Critical dynamics: complexity, universality and "hydro++"



Center for Theoretical Physics



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Outline

• Cumulants evolution: universality lost.

S. Mukherjee, R. Venugopalan and YY, 1506.00645, PRC; 1512.08022, QM proceedings.

• Kibble-Zurek dynamics: universality regained!

S. Mukherjee, R. Venugopalan and YY, 1605.09341.

• *New development*: Hydro + +.

M. Stephanov and YY, in preparation.

Cumulants evolution

Observables: $\langle (\delta N)^n \rangle \sim \kappa_n \sim \langle (\delta \sigma)^n \rangle$

Universality:
$$\kappa_n^{\text{eq}} \sim \xi_{\text{eq}}^{\#} \times f_n^{\text{eq}}(\theta)$$

• Enhanced non-Gaussian fluctuations.

 $\kappa_4^{\mathrm{eq}} \sim \xi_{\mathrm{eq}}^7$

• Universal pattern in sign.



(equilibrium kurtosis)

Critical dynamics matters

• Critical cumulants are unavoidably out of equilibrium :

$$au_{\sigma} \sim \xi_{\mathrm{e}q}^{z}, \qquad z \approx 3$$

- Non-equilibrium evolution of correlation length (Gaussian cumulants): Berdinkov-Rajagopal, 1999.
- Open question (as of 2014): evolution of non-Gaussian cumulants along trajectories passing the QCD critical regime .

Evolution equations for cumulants

- Critical mode receives "random kicks" from thermalized hard modes => Generalized Langevin equation.
- Averaging out the noise => A set of novel evolution equations for cumulants (S. Mukherjee, R. Venugopalan and YY, 1506.00645, PRC).

$$\frac{d\kappa_n(\tau)}{d\tau} = \frac{1}{\tau_\sigma(\tau)} L_n \left[\kappa_1, \kappa_2, \dots, \kappa_n; \mu_B(\tau), T(\tau), \text{mapping}\right]$$

$$au_{\sigma} \sim \xi_{\mathrm{eq}}^{z}, \qquad z \approx 3$$



• "Sign puzzle" of skewness: "remembrance of things past".



Complexity

- Evolution depends on many non-universal inputs:
 - mapping, location of critical point, width of critical regime.
 - relaxation time.
 - Trajectories in phase diagram.
- The non-equilibrium cumulants look complicated.



0

 $(T_c-T)/\Delta T$

0.7

1.4

-0.7

-1.4

- Is universality lost in complexity?
 - \bullet Which non-universal inputs (collectively denoted by Γ) dominate the dynamics?
 - What has been "memorized"?

Answers to those questions are connected by:

Kibble-Zurek dynamics

Kibble-Zurek dynamics

The formation and evolution of topological defects in cosmological phase transitions. (T.W. Kibble, Physics Reports 67, 183 (1980))

Generalized to vortex generation in superfluids. (W. H. Zurek, "Cosmological experiments in superfluid helium?", Nature 317, 505 (1985))

Kibble-Zurek dynamics in a little bang?

Consider a trajectory passing the critical point.





Comparing relaxation time and "quench" time.

$$\tau_{\sigma}(\tau) > \tau_{\text{que}}^{\xi} \equiv \left|\frac{\xi_{\text{eq}}}{\partial_{\tau}\xi_{\text{eq}}}\right|$$

• An emergent time scale for non-equilibrium dynamics.

 $\tau_{\sigma}(\tau_{\rm KZ}) = \tau_{\rm que}(\tau_{\rm KZ})$

The evolution is frozen.

$$l_{\mathrm{KZ}} = \xi_{\mathrm{eq}}(\tau_{\mathrm{KZ}})$$



- Kibble-Zurek dynamics: non-equilibrium critical evolutions are characterized by $l_{\rm KZ}, \tau_{\rm KZ}$.
- Given bulk evolution and critical equation of state, $l_{\rm KZ}$, $\tau_{\rm KZ}$ can be determined from the definition:

$$\tau_{\sigma}(\tau_{\rm KZ}) = \tau_{\rm que}(\tau_{\rm KZ}) \qquad \qquad l_{\rm KZ} = \xi_{\rm eq}(\tau_{\rm KZ})$$

• An illustrative example: Berdnikov-Rajagopal model revisited.

