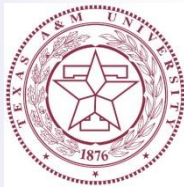


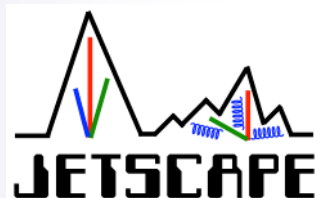
INT Jets+HQ 2017
Seattle, May 11, 2017

Jet Hadronization in Vacuum and in the Medium



Rainer Fries

Texas A&M University



Work in collaboration with

Kyongchol Han

Che Ming Ko

Michael Kordell

Overview

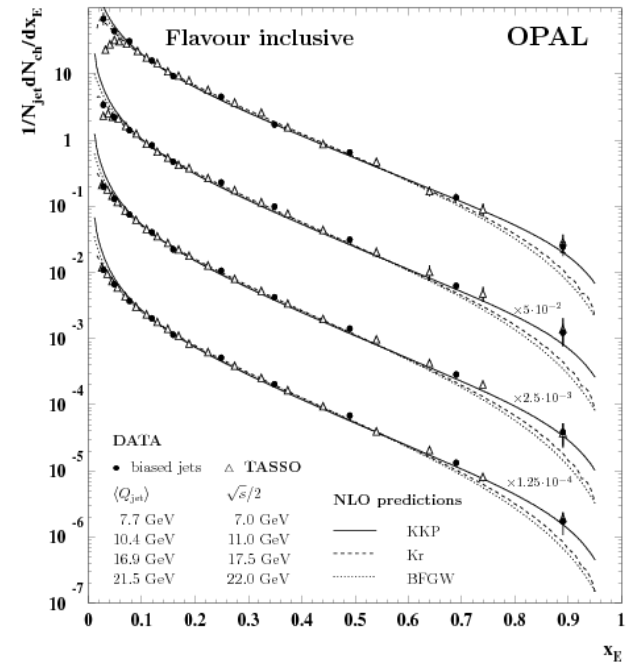
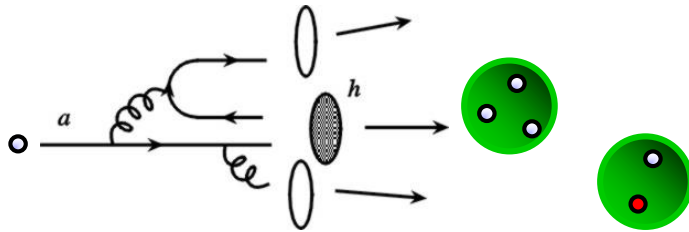
- Hadronization models.
- The JET recombination/fragmentation hybrid model
- Including at medium
- Appendix: Heavy quark recombination+fragmentation



Jet Hadronization

- Fragmentation of a parton from a hard process well-defined and well measured

$$D_{u \rightarrow h} \sim \langle 0 | u_\alpha | h \rangle \langle h | u_\beta | 0 \rangle$$



- Little theoretical and experimental guidance for jet hadronization: n quarks $\rightarrow m$ hadrons
 - Conservation laws and color
 - New generation of sub-jet observables?

Dirt Effects or Interesting Physics?

- Long list of new observables for jet substructure

$$F_N(\{p_i\}) = \sum_{i_1} \sum_{i_2} \dots \sum_{i_N} E_{i_1} E_{i_2} \dots E_{i_N} f_N(\hat{p}_{i_1}, \hat{p}_{i_2}, \dots, \hat{p}_{i_N})$$

All N-tuples N Energies Angular Weighting
(symmetric, vanishes for $\theta_{ij} \rightarrow 0$)

$$v e_n^{(\beta)} = \sum_{\text{all } n\text{-tuples}} (n \text{ energies}) (v \text{ smallest angles})^\beta$$

[Jesse Thaler slides]

.... with varying degree of sensitivity to hadronization.

[A.J. Larkoski, S. Marzani, G. Soyez, J. Thaler: arxiv:1402.2657]

- Hadronization is interesting because
Probe-medium interactions at $T_c = \lim(T \rightarrow T_c)$ probe-medium interactions at $T > T_c$
- The problem: theoretical control over the process.
- But: systematic comparison of p+p, p+A and A+A possible.

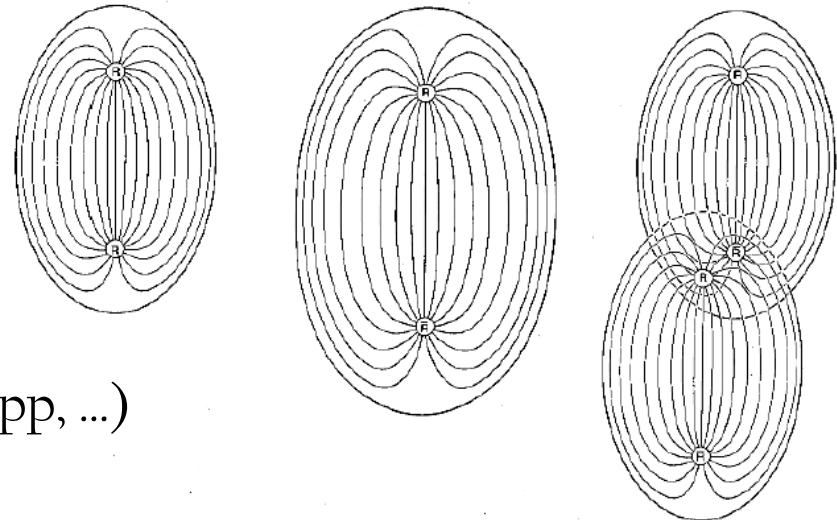


Jet Hadronization: Initial State

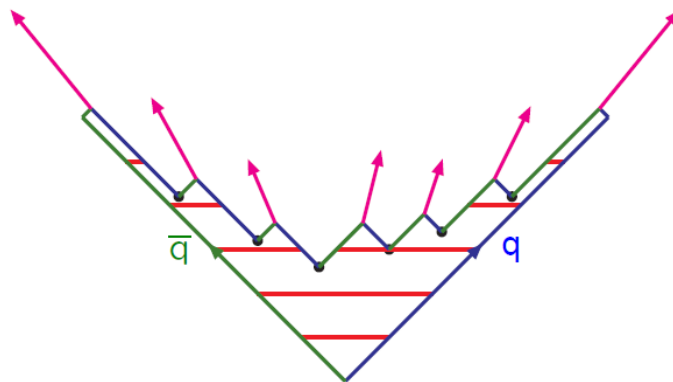
- Working definition: jet = (perturbative) parton shower + hadronization
- Traditional matching point:
 - virtuality for each parton: virtuality cutoff Q_0 .
 - complicated space-time structure
- With medium:
 - Need well-defined space-time picture
 - Parton virtuality and background temperature; might have to add phase space constraints.

String Fragmentation

- Hadronization based on string pictures of hadrons emerging in the early 70s

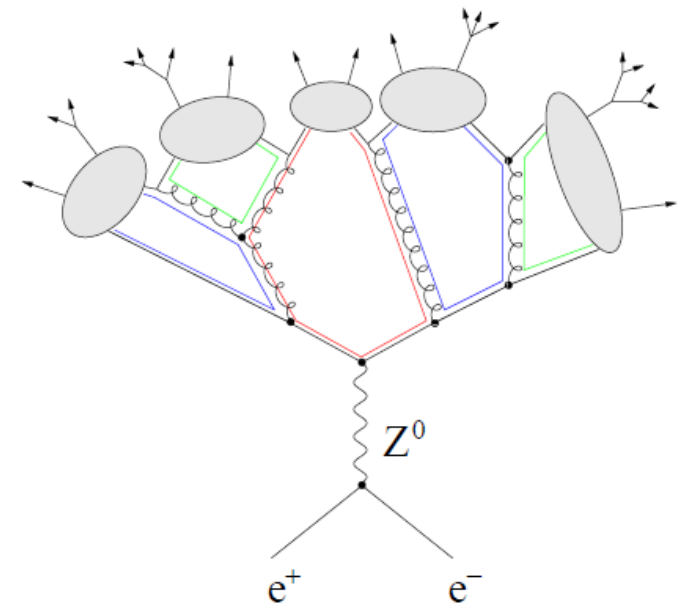


- Lund string model
- JETSET (e^+e^-) \rightarrow PYTHIA (e^+e^- , pp, ...)



