

Errors in Model to Data Comparisons

Data Errors (some examples)

Model Errors (an example)

Discussion ...

Workshop on Precision Spectroscopy of QGP

Institute for Nuclear Theory, UW

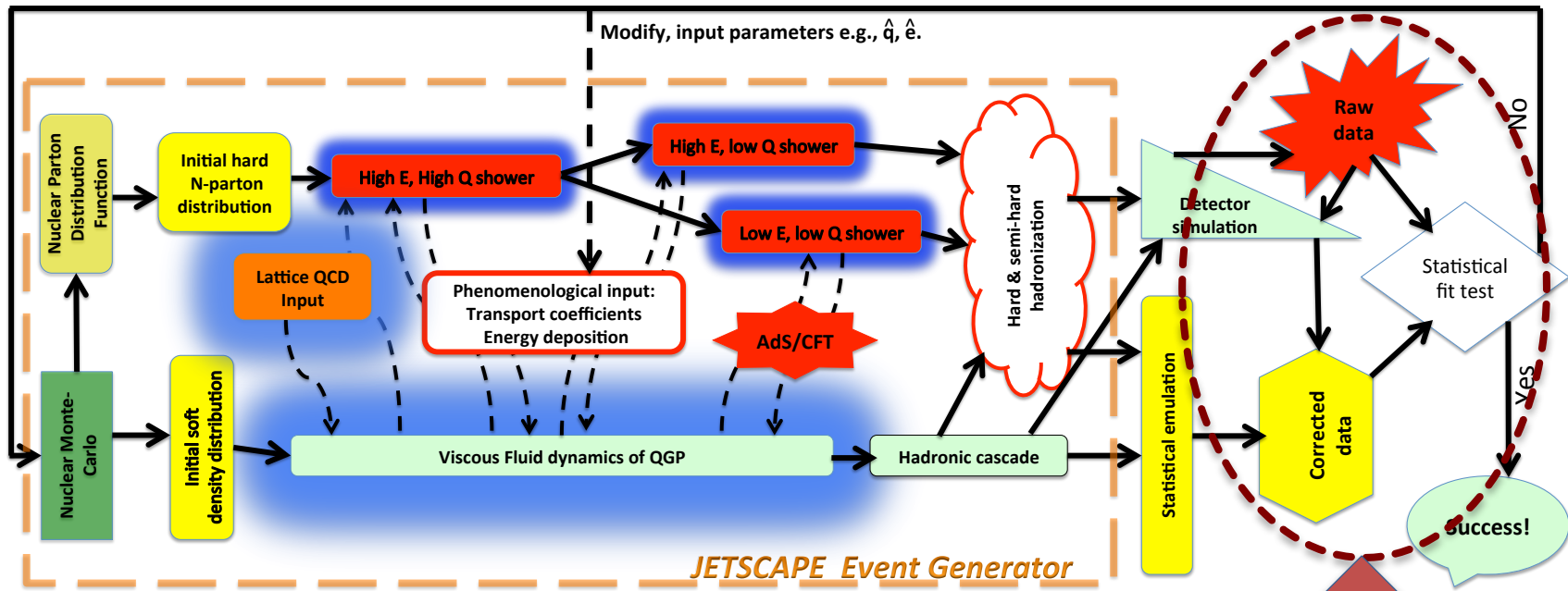
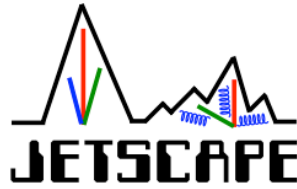
May 25, 2017

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Motivation by



- This talk is about the last stage of the diagram



MADAI treatment of errors

- Seminal Model-to-Data-Comparison by Novak,...Pratt, et al. (14)
<http://link.aps.org/doi/10.1103/PhysRevC.89.034917>
- But underneath the hood of this Ferrari are some squirrels
MADAI errors pegged at 6% (and 3%)

TABLE II. Observables used to compare models to data.

