

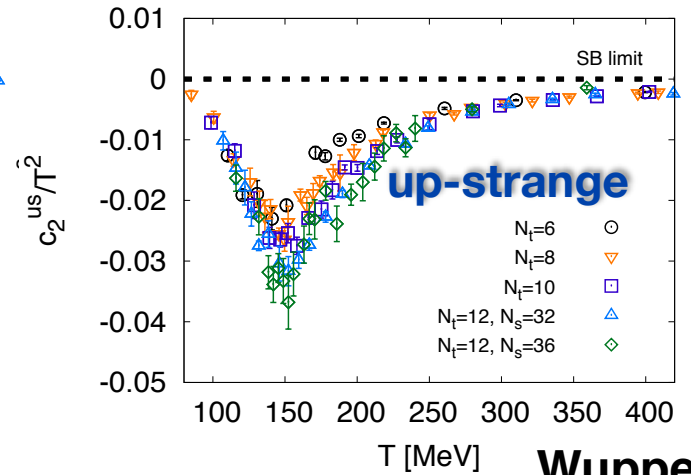
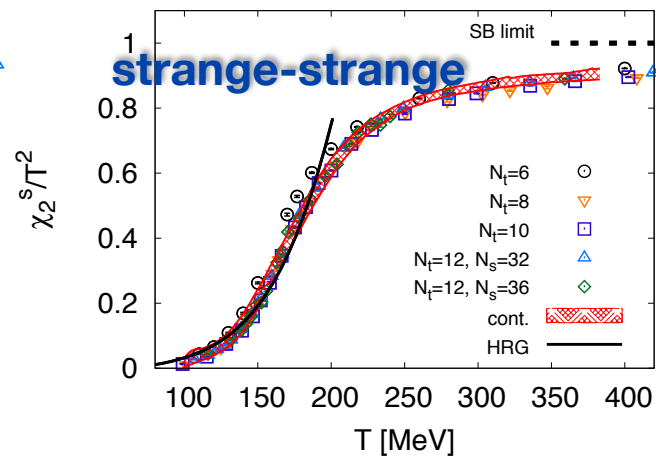
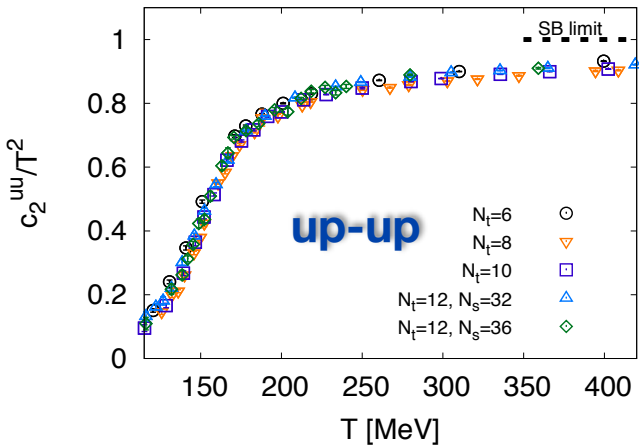
From Susceptibilities to Correlations

Scott Pratt and Clint Young, Michigan State University

- I. Paradigm***
- II. Production***
- III. Propagation***
- IV. Projection***
- V. Practice***
- VI. Predictions***
- VII. Promise***

PARADIGM

Lattice: Susceptibilities & Fluctuations



Wuppertal

$$\frac{1}{V} \langle \delta E \delta E \rangle = T^2 \frac{\partial \epsilon}{\partial T} = T^2 C_V$$

$$\frac{1}{V} \langle \delta P_z \delta P_z \rangle = (P + \epsilon) T = h$$

$$\frac{1}{V} \langle \delta Q_a \delta Q_b \rangle = T \frac{\partial \rho_a}{\partial \mu_b} = \chi_{ab}$$

$$\frac{1}{V} \langle \delta Q_a \delta V \rangle = T \frac{\partial \epsilon}{\partial \mu_a}$$

Grand Canonical Ensemble:

- Energy/charge fluctuate
- heat/particle bath
- infinite equilibration time

Experiment: Correlations

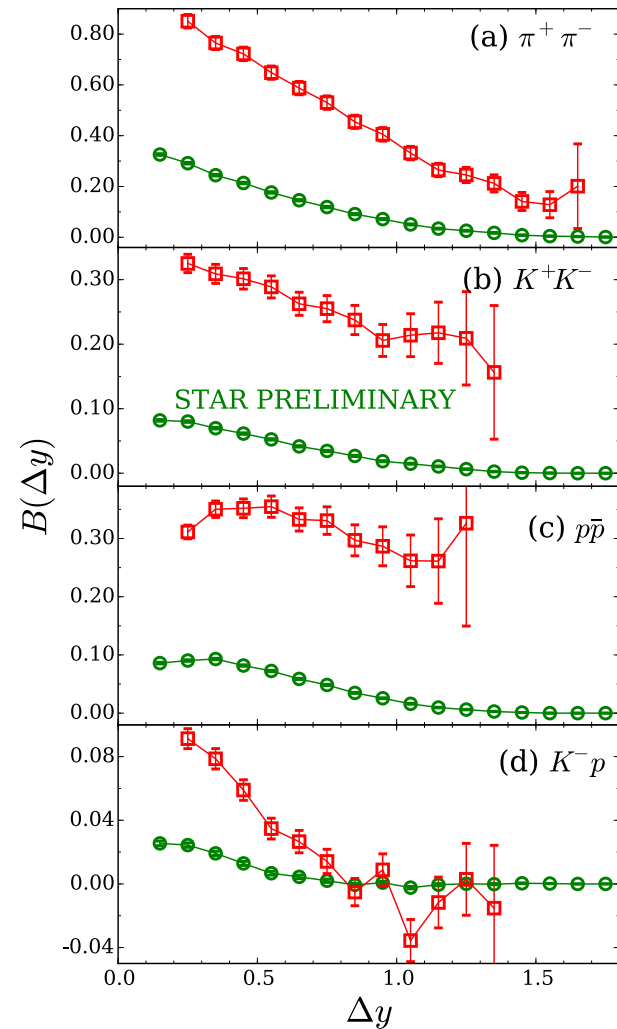
$$C(\delta E(0)\delta E(y))$$

$$C(\delta N_i(0)\delta N_j(y))$$

$$C(\delta E(0)\delta Q_i(y))$$

Final State:

- Charge/Energy/Momentum conserved
- Finite time for quantities to move
- Must translate to momentum space
- δP_z not accessible



