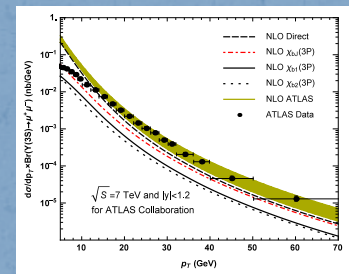
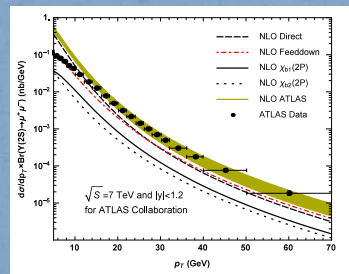
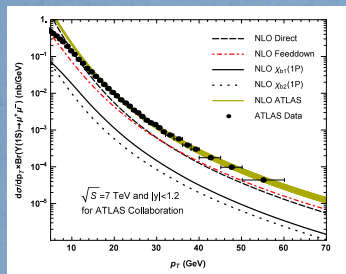
Larger mass, higher scale (smaller coupling) and slower velocity could  
Make Y a better candidate for NRQCD

Y production also allows for more free parameters to allow a description of  
both production and polarization – only Y(3S) has little wiggle room

Han et al,  $p_T > 15$  GeV fit

Lansberg et al

ATLAS  
 $|y| < 1.2$

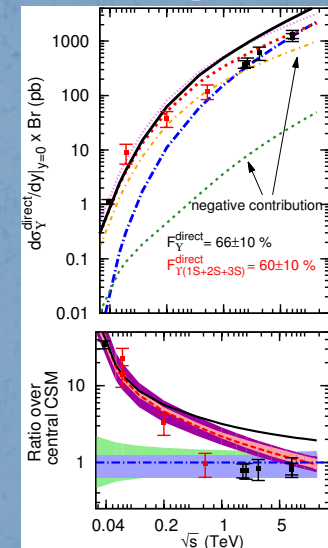
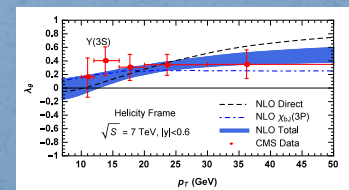
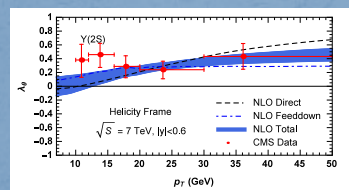
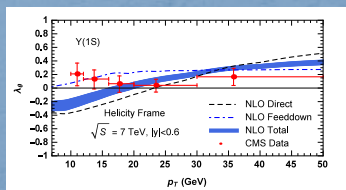


Y(1S)

Y(2S)

Y(3S)

CMS  
 $|y| < 0.6$



Other calculations can give similar agreement with data



# Color Evaporation Model

All quarkonium states treated like heavy quark pairs ( $Q = c, b$ ) below heavy hadron ( $H = D, B$ ) threshold

Color and spin are averaged over in pair cross section so color is 'evaporated' during transition from quark pair to quarkonium without changing kinematics

Distributions for quarkonium family members assumed identical

$$\sigma_Q^{\text{CEM}} = F_Q \sum_{i,j} \int_{4m^2}^{4m_H^2} d\hat{s} \int dx_1 dx_2 f_{i/p}(x_1, \mu^2) f_{j/p}(x_2, \mu^2) \hat{\sigma}_{ij}(\hat{s})$$

Values of quark mass,  $m$ , and scale,  $\mu$ , fixed from NLO calculation of heavy quark pair cross section

Scale factor  $F_Q$  fixed by comparison of  $\sigma_Q^{\text{CEM}}$  to energy dependence of  $J/\psi$  and  $Y$  cross sections,  $\sigma(x_F > 0)$  and  $Bd\sigma/dy|_{y=0}$  for  $J/\psi$ ,  $Bd\sigma/dy|_{y=0}$  for  $Y$ , only one  $F_Q$  for each state of quarkonium family

Spin always summed over so no previous predictions of polarization in CEM

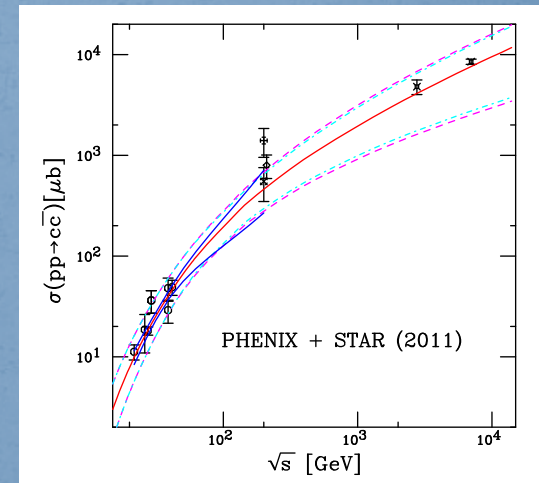
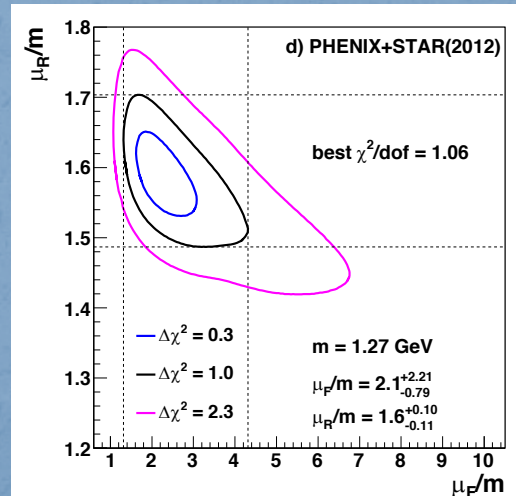
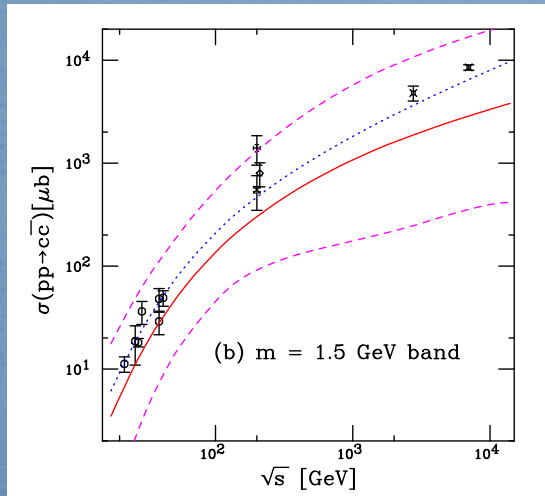
# Fitting charm cross section reduces theoretical uncertainties