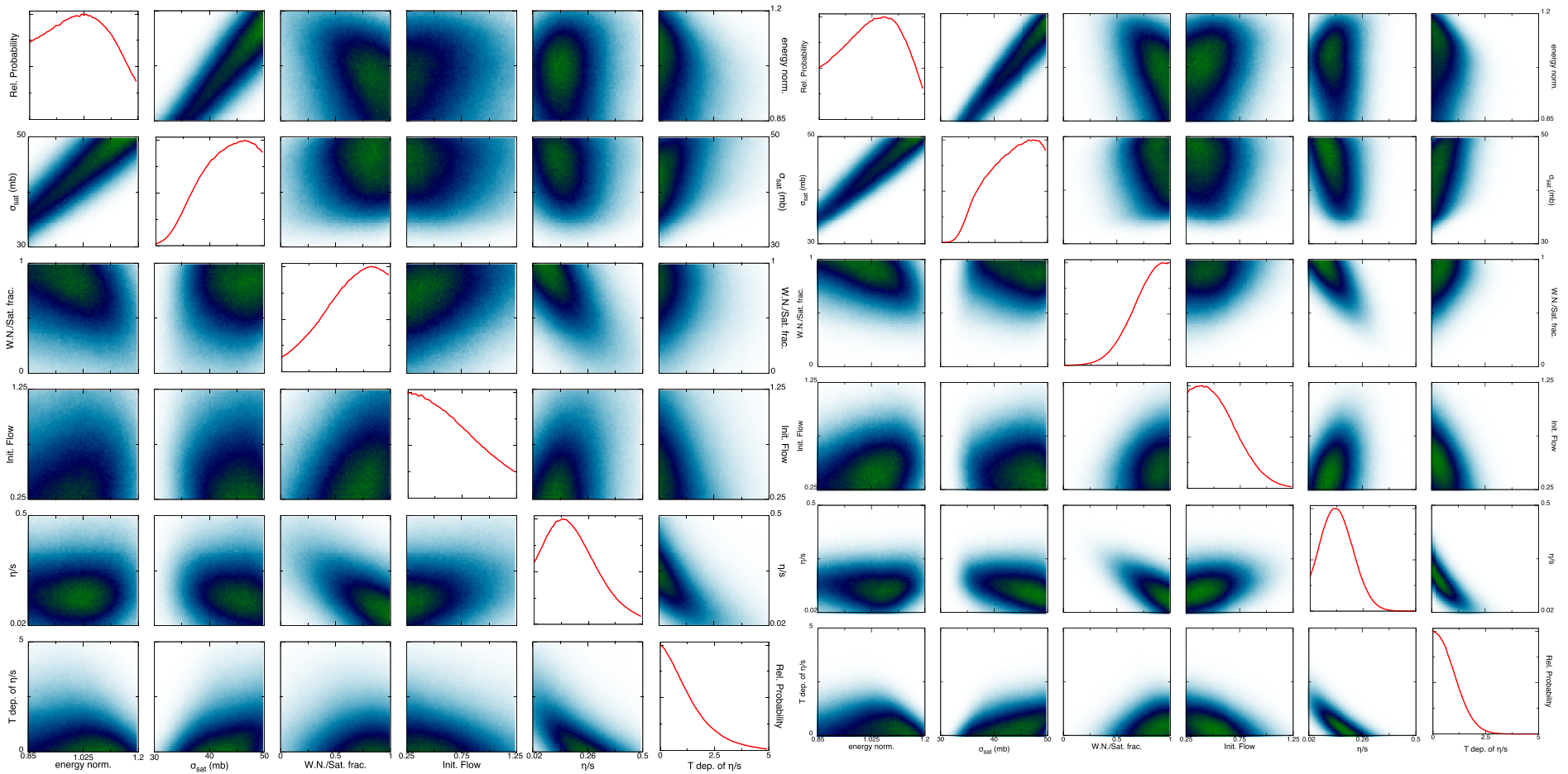
Observable	p_t weighting	Centrality (%)	Collaboration	Uncertainty (%)	Reduced uncertainty
$v_{2,\pi^+\pi^-}$	Average over 11 p_t bins from 160 MeV/c to 1 GeV/c	20–30	STAR ¹ [52]	12	6%
R_{out}	Average over 4 p_t bins from 150–500 MeV/c	0–5	STAR [53]	6	3%
R_{side}	Average over 4 p_t bins from 150–500 MeV/c	0–5	STAR [53]	6	3%
R_{long}	Average over 4 p_t bins from 150–500 MeV/c	0–5	STAR [53]	6	3%
R_{out}	Average over 4 p_t bins from 150–500 MeV/c	20–30	STAR [53]	6	3%
R_{side}	Average over 4 p_t bins from 150–500 MeV/c	20–30	STAR [53]	6	3%
R_{long}	Average over 4 p_t bins from 150–500 MeV/c	20–30	STAR [53]	6	3%
$\langle p_t \rangle_{\pi^+\pi^-}$	200 MeV/c < p_t < 1.0 GeV/c	0–5	PHENIX [54]	6	3%
$\langle p_t \rangle_{K^+K^-}$	400 MeV/c < p_t < 1.3 GeV/c	0–5	PHENIX [54]	6	3%
$\langle p_t \rangle_{p\bar{p}}$	600 MeV/c < p_t < 1.6 GeV/c	0–5	PHENIX [54]	6	3%
$\langle p_t \rangle_{\pi^+\pi^-}$	200 MeV/c < p_t < 1.0 GeV/c	20–30	PHENIX [54]	6	3%
$\langle p_t \rangle_{K^+K^-}$	400 MeV/c < p_t < 1.3 GeV/c	20–30	PHENIX [54]	6	3%
$\langle p_t \rangle_{p\bar{p}}$	600 MeV/c < p_t < 1.6 GeV/c	20–30	PHENIX [54]	6	3%
$\pi^+\pi^-$ yield	200 MeV/c < p_t < 1.0 GeV/c	0–5	PHENIX [54]	6	3%
$\pi^+\pi^-$ yield	200 MeV/c < p_t < 1.0 GeV/c	20–30	PHENIX [54]	6	3%

^aTo account for nonflow correlations, the value of v_2 was reduced by 10% from the value reported in Ref. [52].