Correlations in Equilibrated System

$$\int d\eta \frac{1}{V} \langle \delta Q_a(0) \delta Q_b(\eta) \rangle = \chi_{ab}$$

$$\int d\eta \frac{1}{V} \langle \delta E(0) \delta E(\eta) \rangle = T^2 C_V$$

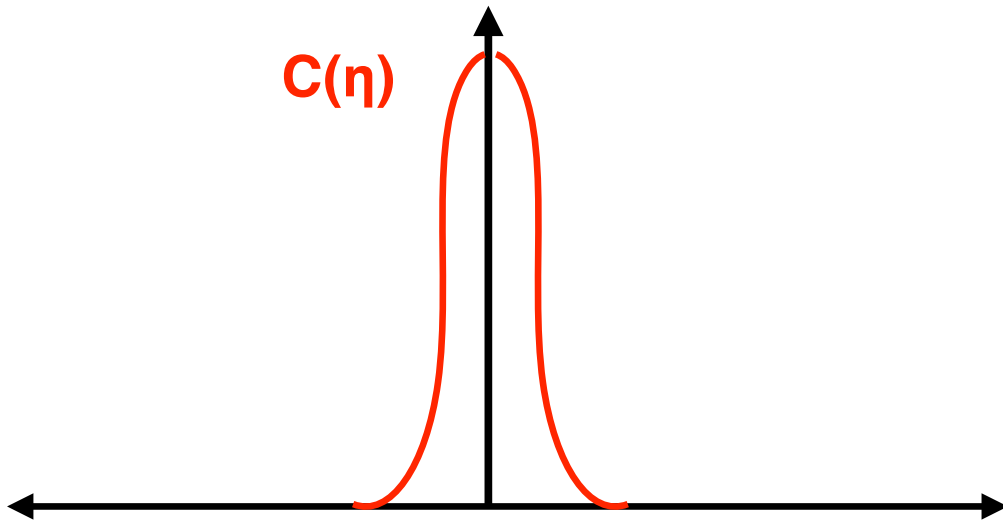
$$\int d\eta \frac{1}{V} \langle \delta P_z(0) \delta P_z(\eta) \rangle = (P + \epsilon)T$$

Correlations are LOCAL

$$\langle \delta Q_a(0) \delta Q_b(\eta) \rangle = \chi_{ab} \delta(\eta)$$

$$\langle \delta E(0) \delta E(\eta) \rangle = T^2 C_V \delta(\eta)$$

$$\langle \delta P_z(0) \delta P_z(\eta) \rangle = (P + \epsilon)T \delta(\eta)$$



Correlations in HI Environment

$$\int d\eta \frac{1}{V} \langle \delta Q_a(0) \delta Q_b(\eta) \rangle = 0$$

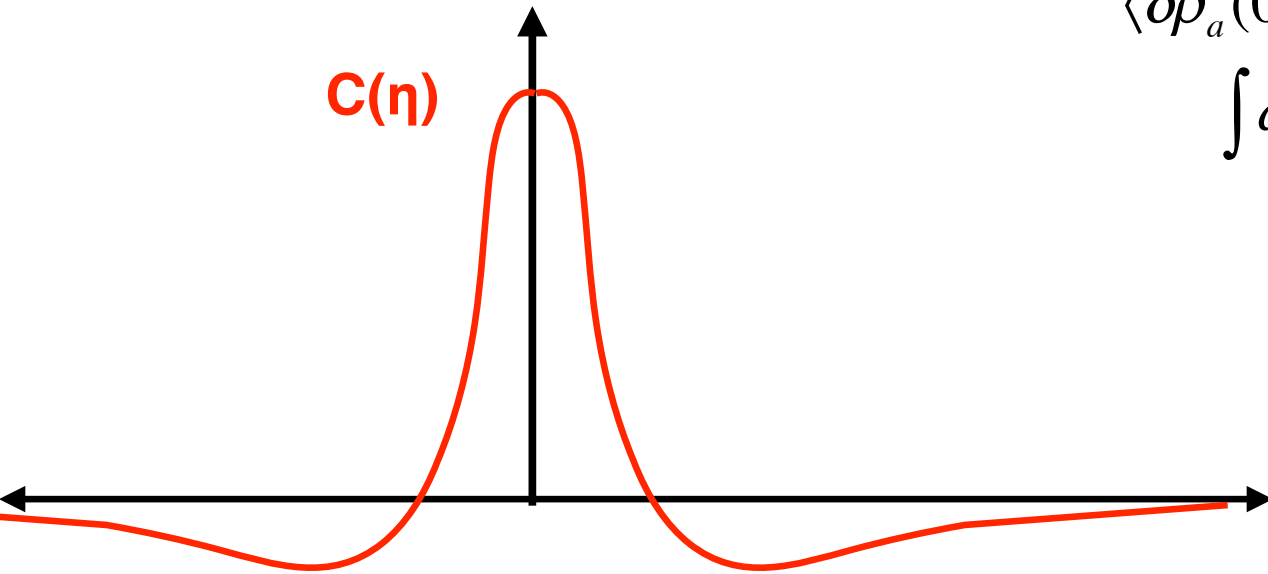
$$\int d\eta \frac{1}{V} \langle \delta E(0) \delta E(\eta) \rangle = 0$$

$$\int d\eta \frac{1}{V} \langle \delta P_z(0) \delta P_z(\eta) \rangle = 0$$

Non-local correlation: C'

$$\langle \delta \rho_a(0) \delta \rho_b(\eta) \rangle = \chi_{ab} \delta(\eta) + C'_{ab}(\eta)$$

$$\int d\eta C'_{ab}(\eta) = -\chi_{ab}$$



Evolution of C'

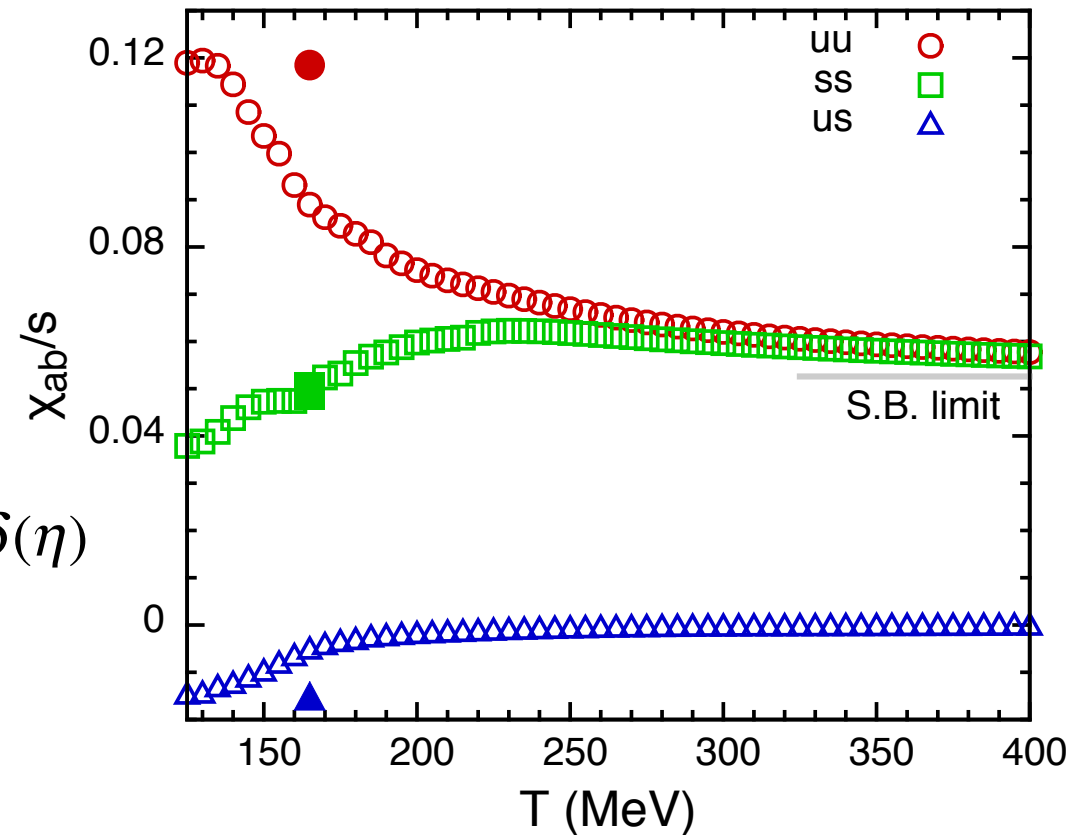
$$\partial C'(\eta, \tau) = \frac{1}{\tau} \partial_{\eta} (\dots) \text{ \{local cons.\}}$$