Fit subset of total charm cross section data to obtain best fit values of  $\mu_F/m$ ,  $\mu_R/m$

$\Delta\chi^2 = 1$  gives uncertainty on scale parameters,  $\Delta\chi^2 = 2.3$  gives one standard deviation on total cross section

LHC results agree well with fit results although not included

No full NNLO cross section, likely to result in large corrections

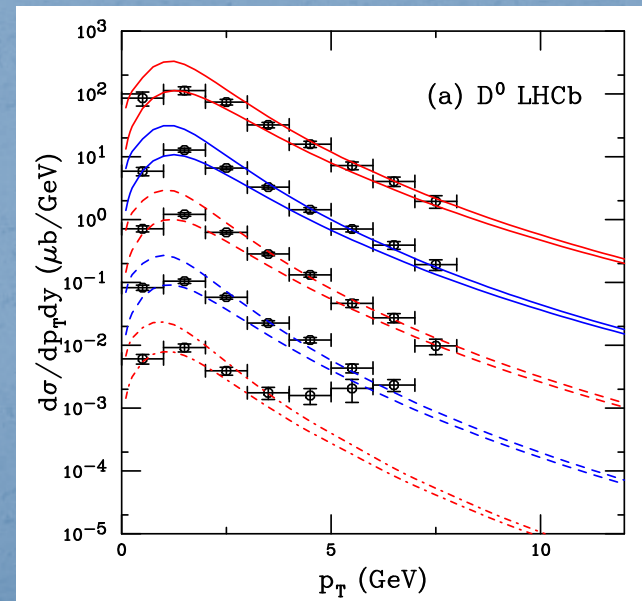
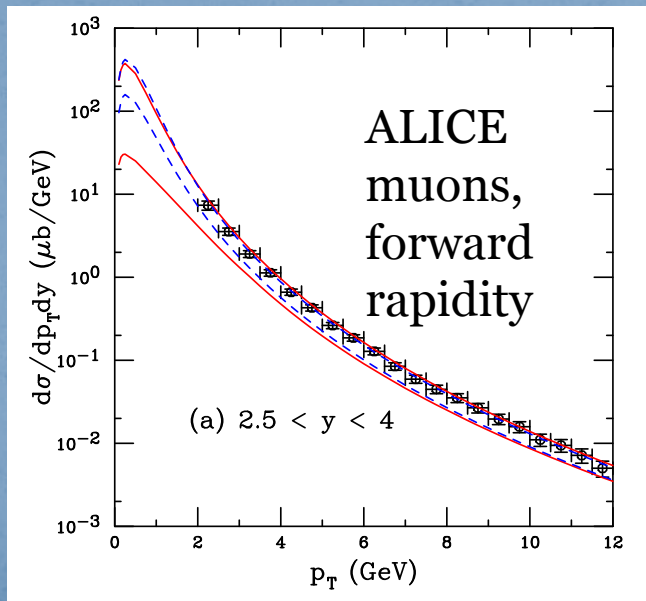




# Open Charm Results at 7 TeV

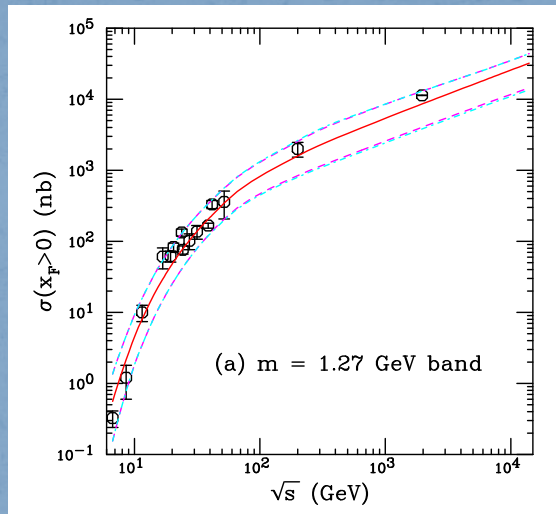
Excellent agreement with FONLL calculations of distributions in 7 TeV p+p collisions

Using results with  $m = 1.27$  GeV and fitted factorization and renormalization scales instead of fiducial variation of scale by factor of two around  $m = 1.5$  GeV reduces uncertainty band on forward rapidity muons and  $D^0$   $p_T$  distributions without reducing agreement with data



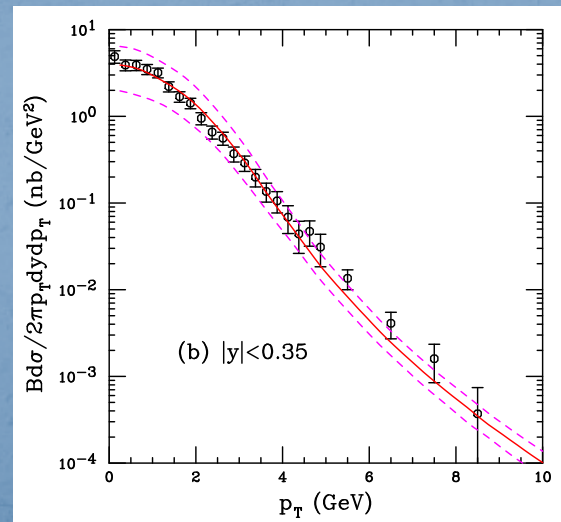
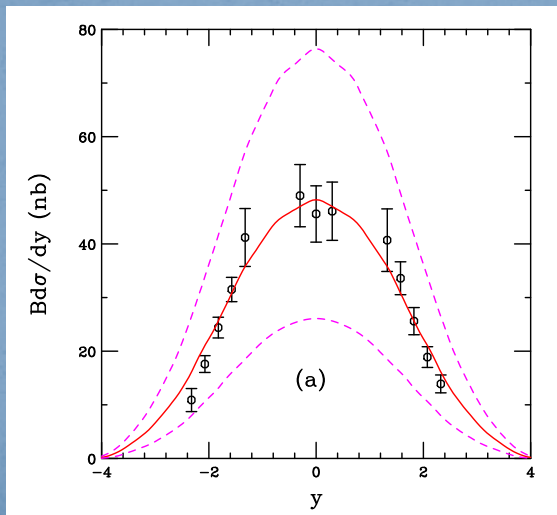
LHCb  
 $D^0$   
 $2 < y < 4.5$ ,  
 $\Delta y = 0.5$   
units