MADAI Error Comparison

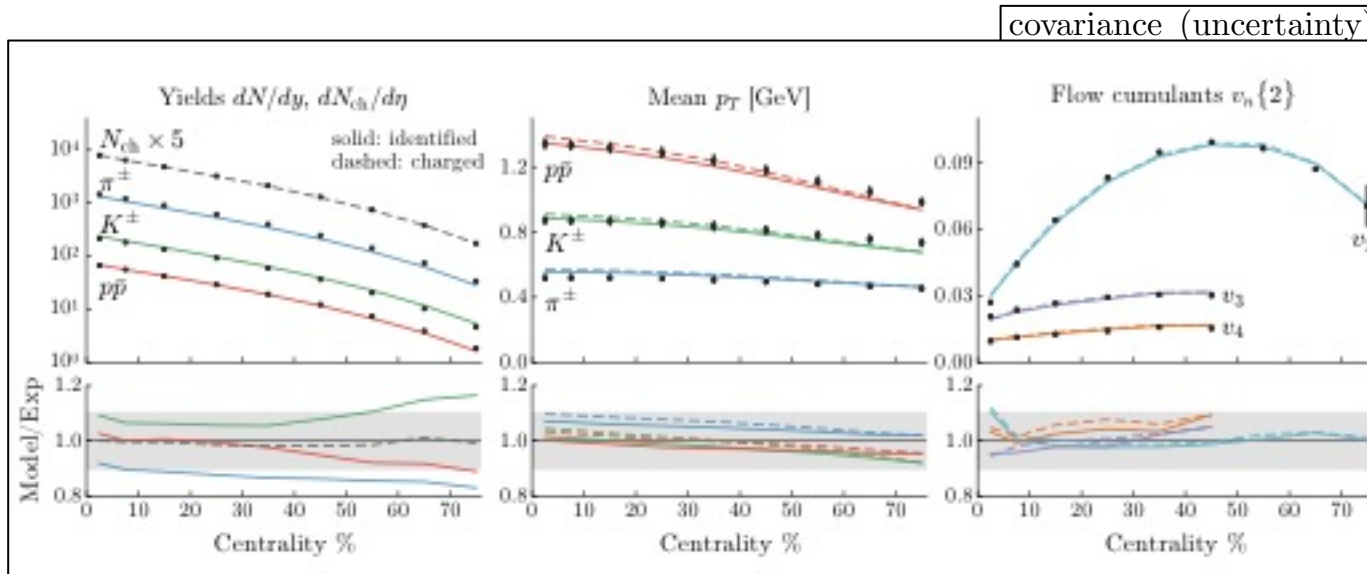
with 6% errors

with 3% errors



The Dukes

- Bernard, ... Bass, et al (16) improved upon work of MADAI
<https://journals.aps.org/prc/abstract/10.1103/PhysRevC.94.024907>
- But not with error treatment – their Tesla still has a few squirrels
 principle component errors pegged at 10%



covariance (uncertainty) matrix. As in previous work
 constant fractional uncertainty on
 so that the covariance matrix

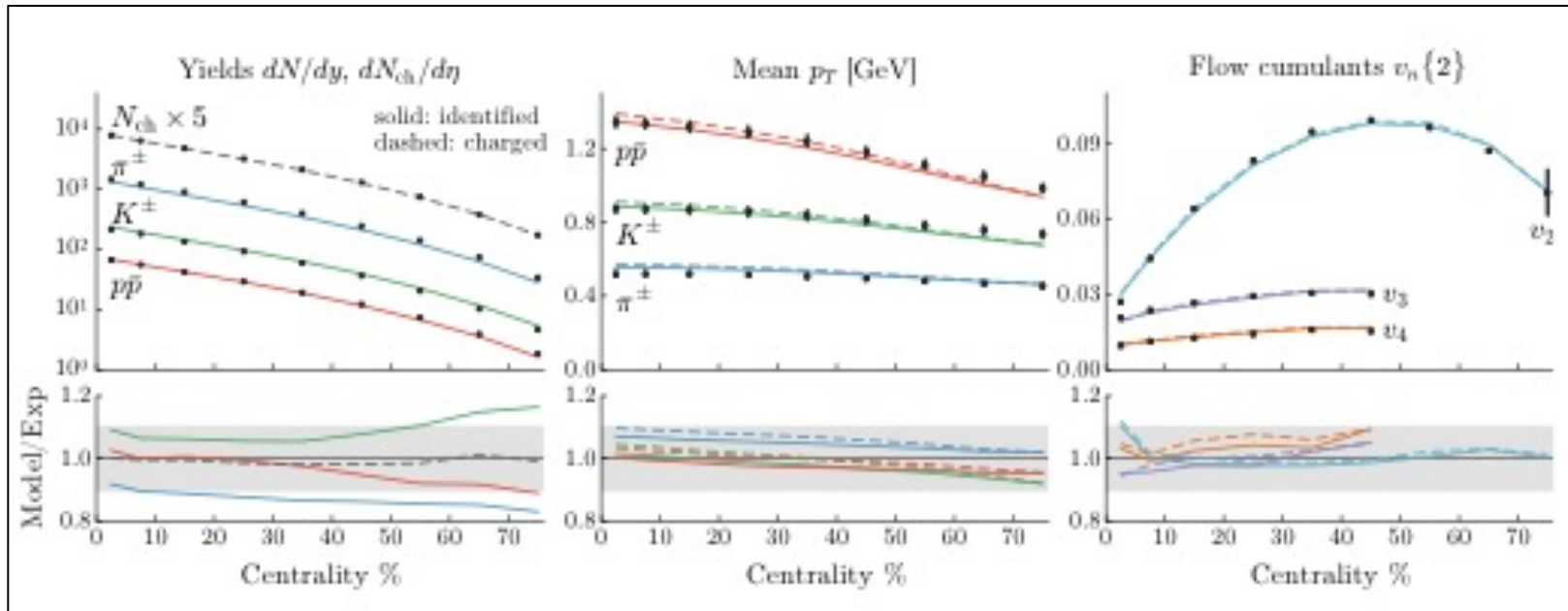
$$\text{diag}(\sigma_z^2 \mathbf{z}_{\text{exp}}), \quad (30)$$