Cluster Hadronization

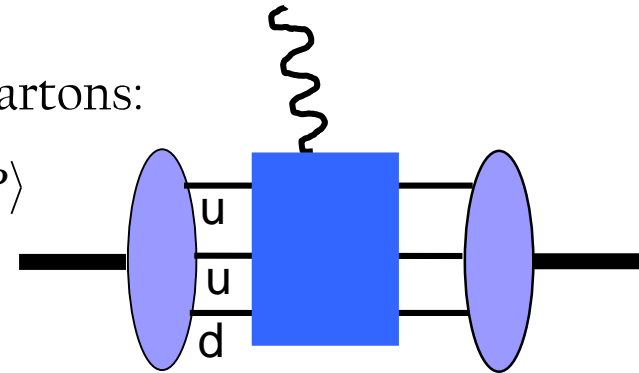
- Force gluon decays into quarks and antiquarks (non-perturbative)
- Local color neutrality: q - q bar form color-neutral clusters.
- Clusters decay into hadrons which can decay further into stable hadrons.
- Quite a few similarities if compared to recombination.



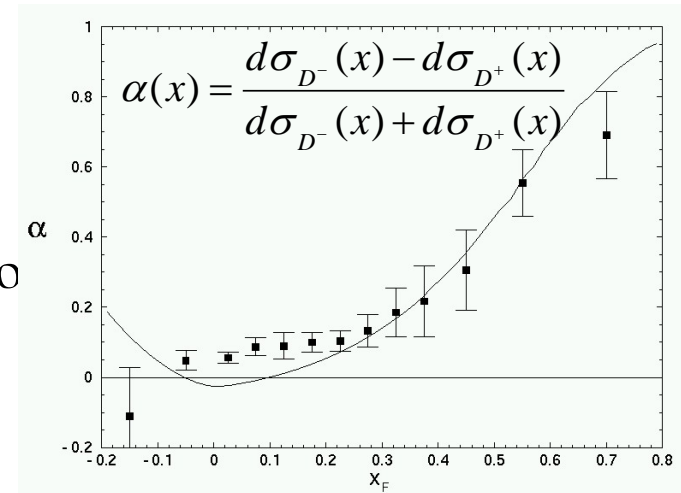
Quark Recombination

- Exclusive processes: recombination of all beam partons:

$$\psi \sim \langle 0 | u_\alpha u_\beta d_\gamma | P \rangle$$



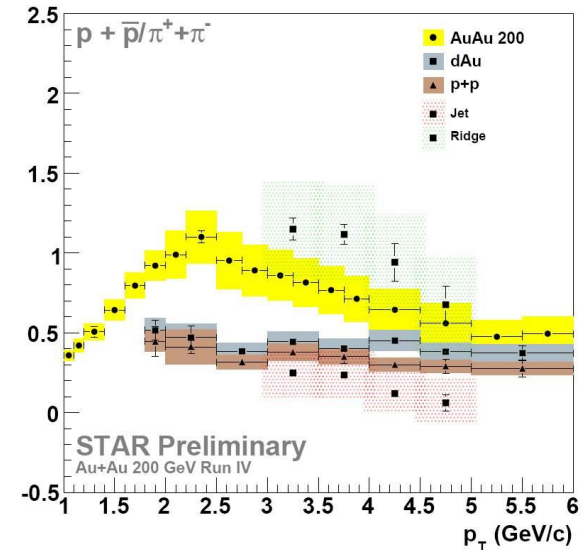
- Leading particle effect: recombination of produced partons with beam partons
- Charm-strange correlations in heavy ion collisions
strangeness enhancement seen in D_s .
- Quark scaling law in HI collisions



The JET Effort on Jet Hadronization

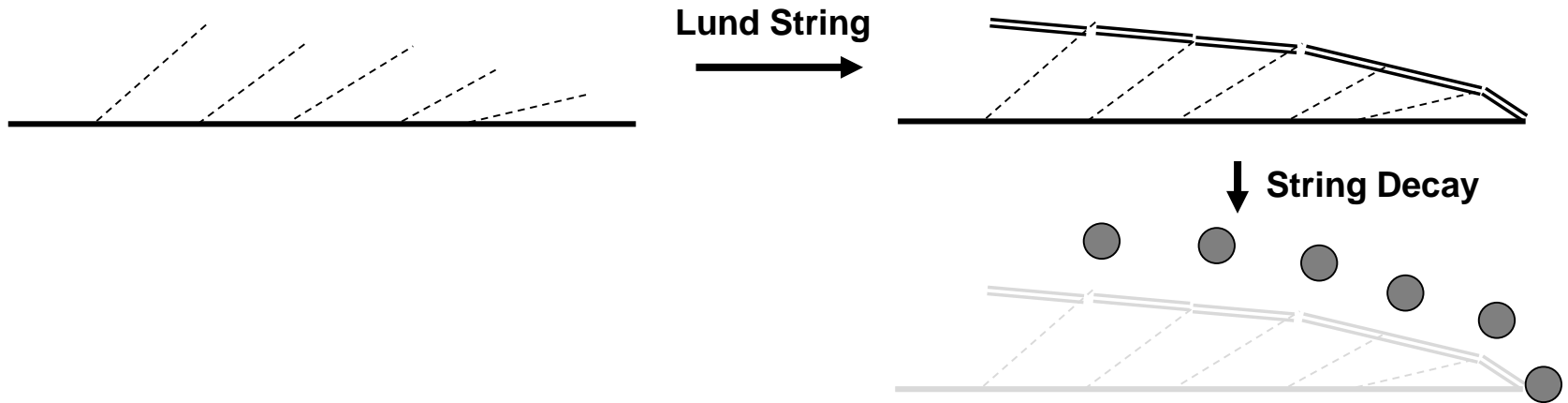
- JET collaboration goal: "... realistic calculations of jet production in a high energy density medium..."
 - Comprehensive modeling of jet showers, including hadron chemistry.
- Find a unified approach to modelling hadronization of jets that includes vacuum and in-medium cases
 - Hybrid model including quark recombination and string fragmentation.
 - Instantaneous recombination of constituent quark degrees of freedom.
 - Include space-time information

[K. Han, R.J.F., C. M. Ko, Phys. Rev. C 93, 045207 (2016)]



Vacuum Baseline

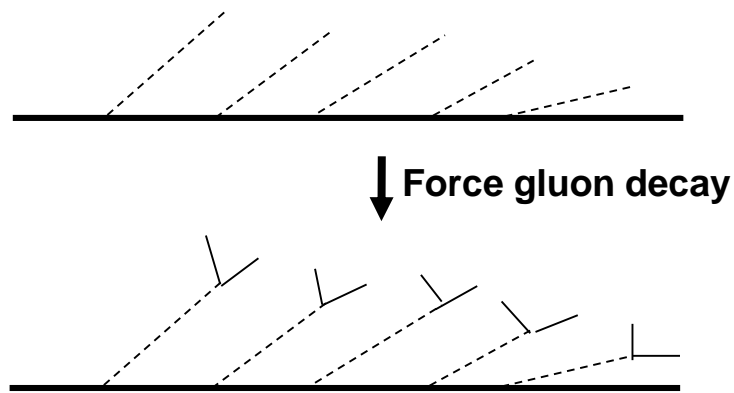
- Parton shower ceases at some virtuality Q_0 .
- Keep string fragmentation for the large- z part of a jet.



- Replace string fragmentation by recombination in the “dense” part of a parton shower.

Prepping Parton Showers

- Perturbative parton showers evolved to a scale \mathcal{Q}_0 .



