

# Equation of state and transport coefficients at finite baryo-chemical potential

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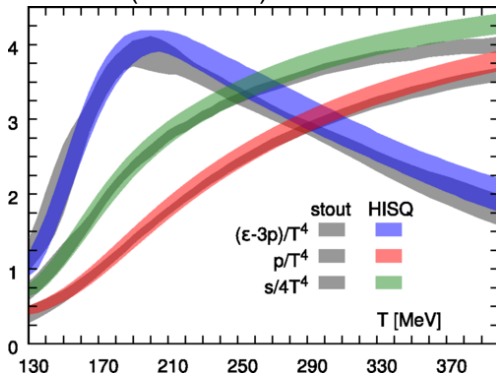
Exploring the QCD Phase Diagram through Energy Scans,  
INT  
October 6<sup>th</sup> 2016

# Outline

- 1 Lattice QCD
- 2 Black Hole Engineering
- 3 Critical Point
- 4 Transport Coefficients
- 5 Freeze-out Line(s)
- 6 Conclusion and Outlook
- 7 Backup
  - Backup

# The Success of Lattice QCD

Equation of State agrees for stout (WB) and HISQ (HotQCD) actions

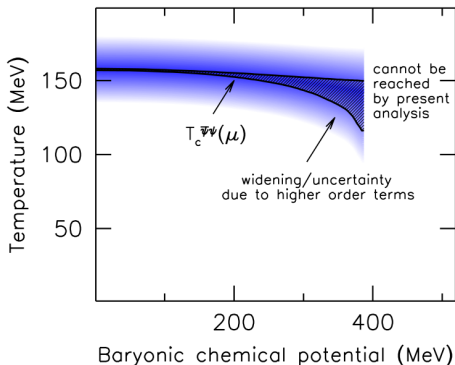


WB Phys.Lett. B730 (2014) 99-104  
HotQCD Phys.Rev. D90 (2014) 094503

# Limitations at Large $\mu_B$ (Sign problem)

Taylor expand pressure in term of  $\mu_B$ , limits results for large  $\mu_B$

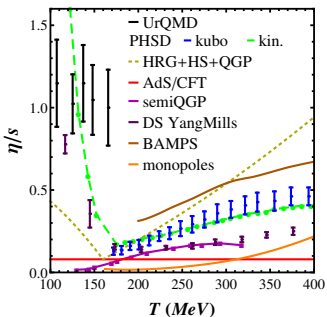
$$\frac{P(\mu_B)}{T^4} = c_0 + c_2 \left(\frac{\mu_B}{T}\right)^2 + c_4 \left(\frac{\mu_B}{T}\right)^4 + c_6 \left(\frac{\mu_B}{T}\right)^6 + \mathcal{O}(\mu_B^8)$$



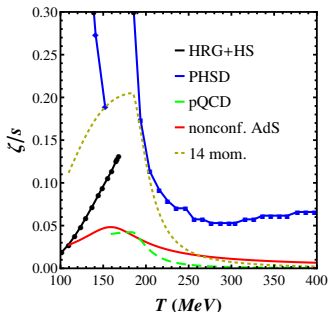
Phys.Lett. B751 (2015) 559-564

# Limitations for transport properties

- Lattice has technical difficulties to compute transport properties
- We're left with a number of models that don't converge...



JNH arXiv:1512.06315 (see for references)



# Filling in the gaps with Black Hole Engineering

What we need...

- Strongly coupled system
- Non-conformal equation of state
  - Equation of State at large baryon chemical potentials
  - Critical Point
- Perfect fluidity
  - Ability to compute transport coefficients near crossover and at large  $\mu_B$