Charge-Charge

charge moves diffusively

$$\partial C'_{ab}(\eta, \tau) = \frac{1}{\tau} \partial_{\eta} \left\{ 2D \partial_{\eta} C'(\eta, \tau) \right\} + S(\tau) \delta(\eta)$$

$$S(\tau) = -\partial_{\tau} \left(\frac{\chi_{ab}}{s} \right)$$



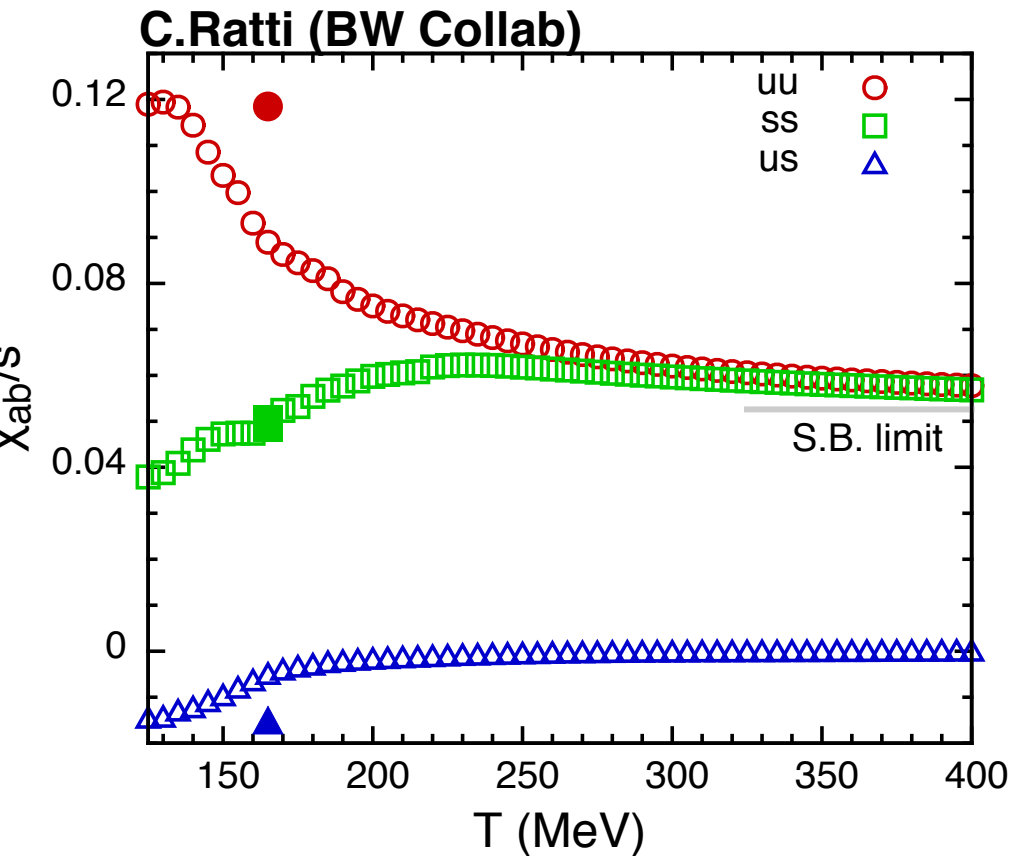
PRODUCTION

Charge Sources

$$\partial_\tau C_{ab}(\tau, \eta) - \frac{2D}{\tau^2} \partial_\eta^2 C_{ab}(\tau, \eta) = \delta(\eta) S_{ab}(\eta)$$

$$S_{ab}(\eta) = -\frac{dS}{d\eta} \frac{d}{d\tau} \left(\frac{\chi_{ab}}{s} \right)$$

**Source determined
by changing susceptibility**



P_x - P_x (transverse momentum)

Source defined by change in T

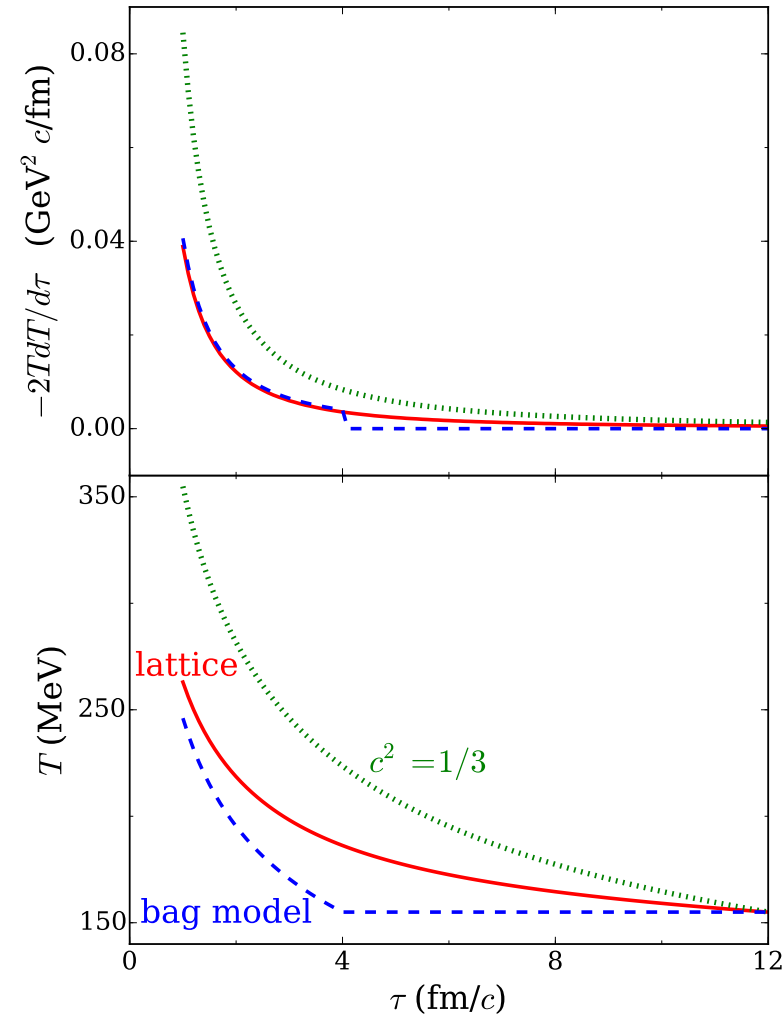
$$C_{xx}(\eta) = \langle \delta P_x(0) \delta P_x(\eta) \rangle$$