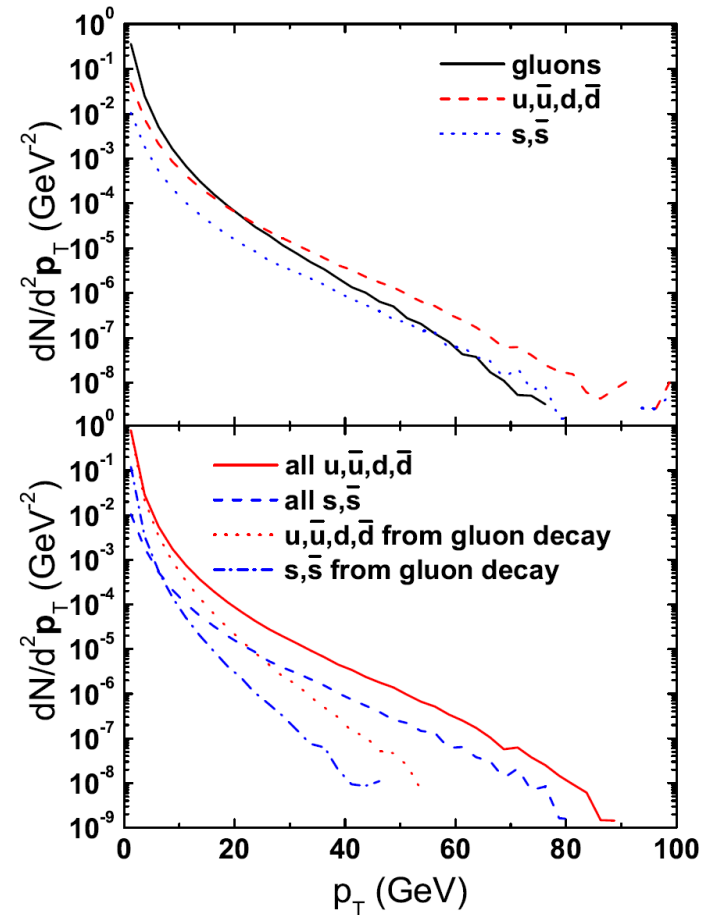
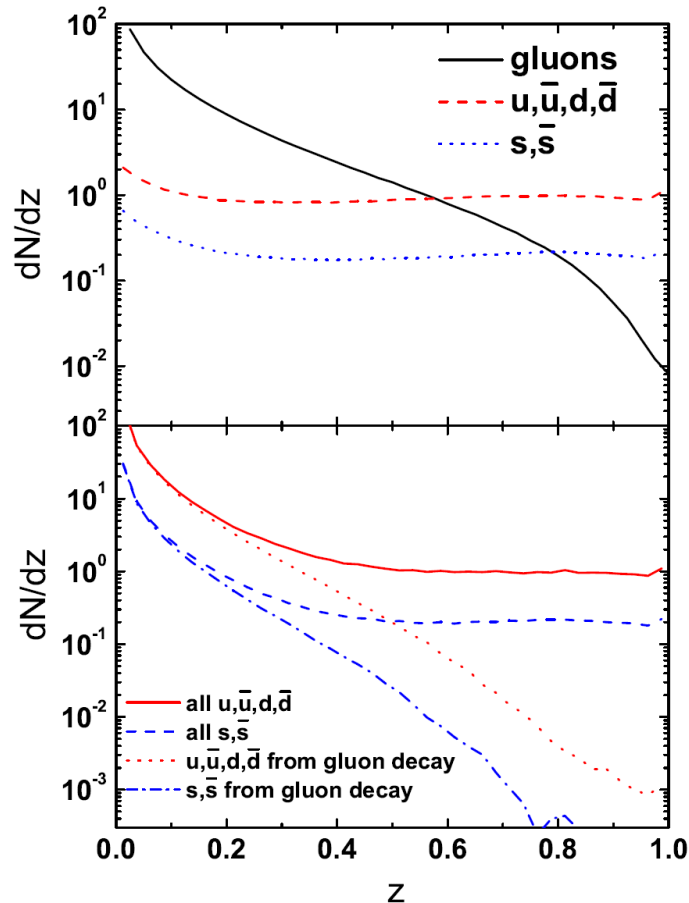
Examples here are PYTHIA showers supplemented by a space-time picture from tracking virtuality of partons which defines an average lifetime until splitting. Here: 100 GeV quark jets.

- Decay gluons with remaining virtualities into quark-antiquark pairs.
- Light and strange quarks, chemistry from phase space

$$\frac{\Gamma(g \rightarrow u\bar{u}, d\bar{d})}{\Gamma(g \rightarrow s\bar{s})} = 2\sqrt{\frac{m^2 - 4m_{u,d}^2}{m^2 - 4m_s^2}} \left(\frac{m^2 + 2m_{u,d}^2}{m^2 + 2m_s^2} \right)$$

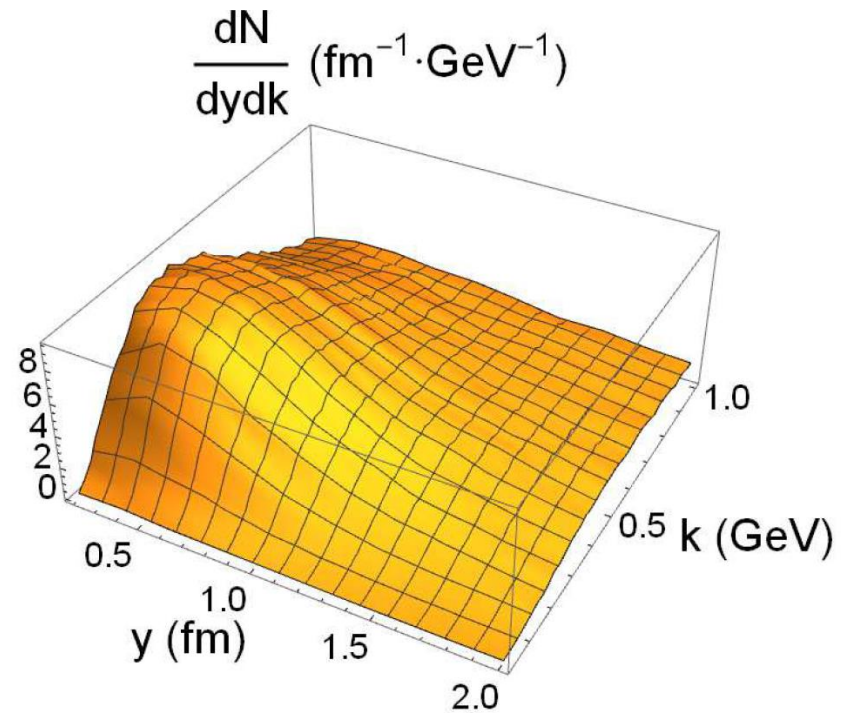
Prepping Parton Showers

- Example: Sample of 10^6 PYTHIA parton showers with $E_{\text{jet}} = 100$ GeV.
- dN/dz and dN/d^2P_T before vs after gluon decay



Prepping Parton Showers

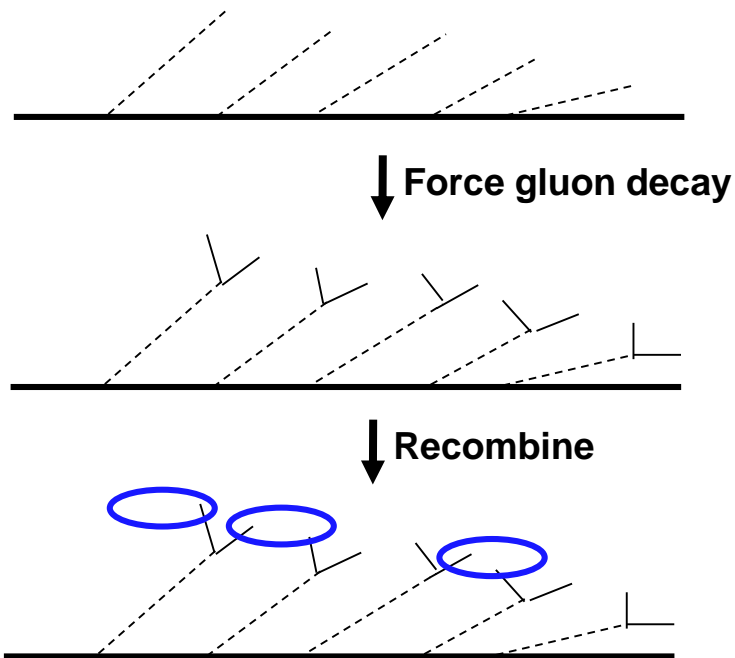
- Distance of quark-antiquark pairs in phase space is the deciding factor for the importance of recombination into mesons.
- Distribution of pair distances in (PYTHIA) parton showers in phase space (in their center of mass frame)
- “Our” PYTHIA jets: most of the jet is relatively “dense” in phase space.
- Long tails exist (\sim large z partons)
- Fundamental test for other jet Monte Carlos: perturbative evolution should not lead to dilute showers, otherwise non-perturbative effects are already dominant.



Recombination Step

- Monte Carlo version of the instantaneous recombination model by Greco, Ko and Levai based on Wigner function formulation

[Greco, Ko & Levai,
PRL 90, 202302 (2003)]



- Evaluate recombination probability for q - q bar pairs and q (q bar) triplets.
- Roll dice to decide if recombination happens.
- Recombine into stable hadrons and resonances, let resonances decay (connections to cluster hadronization).
- Color treated statistically.