# J/ψ Results in CEM



Nelson, Frawley  
& RV

Data are from PHENIX  
at RHIC, 0.2 TeV





# Improved Color Evaporation Model

Relates average final state  $\psi$  momentum,  $\langle p_\psi \rangle$ , to quark pair momentum  $p$

$$\langle p_\psi \rangle = \frac{M_\psi}{M} p + \mathcal{O}(\lambda^2/m_c)$$

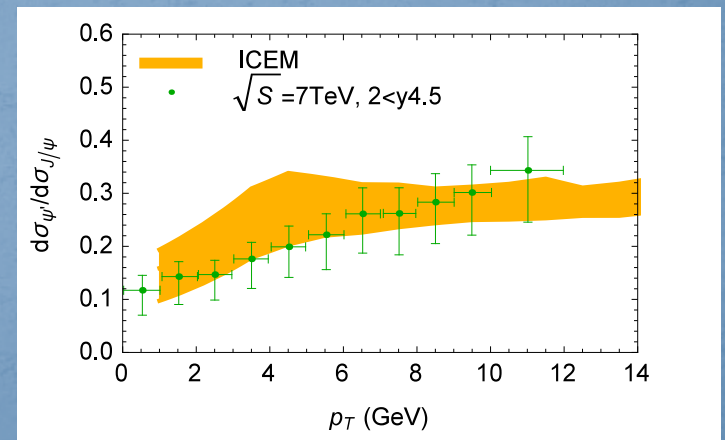
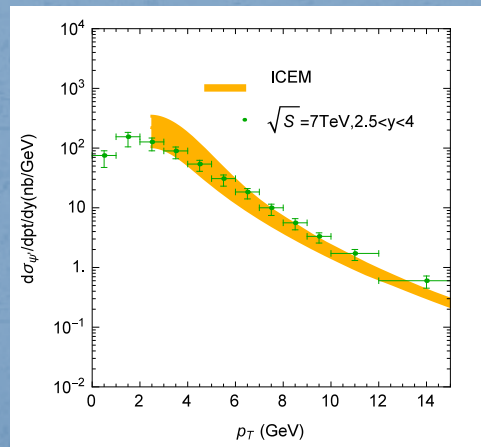
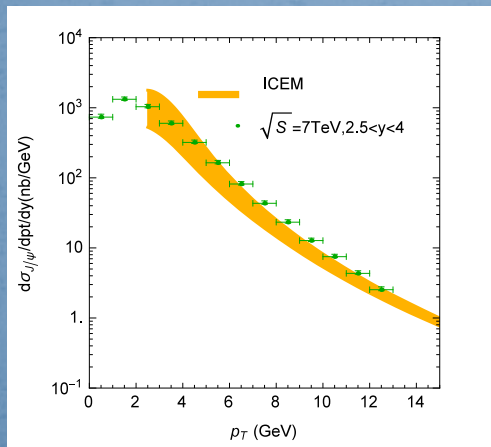
Lower limit on pair mass,  $M$ , has to be larger than  $\langle p_\psi \rangle$ , lower limit on CEM integration has to be increased to  $M_\psi$  so that the transverse momentum distribution becomes

$$\frac{d\sigma_\psi(p)}{dp_T} = F_\psi \int_{M_\psi}^{2m_D} dM \frac{M}{M_\psi} \frac{d\sigma_{c\bar{c}}(M, p')}{dM dp'_T} \Big|_{p'_T=(M/M_\psi)p_T}$$

$J/\psi$

$\psi'$

$\psi'/\psi$  ratio



# Polarization in the CEM

Work done with UC Davis grad student Vincent Cheung

Separate contributions by angular momentum but still integrate over color so no additional parameter fixing is required

First differentiated between longitudinal and transverse polarizations (PRD '17), current work calculates  $\lambda$  for different L and S combinations