$$\delta P_x(\eta) \equiv \int dx dy \delta T_{0x}(\eta)$$

$$\partial_\tau C'_{xx}(\eta, \tau) = \frac{2D}{\tau^2} \partial_\eta^2 C'_{xx}(\eta, \tau) + S_{xx}(\tau) \delta(\eta),$$

$$\begin{aligned} S_{xx}(\tau) &= -\pi R^2 \frac{d}{d\tau} [(P + \epsilon) T \tau] \\ &= -\frac{dS}{d\eta} \frac{d}{d\tau} T^2. \end{aligned}$$

(no transverse flow)



$E_t - E_t$ & $P_z - P_z$

P_z and E_t evolve together

← becomes $E_t - E_t$ correlation

$$C_{TT}(\tau, \eta) = \tau^2 \cosh \eta \langle \delta \tilde{T}_{00}(\tau, 0) \delta \tilde{T}_{00}(\tau, \eta) - \delta \tilde{T}_{0z}(\tau, 0) \delta \tilde{T}_{0z}(\tau, \eta) \rangle \\ + 2\tau^2 \sinh \eta \langle \delta \tilde{T}_{00}(\tau, 0) \delta \tilde{T}_{0z}(\tau, \eta) \rangle.$$

$$C_{LT}(\tau, \eta) = \tau^2 \langle \delta \tilde{T}_{00}(\tau, 0) \delta \tilde{T}_{00}(\tau, \eta) - \delta \tilde{T}_{0z}(\tau, 0) \delta \tilde{T}_{0z}(\tau, \eta) \rangle \sinh \eta + 2\tau^2 \langle \delta \tilde{T}_{00}(\tau, 0) \delta \tilde{T}_{0z}(\tau, \eta) \rangle \cosh \eta$$

$$C_{LL}(\tau, \eta) = \tau^4 \langle \delta \tilde{T}_{0z}(\tau, 0) \delta \tilde{T}_{0z}(\tau, \eta) \rangle$$

← becomes $P_z - P_z$ correlation

E_t-E_t & P_z-P_z

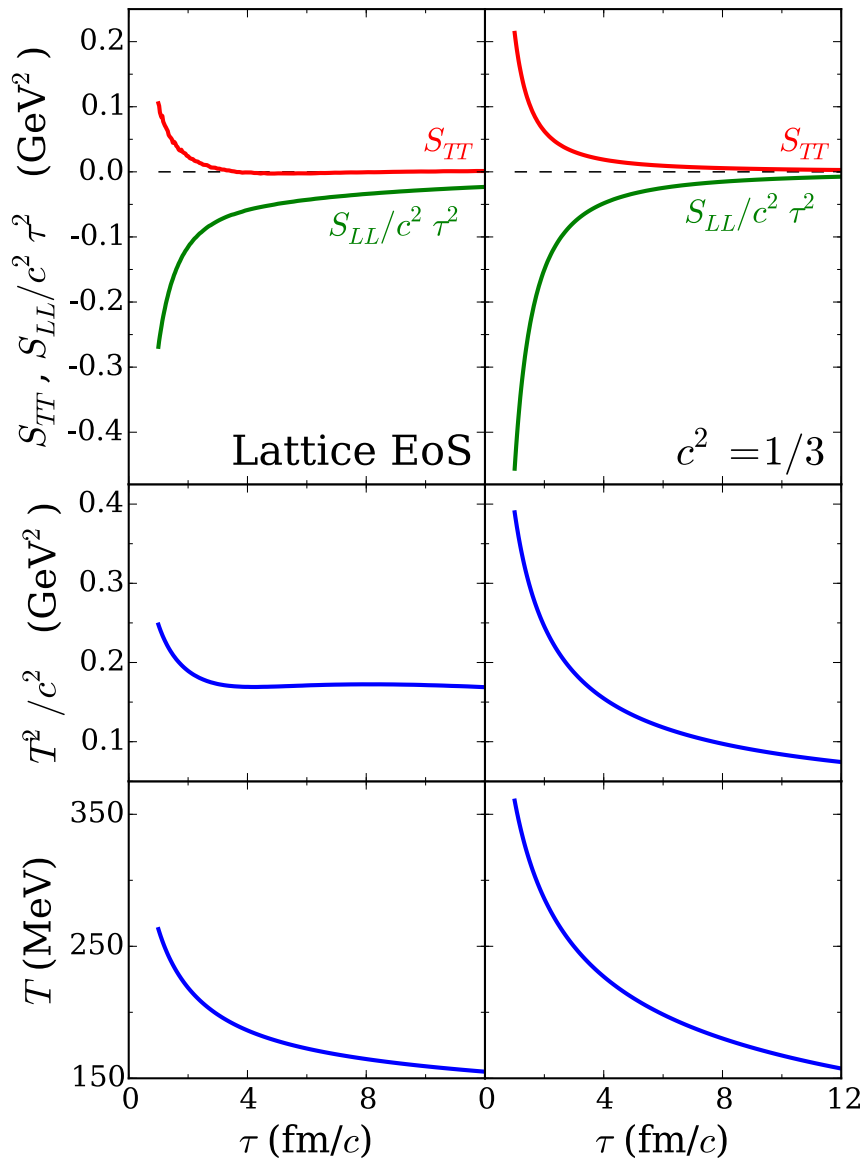
T^2/c^2 related to specific heat

$$S_{TT}(\tau) = -\frac{dS}{d\eta} \partial_\tau \left\{ T^2 \left(\frac{1}{c^2} - 1 \right) \right\}$$

$$S_{LL}(\tau) = -\frac{dS}{d\eta} \partial_\tau \left\{ \tau^2 T^2 \right\},$$

$$S_{LT} = 0$$

Same as P_x-P_x source



E_t-E_t & P_z-P_z

T^2/c^2 related to specific heat

$$S_{TT}(\tau) = -\frac{dS}{d\eta} \partial_\tau \left\{ T^2 \left(\frac{1}{c^2} - 1 \right) \right\}$$

$$S_{LL}(\tau) = -\frac{dS}{d\eta} \partial_\tau \left\{ \tau^2 T^2 \right\},$$

$$S_{LT} = 0$$

Same as P_x-P_x source

PROPAGATION

P_x - P_x (transverse momentum)

