



INT Program INT-17-1b

Precision spectroscopy of QGP  
properties with jets and heavy quarks

May 31<sup>st</sup> 2017

# Quarkonium production in proton-nucleus collisions

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properties with jets and heavy quarks

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## Quarkonium production in

- p-p (at LHC)
- p-A collisions (at both RHIC and LHC)

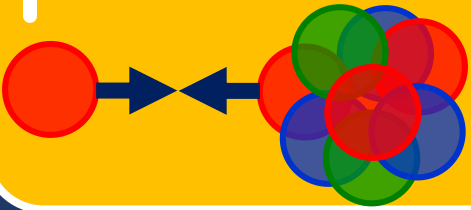
# quarkonium in pp, pA, AA

p-p



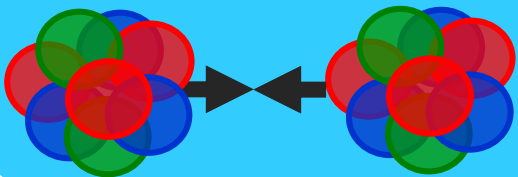
“vacuum” reference for AA, pA,  
genuine pp physics program

p-A



cold nuclear matter effects:  
shadowing/CGC, energy loss...

A-A



hot matter effects:  
regeneration vs suppression

Nuclear modification factor:

$$R_{AA}^{J/\psi} = \frac{Y_{AA}^{J/\psi}}{\langle T_{AA} \rangle \sigma_{pp}^{J/\psi}}$$

Medium effects quantified  
comparing AA (pA) quarkonium  
yield with the pp cross section,  
scaled by a geometrical factor  
(from Glauber model)

- no medium effects  $\rightarrow R_{AA} = 1$
- hot/cold matter effects  $\rightarrow R_{AA} \neq 1$

p-p collisions



# J/ψ production in pp collisions at LHC

➔ Focussing just on LHC results...RHIC covered in Zebo's talk!

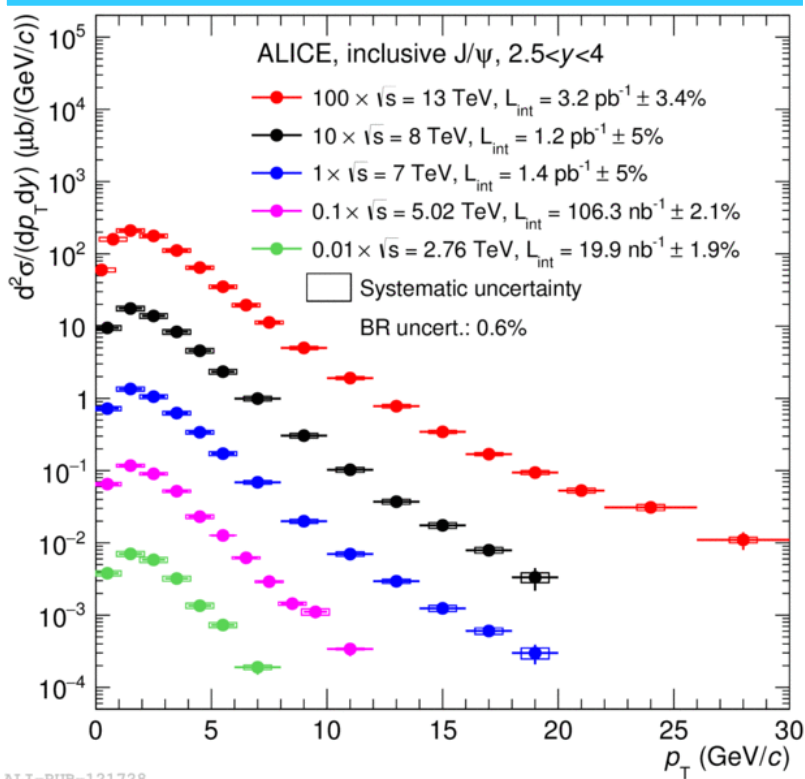
➔ Quarkonium production now measured at LHC in many collision systems

$\sqrt{s}$ (TeV)	2.76	5	7	8	13
pp	X	X	X	X	X

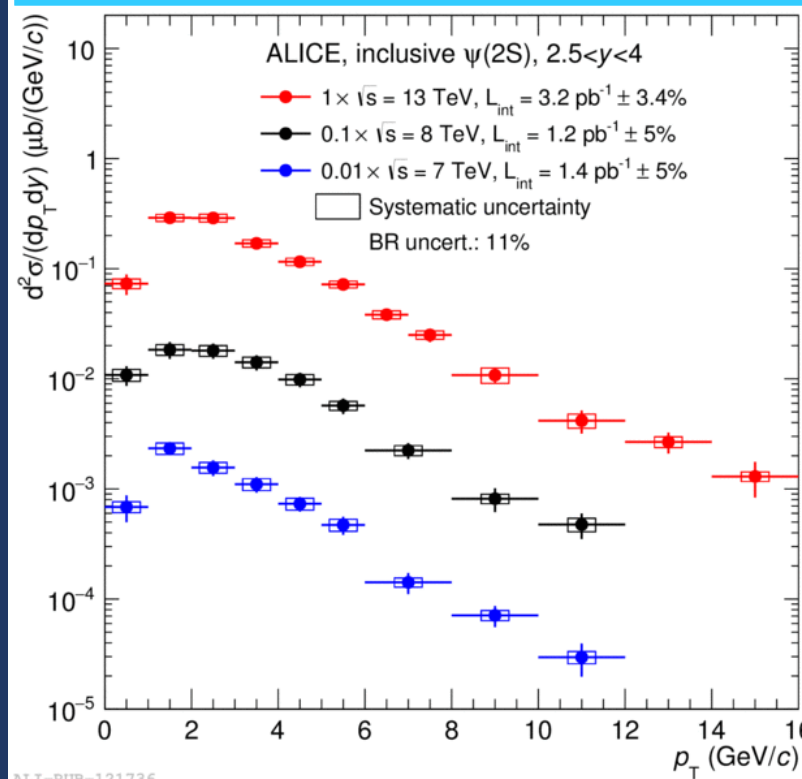
not exhaustive  
➔ selection of results on:

- $\gamma, p_T$  differential cross sections
- Double ratios
- Non prompt J/ψ fraction
- Self-normalized yields vs event multiplicity

J/ψ

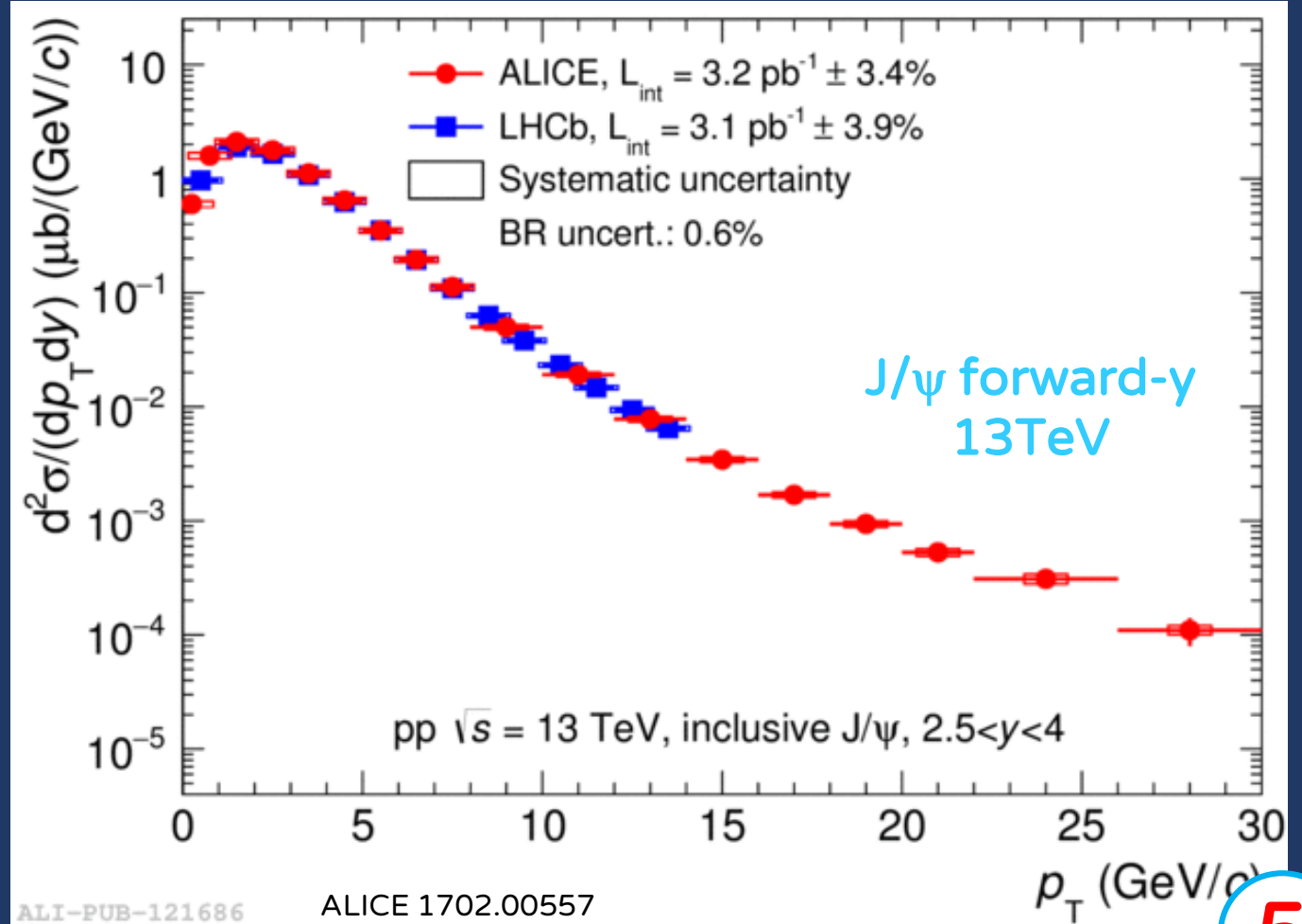
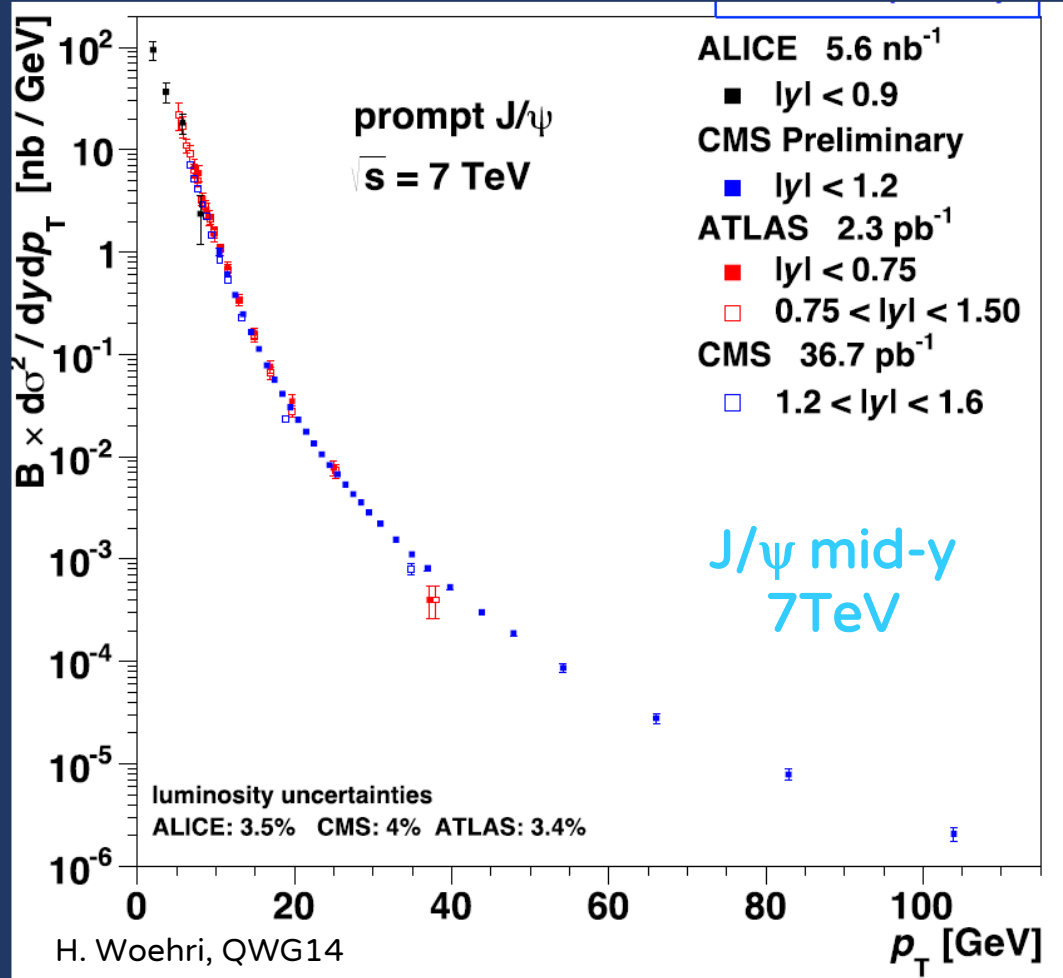


ψ(2S)



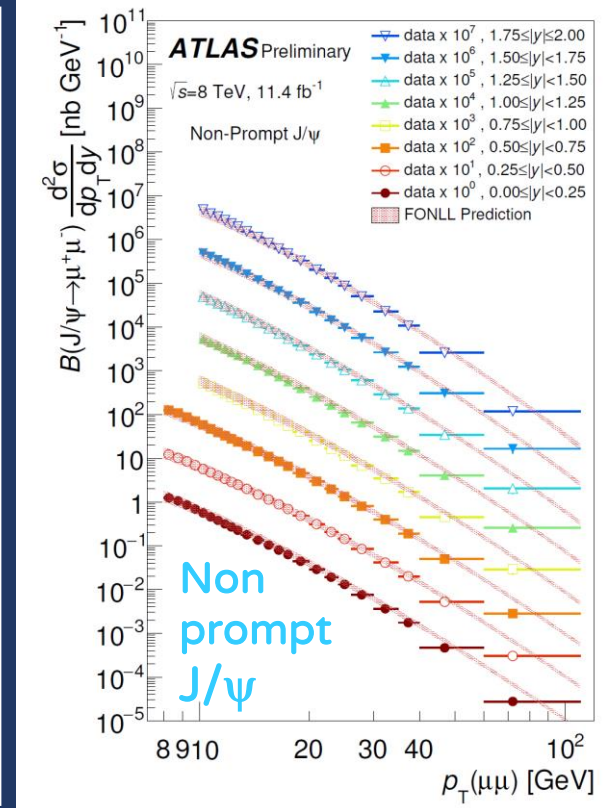
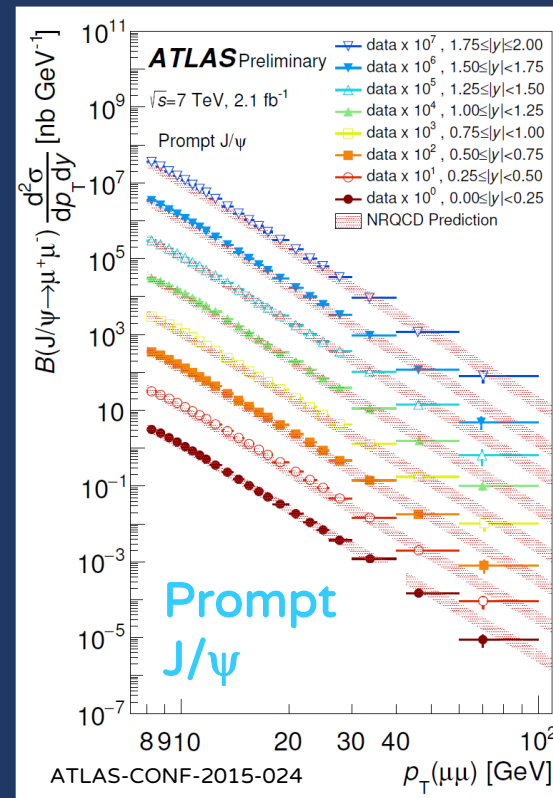
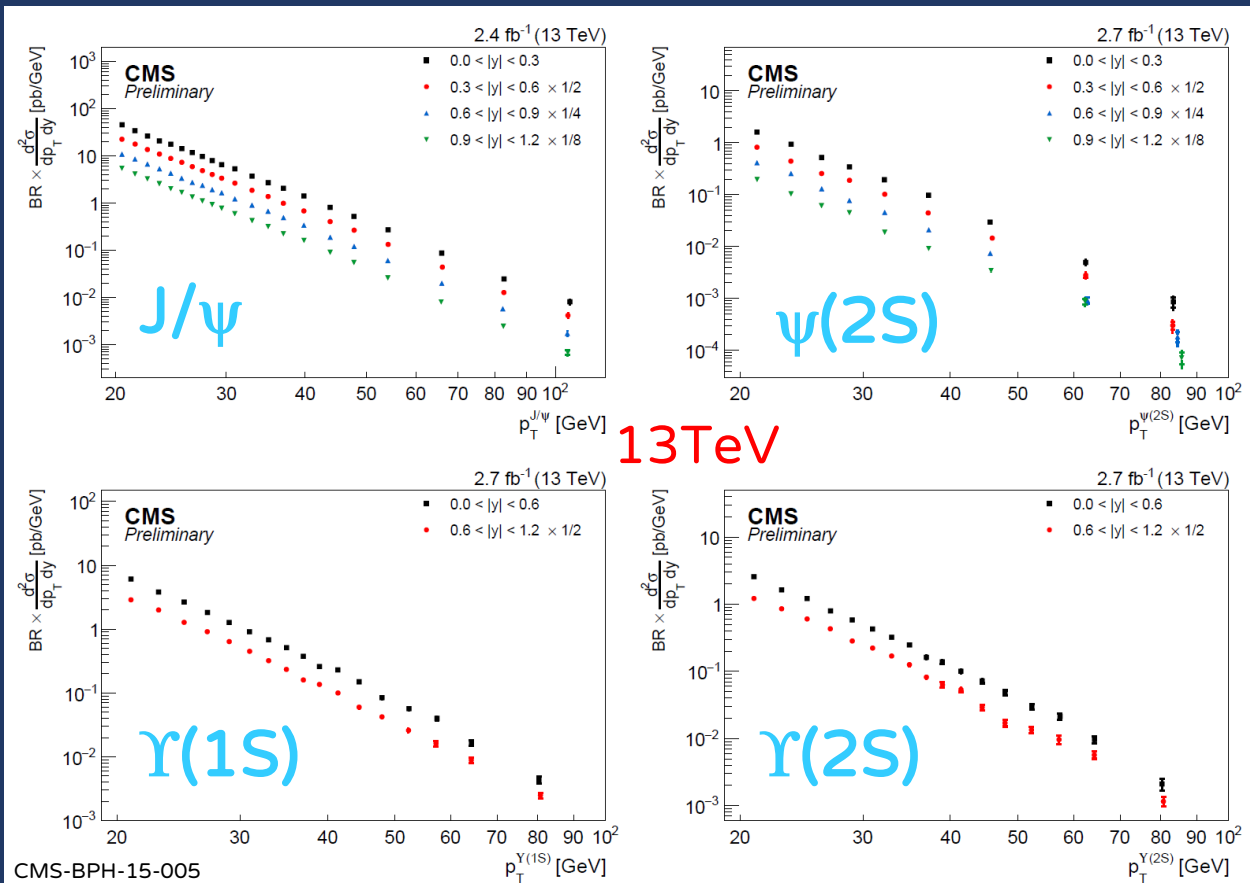
# Comparison between experiments

➔ usually, good agreement between experiments in common kinematic regions  
➔ no hint for significant discrepancies



# Quarkonium $p_T$ cross sections

→  $p_T$  range significantly extended by CMS and ATLAS measurements → up to 120 GeV!



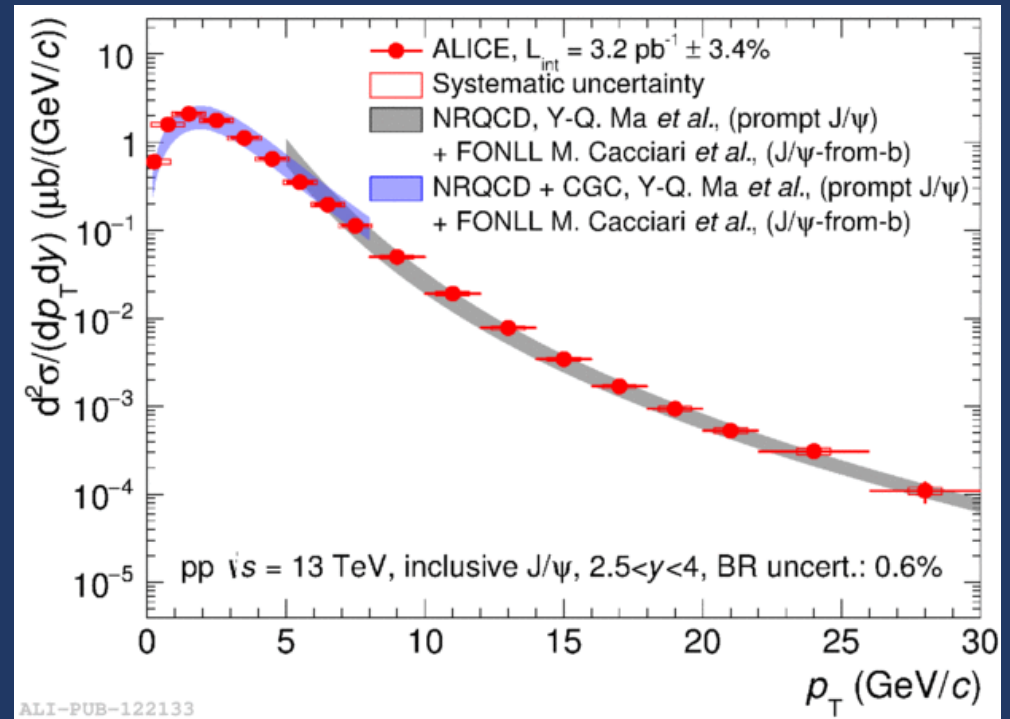
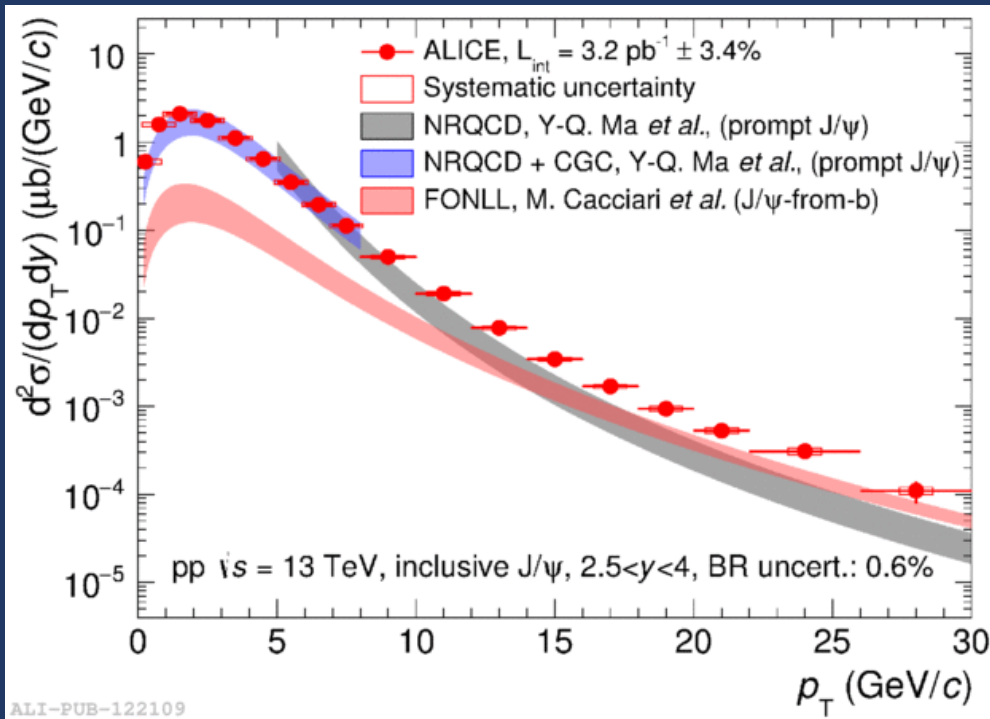
- NLO NRQCD describes prompt  $J/\psi$  production
- FONLL describes non-prompt contribution, but gives slightly harder  $p_T$  spectra



# Inclusive $J/\psi$ production

➔ ALICE inclusive  $J/\psi$  production is described, down to zero  $p_T$ , by a sum of:

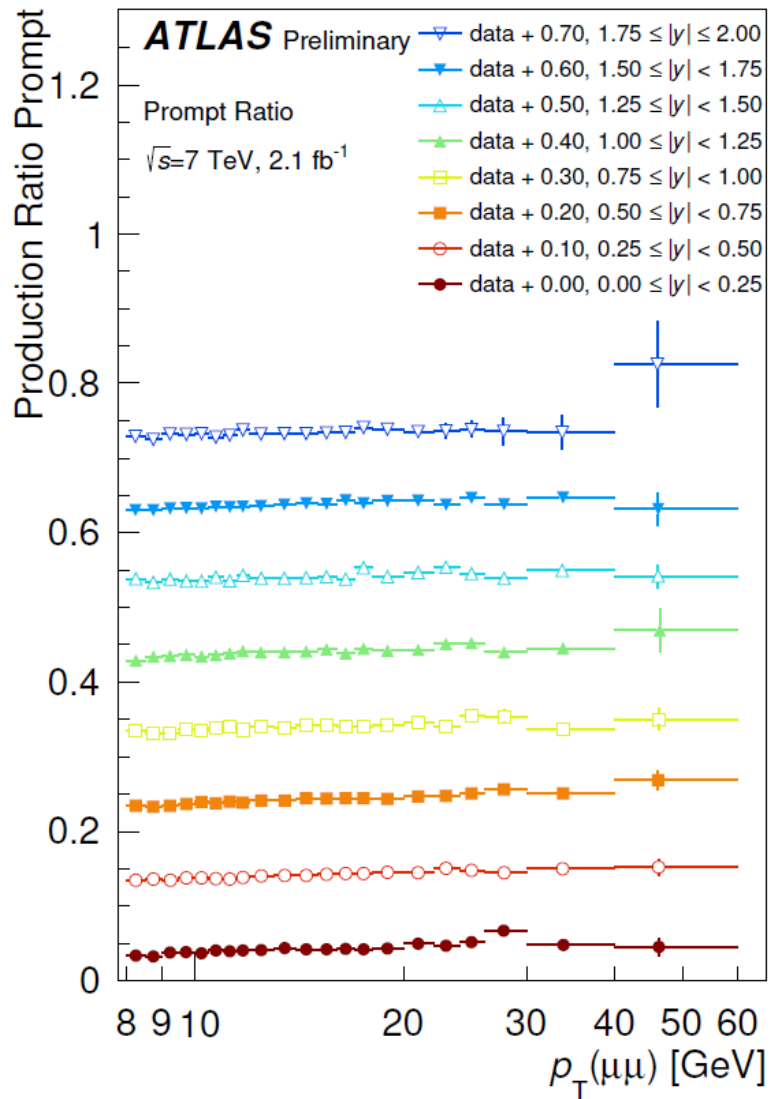
- (NLO) NRQCD for the prompt contribution at intermediate and high  $p_T$
- NRQCD + CGC for prompt  $J/\psi$  at low  $p_T$
- FONLL for  $J/\psi$  from B decay