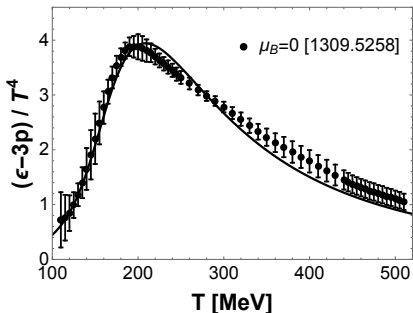
One alternative

Gauge/gravity duality - Black Hole Engineering



# non-conformal Equation of State

$$S = \frac{1}{2\kappa^2} \int_{\mathcal{M}_5} d^5x \sqrt{-g} \left[ \mathcal{R} - \frac{1}{2} (\partial_\mu \phi)^2 - \underbrace{V(\phi)}_{\substack{\phi \neq \text{const} \\ (\text{nonconformal})}} - \underbrace{\frac{f(\phi)}{4} F_{\mu\nu}^2}_{\mu_B \neq 0} \right]$$



$\kappa^2$  gravitational constant  
 $V(\phi)$  dilaton potential  
 $f(\phi)$  Maxwell-Dilaton coupling  
 all fixed to lattice data at  $\mu_B = 0$

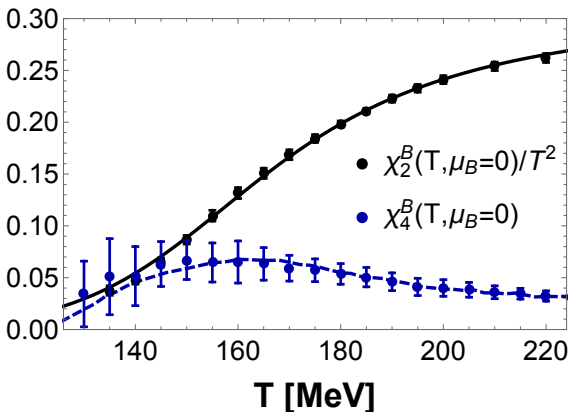
→ non-conformality!  
 Allows for  $\zeta/s > 0!$

See also: DeWolfe, Gubser, Rosen  
 PRD83(2011)086005; PRD84(2011)126014  
 Rougemont et al JHEP1604(2016)102;  
 Rougemont, Noronha, JNH, PRL115(2015)no.20,202301



# Baryon susceptibilities

Derivatives of the pressure  $\chi_n^B = \partial^n p / \partial \mu_B^n = \partial^{n-1} \rho / \partial \mu_B^{n-1}$



R. Rougemont, J. Noronha, JNH, PRL115(2015)no.20,202301

# Perfect Fluidity

Shear viscosity to entropy  
density\*

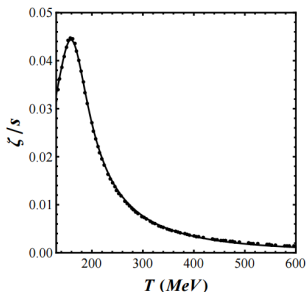
$$\frac{\eta}{s} = \frac{1}{4\pi}$$

Kovtun, Son, Starinets, 2005

\*The AdS/CFT bound that wasn't:

- Magnetic field violates KSS limit  
PRD90(2014)no.6,066006
- $\eta/s(T)$  higher-order  
derivatives of the action  
PRD77(2008)126006

non-conformal EOS needed for  
bulk viscosity

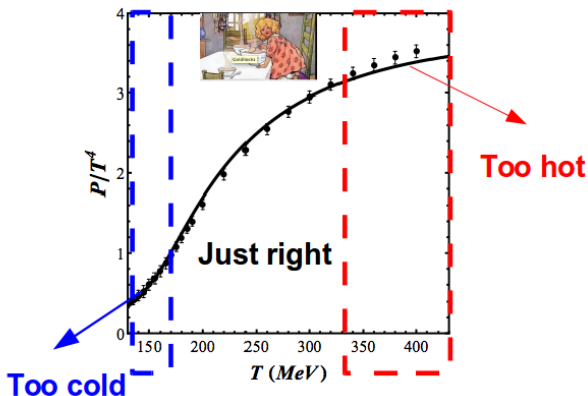


+13 more Israel-Stewart  
transport coefficients near  $T_c$   
S. Finazzo, R. Rougemont, H. Marrochio, J. Noronha,  
JHEP 1502 (2015) 051