Spread is diffusive, D related to shear viscosity

$$C_{xx}(\eta) = \langle \delta P_x(0) \delta P_x(\eta) \rangle$$

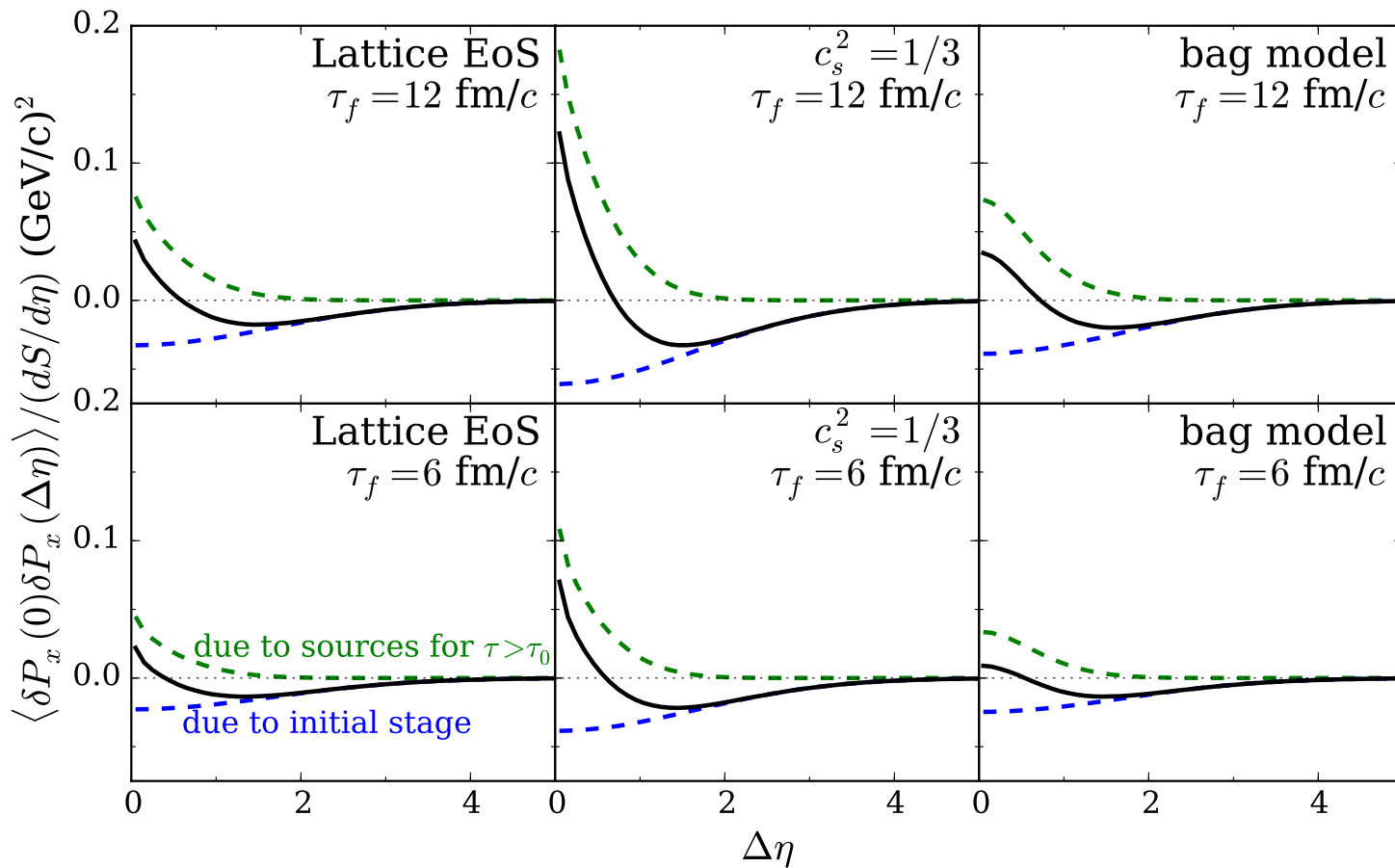
$$\delta P_x(\eta) \equiv \int dx dy \delta T_{0x}(\eta)$$

$$\partial_\tau C'_{xx}(\eta, \tau) = \frac{2D}{\tau^2} \partial_\eta^2 C'_{xx}(\eta, \tau) + S_{xx}(\tau) \delta(\eta),$$

$$\begin{aligned} S_{xx}(\tau) &= -\pi R^2 \frac{d}{d\tau} [(P + \epsilon) T \tau] \\ &= -\frac{dS}{d\eta} \frac{d}{d\tau} T^2. \end{aligned}$$

$$D = \frac{\eta_s}{P + \epsilon}$$

P_x - P_x (transverse momentum)



Result depends on:

- **Initial C'**
- **Eq. of State**
- **Viscosity**

E_t-E_t & P_z-P_z

$$C_{TT}(\tau, \eta) = \tau^2 \cosh \eta \langle \delta \tilde{T}_{00}(\tau, 0) \delta \tilde{T}_{00}(\tau, \eta) - \delta \tilde{T}_{0z}(\tau, 0) \delta \tilde{T}_{0z}(\tau, \eta) \rangle \\ + 2\tau^2 \sinh \eta \langle \delta \tilde{T}_{00}(\tau, 0) \delta \tilde{T}_{0z}(\tau, \eta) \rangle.$$

$$C_{LT}(\tau, \eta) = \tau^2 \langle \delta \tilde{T}_{00}(\tau, 0) \delta \tilde{T}_{00}(\tau, \eta) - \delta \tilde{T}_{0z}(\tau, 0) \delta \tilde{T}_{0z}(\tau, \eta) \rangle \sinh \eta + 2\tau^2 \langle \delta \tilde{T}_{00}(\tau, 0) \delta \tilde{T}_{0z}(\tau, \eta) \rangle \cosh \eta$$

$$C_{LL}(\tau, \eta) = \tau^4 \langle \delta \tilde{T}_{0z}(\tau, 0) \delta \tilde{T}_{0z}(\tau, \eta) \rangle$$

Correlation mixes up various components

$$\tau \partial_\tau C'_{TT} + 2\partial_\eta \left\{ \frac{(1+c^2)}{2} \cosh \eta (\cosh \eta C_{LT} - \sinh \eta C_{TT}) + c^2 \sinh \eta (\cosh \eta C_{TT} - \sinh \eta C_{LT} + (1/\tau^2) C_{LL}) \right. \\ \left. - \sinh \eta (1/\tau^2) C_{LL} \right\} = (S_{TT}/A) \delta(\eta),$$