➔ Similar description holds for all the  $pp$  energies

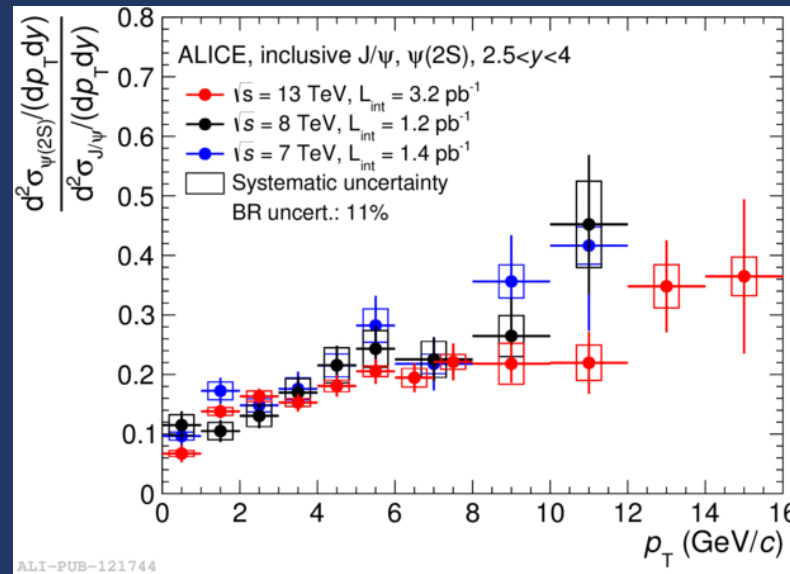


# $\psi(2S)/J/\psi$ in pp

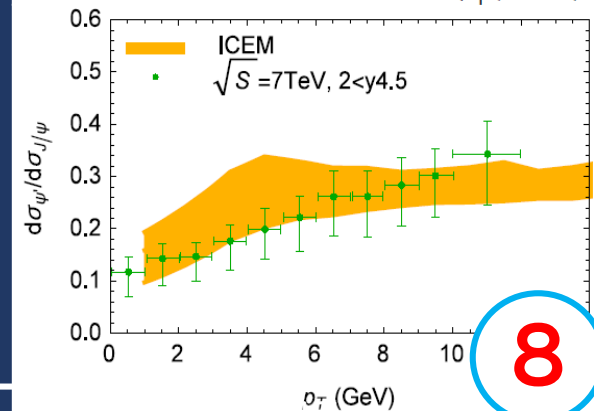
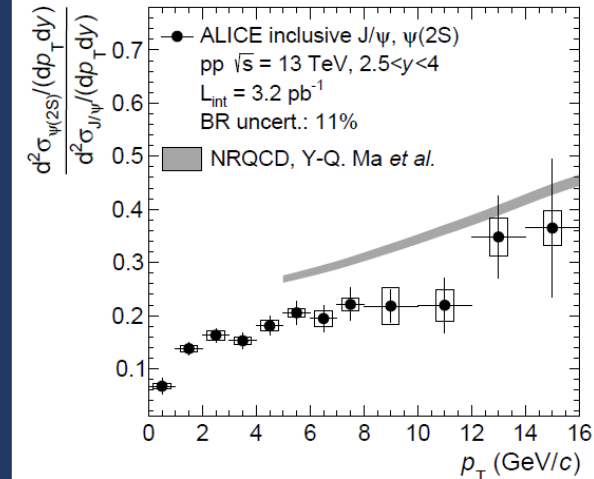
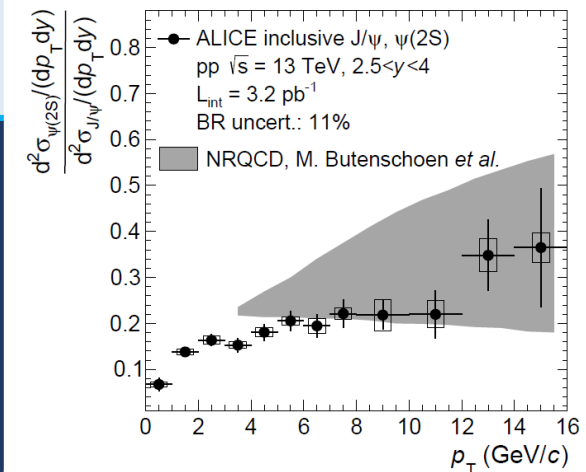


ATLAS-CONF-2015-024

➔ Ratio  $\psi(2S)/\psi$  shows an increase towards high  $p_T$ , with no energy and no  $y$ -dependence

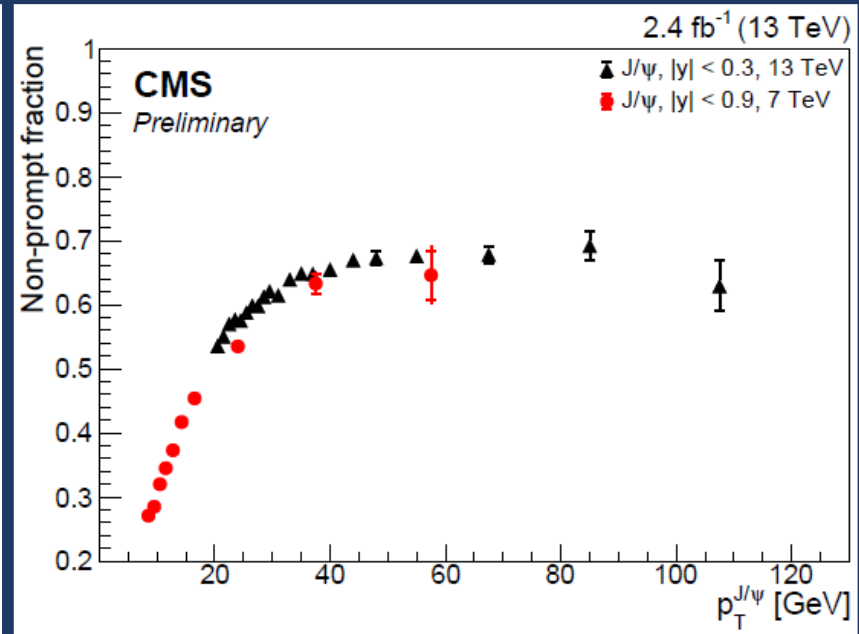
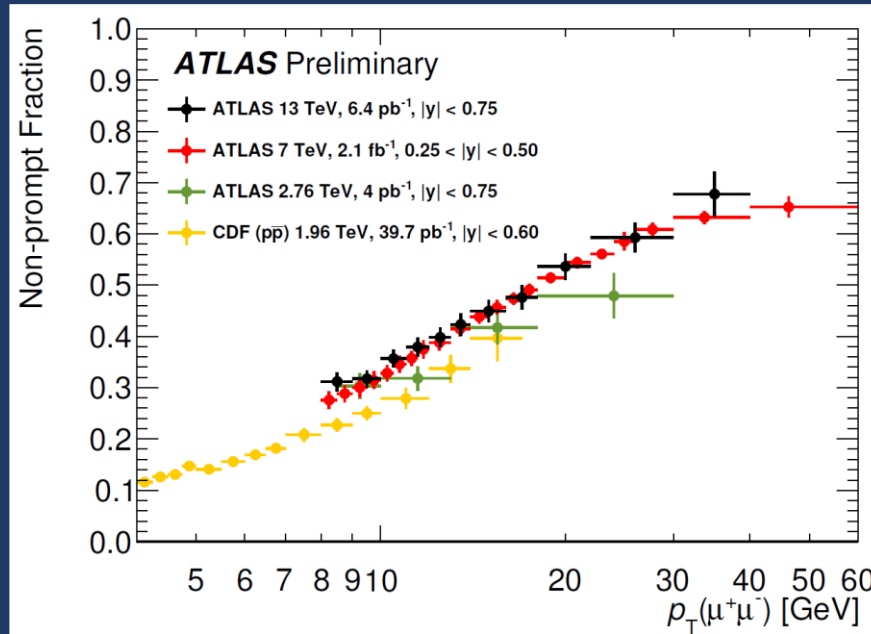
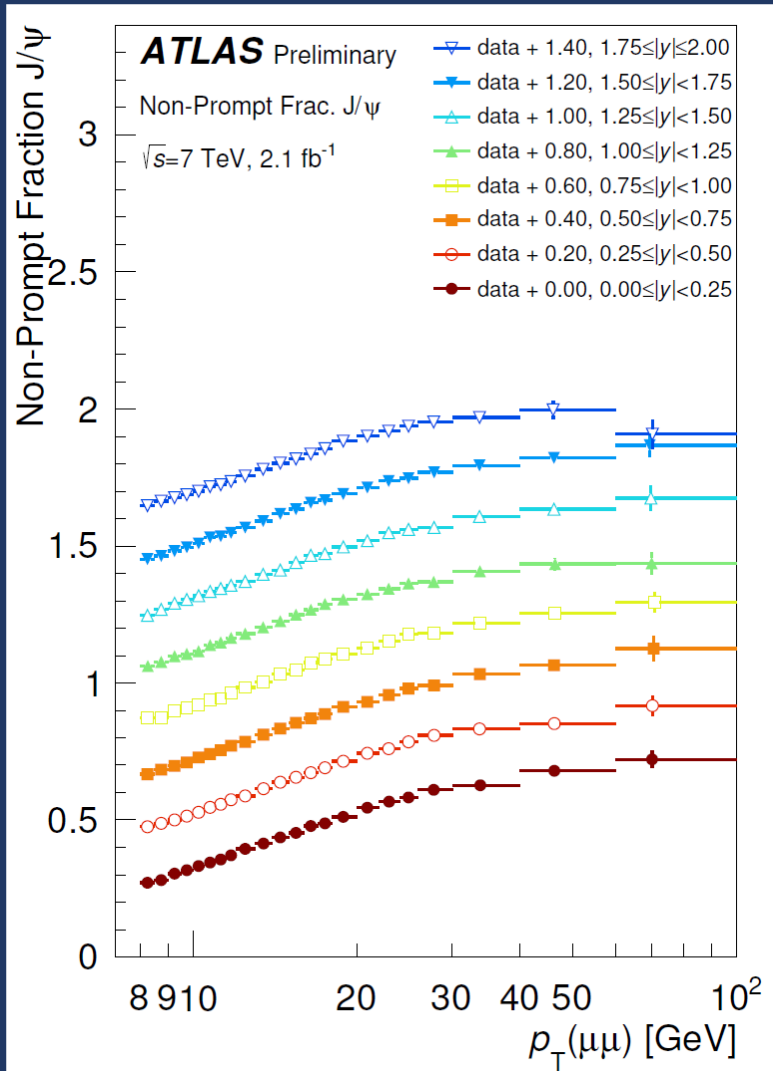


➔  $\psi(2S)/\psi$  is an interesting testing ground for models because of the error cancellation in both data and theory



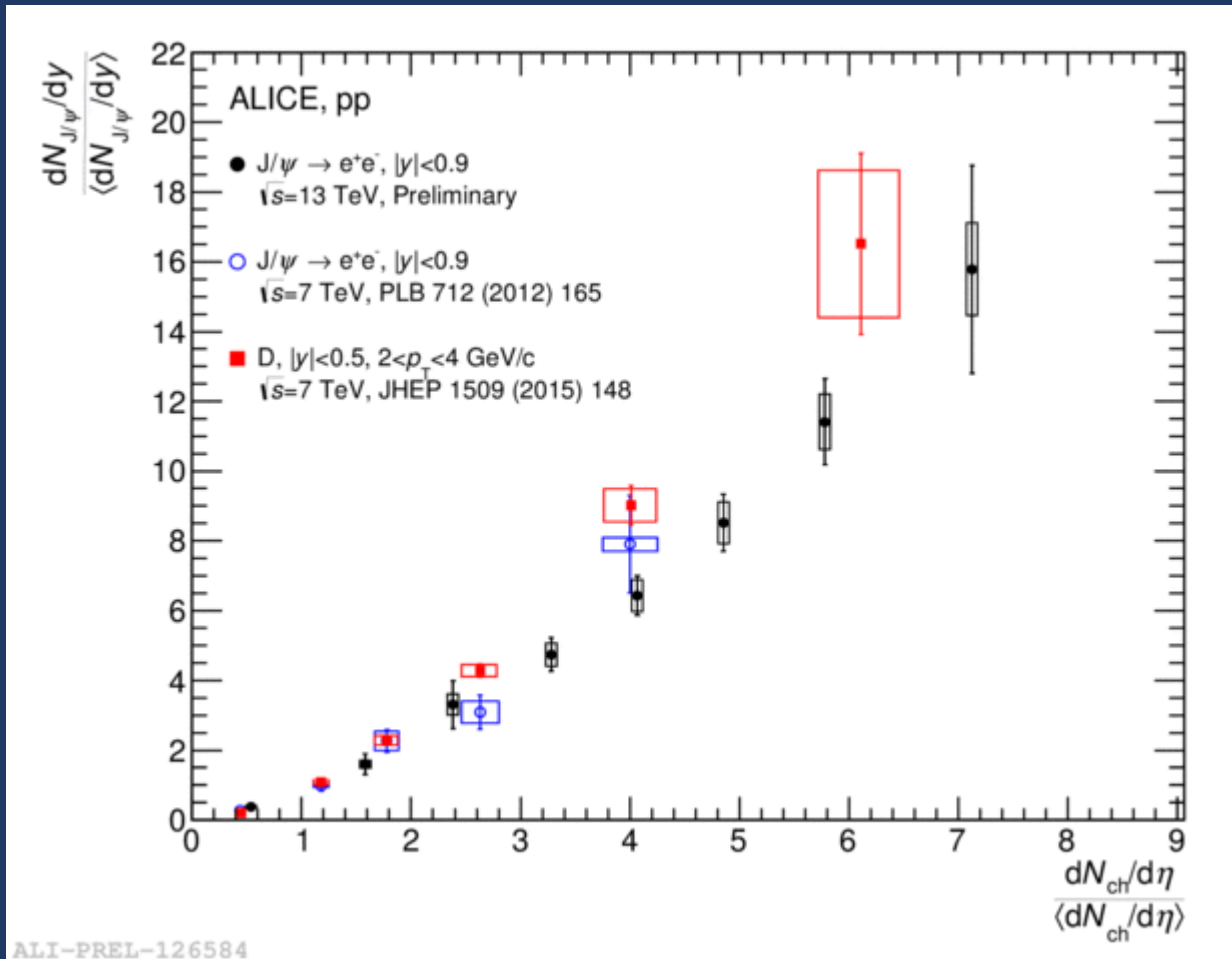
# B feed down contribution

➔ Non-prompt fraction increases steadily, with  $p_T$ , from 10 to 60%, with no significant variation with  $y$ , and then saturates



➔ No  $\sqrt{s}$  dependence is observed between 7 and 13 TeV, while a difference is visible wrt lower energies

# J/ψ production vs hadronic multiplicity



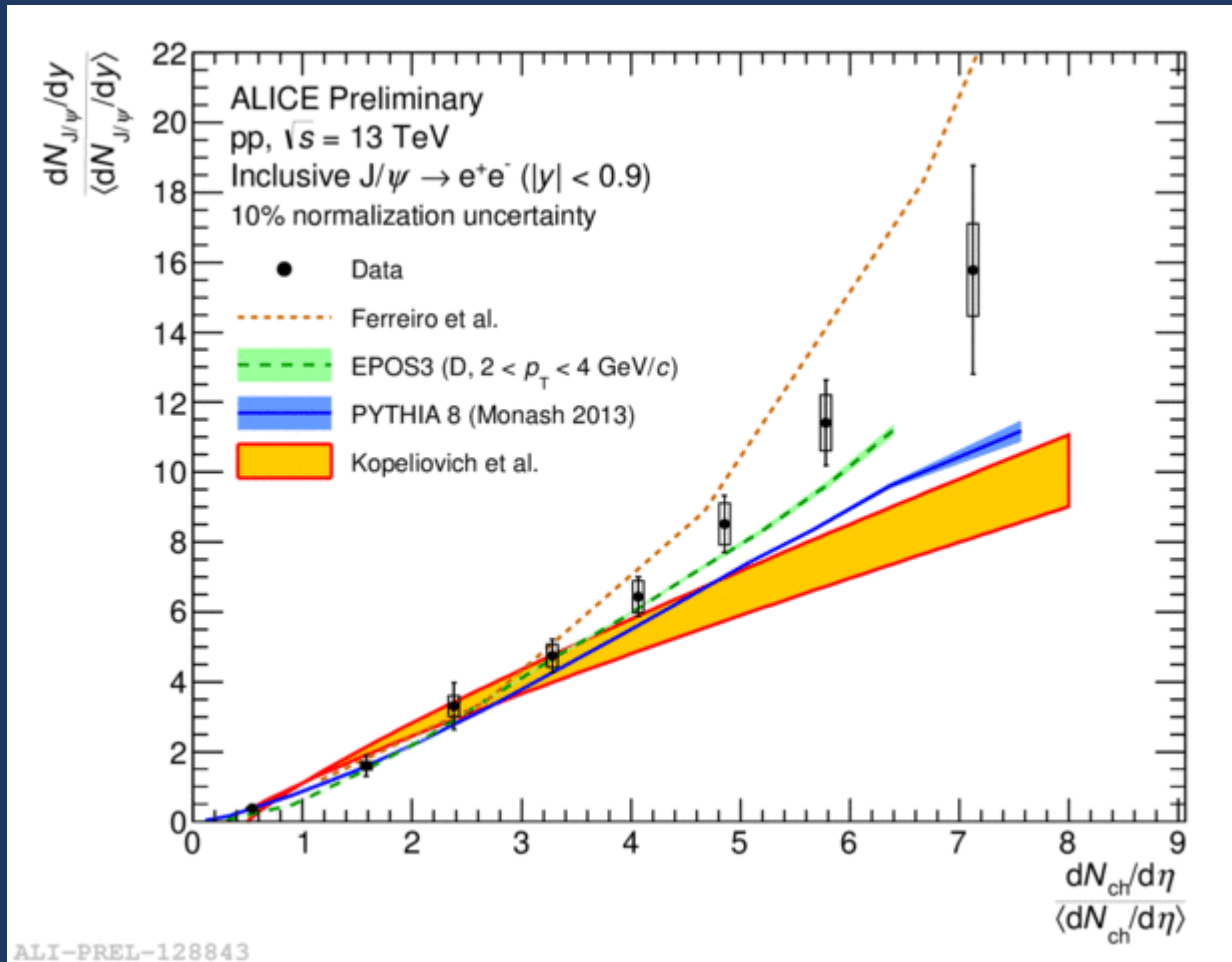
➔ Increase of J/ψ yield with event multiplicity observed at 7 and 13 TeV

➔ Stronger than linear increase, reaching up to 15 times the average J/ψ value, at a multiplicity of about 7 times the mean value

➔ No significant energy dependence

➔ Similar rise for open and closed charm (caveat: different p<sub>T</sub> and y range)

# J/ψ production vs hadronic multiplicity

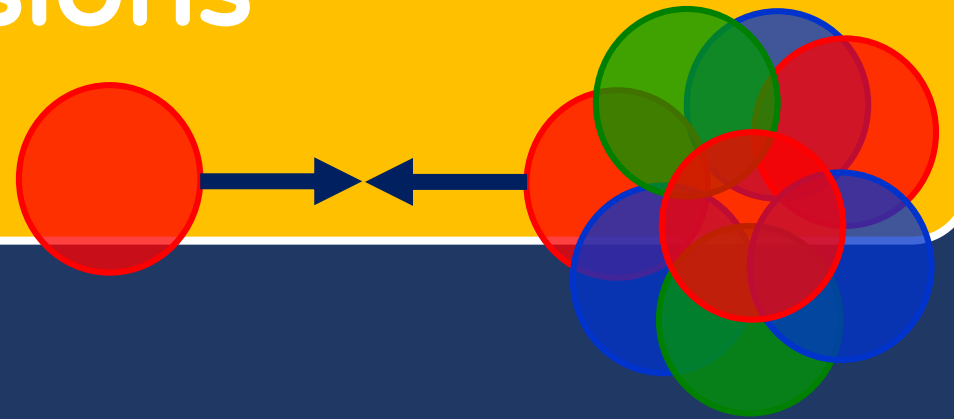


➔ Models describe the lower multiplicity data, while they deviate at high multiplicity

➔ Models attribute the observed behavior to different underlying processes:

- EPOS3 and PYTHIA: include MPI
- Kopeliovich: high multiplicities reached via contribution of higher Fock states
- Percolation: mimic MPI via interactions of colour sources with finite spatial extension

# p-A collisions



# Experimental pA landscape

Facility	Experiment	System	$\sqrt{s_{NN}}$ (GeV)	$Y_{cms}$ range	Data taking
SPS	NA50	p-Be,Al,Cu,Ag,W,Pb	27	$-0.4 < y < 0.6$	1996-2000
			29	$-0.5 < y < 0.5$	
	NA60	p-Be,Al,Cu,In,W,Pb,U	17	$0.3 < y < 0.8$	2004
			27	$-0.1 < y < 0.3$	
FNAL	E866	p-Be, Fe, W	39	$-0.6 < y < 2.5^*$	~1996
HERA	HERA-B	p-C, Ti, W	42	$-1.5 < y < 0.8$	2002
RHIC	PHENIX, STAR	d-Au	200	$-2.2 < y < 2.4$	>2003
		p-Al, Au		$1.2 <  y  < 2.2$	2015
LHC	ALICE	p-Pb	5020 8016	$-4.46 < y < 3.53$	2013 2016
	ATLAS			$-2.87 < y < 1.94$	
	CMS			$-2.87 < y < 1.93$	
	LHCb			$-5.0 < y < -2.5$ $1.5 < y < 4.0$	

➔ a large wealth of data has been collected in pA/dA collisions, in parallel with QGP studies in heavy-ion collisions

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Fixed target experiments:  
Data collected on several A targets



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Collider experiments

- usually p vs a single beam specie
- forward and backward y range might be covered

# How charmonium is studied in pA?

- Varying the amount of nuclear matter crossed by cc pair  
e.g. studying  $J/\psi$  production vs.  $A$  or centrality
- Selecting the kinematics of quarkonium states  
e.g. selecting events where the resonance is formed inside or outside the nucleus
- Comparing the behavior of different resonances

$$1 \quad \sigma_{J/\psi}^{pA} = \sigma_{J/\psi}^{pp} A e^{-\langle \rho L \rangle} \sigma_{abs}$$

the larger  $\sigma_{abs}$ ,  
the more  
important the  
CNM effects

$$2 \quad \sigma_{J/\psi}^{pA} = \sigma_{J/\psi}^{pp} A^\alpha$$

$\alpha = 1$   
→ no nuclear effects  
 $\alpha \neq 1$   
→ nuclear effects

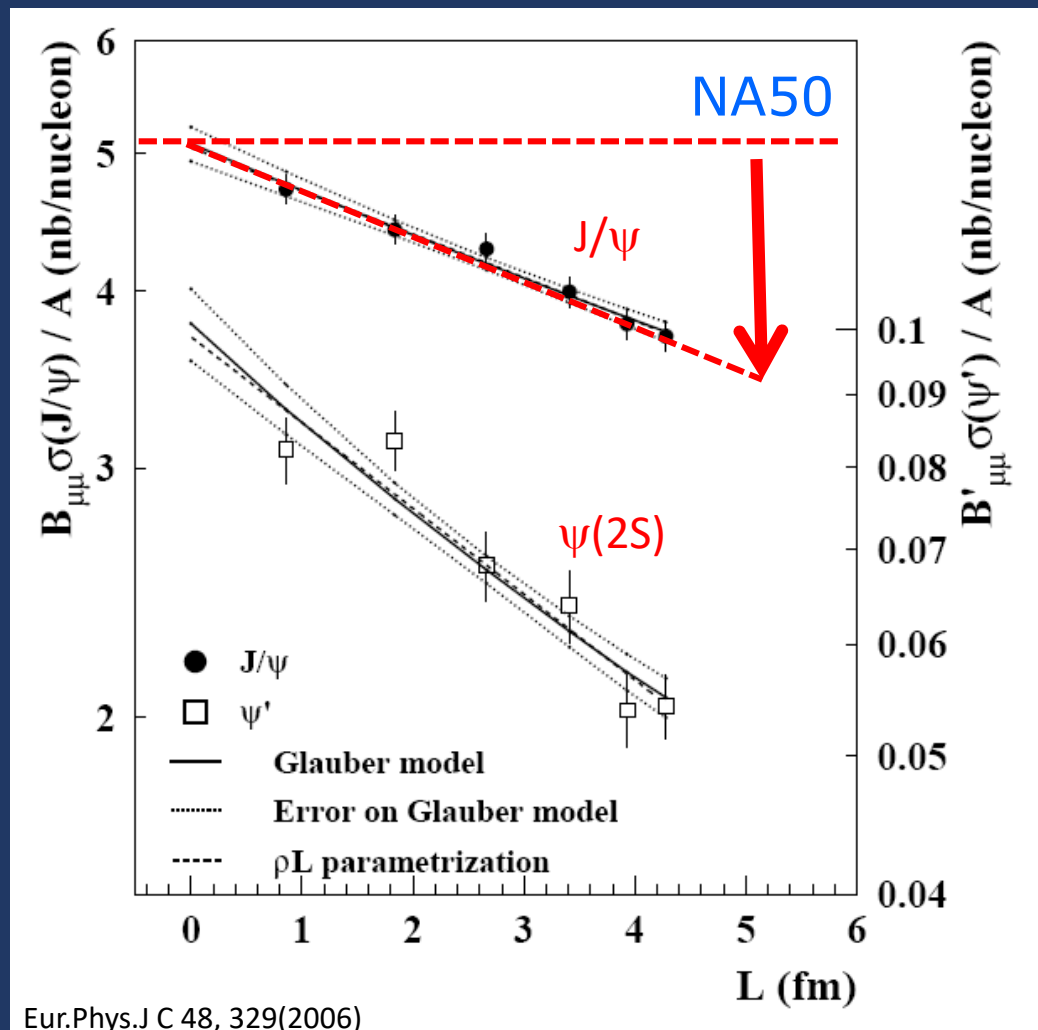
$$3 \quad R_{J/\psi}^{pA} = \frac{\sigma_{J/\psi}^{pA}}{A \sigma_{J/\psi}^{pp}}$$

$R_{pA} = 1$   
→ no nuclear effects  
 $R_{pA} \neq 1$   
→ nuclear effects

→  $\sigma_{abs}$  and  $\alpha$  are “effective” quantities  
which quantify the size of CNM effects

# Results from SPS

➔ A significant reduction of charmonium yields per NN collision is observed



Eur.Phys.J C 48, 329(2006)

➔ reduction interpreted as due to “nuclear absorption” of the cc pair in medium

➔ stronger absorption for the less bound state ψ(2S) at mid-γ

Nucleus crossing time comparable or larger than charmonium formation time:

→ fully formed resonances traverse the nucleus

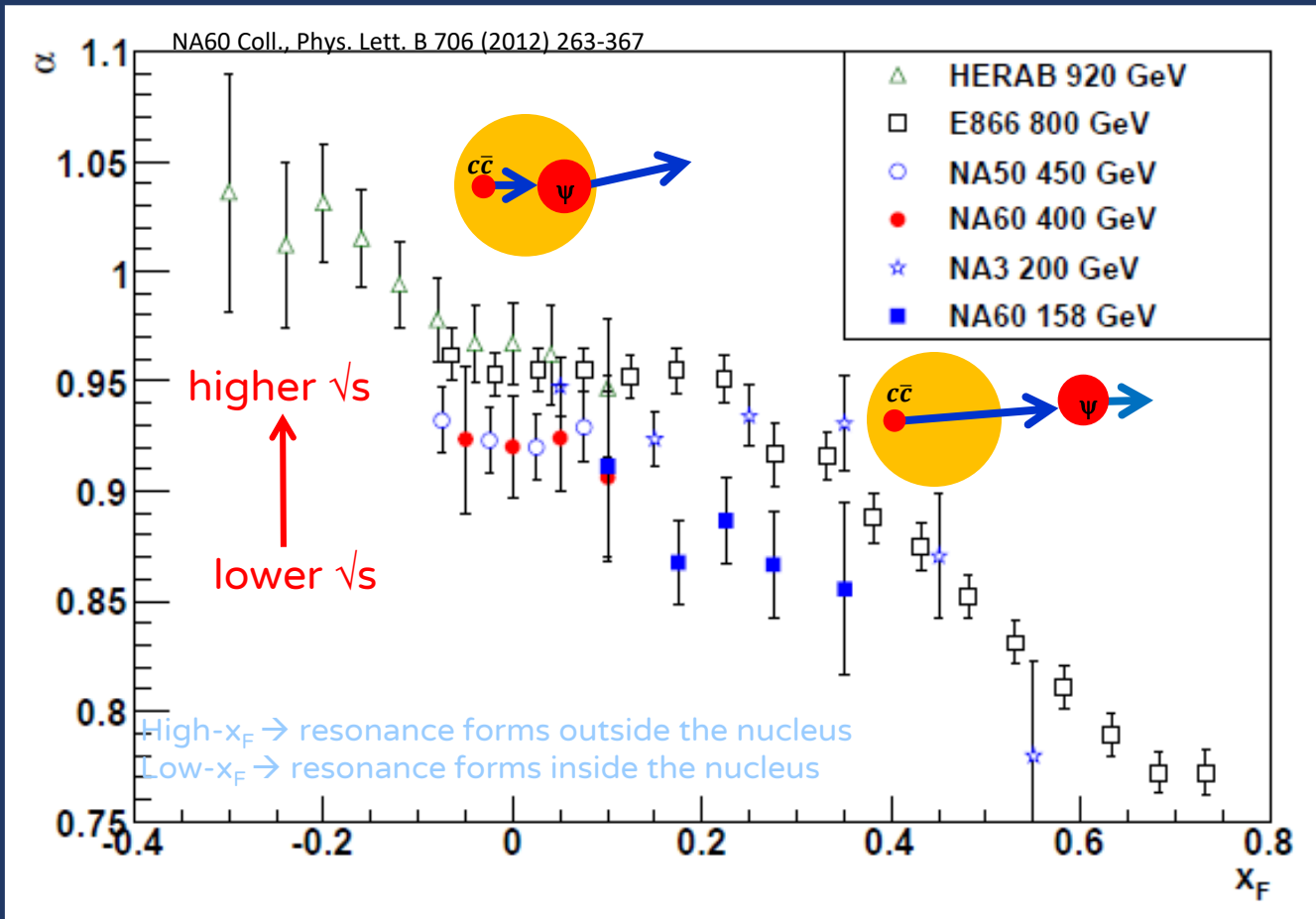
➔ Fitting with  $\sigma_{J/\psi}^{pA} = \sigma_{J/\psi}^{pp} \cdot A \cdot e^{-(\rho L) \sigma_{abs}}$

$$\sigma_{abs} J/\psi = 4.5 \pm 0.5 \text{ mb}$$

$$\sigma_{abs} \psi(2S) = 8.3 \pm 0.9 \text{ mb}$$

# $J/\psi$ as a function of $x_F$

Compilation of  $J/\psi$  results obtained in several fixed target experiments

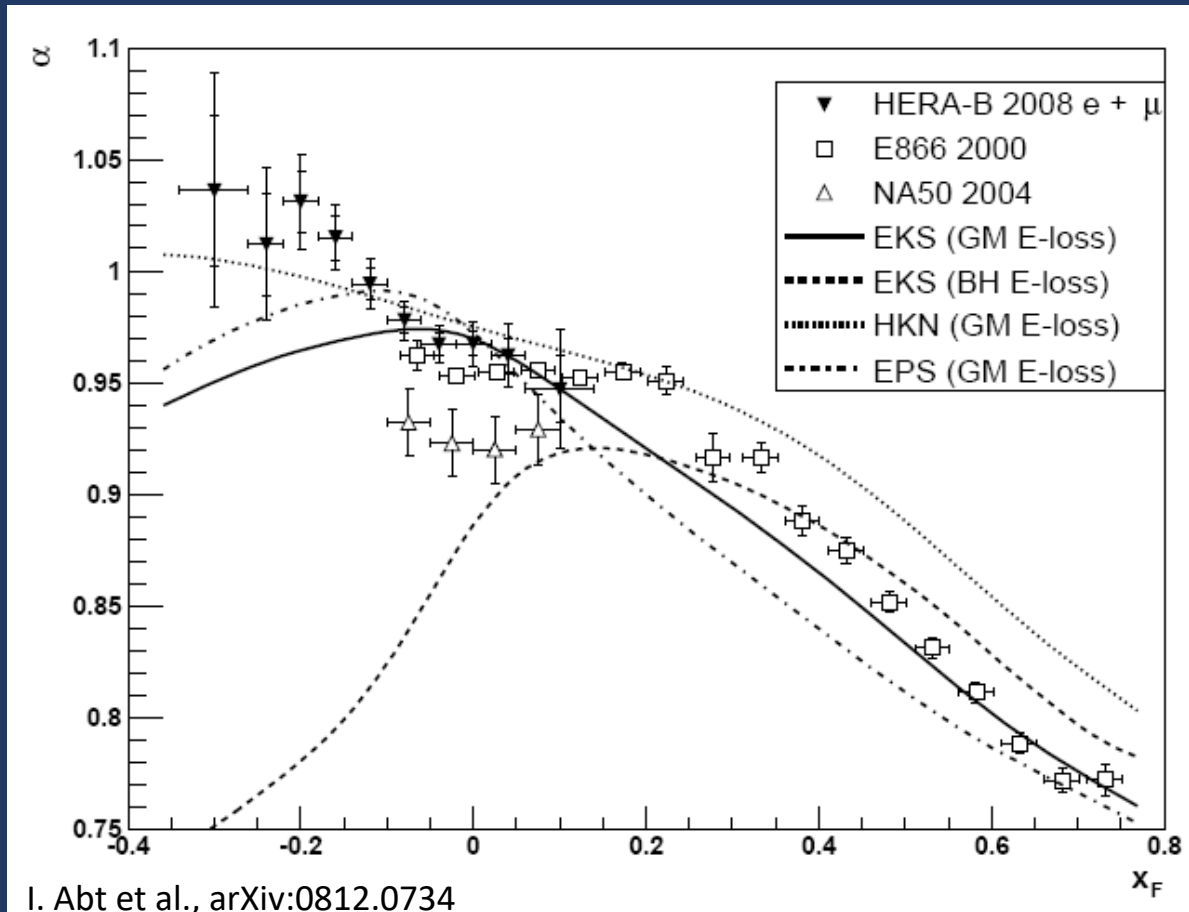


$J/\psi$  yield in pA is modified with respect to pp collisions, with a strong kinematic dependence

- $\alpha$  strongly decreases with  $x_F$
- for a fixed  $x_F$ , CNM are stronger at lower  $\sqrt{s}$

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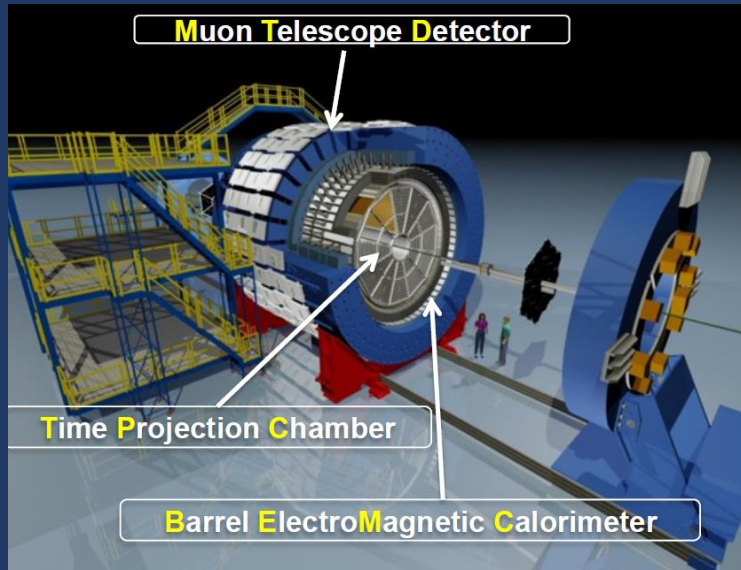
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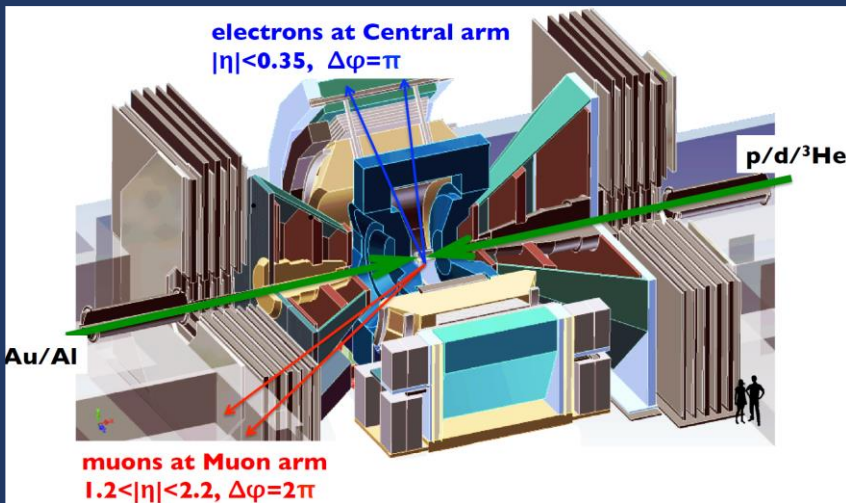
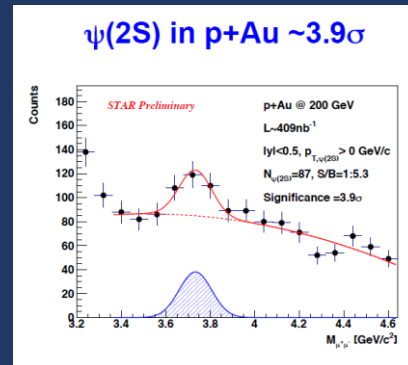
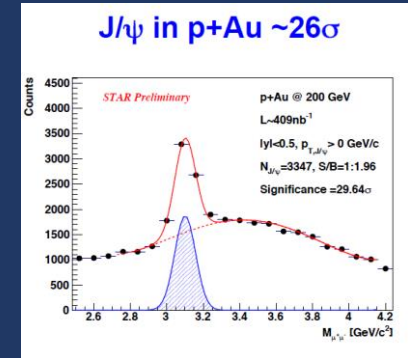
Theoretical description over the full  $x_F$  range still difficult!

# pA-dA data taking at RHIC



➔ Several collision systems and energy investigated at RHIC:

$\sqrt{s}$ (GeV)	pp	p-Al	p-Au	d-Au	$^3\text{He}$ -Au
510	X				
200	X	X	X	X	X
130				X	
62.4	X			X	
39				X	
27				X	
20				X	



➔ Significant improvements:

STAR

→ dimuon trigger with MTD, enhancing  $J/\psi$  and  $\Upsilon$  capabilities

PHENIX

→ VTX and FVTX improving tracking and vertexing

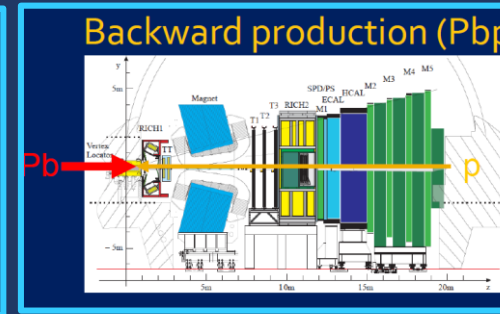
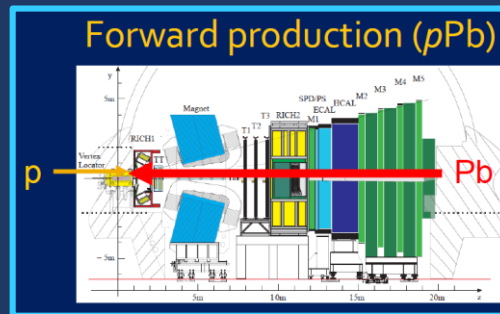
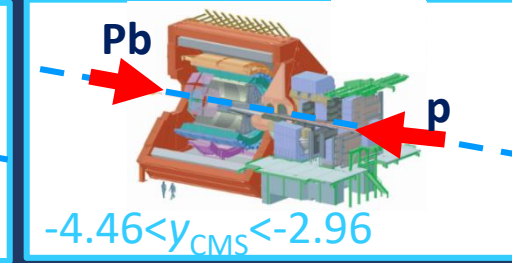
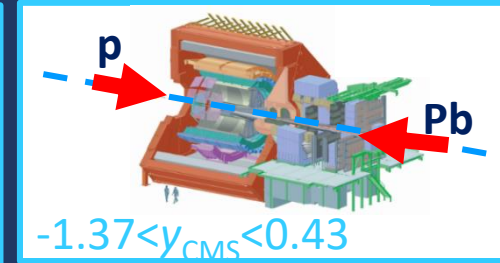
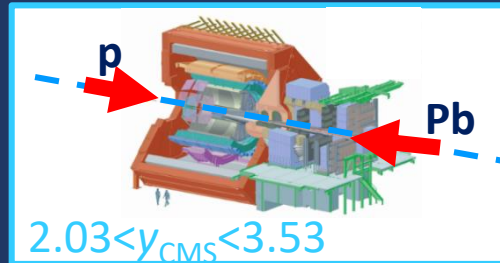


# pA data taking at LHC

➔ pPb collisions at  $\sqrt{s_{NN}} = 5.02$  and 8.16 TeV

➔ ALICE and LHCb data are collected with two beam configurations: p-Pb and Pb-p, with  $\Delta y = 0.465$

	$\sqrt{s_{NN}}$ (TeV)	fw-y	mid-y	bck-y
ALICE	5.02	$5\text{nb}^{-1}$	$51\mu\text{b}^{-1}$	$6\text{nb}^{-1}$
	8.16	$9\text{nb}^{-1}$	n.a.	$13\text{nb}^{-1}$
LHCb	5.02	$1\text{nb}^{-1}$		$0.5\text{nb}^{-1}$
	8.16	$14\text{nb}^{-1}$		$21\text{nb}^{-1}$
CMS	5.02		$35\text{nb}^{-1}$	
	8.16		n.a.	
ATLAS	5.02		$28\text{nb}^{-1}$	
	8.16		$0.2\text{pb}^{-1}$	



➔ Significant statistics increase between Run1 and Run2



# pA results from RHIC and LHC

➔ All quarkonium states have been extensively studied in pA (dA) collisions

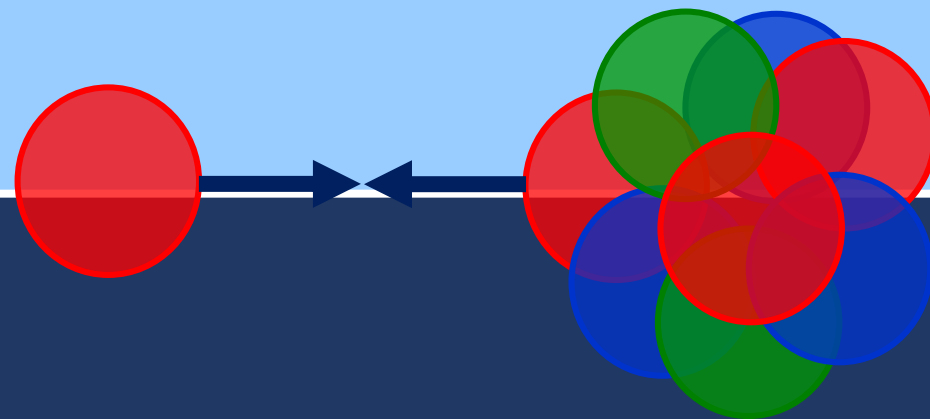
	J/ψ				ψ(2S)				Υ			
	pT	y	centrality	multip.	pT	y	centrality	multip.	pT	y	centrality	multip
$\sqrt{s_{NN}} = 5.02 \text{ TeV}$												
ALICE	x	x	x	x	x	x	x		x	x		
ATLAS	x	x	x	x	x	x	x	x	x	x	x	x
CMS	x	x		x	x	x						x
LHCb	x	x			x	x			x	x		
$\sqrt{s_{NN}} = 200 \text{ GeV}$												
PHENIX	x	x	x			x	x	x		x		
STAR	x	x	x							x		

# pA results from RHIC and LHC

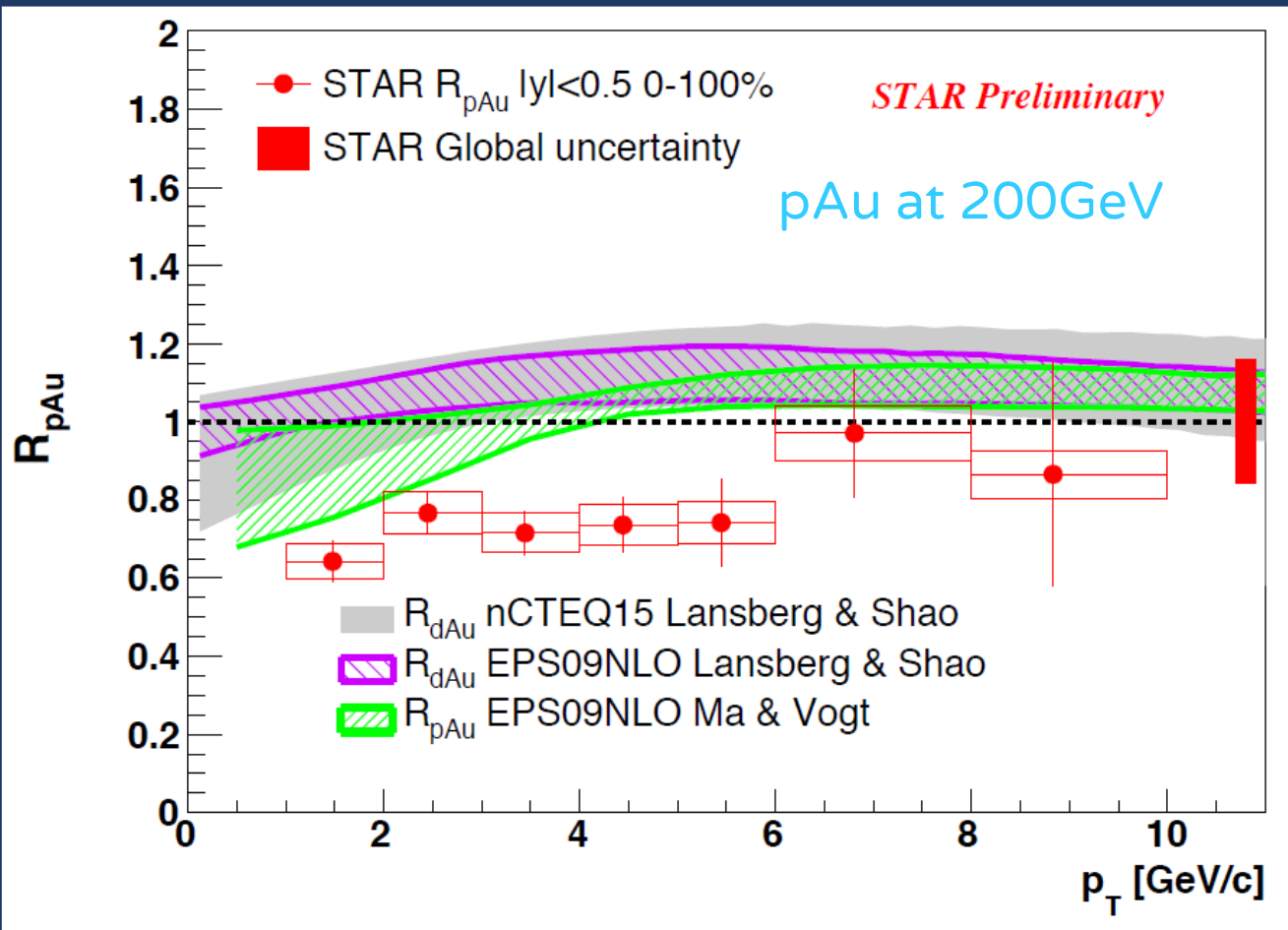
➔ Analysis of the pA data at  $\sqrt{s_{NN}} = 8.16\text{TeV}$  still at the beginning...

	J/ψ				ψ(2S)				Υ			
	pT	y	centrality	multip.	pT	y	centrality	multip.	pT	y	centrality	multip
$\sqrt{s_{NN}} = 8.16\text{ TeV}$												
ALICE	x	x										
ATLAS												
CMS												
LHCb	x	x										

p-A collisions:  
J/ψ



# $J/\psi R_{pA}$ at RHIC

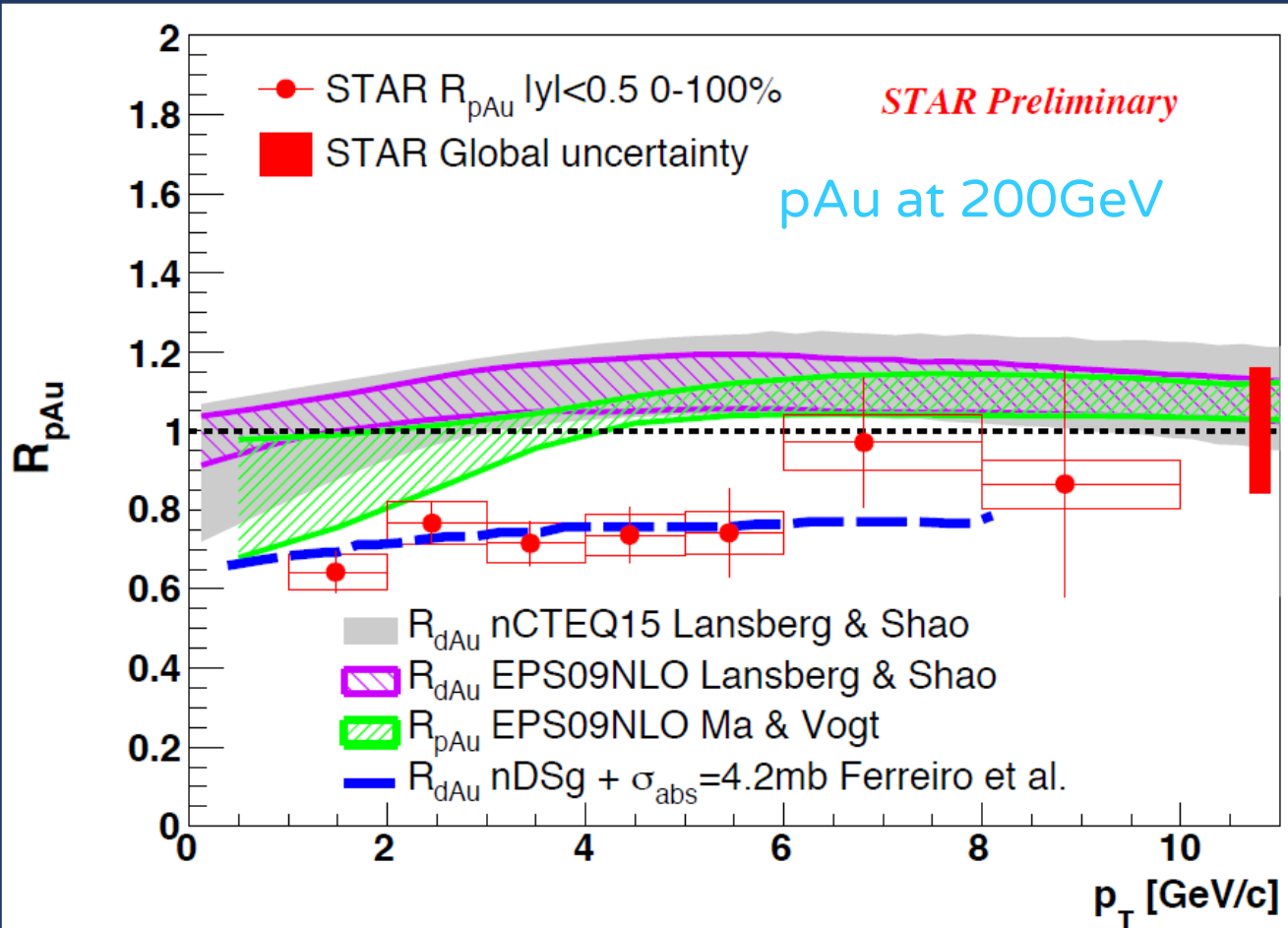


→  $J/\psi R_{pA}$  shows a slightly increasing trend towards high  $p_T$

→ Shadowing models predicts  $R_{pA}$  slightly higher than unity

→ Is there room for other CNM effects on top of shadowing?

# $J/\psi R_{pA}$ at RHIC

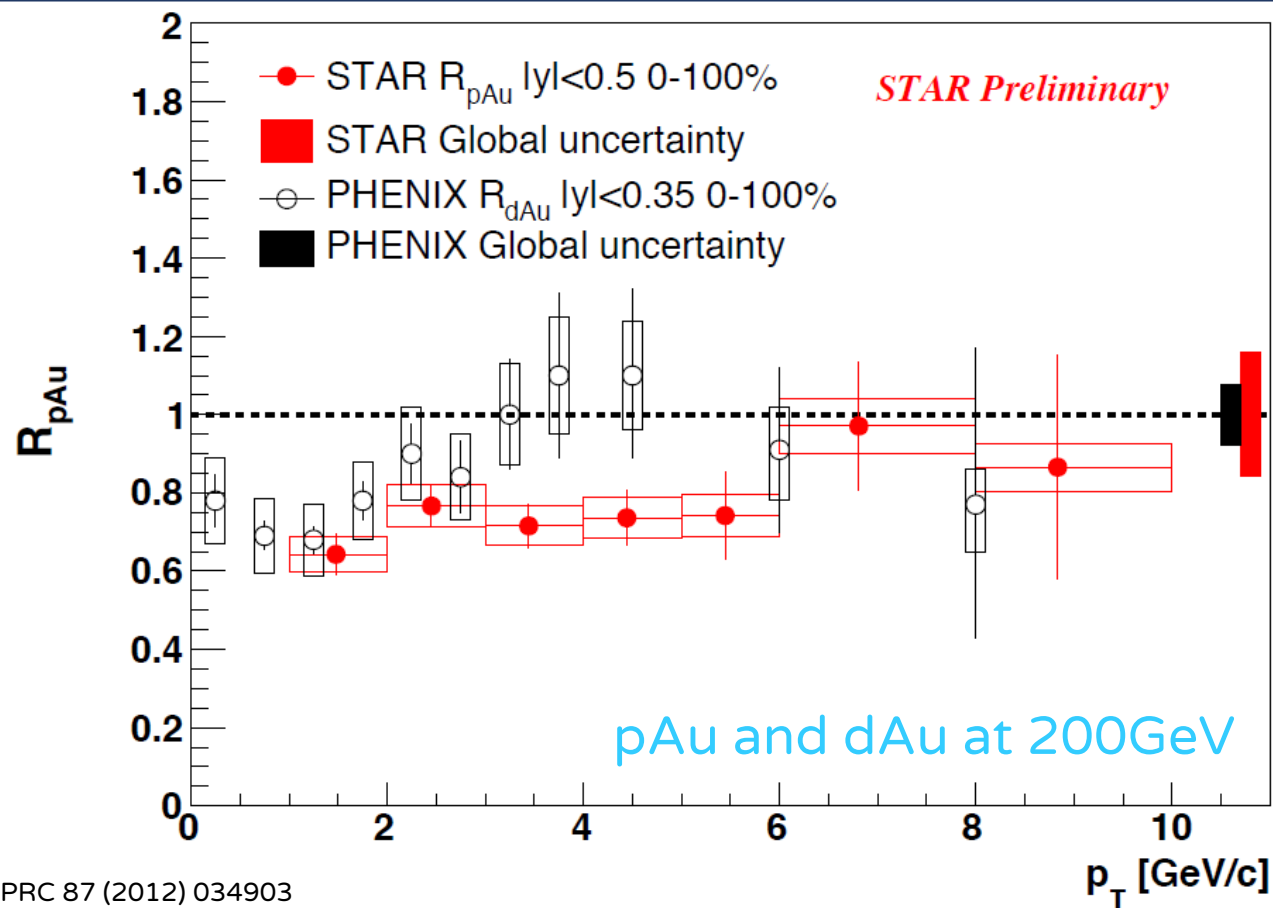


➔  $J/\psi R_{pA}$  shows a slightly increasing trend towards high  $p_T$

➔ Shadowing models predicts  $R_{pA}$  slightly higher than unity

➔ Data seems to allow the inclusion of an additional contribution, as the cc break up in medium, on top of shadowing