[K. Han, R.J.F., C. M. Ko, Phys. Rev. C 93, 045207 (2016)]

Recombination Step

- Well-known Wigner function coalescence formulas:

$$\frac{dN_M}{d^3\mathbf{p}_M} = g_M \int d^3\mathbf{x}_1 d^3\mathbf{p}_1 d^3\mathbf{x}_2 d^3\mathbf{p}_2 f_q(\mathbf{x}_1, \mathbf{p}_1) f_{\bar{q}}(\mathbf{x}_2, \mathbf{p}_2) \times W_M(\mathbf{y}_1, \mathbf{k}_1) \delta^{(3)}(\mathbf{P}_M - \mathbf{p}_1 - \mathbf{p}_2), \quad (3)$$

[RJF, V. Greco, P. Sorensen, Ann. Rev. Nucl. Part. Sci. 58, 177 (2008)]

$$\frac{dN_B}{d^3\mathbf{p}_B} = g_B \int d^3\mathbf{x}_1 d^3\mathbf{p}_1 d^3\mathbf{x}_2 d^3\mathbf{p}_2 d^3\mathbf{x}_3 d^3\mathbf{p}_3 f_{q_1}(\mathbf{x}_1, \mathbf{p}_1) \times f_{q_2}(\mathbf{x}_2, \mathbf{p}_2) f_{q_3}(\mathbf{x}_3, \mathbf{p}_3) W_B(\mathbf{y}_1, \mathbf{k}_1; \mathbf{y}_2, \mathbf{k}_2) \times \delta^{(3)}(\mathbf{P}_B - \mathbf{p}_1 - \mathbf{p}_2 - \mathbf{p}_3), \quad (4)$$

- Can be turned into a formula for recombination probability (here meson)

$$\overline{W}_M(\mathbf{y}, \mathbf{k}) = \int d^3\mathbf{x}'_1 d^3\mathbf{k}'_1 d^3\mathbf{x}'_2 d^3\mathbf{k}'_2 \times W_q(\mathbf{x}'_1, \mathbf{k}'_1) W_{\bar{q}}(\mathbf{x}'_2, \mathbf{k}'_2) W_M(\mathbf{y}', \mathbf{k}').$$

- Evaluated at equal time in the pair or triplet rest frame.
- Throw dice to accept or reject a pair or triplet for recombination.



Recombination Step

- Bound state Wigner function derived from harmonic oscillator wave functions (L_n =Laguerre polynomials).

$$W_n(u) = 2(-1)^n L_n \left(\frac{4u}{\hbar\omega} \right) e^{-2u/\hbar\omega} \quad u = \frac{\hbar\omega}{2} \left(\frac{x^2}{\sigma^2} + \sigma^2 k^2 \right)$$

- For the probabilities to be positive definite, need proper q, \bar{q} Wigner functions: Introduce Husimi smearing by representing quarks by proper wave packets of width δ .
- For $\sigma^2 = 2\delta^2$ the result for the overlap of wave packets and Wigner function is extremely simple. The probability densities for the n -th excited states are

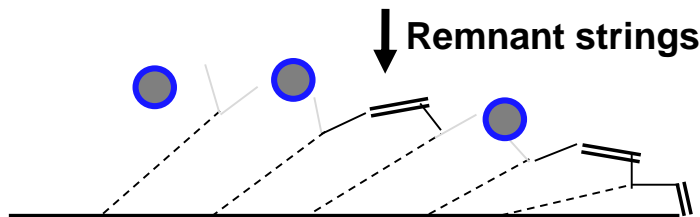
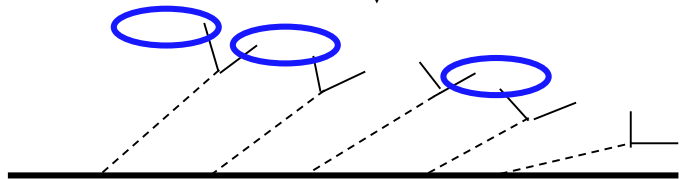
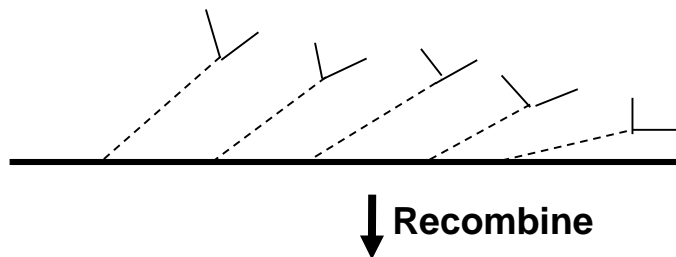
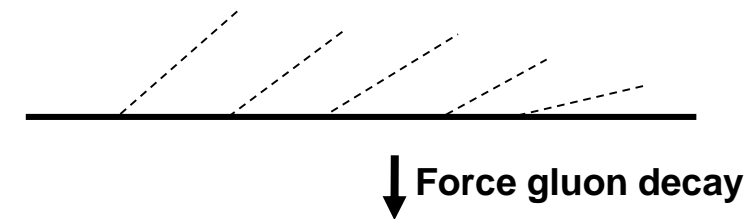
$$\bar{W}_{M,n}(\mathbf{y}, \mathbf{k}) = \frac{v^n}{n!} e^{-v} \quad v = \frac{1}{2} \left(\frac{\mathbf{y}^2}{\sigma_M^2} + \mathbf{k}^2 \sigma_M^2 \right)$$

- The true shape and size of the input wave packets are not known.
- Hadron wave function widths fixed by measured charge radii.

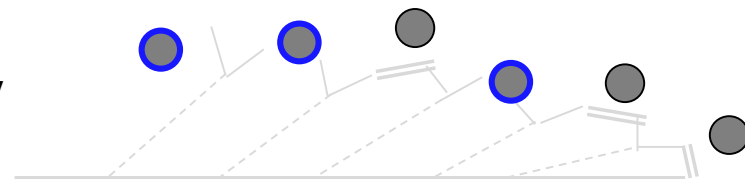


Remnant Strings

- Naturally there are remnant quarks and antiquarks which have not found a recombination partner.



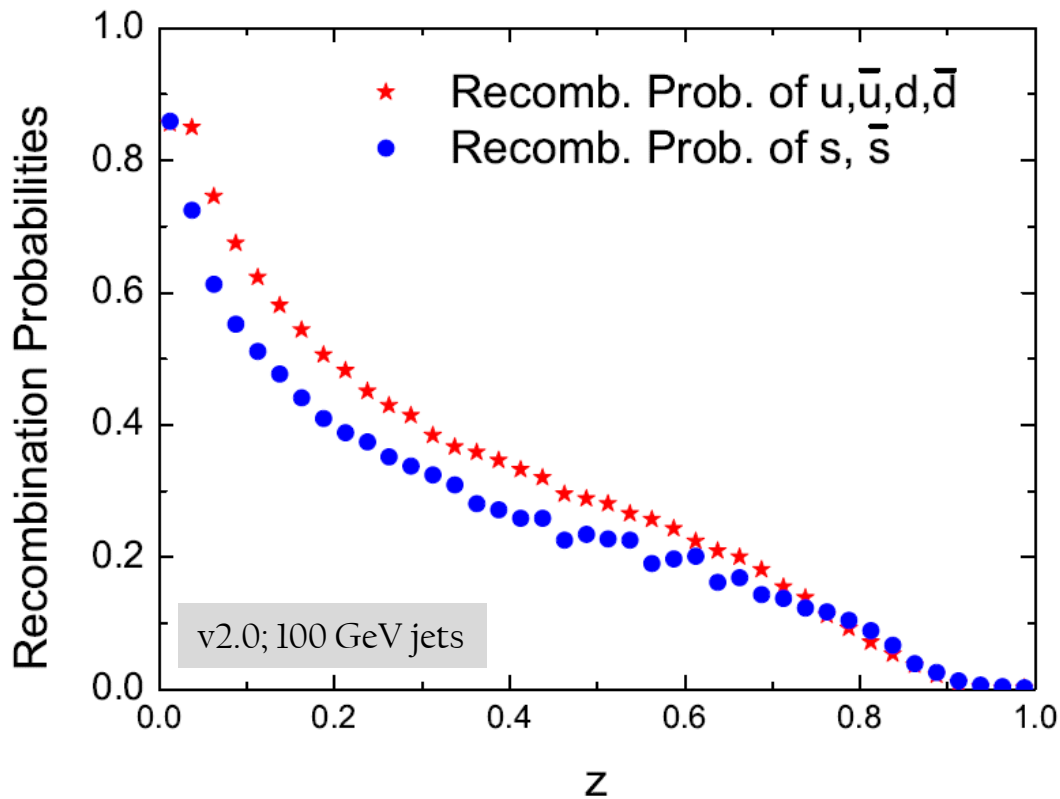
String Decay



- Why? No confinement in parton shower, quarks can get far away.
- In reality: colored object needs to stay connected.
- Return these partons to PYTHIA to connect them with remnant strings.

