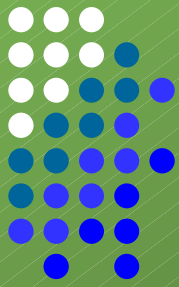


Neutrino Oscillations and the MaVaN Model

Kevin Weil

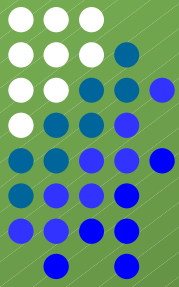
UW REU

Summer 2004



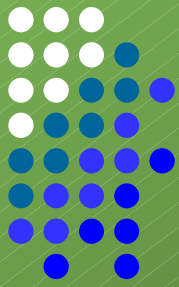
Outline

- Massive neutrinos and the Standard Model
- Neutrino oscillations
 - The standard picture
 - The MaVaN picture
- Outlook



Outline

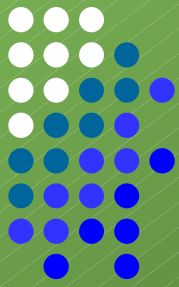
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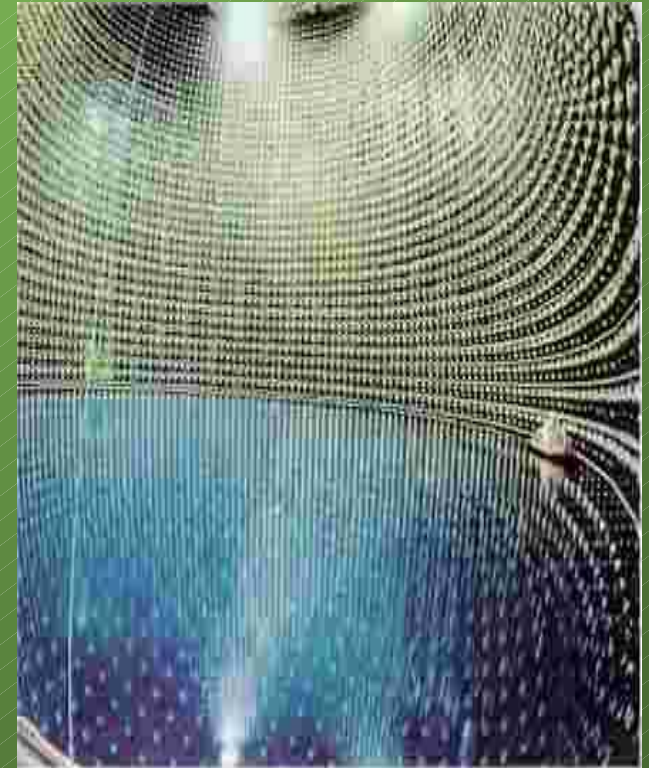
The Standard Model

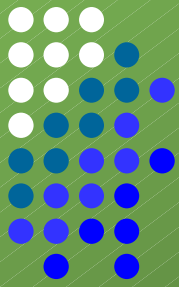
- According to the Standard Model, all three neutrino masses are zero
- Nonzero masses can work, but require an extension of the model
- Difficult to measure masses because neutrinos rarely interact
 - A neutrino of moderate energy can penetrate many light years of lead!
 - They're passing through us *right now*

Super-Kamiokande



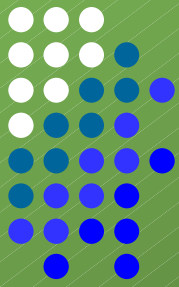
- 50,000 ton underground water tank
- Photomultiplier tubes see results of electron neutrino interactions
 - But not the actual neutrinos
- Fewer results than expected
 - Resolution: neutrino oscillations, which can only happen with nonzero mass





Outline

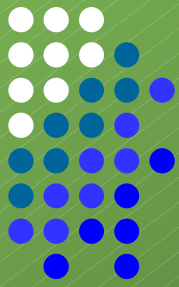
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Neutrino Oscillations

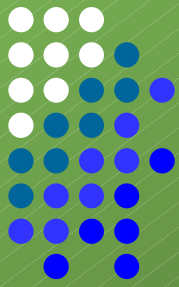
- Consider only electron and muon neutrinos
 - “Flavor eigenstates” $|\nu_e\rangle, |\nu_\mu\rangle$
- Different from the eigenstates of the Hamiltonian
 - “Mass eigenstates” $|\nu_1\rangle, |\nu_2\rangle$
- Related by a (vacuum) mixing angle θ_0
 - To get from one basis to the other, multiply by a unitary transformation $U(\theta_0)$:

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix} = \begin{pmatrix} \cos(\theta_0) & \sin(\theta_0) \\ -\sin(\theta_0) & \cos(\theta_0) \end{pmatrix} \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \end{pmatrix}$$



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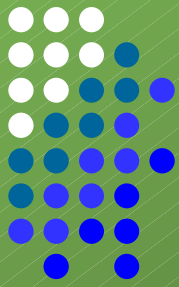
A Quick Derivation (1 of 4)

- In the mass eigenstate basis, the mass matrix is

$$M_{\text{mass}} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix}$$

- In the flavor eigenstate basis, it is then

$$\begin{aligned} M_{\text{flavor}} &= U(\theta_0) M_{\text{mass}} U^\dagger(\theta_0) \\ &= (m_2^2 - m_1^2) \begin{pmatrix} -\cos(2\theta_0) & \sin(2\theta_0) \\ \sin(2\theta_0) & \cos(2\theta_0) \end{pmatrix} + \frac{1}{2} (m_2^2 + m_1^2) \mathbb{1} \end{aligned}$$



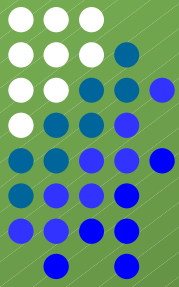
A Quick Derivation (2 of 4)

- In the non-relativistic limit, $\sqrt{p^2 + m^2} \approx p + \frac{m^2}{2p}$
- The kinetic energy is then

$$\begin{aligned} T &= p\mathbb{1} + \frac{1}{2p}M_{\text{flavor}}^2 \\ &= \frac{1}{4p} (m_2^2 - m_1^2) \begin{pmatrix} -\cos(2\theta_0) & \sin(2\theta_0) \\ \sin(2\theta_0) & \cos(2\theta_0) \end{pmatrix} + \frac{1}{4p} (4p^2 + m_2^2 + m_1^2) \mathbb{1} \end{aligned}$$

- Wolfenstein (1978) derives potential term (MSW effect)
 - Matter almost entirely first-generation leptons and quarks
 - Weak charged current interactions single out the electron neutrino component

$$V_{\text{MSW}} = \sqrt{2}G_F n_e \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$



A Quick Derivation (3 of 4)

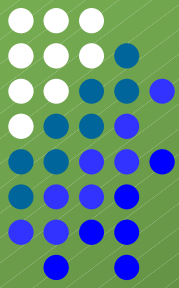
- Define $\delta m^2 = m_2^2 - m_1^2$
- Drop terms proportional to the identity
- Define $\Omega = \frac{2\sqrt{2}G_F n_e E}{\delta m^2}$

$$H_{\text{eff}} = \frac{\delta m^2}{4} \begin{pmatrix} -(\cos(2\theta_0) - \Omega) & \sin(2\theta_0) \\ \sin(2\theta_0) & \cos(2\theta_0) - \Omega \end{pmatrix}$$

- After some algebra and trigonometry:

$$\sin^2(2\theta_m) = \frac{\sin^2(2\theta_0)}{\sin^2(2\theta_0) + (\cos(2\theta_0) - \Omega)^2}$$

- Mixing angle changes in matter!



A Quick Derivation (4 of 4)