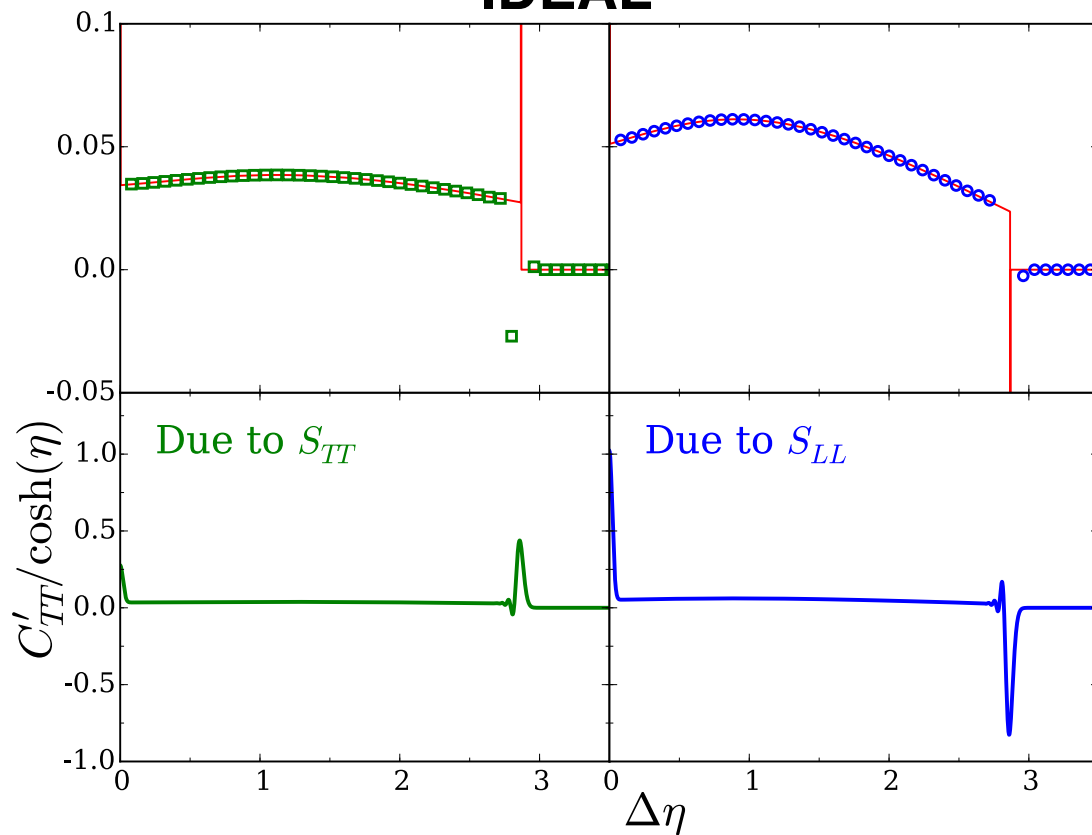
$$\tau \partial_\tau C'_{LT} + 2\partial_\eta \left\{ c^2 \cosh \eta (\cosh \eta C'_{TT} - \sinh \eta C_{LT} + (1/\tau^2) C'_{LL}) - \cosh \eta (1/\tau^2) C'_{LL} \right. \\ \left. + \frac{(1+c^2)}{2} \sinh \eta (\cosh \eta C'_{LT} - \sinh \eta C'_{TT}) \right\} = (S_{LT}/A) \delta(\eta),$$

$$\partial_\tau C'_{LL} + \tau^2 c^2 \partial_\eta \{ -\cosh \eta C'_{LT} + \sinh \eta C'_{TT} \} = (S_{LL}/A) \delta(\eta).$$