Use Improved CEM to separate dependence of e.g. 1S and 2S states

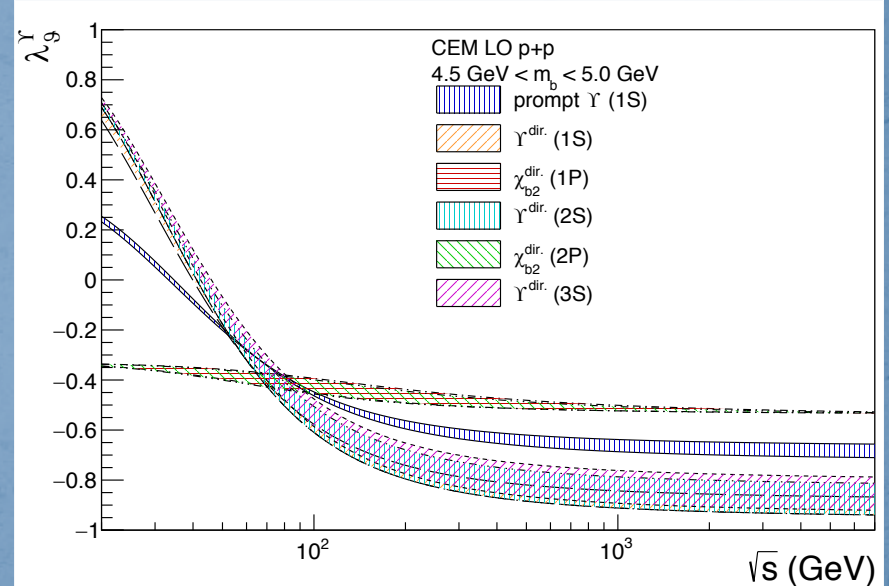
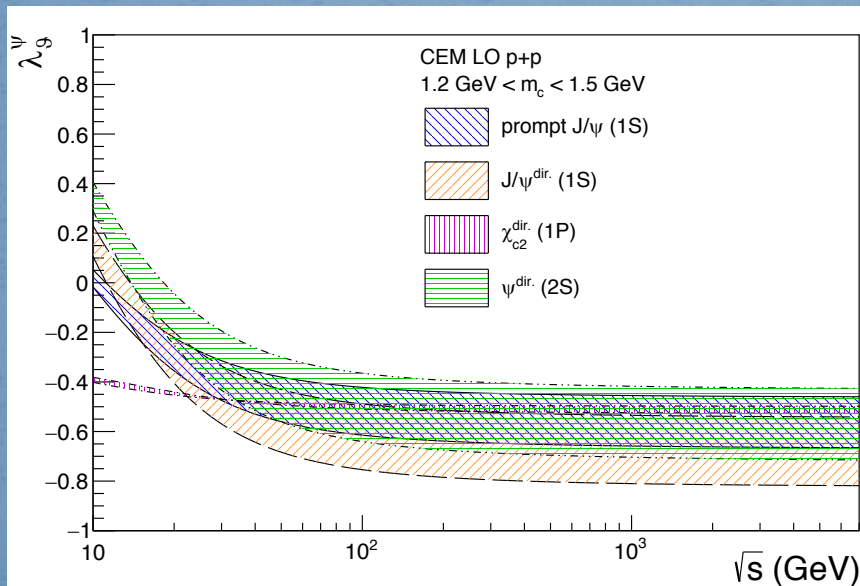
LO so far, thus results only available for center of mass energy and rapidity/ $x_F$ , for his dissertation, Vincent will calculate the  $p_T$  dependence of the polarization



# Polarization in the CEM: energy dep

Dependence of  $\lambda$  on center of mass energy, bands show quark mass dependence, largest source of uncertainty in calculation

$\chi_c$  and  $\chi_b$  states show smallest variation with energy and quark mass



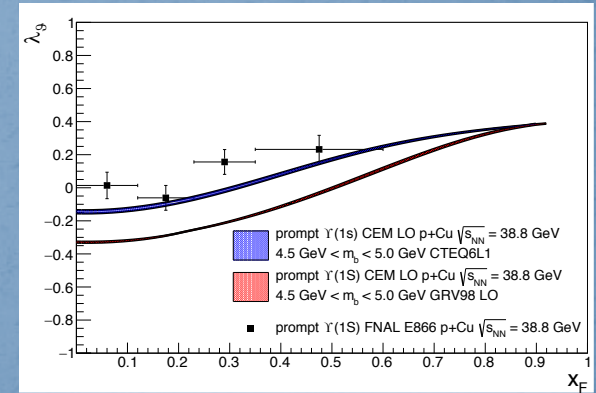
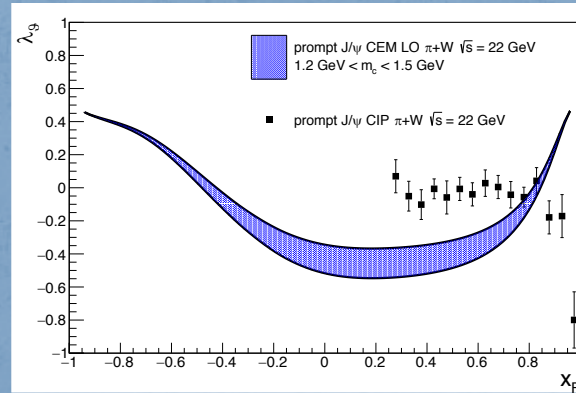
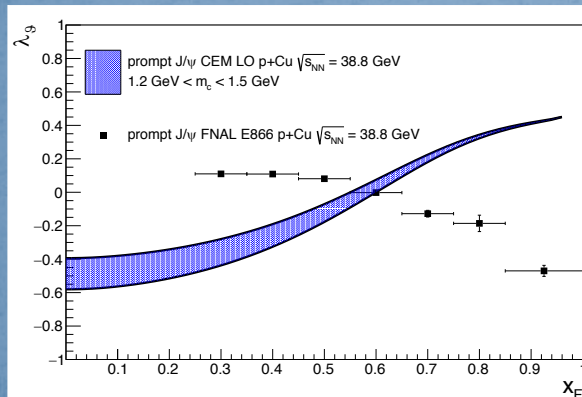
# Polarization in the CEM: $x_F$ dep

Comparison of calculations with E866 p+Cu J/ $\psi$  and  $\Upsilon$  data at 38.8 GeV and CIP  $\pi + W$  J/ $\psi$  data at 22 GeV

Calculations are LO CEM so no  $p_T$  is included thus there is a small kinematic mismatch between calculations

$\pi+A$  and  $p+A$  are quantitatively different at forward  $x_F$ , depend on PDF and energy

$\Upsilon$  agreement is rather good

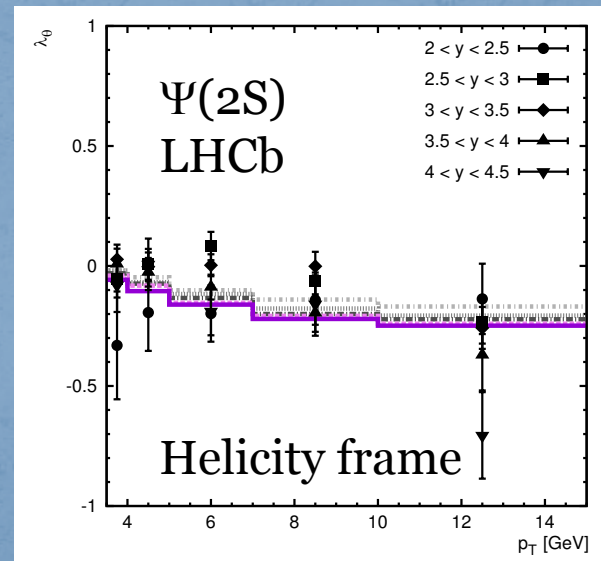
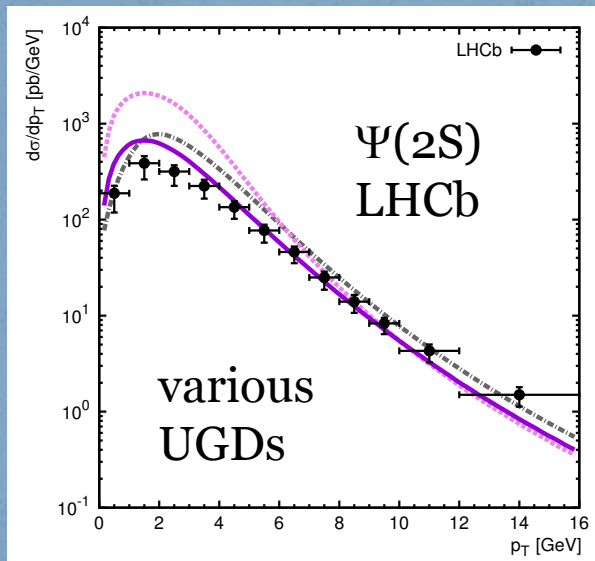