Scaling with length is not enough



- Rescale Gaussian cumulants determined from Berdnikov-Rajagopal model by $\ensuremath{\mathit{l}_{\rm KZ}}^2$.
- The peak value looks universal, but time-dependence does not.
- A step forward: let us rescale time by $\tau_{\rm KZ}$!



• We illustrated the existence of a scaling function:

$$\kappa_2(\tau;\Gamma) \sim l_{\text{KZ}}^2(\Gamma) \underbrace{f_2(\tau/\tau_{\text{KZ}}(\Gamma))}_{\text{Universal}} \quad (\Gamma: \text{non-universal inputs})$$

NB: the study of non-equilibrium scaling function is a new frontier in condensed matter cummunity: new physics in an old paper!

The Kibble-Zurek Problem: Universality and the Scaling Limit

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> Steven S. Gubser and S. L. Sondhi Department of Physics, Princeton University, Princeton, NJ 08544 (Dated: September 20, 2012)

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PHYSICAL REVIEW LETTERS

week ending 6 JULY 2012

Nonequilibrium Dynamic Critical Scaling of the Quantum Ising Chain

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We solve for the time-dependent finite-size scaling functions of the one-dimensional transverse-field Ising chain during a linear-in-time ramp of the field through the quantum critical point. We then simulate Mott-insulating bosons in a tilted potential, an experimentally studied system in the same equilibrium universality class, and demonstrate that universality holds for the dynamics as well. We find qualitatively

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Universality in the Dynamics of Second-Order Phase Transitions

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New ingredients when applied to heavy-ion collisions:

- Extending scaling hypothesis for non-Gaussian cumulants.
- Applying non-equilibrium scaling for trajectories away from the critical point:

A new realization of KZ dynamics

• For a generic trajectory near the crossover line, $\theta(\tau)$ flips sign.



• Relaxation time remains finite, but θ changes too fast!



• A new non-equilibrium scaling variable: ("memory of sign"). $\theta_{\rm KZ} = \theta(-\tau_{\rm KZ})$

• Non-equilibrium skewness: $\theta_{\rm F} > 0, \theta_{\rm KZ} < 0$

 Generalized scaling hypothesis (S. Mukherjee, R. Venugopalan and YY, 1605.09341):

$$\kappa_n(\tau;\Gamma) \sim l_{\mathrm{KZ}}^{\#}(\Gamma) f_n(\tilde{t},\theta_{\mathrm{KZ}})$$

- Testing scaling hypothesis: different trajectories, same $\theta_{\rm KZ}$.
- Equilibrium cumulants, τ_{KZ} , l_{KZ} are different for those trajectories.
- Expectation from the scaling hypothesis: scaling functions are independent of trajectories.





"Trilogy of Universality": universality lost, universality regained, universality discovered (to be composed)

Hydro++



A consistent description of flow and fluctuation observables is required.



 Hydro + + : a general framework for hydro with a slow mode (e.g. hydro + axial charge density, M. Stephanov, Ho-Ung Yee and YY 1501.00222).

- Applications to critical dynamics: the additional slow mode is the critical mode.
- Generalized entropy can be constructed as:

$$s(\epsilon, n; \sigma) = s_0(\epsilon, n) + \frac{1}{2!}a_2(\epsilon, n) \left[\sigma - \sigma_0(\epsilon, n)\right]^2 + \frac{1}{3!}a_3(\epsilon, n) \left[\sigma - \sigma_0(\epsilon, n)\right]^3 + \dots$$

(completely fixed by Ising model and mapping relation)

• The 6th equation (can be matched to evolution equation for the first cumulant):

$$u^{\mu}\partial_{\mu}\,\sigma = -\gamma\,\partial_{\sigma}\,s(\epsilon,n;\sigma)$$

• Illustrative example: Bjorken expansion (M.Stephanov and YY in preparation).

From better cartoon towards quantitive studies



Scaling, 16'

non-equilibrium, 15'

Equilibrium, II'

Hydro ++: next stage, new physics, let us do it together!