# $J/\psi R_{pA}$ at RHIC: pAu vs dAu

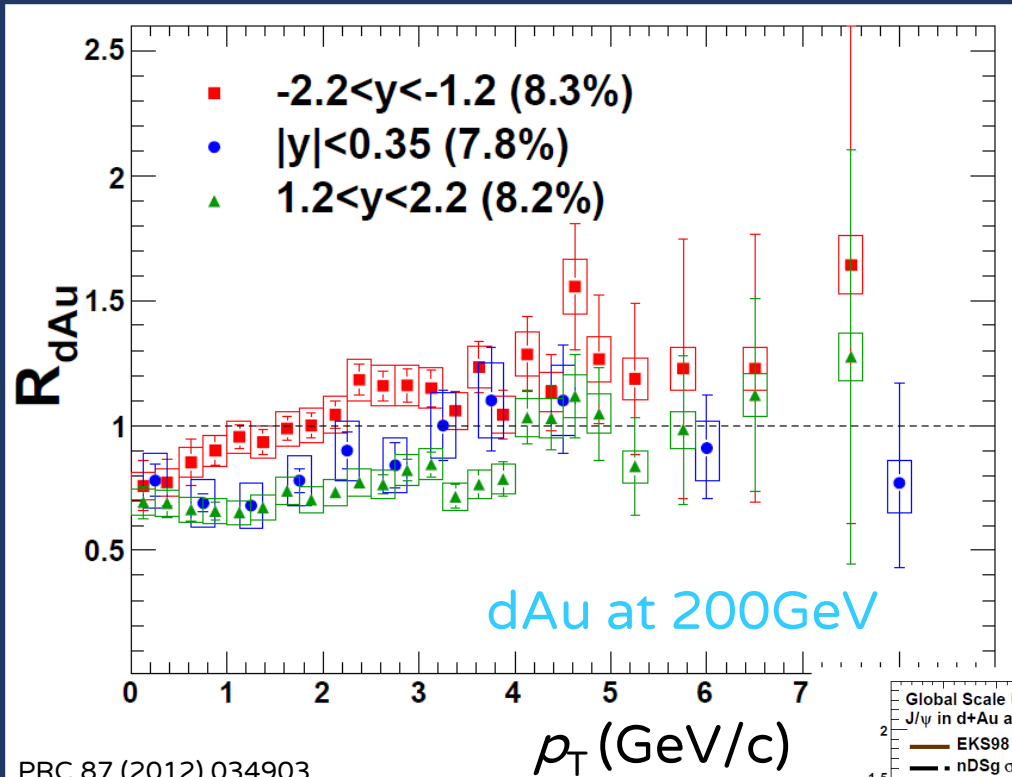


→  $R_{dAu}$  pattern is consistent, within uncertainties, with  $R_{pAu}$  at the same energy

→ (rather) similar CNM effects in pAu and dAu

→ ...but  $R_{dA}$  may be increasing faster with  $p_T$  (at  $p_T \sim 3.5-5\text{GeV}$ , significance is  $1.4\sigma$ )

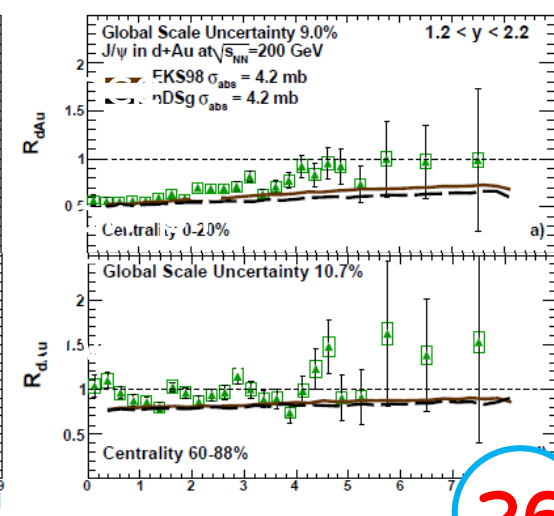
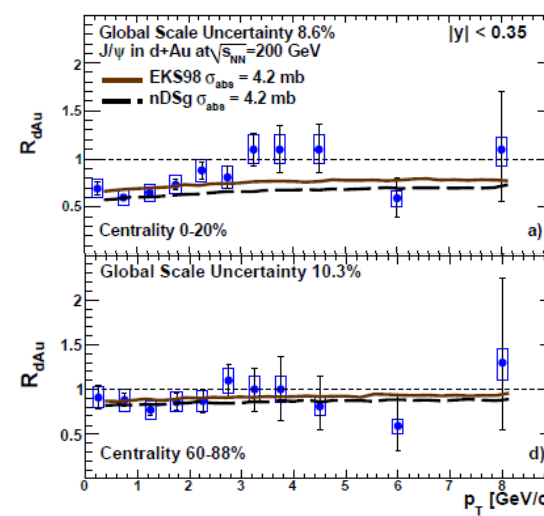
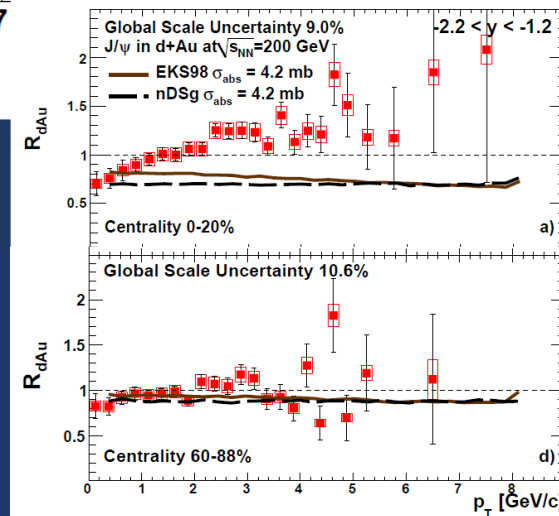
# $J/\psi$ $R_{dAu}$ at RHIC



- ➔  $R_{dAu}$  shows a  $p_T$  and  $y$  dependent trend:
- CNM effects are stronger at low  $p_T$
  - $R_{dAu}$  approaches unity at high  $p_T$
  - CNM effect are more sizeable at forward- $y$

➔ Models based on shadowing +  $c\bar{c}$  breakup describe  $R_{dAu}$ , except for central collisions and negative  $y$

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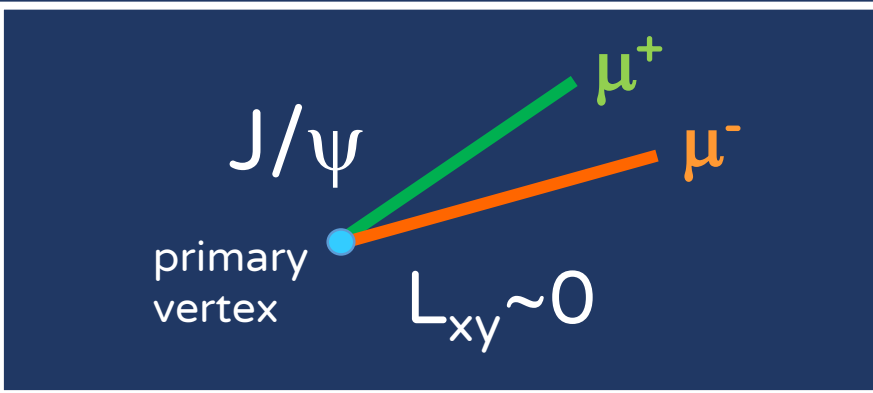




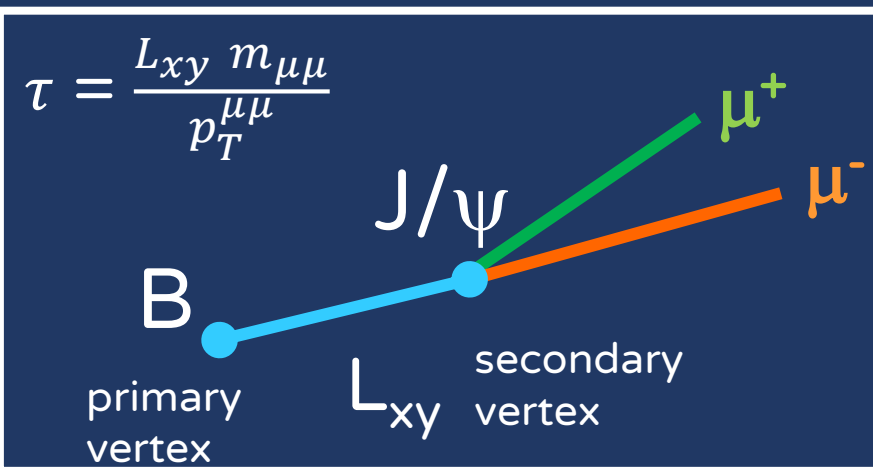
# prompt and non-prompt J/ψ at LHC

→ prompt and non-prompt J/ψ are separated through 2D fit to mass and pseudo-proper decay time

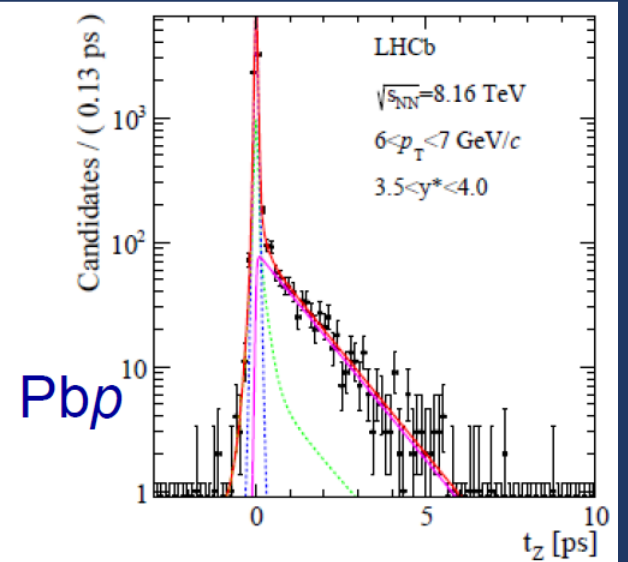
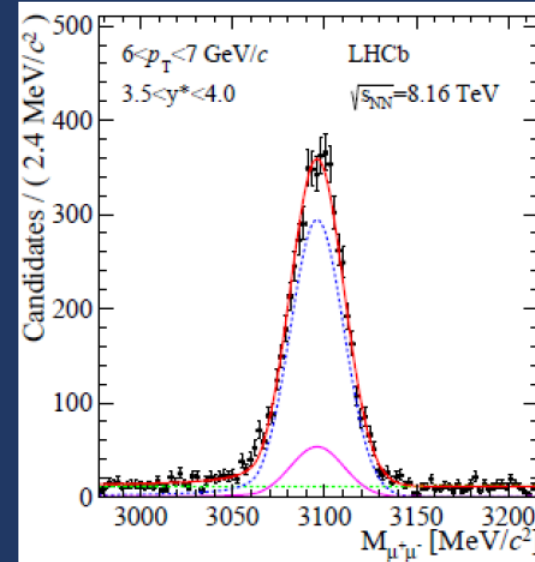
Prompt J/ψ



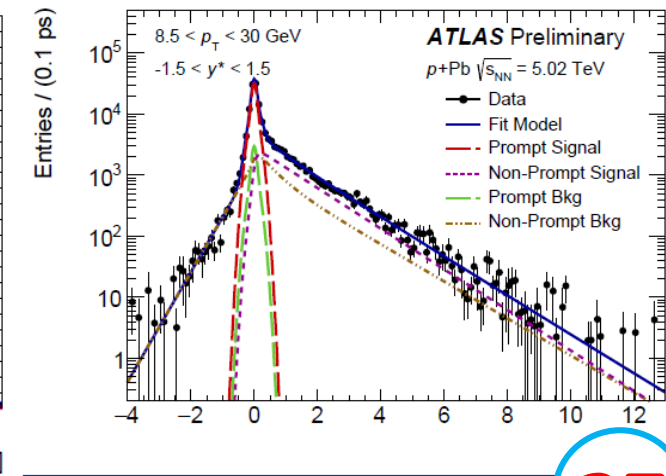
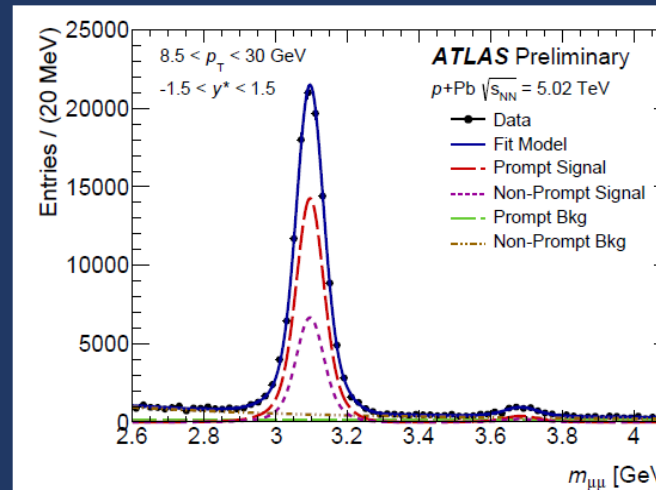
non Prompt J/ψ



$$\tau = \frac{L_{xy} m_{\mu\mu}}{p_T^{\mu\mu}}$$



LHCb-PAPER-2017-014

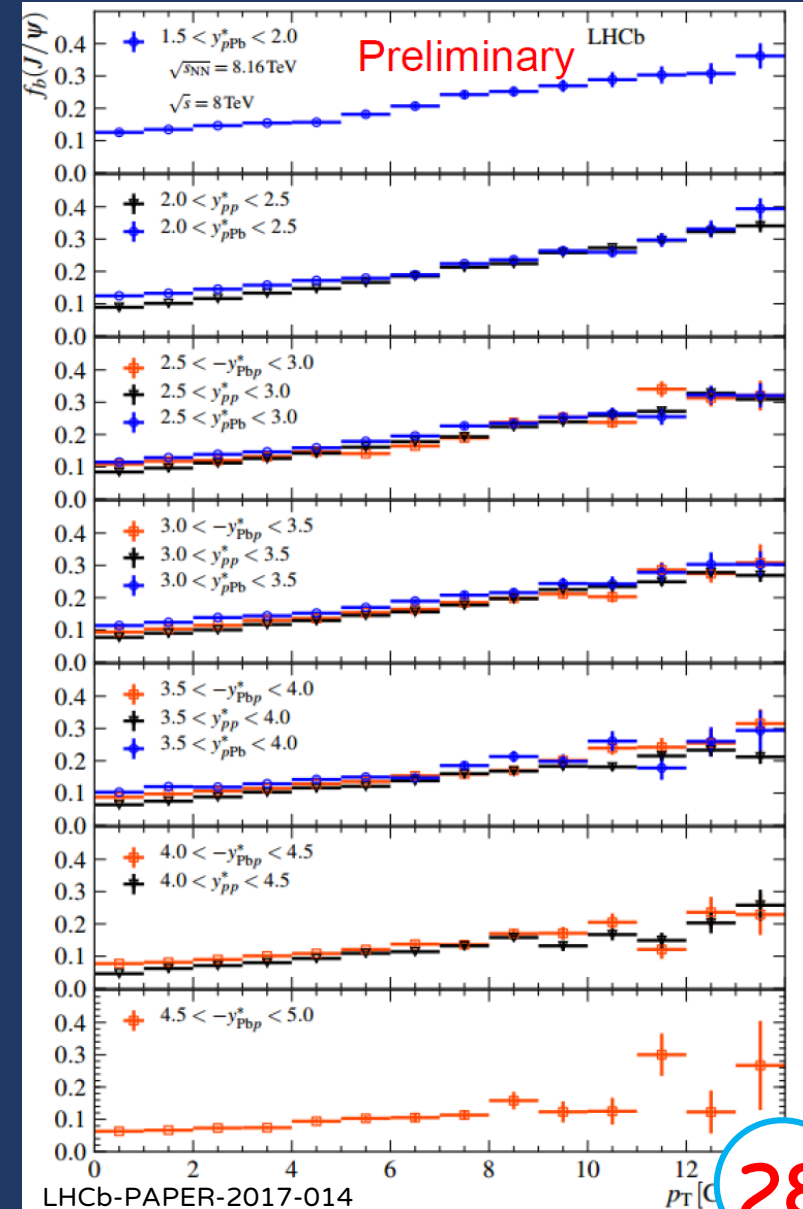
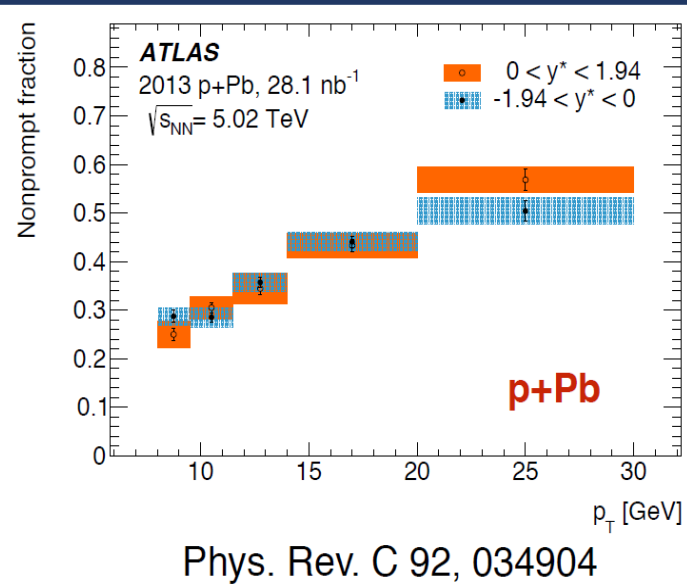
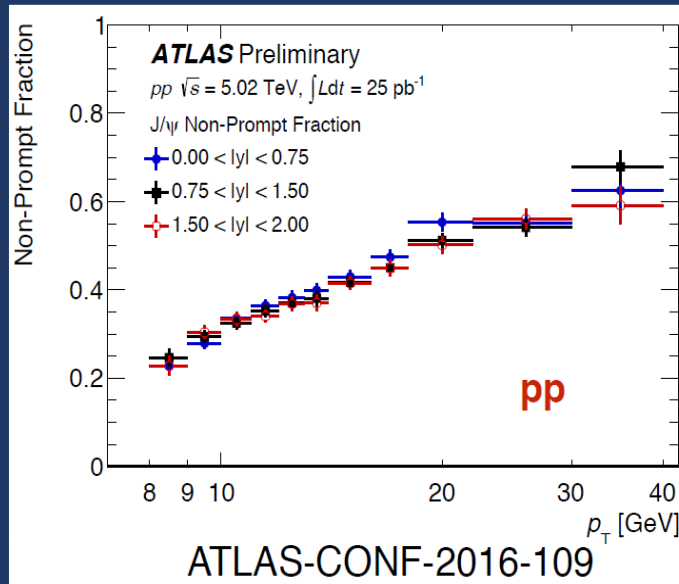


# Fraction of $J/\psi$ from B

➔ Fraction of  $J/\psi$  from B can be evaluated as

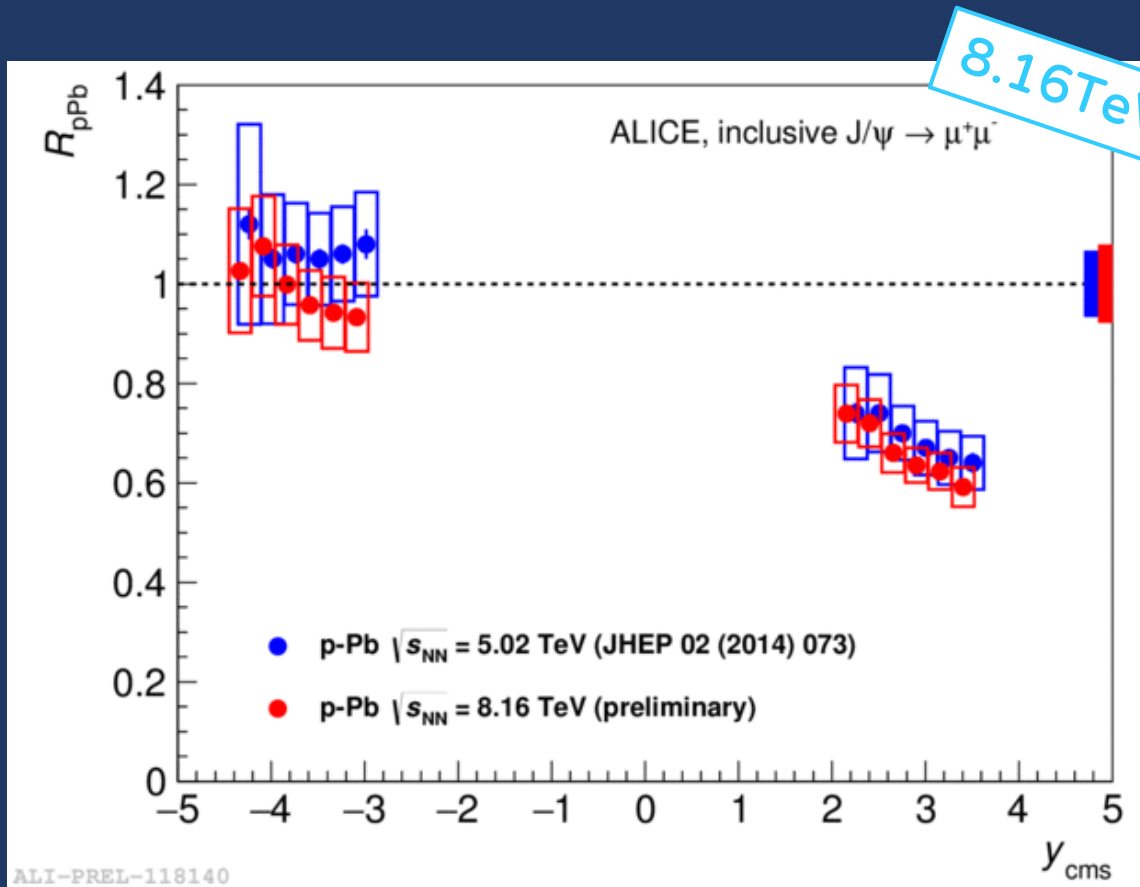
$$F_B = \frac{\sigma_{\text{non prompt } J/\psi}}{\sigma_{\text{prompt } J/\psi} + \sigma_{\text{non prompt } J/\psi}}$$

➔ Similar  $p_T$  dependence in pp, pPb and PbP



➔  $F_B$  increases from 10% at low  $p_T$  up to 40-60% at high  $p_T$ , with a weak  $y$  dependence

# $J/\psi$ $R_{pA}$ at $\sqrt{s_{NN}} = 5.02$ and $8.16$ TeV



CERN-ALICE-PUBLIC-2017-001

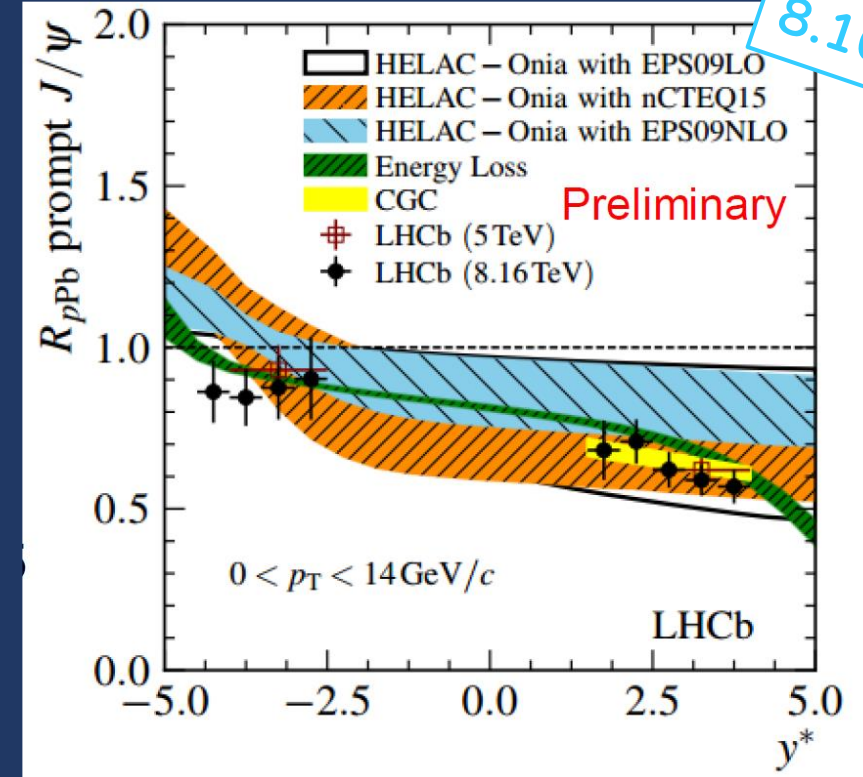
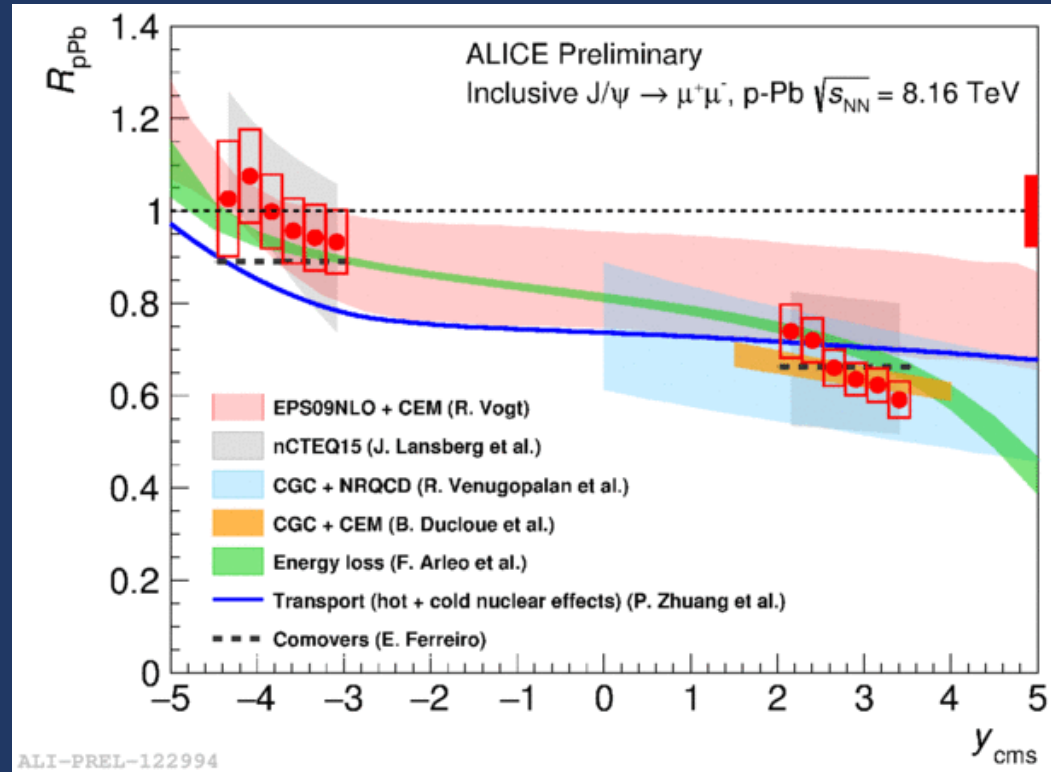
→ Clear  $J/\psi$  suppression at forward- $y$ , while  $R_{pA}$  is compatible with unity at backward- $y$

→  $R_{pA}$  compatible at  $\sqrt{s_{NN}} = 5.02$  and  $8.16$  TeV, even if  $x$  coverage is slightly different

	$\sqrt{s_{NN}} = 8.16$ TeV, $p_T^{J/\psi} = 0$	$\sqrt{s_{NN}} = 5.02$ TeV, $p_T^{J/\psi} = 0$
p-going	$1.1 \cdot 10^{-5} < x < 5 \cdot 10^{-5}$	$2 \cdot 10^{-5} < x < 8 \cdot 10^{-5}$
Pb-going	$7.3 \cdot 10^{-3} < x < 3.3 \cdot 10^{-2}$	$1 \cdot 10^{-2} < x < 5 \cdot 10^{-2}$