# $k_T$ Factorization for Quarkonium

Uses offshell color singlet and color octet matrix elements in NRQCD but with unintegrated gluon distributions to probe lower  $p_T$  without resummation

Fits to color octet LDMEs give smaller values than NRQCD with collinear factorization, better agreement with polarization data

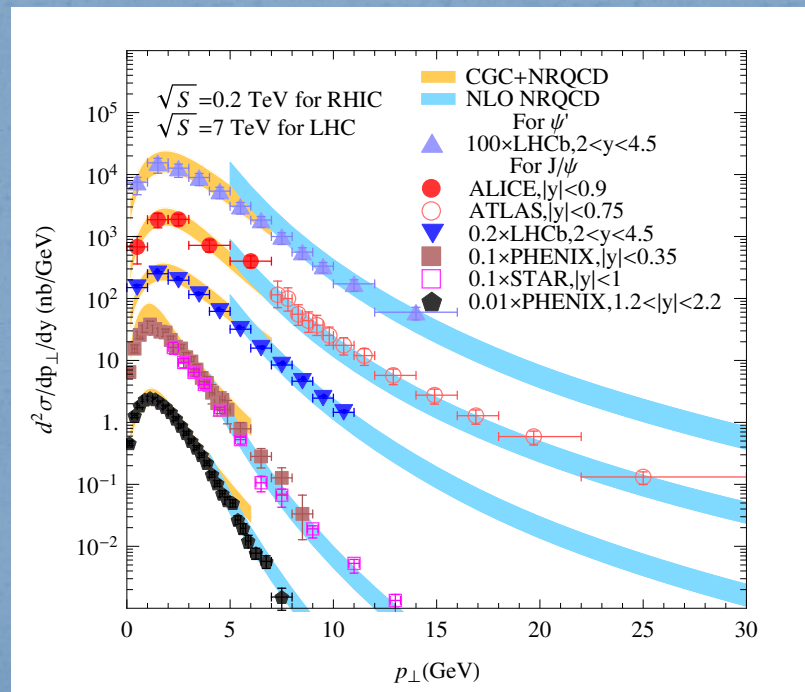


Baronov et al.

# Color glass condensate and NRQCD

Uses saturation model of gluon distributions in the proton with NRQCD  
color octet LDMEs

Saturation physics at low  $p_T$ , normal collinear factorization at high  $p_T$ , matching at intermediate  $p_T$



Ma and Venugopalan



# Summary of Quarkonium in p+p

- Production mechanism still not settled after more than 40 years
- NRQCD has long appeared promising but has many difficulties still remaining
- $k_T$  factorization and color glass model can address low  $p_T$ , different mix of LDMEs
- New recent work on color evaporation may be helpful

# When the nucleus is a target: p+Pb collisions at the LHC

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Nuclear matter effects quantified as ‘nuclear suppression factor’

$$R_{pA}(p_T, y) = \frac{1}{A} \frac{d\sigma_{pA}/dp_T dy}{d\sigma_{pp}/dp_T dy}$$



# Parton Density in the Initial State

Collinear factorization (DGLAP evolution): parton densities in the nucleus are modified based on global analyses of all data over a wide range of momentum fractions

- Nuclear DIS (electron, muon and neutrino-induced)
- Drell-Yan
- $\pi^0$  distributions
- High  $p_T$  jets (new, p+Pb 5 TeV data)
- $W^+$ ,  $W^-$  and  $Z^0$  production (new, p+Pb 5 TeV data)

Global analyses available from various groups: Eskola et al. (EKS98, EPS09, EPPS16 – latest); nDS, nDSg, DSSZ; nCTEQ sets; HKN sets

Saturation, Color Glass Condensate: assumes  $k_T$  ordering and evolution in  $x$ , important at low  $x$  and low  $Q^2$ ,  $Q^2 < Q_{\text{sat}}^2$

At high gluon density, recombination of gluons,  $2 \rightarrow 1$ , competes with gluon emission

$Q_{\text{sat}}$  depends on center of mass energy,  $x$ , expected to grow as  $A^{1/3}$  for nuclei  
Hybrid models used to interpolate between low and high  $x$  regimes



# Cold Matter Energy Loss

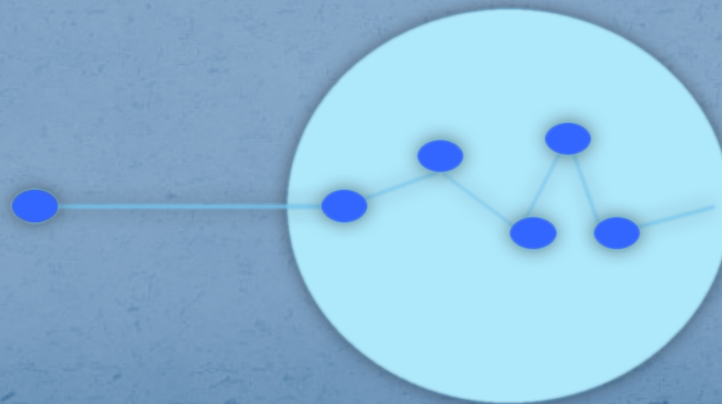
Energy loss in medium: Both initial state (before hard scattering) and final state (after hard scattering) have been considered

$$R_{pA} < 1 \text{ (forward rapidity, high } p_T)$$

Cronin effect: Increase in average transverse momentum of the final state due to multiple scattering in the medium

$$R_{pA} > 1 \text{ (backward rapidity, low } p_T)$$

Energy loss and Cronin are intertwined and effectively one can cause the other: a loss at high momentum can result in enhancement at low