present study. This is a simple
 vatively account for the various
 the experimental data, model
 r predictions. It certainly limi-
 tative uncertainties in the final
 d is an obvious target for im-

provement in future studies.

The Dukes

- Q: Why were relative errors set to 10% ?
- A: That's what was needed !



Comparison of most probable model results to ALICE data

An Earlier Foray into Model-2-Data-Comparisons

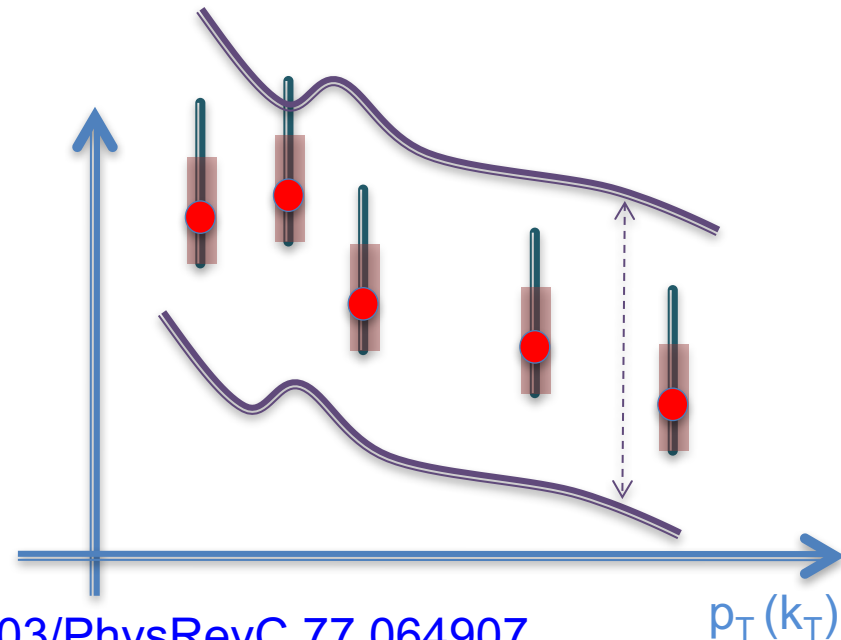
- CHIMERA = Comprehensive Heavy Ion Model Evaluation & Reporting Algorithm
- Non-Bayesian generation χ^2 map in T- η /s space with simultaneous comparison PHENIX and STAR spectra, flow, HBT
<https://dx.doi.org/10.1103/PhysRevC.87.044901> (RAS 2013)
- Just a Subaru, but no squirrels under the hood
- χ^2 evaluations performed using full statistical and systematic errors *as reported by PHENIX and STAR*

Systematic Errors in CHIMERA

- Evaluate χ^2 from model & data, accounting for type A & B errors

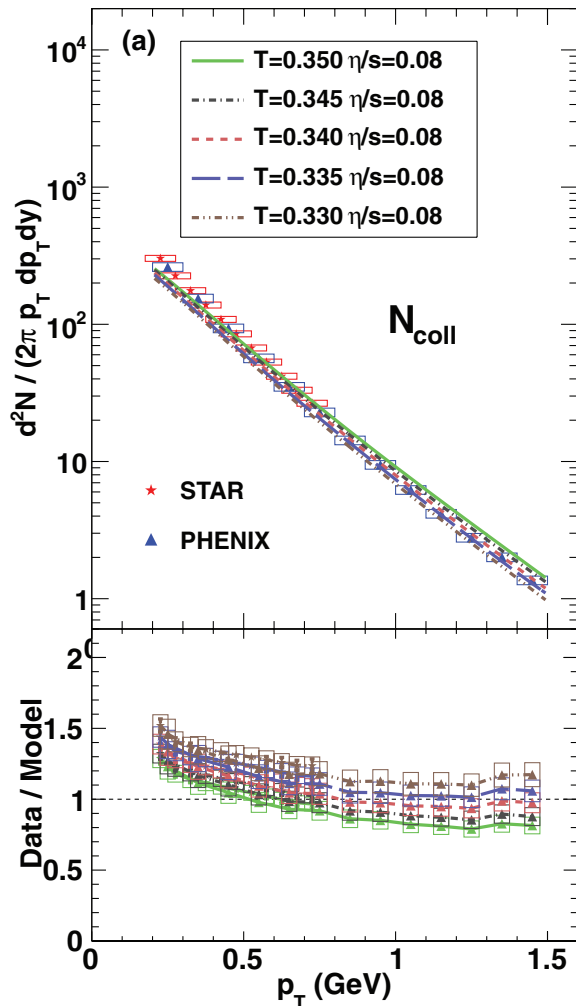
- A type: uncorrelated (σ_i)
- B type: correlated frac. (σ_b)
- C type: normalization (σ_c)
- D type: correlated tilt (not considered)

$$\bar{\chi}^2(\epsilon_b, \epsilon_c, p) = \left[\left(\sum_{i=1}^n \frac{(y_i + \epsilon_b \sigma_{b_i} + \epsilon_c y_i \sigma_c - \mu_i(p))^2}{\bar{\sigma}_i^2} \right) + \epsilon_b^2 + \epsilon_c^2 \right]$$