# Transition Region-Goldilocks Zone

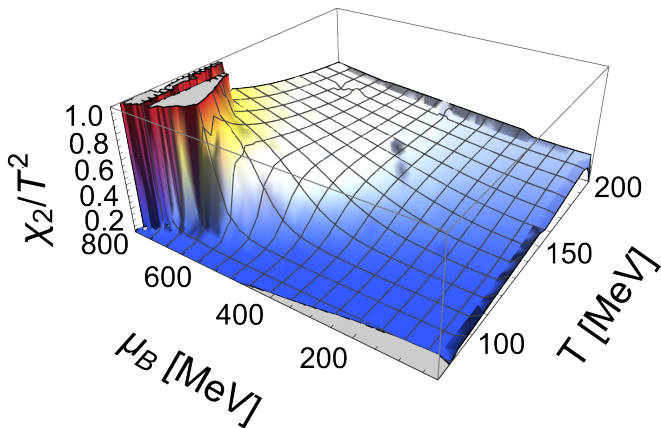
Only concerned with the transition region

“Holographic Goldilocks”



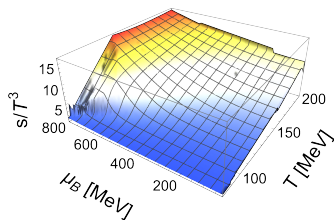
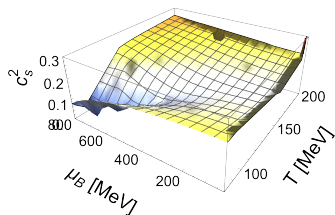
# Critical Point

Critical Point ( $T = 90$  MeV and  $\mu_B = 725$  MeV) emerges naturally from the theory

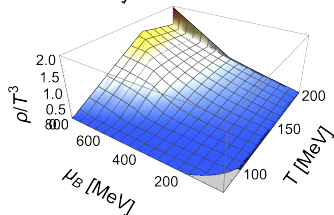


# Equation of State at finite $\mu_B$

Everything at  $\mu_B > 0$  is a prediction

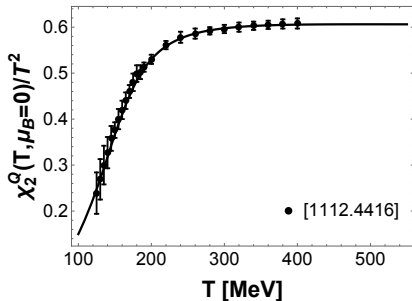
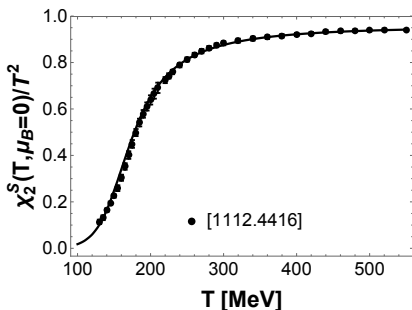


Critical behavior very sensitive  
to  $\mu_B = 0$  Lattice QCD input



# Strangeness and Electric Charge

- Since  $\mu_S < \mu_B$  and  $\mu_Q \ll \mu_B$ , assume  $\mu_S \sim \mu_B \sim 0$
- Caveat: only valid when  $\mu_S/\mu_B$  and  $\mu_Q/\mu_B$  are small



# Viscosity vs. Dynamic Universality Class

## No Critical Point

- Calculations possible within HRG, Transport etc
- General  $\downarrow$  in viscosity as  $\mu_B \uparrow$

## Dynamical CP phenomena Review

Hohenberg and Halperin, Rev. Mod. Phys. 49, 435

## Critical Point Universality Class H

- 3D Ising Model-  
Mixing between  
chiral condensate  
and baryon density

Son and Stephanov PRD70  
(2004) 056001

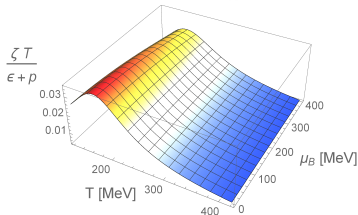
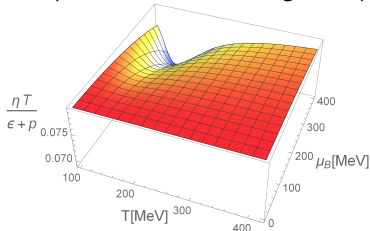
- Divergence  $\uparrow$  in viscosity as approaching CP
- See Stephanov and Yin's talks