$2 \rightarrow 1$  kinematics,  $p_T = 0$

# $J/\psi R_{pA}$ vs rapidity: theory comparison

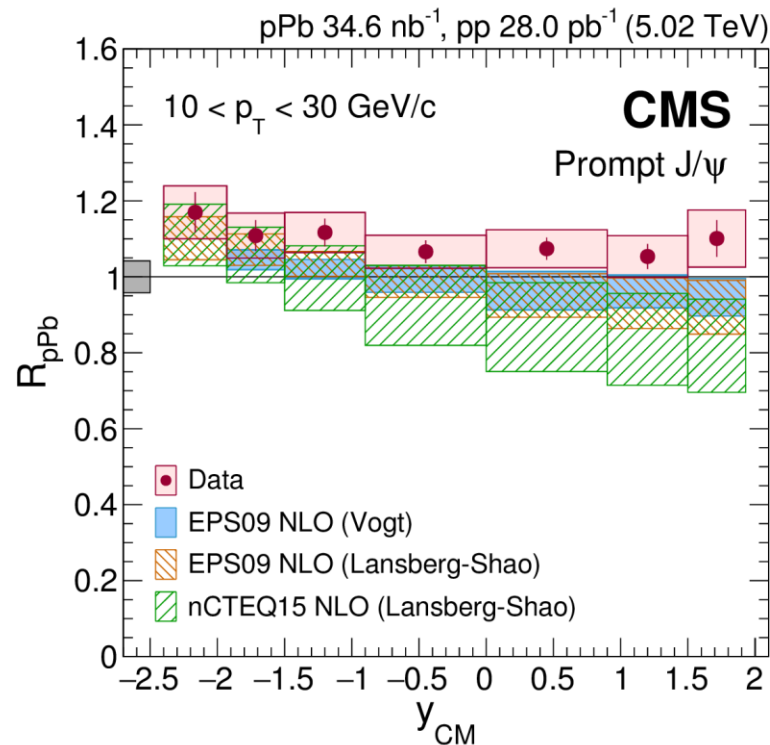
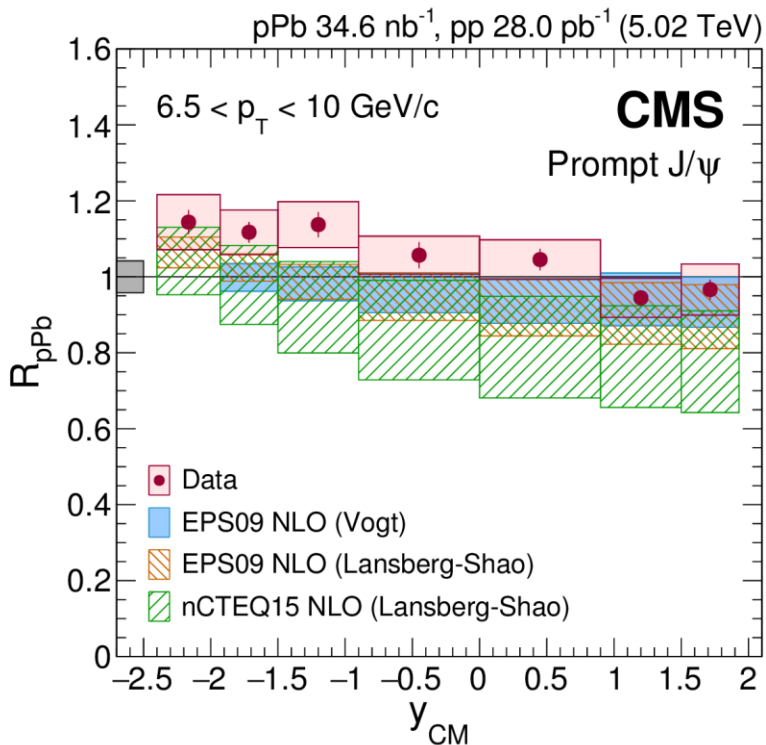


- ➡ Good agreement between ALICE and LHCb data
- ➡ Results described by models based on shadowing and/or energy loss

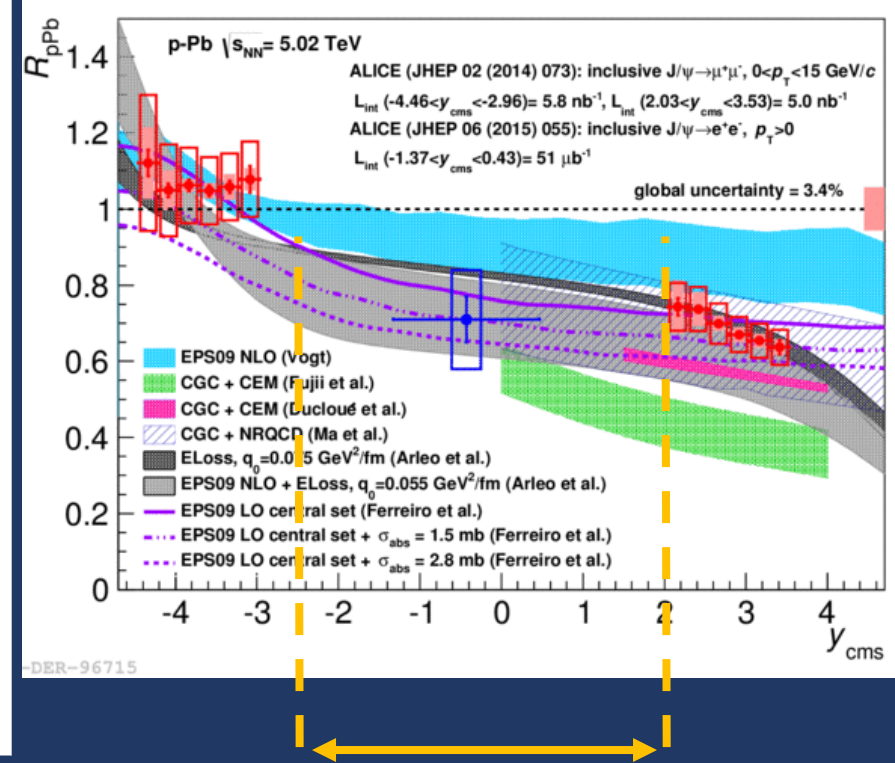
Size of theory uncertainties (mainly shadowing) still limits a more quantitative comparison

# Rapidity dependence of $J/\psi$ $R_{pA}$

CMS: high  $p_T$



ALICE: low  $p_T$

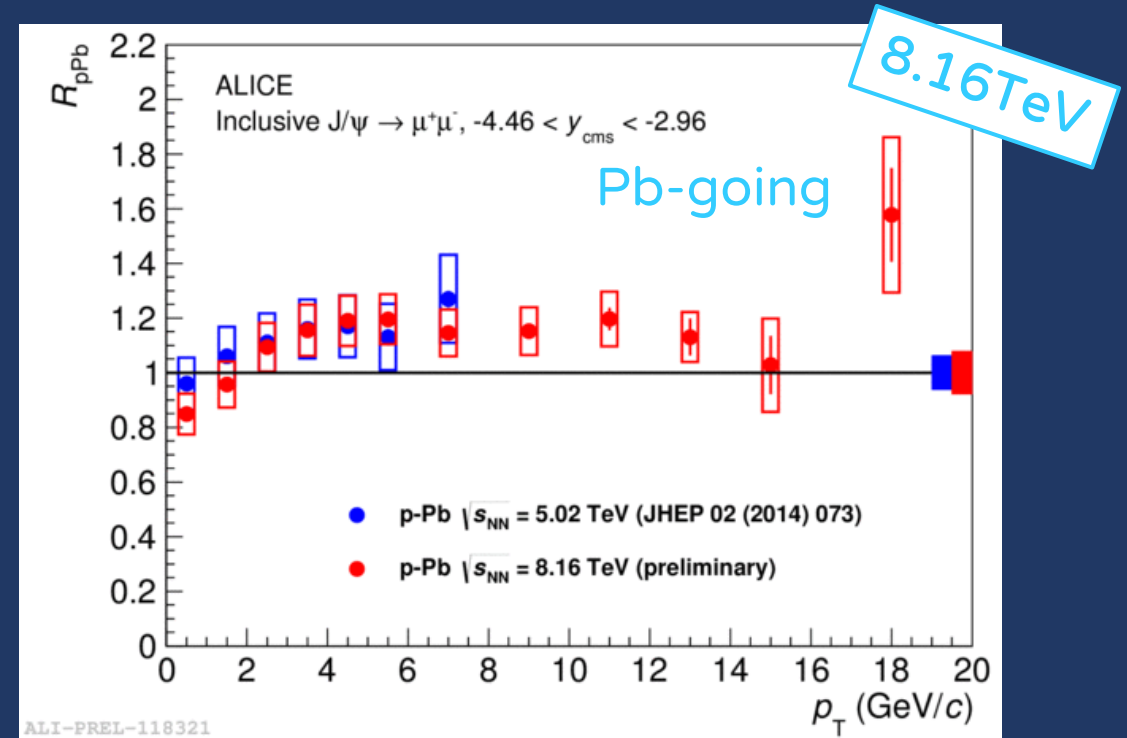
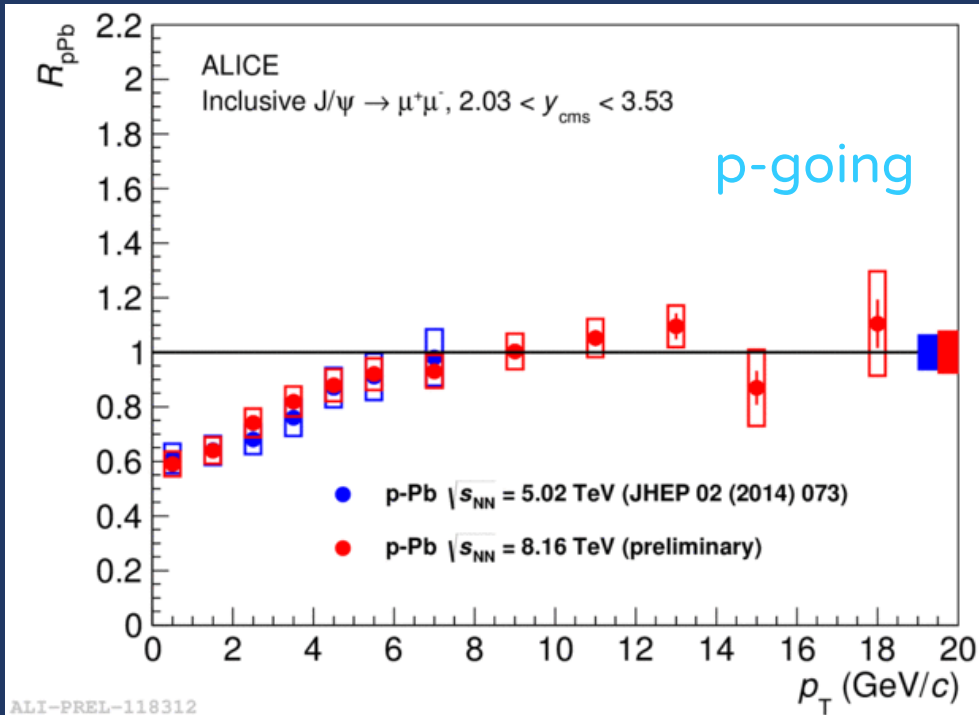


- ➔ Mid- $y$   $R_{pA}$  for high  $p_T$   $J/\psi$  is slightly higher than unity
- ➔ Shadowing implementations tend to underestimate the size of CNM effects

ALICE maximum  $p_T$  reach at mid  $y$  is 10 GeV/c  $\rightarrow R_{pA}$ , significantly smaller than the CMS one, reflects the  $p_T$  coverage



# $p_T$ dependence of $J/\psi$ $R_{pA}$

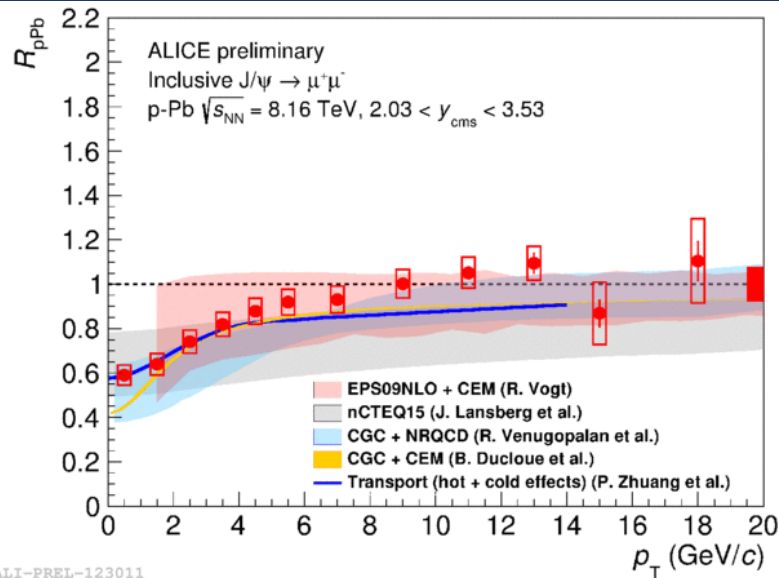


➔  $p_T$  coverage extended up to 20 GeV/c in Run2

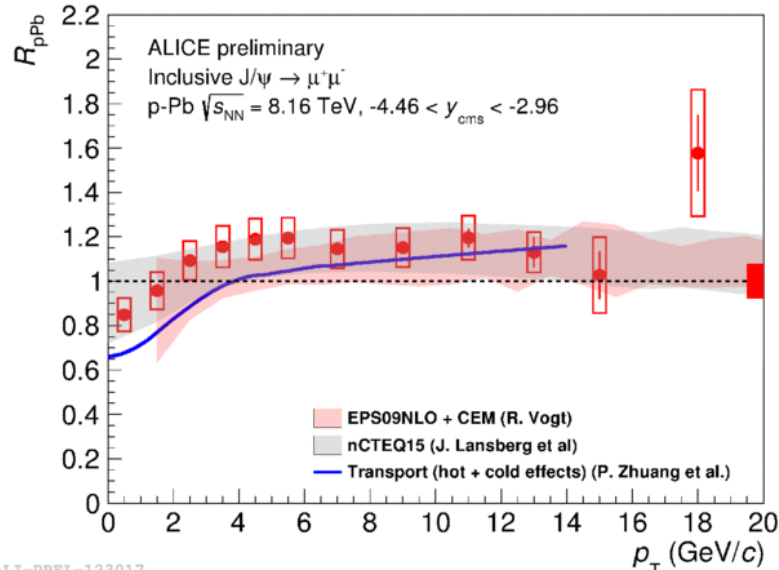
- p-going:  $R_{pA}$  increases with  $p_T$
- Pb-going:  $R_{pA}$  rather constant

➔ The strong  $J/\psi$  suppression observed in Pb-Pb data at high  $p_T$  cannot be due to CNM effects

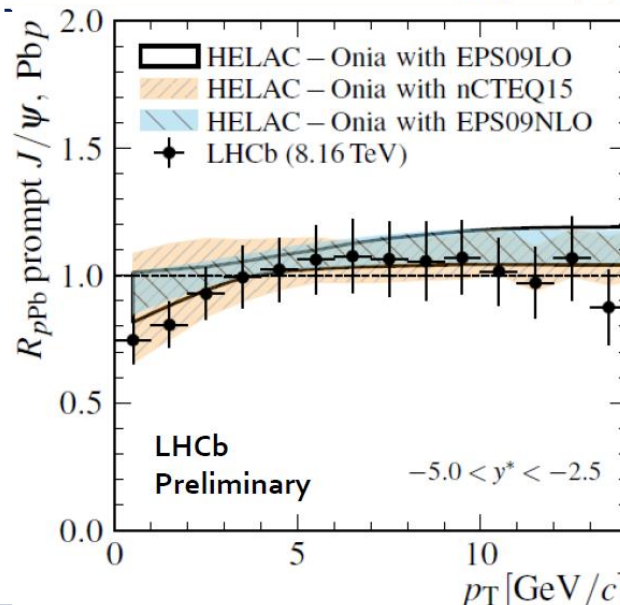
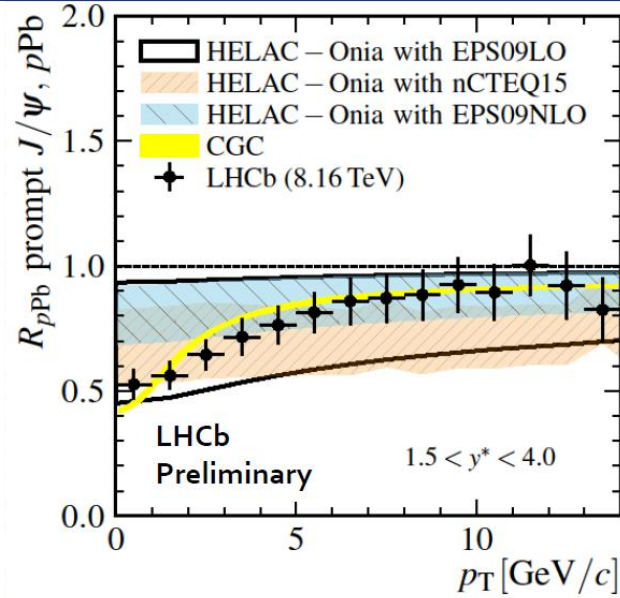
# $p_T$ dependence of $J/\psi R_{pA}$ : theory comparison



ALI-PREL-123011



ALI-PREL-123017

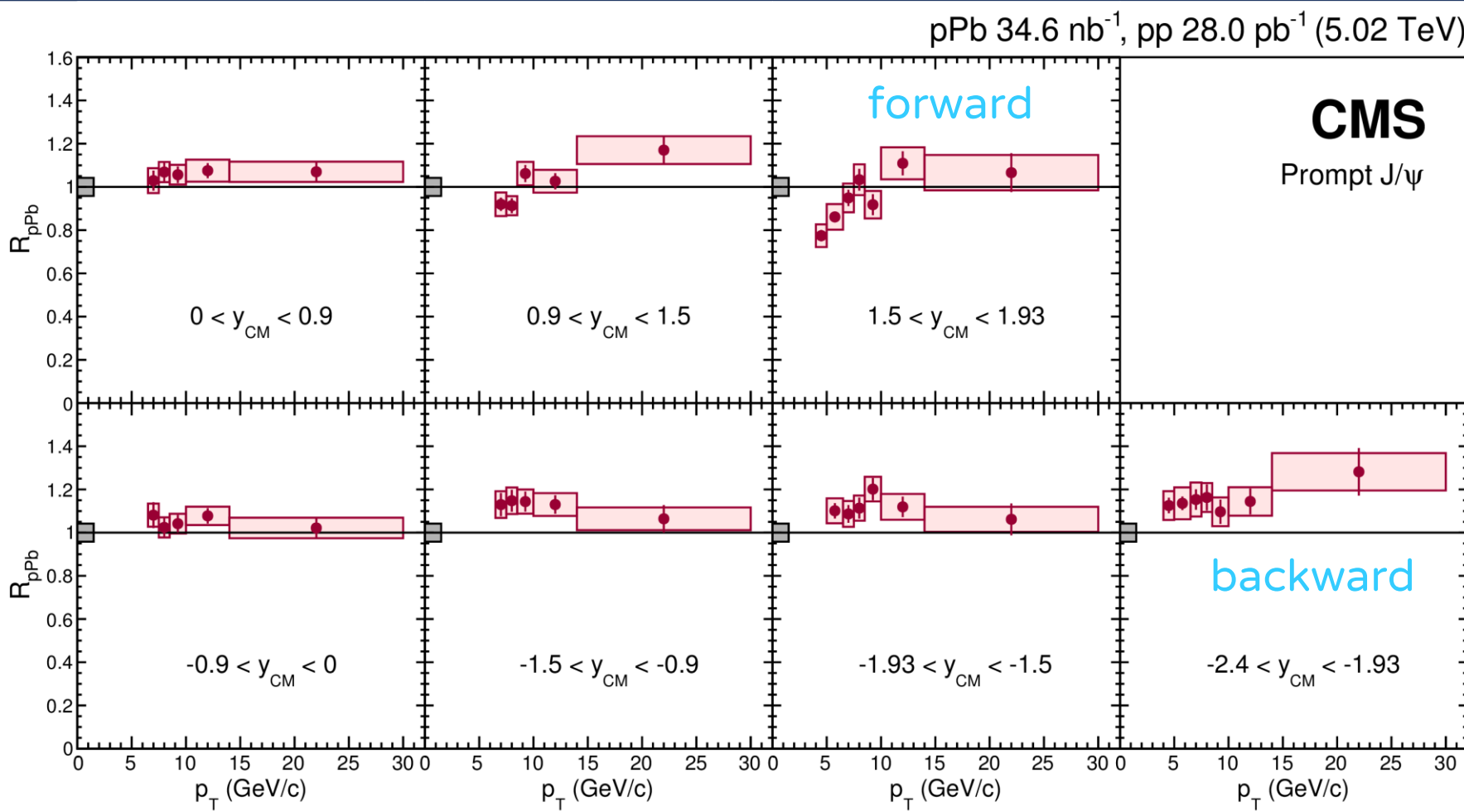


→ Slightly different  $y$  coverage in ALICE and LHCb, but rather similar  $p_T$  dependences

→ Shadowing and energy loss models describe the  $R_{pA}$  trend vs  $p_T$

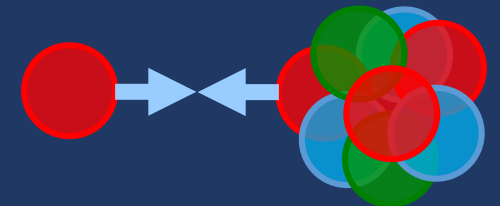


# $J/\psi R_{pA}$ at high $p_T$



$R_{pA}$  of high  $p_T$  prompt  $J/\psi$  shows

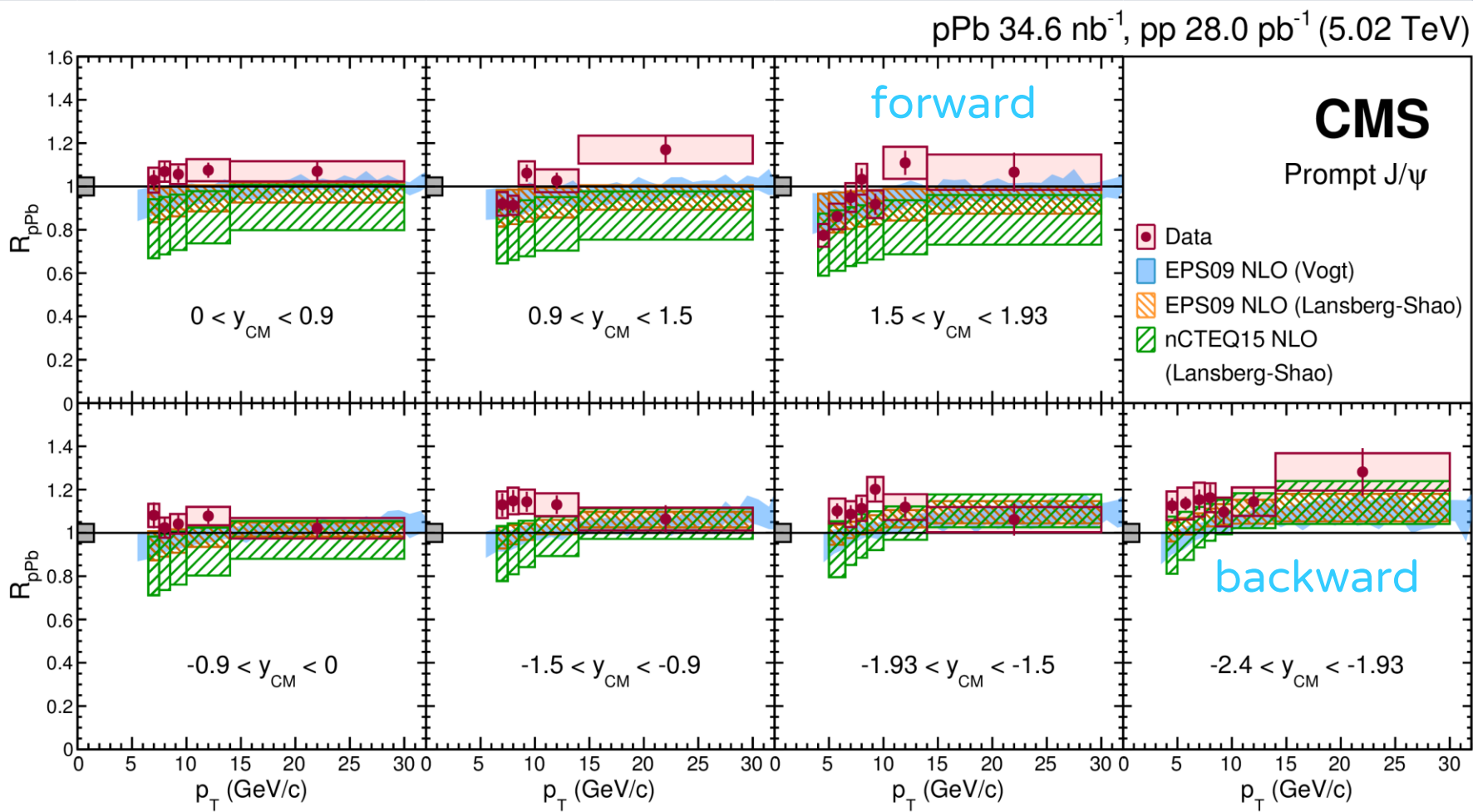
- values slightly higher than unity at mid and backward rapidity
- hint for stronger CNM effects at the edges of the  $y$  domain



forward

backward

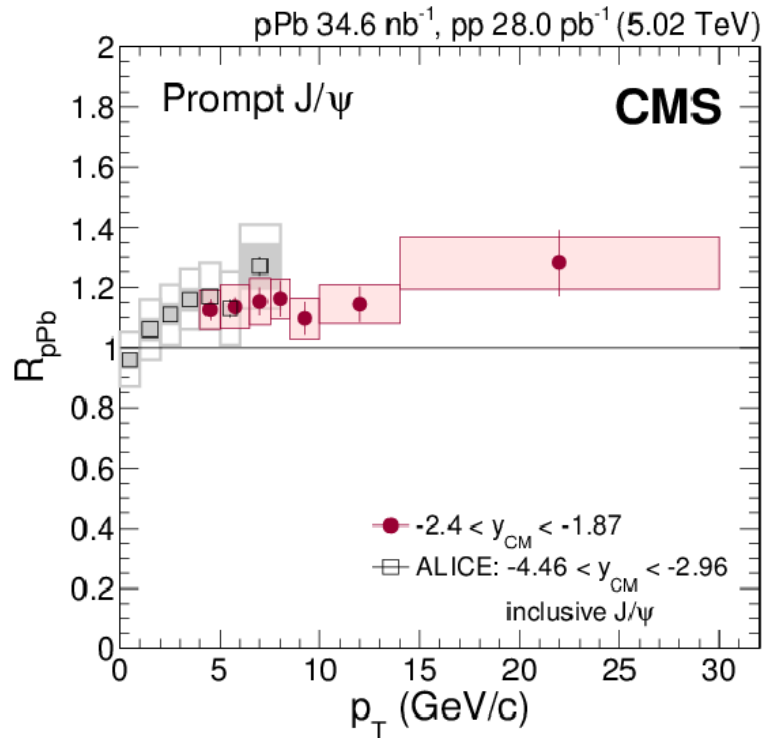
# $J/\psi R_{pA}$ at high $p_T$



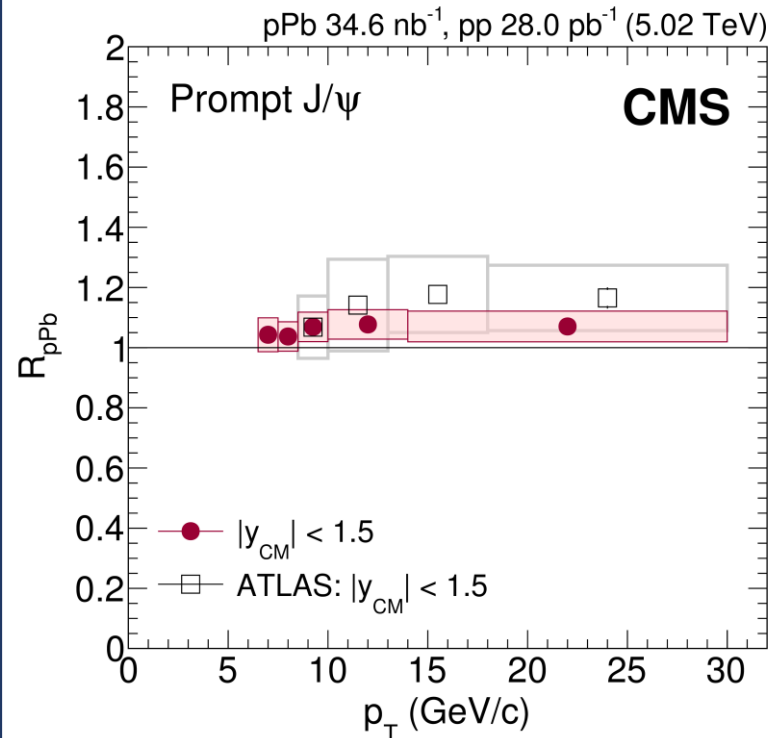
➔ Different shadowing implementations describe the data trend (even if slightly at the lower edge)