Error definitions based on <https://dx.doi.org/10.1103/PhysRevC.77.064907>

How well did it work?



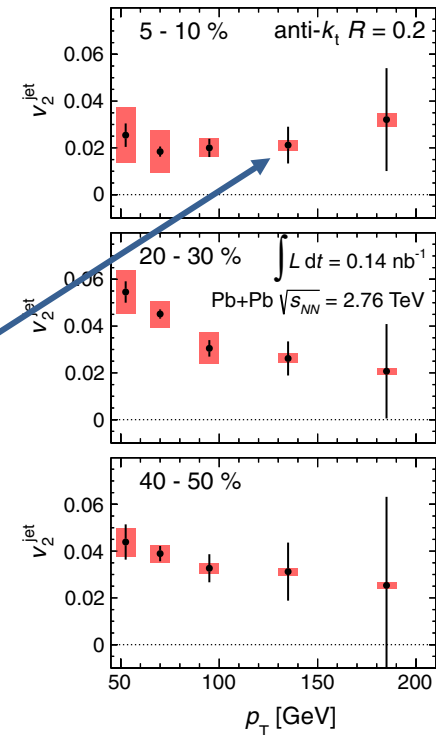
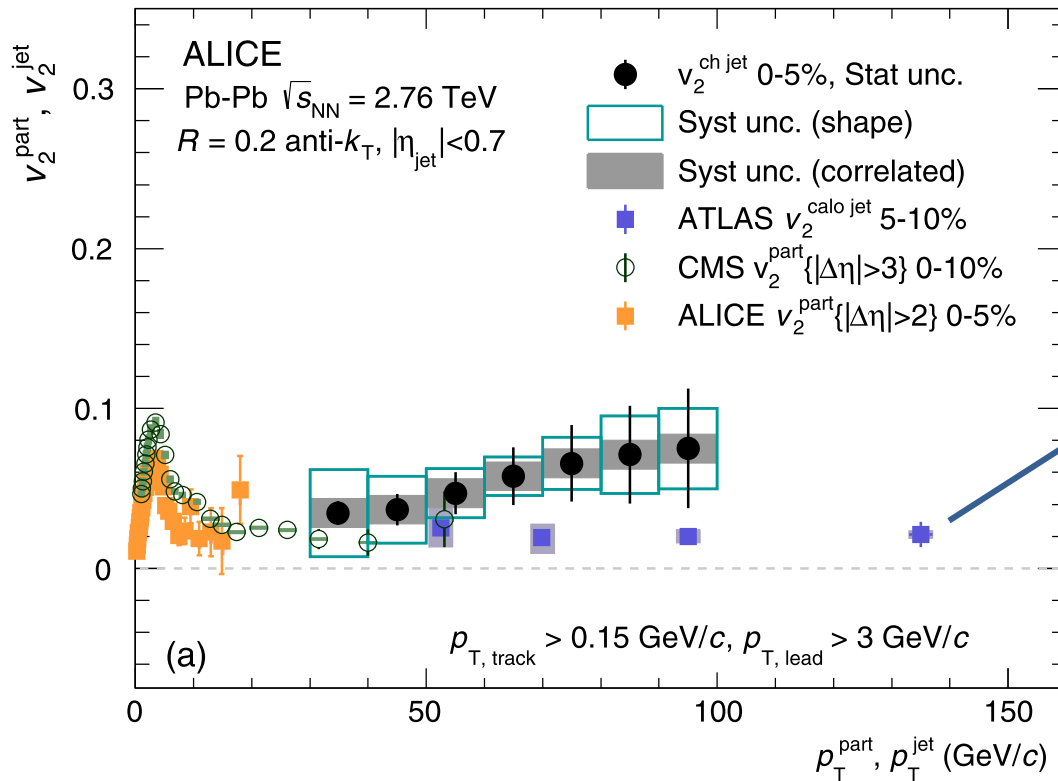
T_{cent} (GeV)	$\chi_{ndf}^2 N_{\text{coll}}$	
	PHENIX	STAR
0.350	13.8	2.30
0.345	2.77	1.75
0.340	15.7	8.15
0.335	75.9	6.87
0.330	60.3	6.94

- Separate systematic errors allowed **STAR** and **PHENIX** data to shift independently to achieve reasonable χ^2 values

Jet errors will be more challenging?