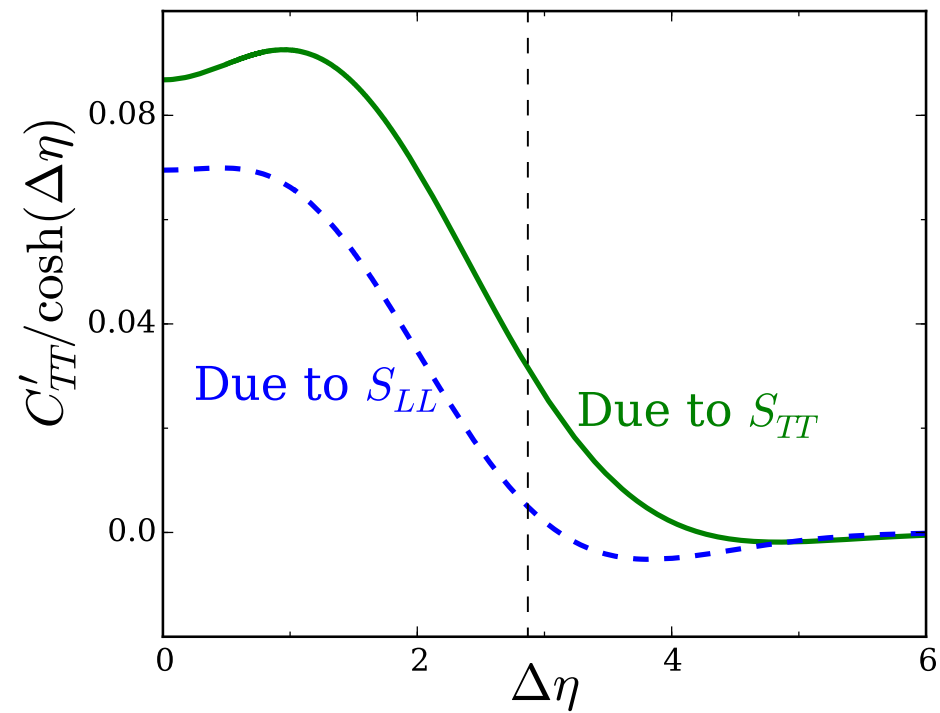
E/P_z Evolution is COMPLICATED

Evolve Pulse from $\tau=1$ fm/c to $\tau=12$ fm/c

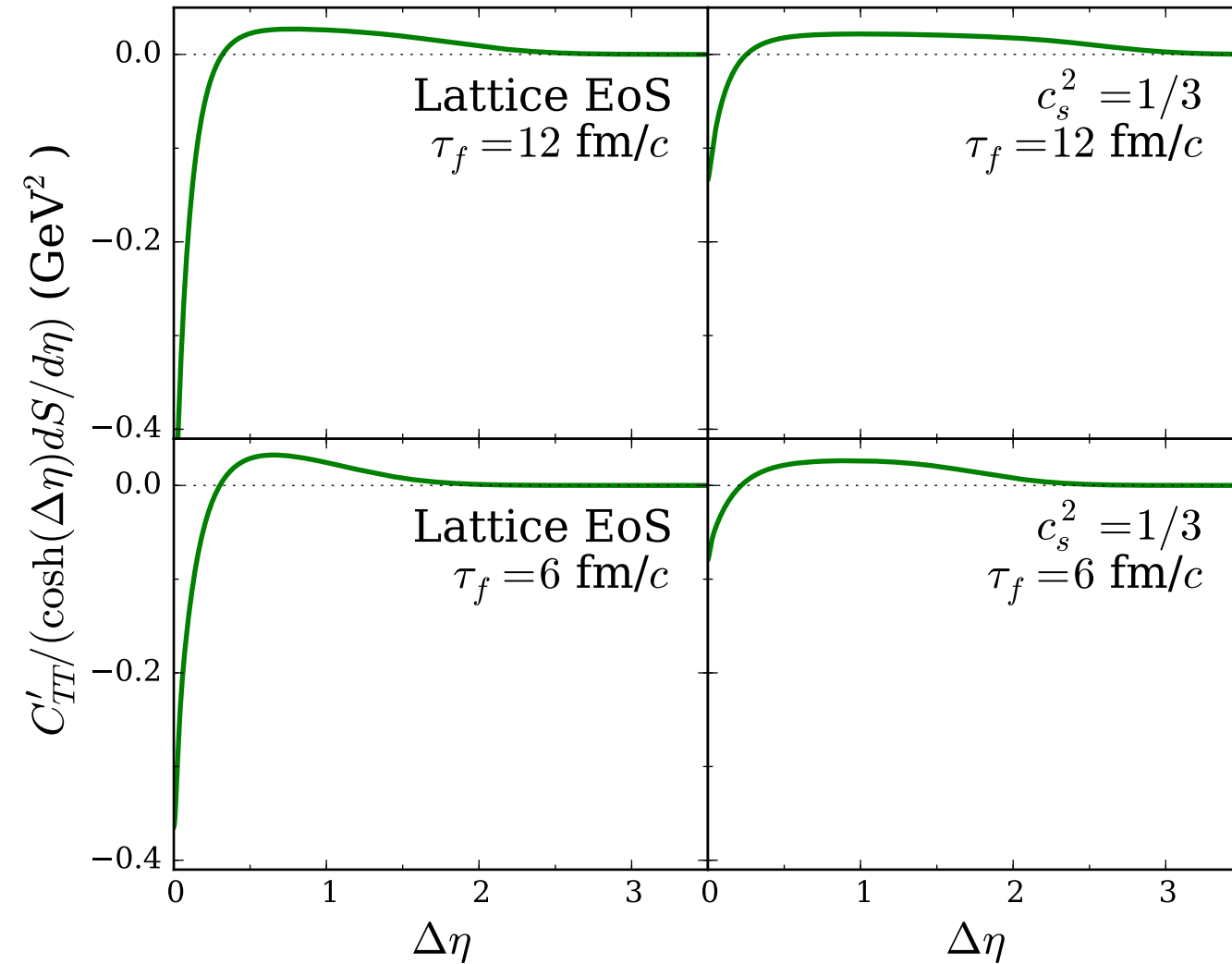
IDEAL



With Viscosity



E_t-E_t & P_z-P_z



- **Ignores initial source & transverse expansion, What about jets?**
- **Sensitive to EoS, but COMPLICATED**

PROJECTION

Projecting $C(\Delta\eta)$ to $C(\Delta y)$ either final state or cascade

$$f_h(\mathbf{p}) = f_h^{(0)} e^{\lambda_{a,\mu} \delta j_{h,a}^\mu(\mathbf{p})},$$

$$j_{h,a}^\mu(\mathbf{p}) = q_{h,a} \frac{p^\mu}{E_p}, \quad \text{Alter } f(p) \text{ by Lagrange Multipliers } \lambda_a$$

$$\delta f_h(\mathbf{p}) \approx f_h^{(0)}(\mathbf{p}) \lambda_{a,\mu} \delta j_{h,a}^\mu(\mathbf{p})$$

q_a could refer to any “charge”
including E or P

$$\chi_{ab}^{\mu\nu} = \sum_h \int \frac{d^3 p}{(2\pi)^3} f_h^{(0)}(\mathbf{p}) q_{h,a} \frac{p^\mu}{E_p} \frac{p^\nu}{E_p} q_{h,b}$$

$$\lambda_a^\mu = (\chi^{-1})_{ab}^{\mu\nu} \delta j_{\nu,b}$$

Projecting $C(\Delta\eta)$ to $C(\Delta y)$ either final state or cascade

$\delta\rho d\Omega_0 + \delta j_i d\Omega_i$ is charge that moves through
hypersurface $d\Omega^\mu$

$$\delta dN_h = d\Omega_0 \frac{d^3 p}{(2\pi)^3} f^{(0)}(\mathbf{p}) q_{h,a} \chi_{ab}^{-1} \delta\rho_b,$$

$$+ d\Omega_i \frac{d^3 p}{(2\pi)^3 E_p} f^{(0)}(\mathbf{p}) q_{h,a} \frac{p_i}{E_p} \left(\chi^{(J)} \right)_{ij,ab}^{-1} \delta j_{b,j}$$

$$\chi_{ij,ab}^{(J)} = \delta_{ij} \chi_{ab}^{(J)},$$

$$= \delta_{ij} \frac{1}{3} \sum_h \int \frac{d^3 p}{(2\pi)^3} f_h^{(\text{eq})} q_{h,a} q_{h,b} \frac{|\mathbf{p}|^2}{E_p^2}$$

**Works for all species,
any hyper-surface, $d\Omega$**

PRACTICE

Algorithm I: Solve Correlations

$$\partial_{\tau} C_{ab}(\tau, \eta) - \frac{1}{\tau} \partial_{\eta} \{ \dots \dots \dots \} = S_{ab}(\tau) \delta(\eta)$$

- **Reasonable for Bjorken geometry**
- **Impractical for transverse flow or boost invariance**
- **For BES, Q_a , P_x , P_z , E_t all mix**

Algorithm II: Propagate Q_a & Q_b separately

$$d\sigma_{ab} = d^4x S_{ab}(x)$$

$$d\sigma_{ab} \rightarrow d\tilde{\sigma}_{ab} = \begin{pmatrix} \tilde{\sigma}_{11} & 0 & 0 \\ 0 & \tilde{\sigma}_{22} & 0 \\ 0 & 0 & \tilde{\sigma}_{33} \end{pmatrix}$$

$$d\tilde{q}_a^{(1)} = (\pm) \sqrt{|d\tilde{\sigma}_{aa}|},$$

$$d\tilde{q}_b^{(2)} = d\tilde{\sigma}_{aa} / d\tilde{q}_a^{(1)}$$

$$dq_a = U_{ab}^{-1} d\tilde{q}_b$$

- Find source at each point in space-time

- diagonalize

Make charges in new basis with random signs

- Return to original basis

- Propagate (1) and (2) in separate evolutions
- Correlate (1) with (2)
- Add correlations from smooth distribution to account for local correlation

Algorithm II: Propagate Q_a & Q_b separately

SOME CHOICES

- **All source points together or one-at-a-time
one-at-a-time, no noise but need to sample $S(x)$**
- **For purely diffusive, random walk works great,
label partners for (1) and (2)
no need to rerun hydro evolution**

Algorithm III: Thermal Noise

Kapusta, Mueller, Young, Stephanov...