# Comparison among experiments

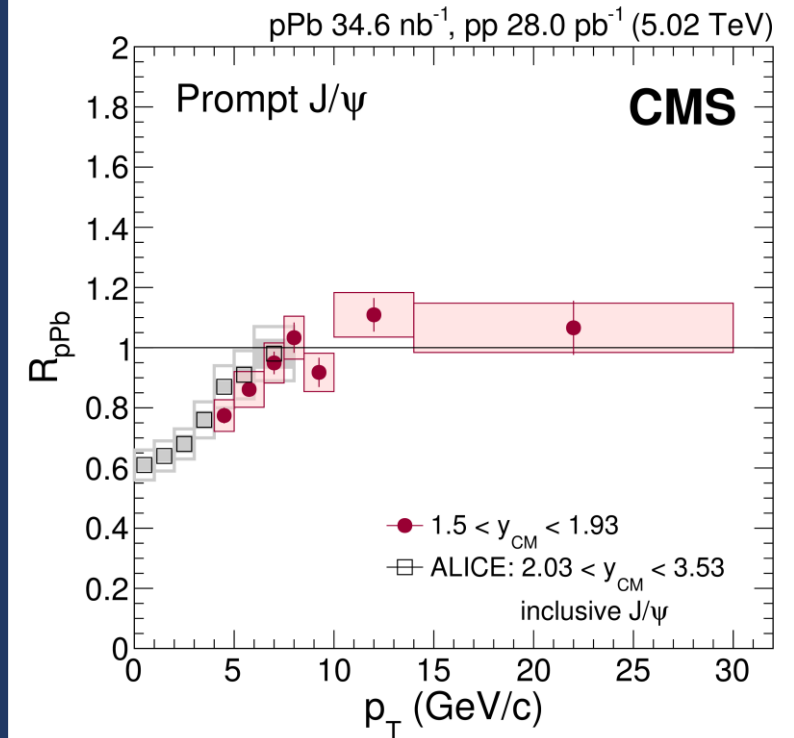
## Backward-y



## Mid-y



## Forward-y



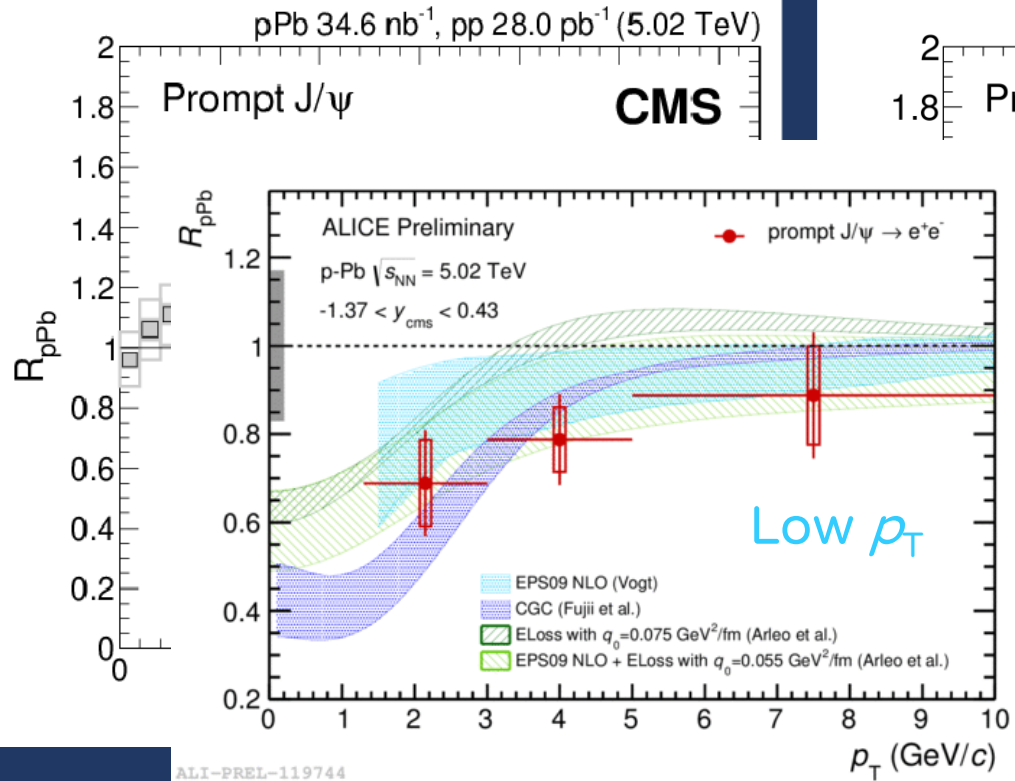
➔ Compilation of results from different experiments shows

- Good compatibility in close rapidity ranges
- Broad  $p_T$  coverage from 0 to 30 GeV/c

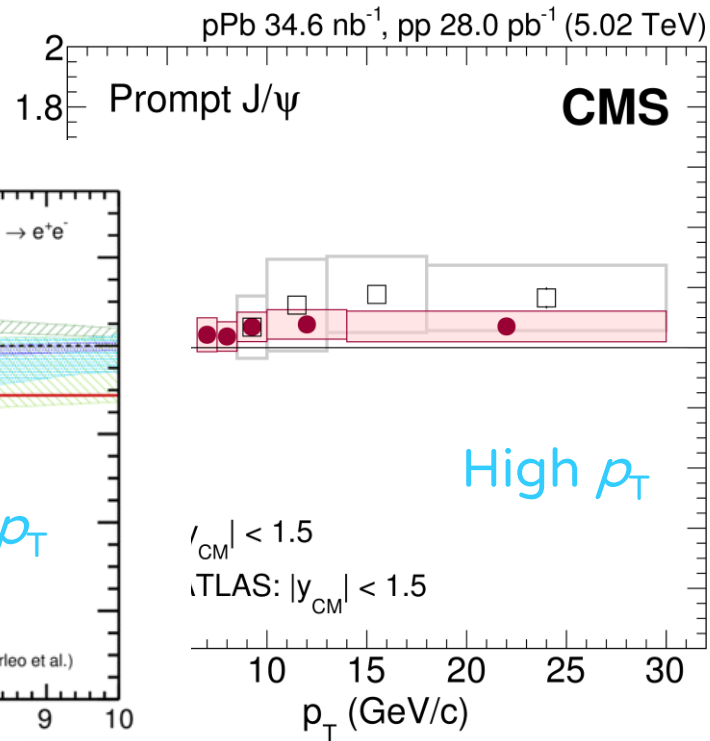
➔ Pattern confirms strongest CNM effects at low  $p_T$  and forward  $y$

# Comparison among experiments

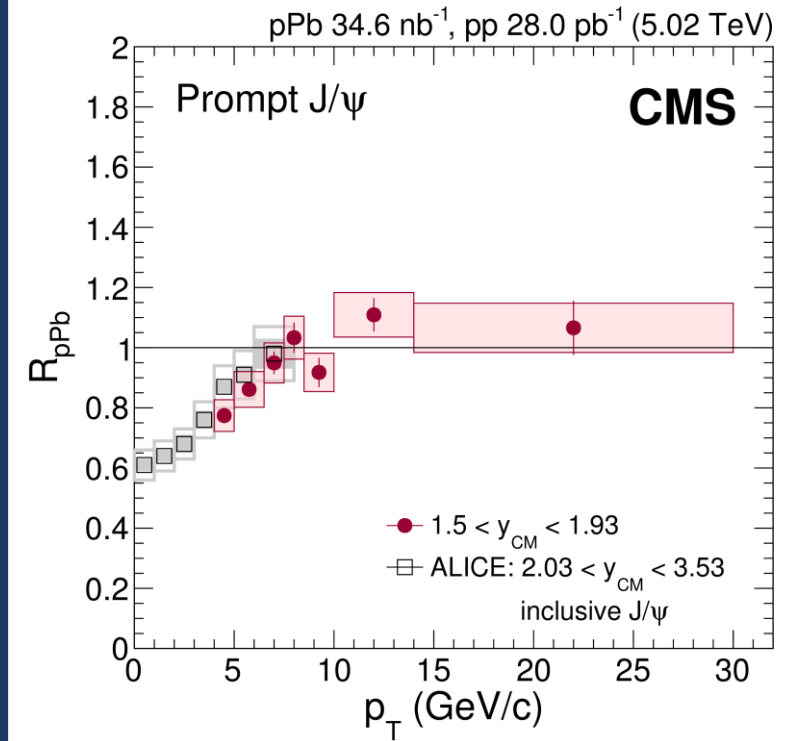
Backward-y



Mid-y

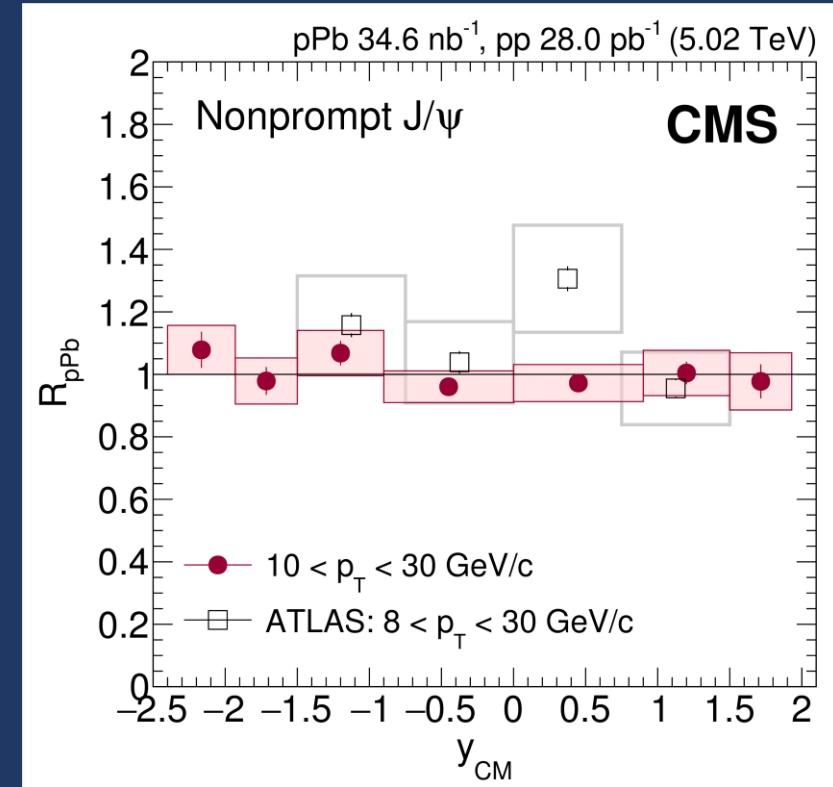
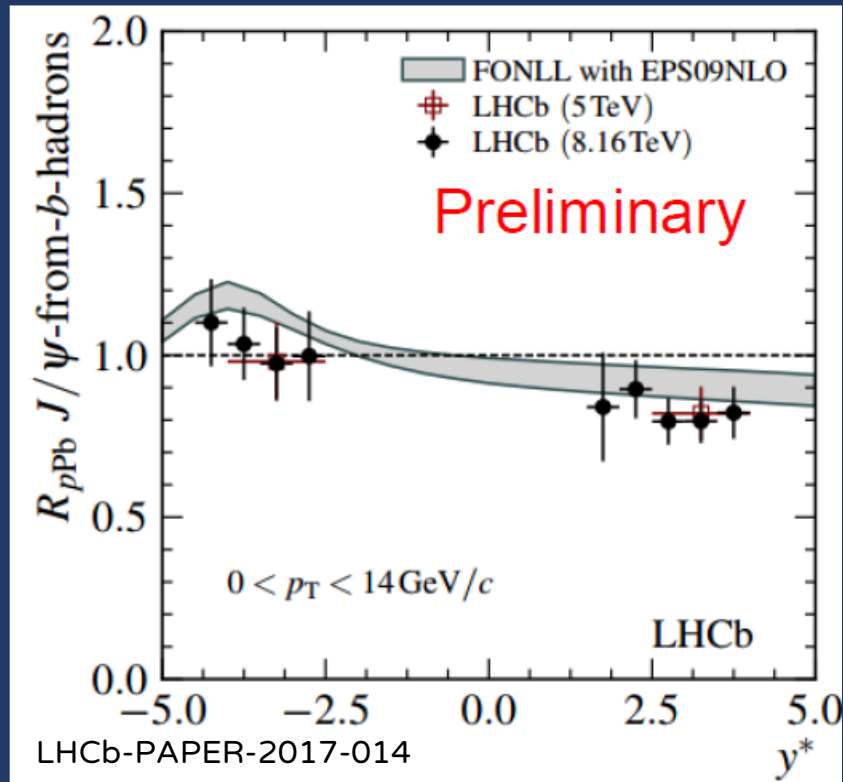


Forward-y



➔ Same decreasing trend towards low  $p_T$  observed also at mid-rapidity

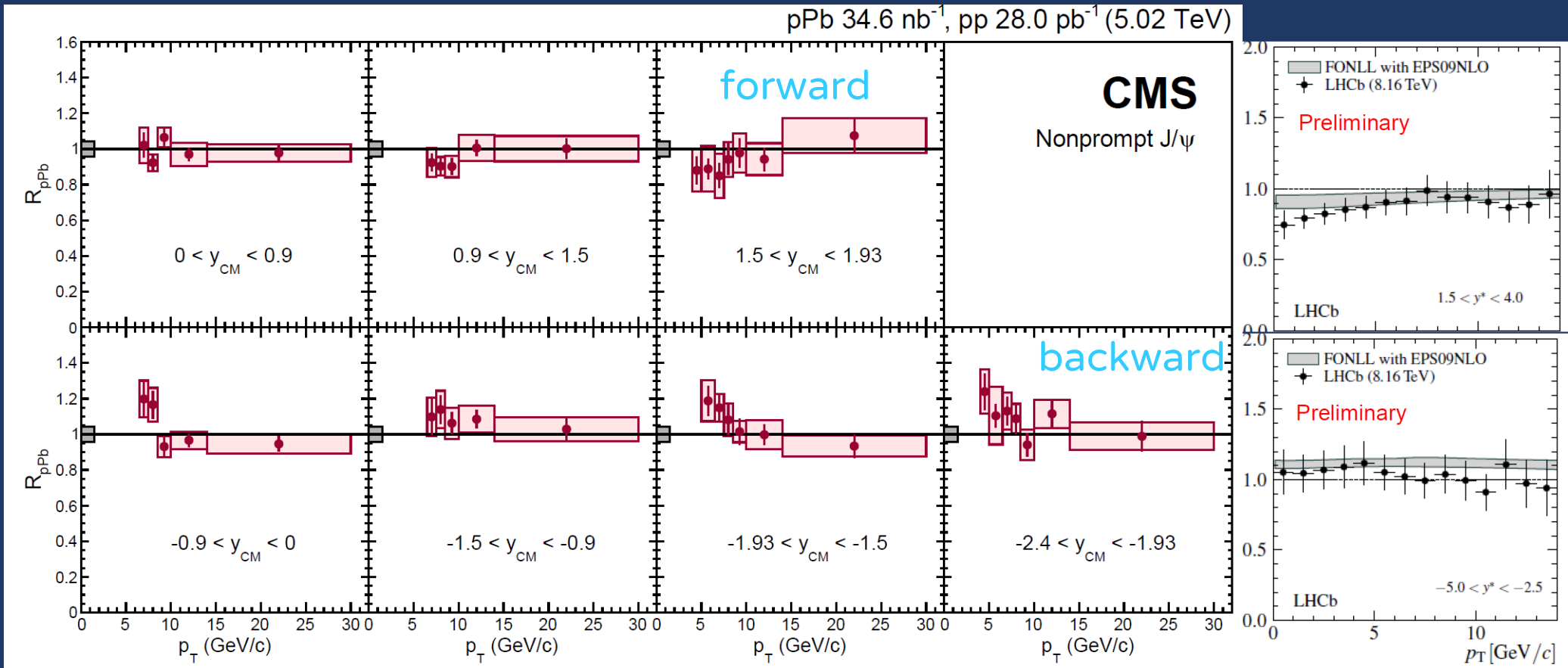
# $R_{pA}$ of non-prompt $J/\psi$



- ➔ Agreement between 5.02 and 8.16 TeV results
- ➔ Small CNM effects on non-prompt  $J/\psi$
- ➔ Overall agreement with FONLL+EPS09NLO

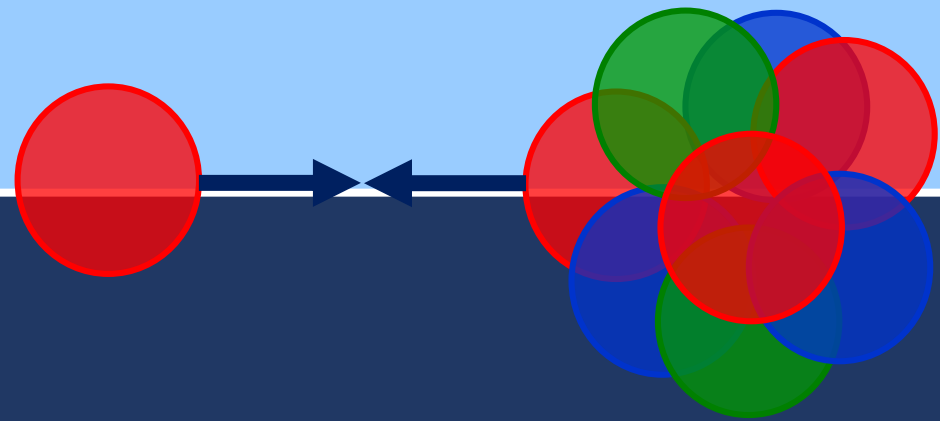
- ➔ Complementary  $y$  range covered by ATLAS, CMS and LHCb (but a different  $p_T$  range)
- ➔ Consistent ATLAS and CMS results
- ➔ no  $y$  dependence for high  $p_T$  non-prompt  $J/\psi$

# $R_{pA}$ of non-prompt $J/\psi$



- ➔ complementary LHCb and CMS results
- ➔ high  $p_T$  non-prompt  $J/\psi$  show negligible CNM effects
- ➔ weak  $p_T$  dependence, significant only in the LHCb domain, where it reaches 30%

p-A collisions:  
 $\psi(2S)$

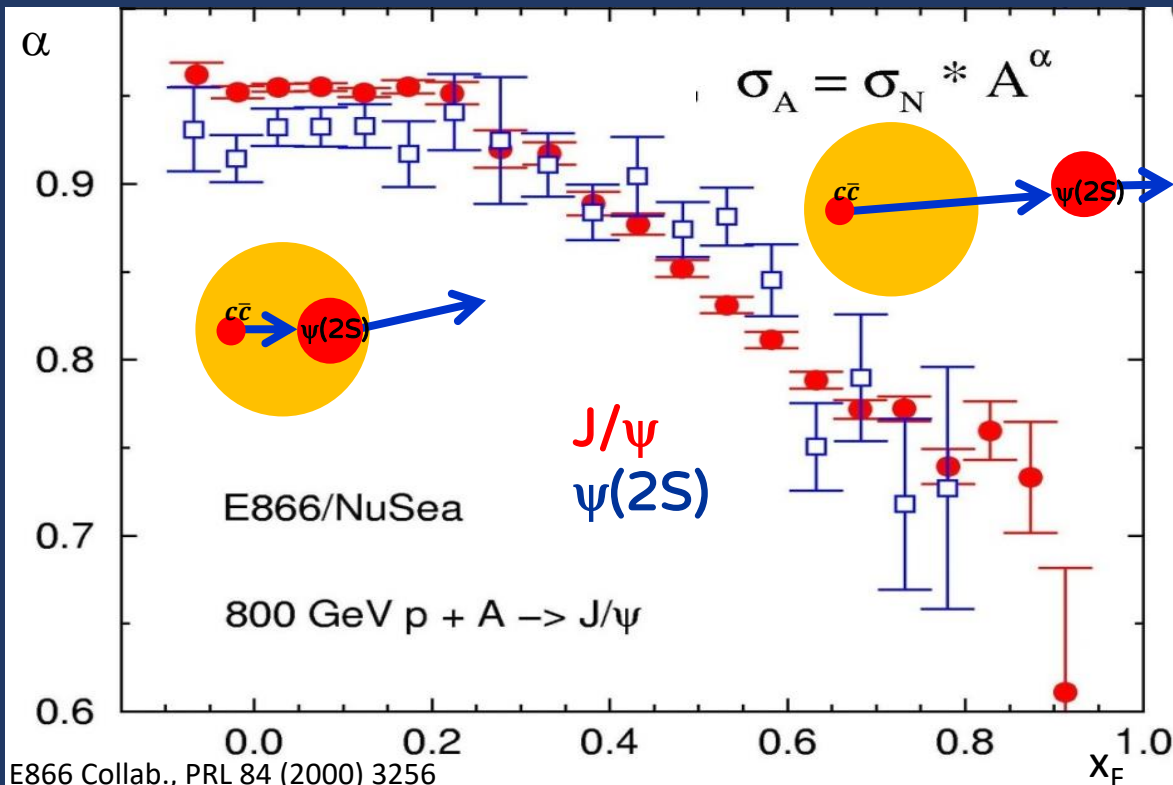




# $\psi(2S)$ as a function of $x_F$

Being more weakly bound than the  $J/\psi$ , the  $\psi(2S)$  is an interesting probe to have further insight on the charmonium behaviour in pA

$\psi(2S)$  production is modified by CNM effects depending on its kinematic



**mid- $y$  ( $x_F \sim 0$ ):**  $\psi(2S)$  suppression stronger than  $J/\psi$  one,  $\rightarrow$  break-up of fully formed resonance traversing the nucleus

charmonium formation time < crossing time

**fw- $y$  (high  $x_F$ ):** suppression becomes roughly identical  $\rightarrow$  dominated by energy loss

charmonium formation time > crossing time

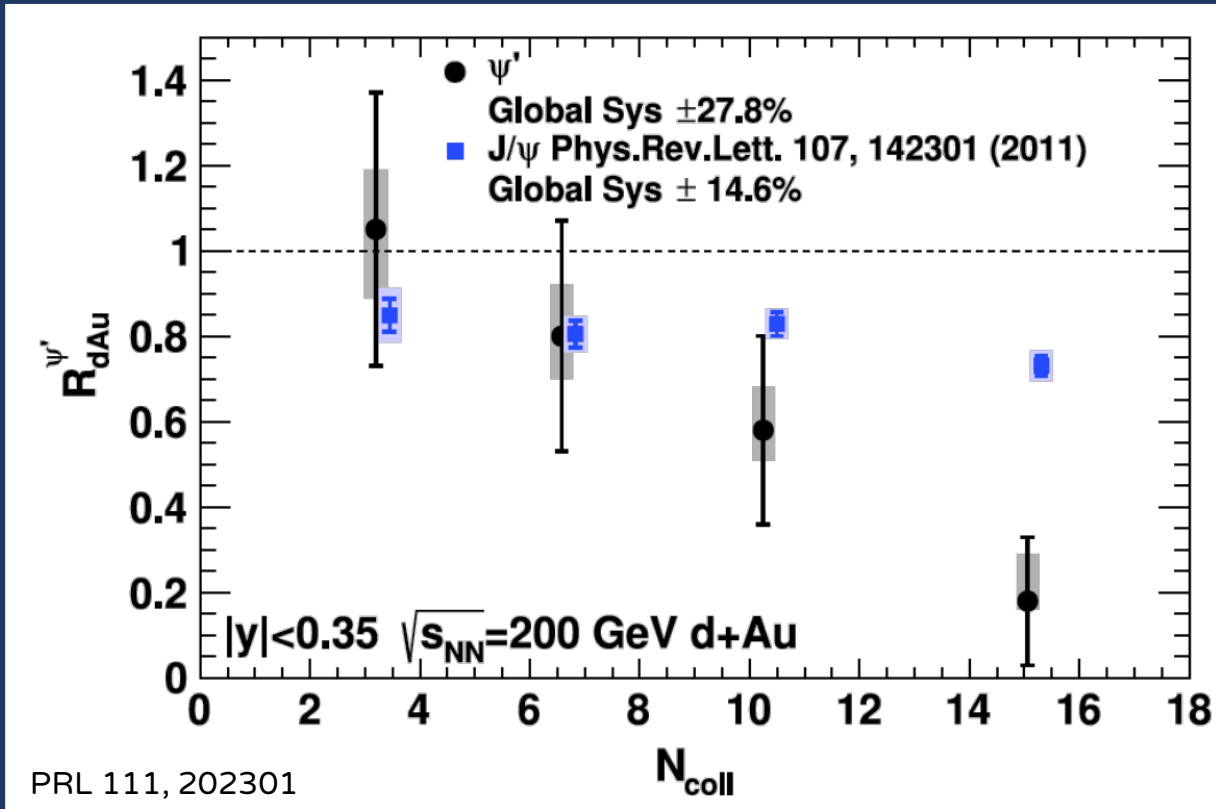
Experiment	$\sqrt{s}$ (GeV)	$\tau$ (fm/c)
NA50/60	17-27	0.29-0.34
E866	38.8	0.28-0.02
HERA-B	41.6	0.28-0.003

$$\tau_c = \frac{\langle L \rangle}{(\beta_z \gamma)}$$



# $\psi(2S)$ in pA collisions

→ at RHIC and LHC energies,  $\psi(2S)$  suppression stronger than the  $J/\psi$  one



→ unexpected because time spent by the  $c\bar{c}$  pair in the nucleus ( $\tau_c$ ) is shorter than charmonium formation time ( $\tau_f$ )  
→ break up of the fully formed resonance in the nuclear medium should not play a role

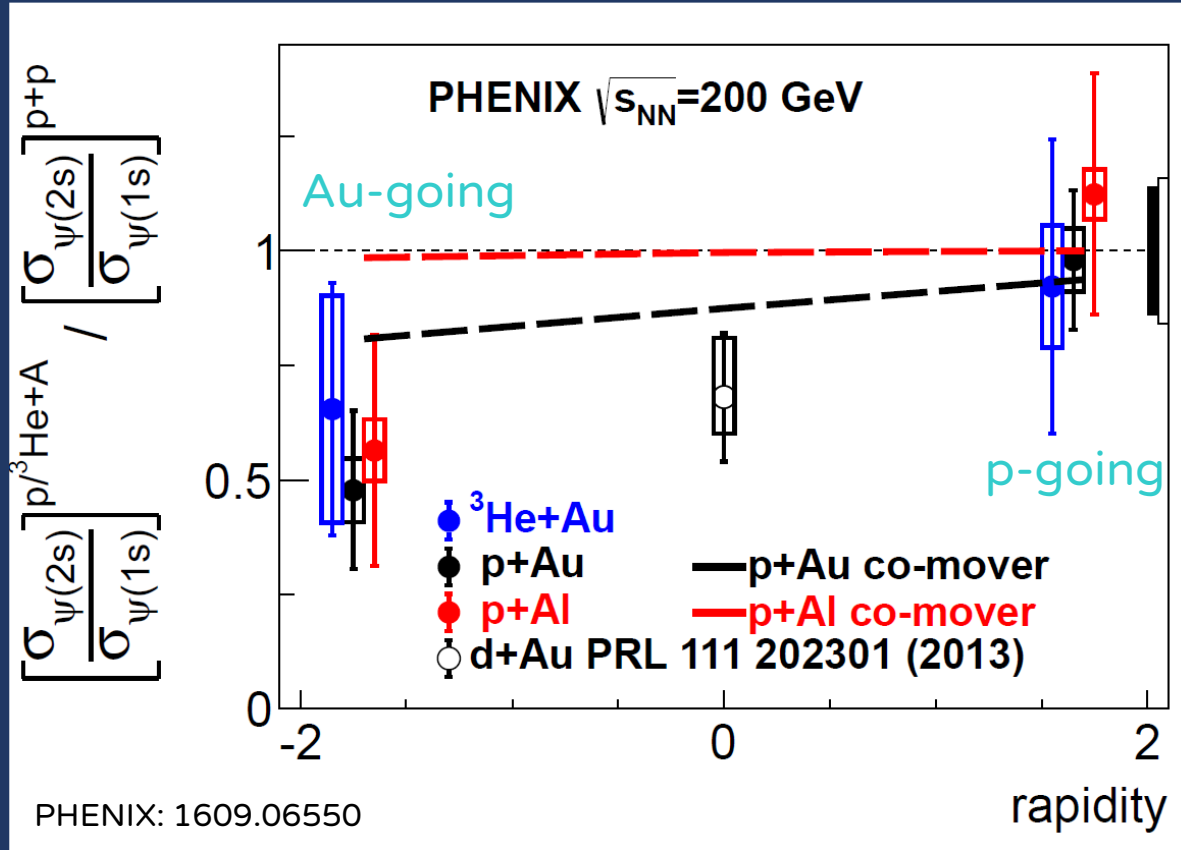
	$\sqrt{s}$ (GeV)	$\tau$ (fm/c)
PHENIX	200	fw-y: 0.0035 bck-y: 0.28
ALICE	2760	fw-y: $10^{-4}$ - $10^{-5}$ bck-y: 3-7 $10^{-2}$