<http://dx.doi.org/10.1016/j.physletb.2015.12.047>

<https://doi.org/10.1103/PhysRevLett.111.152301>



- One cannot directly compare ALICE results to ATLAS (without a model)
- But one can ask whether deviation from zero is significant

ALICE $v_2^{\text{ch,jet}}$ Error Analysis

	p_T^{chjet} (GeV/c)	Uncertainty on $v_2^{\text{ch jet}}$					
		30-40	60-70	80-90	30-40	60-70	80-90
	Centrality (%)	0-5			30-50		
Shape	Unfolding	0.017	0.012	0.016	0.016	0.011	0.015
	p_T^{chjet} -measured	0.013	\ll stat	\ll stat	0.024	\ll stat	\ll stat
	$\rho_{\text{ch}}(\varphi)$ fit	0.015	\ll stat	0.016	\ll stat	\ll stat	\ll stat
Total		0.027	0.012	0.023	0.029	0.011	0.015
Correlated	Tracking	0.009	0.009	0.009	0.007	0.007	0.007
	p_T^{chjet} -unfolded	\ll stat	\ll stat	\ll stat	\ll stat	\ll stat	\ll stat
Total		0.009	0.009	0.009	0.007	0.007	0.007

χ^2 for the hypothesis $v_2^{\text{ch jet}} = \mu_i$ is calculated by minimizing

$$\tilde{\chi}^2(\varepsilon_{\text{corr}}, \varepsilon_{\text{shape}}) = \left[\left(\sum_{i=1}^n \frac{(v_{2,i} + \varepsilon_{\text{corr}}\sigma_{\text{corr},i} + \varepsilon_{\text{shape}} - \mu_i)^2}{\sigma_i^2} \right) + \varepsilon_{\text{corr}}^2 + \frac{1}{n} \sum_{i=1}^n \frac{\varepsilon_{\text{shape}}^2}{\sigma_{\text{shape},i}^2} \right] \quad (16)$$

also based on <https://dx.doi.org/10.1103/PhysRevC.77.064907>, yields $p(m=0) = 0.009$, ask Redmer for more details...

Question for the rest of us: is there a better approach?

Why not produce the full co-variance matrix?

- See upcoming publication by NIFFTE = Neutron Induced Fission Fragment Tracking Experiment <http://niffte.calpoly.edu>
- Co-variant inputs include
 - Beam pile-up
 - Target contamination
 - Tracking efficiency
 - Analysis cuts
- But what would we do with it?

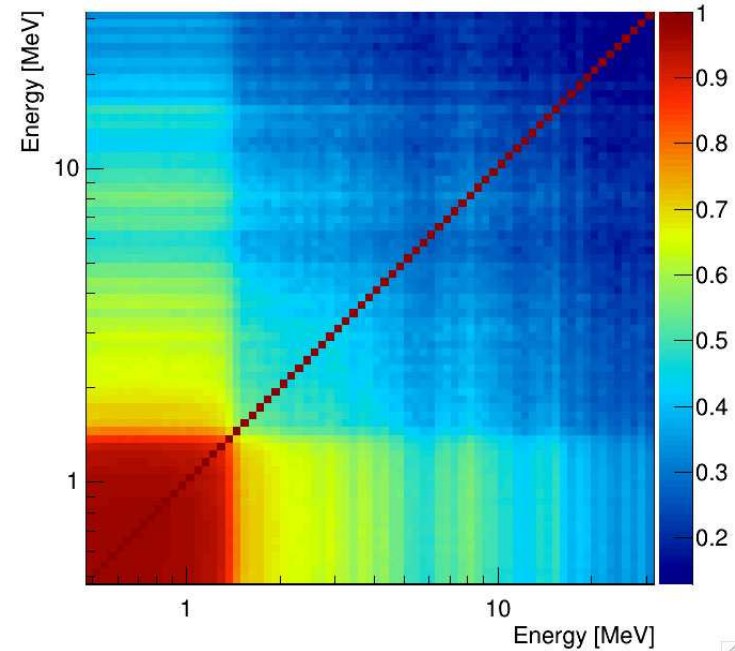
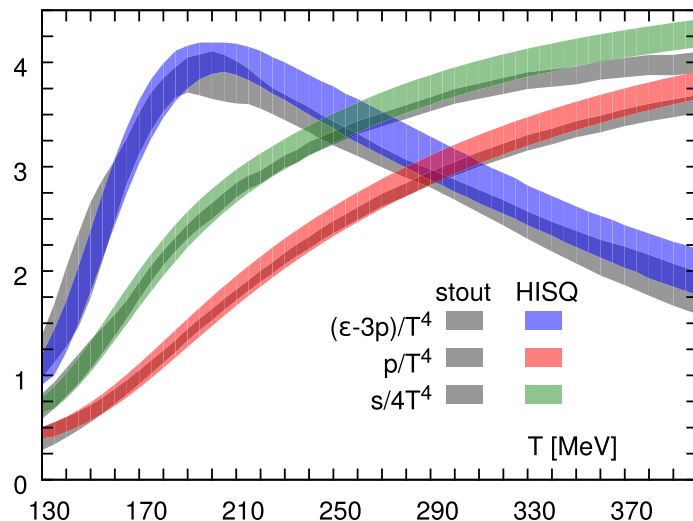


FIG. 16. The $^{238}\text{U}(n,f)/^{235}\text{U}(n,f)$ correlation matrix measured in this work. At low neutron energy, the contaminant correction becomes the largest source of uncertainty, resulting in a large correlated region in the correlation matrix. The contaminant correction is a fixed value at all energies and as the ratio becomes small at low energy, a large relative uncertainty results. The z-axis represents the value of the correlation matrix elements.