## Critical Point Universality Class B

- Black Hole Engineering
- Currently B conserved, working on S & Q!
- $\downarrow$  in viscosity as approaching CP
- Original AdS/CFT CP: Phys.Rev. D78 (2008) 106007

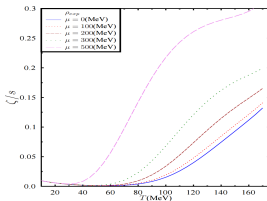
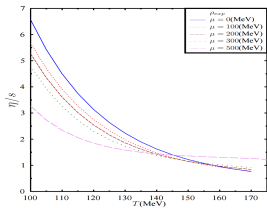
# Do we still have perfect fluidity at finite $\mu_B$ ?

CP (Class B- no divergence)



Rougemont, Noronha, JNH, Ratti to appear shortly

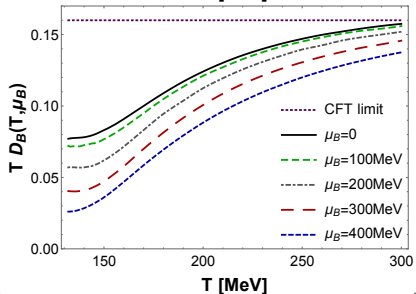
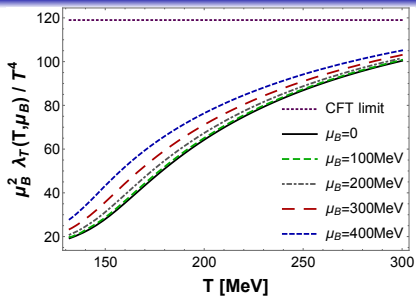
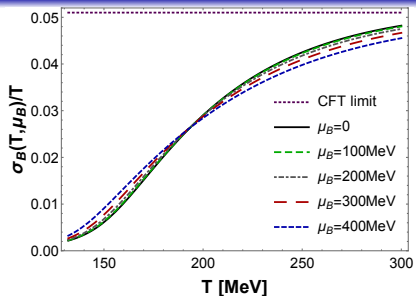
HRG/No CP



Kadam, Mishra Nucl.Phys. A934 (2014) 133-147  
See also Denicol, Jeon, Gale, Noronha Phys.Rev. C88 (2013) no.6, 064901



# Baryon Transport Coefficients



## DC Conductivity

$$\sigma_B = -\frac{\Lambda}{2\kappa^2 \phi_A^{1/\nu}} \lim_{\omega \rightarrow 0} \frac{h(r)f(\phi)e^{2A(r)} \text{Im}[a^*(r, \omega)a'(r, \omega)]}{\omega}$$