McGlinchey, Frawley, Vogt PRC 87 054910

→ shadowing and energy loss, almost identical for  $J/\psi$  and  $\psi(2S)$ , do not account for the different suppression

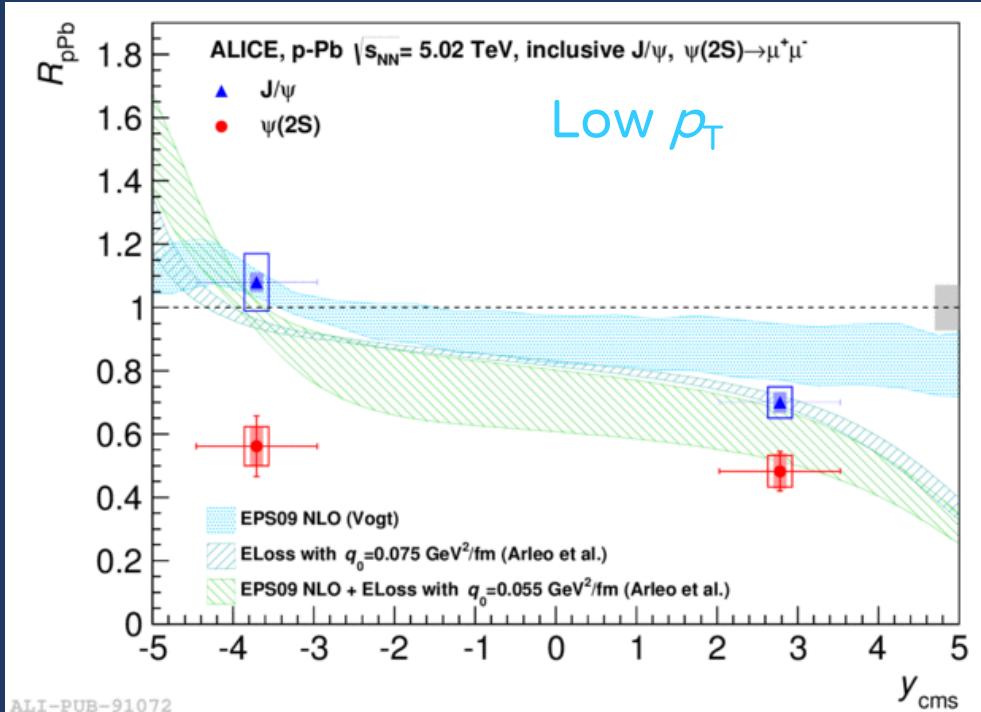
# $\psi(2S)$ in pA collisions at RHIC

➔ at RHIC and LHC energies,  $\psi(2S)$  suppression stronger than the  $J/\psi$  one

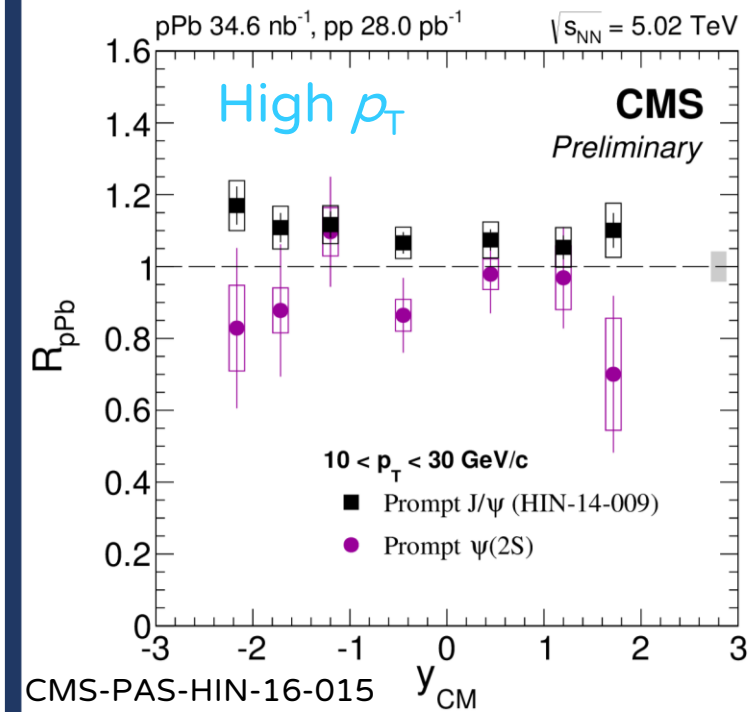
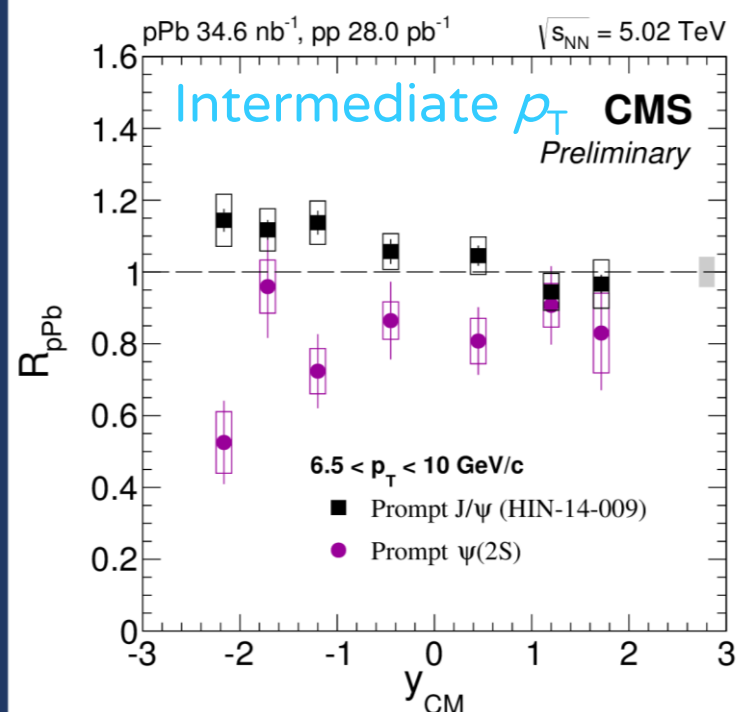


- ➔ Suppression is more important at backward-y
- ➔ Final state effects needed to explain the behaviour

# $\psi(2S)$ RpA at LHC



ALI-PUB-91072



➔ Stronger  $\psi(2S)$  suppression with respect to the  $J/\psi$  one, mainly:

- at backward rapidity
- at the low(est)  $p_T$

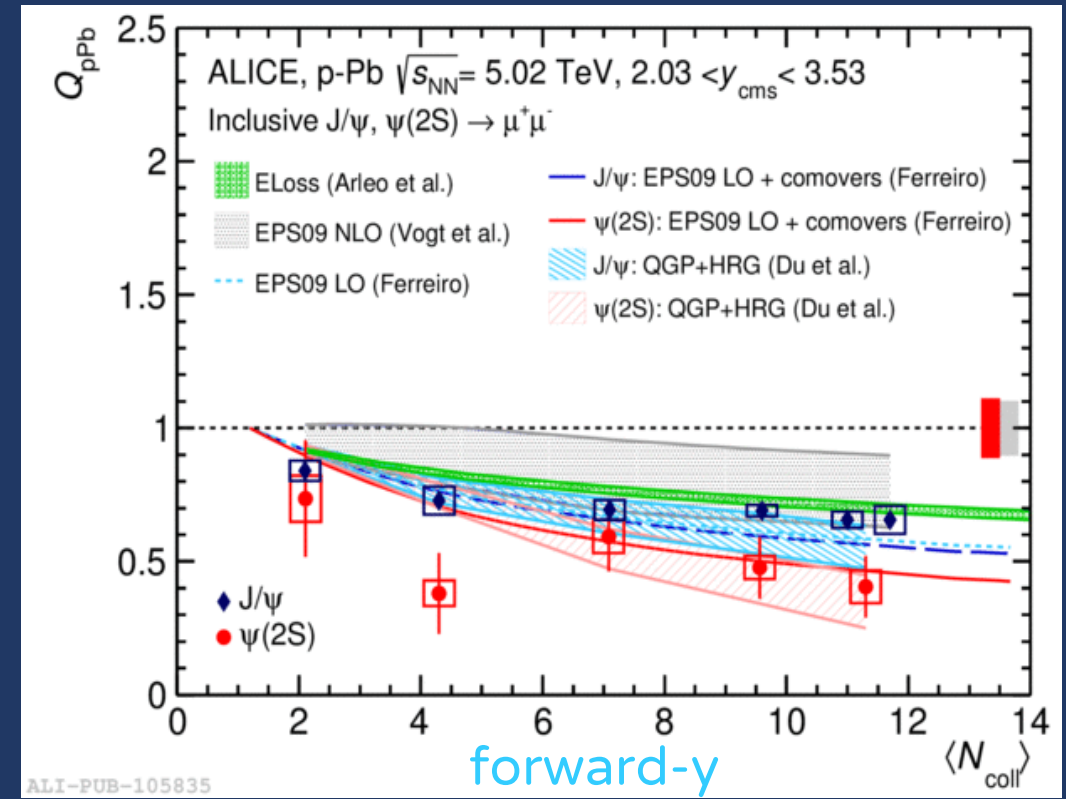
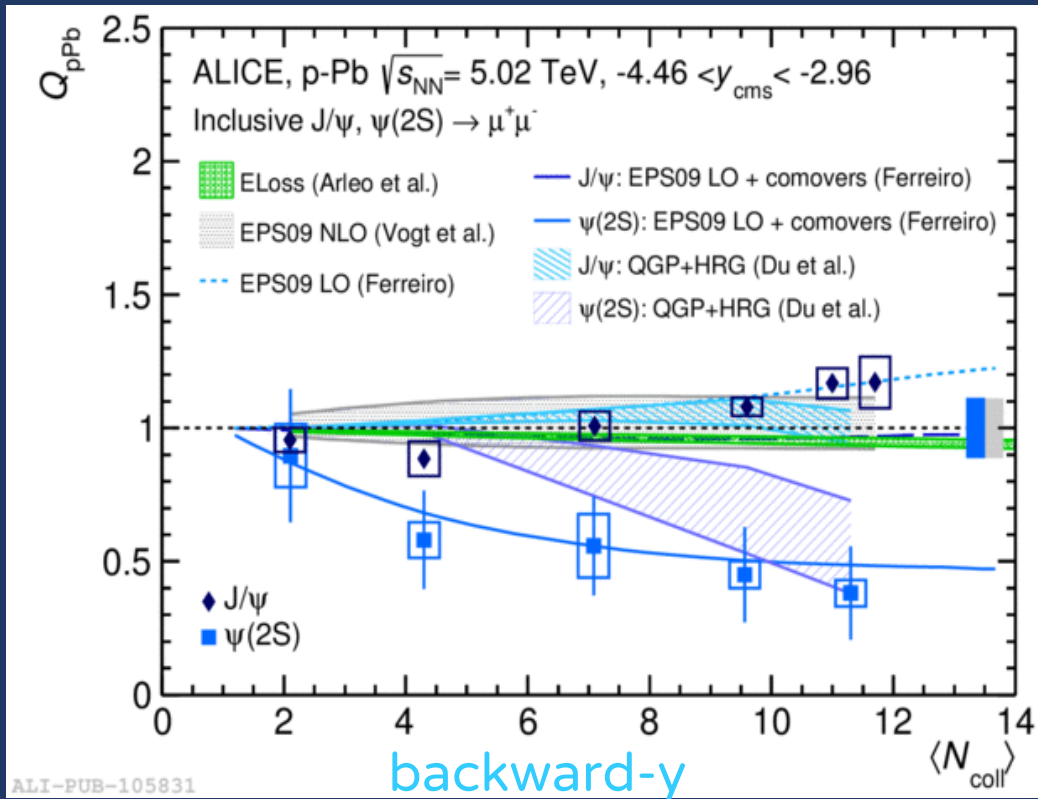
➔ Backward- $y$ : size of  $\psi(2S)$  suppression rather similar between ALICE and CMS

➔ Forward- $y$ : suppression is more important in the ALICE  $p_T$  range

➔ Comparison with models in the high  $p_T$  range would be interesting!

# $\psi(2S)$ in pA collisions at LHC

➔ Suppression is stronger in central collisions

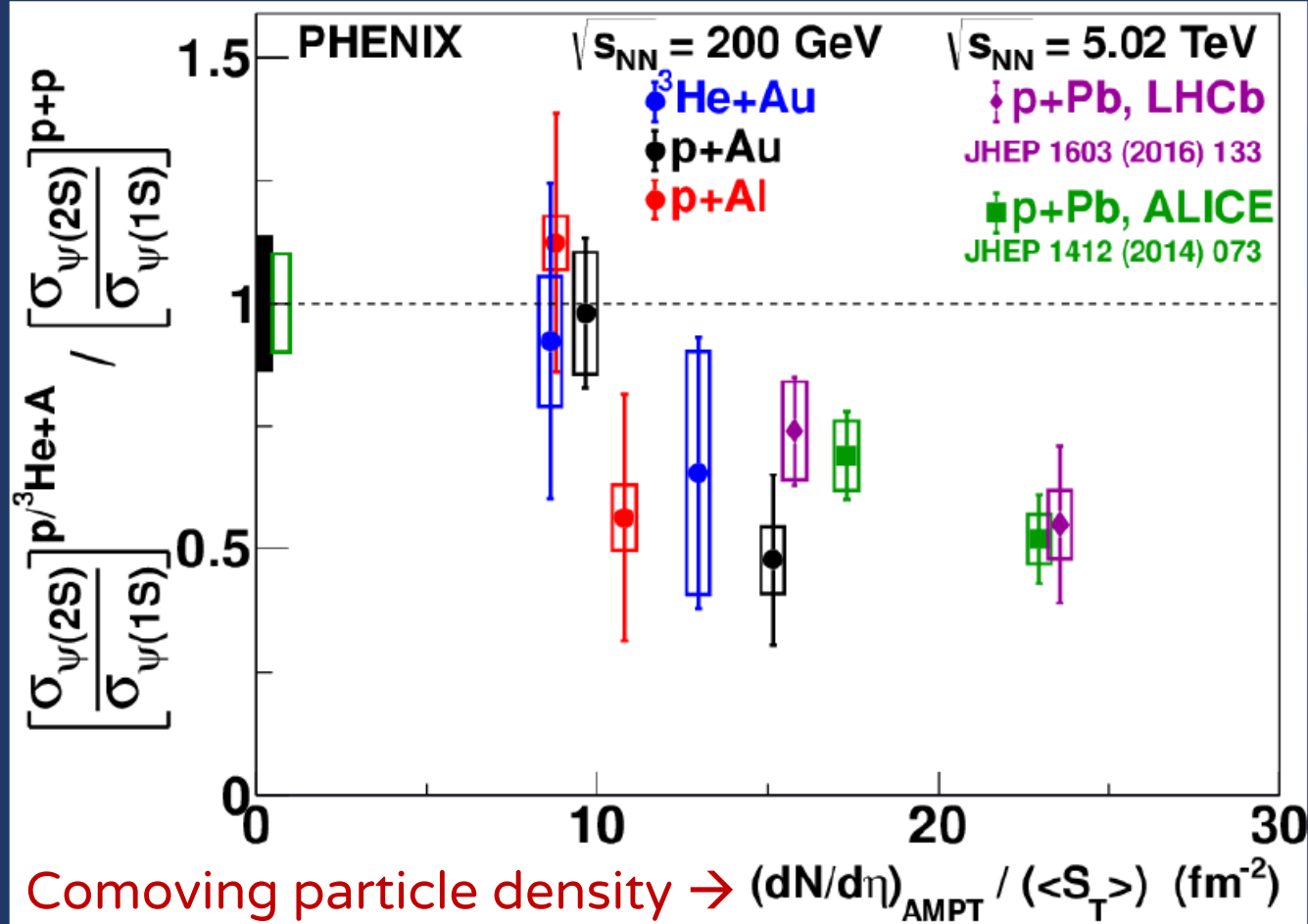


➔ Clear evidence for a stronger  $\psi(2S)$  suppression, wrt  $J/\psi$ , at backward rapidity

➔ QGP+hadron resonance gas (Rapp) or comover (Ferreiro) models describe the stronger  $\psi(2S)$  suppression

# $\psi(2S)$ in pA at RHIC and LHC

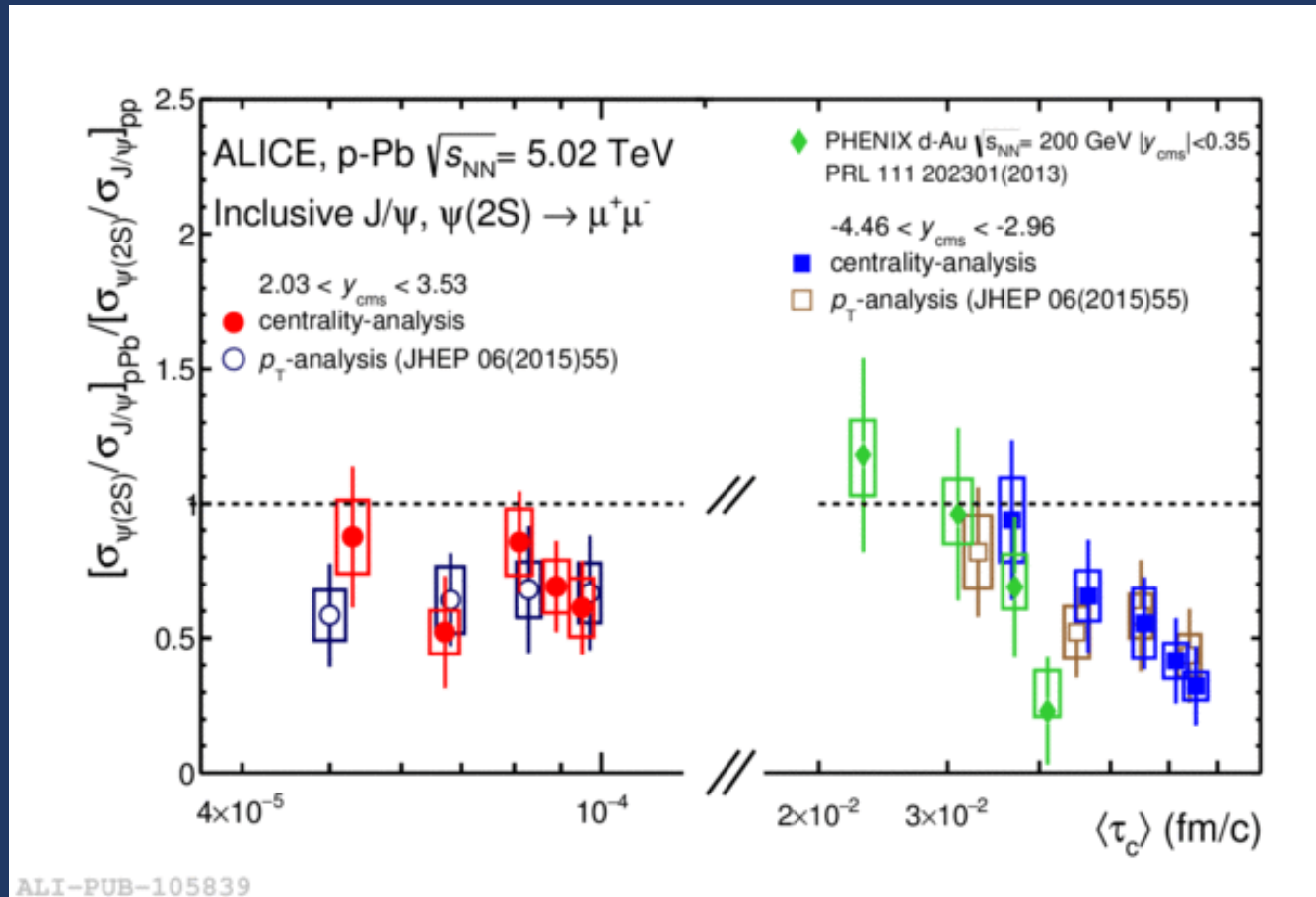
Under the assumption that the  $\psi(2S)$  is suppressed by final state effects, RHIC and LHC results can be compared in terms of comoving particle densities



- ➔ Largest comoving particle density reached at LHC, at backward-y
- ➔ Backward-y RHIC data reach a similar comoving particle density as forward LHC
- ➔ Double ratio decreases, increasing the comoving particle density  
→ Consistent with the  $\psi(2S)$  break-up in final state interactions

# $\psi(2S)$ in pA at RHIC and LHC

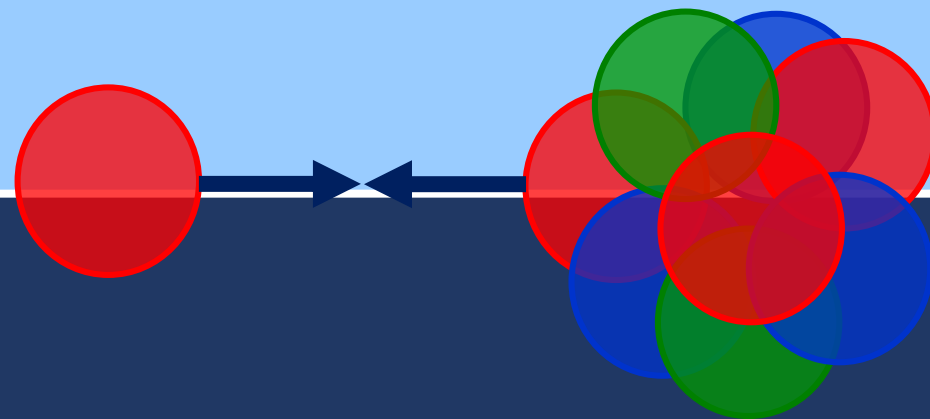
Relative  $J/\psi$  and  $\psi(2S)$  suppression is studied as a function of the crossing times



- ➔ Within uncertainties a scaling between RHIC and LHC is observed
- ➔ At backward-y, where the largest  $\tau_c$  are reached, a decreasing trend is observed
  - ➔ Even if the quarkonium formation time is larger than  $\tau_c$ , is a fraction of the cc pairs hadronizing inside the nucleus?
- ➔ At forward-y,  $\tau_c \ll \tau_f$  by 2-3 order of magnitude.
  - ➔ In principle no final state effects related to cold nuclear matter can play a role, only comover interaction

p-A collisions:

$\Upsilon$

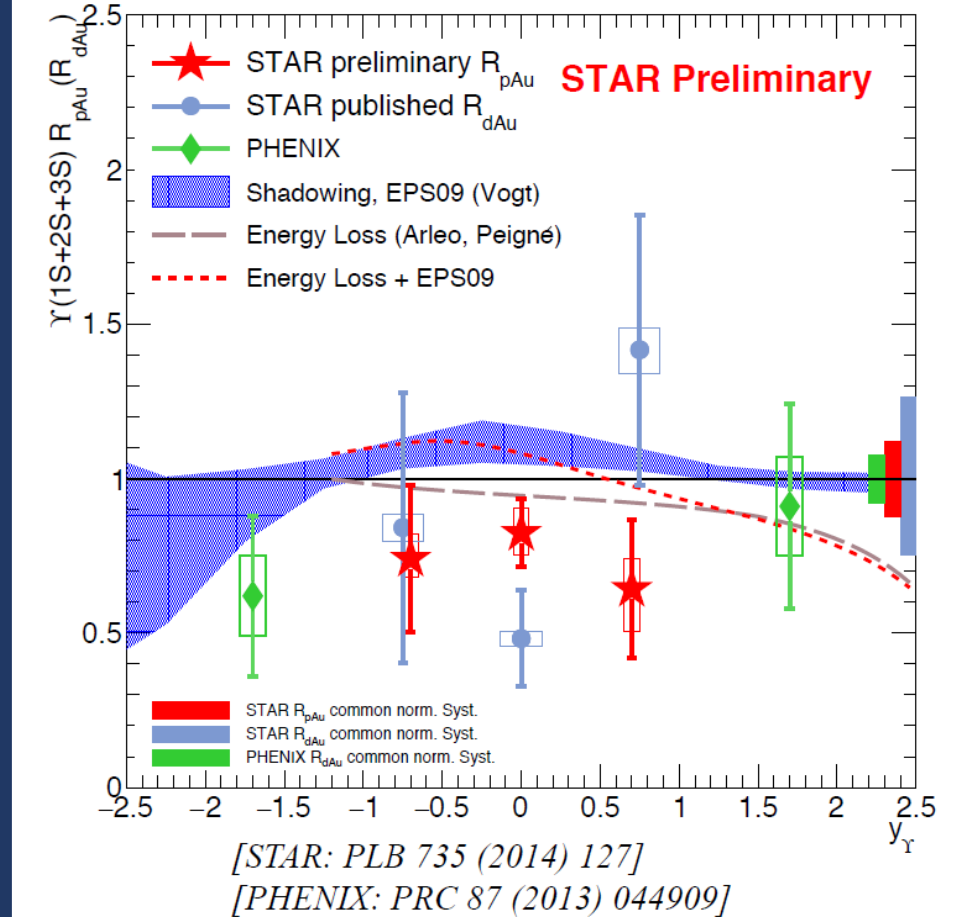
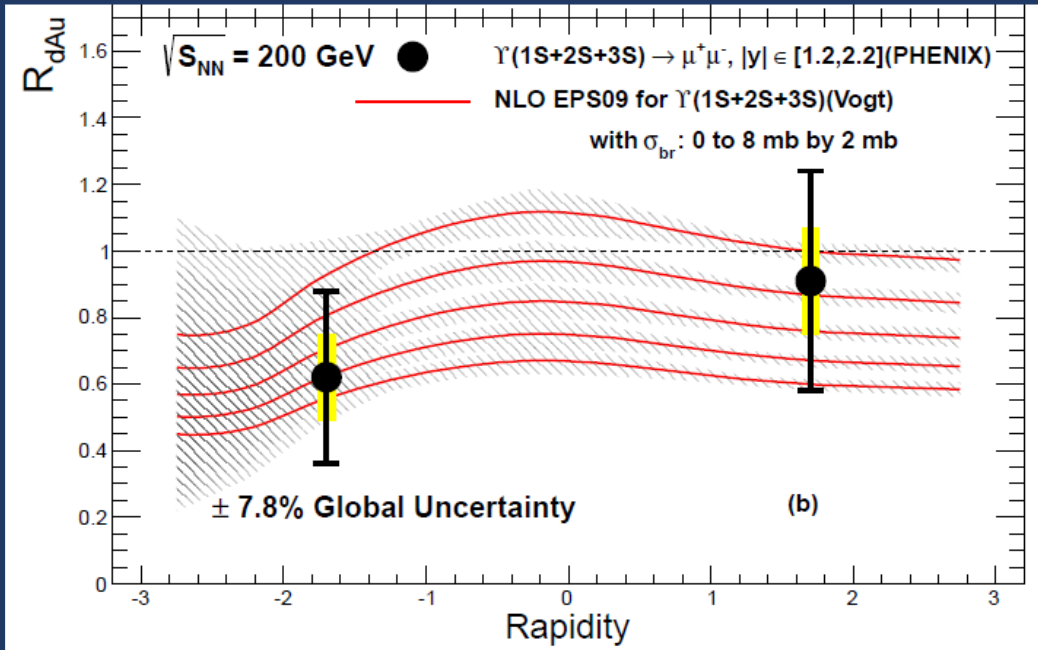