Modeling Errors : Example from the Lattice

- From 2008—2014 HotQCD and Wuppertal-Budapest Collaborations spent 100s of millions of core hours calculating the QGP EoS for insertion into hydrodynamic code

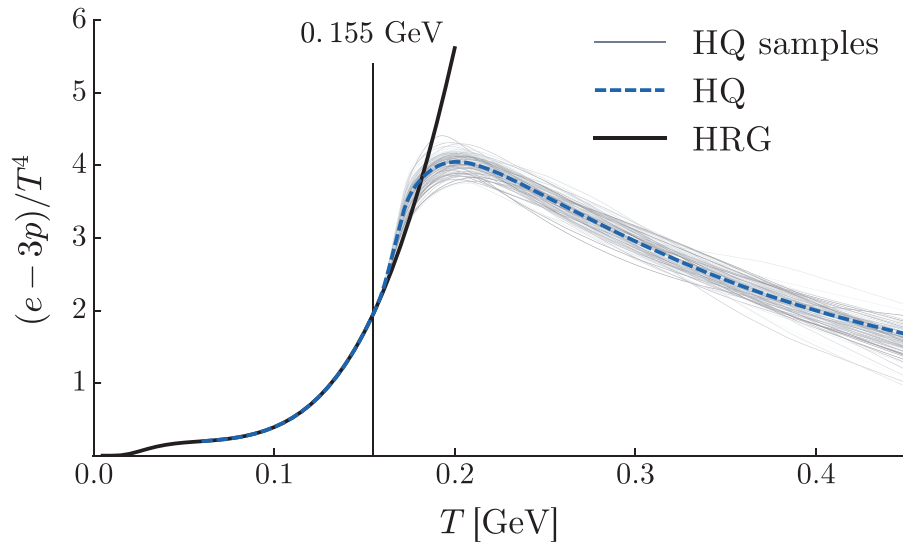


<https://dx.doi.org/10.1103/PhysRevD.90.094503>
<http://linkinghub.elsevier.com/retrieve/pii/S0370269314000197>

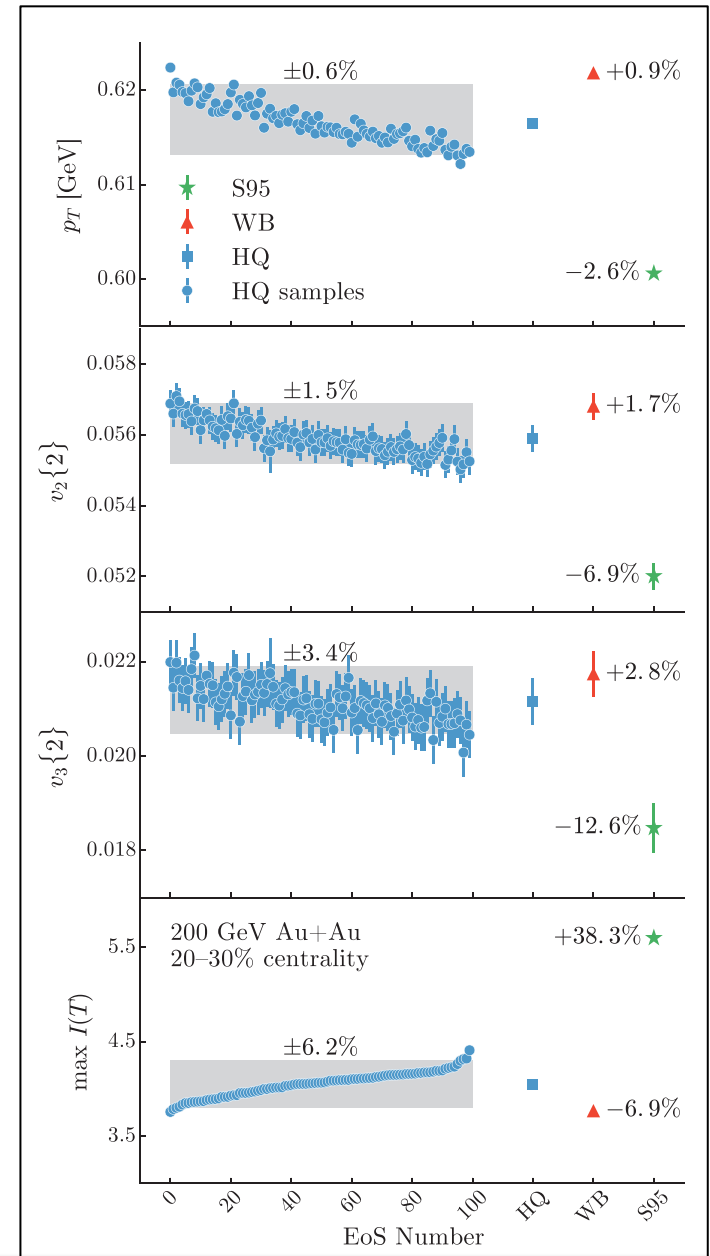
- In 2016, S. Moreland and RAS sought to answer
 - Does it matter which EoS result is used for hydro?
 - How much variation is within the systematic error bands?
 - Will we ever need to repeat these calculations on finer lattices?