# Additional Cold Matter Effects present for Quarkonium: Size Matters

## Nuclear Absorption:

- After heavy flavor pair produced, it can break up due to interactions with nucleons
- Relevant for regions of phase space where quarkonium state is produced in matter, e.g. backward rapidity at the LHC and RHIC

## Comovers:

- Quarkonium states break up due to interactions with produced particles
- More loosely bound states are more likely to break up
- Effect increases with collision centrality (comover density)

Both absorption and comover interaction cross sections expected to depend on quarkonium size

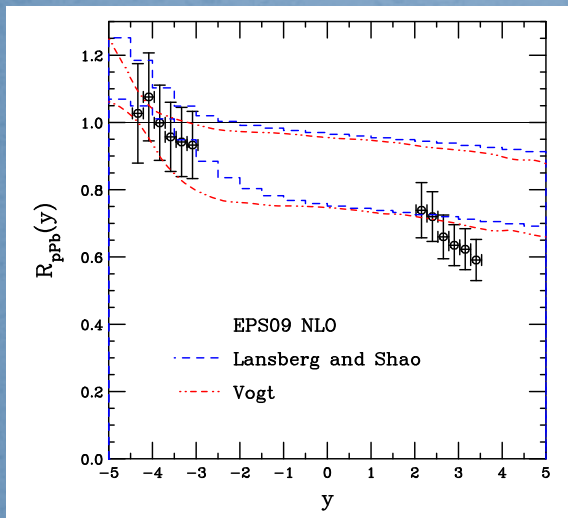
$$\sigma_C/\sigma_C \propto (R_C/R_C)^2$$

# Suppression in p+Pb at 8 TeV: $y$ dep

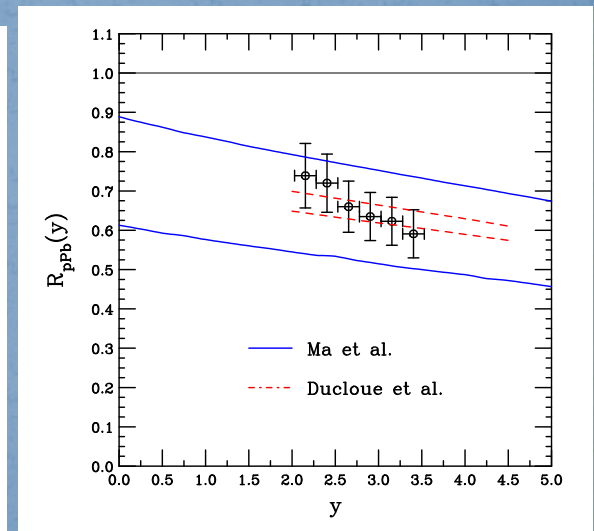
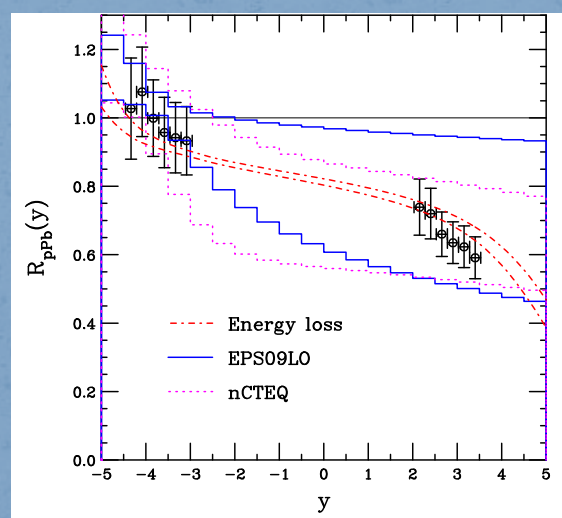
All calculations do a reasonable job of describing preliminary ALICE data

EPS09 NLO is marginal at forward rapidity due to difference in low  $x$  behavior of CTEQ6M and CTEQ61L

CGC+NRQCD band is larger because different color states shown separately



Collinear factorization: shadowing only  
and energy loss only



CGC+CEM (Ducloue et al)  
CGC+NRQCD (Ma et al)

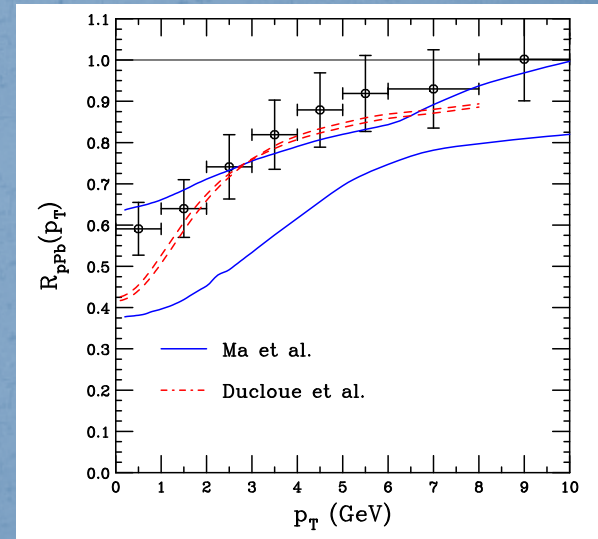
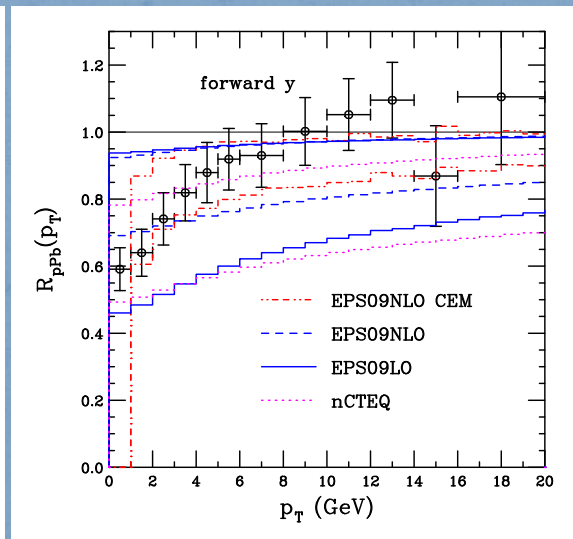
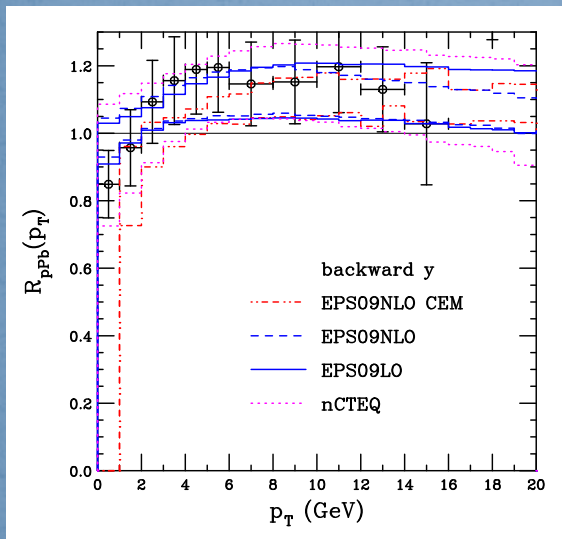