$$\text{Baryon diffusion } D_B = \sigma_B / \chi_2^B$$

Thermal conductivity

$$\lambda_T = (\sigma_B / T) [(\epsilon + p) / \rho]^2$$

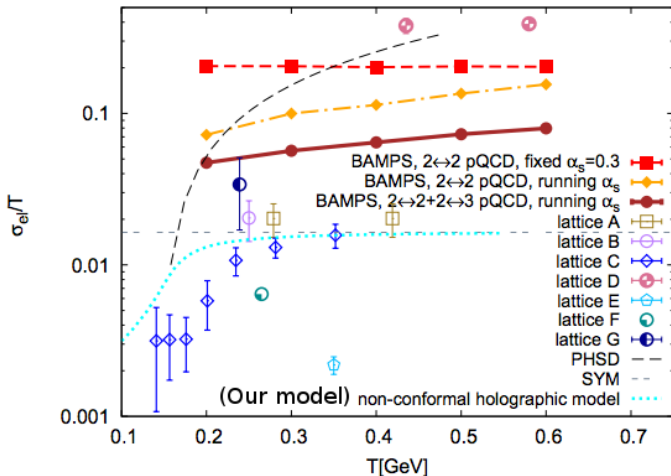
Rougemont, Noronha, JNH

Phys.Rev.Lett. 115 (2015) no.20,

202301

# Electric Charge from various studies at

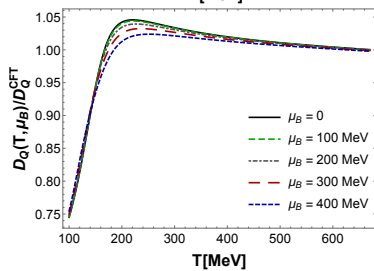
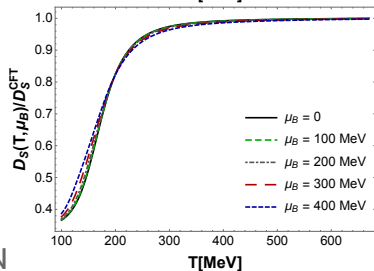
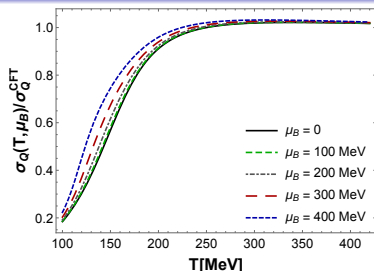
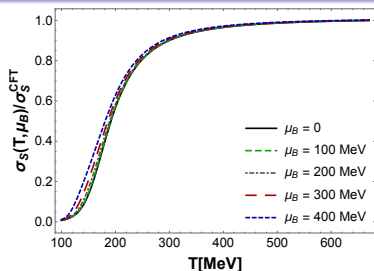
$$\mu_B = \mu_S = \mu_Q = 0$$



Greif et al Phys.Rev. D90 (2014) no.9, 094014 (see citations within)

# Strange and Electric Conductivity Transport Coefficients

(to appear soon)



# Vorticity Viscous Coupling (only at $\mu_B = 0$ so far)

Israel-Stewart 2<sup>nd</sup> order terms:

- Shear coupling terms:

$$-\frac{\lambda_2}{\eta} \pi_\lambda^{\langle \nu} \Omega^{\mu \rangle \lambda}$$

$$-\lambda_3 \Omega_\lambda^{\langle \mu} \Omega^{\nu \rangle \lambda}$$

- Bulk coupling term:

$$+\lambda_3 \Omega_{\mu\nu} \Omega^{\mu\nu}$$

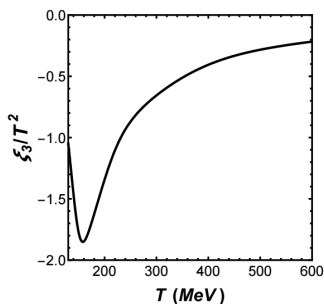
In red → violates causality

Denicol et al Phys.Rev. D85 (2012) 114047; Phys.Rev.

D89 (2014) no.7, 074005; Finazzo et al JHEP 1502

(2015) 051

Vorticity puzzle: including vorticity in hydro while preserving causality? Models beyond hydro?



$$\lambda_2 = -\ln \frac{2\eta}{\pi T} \text{ and}$$

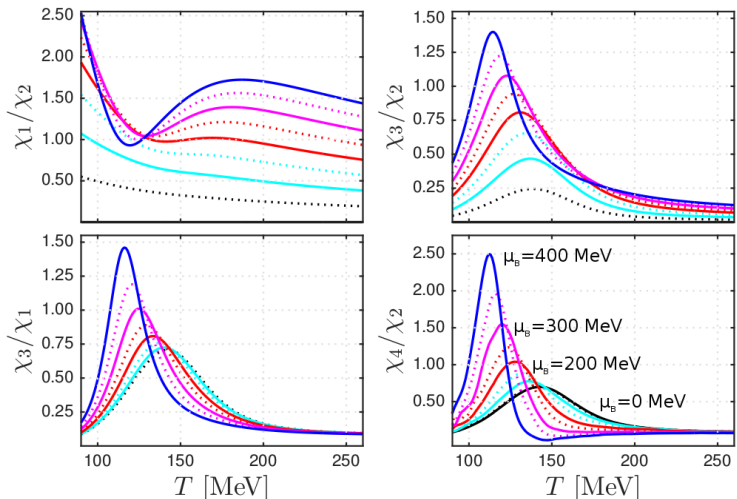
$$\xi_2 = 2\eta\tau_\pi c_s^2 \left( \frac{1}{3} - c_s^2 \right)$$