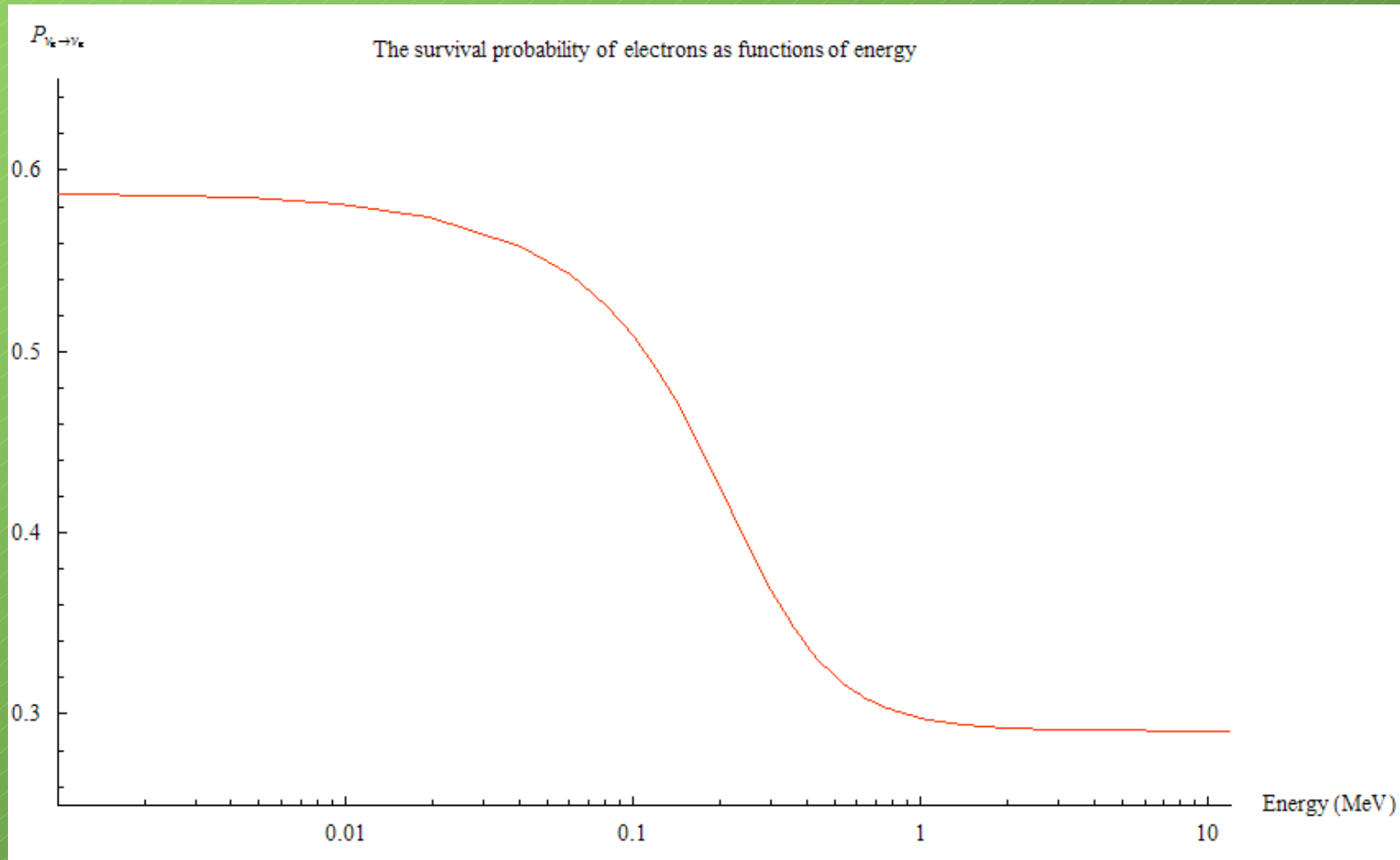
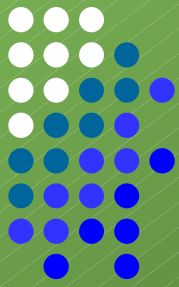
- From QM,

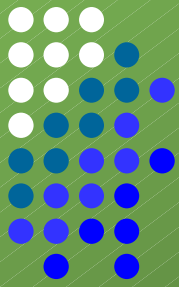
$$|\nu_j(t)\rangle = |\nu_j(0)\rangle \exp\left(-i \int_0^t E_j(t') dt'\right)$$

- Use the adiabatic approximation (essentially assuming density varies slowly):

$$\begin{aligned} P_{\nu_e \rightarrow \nu_e} &= |\langle \nu_e(T) | \nu_e(0) \rangle|^2 \\ &= \cos^2(\theta_m) \cos^2(\theta_0) + \sin^2(\theta_m) \sin^2(\theta_0) + \\ &\quad \frac{1}{2} \sin(2\theta_m) \sin(2\theta_0) \cos\left(\int_0^T (E_2(t') - E_1(t')) dt'\right) \\ &\rightarrow \frac{1}{2} + \frac{1}{2} \cos(2\theta_m) \cos(2\theta_0) \end{aligned}$$

No More Derivations!

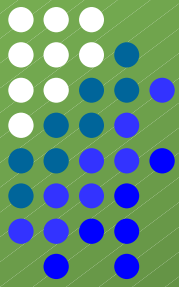




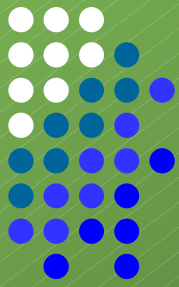
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Mass Varying Neutrinos (MaVaNs)



- Add a new scalar field, the “acceleron” A
- Postulate a heavy sterile (dark) neutrino
 - Mass dependent on expectation value $\langle h A_i \rangle$, which is a function of n_e
- Physical justification: measured dark energy density and neutrino energy density are similar
 - MaVaNs can help explain this without fine tuning
 - Many more cosmological justifications - see [hep-ph/0309800](https://arxiv.org/abs/hep-ph/0309800)
- But do they agree with experimental results?