# Suppression in p+Pb at 8 TeV: $p_T$ dep

All calculations do a reasonable job of describing preliminary ALICE data

Shadowing uncertainty bands are smaller vs.  $p_T$  at backward rapidity

CGC+NRQCD and CGC+CEM calculations have different curvature at low  $p_T$



Collinear factorization: shadowing only  
and energy loss only (RV, Lansberg and Shao)

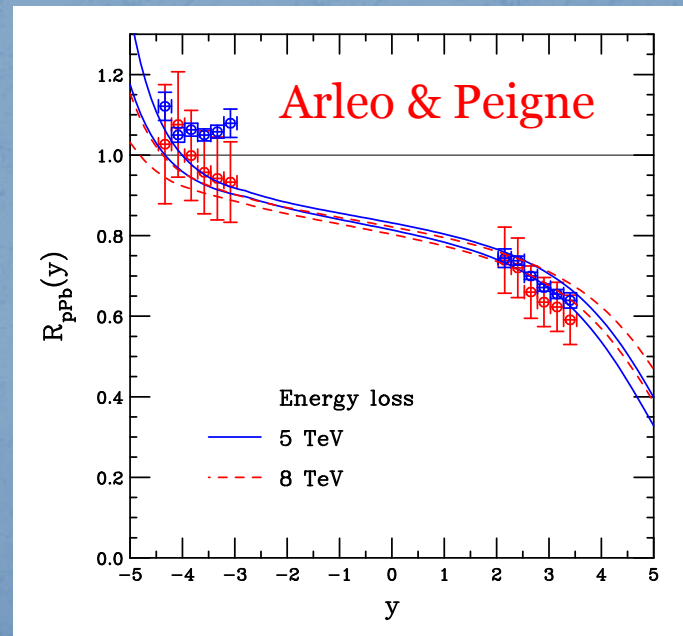
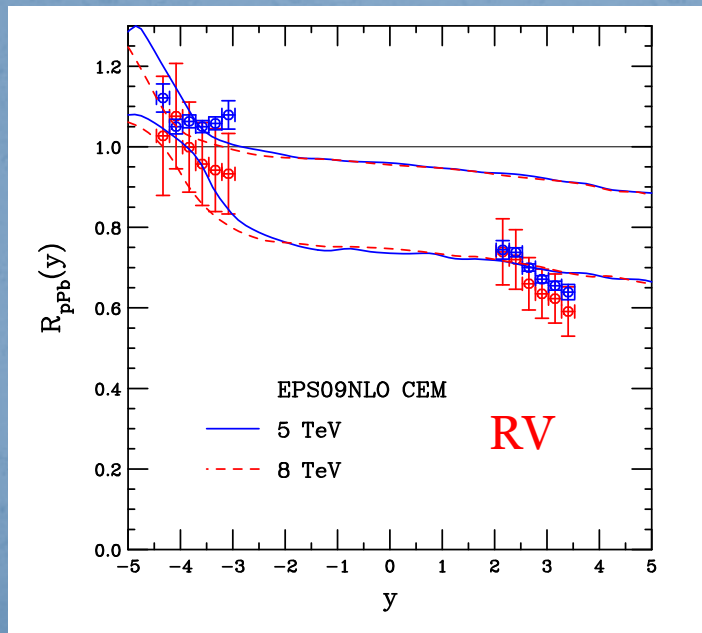
CGC+CEM (Ducloue et al)  
CGC+NRQCD (Ma et al)

# 5.02 TeV vs. 8.16 TeV

Comparison is actually 5 vs. 8 TeV, results are shown for cases where the same input models were used in both cases

Only small differences seen in calculations at the two energies, EPS09 NLO CEM is mostly different at backward rapidity, shadowing is maximal at forward  $y$

Data are also rather similar, perhaps more dependence on  $y$  in backward region