- Add external “currents” that are correlated

$$\langle j_i^{(n)}(x_1) j_k^{(n)}(x_2) \rangle = 2 \sigma T \delta_{ik} \delta^4(x_1 - x_2)$$

conductivity

- Solve one-particle evolution
- Generates local “delta” function

$$C(r_1 - r_2) = \chi \delta(r_1 - r_2) + \dots$$

- Correlate event with itself
- Must avoid double counting of local correlation
- Identical to previous methods for small mesh size

What's best?

For purely diffusive evolution of charges:

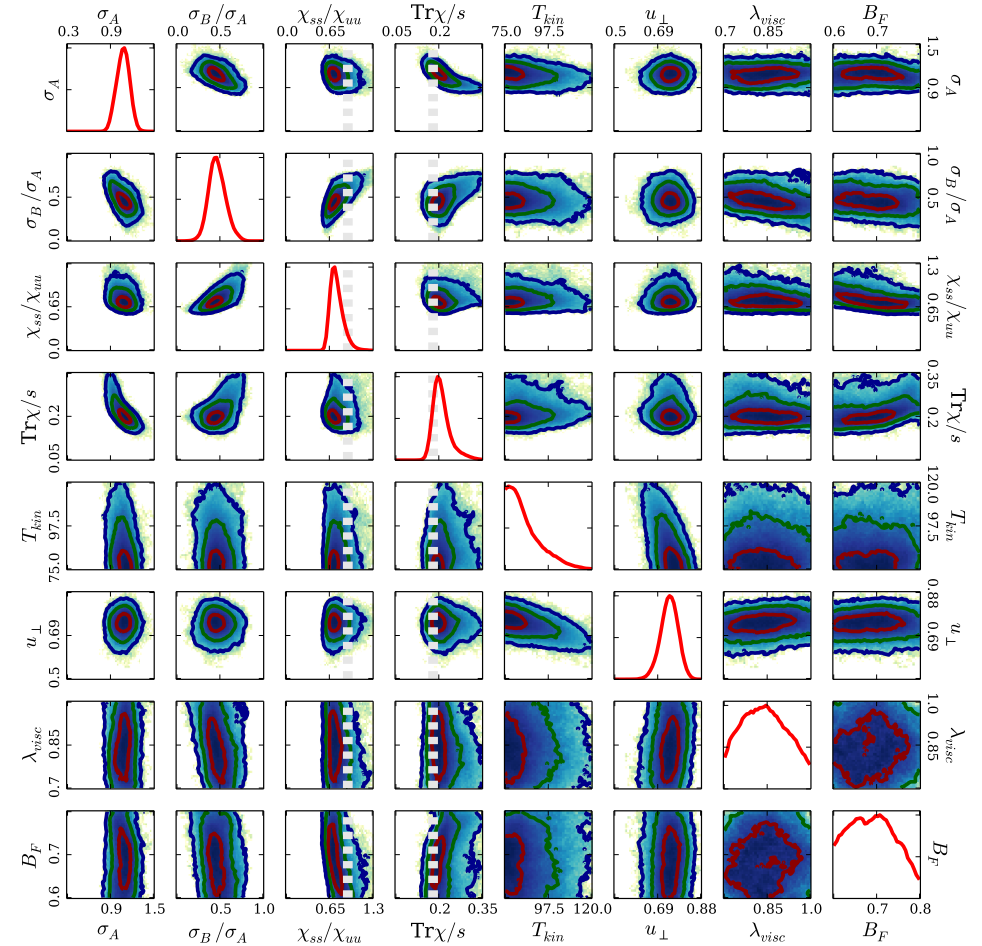
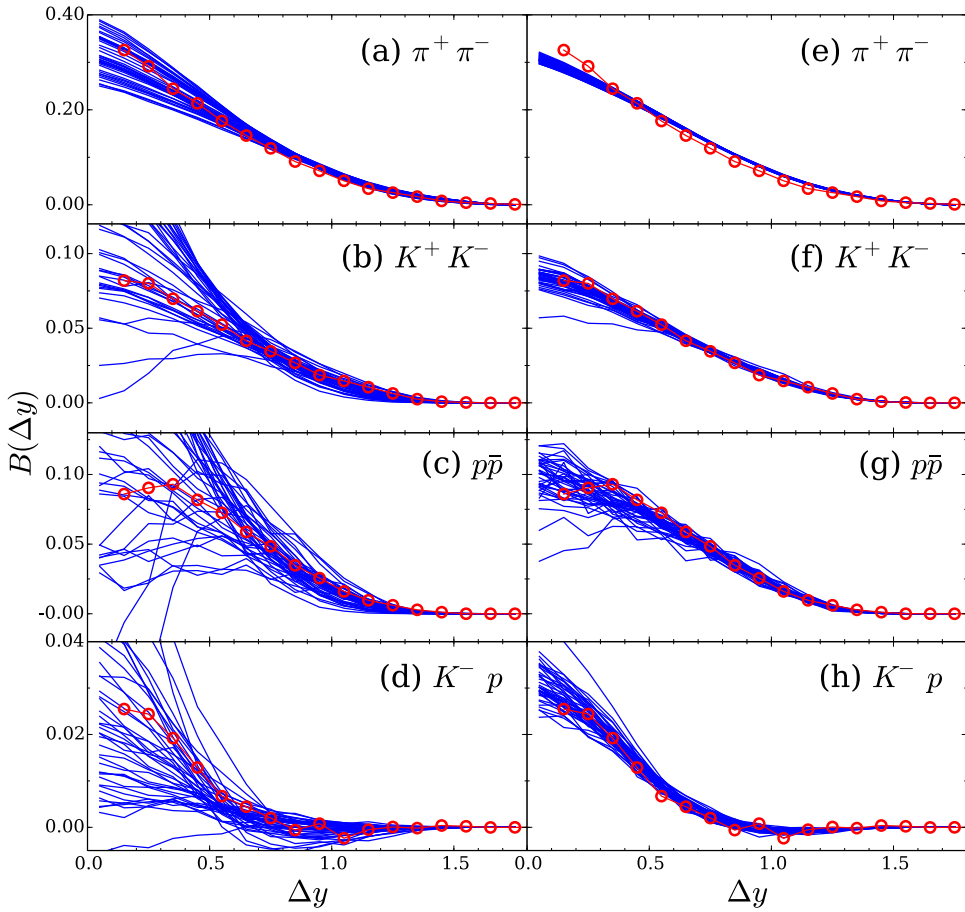
- **Algorithm II with random walk**

For energy/momentum:

- **Not obvious**
- **Both II and III require thousands of calculations to overcome noise or sample S .**

PREDICTIONS

STAR BALANCE FUNCTIONS VS. TWO-SURGE MODEL



Demonstrates ~ chemical equilibrium of QGP phase!!!

PROMISE

Charge Correlations

- 6-D measurement, $C(P,Q)$
- Chemical evolution
- Determine Diffusion Constant
- Understand initial thermalization of quarks

Energy Momentum

- Much more difficult to model
- Extra handle on EoS and viscosity
- Phase transition signal
- Understand initial-state correlations / jets

Limitations

- Assume instantaneous local equilibration
- What about $N>2$ correlations?