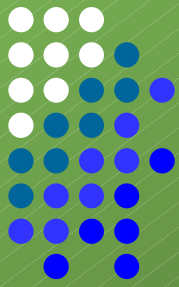


My Last Few Weeks

- Simplify model by integrating out heavy sterile neutrino
 - Assume $M_{\text{sterile}} = K n_e^r$
- New Hamiltonian is

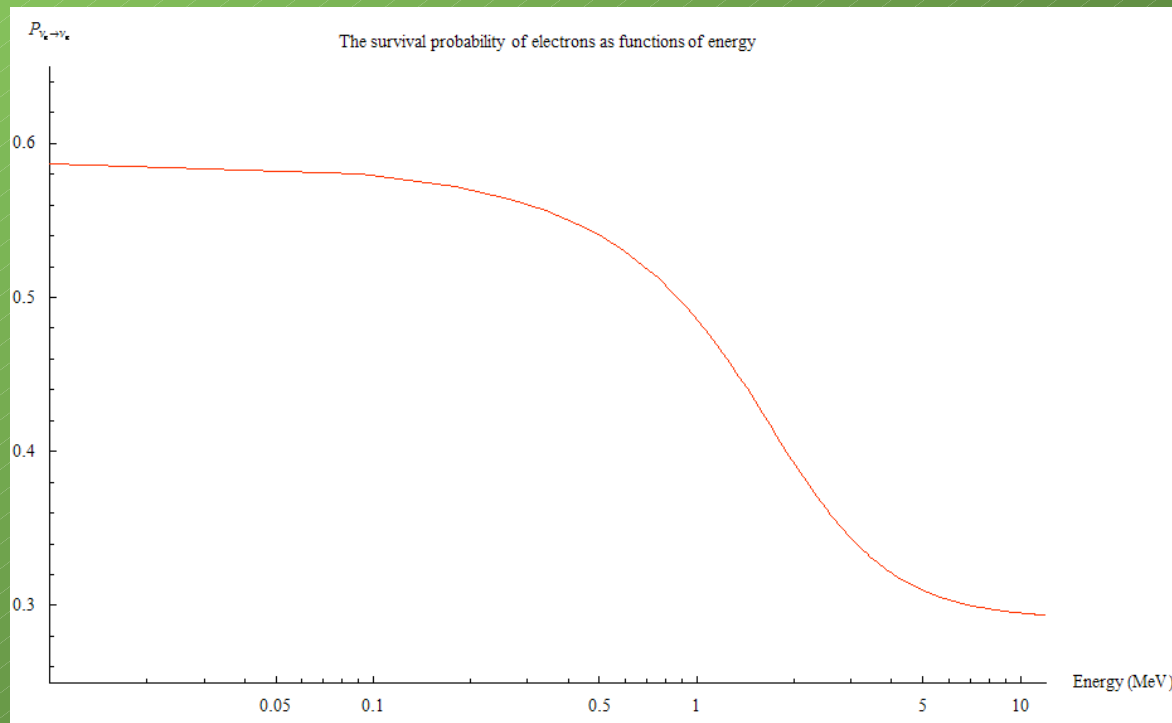
$$H_{\text{MaVaN}} = \frac{1}{2E} \frac{m_D^4}{K^2} n_e^{-2r} \begin{pmatrix} \frac{2\sqrt{2}K^2 G_F n_e^{2r+1} E}{m_D^4} + \sin^2(\theta) & \sin(\theta) \cos(\theta) \\ \sin(\theta) \cos(\theta) & \cos^2(\theta) \end{pmatrix}$$

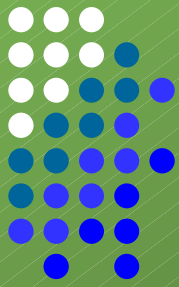
- Radically different electron density dependence



Reproducing Measurements

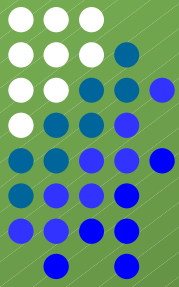
- Theory of neutrino masses has a new basis
 - Must reproduce experimental results to be useful
- We can!





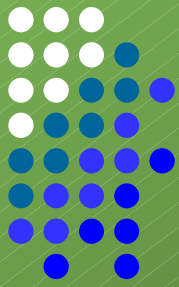
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Outlook

- This is a very positive result -- no a priori reason that MaVaNs should reproduce experimental results
- Experiments like KamLAND provide further constraints to test
- Lots to explore!



Thanks

- To my two great advisors, Ann Nelson and Neal Weiner
- To my fellow REU students
- To UW Physics
- To the NSF