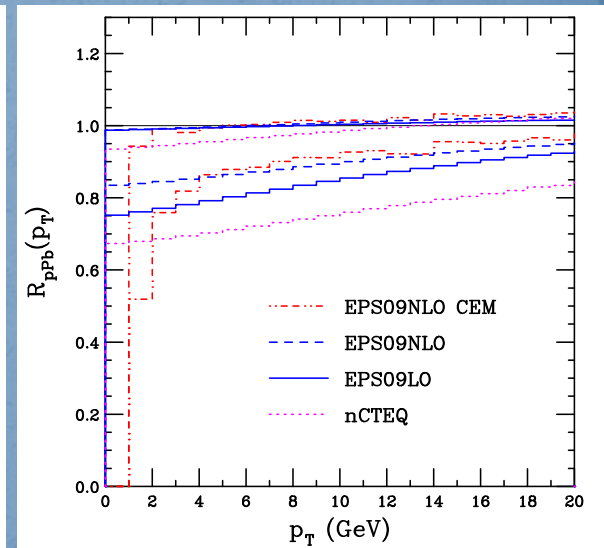
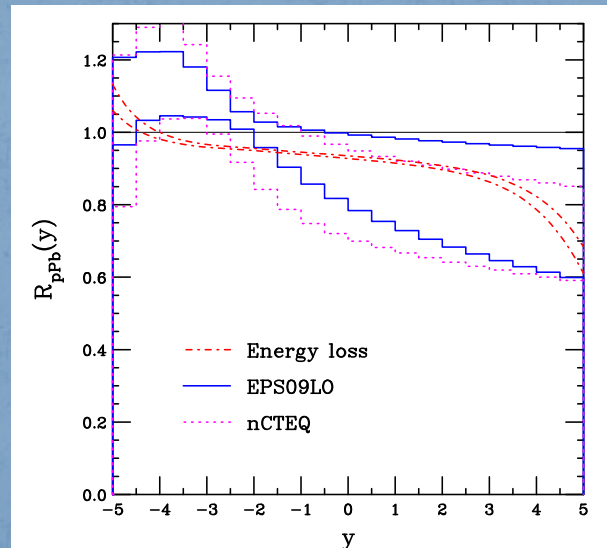
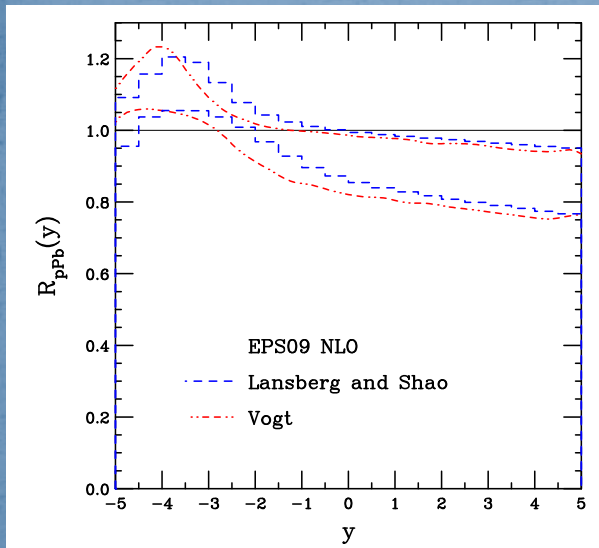




# Predictions for $Y(1S)$ inclusive

Uncertainty bands are smaller for Upsilon results because mass scale is larger, more evolution of nPDFs, somewhat higher  $x$  as well

All calculations are within uncertainties of each other

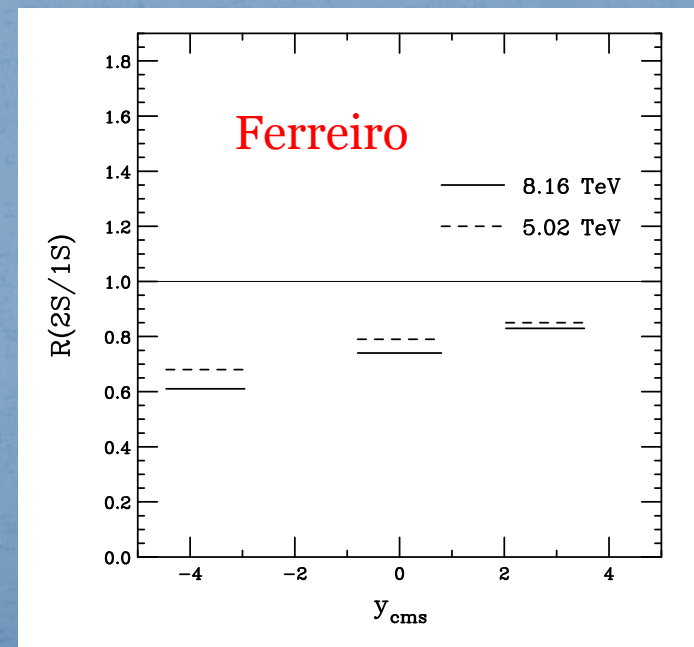
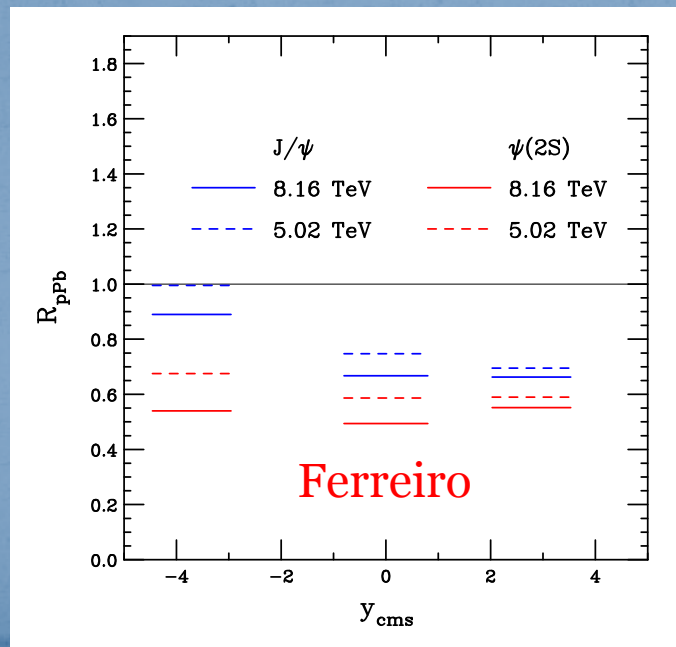


RV, Landsberg and Shao, Arleo and Peigne

# Suppression by comovers

Left side compares  $R_{pPb}$  in different rapidity regions for the two energies, biggest difference is at backward rapidity, at forward rapidity, difference is negligible

Right side shows double ratio,  $\psi(2S)/\psi(1S)$ , for the two energies, same trend seen





# p+Pb Quarkonium Summary

- Minimum bias results likely too dilute to apply hot matter models but high multiplicity central p+Pb events may be more relevant
- Multiple models can explain the trends in the quarkonium data
- Higher precision data are needed to separate effects and eliminate models – as ever the case