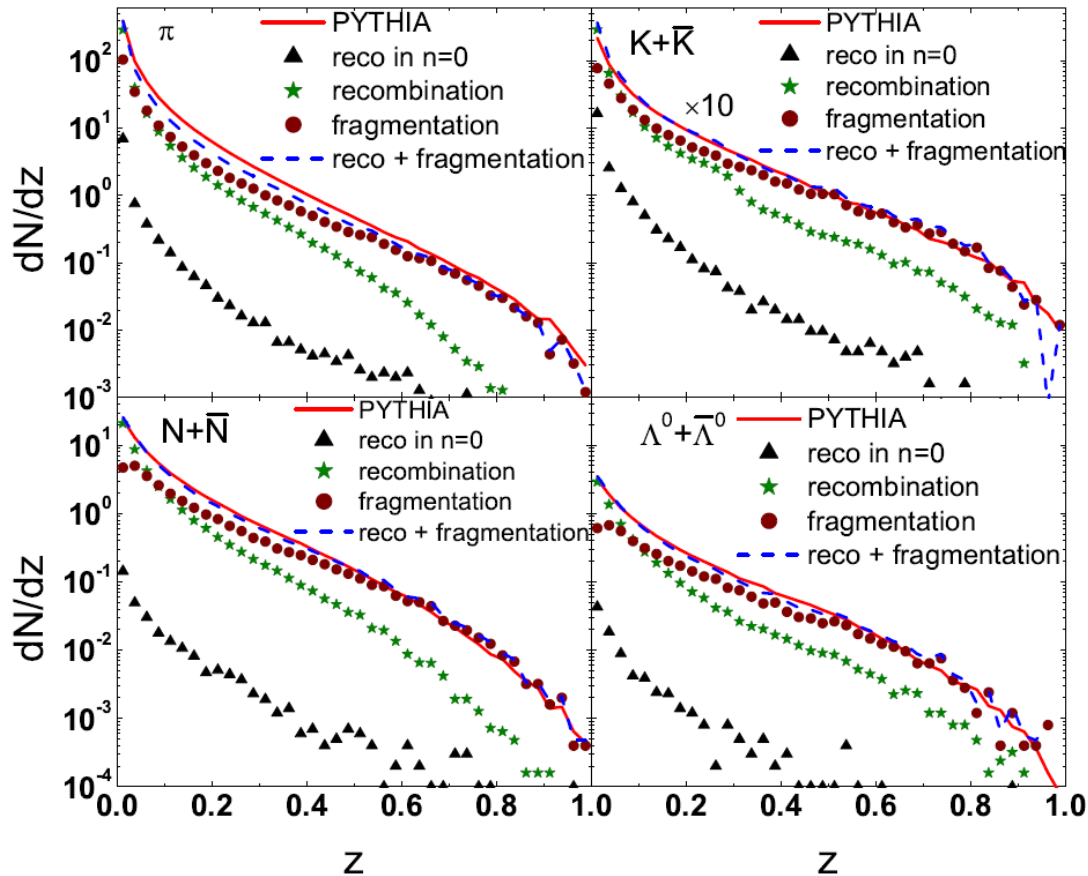
Remnant Strings

- Check on recombination probability (100 GeV PYTHIA vacuum jets)



Results (Vacuum)

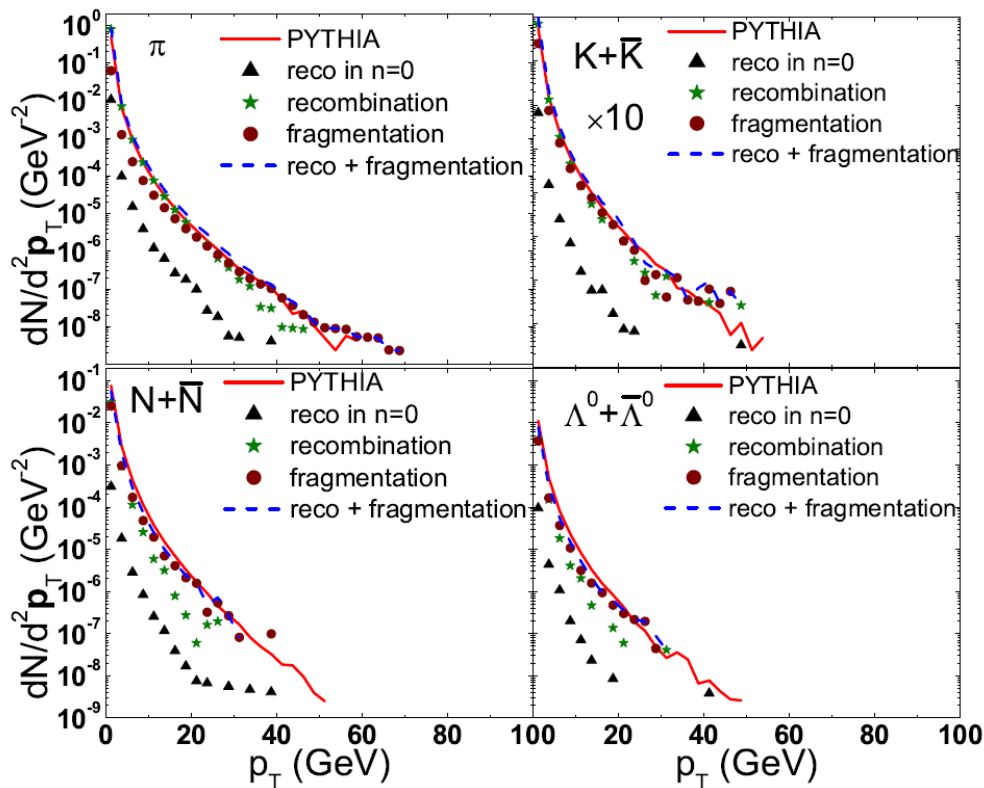
- Longitudinal structure: dN/dz of stable particles compared to PYTHIA string fragmentation (e^+e^-).



v2.0; 100 GeV jets

Results (Vacuum)

- Transverse structure: dN/d^2p_T of stable particles compared to PYTHIA string fragmentation ($e+e^-$).

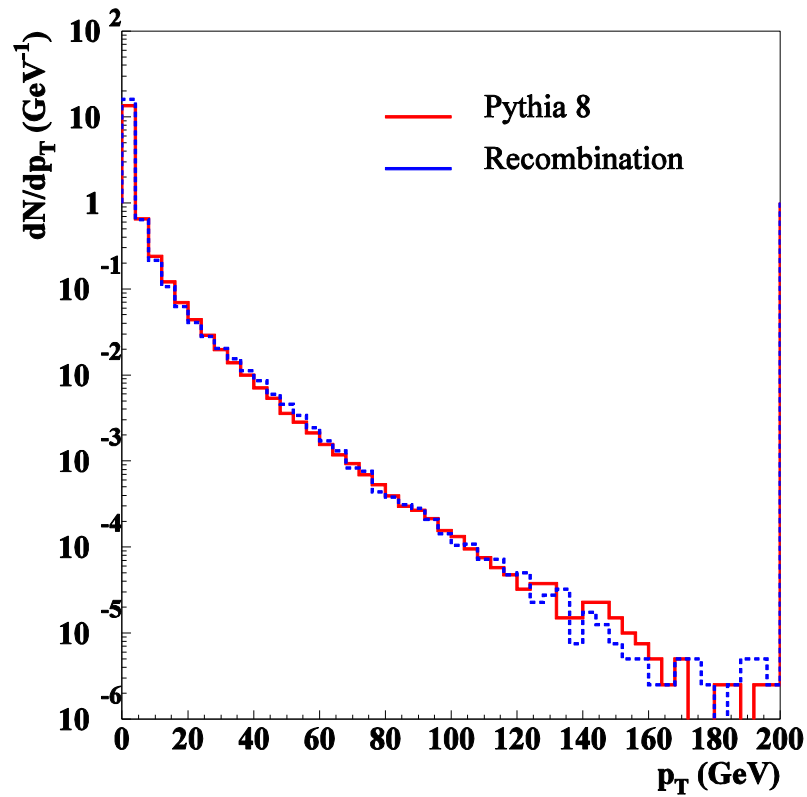


- Generally good agreement with pure string fragmentation.
- No precision tuning to data.

v2.0; 100 GeV jets

Results (Vacuum)

- More complicated systems: dN/dp_T of hadrons in p+p.