# $\Upsilon$ in pA and dA at RHIC

➔  $\Upsilon(1S+2S+3S)$  measured in both pA and dA

- Large uncertainties for  $R_{dAu}$  prevent a clear understanding of the  $\Upsilon$   $y$ -evolution
- Need for bb break-up in nuclear matter not obvious



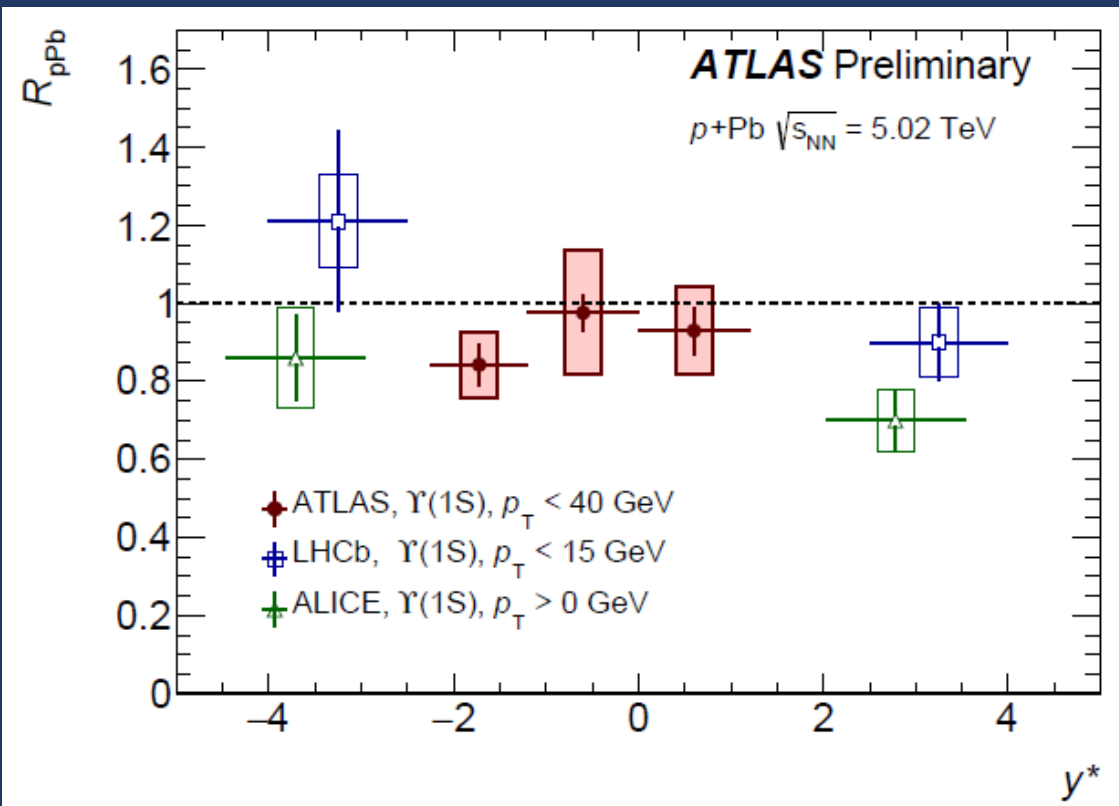
➔ New STAR  $R_{pA}$  ( $\sim 0.8$ ), with improved precision wrt  $R_{dAu}$ , suggest cold nuclear matter effects on  $\Upsilon(1S+2S+3S)$  at mid- $y$



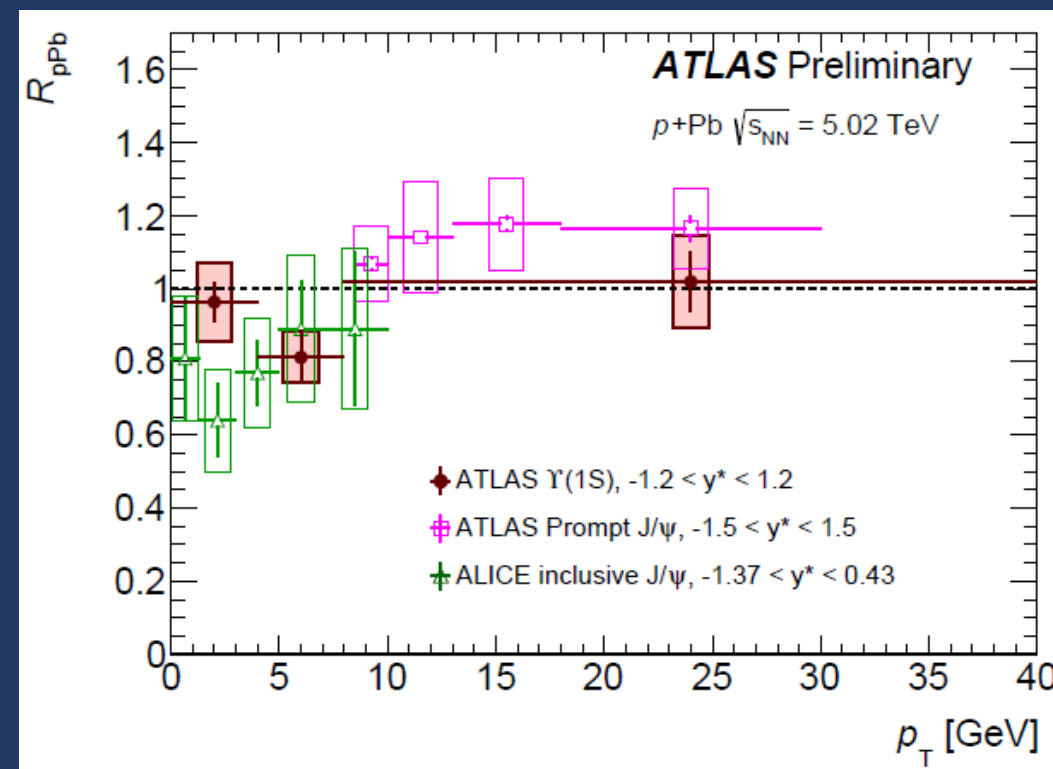
# RpA $\Upsilon(1S)$ at LHC

Wide  $y$  coverage explored by ALICE, ATLAS and CMS

ALICE and LHCb compatible within uncertainties, but LHCb values systematically larger

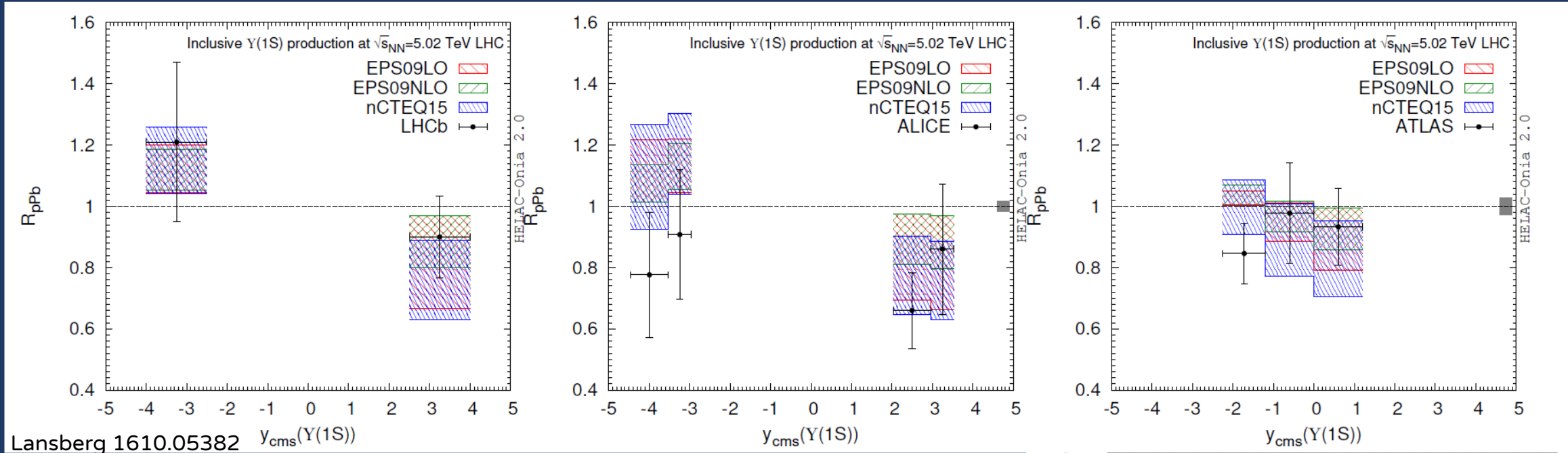


Hint for a stronger  $\Upsilon(1S)$  suppression at forward- $y$  in ALICE data, while mid- and backward- $y$  is compatible with no modification



$p_T$  dependence not conclusive so far

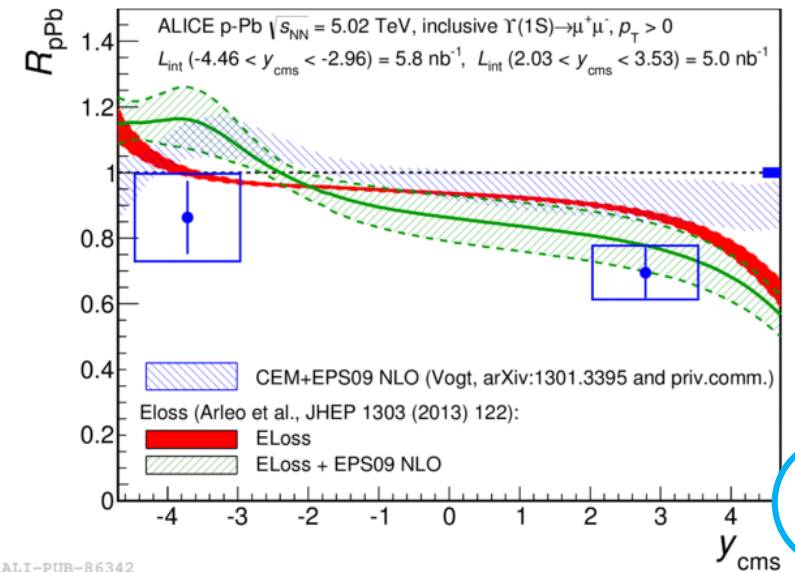
# $R_{pA} \Upsilon(1S)$ at LHC



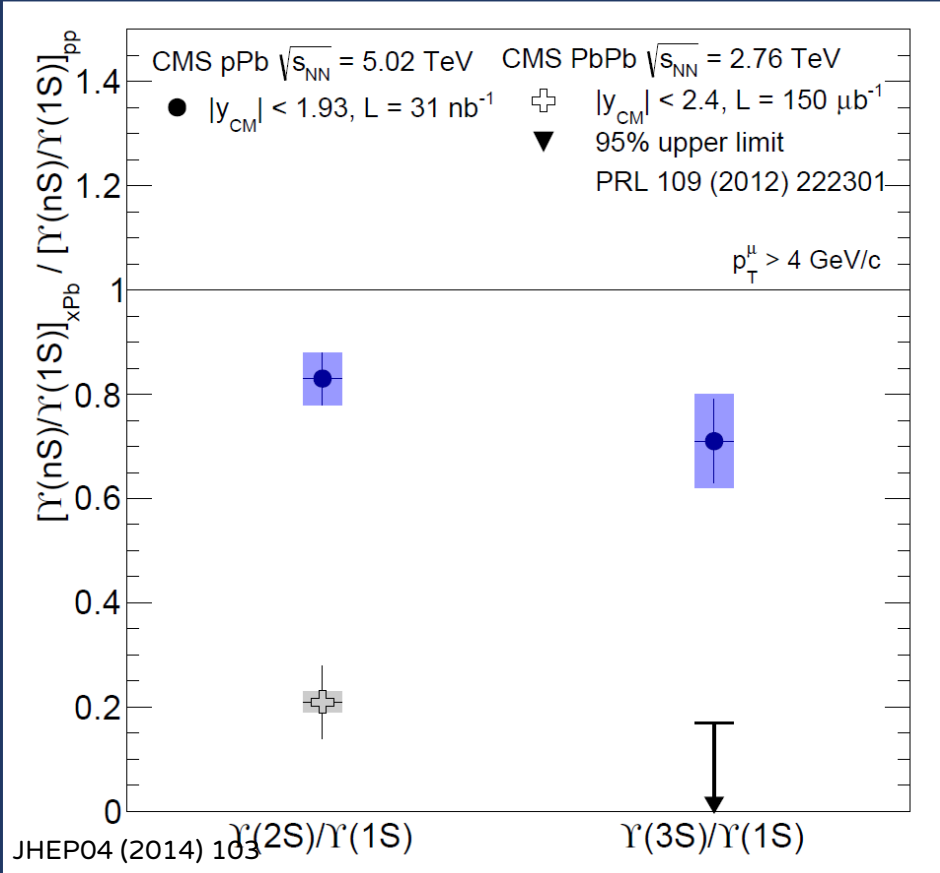
Lansberg 1610.05382

➔  $R_{pA}$  compatible, within uncertainties, with modifications due to shadowing and/or energy loss

➔ At backward- $y$ , models tend to be closer to LHCb measurement, while they overestimate the ALICE  $R_{pA}$



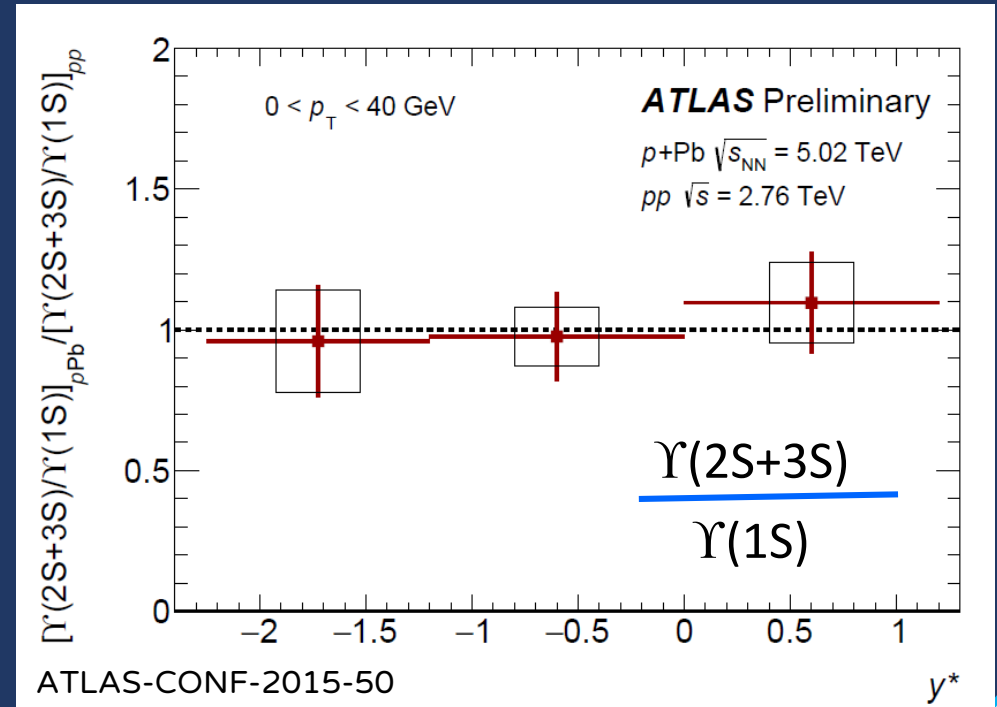
# Excited $\Upsilon$ states



➔ Excited bottomonium states are more suppressed than the ground state already in pA collisions

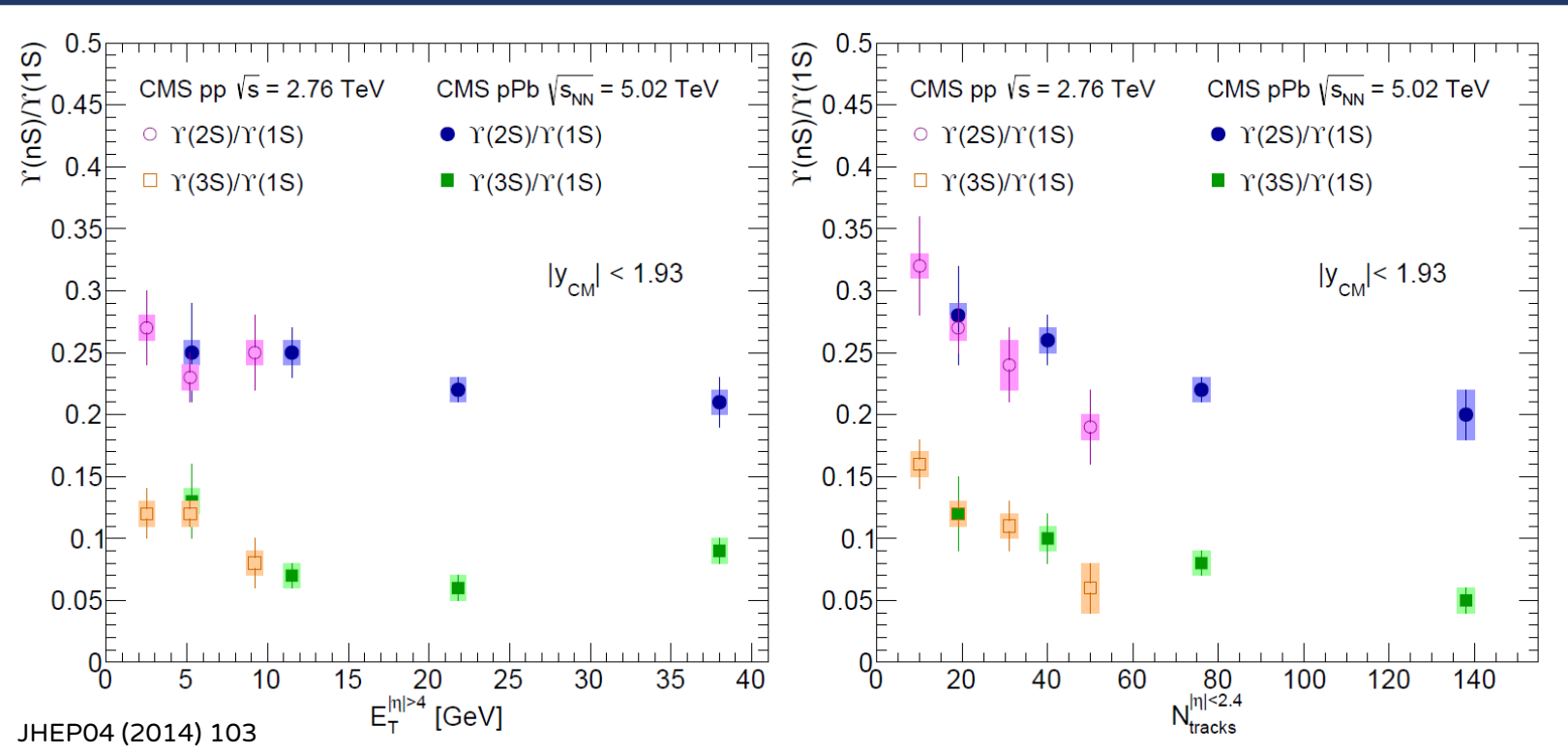
- ➔ Similar initial state effects for all the  $\Upsilon$  states
- ➔ Suggestive of final state effects at play on 2S and 3S states?

➔ ATLAS suggests no difference between excited and ground states, even is still compatible with CMS within uncertainties



# Excited $\Upsilon$ states

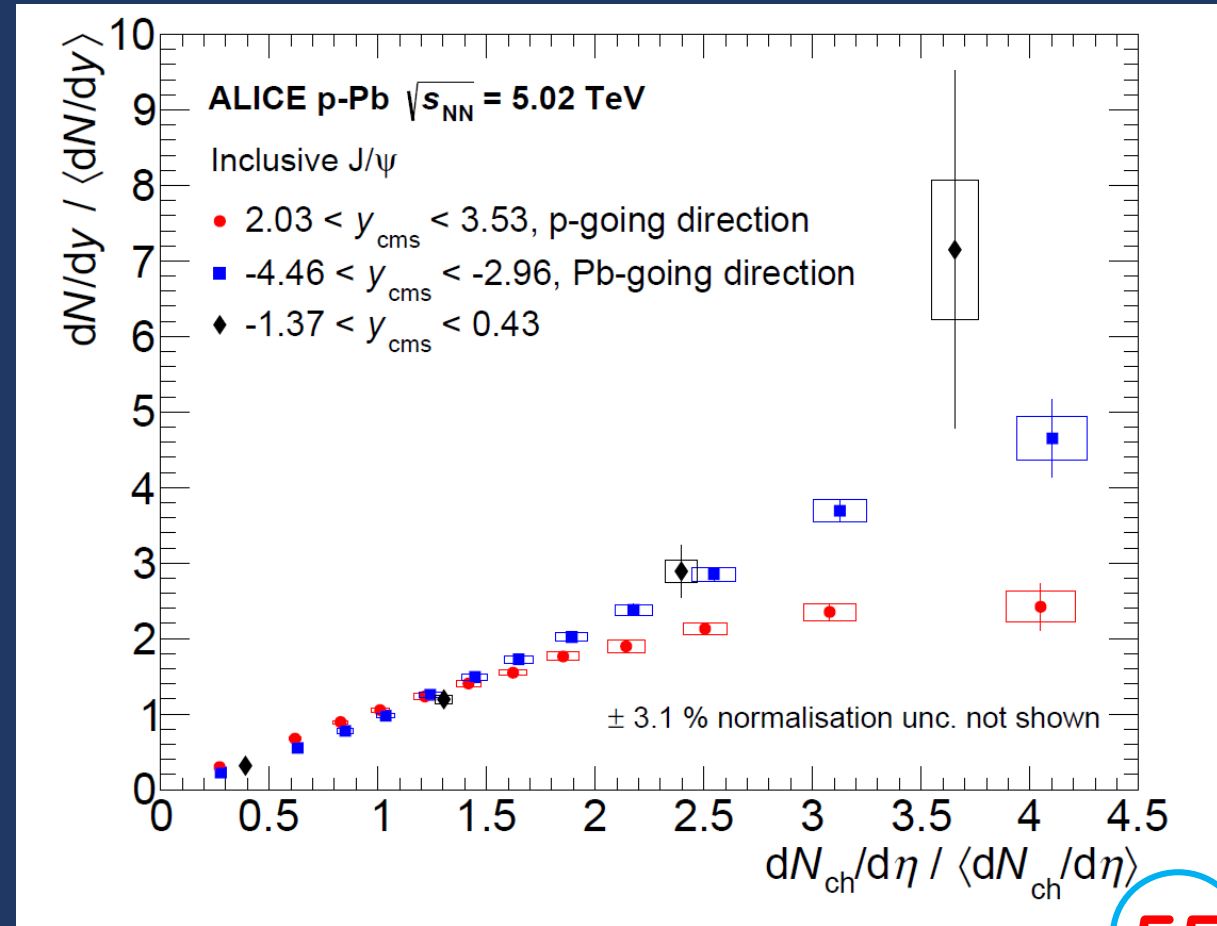
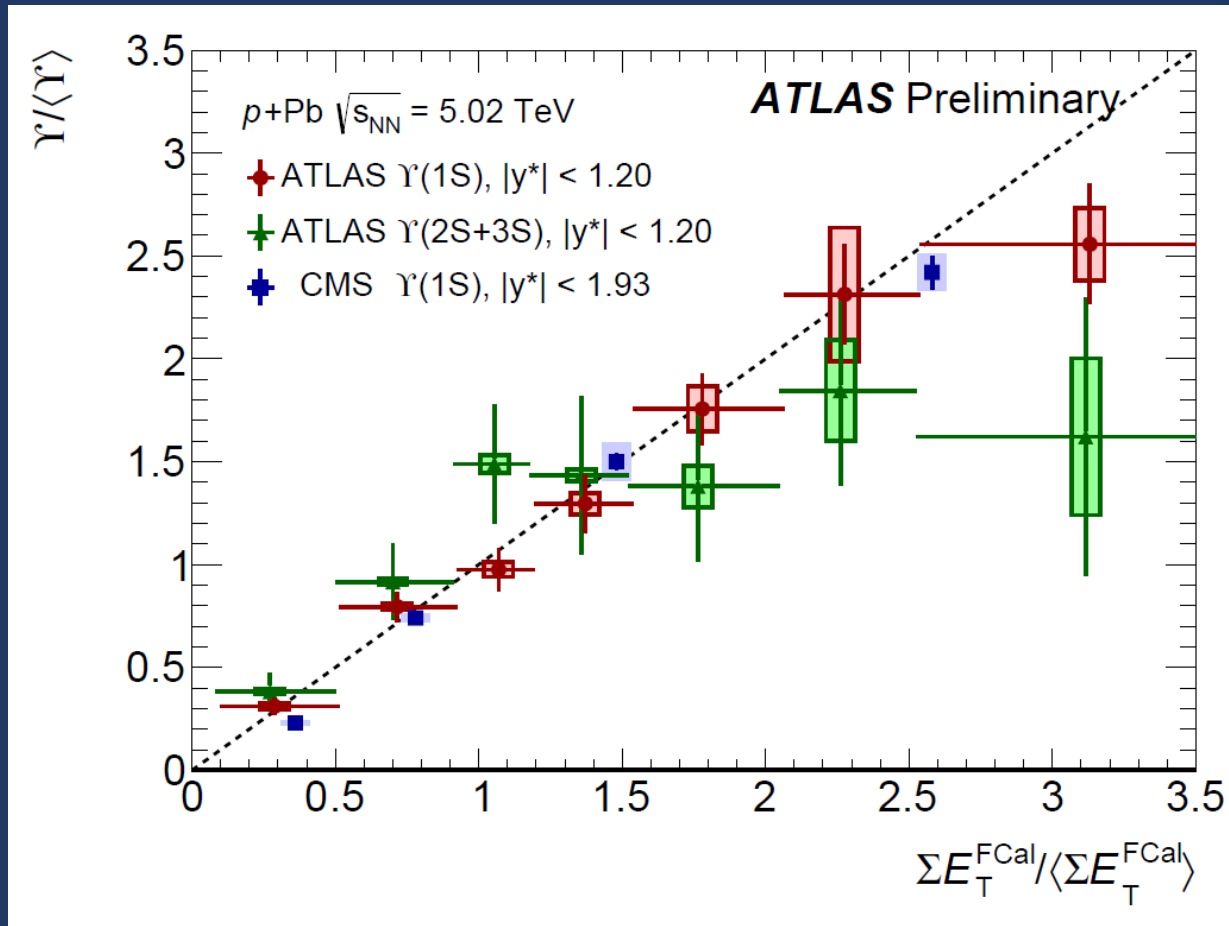
➔ Suppression increases with event multiplicity, but with a different trend depending on the adopted estimator



Larger number of particles produced with ground state or suppression of excited states?

# Self normalized ratios vs event activity

- ➔ Increase of the self-normalized yields vs event activity
- ➔ Compatible trends observed by ATLAS and CMS

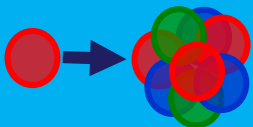


# Conclusions

Several quarkonium states now accessible with high precision in p-A and d-A

pp 

Cross-sections, double ratio and non-prompt contributions precisely measured over a broad kinematic range

pA 

Interplay of shadowing and energy loss describes  $J/\psi$  and  $\Upsilon$  production in p-Pb  
Stronger suppression observed on  $\psi(2S)$  due to QGP-like effects in pA

Many new results still to come....

Thanks!