Finazzo et al JHEP 1502 (2015) 051

# What defines the transition region?

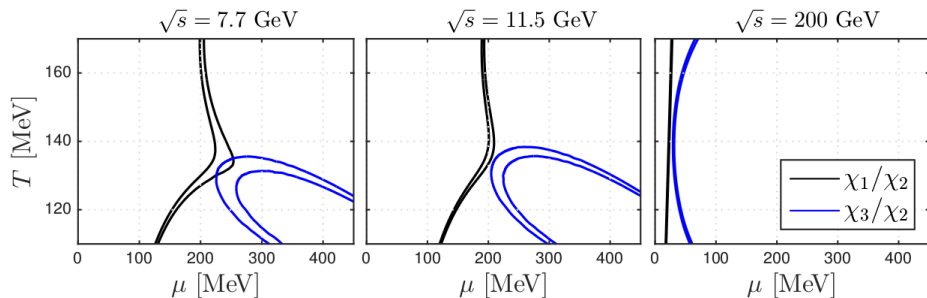
- Determine the **inflection point of the susceptibilities**  $\chi_n^{BSQ}$ 's across  $\mu_B$
- Determine the **inflection point of the transport coefficients** across  $\mu_B$ 
  - Should the inflection point of transport coefficients match the chemical freeze-out line?
- **Susceptibilities compared to experiments:** Compare derivatives of the pressure e.g.  $\chi_2^B/\chi_1^B$  to experimental data
  - All caveats from Claudia Ratti's talk

# Susceptibilities in the black hole model



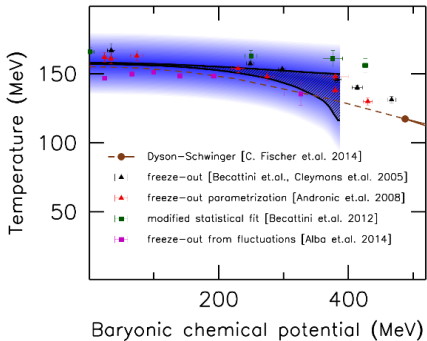
- Compare  $\chi_1/\chi_2$  and  $\chi_3/\chi_2$  to (net-p) data from STAR  
[STAR] Phys. Rev. Lett. 112 (2014) 032302

# Extraction of the freeze-out line from susceptibilities

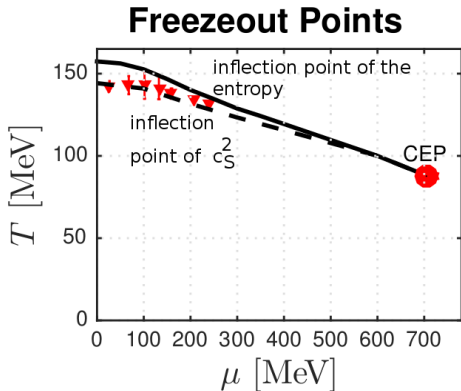


- Freeze out points  $[T - \mu_B]$  are extracted from the line made by the closer points between  $\chi_1/\chi_2$  and  $\chi_3/\chi_2$

# Freeze-out line from the blackhole model



R. Bellwied et. al., Phys. Lett. B **751**  
(2015) 053



Black Hole Model





# Conclusions and Outlook

- Black hole engineering provides a strongly interacting theory, non-conformal EoS that matches lattice, and calculable transport coefficients
- Critical Point arises at  $T = 90$  MeV and  $\mu_B = 725$  MeV  
**sensitive to Lattice data at  $\mu_B = 0$ !**
- Near crossover,  $\mu_B \geq 0$ , transport coefficients are suppressed compared to conformal field theory
- Freeze-out line compared to experimental data correlated with minimum of  $c_s^2$
- Theory work is needed! Inclusion of multiple nonzero chemical potentials