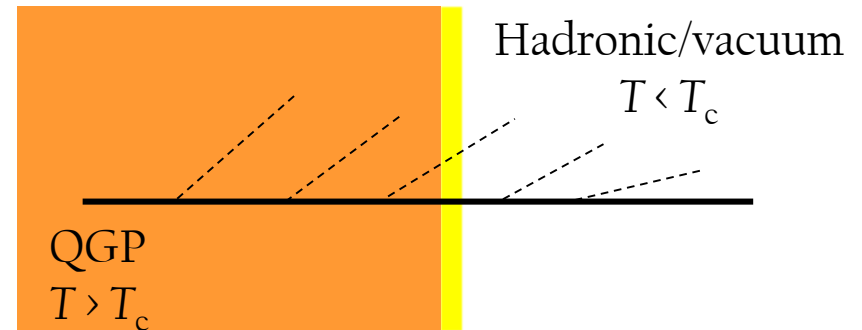


v1.0

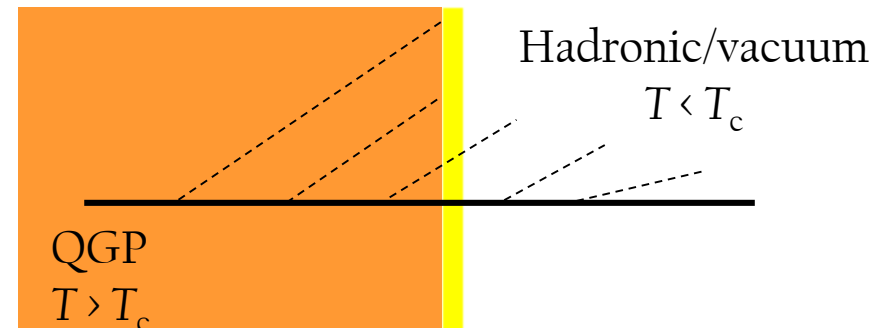
Credit: Tan Luo

Adding a Medium

- The space-time picture is complicated and it matters!
- All relevant partons have to be on the surface of the QGP or outside the QGP to hadronize.
- Propagate all shower partons to the hadronization hypersurface, or make them part of the medium.



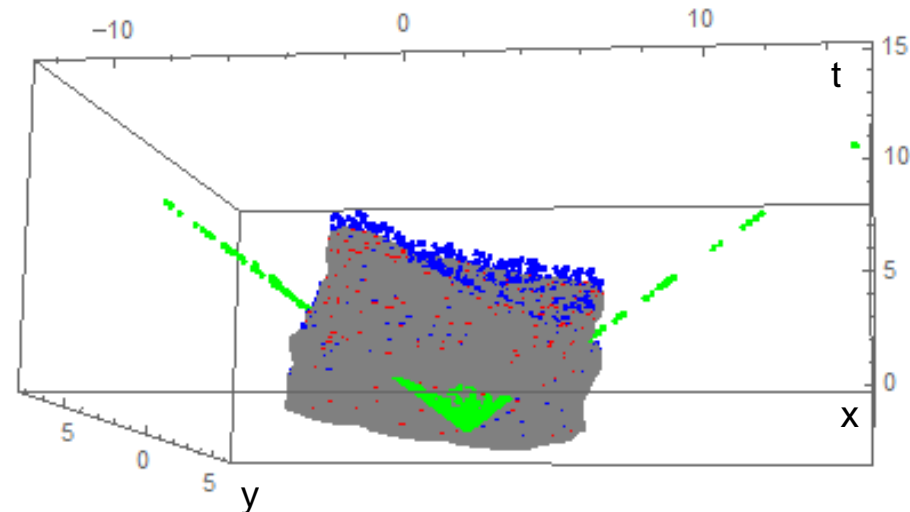
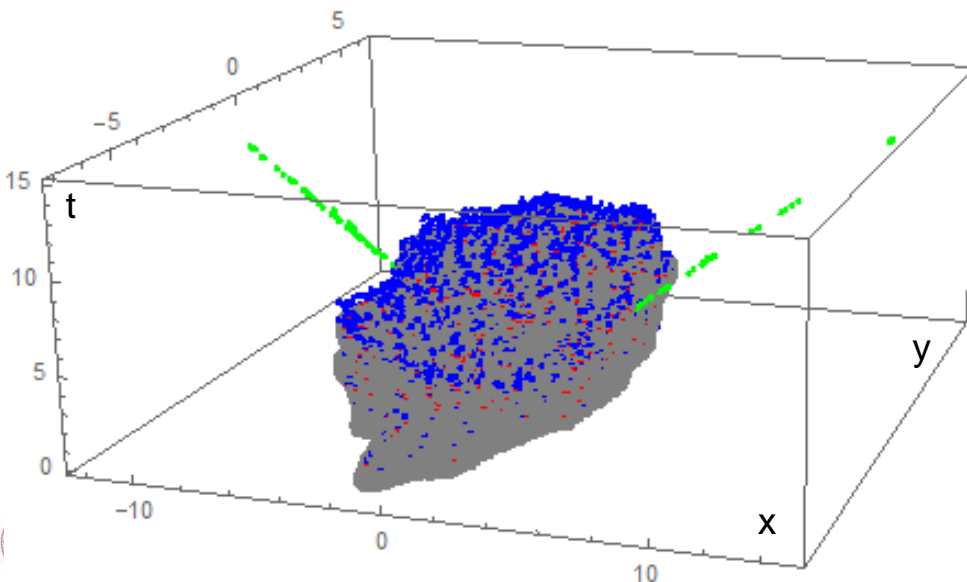
Hadronization
 $T \sim T_c$