Propagating Errors in the EoS

<http://link.aps.org/pdf/10.1103/PhysRevC.93.044913>



- HotQCD sys. errors calculated with simultaneous spline interpolation with continuum extrapolation. Errors impact observables by 3% or less.



And for Dinesh ...

- Bayesian determination of EoS,
- Pratt, Sangaline, Sorenson, Wang (2016)
<https://doi.org/10.1103/PhysRevLett.114.202301>

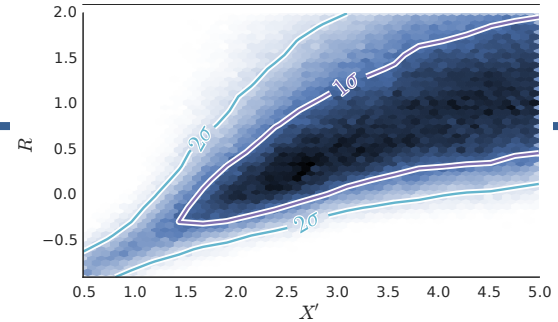
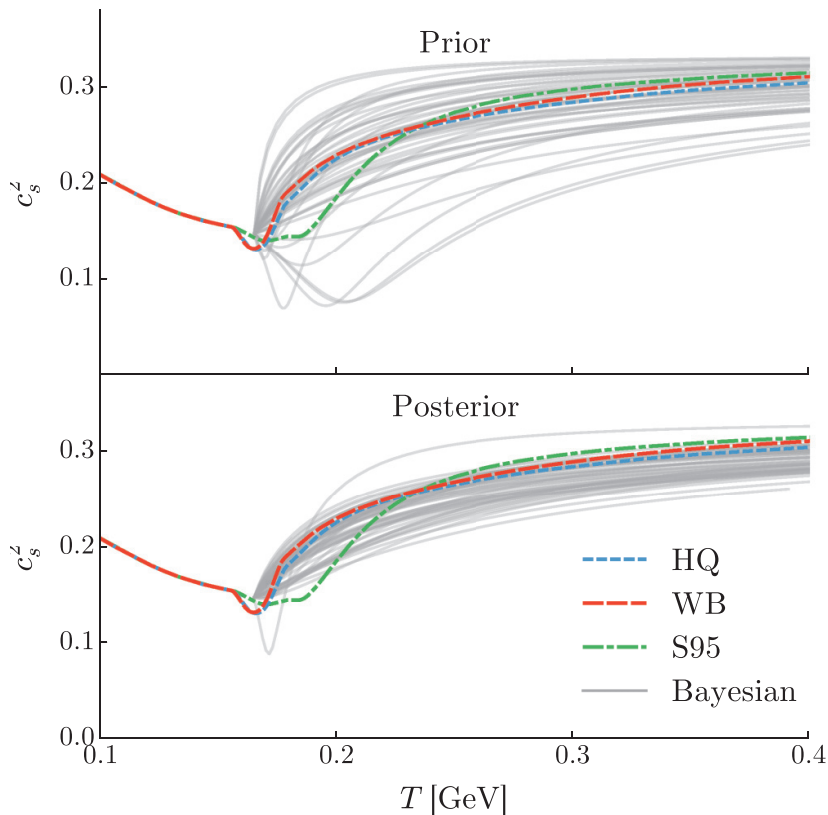


FIG. 4. The posterior likelihood for the two parameters that describe the equation of state, X' and R , have a preference to be along the diagonal. This shows that experiment constrains some integrated measure of the overall stiffness of the equation of state, i.e. a softer equation of state just above T_c is consistent with the data if it is combined with a more rapid stiffening at higher temperature.

$$c_s^2(\epsilon) = c_s^2(\epsilon_h) + \left(\frac{1}{3} - c_s^2(\epsilon_h) \right) \frac{X_0 x + x^2}{X_0 x + x^2 + X'^2}, \quad (2)$$

$$X_0 = X' R c_s(\epsilon) \sqrt{12}, \quad x \equiv \ln \epsilon / \epsilon_h,$$

where ϵ_h is the energy density corresponding to $T = 165$ MeV. The two parameters R and X' describe the behavior of the speed of sound at energy densities above ϵ_h . Whereas R describes how the speed of sound rises or falls for small x , X' describes how quickly the speed of sound eventually approaches $1/3$ at high temperature. Once given $c_s^2(\epsilon)$, thermodynamic relations provide all other representations of the equation of state. Runs were performed for $0.5 < X' < 5$, and with $-0.9 < R < 2$. In the limit $R \rightarrow -1$ the speed of sound will have a minimum of zero.

Conclusions

- Application of Bayesian methods to determine properties of QGP has progressed rapidly, treatment of errors has not
- Proper treatment of experimental errors is necessary if we are to compare to similar observables from different experiments
- Approach defined in [PhysRevC.77.064907](#) appears to be the one most followed
- Can we (must we) do better ?
- Assigning epistemic uncertainty to models is another challenge
- Future INT workshop for theory, experiment, and statistics ?