Hadronization
 $T \sim T_c$

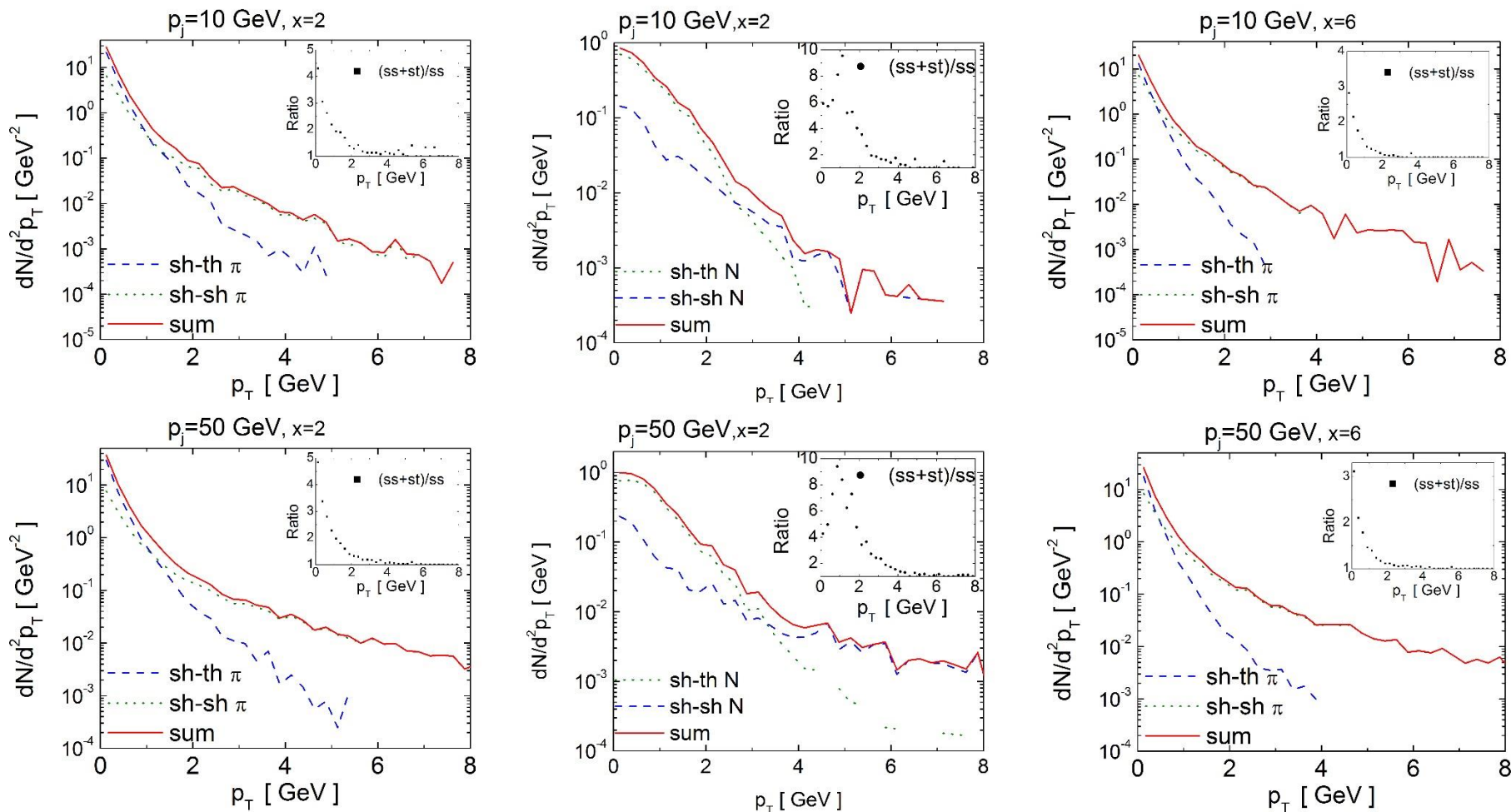
Adding a Medium

- iEBE (Ohio State) event-by-event hydro with sampled thermal partons on the $T=T_c$ hypersurface.
- As input for recombination we need to sample $f(x, p)$, not $f(x, p)p \cdot \sigma/E$
- Plots: 500 PYTHIA (vacuum!) showers emerging from the center embedded into an iEBE event
 - blue = sampled thermal partons; green = shower; grey = hypersurface



Early Results With Medium

- PYTHIA showers propagated to the $T=T_c$ hypersurface ($E < 600$ MeV quarks assumed lost); sh-sh includes remnant fragmentation.



Some Results With Medium

- Medium effects can be seen for small momenta < 10 GeV.
- Shower-thermal recombination takes over from pure shower-shower recombination; remnant strings are still important.
- Importance of shower-thermal vs shower-shower: $N > K > \pi$. Large for nucleons.
- Caveat: these are not proper in-medium showers.



Current Work

- TReC Code versions

v1.0: quark jets in vacuum; internal only

v2.0: quark jets in vacuum; H.o. Wigner functions; described in PRC 93, 045207 (2016)

v2.1: basic in-medium functionality added

v2.1B: Berkeley/CCNU custom version

See talk by Tan Luo

v2.1T: Texas A&M updates and bugfixes

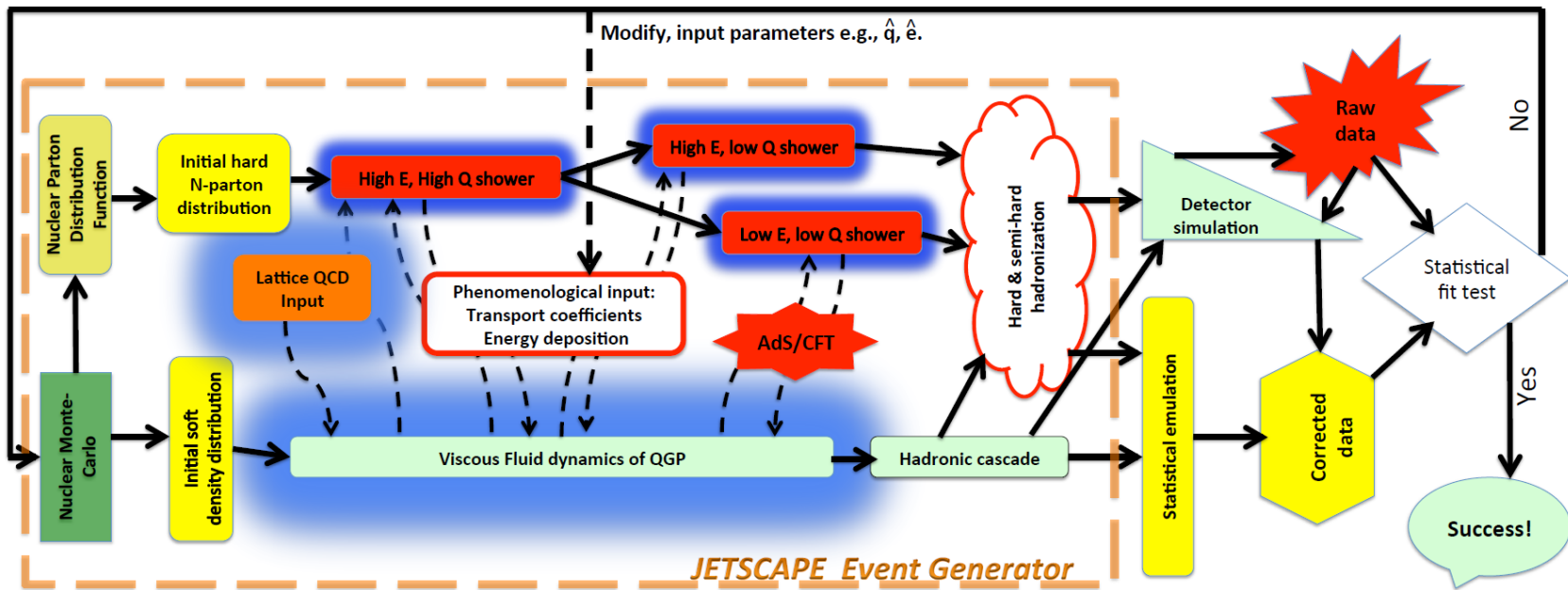
v2.2: new, improved release

v3.0: future JETSCAPE version



Current Work

- The JETSCAPE framework



Summary

- We have developed an event-by-event hybrid hadronization module for jet Monte Carlos.
- Quark recombination including resonances, supplemented by string fragmentation.
- Medium effects by sampling hydro event-by-event.
- v2.2 will be available soon.
- Systematic studies of jet medium effects with in-medium jet MCs are ongoing.



Appendix: Heavy Quark Hadronization

Work in collaboration with
Min He
Ralf Rapp

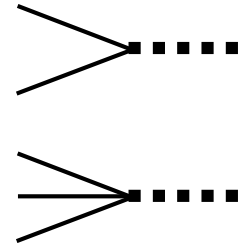


Resonance Recombination

- Early recombination models:

- Instantaneous projection of states
- $2 \rightarrow 1$ & $3 \rightarrow 1$ processes: no energy conservation

$$N_M = \int \frac{d^3 P}{(2\pi)^3} \langle M; P | \rho | M; P \rangle$$



- Resonance recombination:

- Mesons appear as resonances of quark-antiquark scattering
- Described by Boltzmann equation, start with ensemble of quarks/antiquarks

$$\frac{\partial}{\partial t} f_M(t, \vec{p}) = -\frac{\Gamma}{\gamma_p} f_M(t, \vec{p}) + g(\vec{p})$$



[Ravagli & Rapp PLB 655 (2007)]
 [Ravagli, van Hees & Rapp, PRC 79 (2009)]

- Breit-Wigner resonance cross sections:

$$\sigma(s) = C_M \frac{4\pi}{k^2} \frac{(\Gamma m)^2}{(s - m^2)^2 + (\Gamma m)^2}$$

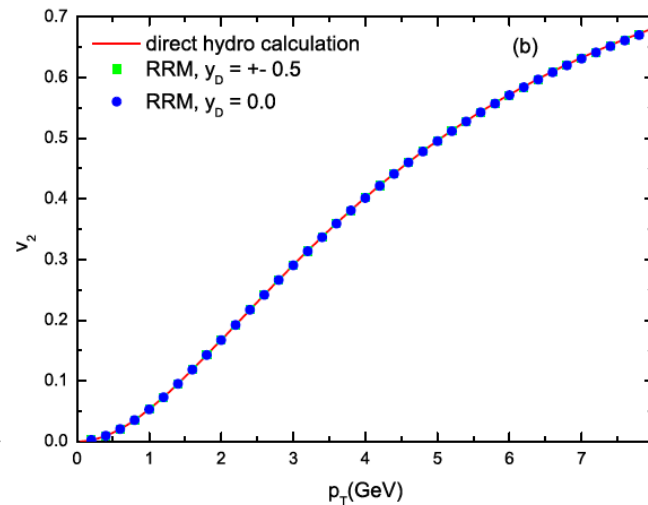
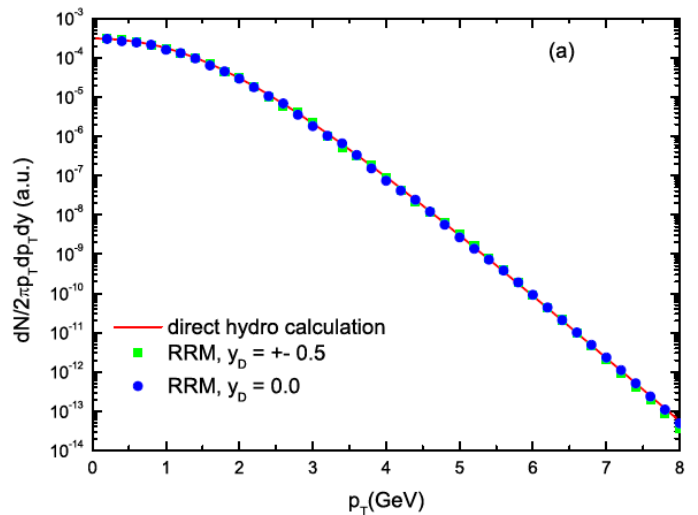
- Long-time limit:
$$E \frac{dN_M}{d^3 P} = \frac{E\gamma}{8(2\pi)^3 \Gamma} \int \frac{d^3 x d^3 p_{rel}}{(2\pi)^3} f_a(x, p_1) f_a(x, p_2) \sigma(s) v_{rel}(P, p_{rel})$$

- Conserves energy and momentum, should be able to attain equilibrium.
- Compatible with the picture of a strongly interacting medium.



Test of Equilibrium Limit

- Energy conservation + detailed balance + equilibrated quark input \rightarrow equilibrated hadrons!
- Can be applied to any hadronization hypersurface Σ .
- Numerical tests: compare blast wave hadrons at $T_c - \varepsilon$ to hadrons coalesced from quarks of the same blast wave at $T_c + \varepsilon$: heavy-light system at AZHYDRO hypersurface



[M. He, RJF and R. Rapp, PRC 82 (2010)]

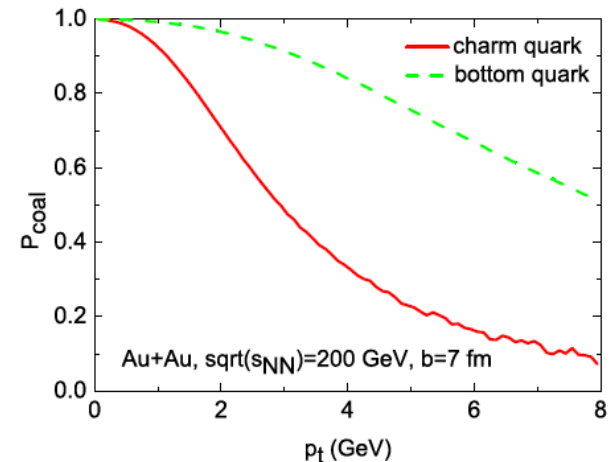
Implementation for Heavy Quarks

- Resonance recombination ideal for systems close to equilibrium.
- How to decide recombination vs fragmentation rate?

Q-q recombination rate \sim Q-q in medium scattering rate!

$$P_{\text{coal}}(p) = \Delta\tau_{\text{res}} \Gamma_Q^{\text{res}}(p)$$

- Apply fragmentation with probability $1 - P_{\text{coal}}(p)$.
- Hence consistent with in-medium dynamics.
 - Low momenta = recombination dominated (co-moving thermal partons!)
 - High momenta = fragmentation dominated (no co-moving thermal partons)
- Usually two extreme assumptions about $\Delta\tau$:
 - Corresponding to $P_{\text{coal}} = 1$ or $1 - e^{-1}$ at $p=0$.
- Total recombination probability averaged over fluid cells in lab frame:



The D_s as a Signature

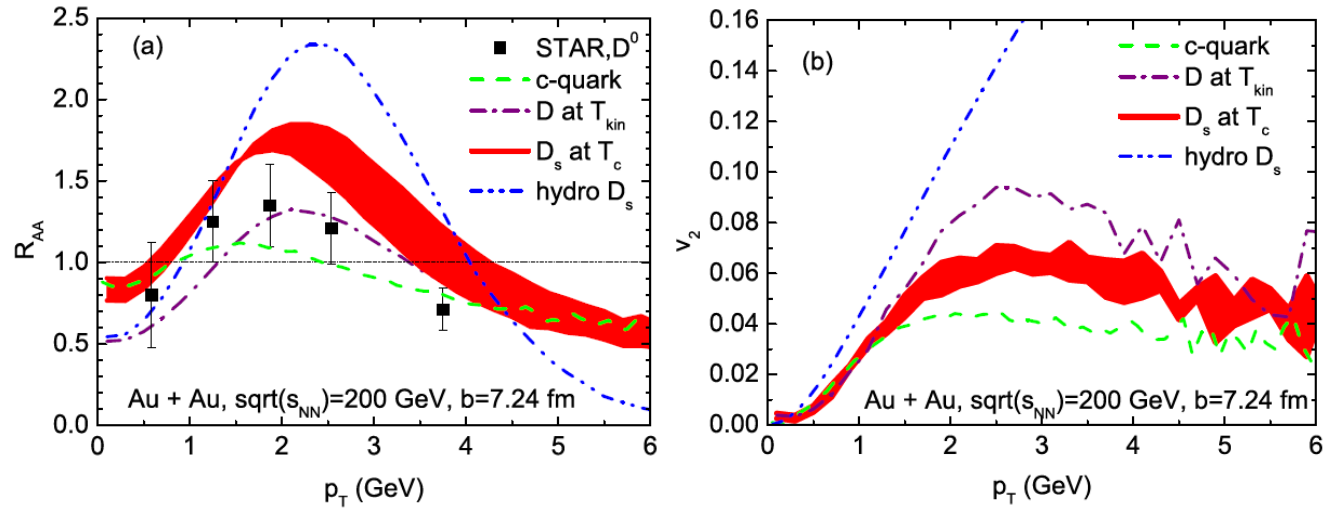
- D_s = charm-strange bound state.
- Signature 1: D_s vs $D R_{AA}$ is a measure for strength of recombination vs fragmentation.
 - Charm in D_s and D suffer from same drag and diffusion up to T_c .
 - If charm fragments: D_s/D as in p+p.
 - If charm recombines: D_s picks up enhanced strangeness $\rightarrow D_s$ enhanced.
 - Numerical check: hadronic phase does not destroy this signal.
- Signature 2: D_s vs $D v_2$ can measure the relative strength of D_s vs D interactions in the hadronic phase.
 - D_s is an analogue to multi-strange hadrons in the light sector.
 - If there is early freeze-out it can be read off from the D_s vs $D v_2$.

[M. He, RJF and R. Rapp, Phys. Rev. Lett 110, 112301 (2013)]



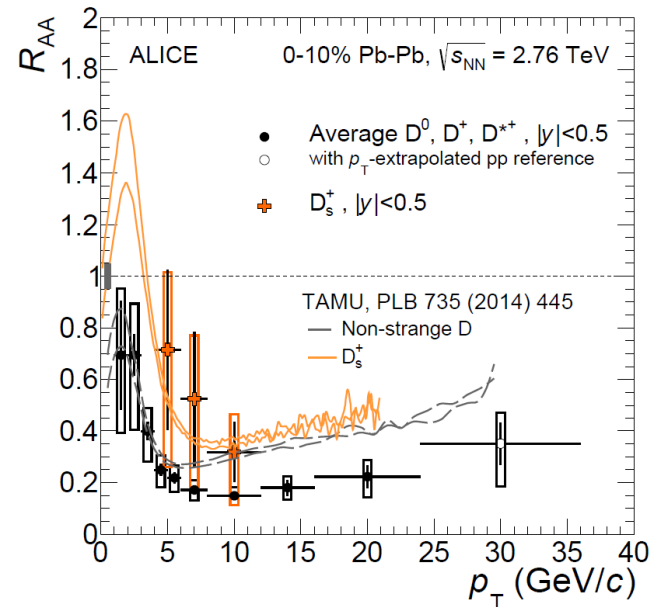
Results

- RHIC:



[M. He, RJF and R. Rapp, Phys. Rev. Lett 110, 112301 (2013)]

- D_s enhancement seen but not yet statistically significant



Backup



Recombination Step

- Parameters (harmonic oscillator WF case)

TABLE I. Table of measured charge radii R (from Ref. [21]), widths σ_M (and σ_B), and statistical factor g for all hadrons used in this calculation.

Hadron	R [fm]	σ_M (and σ_B) [fm]
π	0.67	1.09
ρ	–	1.09
K	0.56	1.10
K^*	–	1.10
N	0.88	1.24
Δ	–	1.24
Λ	–	1.15

Towards a Full Picture

- Some older results calculated using v1.0 and a blast wave medium reproduces experimental p/ π ratio (jet, jet-medium and thermal medium itself included)

